

Cosmic chronometers to calibrate the ladders and measure the curvature of the Universe. A model-independent study

MNRAS 523 (2023) 3, 3406-3422

Arianna Favale
(arianna.favale@roma2.infn.it)

with Adrià Gómez-Valent and Marina Migliaccio

August 29th, 2023
Cosmology 2023 in Miramare, Trieste



SAPIENZA
UNIVERSITÀ DI ROMA



TOR VERGATA
UNIVERSITÀ DEGLI STUDI DI ROMA



Tensions in Cosmology

Cosmological and astrophysical data have recently been found to be in some tension (2σ or larger) with the **standard Λ CDM model**, as specified by the *Planck* 2018 parameter values

Among others, there are

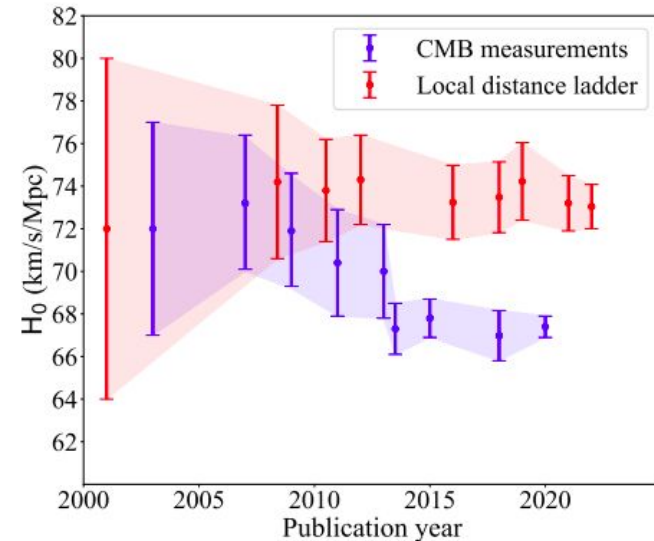
- *Hubble tension* ($\approx 5\sigma$) $\rightarrow H_0$
- *Growth tension* ($2-3\sigma$) $\rightarrow \sigma_8$, $S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$
- *CMB anisotropy anomalies* ($2-3\sigma$)

Tensions in Cosmology

Cosmological and astrophysical data have recently been found to be in some tension (2σ or larger) with the **standard Λ CDM model**, as specified by the *Planck* 2018 parameter values

Among others, there are

- **Hubble tension ($\approx 5\sigma$)** $\rightarrow H_0$ \rightarrow age and total energy content of the Universe
- Growth tension (2-3 σ) $\rightarrow \sigma_8$, $S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$
- CMB anisotropy anomalies (2-3 σ)



Tensions in Cosmology

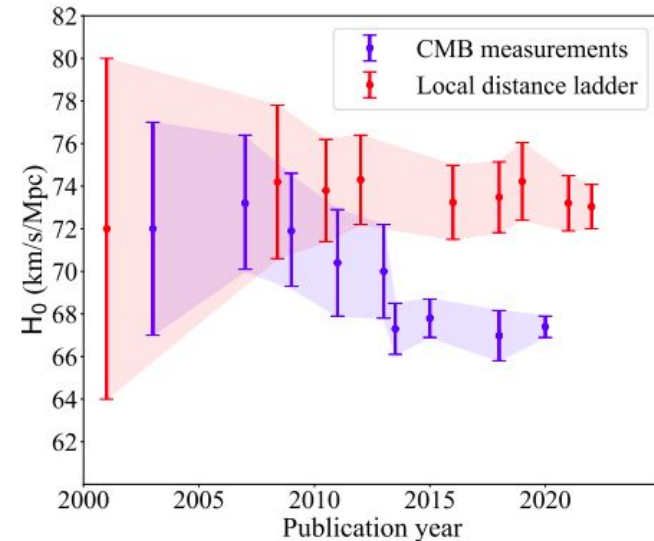
Cosmological and astrophysical data have recently been found to be in some tension (2σ or larger) with the **standard Λ CDM model**, as specified by the *Planck* 2018 parameter values

Among others, there are

- **Hubble tension ($\approx 5\sigma$)** $\rightarrow H_0$ \rightarrow age and total energy content of the Universe
- Growth tension (2-3 σ) $\rightarrow \sigma_8$, $S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$
- CMB anisotropy anomalies (2-3 σ)

Systematics in the data ?

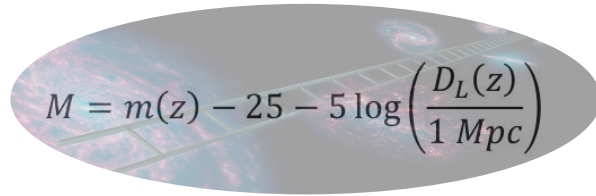
New physics beyond flat Λ CDM ?



The cosmic distance ladder

Local distance ladder

model-independent estimates of H_0 in the local Hubble flow ($z < 0.15$) through the calibration of the absolute magnitude, M , of standard candles such as Supernovae of type Ia (SNIa) at $z < 0.02$

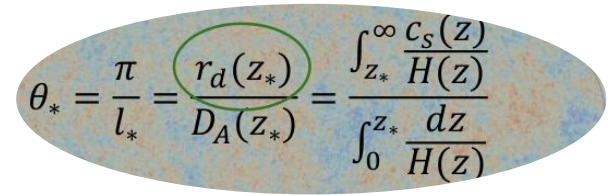


The equation $M = m(z) - 25 - 5 \log \left(\frac{D_L(z)}{1 \text{ Mpc}} \right)$ is displayed inside an oval with a background image of a galaxy cluster. The equation relates the absolute magnitude M to the apparent magnitude $m(z)$ and the luminosity distance $D_L(z)$ in megaparsecs.

$$M = m(z) - 25 - 5 \log \left(\frac{D_L(z)}{1 \text{ Mpc}} \right)$$

Inverse distance ladder

model-dependent estimates of H_0 : the standard ruler r_d , i.e. the sound horizon at the baryon-drag epoch, is used to calibrate cosmic distances and so the expansion rate $H(z)$



