

# Direct detection of low-mass dark matter with strong matter interactions

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**Partikeldagarna 2019**

Linköping University, 02.10.2019

In collaboration with

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- Chris Kouvaris
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Based on

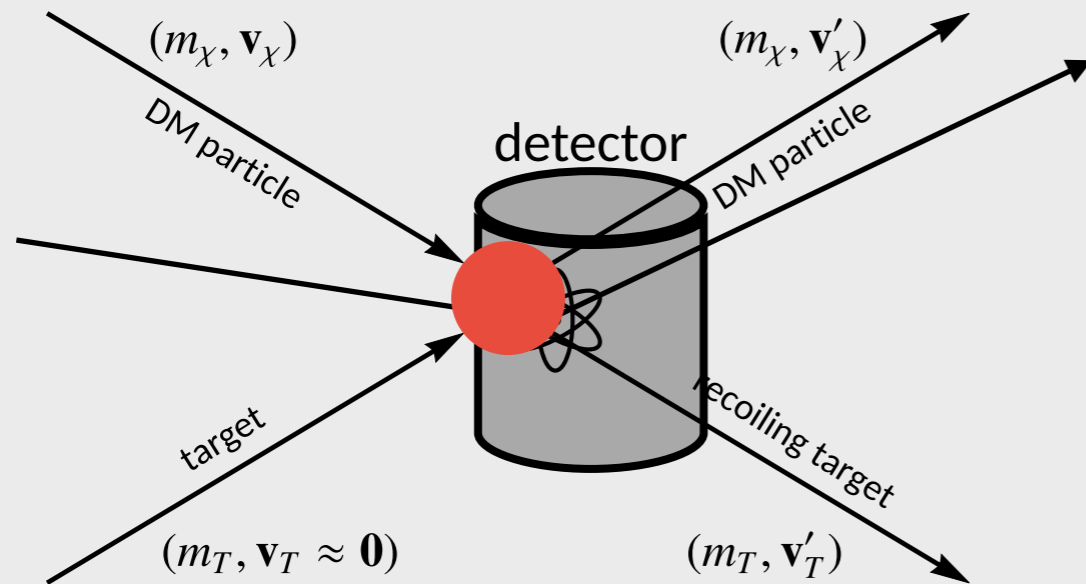
- [arXiv:1905.06348]
- [arXiv:1802.04764]



**CHALMERS**  
UNIVERSITY OF TECHNOLOGY

# Direct Detection of Dark Matter

Basic idea: Look for the aftermath of a DM-atom collision in a detector.



Nuclear recoils as observable for GeV-scale DM searches.

M.W. Goodman and E. Witten, Phys.Rev. D31 (1985) 3059  
 I. Wasserman, Phys. Rev. D33 (1986) 2071  
 A.K. Drukier et al., Phys. Rev. D33 (1986) 3495

Event spectrum:

$$\frac{dR}{dE_R} = N_T \frac{\rho_\chi}{m_\chi} \iiint d^3\mathbf{v} v f_\chi(\mathbf{v}) \frac{d\sigma_N}{dE_R} \Theta(v - v_{\min}(E_R)).$$

Detector size

Astrophysics

Particle physics

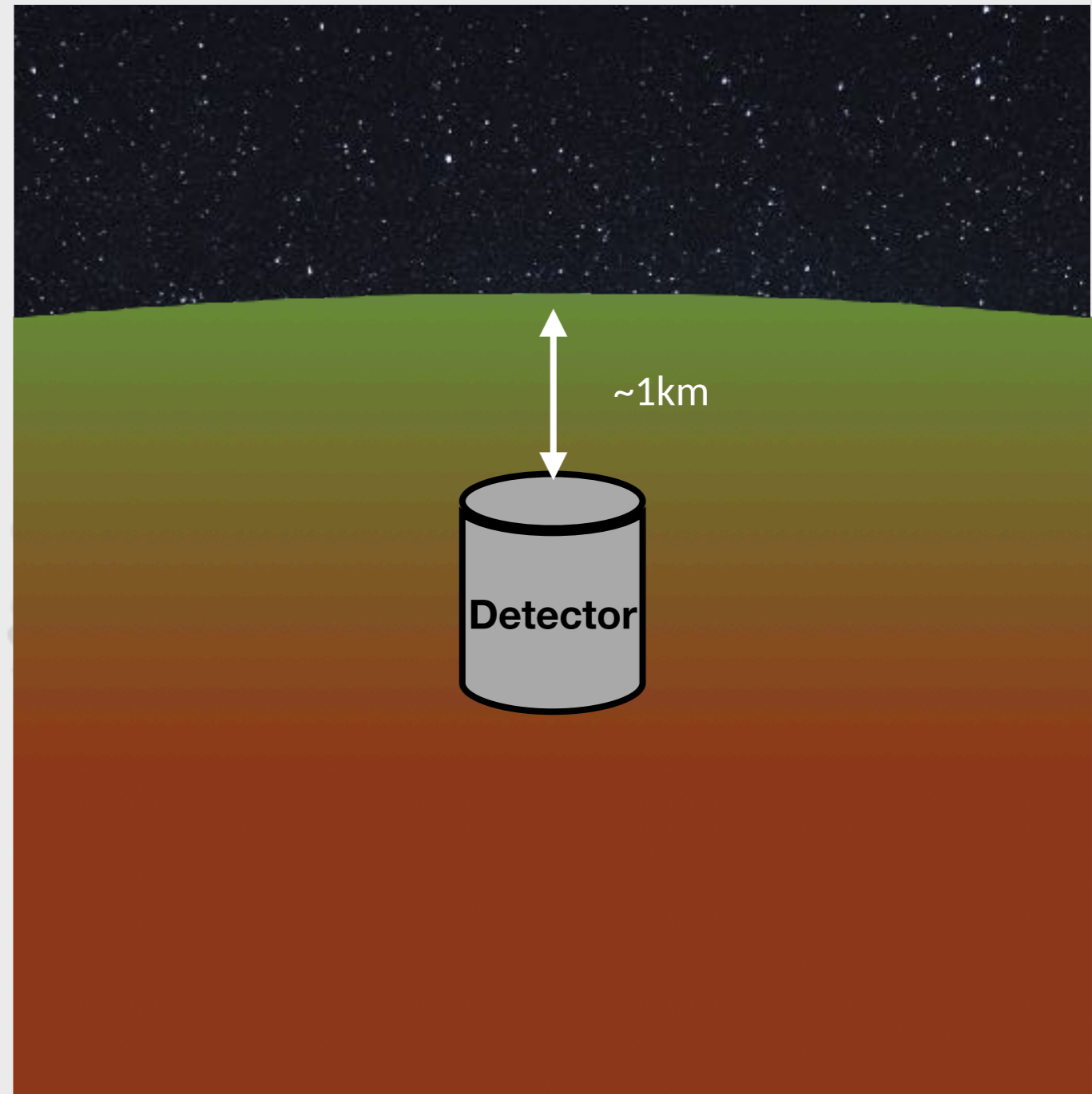
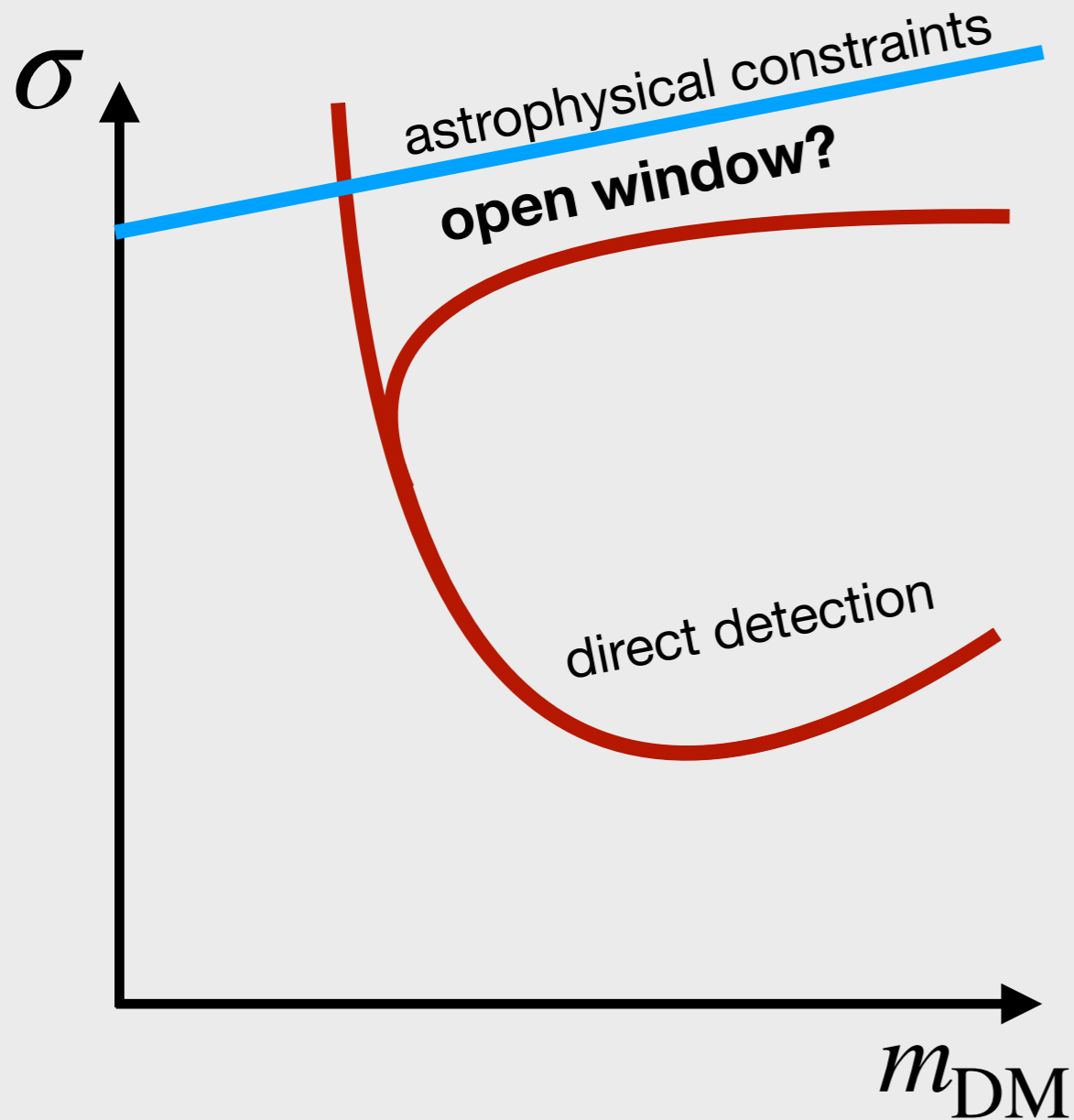
Kinematics

