

# Finite-size effects on the leading electromagnetic corrections to the hadronic vacuum polarisation

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# Overview

1 Why is the  $(g - 2)_\mu$  interesting?

2 Results and implications

3 Conclusions

# The anomalous magnetic moment of the muon

- Gyromagnetic ratio  $g_\mu$  defined through  $\mathbf{M} = g_\mu(e/2m_\mu)\mathbf{S}$
- Anomalous magnetic moment  $a_\mu$

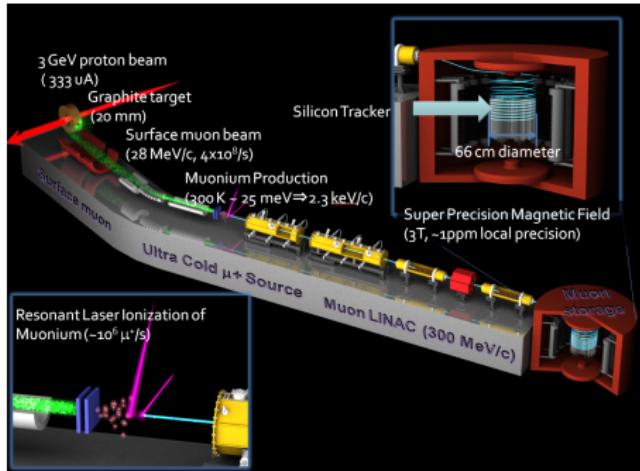
$$a_\mu = \frac{g_\mu - 2}{2}$$

- Deviation from Dirac prediction  $g_\mu = 2$
- Latest experimental value is (BNL)

$$a_\mu = 11\,659\,208.9(5.4)_{\text{stat}}(3.3)_{\text{sys}} \cdot 10^{-10} [0.54 \text{ ppm}]$$

- Roughly  $3\sigma$  deviation between experiment and theory
- Aim to improve error by factor of four to 0.14 ppm

# The anomalous magnetic moment of the muon



J-PARC

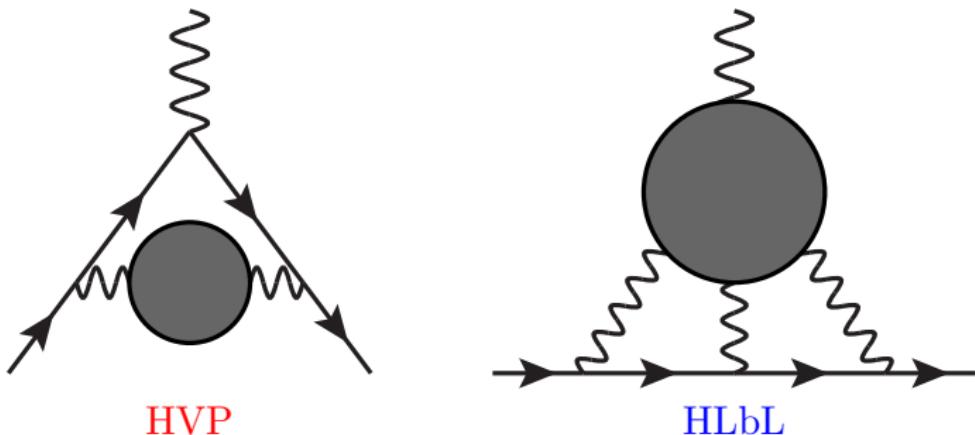


Fermilab

# On the theoretical side

$$a_\mu = a_\mu^{\text{QED}} + a_\mu^{\text{EW}} + a_\mu^{\text{had}},$$

$$a_\mu^{\text{had}} = \color{red}{a_\mu^{\text{HVP}}} + \color{blue}{a_\mu^{\text{HLbL}}}$$



- The hadronic vacuum polarisation (HVP) can be calculated in various ways

- Vector 2-point function

$$\Pi^{\mu\nu}(q^2) = i \int d^4x e^{iq\cdot x} \langle 0 | T(j^\mu(x) j^{\nu\dagger}(0)) | 0 \rangle$$

$$\Pi^{\mu\nu}(q^2) = (q^\mu q^\nu - q^2 g^{\mu\nu}) \Pi(q^2)$$

$$\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$$

- Dispersion relations:  $\text{Im } \hat{\Pi} \sim \sigma(e^+e^- \rightarrow \text{hadrons})$
- Lattice is competitive at the same accuracy