The equation $\theta_* = \frac{\pi}{l_*} = \frac{r_d(z_*)}{D_A(z_*)} = \frac{\int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz}{\int_0^{z_*} \frac{dz}{H(z)}}$ is displayed inside an oval with a background image of a galaxy cluster. The sound horizon $r_d(z_*)$ is circled in green. The equation relates the angular size θ_* of a standard ruler r_d at redshift z_* to the angular diameter distance D_A and the Hubble expansion rate $H(z)$.

$$\theta_* = \frac{\pi}{l_*} = \frac{r_d(z_*)}{D_A(z_*)} = \frac{\int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz}{\int_0^{z_*} \frac{dz}{H(z)}}$$



The **Hubble tension** can also be recast in a tension in the **calibrators of the local** and **inverse distance ladders**

A model-independent path to address the Hubble tension



We need **alternative ways** of calibrating the ladders

- Using **low-redshift** observations that are free from the main drivers of the current tensions: Cosmic Chronometers (**CCH**), **SN Ia** and Baryon Acoustic Oscillations (**BAO**)
- Performing **model-independent consistency tests** of low- z data sets
- Tool: Gaussian Processes → reconstruction of cosmological functions in an agnostic way

Low- z data sets for model-independent analysis

Cosmic chronometers (CCH)

32 direct measurements of $H(z)$ using massive and passively-evolving galaxies in $z = [0.07, 1.965]$

→ rely only on the Cosmological Principle (CP) and General Relativity

Supernovae of type Ia (SNIa)

Pantheon+ compilation: 1701 light curves of 1550 spectroscopically confirmed SNIa in $z = [0.001, 2.26]$ coming from 18 different surveys

→ We use only SNIa free from the local distance ladder calibration: 1624 data points

Baryon Acoustic Oscillations (BAO)

11 data points in $z = [0.12, 1.48]$ from galaxy surveys (e.g., WiggleZ, BOSS, DES Y3, 6dFGS+SDSS), given in terms of angular BAO and radial BAO

Gaussian Processes (GPs)



How do we reconstruct the trend of our cosmological quantities in a **model-independent** way?

Gaussian Processes (GPs)



How do we reconstruct the trend of our cosmological quantities in a **model-independent** way?

Gaussian Processes: a bayesian tool designed to perform **data-driven** reconstruction of smooth trends from Gaussian distributed data

$$f(x) \sim GP(\mu(x), D[K(x, \tilde{x}), C])$$

Advantages:

- No need for a model – **agnostic estimates** of cosmological quantities
- Performs reconstruction of functions **where we don't have data points**
- **Few assumptions:** Choice of a **kernel function**

A model-independent path to address the Hubble tension



Test some of the basic assumptions behind Λ CDM, which are usually taken for granted:

- 1) the constancy of the SNIa absolute magnitude and
- 2) the homogeneity property of the universe

A model-independent path to address the Hubble tension

Test some of the basic assumptions behind Λ CDM, which are usually taken for granted:

- 1) the constancy of the SNIa absolute magnitude and 2) the homogeneity property of the universe

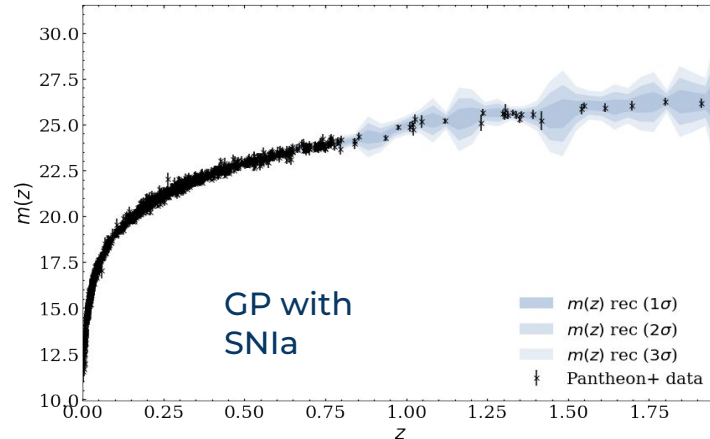
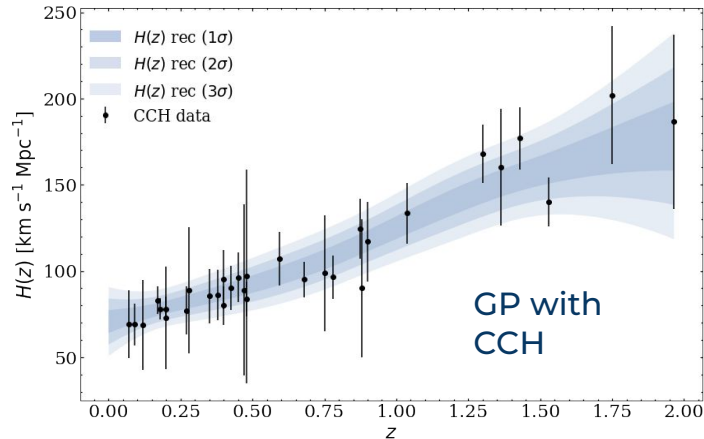


$$\Omega_k = 0$$

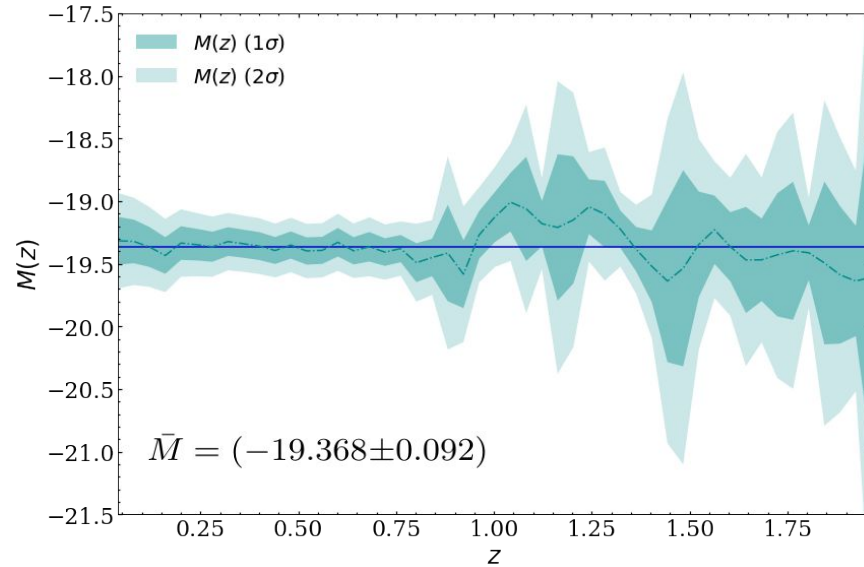
$$D_L(z) = c(1+z) \int_0^z \frac{dz'}{H(z')}$$