For sub-GeV DM, one can also look for DM-electron interactions.

R. Essig et al., Phys. Rev. D85 (2012) 076007

# Direct detection of strongly interacting DM

Goodman and Witten, Phys.Rev. D31 (1985) 3059  
Starkman et al, Phys.Rev. D41 (1990) 3594

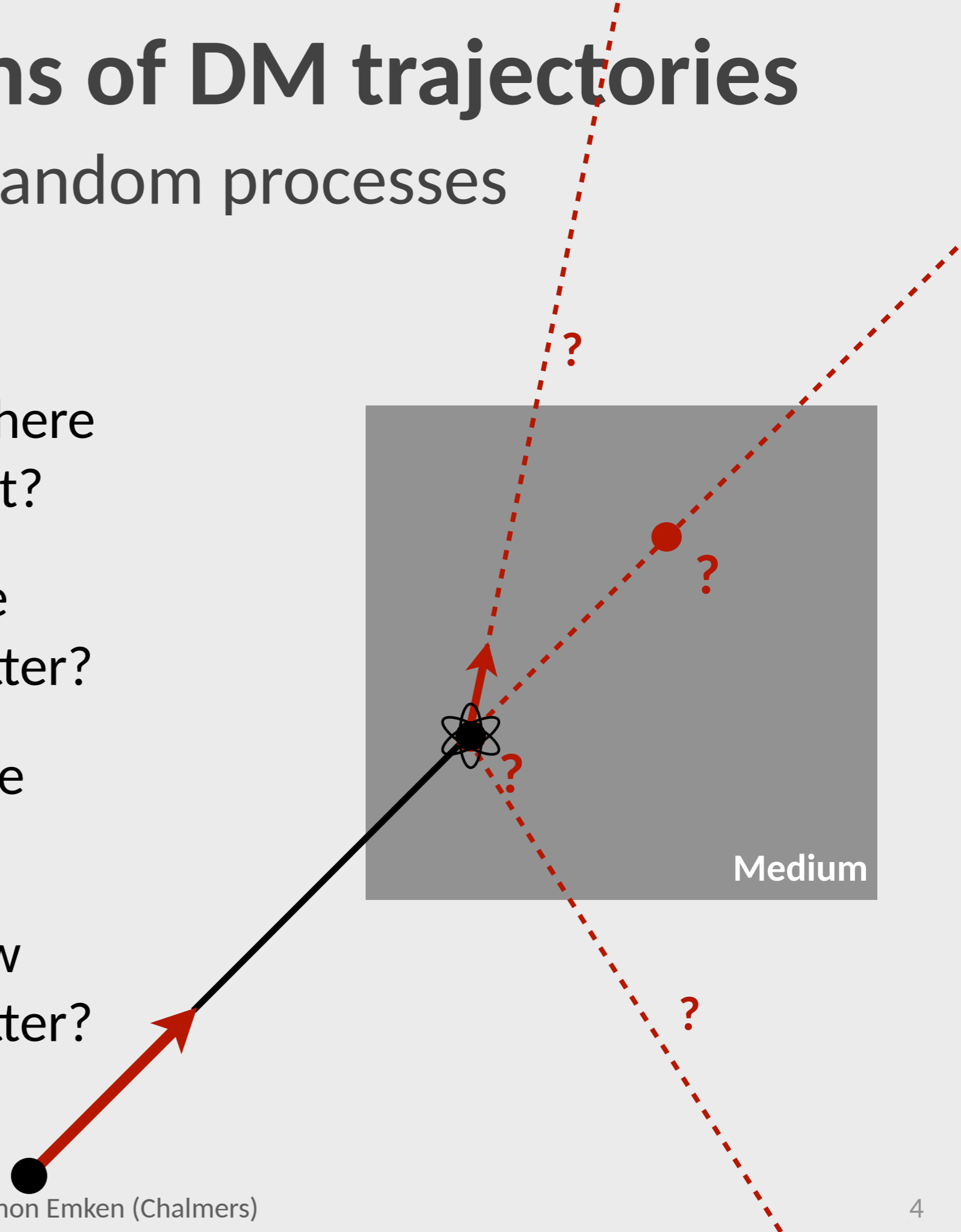


# MC simulations of DM trajectories

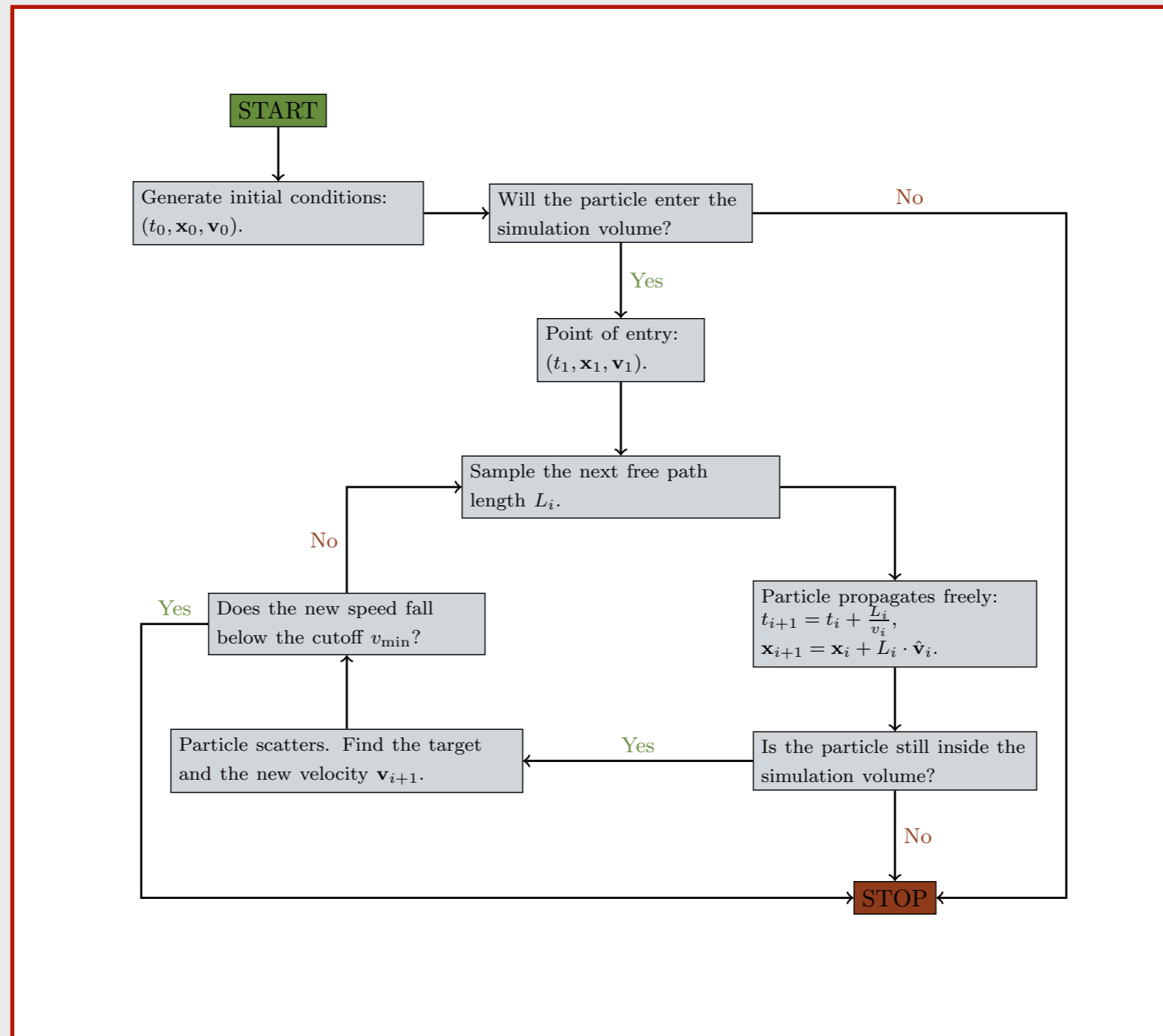
The fundamental random processes

1. **Initial Conditions:** Where does the particle start?
2. **Free distance:** Where does the particle scatter?
3. **Target:** What does the particle scatter on?
4. **Scattering angle:** How does the particle scatter?

