

Fingerprints of freeze-in dark matter in an early matter-dominated era

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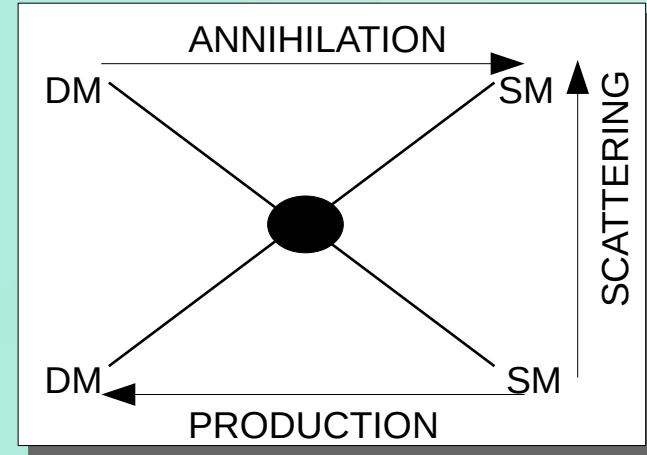
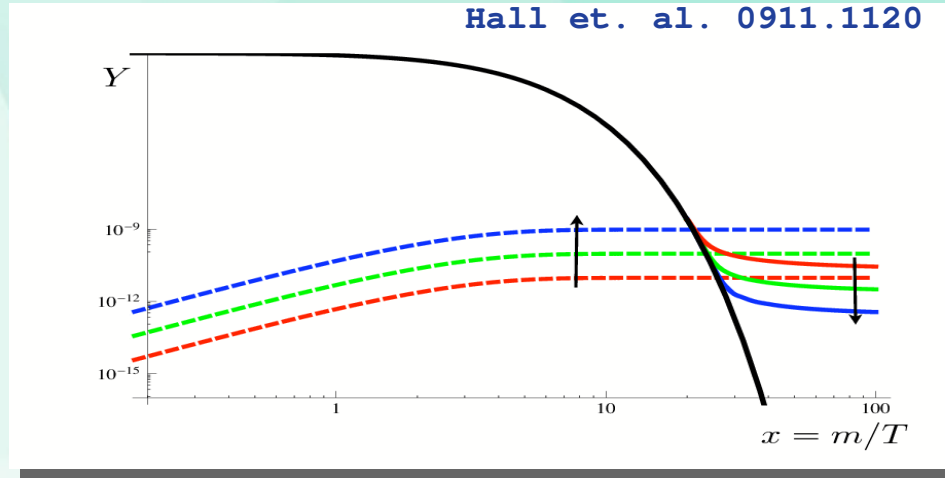
Based on 2204.03670
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Freeze-out vs. Freeze-in



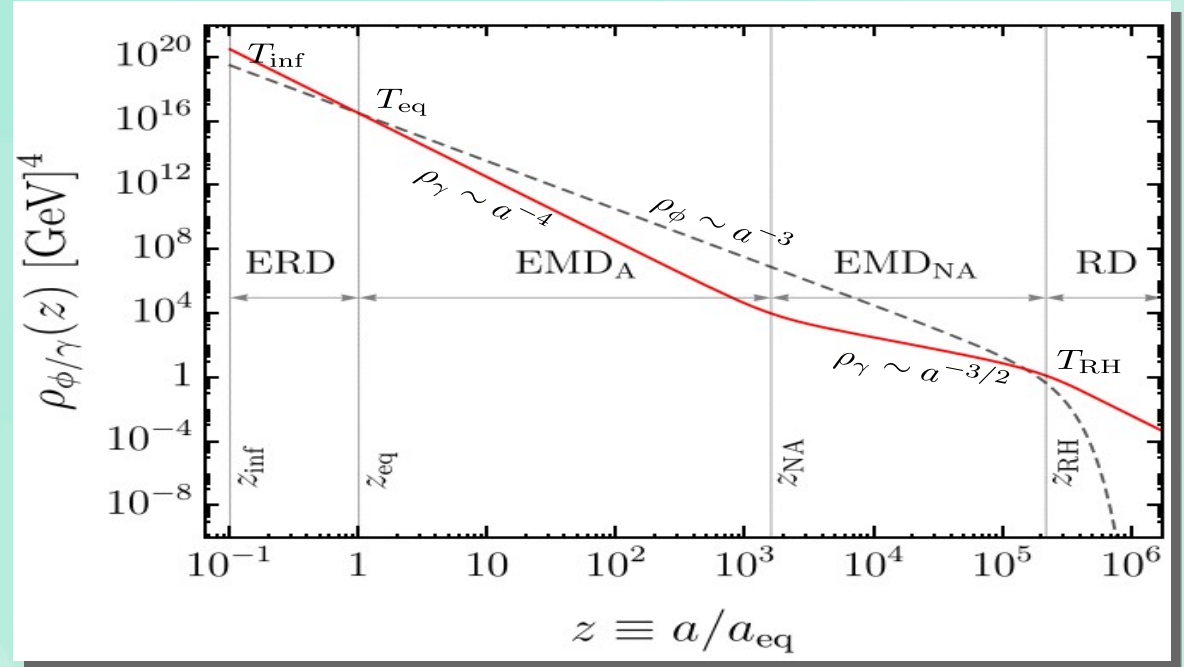
Freeze-out	Freeze-in
DM-SM in thermal equilibrium, Large coupling required	DM-SM never in thermal equilibrium, Extremely small coupling
At high temperature DM has thermal abundance	Initial abundance of DM at the end of inflation is negligible
Cosmology at high temperature is irrelevant	Cosmology from the end of inflation till today impacts relic

