

The muon g - 2 short-distance

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# SHORT DISTANCE CONTRIBUTIONS TO THE MUON g - 2



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## Introduction

- Magnetic moment:  $\vec{\mu} = g \frac{q}{2m} \vec{S}$
- For angular momentum: g = 1
- Dirac equation: g = 2
- Structure and/or QFT give different values
- Anomaly:  $a = \frac{g-2}{2}$ • QED (Schwinger):  $a = 1 + \frac{\alpha}{2\pi} + \dots$



Jacob Bourjaily/Wikipedia



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# Why do we do this?

The muon  $a_{\mu}=rac{g_{\mu}-2}{2}$  will be measured more precisely





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#### Fermilab



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- Experiments done for electron and muon: very precise
- $a_e^{\exp} = 115965218.073(0.028) \ 10^{-11}$
- Discrepancy 2.4 $\sigma$ :  $\alpha$  from Cs-atom interferometry or  $a_e$  $\alpha^{-1} = 137.0359990460(270)(Cs)$  $\alpha^{-1} = 137.0359991496(330)(a_e)$
- Both errors dominated by experiment (and largely systematic)
- Difference not relevant for QED part of  $a_{\mu}$

#### Introduction

- Experiments done for electron and muon: very precise
- $a_e^{\exp} = 115965218.073(0.028) \ 10^{-11}$
- $a_{\mu}^{\exp} = 116592089(54)(33) \ 10^{-11}$
- $a_{\mu}^{\text{the}} = 116591810(43) \ 10^{-11}$
- Discrepancy 3.7*σ*: 279(76) 10<sup>-11</sup>
- White paper: arXiv: 2006.04822 Phys. Rep. 887 (220) 1-166
   JB,N.Hermansson-Truedsson,S.Leupold,A.Rodríguez-Sánchez (132 authors)
   Large theory collaboration, error: consensus
   Numbers (except the new ones here) taken from there
- Impressive agreement theory experiment about 13 digits in a<sub>e</sub> and about 9 in a<sub>μ</sub>

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White paper



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# Introduction



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- $a_{\mu}^{SM}=a_{\mu}^{QED}+a_{\mu}^{EW}+a_{\mu}^{had}$
- $a_{\mu}^{QED} = 116584718.931(104) \ 10^{-11}$

up to 4 loops essentially analytically, 5-loops numerically, 6-loops estimate (main uncertainty)

•  $a_{\mu}^{EW} = 153.6(1.0) \ 10^{-11}$ 

Done to two-loop, main uncertainty from long-distance hadronic contributions (via anomaly)

•  $a_{\mu}^{\exp} - a_{\mu}^{QED} - a_{\mu}^{EW} = 7216(63) \ 10^{-11}$ 

## Hadronic contributions





- Muon and photon lines, representative diagrams
- The blobs are hadronic contributions
- There are higher order contributions of both types: known accurately enough
- $a_{\mu}^{HVP} = 6845(40) \ 10^{-11} \ (LO+NLO+NNLO)$

• 
$$a_{\mu}^{HLbL} = 92(18) \ 10^{-11} \ (LO+NLO)$$

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# HLbL: the main object to calculate



- Muon line and photons: well known
- The blob: fill in with hadrons/QCD
- Trouble: low and high energy very mixed
- q<sub>4</sub> always at zero
- Double counting needs to be avoided: hadron exchanges versus quarks



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• "Long distance": under good control

- Dispersive method: Berne group around G. Colangelo
- $\pi^0$  (and  $\eta, \eta'$ ) pole: 93.8(4.0) 10<sup>-11</sup>
- Pion and kaon box (pure): -16.4(2) 10<sup>-11</sup>
- $\pi\pi$ -rescattering (include scalars below 1 GeV):-8(1) 10<sup>-11</sup>
- Charm (beauty, top) loop: 3(1) 10<sup>-11</sup>
- "Short and medium distance"
  - Axial vector: 6(6) 10<sup>-11</sup>
  - Short-distance: 15(10) 10<sup>-11</sup>
- Clearly the last item needs improvement

## Short-distance



- =  $\Pi^{\mu\nu\lambda\sigma}(q_1, q_2, q_3)$
- Actually we really need  $\frac{\delta\Pi^{\mu
  u\lambda\sigma}(q_1,q_2,q_3)}{\delta q_{4
  ho}}$
- Mixed short-distance:  $q_4$  at zero,  $q_1^2, q_2^2, q_3^2$  large •  $q_i^2 = -Q_i^2$