**Repeat steps 2.-4.**



# MC simulation algorithm



$t=0$  s



Emken 2016

# MC simulation of the overburden of detectors

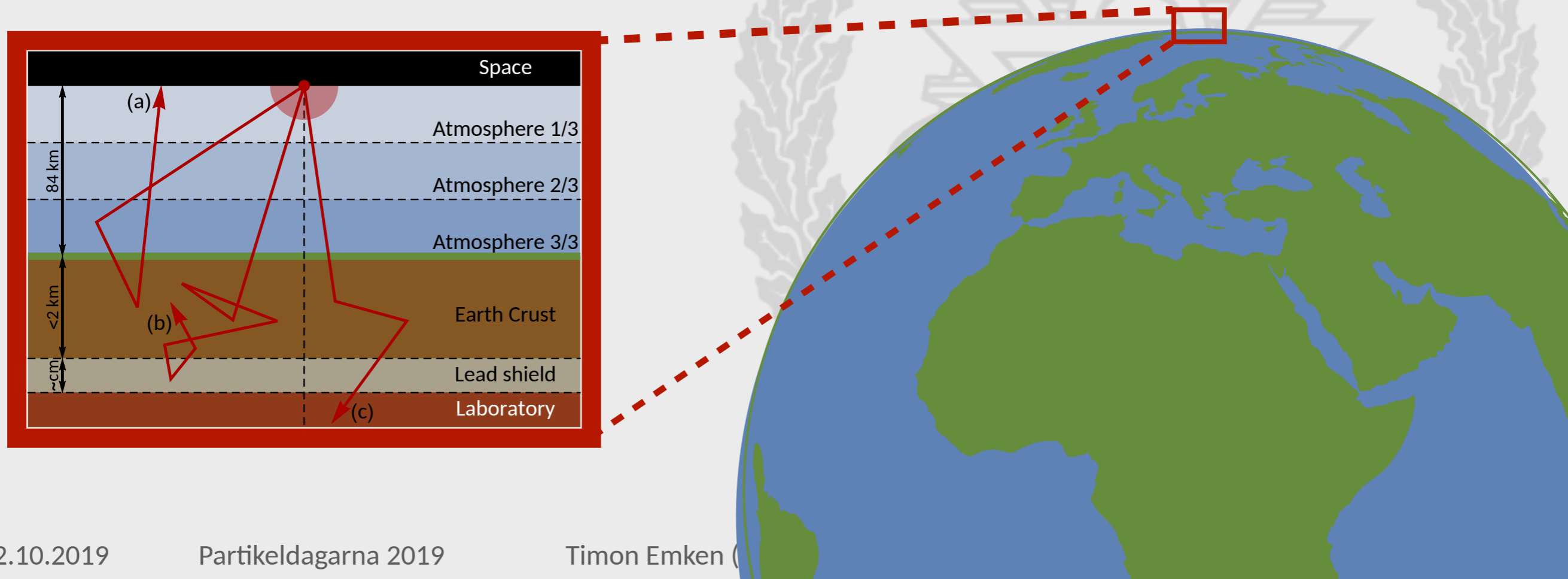
TE, C. Kouvaris, I.M. Shoemaker, Phys.Rev. D96 (2017) no.1, 015018

M.S. Mahdawi, G.R. Farrar, JCAP 1712 (2017) 004

TE, C. Kouvaris, Phys.Rev. D97 (2018) no.11, 115047

M.S. Mahdawi, G.R. Farrar, JCAP 1810 (2018) no.10, 007

To find the **critical** cross-section, where a given experiment loses sensitivity to strongly interacting DM, we only simulate the overburden, not the entire Earth.