Early matter domination

BSM motivations for EMD

- Meta-stable matter fields
- Oscillating scalar fields
- Moduli
- SUSY condensates
- Dilaton
- Q-balls
- Curvaton

$$\begin{aligned}\dot{\rho}_\phi + 3(1 + \omega)H\rho_\phi &= -(1 + \omega)\Gamma_\phi\rho_\phi \\ \dot{\rho}_\gamma + 4H\rho_\gamma &= (1 + \omega)\Gamma_\phi\rho_\phi \\ H &= \frac{1}{\sqrt{3}M_p} \sqrt{\rho_\phi + \rho_\gamma}\end{aligned}$$



Constraints from BBN: $T_{RH} \gtrsim \text{few MeV}$

Evolution is dependent on the dissipation rate

Matter dissipation rate

In general depends on the temperature and the expansion of the universe

$$\Gamma_\phi = \hat{\Gamma} \left(\frac{T}{T_{\text{eq}}} \right)^n \left(\frac{a}{a_{\text{eq}}} \right)^k$$

Examples:

**Oscillating scalar fields
with $V(\phi) \sim \phi^p$ potential**

$$\Gamma_{\phi \rightarrow f\bar{f}} \propto m_\phi(t) \propto a^{-3(p-2)/(p+2)}, \quad (\text{for fermionic decay}),$$

$$\Gamma_{\phi \rightarrow bb} \propto m_\phi^{-1}(t) \propto a^{3(p-2)/(p+2)}, \quad (\text{for bosonic decay}),$$

Garcia et. al. 2012.10756

Moduli decay:

