

A definitive test of the cold dark matter hypothesis

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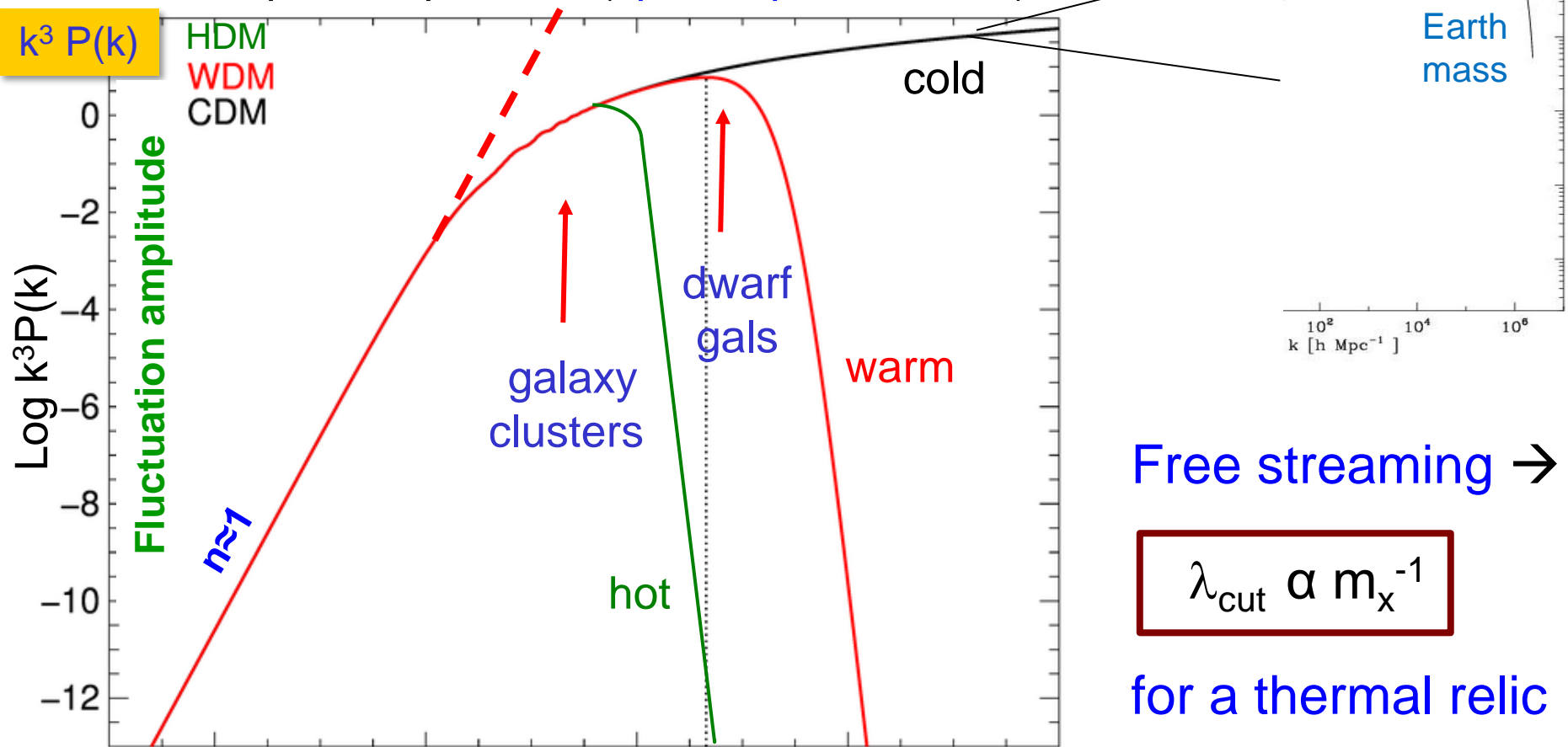
Non-baryonic dark matter candidates

From the early 1980s:

Type	example	mass
hot	neutrino	few tens of eV
warm	sterile ν	keV-MeV
cold	axion neutralino	10^{-5} eV - 100 GeV