$$M(z) = m(z) - 25 - 5 \log \left(\frac{D_L(z)}{1 \text{ Mpc}} \right)$$



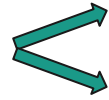
1) Testing the constancy of M



with the CCH and Pantheon+ data sets there is no significant statistical preference for the evolution of $M(z)$

2) Testing the Cosmological Principle

$$\Omega_k(z) \longrightarrow D_A(z) = \frac{D_L(z)}{(1+z)^2}$$



1 $D_L(z) = 10^{m(z) - M(z) - 25}$

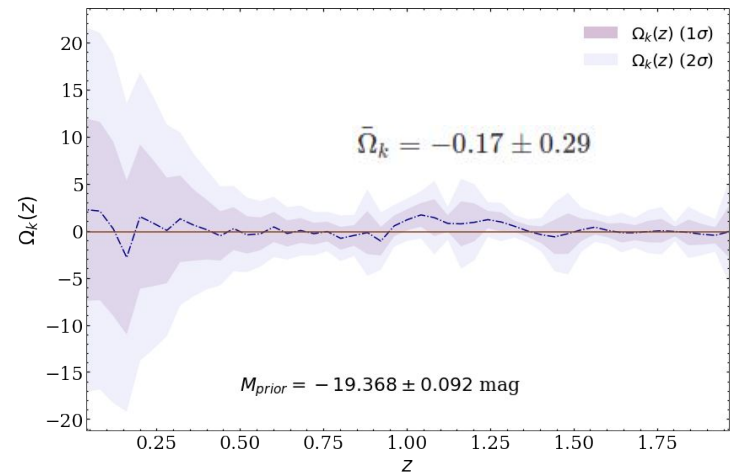
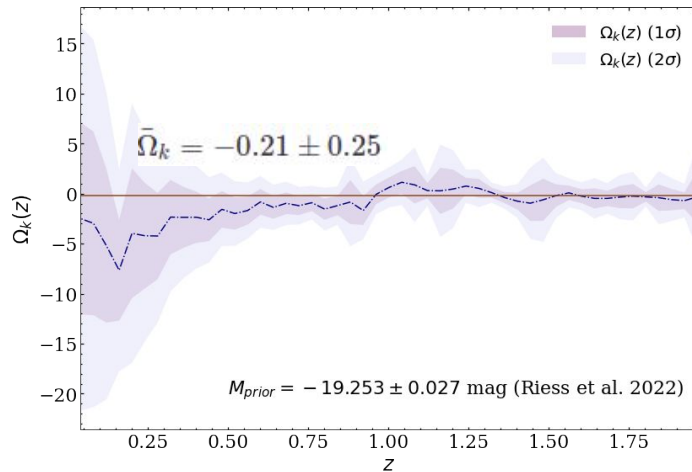
2 $D_L(z) = \frac{c(1+z)}{\sqrt{\Omega_k} H_0^2} \sinh \left(\sqrt{\Omega_k} H_0^2 \int_0^z \frac{dz'}{H(z')} \right)$

2) Testing the Cosmological Principle

$$\Omega_k(z) \longrightarrow D_A(z) = \frac{D_L(z)}{(1+z)^2}$$

1 $D_L(z) = 10^{m(z) - M(z) - 25}$

2 $D_L(z) = \frac{c(1+z)}{\sqrt{\Omega_k} H_0^2} \sinh \left(\sqrt{\Omega_k} H_0^2 \int_0^z \frac{dz'}{H(z')} \right)$



no hint for a violation of the CP from the CCH and Pantheon+ data sets

Consistency test for the BAO data

Angular and radial data: $r_d H(z)$, $D_V(z)/r_d$, $D_M(z)/r_d$

$$\Omega_k = 0$$

$$D_L(z) = c(1+z) \int_0^z \frac{dz'}{H(z')} \leftarrow \text{GP with CCH}$$

$$\left\{ \begin{array}{l} D_A(z) = D_L(z)/(1+z)^2 \\ D_M(z) = (1+z)D_A(z) \\ \frac{D_V(z)}{r_d} = \frac{1}{r_d} \left[D_M^2(z) \frac{cz}{H(z)} \right]^{\frac{1}{3}} \end{array} \right.$$

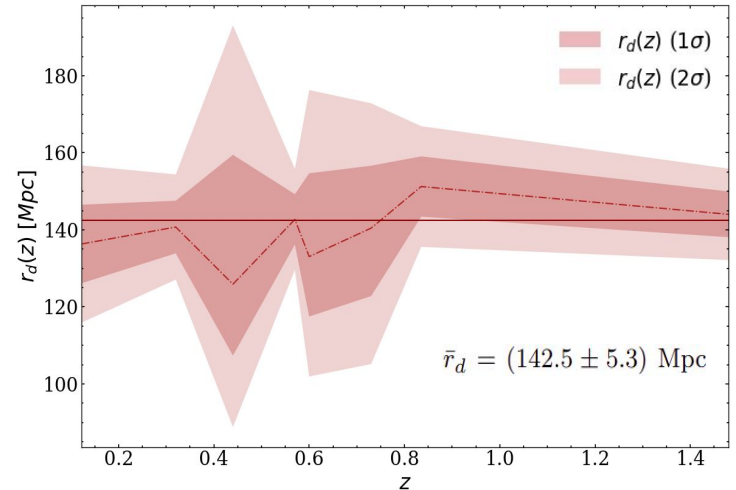
Consistency test for the BAO data

Angular and radial data: $r_d H(z)$, $D_V(z)/r_d$, $D_M(z)/r_d$

$$\Omega_k = 0$$

$$D_L(z) = c(1+z) \int_0^z \frac{dz'}{H(z')} \leftarrow \text{GP with CCH}$$

$$\left\{ \begin{array}{l} D_A(z) = D_L(z)/(1+z)^2 \\ D_M(z) = (1+z)D_A(z) \\ \frac{D_V(z)}{r_d} = \frac{1}{r_d} \left[D_M^2(z) \frac{cz}{H(z)} \right]^{\frac{1}{3}} \end{array} \right.$$