$$\Gamma_\phi \propto \frac{T^3}{M_p^2}$$

Bodeker, hep-ph/0605030

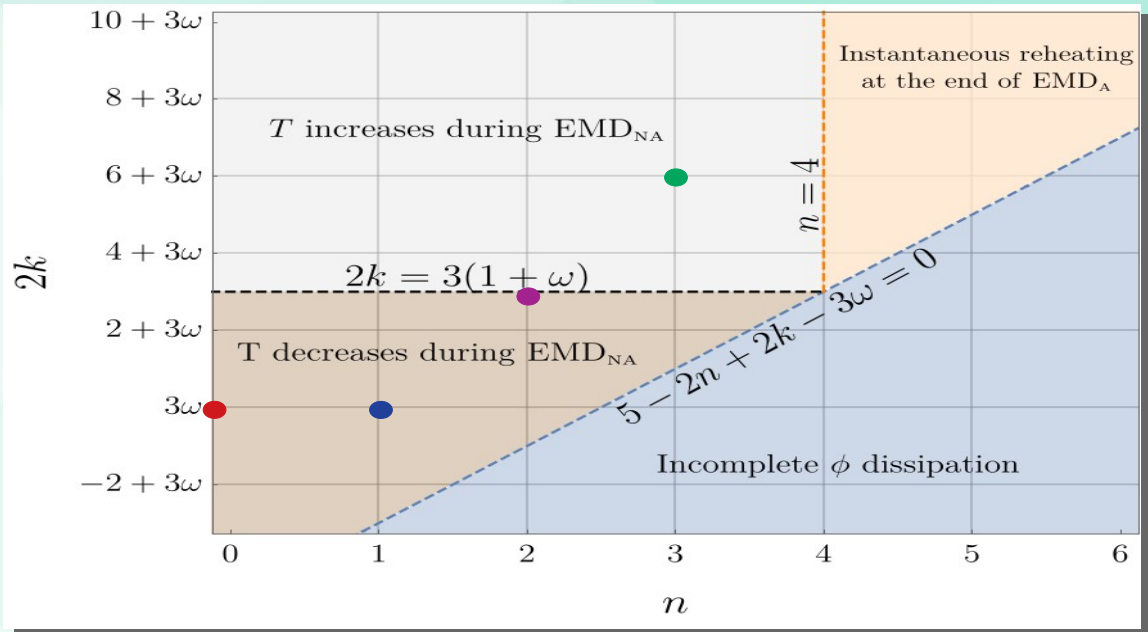
More Examples:

Γ_ϕ	(n, k, ω)	$T(z)$ during EMD _{NA}
const.	(0, 0, 0)	decreases with z
T	(1, 0, 0)	decreases with z
$\langle \phi \rangle^{-2}$	(0, 3, 0)	increases with z
$\frac{T^3}{\langle \phi \rangle^2}$	(3, 3, 0)	increases with z
$\frac{T^2}{\langle \phi \rangle}$	(2, 3/2, 0)	remains constant
$\frac{T^2}{\langle \phi \rangle}$	(2, 6/5, 1/5)	decreases with z

Mukaida et. al. 1208.3399, 1212.4985

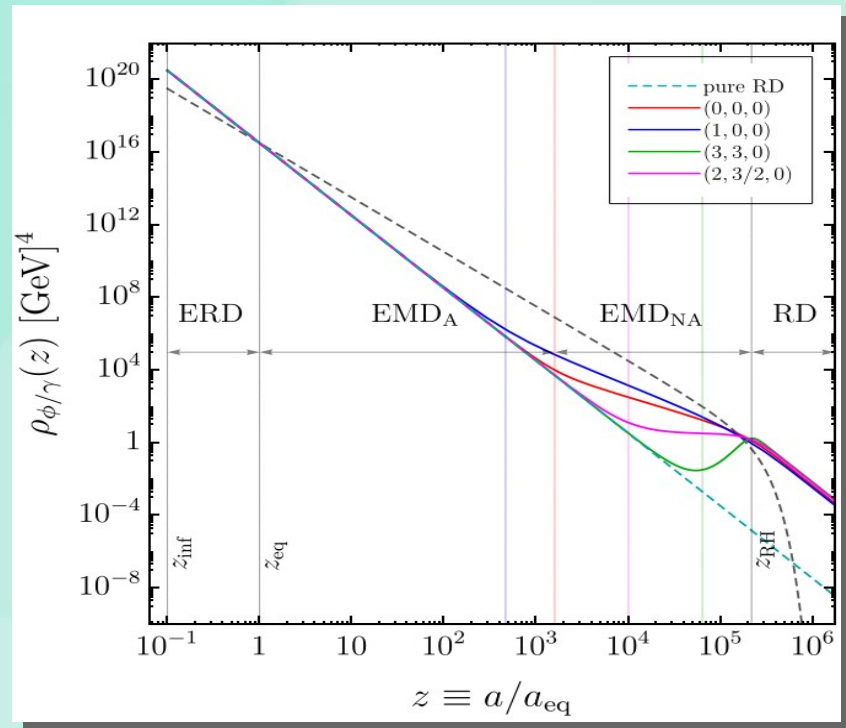
Drewes, 1406.6243

Co et. al. 2007.04328



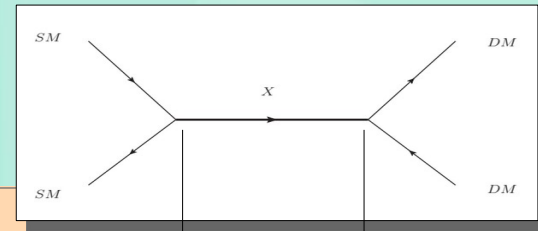
$$\Gamma_\phi = \hat{\Gamma} \left(\frac{T}{T_{\text{eq}}} \right)^n \left(\frac{a}{a_{\text{eq}}} \right)^k$$