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- Quarkloop Quarkloop constituent SD: naive SD: correct SD: numerical
- SD: perturbative

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 $a_4=0$ 

- Use (constituent) quark loop
- Used for full estimates since the beginning (1970s)
- Used for short-distance estimates with mass as a cut-off JB, Pallante, Prades, 1996

• We recalculated:



- In agreement with quarkloop formulae from Hoferichter, Stoffer, private communication
- In agreement with known numerics



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SD: perturbativ

# Quarkloop: *u*, *d*, *s*





- $M_Q$  provides an infrared cut-off,  $M_Q 
  ightarrow 0$  divergent
- About  $12 \times 10^{-11}$  from above 1 GeV for  $M_Q = 0.3$  GeV
- About  $17 \times 10^{-11}$  from above 1 GeV for  $M_Q = 0$

# Quarkloop



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Conclusions

#### • Is it a first term in a systematic OPE?

- OPE has been used as constraints on specific contributions
  - $\pi^0 \gamma^* \gamma^*$  asymptotic behaviour
  - Constraints on many other hadronic formfactors
  - $Q_1^2 pprox Q_2^2 \gg Q_3^2$  Melnikhov, Vainshtein 2003
- JB, N. Hermansson-Tuedsson, A. Rodríguez-Sánchez, Phys.Lett. B798 (2019) 134994[arxiv:1908.03331]
   +Laub, JHEP 10 (2020) 203 [arxiv:2008.13487]
   +to be published (soon?)

# Short-distance: first attempt

$$\Pi^{\mu\nu\lambda\sigma} = -i \int d^4x d^4y d^4z e^{-i(q_1\cdot x + q_2\cdot y + q_3\cdot z)} \left\langle T\left(j^{\mu}(x)j^{\nu}(y)j^{\lambda}(z)j^{\sigma}(0)\right)\right\rangle$$

- Usual OPE: x, y, z all small
- First term in the expansion is the quark-loop no problem with  $\partial/\partial q_4^{\rho}$  and  $q_4 \rightarrow 0$

*p* in loop  $\Rightarrow$  no singular propagators:

• Next term problems: no loop momentum;

 $q_4 
ightarrow 0$  propagator diverges:



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# Short-distance: correctly

- Similar problem in QCD sum rules for electromagnetic radii and magnetic moments
- loffe, Smilga, Balitsky, Yung, 1983
- For the q<sub>4</sub>-leg use a constant background field and do the OPE in the presence of that constant background field
- Use radial gauge:  $A_4^{\lambda}(w) = \frac{1}{2}w_{\mu}F^{\mu\lambda}$ whole calculation is immediately with  $q_4 = 0$ .
- First term is exactly the usual quark loop (quark masses: next order)





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# Short-distance: next term(s)

- Do the usual QCD sum rule expansion in terms of vacuum condensates
- There are new condensates, induced by the constant magnetic field:  $\langle \bar{q}\sigma_{\alpha\beta}q\rangle \equiv e_q F_{\alpha\beta}X_q$
- Lattice QCD Bali et al., arXiv:1206.4205  $X_u = 40.7 \pm 1.3$  MeV,
- Only starts at  $1/Q^2$  via  $m_q X_q$  corrections to the leading quark-loop result
- $X_q$  and  $m_q$  are very small, only a very small correction
- X<sub>q</sub>: contain IR divergent perturbative parts, combine with the m<sup>2</sup><sub>q</sub> corrections from the quark-loop consistently
- Next order: very many condensates contribute, lots of IR mixing and redefinitions.



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# Short-distance: numerical results