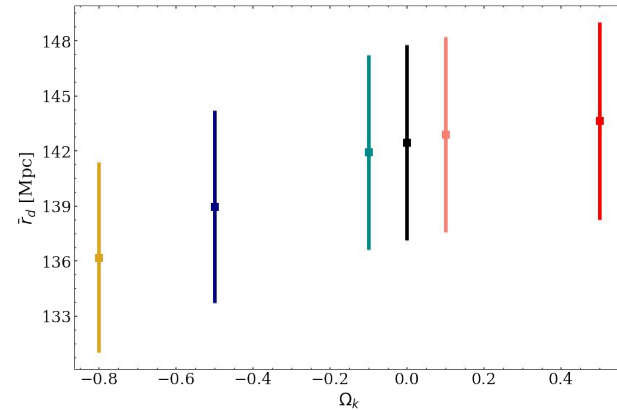
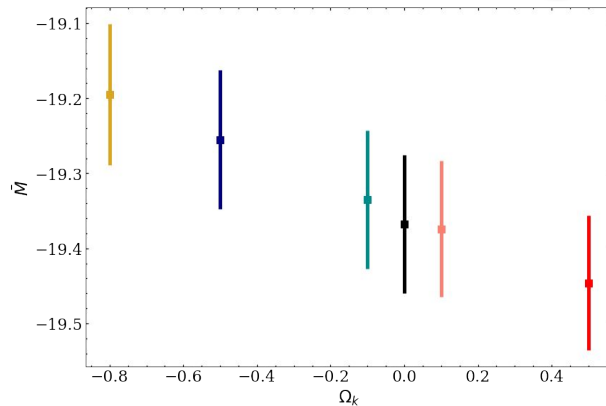


the values of r_d at different redshifts are statistically consistent with each other

Prior dependence

All the **central values** that we find in these analyses depend on the **subjective choice of the priors**:

$$D_L(z) = \frac{c(1+z)}{\sqrt{\Omega_k} H_0} \sinh \left(\sqrt{\Omega_k} H_0 \int_0^z \frac{dz'}{H(z')} \right)$$



➔ A more robust analysis has to be carried out in a **multi-dimensional** parameter space to **jointly** constrain M, Ω_k, r_d

Joint analysis

We combine the CCH, SNIa and BAO data sets to obtain joint constraints in the parameter space (M, Ω_k, r_d) through a grid-search method based on a χ^2 statistics

$$\chi_{\mu,i}^2 = \sum_{k,l=1}^{1624} [m(z_k) - m_{rec,\mu,i}(z_k)] C_{kl}^{-1} [m(z_l) - m_{rec,\mu,i}(z_l)] \quad (\text{CCH+SNIa})$$

$$\chi_{\mu,i}^2 = \sum_{n,j=1}^{11} [x(z_n) - x_{rec,\mu,i}(z_n)] C_{nj}^{-1} [x(z_j) - x_{rec,\mu,i}(z_j)] \quad (\text{CCH+BAO})$$

By combining the 2D analyses we get: $\chi^2(M, \Omega_k, r_d) = \chi^2(M, \Omega_k) + \chi^2(\Omega_k, r_d)$ (CCH+SNIa+BAO)

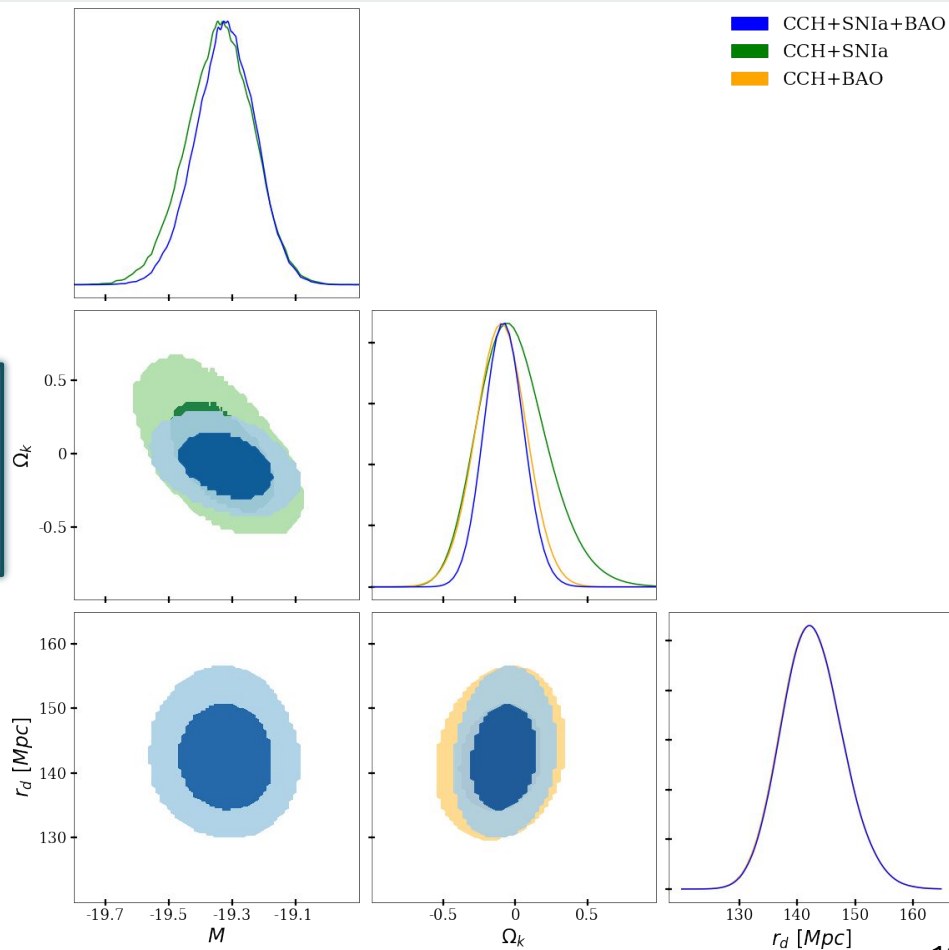
and we associate a weight to each grid point: $w_\mu \propto B_\mu \exp(-\bar{\chi}_\mu^2/2) \underbrace{\sum_{i=1}^N \exp(-[\chi_{\mu,i}^2 - \bar{\chi}_\mu^2]/2)}_{\equiv f_\mu}$