Epoch	z	$T(z)$	$H(z)$
ERD	$z_{\text{inf}} < z < 1$	$\frac{T_{\text{eq}}}{z}$	$\sqrt{\frac{\rho_\gamma(T_{\text{eq}})}{3M_p^2}} z^{-2}$
EMD_A	$1 < z < z_{\text{NA}}$	$\frac{T_{\text{eq}}}{z}$	$\sqrt{\frac{\rho_\gamma(T_{\text{eq}})}{3M_p^2}} z^{-\frac{3}{2}(1+\omega)}$
EMD_{NA}	$z_{\text{NA}} < z < z_{\text{RH}}$	$T_{\text{RH}} \left(\frac{z}{z_{\text{RH}}} \right)^{\frac{\delta-8+2n}{8-2n}}$	$\sqrt{\frac{\rho_\gamma(T_{\text{eq}})}{3M_p^2}} z^{-\frac{3}{2}(1+\omega)}$
RD	$z_{\text{RH}} < z$	$T_{\text{eq}} z_{\text{RH}}^{\frac{1-3\omega}{4}} z^{-1}$	$\frac{\sqrt{\rho_\gamma(T_{\text{RH}})}}{\sqrt{3}M_p} \left(\frac{z}{z_{\text{RH}}} \right)^{-2}$



Freeze-in production rate

$$\dot{n}_\chi + 3Hn_\chi = -\langle\sigma v\rangle (n_\chi^2 - n_\chi^{\text{eq}2}) = R(T)$$

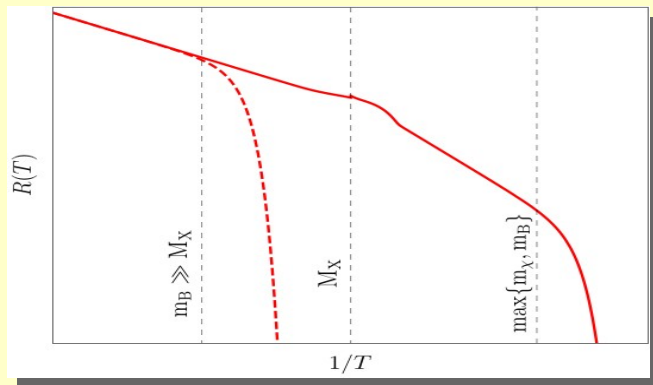


$$g_{X\text{-SM}} \ll 1 \quad g_{X\text{-DM}} \sim \mathcal{O}(1)$$

$$R(T) \sim \frac{T^{p+4}}{M_X^4} e^{-\frac{2\max\{m_X, m_B\}}{T}}$$

T_{inf}

M_X

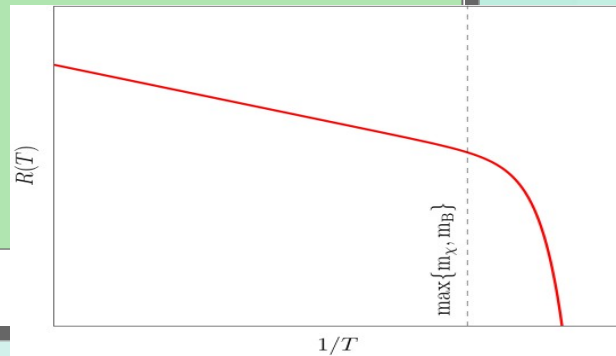


$$R(T) \sim \begin{cases} T^p e^{-\frac{2m_B}{T}}, & m_B \gg M_X \\ T^p, & m_B \ll M_X \ll T \\ \frac{\pi T M_X^p}{\Gamma_X} K_1\left(\frac{M_X}{T}\right), & m_B \ll M_X \sim T \\ \frac{T^{p+4}}{M_X^4} e^{-\frac{2m_B}{T}}, & m_B, T \ll M_X \end{cases}$$