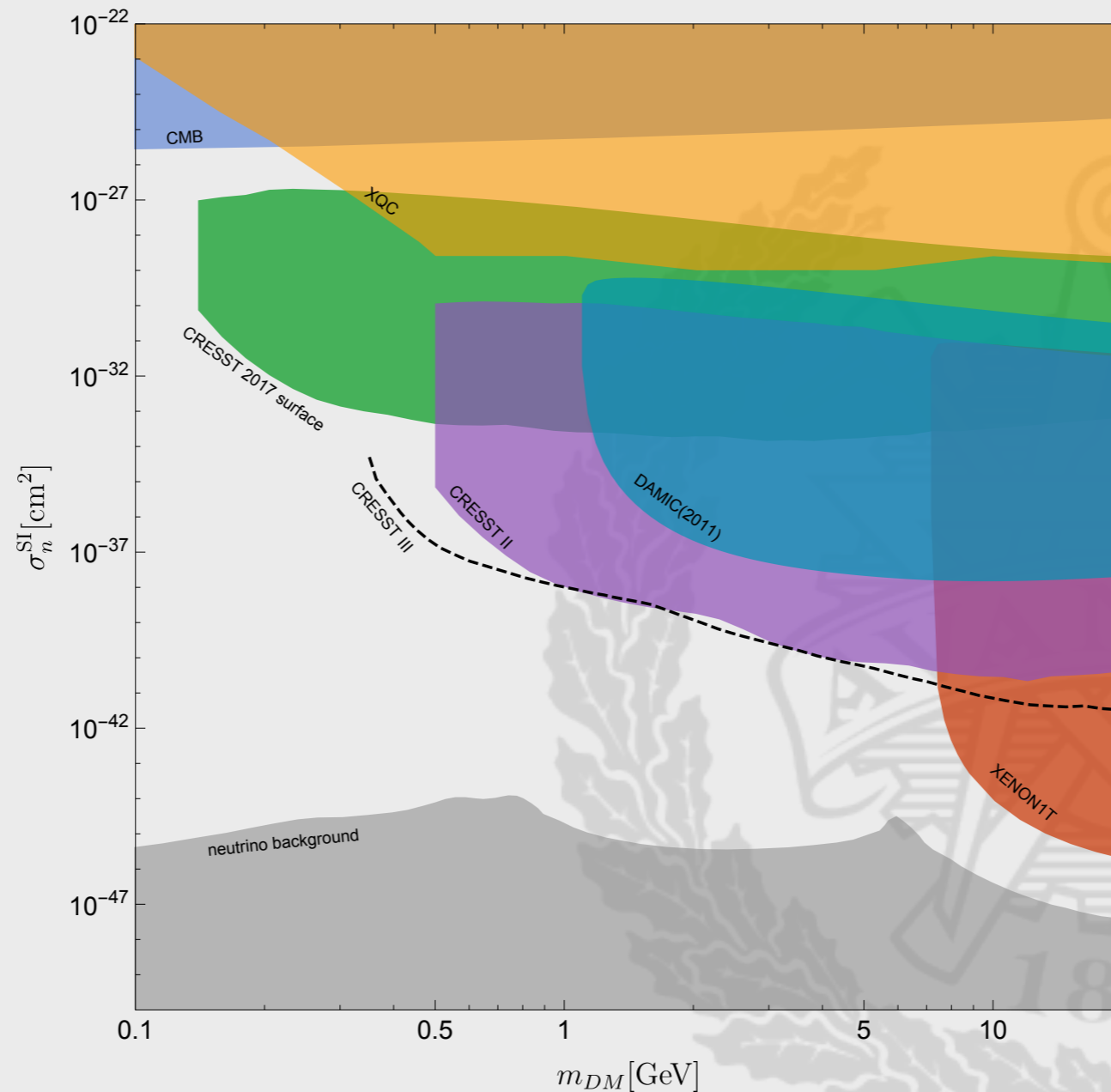


# Shielding **vs.** Detection





# Constraints on the DM-nucleon scattering cross-section



# Including DM-electron scatterings

The incoming DM flux gets attenuated by

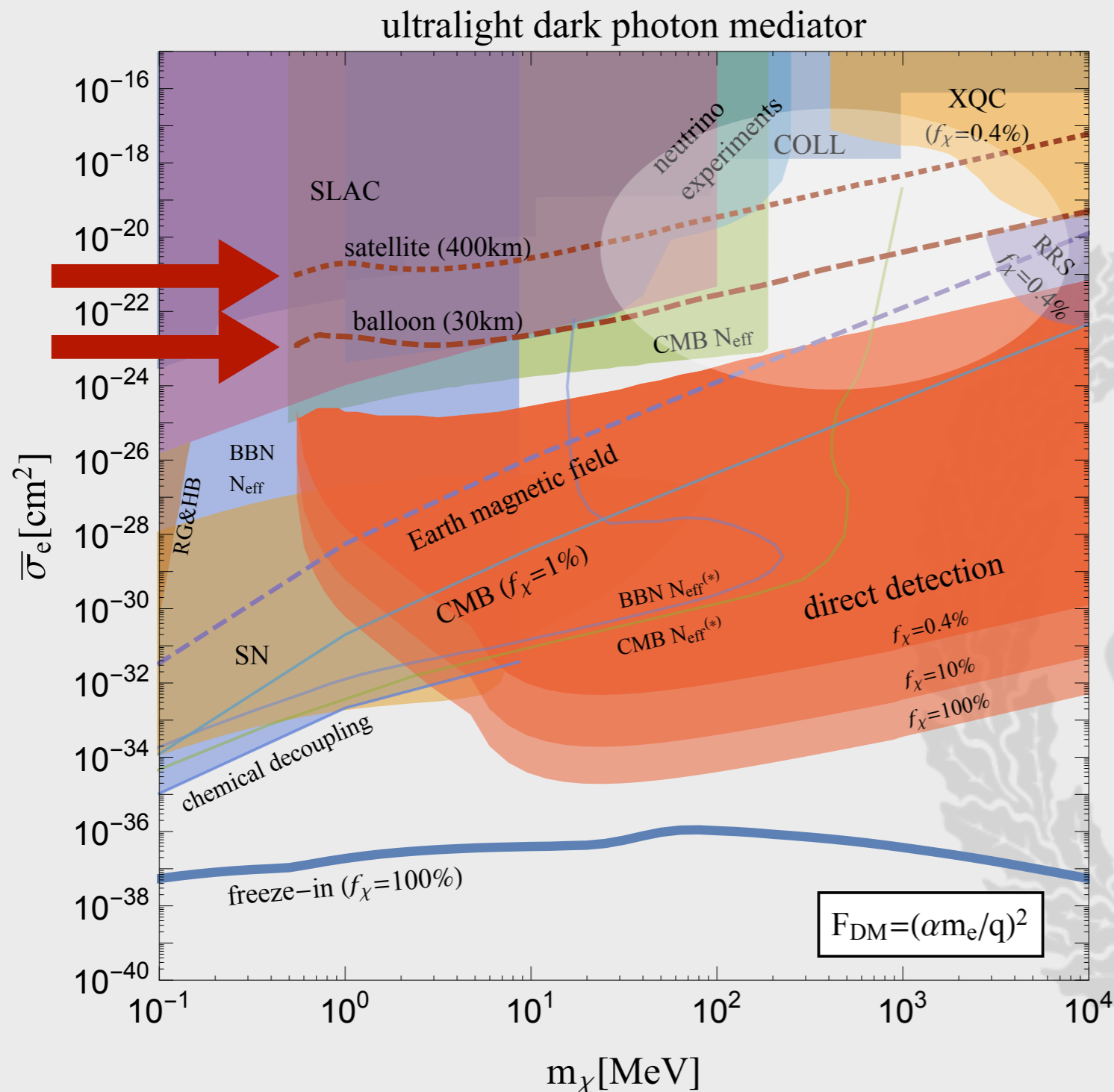
1. Elastic nuclear scatterings.
2. Elastic DM-electron scatterings.
3. Inelastic DM-electron scatterings (ionizations/excitations).

detection process  $\neq$  attenuation/stopping process



We need a model.

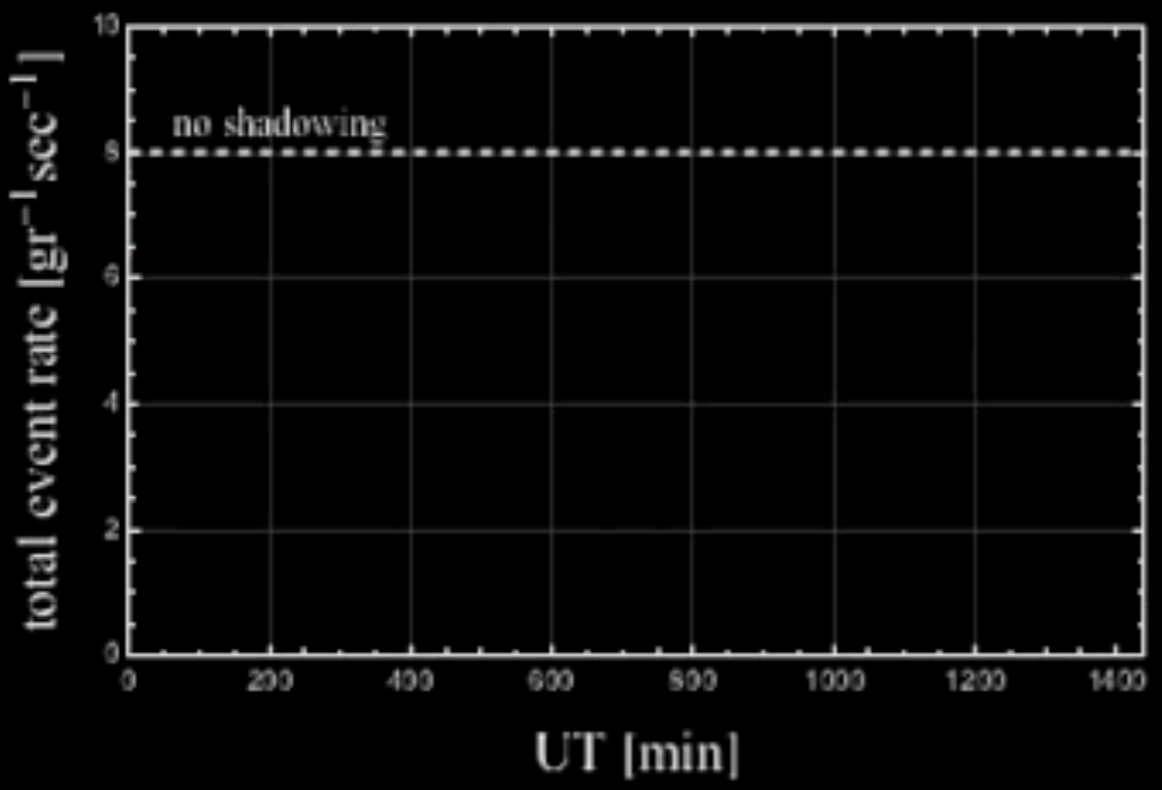
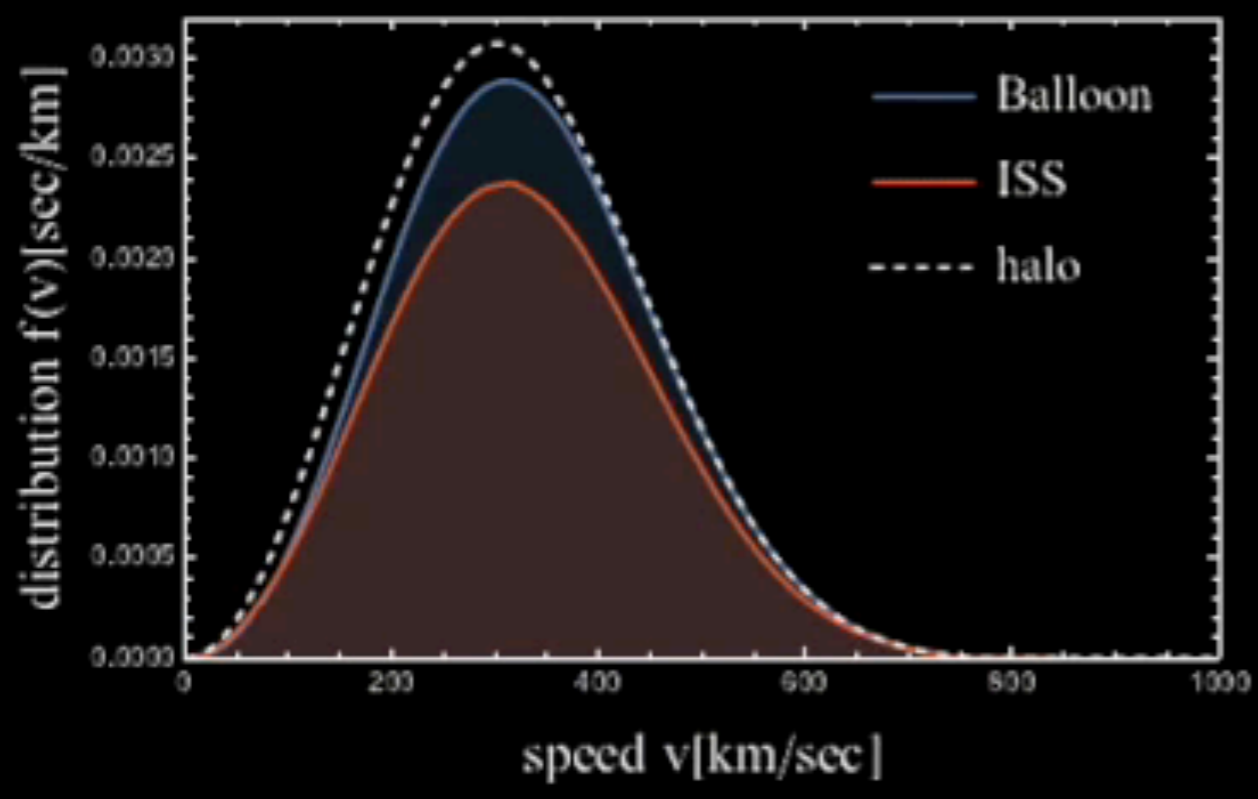
# Constraints on a sub-dominant component of strongly interacting DM