$\longrightarrow \chi_{\mu,eff}^2 = \bar{\chi}_\mu^2 - 2 \ln(B_\mu f_\mu)$

Joint analysis - Results

■ CCH+SN Ia+BAO
■ CCH+SN Ia
■ CCH+BAO

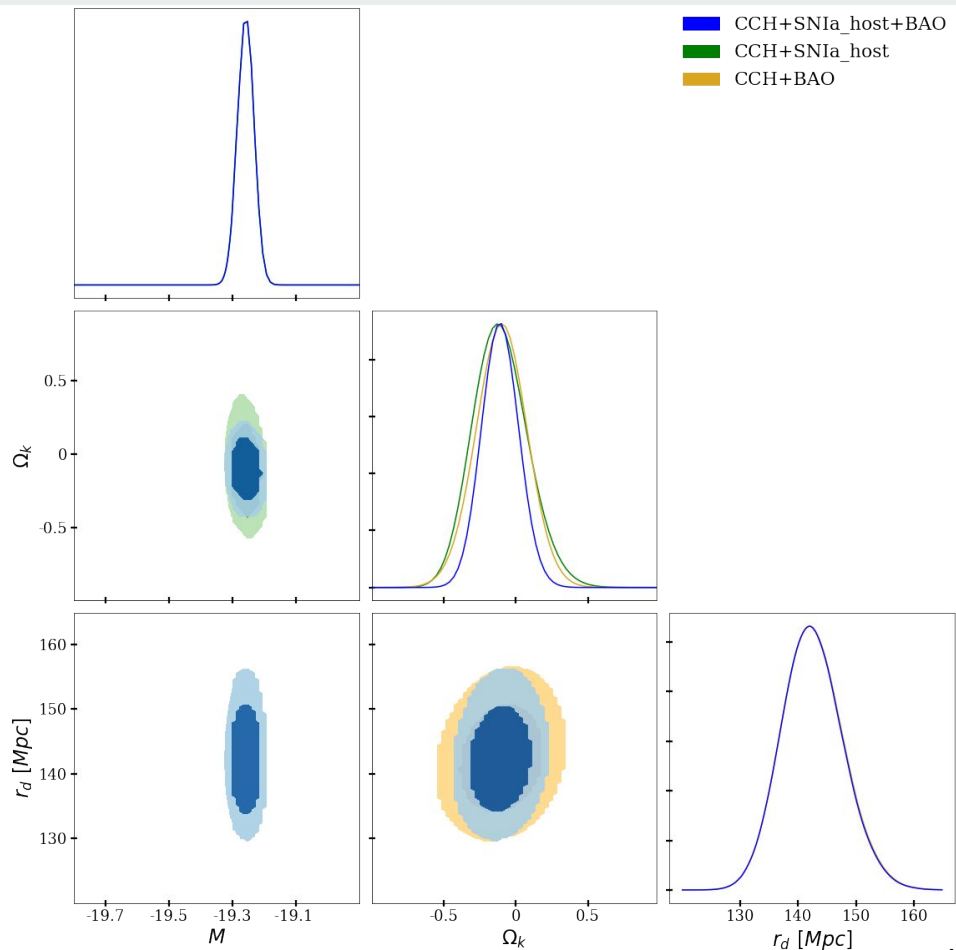
	CCH+SN Ia	CCH+BAO	CCH+SN Ia+BAO
M [mag]	$-19.344^{+0.116}_{-0.090}$		$-19.314^{+0.086}_{-0.108}$
Ω_k	$-0.07^{+0.27}_{-0.21}$	-0.10 ± 0.18	$-0.07^{+0.12}_{-0.15}$
r_d [Mpc]		$141.9^{+5.6}_{-4.9}$	142.3 ± 5.3



Joint analysis - Results

Including the SNIa located in the **Cepheid host galaxies** employed by **SHoES** to calibrate the SNIa in the second rung of the cosmic distance ladder has a little impact on the results:

CCH+SNIa_host+BAO	
M [mag]	$-19.252^{+0.024}_{-0.036}$
Ω_k	$-0.10^{+0.12}_{-0.15}$
r_d [Mpc]	$141.9^{+5.6}_{-4.9}$