| Order            | Contribution          | $Q_{ m min}=1{ m GeV}$ | $Q_{\min} = 2 \mathrm{GeV}$ |  |
|------------------|-----------------------|------------------------|-----------------------------|--|
| $1/Q_{ m min}^2$ | quark-loop            | $1.73\cdot10^{-10}$    | $4.35\cdot10^{-11}$         |  |
| $1/Q_{ m min}^4$ | quark-loop, $m_q^2$   | $-5.7\cdot10^{-14}$    | $-3.6\cdot10^{-15}$         |  |
|                  | X <sub>2,m</sub>      | $-1.2\cdot10^{-12}$    | $-7.3 \cdot 10^{-14}$       |  |
| $1/Q_{\min}^6$   | $X_{2,m^3}$           | $6.4 \cdot 10^{-15}$   | $1.0\cdot10^{-16}$          |  |
| ,                | $X_3$                 | $-3.0\cdot10^{-14}$    | $-4.7 \cdot 10^{-16}$       |  |
|                  | <i>X</i> <sub>4</sub> | $3.3\cdot10^{-14}$     | $5.3 \cdot 10^{-16}$        |  |
|                  | $X_5$                 | $-1.8\cdot10^{-13}$    | $-2.8 \cdot 10^{-15}$       |  |
|                  | $X_6$                 | $1.3\cdot10^{-13}$     | $2.0 \cdot 10^{-15}$        |  |
|                  | <i>X</i> <sub>7</sub> | $9.2\cdot10^{-13}$     | $1.5\cdot10^{-14}$          |  |
|                  | $X_{8,1}$             | $3.0\cdot10^{-13}$     | $4.7 \cdot 10^{-15}$        |  |
|                  | X <sub>8,2</sub>      | $-1.3\cdot10^{-13}$    | $-2.0 \cdot 10^{-15}$       |  |

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•  $Q_1, Q_2, Q_3 \ge Q_{\min}$ 

- Nonperturbative short-distance corrections are small
- Suppression by small quark masses or small condensates
- Nonperturbative short-distance corrections are small

# Short-distance: $1/Q_{\min}^2$

• Can we understand scaling with  $Q_{\min}$ ?

• 
$$a_{\mu} = \frac{2\alpha^3}{3\pi^2} \int_0^\infty dQ_1 dQ_2 Q_1^3 Q_2^3 \int_{-1}^1 d\tau \sqrt{1-\tau^2} \sum_{i=1,12} \hat{T}_i \overline{\Pi}_i$$

- Do  $Q_i 
  ightarrow \lambda Q_i$
- overall factor goes as  $\lambda^8$
- Quark loop has no scale thus  $\hat{\Pi}_i$  scale with their dimension  $\hat{\Pi}_1, \hat{\Pi}_4 \sim \lambda^{-4}, \qquad \hat{\Pi}_7, \hat{\Pi}_{17}, \hat{\Pi}_{39}, \hat{\Pi}_{54} \sim \lambda^{-6}$
- $\Longrightarrow \overline{\Pi}_{1,\dots,4} \sim \lambda^{-4}$   $\overline{\Pi}_{5,\dots,12} \sim \lambda^{-6}$
- Expand the  $T_i$  for  $Q_i \gg m_{\mu}$ :  $T_1 \sim m_{\mu}^4$ ,  $T_{i\neq 1} \sim m_{\mu}^2$  $T_1 \sim \lambda^{-8}$ ,  $T_{2,3,4} \sim \lambda^{-6}$ ,  $T_{5,...,12} \sim \lambda^{-4}$
- Put all together: quark-loop scales as  $a_{\mu}^{
  m SD~ql}\sim\lambda^{-2}$
- $m_q X_q$  adds an overall factor  $\Longrightarrow a_\mu^{{
  m SD} X_q} \sim \lambda^{-4}$



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- Representative diagram:
- Use method of master integrals: <sup>44</sup>disadvantage: large numerical cancellations between integrals

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- All integrals are known
- Infrared and UV divergences in individual diagrams
- Dimensional regularization:  $d = 4 2\epsilon$
- All  $1/\epsilon^3, 1/\epsilon^2, 1/\epsilon^2$  cancel
- Several independent calculations that agree
- Find some typos in integral papers (I hate signs)

## Perturbative corrections



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- Preliminary results
   a<sup>HLbL SD gluonic</sup><sub>μ</sub> = -1.7 10<sup>-11</sup>
- $Q_{\min}=1$  GeV,  $lpha_{\mathcal{S}}=0.33$
- Main uncertainty: how do we handle  $\alpha_S$
- No sign that it is very large (about -10%)

#### Conclusions

- We have shown that the massless quarkloop really is the first term of a proper OPE expansion for the HLbL
- We have shown how to properly go to higher orders
- We have calculated the next two terms in the OPE
  - NLO: suppressed by quark masses and a small  $X_q$
  - NNLO: large number of induced condensates but all small
  - Numerically not relevant at the present precision
- Gluonic corrections about -10%
- Why do this: matching of the sum over hadronic contributions to the expected short distance domain
- Finding the onset of the asymptotic domain

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