# Finite volume corrections

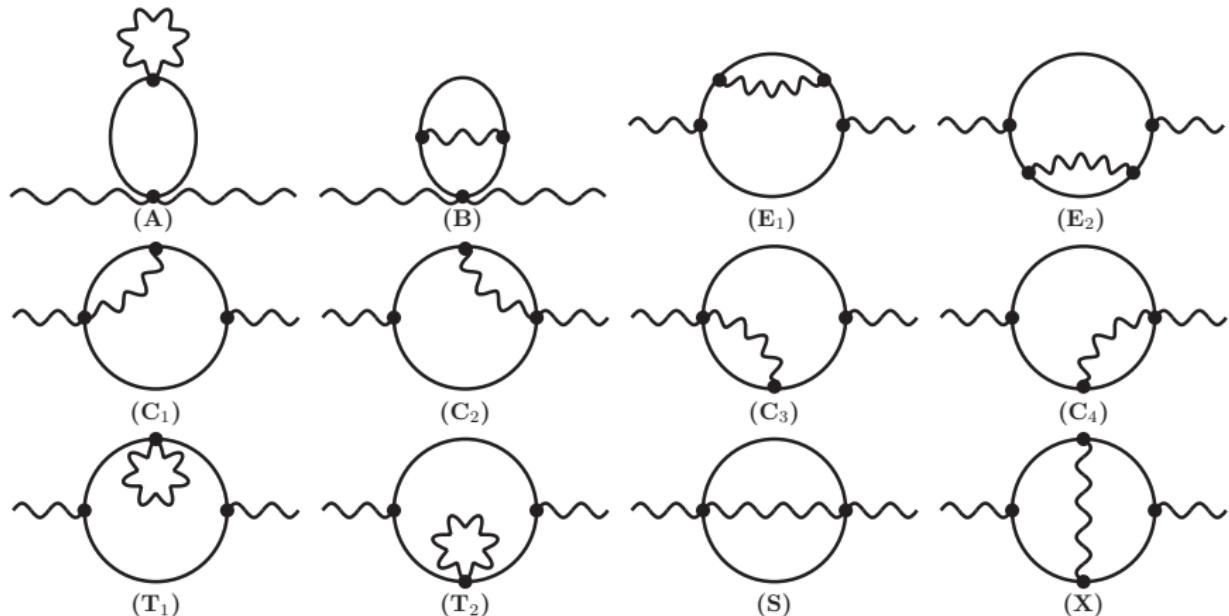
- Lattice is done in finite volume
- Produces finite volume effects depending on lattice size  $L$
- Depend on particle properties

Massive particles:  $e^{-mL}$

Massless particles:  $\frac{1}{L^a}$

- Electromagnetic corrections to the HVP potentially dangerous
- How big are the effects?  $a \geq 2$  but what is it really?

# Diagrams at NLO



- Scalar QED: Pure pion loops
- Compare with lattice simulations and vegas MC evaluation of loop integrals in lattice perturbation theory

# Actual calculation

- 2-loop integrals

$$\int \frac{d^4 k}{(2\pi)^4} \frac{d^4 \ell}{(2\pi)^4} \tilde{\pi}(q^2, k, \ell)$$

- Do energy integrals analytically
- On the lattice momenta are discretised:  $\mathbf{k} = (2\pi/L)\mathbf{n}$
- Taylor expand in  $1/L$
- Want the finite volume corrections, so for each diagram  $U$  we need to calculate