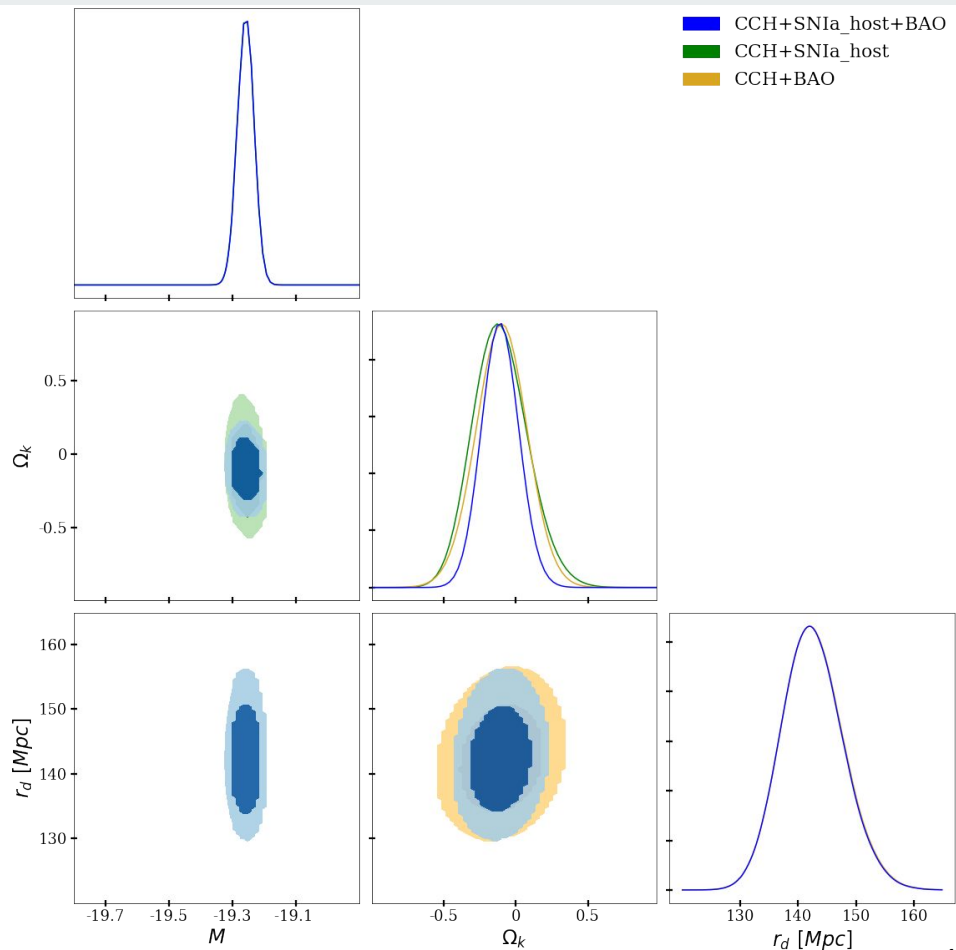


Joint analysis - Results

Including the SNIa located in the **Cepheid host galaxies** employed by **SHoES** to calibrate the SNIa in the second rung of the cosmic distance ladder has a little impact on the results:

	CCH+SNIa	host+BAO
M [mag]	$-19.252^{+0.024}_{-0.036}$	
Ω_k	$-0.10^{+0.12}_{-0.15}$	
r_d [Mpc]	$141.9^{+5.6}_{-4.9}$	

→ $M^{R22} = (-19.253 \pm 0.027)$ mag
(Riess+, 2022)



H_0 measurement

We make use of the cosmographical expansion

$$D_L(z) = \frac{cz}{H_0} \left[1 + \frac{z}{2} (1 - q_0) \right] + \mathcal{O}(z^3)$$

Employing as a prior our CCH+SN Ia+BAO constraint on M and the apparent magnitudes of the SN Ia in the Hubble flow ($0.023 < z < 0.15$)

$$M(z) = m(z) - 25 - 5 \log \left(\frac{D_L(z)}{1 \text{ Mpc}} \right)$$

We find: $H_0 = (71.5 \pm 3.1) \text{ km s}^{-1} \text{ Mpc}^{-1}$

H_0 measurement

We make use of the cosmographical expansion

$$D_L(z) = \frac{cz}{H_0} \left[1 + \frac{z}{2} (1 - q_0) \right] + \mathcal{O}(z^3)$$

Employing as a prior our CCH+SN Ia+BAO constraint on M and the apparent magnitudes of the SN Ia in the Hubble flow ($0.023 < z < 0.15$)

$$M(z) = m(z) - 25 - 5 \log \left(\frac{D_L(z)}{1 \text{ Mpc}} \right)$$

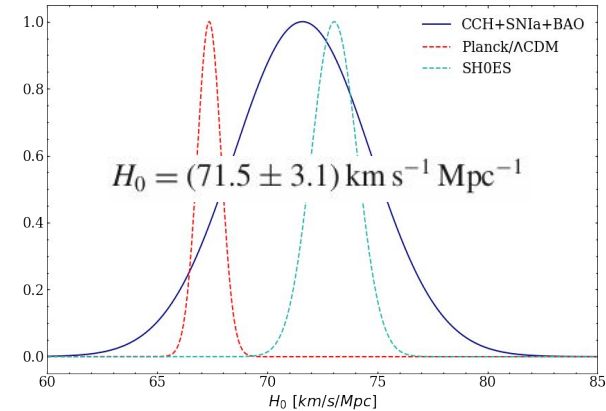
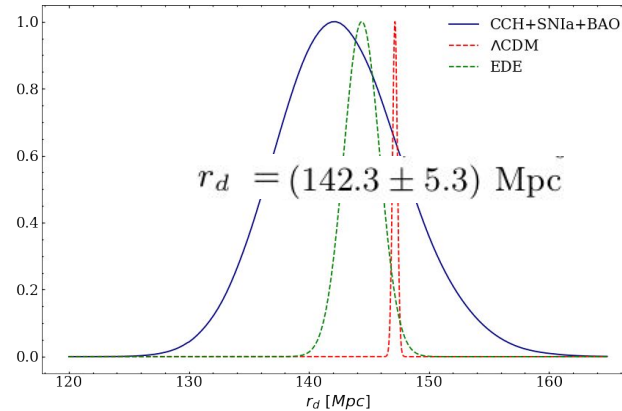
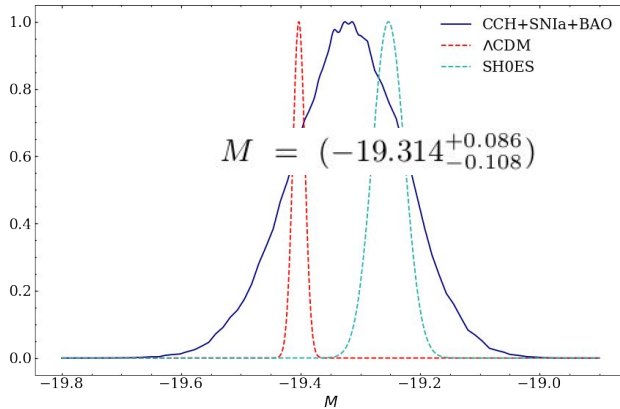
We find:

$$H_0 = (71.5 \pm 3.1) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Main results & Conclusions