$$M_X = 2m_X$$

$$R(T) \sim T^p e^{-\frac{2\max\{m_X, m_B\}}{T}}$$

m_X

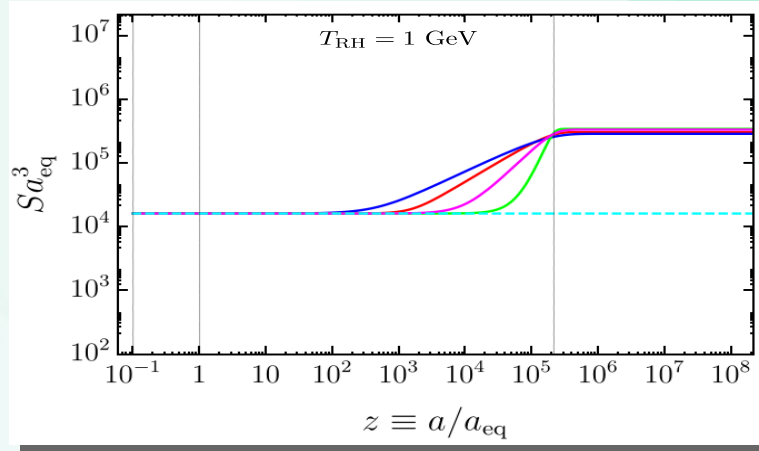


Freeze-in DM yield during EMD

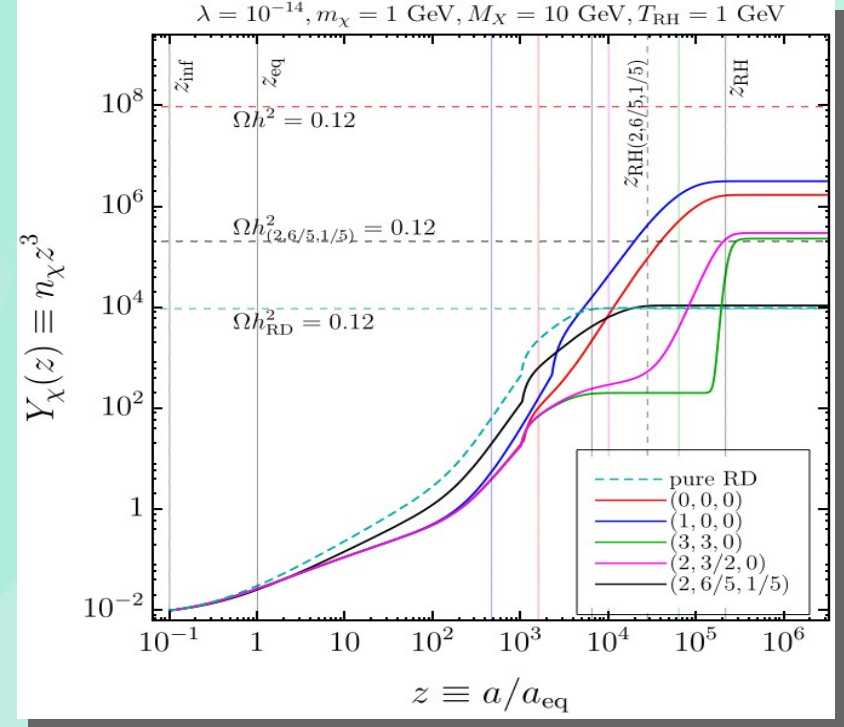
- DM yield dilutes due to entropy production
- Non-trivial temperature evolution also changes the DM production rate during non-adiabatic EMD

$$\frac{\Omega h^2}{\Omega h_{RD}^2} = \frac{Y_\chi(z_0)}{Y_\chi^{RD}(z_0)} \left(\frac{z_0^{RD}}{z_0} \right)^3 = \frac{Y_\chi(z_0)}{Y_\chi^{RD}(z_0)} \left(\frac{T_{RH}}{T_{eq}} \right)^{\frac{1-3\omega}{1+\omega}}$$