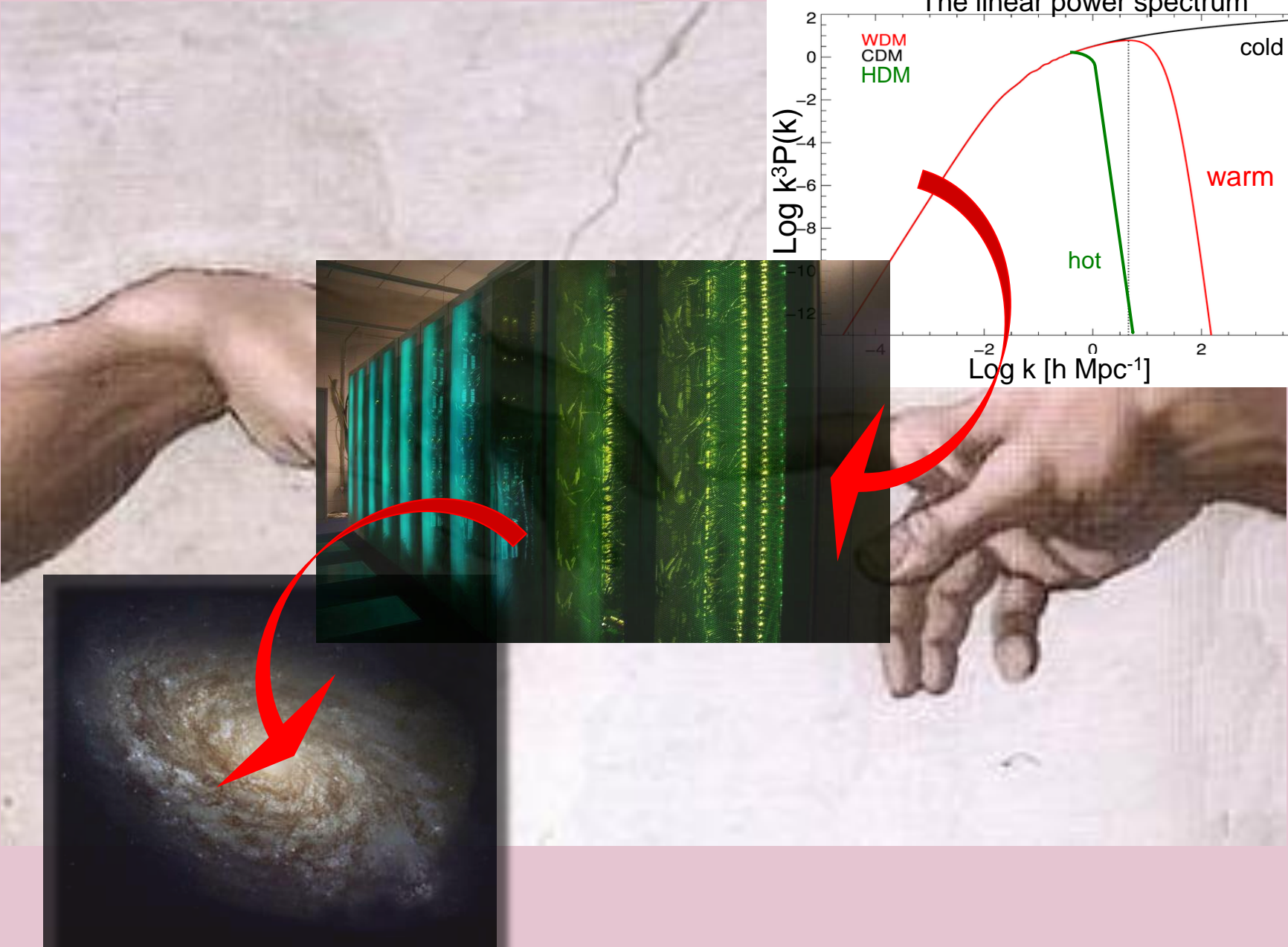
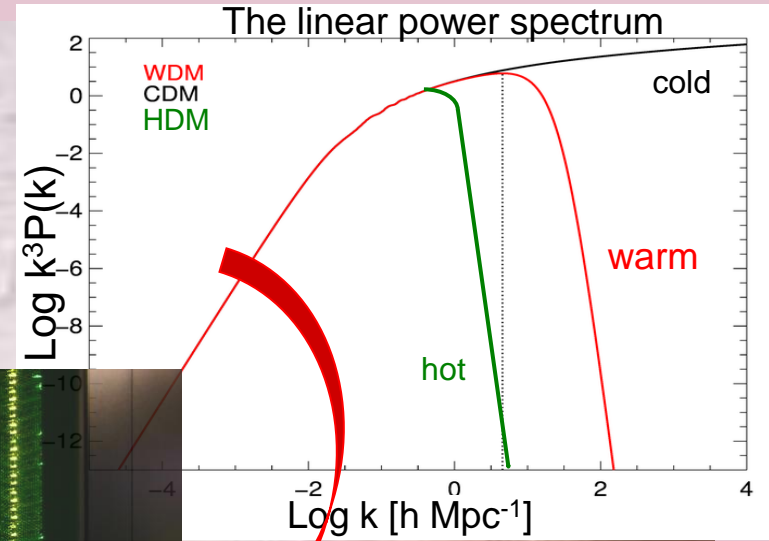
The dark matter power spectrum