- ★ The CCH and Pantheon+ data sets do not point to a time evolution of the SNIa intrinsic luminosity nor a breaking of the homogeneity of the universe at large scales. Both, M and Ω_k are compatible at 68% C.L. with a constant.
- ★ Under the precision of the CCH, we find that the BAO data are all compatible with a common value of r_d .
- ★ Motivated by these results, we have jointly constrained with CCH, SNIa and BAO the constant values of Ω_k , M and r_d by applying a model-independent method which is also independent from the first rungs of the cosmic distance ladder employed by SHoES and the CMB data from Planck, i.e. from the main data sets involved in the Hubble tension.

Main results & Conclusions



- ★ Nevertheless, the uncertainties we find are still too large to arbitrate the tension yet.
- ★ The proposed method will find interesting applications in the future with upcoming data (e.g., SNIa from the Vera C. Rubin Observatory's LSST and BAO from Euclid and DESI)
- ★ It will serve both as a discriminator of models beyond Λ CDM and an independent means of testing the calibration of the cosmic distance ladders.



Thank you for your attention!

Backup slides

About Gaussian Processes

$$\bar{\mathbf{f}}^* = \boldsymbol{\mu}^* + K(\mathbf{X}^*, \mathbf{X})[K(\mathbf{X}, \mathbf{X}) + C]^{-1}(\mathbf{y} - \boldsymbol{\mu})$$

$$\text{cov}(\mathbf{f}^*) = K(\mathbf{X}^*, \mathbf{X}^*) - K(\mathbf{X}^*, \mathbf{X})[K(\mathbf{X}, \mathbf{X}) + C]^{-1}K(\mathbf{X}, \mathbf{X}^*)$$

Kernel functions:



Hyperparameters:

$$K(x, \tilde{x}) = \sigma_f^2 \exp\left[-\frac{(x-\tilde{x})^2}{2l^2}\right]$$

- l the characteristic length scale of significant changes in $f(x)$
- σ_f the variance, amplitude of significant changes in $f(x)$

$$K(x, \tilde{x}) = \sigma_f^2 \exp\left[\left(-\frac{\sqrt{3}|x-\tilde{x}|}{l}\right)\left(1 + \frac{\sqrt{3}|x-\tilde{x}|}{l}\right)\right] \quad \text{Matérn 3/2 covariance function}$$

$$\ln \mathcal{L} = \ln p(\mathbf{y}|\mathbf{X}, \sigma_f, l) = -\frac{1}{2}(\mathbf{y} - \boldsymbol{\mu})^T [K(\mathbf{X}, \mathbf{X}) + C]^{-1}(\mathbf{y} - \boldsymbol{\mu}) - \frac{1}{2} \ln |K(\mathbf{X}, \mathbf{X}) + C| - \frac{n}{2} \ln 2\pi$$

Backup slides

Weighted mean

$$\mathcal{L}(M) = \mathcal{N} \exp \left[-\frac{1}{2} \sum_{i,j=1}^{n_p} (M - \bar{M}_i)(M - \bar{M}_j)(C^{-1})_{ij} \right]$$

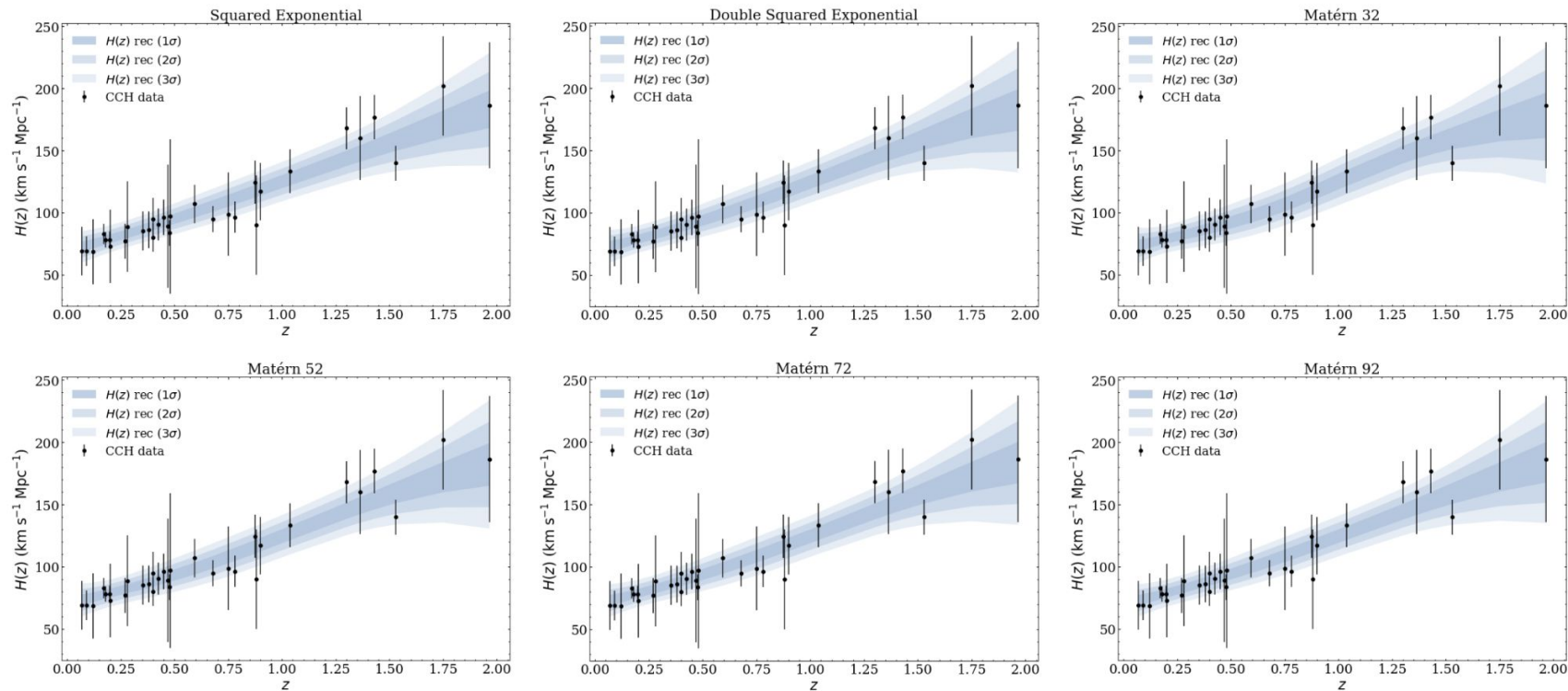
$$\mathcal{L}(M) = \tilde{\mathcal{N}} \exp \left[-\frac{1}{2} \left(\sum_{i,j=1}^{n_p} A_{ij} \right) \left(M - \frac{\sum_{i,j=1}^{n_p} \bar{M}_i A_{ij}}{\sum_{i,j=1}^{n_p} A_{ij}} \right)^2 \right]$$