$$\sim 10^{-2} - 10^{-3}$$



$$\frac{dY_\chi(z)}{dz} = \frac{z^2 R(T(z))}{H(z)}$$

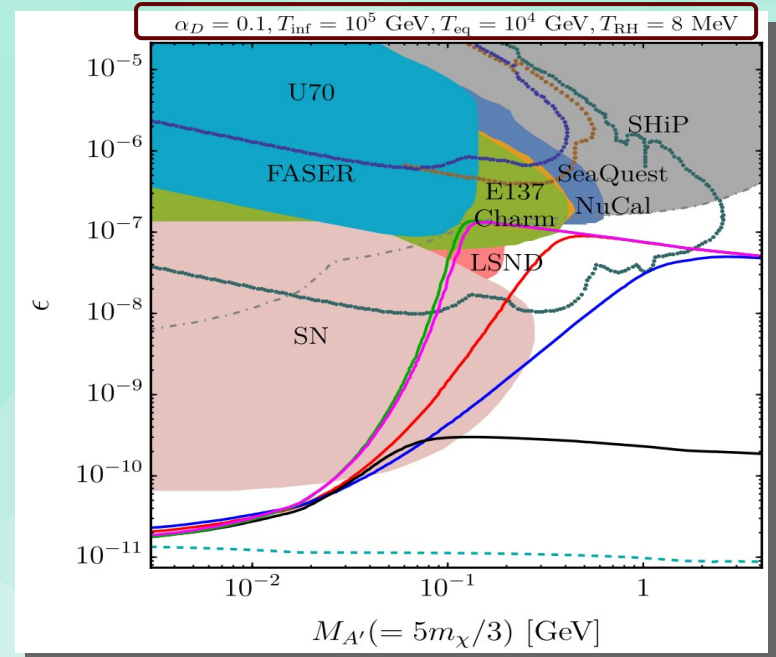
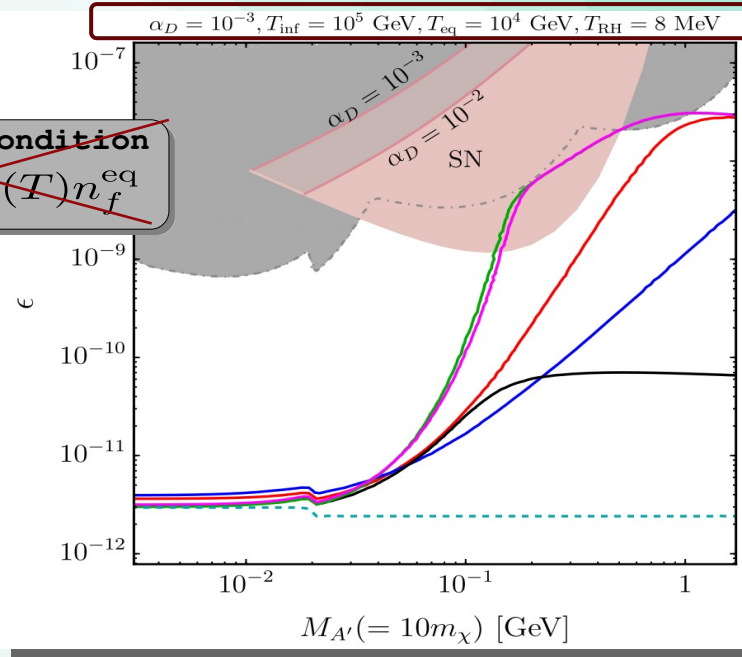


Larger coupling is required to saturate freeze-in relic in presence of EMD epoch

Dark photon portal

$$\mathcal{L} = -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} - \frac{\epsilon}{2}F'_{\mu\nu}F^{\mu\nu} + \frac{1}{2}M_{A'}^2 A'_\mu A'^\mu + \bar{\chi}(i\partial - m_\chi)\chi + g_D \bar{\chi}\gamma^\mu A'_\mu \chi$$

~~Freeze-in condition~~
 ~~$R(T) \ll H(T)n_f^{\text{eq}}$~~



Parameter space satisfying observed relic is accessible to experiments in the presence of an early matter dominated era

Summary

- Freeze-in DM relic depends on the **non-standard epochs of cosmology** at high temperatures
- An epoch of **pre-BBN early matter domination** leads to **freeze-in with larger couplings**
- Details depend crucially on the **temperature and expansion dependent dissipation rate** of the dominating matter field
- **Dark photon portal dark matter** model may come under the **experimental radar** in presence of early matter domination

**THANK
YOU**