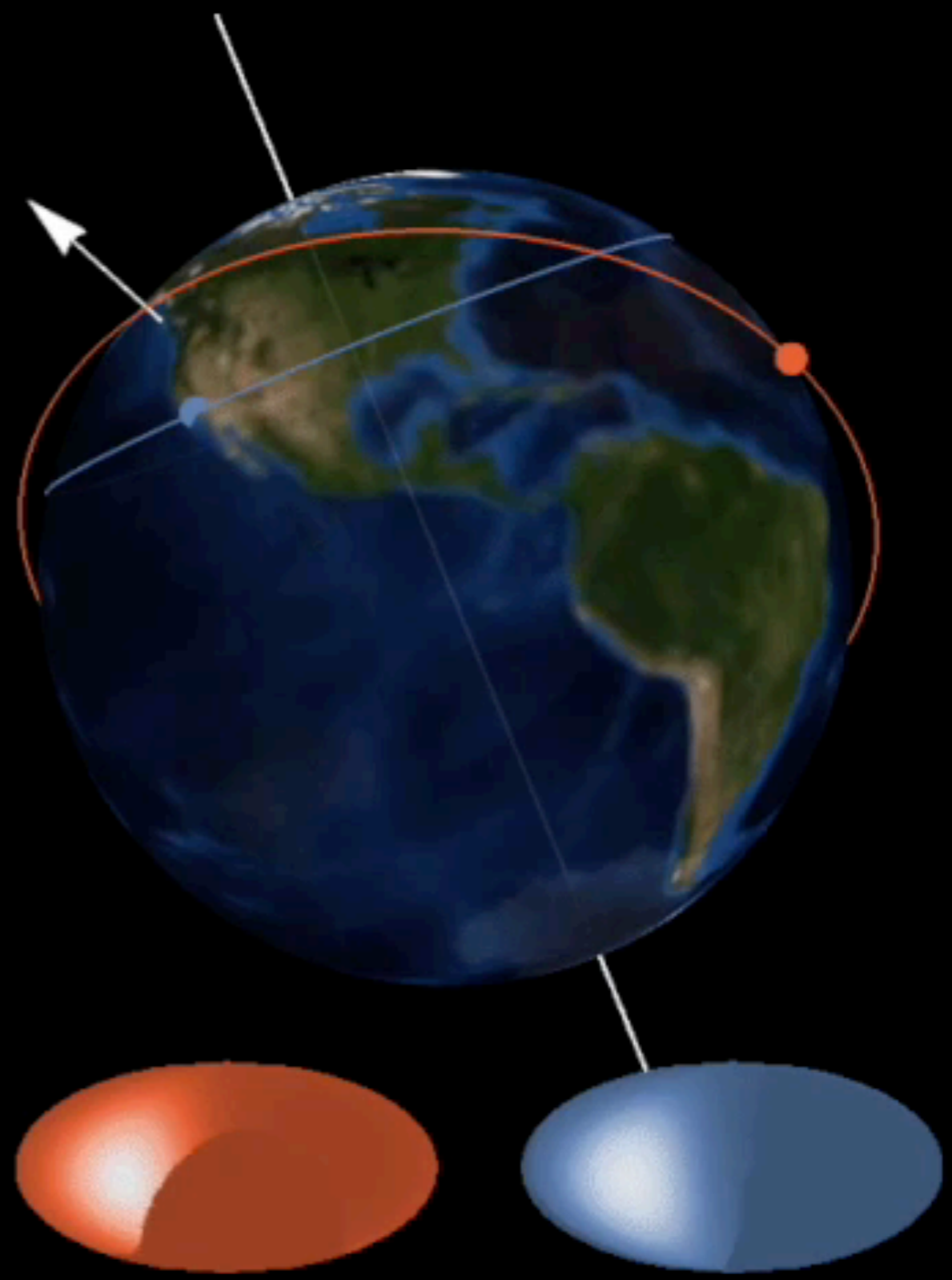
- So is there an open window in parameters space?
- Probably not for milli-charged DM.
- Definitely not for
  - $f_\chi = 100\%$
- Yes, under certain conditions:
  - Sub-dominant component.
    - $f_\chi < 0.4\%$
  - Ultralight, but not massless mediator.
  - Small dark gauge coupling.

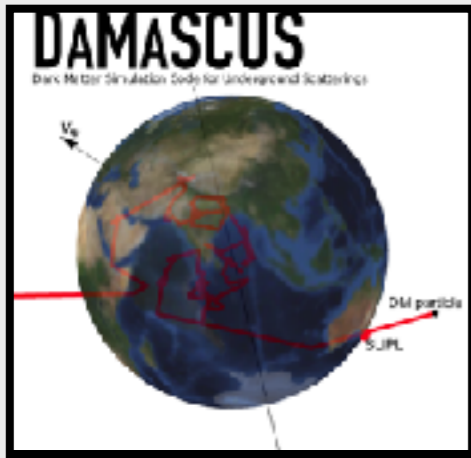
# Balloon/satellite experiments

Emken 2018



00:00

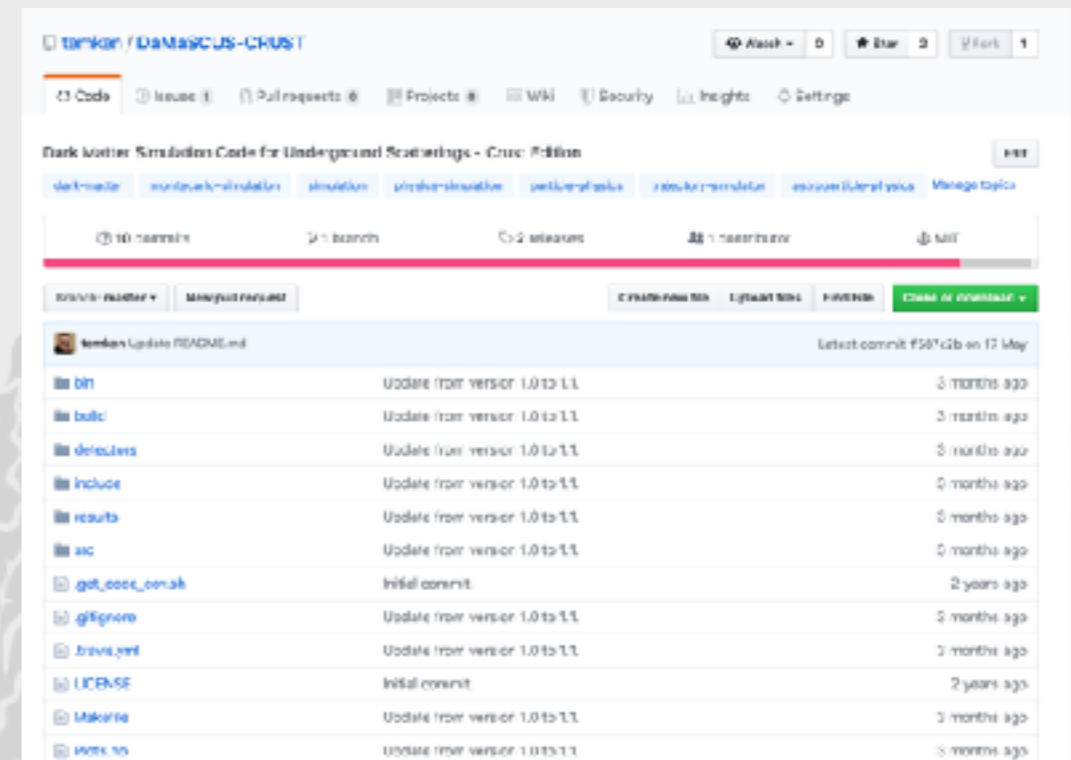




# The DaMaSCUS code

Dark Matter Simulation Code for Underground Scatterings

- Written in C++.
- Fully parallelized with *MPI*.
- Results were generated on the ABACUS2.0 supercomputer.
- The code is public.