$$A_{ij} \equiv (C^{-1})_{ij} \quad C_{ij} = \frac{1}{N_{real}} \sum_{\mu=1}^{N_{real}} (M_{\mu,i} - \bar{M}_i)(M_{\mu,j} - \bar{M}_j)$$

$$\bar{M} = \frac{\sum_{i,j=1}^{n_p} \bar{M}_i A_{ij}}{\sum_{i,j=1}^{n_p} A_{ij}} \quad ; \quad \sigma^2 = \frac{1}{\sum_{i,j=1}^{n_p} A_{ij}}$$

Backup slides

GP Kernels performance test on the reconstruction of the Hubble function with CCH

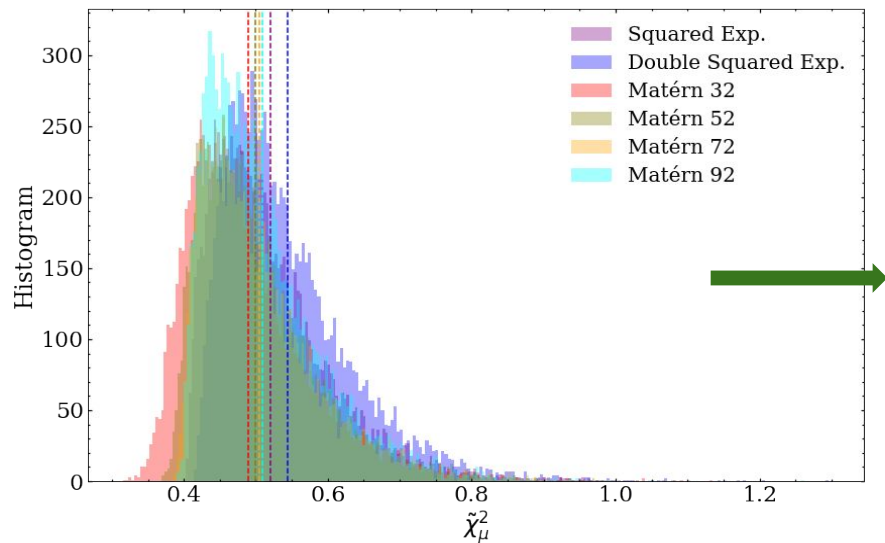


Backup slides

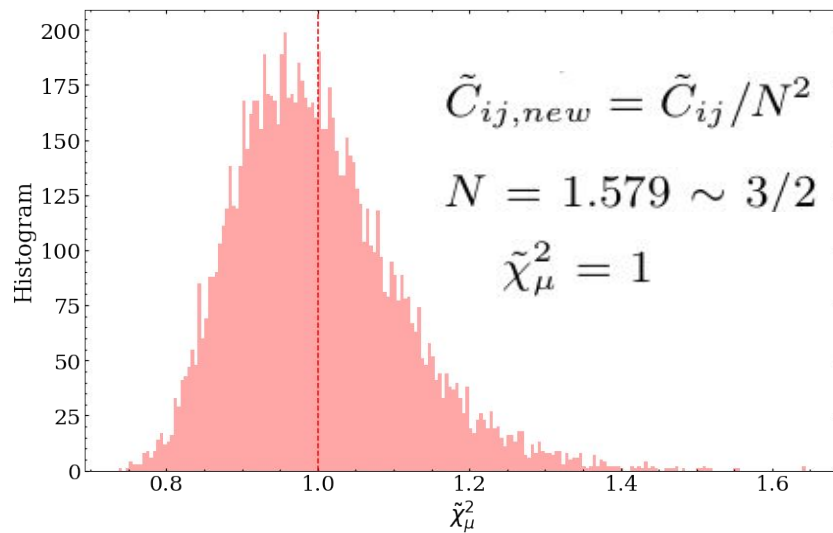
$$\chi_{\mu}^2 = \sum_{i,j=1}^{32} [H(z_i) - H_{rec,\mu}(z_i)] \tilde{C}_{ij}^{-1} [H(z_j) - H_{rec,\mu}(z_j)]$$

$$P_{\tilde{\chi}_{K_i}^2 < \tilde{\chi}_{K_j}^2} = \frac{1}{1 + P_j/P_i}$$

Kernels	P_{SE}/P_j
SE vs DSE	1.42
SE vs M32	0.62
SE vs M52	0.72
SE vs M72	0.82
SE vs M92	0.83



Considering smaller uncertainties in CCH data



Backup slides

Results

$$M = (-19.326^{+0.050}_{-0.068})$$

$$\Omega_k = 0.10^{+0.12}_{-0.15}$$

$$r_d = (142.6^{+3.9}_{-3.5}) \text{ Mpc}$$