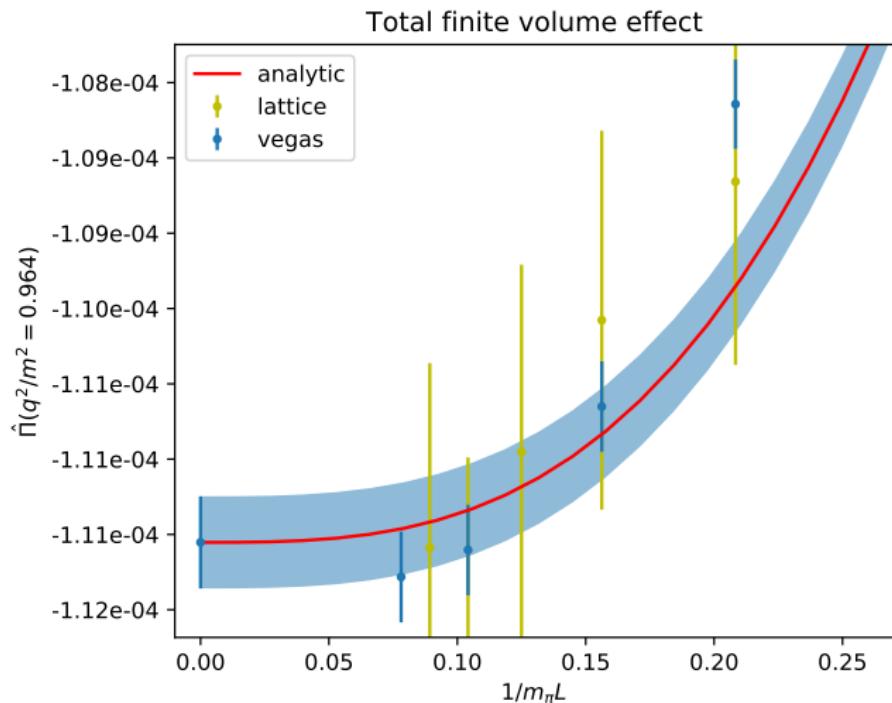
$$\Delta \hat{\Pi}_U(q^2) = \left( \sum_{\mathbf{n}}' - \int d^3 \mathbf{n} \right) \int d^3 \ell \hat{\pi}_U(q^2, \mathbf{n}) + \mathcal{O}\left(e^{-m_\pi L}\right)$$

# Finite volume corrections

$$\Delta \hat{\Pi}(q^2) = \frac{c_0}{m_\pi^3 L^3} \left( -\frac{16}{3} \Omega_{0,3} - \frac{5}{3} \Omega_{2,2} + \frac{40}{9} \Omega_{2,3} - \frac{3}{8} \Omega_{4,1} + \frac{7}{6} \Omega_{4,2} + \frac{8}{9} \Omega_{4,3} \right) + \mathcal{O}\left(\frac{1}{L^4}, e^{-m_\pi L}\right)$$

- Universality: adding form factors yields same result
- Leading order is  $1/L^3$ : Neutral current and photon far away sees no charge – dipole effect
- EM FV small for ordinary lattice sizes
- $c_0$  depends on lattice formulation of QED

# Diagrams

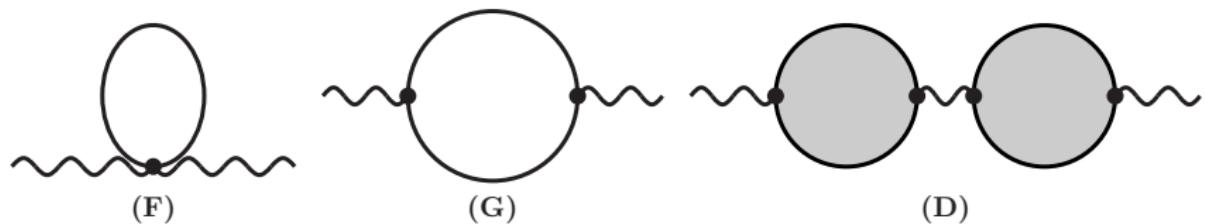


# Conclusions

- Finite volume corrections smaller than anticipated:  $c_0/(m_\pi L)^3$
- Need not worry about these corrections at normal lattice sizes
- Hadronic vacuum polarisation can be evaluated entirely on the lattice

# Back-up slides

# LO and NLO disconnected diagrams



# Charged currents

$$\begin{aligned}\Delta \hat{\Pi}_{\text{charged}}(q^2) = & \frac{1}{m^3 L^3} \left( -\frac{13}{24} \Omega_{2,2} + \frac{20}{9} \Omega_{2,3} - \frac{15}{64} \Omega_{4,1} + \frac{7}{24} \Omega_{4,2} + \frac{4}{9} \Omega_{4,3} \right) \\ & + \frac{c_1}{m^2 L^2 \pi} \left( -\frac{8}{3} \Omega_{-1,3} + \Omega_{1,2} + \frac{8}{3} \Omega_{1,3} + \frac{1}{8} \Omega_{3,1} \right) \\ & + \mathcal{O}\left(\frac{1}{L^4}, e^{-mL}\right)\end{aligned}$$

# Lattice numbers

- Time extent  $T = 128$
- Eight different volumes with  $L$  between 16 and 64
- $am_0 = 0.2$
- $q^2 = 0.964 m_0^2$
- 40 000 photon field configurations for  $L \leq 48$ , 10 000 for  $L = 56$  and  $L = 64$