<http://github.com/temken/>



**Thank you!**



# Backup Slides





## Rare event simulation I

# Importance Sampling

M.S. Mahdawi, G.R. Farrar, JCAP 1712 (2017) 004

- Rare event technique, which modifies the PDFs of the simulation.

$$\langle Y \rangle_I = \int_I dx Y(x) f(x) = \int_I dx Y(x) \frac{f(x)}{\hat{g}(x)} \hat{g}(x)$$

- Try to “mimic” the successful runs by introducing a bias into the simulations.
- Compensate by a weight factor.

$$f_\lambda(x) = \frac{1}{\lambda} \exp\left(-\frac{x}{\lambda}\right)$$
$$g_\lambda(x) = \frac{1}{(1 + \delta_\lambda)\lambda} \exp\left(-\frac{x}{(1 + \delta_\lambda)\lambda}\right)$$

$$f_\theta(\cos \theta) = \frac{1}{2}$$
$$g_\theta(\cos \theta) = \frac{1 + \delta_\theta \cos \theta}{2}$$

Statistical weight:  $w_{\lambda,i} = \frac{f_\lambda(l_i)}{g_\lambda(l_i)}$

# Rare event simulation II

## Geometric Importance Splitting

- “More interesting” particles get split into copies.
- Requires the definition of an importance function,

$$I : \mathbb{R}^3 \rightarrow \mathbb{R}$$

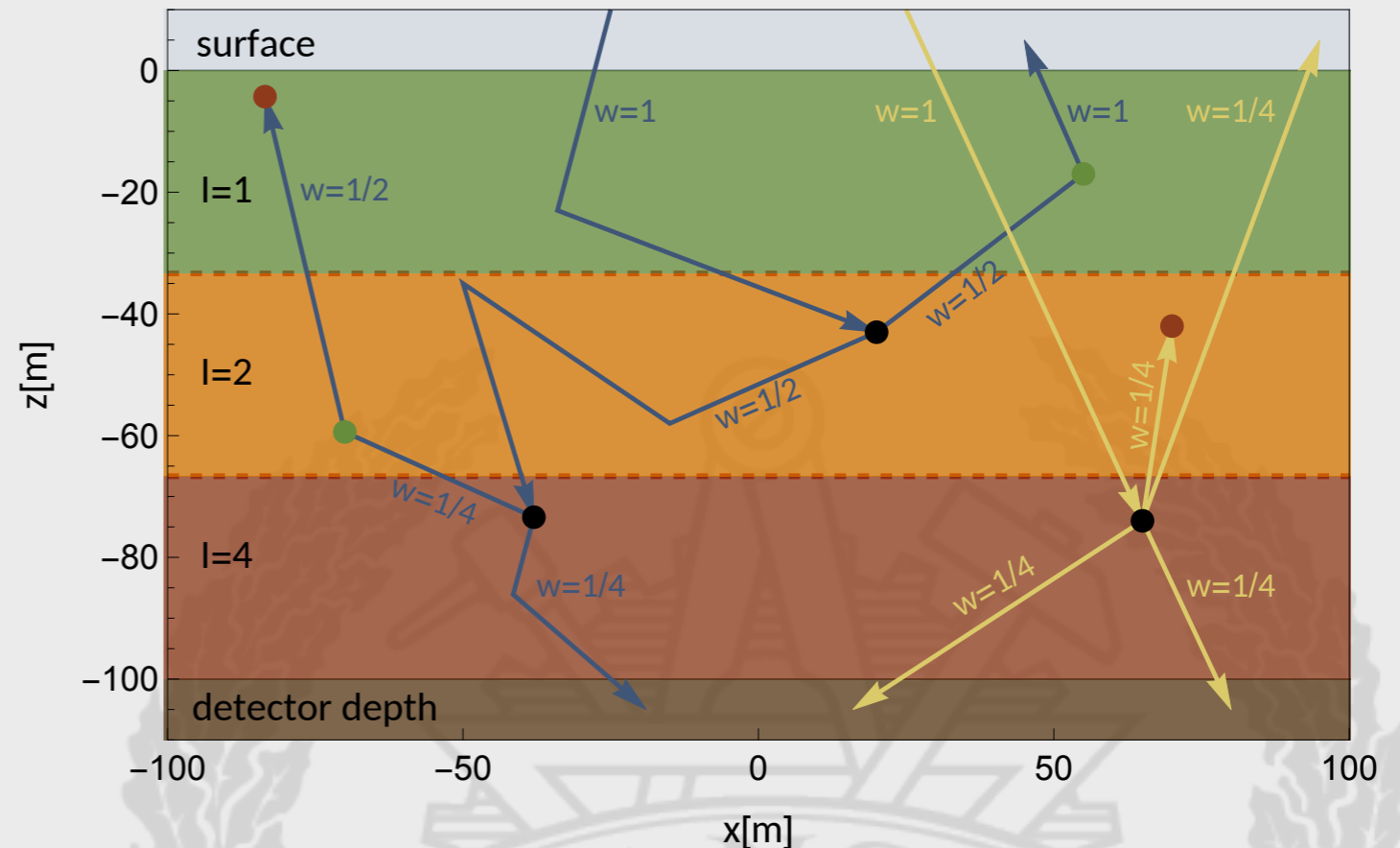
- If the importance increases,

$$\nu \equiv \frac{I_{i+1}}{I_i} > 1$$

- the particle gets split into

$$n = \begin{cases} \nu, & \text{if } \nu \in \mathbb{N}, \\ \lfloor \nu \rfloor, & \text{if } \nu \notin \mathbb{N} \wedge \xi \geq \Delta, \\ \lfloor \nu \rfloor + 1, & \text{if } \nu \notin \mathbb{N} \wedge \xi < \Delta, \end{cases}$$

copies.



● splitting ● Russian Roulette survival ● Russian Roulette kill

New statistical weight

$$w_{i+1} \equiv \frac{w_i}{n}$$

- Otherwise: **Russian Roulette**

# Including DM-electron scatterings

## The Dark Photon Model

- Extend the SM by a DM particle and a U(1) gauge group with kinetic mixing.

$$\mathcal{L}_D = \bar{\chi}(i\gamma^\mu D_\mu - m_\chi)\chi + \frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} + m_{A'}^2 A'_\mu A'^\mu + \epsilon F_{\mu\nu}F'^{\mu\nu}$$

- For kinetic mixing with the photon, the DM couples to electric charge.

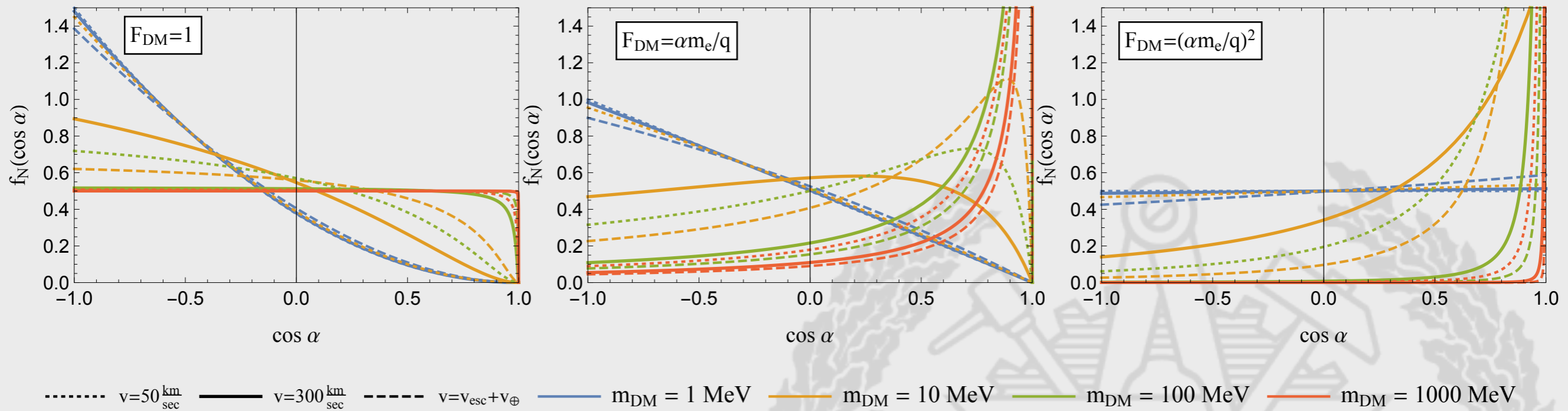
$$\frac{d\sigma_N}{dq^2} = \frac{\sigma_p}{4\mu_{\chi p}^2 v_\chi^2} F_{\text{DM}}(q)^2 F_N(q)^2 Z^2$$

- Hierarchy between the DM-proton and DM-electron cross section:

$$\frac{\sigma_p}{\sigma_e} = \left( \frac{\mu_{\chi p}}{\mu_{\chi e}} \right)^2$$

S.K. Lee et al, PRD92 (2015) 083517

# New scattering kinematics



DM form factor

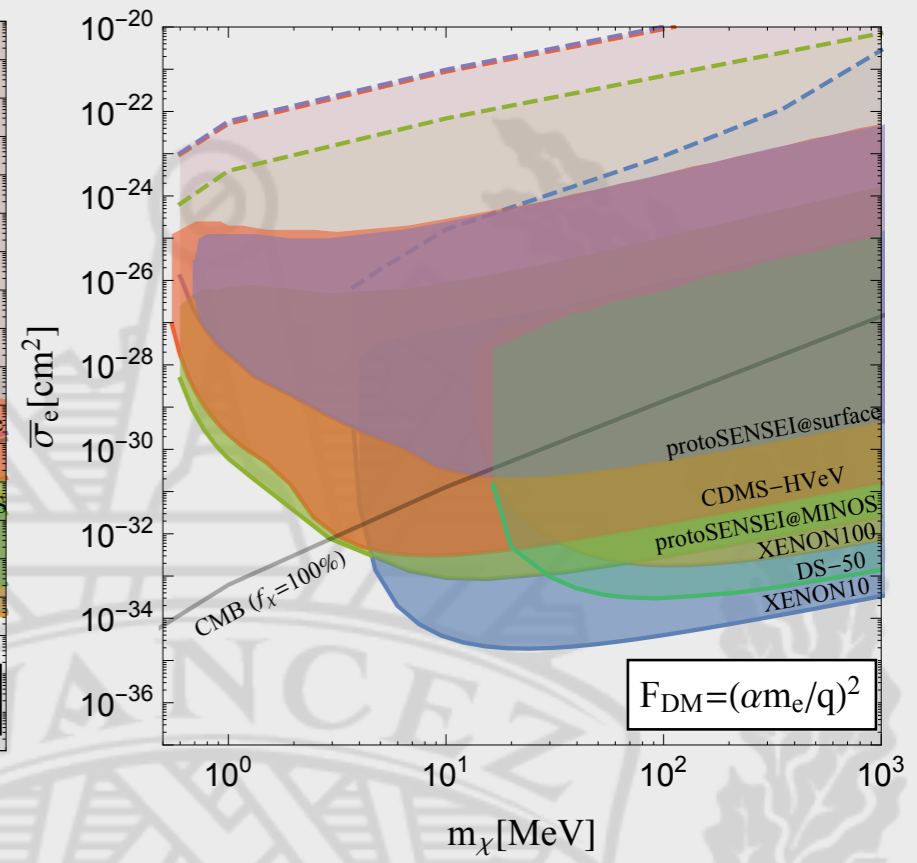
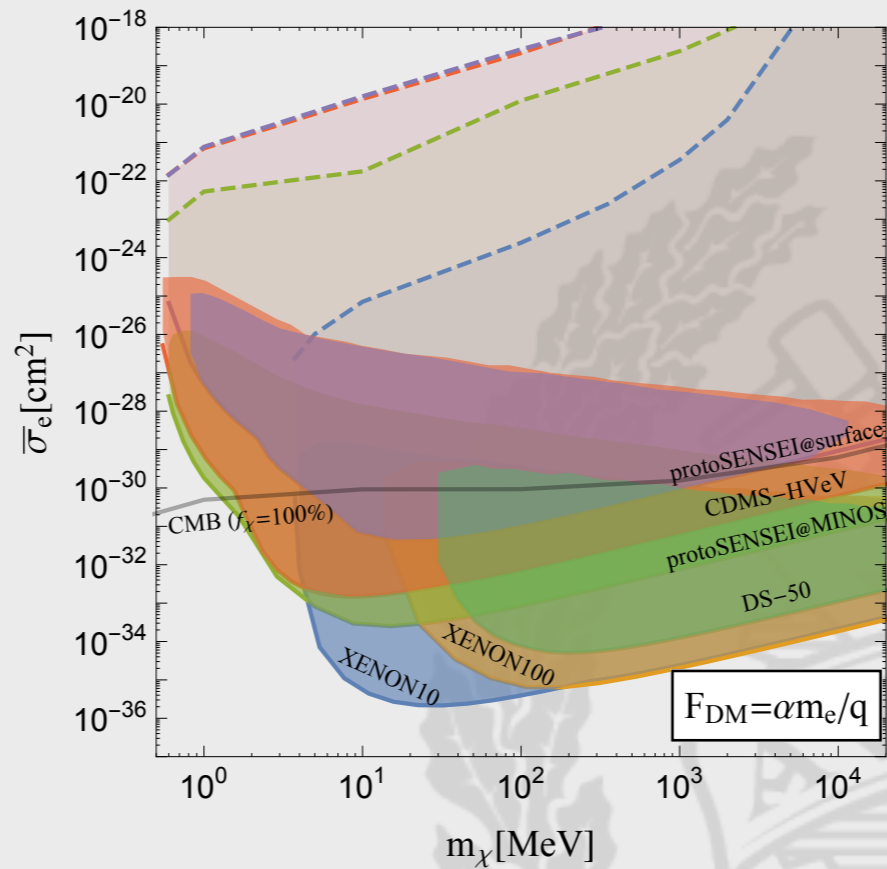
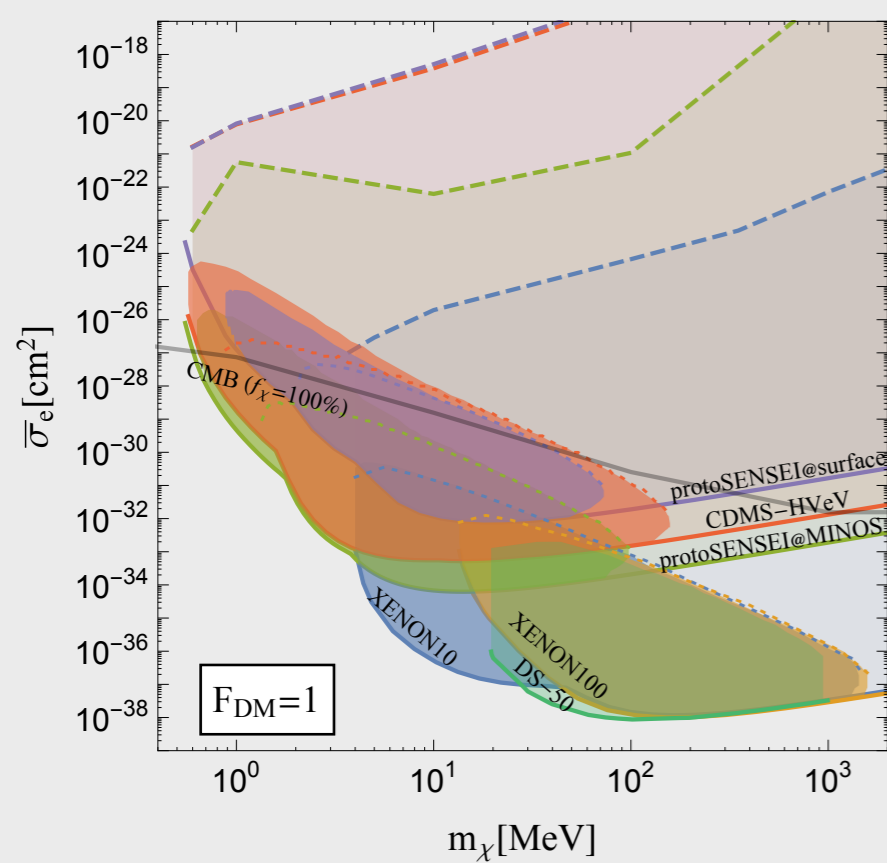
VS

Charge screening

$$F_{\text{DM}}(q) = \begin{cases} 1, & \text{for heavy mediator,} \\ \frac{q_{\text{ref}}}{q}, & \text{for ED interactions,} \\ \left(\frac{q_{\text{ref}}}{q}\right)^2, & \text{for light mediator.} \end{cases}$$

$$F_A(q) = \frac{a^2 q^2}{1 + a^2 q^2}$$

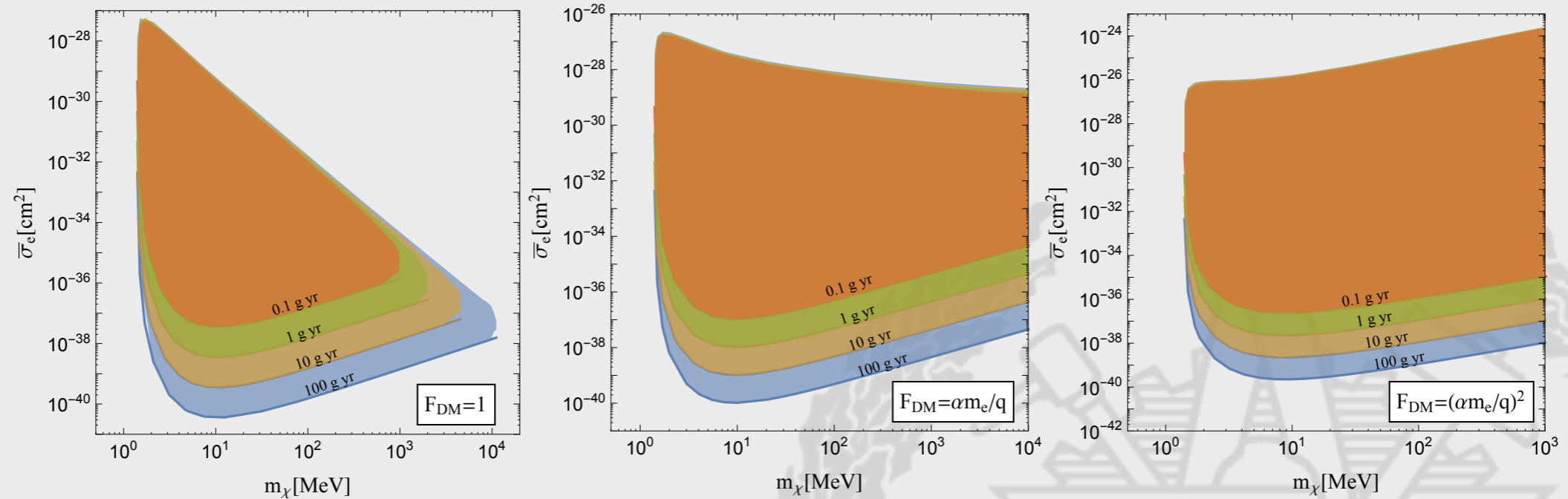
# DM-Electron Scattering constraints



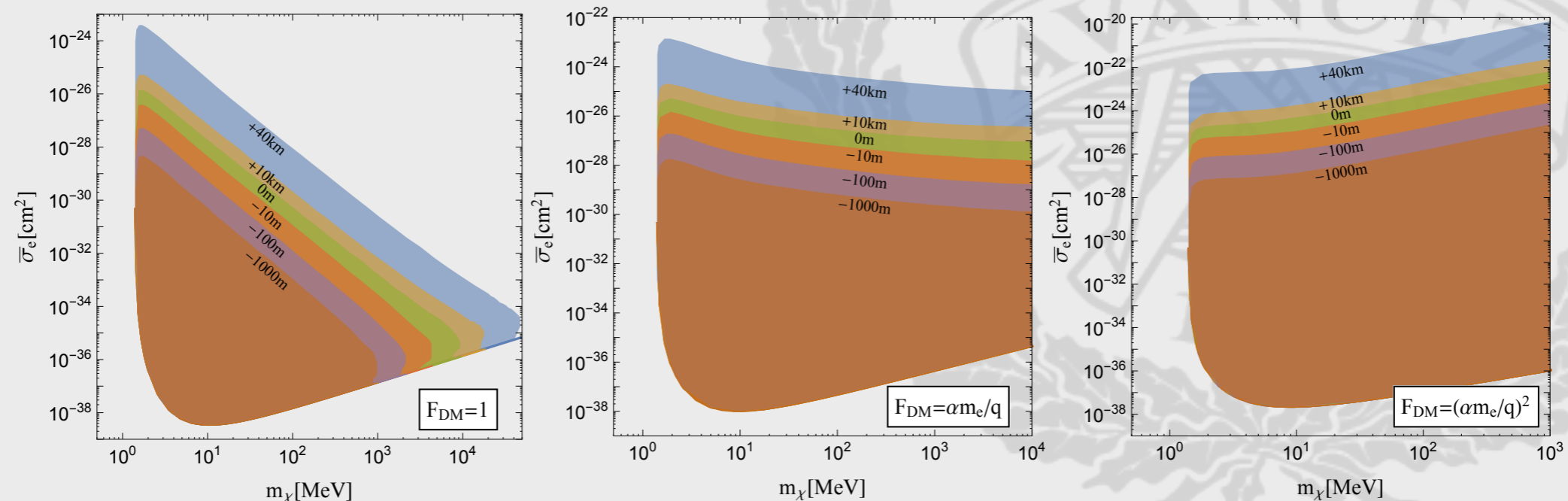


# How to push towards **stronger** interactions

- Higher Exposures



- Shallower Laboratories



# Kernel Density Estimation (KDE)

A non-parametric method to estimate an unknown PDF based on data.

For a data set  $\{x_1, x_2, \dots, x_N\}$  we can estimate the PDF via

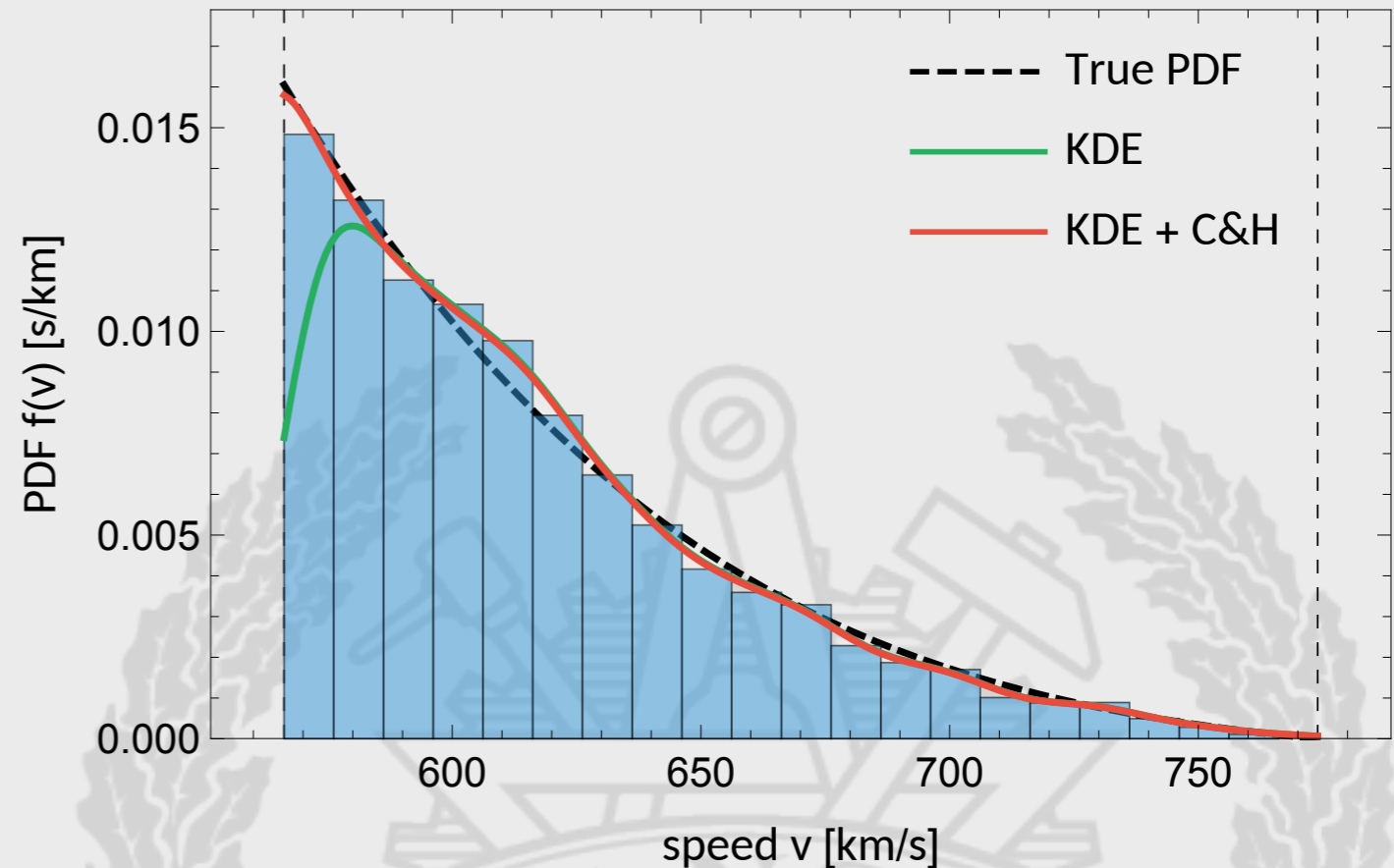
$$\hat{f}_h(x) = \frac{1}{h} \sum_{i=1}^N K\left(\frac{x - x_i}{h}\right).$$

E.g. with a Gaussian Kernel,

$$K(x) = \frac{1}{\sqrt{2\pi}} \exp\left(\frac{-x^2}{2}\right).$$

We set the bandwidth  $h$  using Silverman's rule of thumb,

$$h = \left(\frac{4}{3N}\right)^{1/5} \hat{\sigma}.$$



The bias at the domain's boundary has to be compensated, e.g. by a pseudo-data method by Cowling and Hall

R. Karunamuni, T. Alberts, *Statistical Methodology* 2 (2005), 191

A. Cowling, P. Hall, *Journal of the Royal Statistical Society*, B58 (1996), 551