The linear power spectrum ("power per octave")



These possibilities can be tested with astrophysics

Non-linear evolution



Non-linear evolution: simulations

Assumption about content of Universe → Initial conditions

Relevant equations:

Collisionless Boltzmann;
 Poisson; Friedmann eqns;
 Radiative hydrodynamics
 Subgrid astrophysics



How to make a virtual universe



Hot dark matter

$$m_\nu = 30 \text{ eV} \rightarrow \Omega_m = 1$$

1981

HAS THE NEUTRINO A NON-ZERO REST MASS?
(Tritium β -Spectrum Measurement)

V. Lubimov, E. Novikov, V. Nozik, E. Tretyakov
Institute for Theoretical and Experimental Physics, Moscow, U.S.S.R.

V. Kosik
Institute of Molecular Genetics, Moscow, U.S.S.R.

ABSTRACT

The high energy part of the β -spectrum of tritium in the molecule was measured with high precision by a toroidal β -spectrometer. The results give evidence for a non-zero electron anti-neutrino mass.

Fifty years ago Pauli introduced the neutrino to explain the β -spectrum shape. Pauli made the first estimate of the neutrino mass ($E_\beta \text{ max} \approx$ nuclei mass defect): it should be very small or maybe zero. Up to now the study of the β -spectrum shape is the most sensitive, direct method of neutrino mass measurement.

For allowed β -transitions, if $M_\nu = 0$, then $S \approx (E - E_0)^2$. The Kurie plot is then a straight line with the only kinematic parameter being $E_k = E_0$ (total β -transition energy). If $M_\nu \neq 0$, then $S \approx (E_0 - E) \sqrt{(E_0 - E)^2 - M_\nu^2}$. The Kurie plot is then distorted, especially near the endpoint.

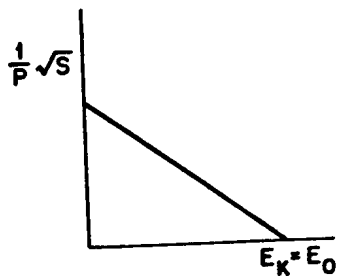


Fig. 1. Kurie plot for $M_\nu = 0$.

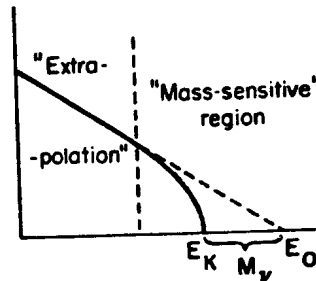


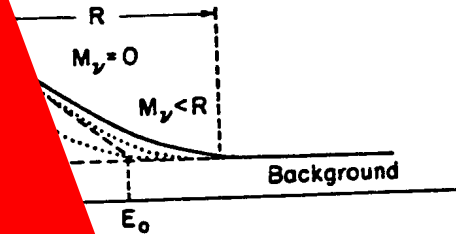
Fig. 2. Kurie plot for $M_\nu \neq 0$.

The method for the neutrino mass measurement is to obtain E_0 from the extrapolation and obtain E_k from the spectrum intercept. Then $M_\nu = E_0 - E_k$. Qualitatively, $M_\nu \neq 0$ if the β -spectrum near the endpoint runs below the extrapolated curve.

* Paper presented by Oleg Egorov.



things are more complicated. The apparatus resolution strongly affects the spectrum endpoint and rather the spectrum slope.



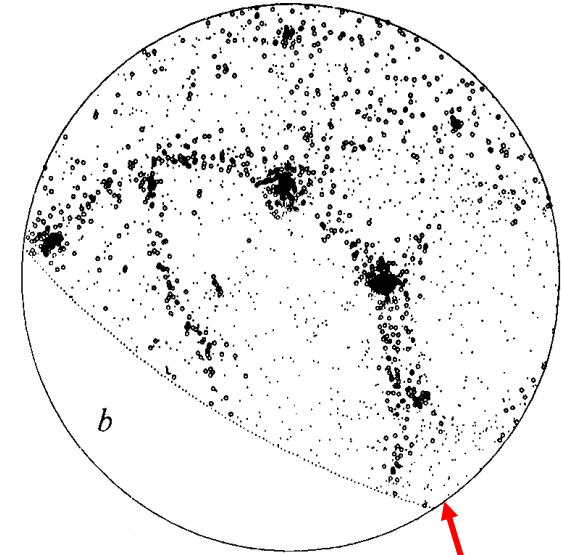
Realistic Kurie plot.

extrapolation. However, we are unable to determine M_ν , then once again the lack of counts near the endpoint indicate that $M_\nu \neq 0$. If $M_\nu \leq R$, the changes due to M_ν and the influence of R are indistinguishable. For M_ν determination the knowledge of R is compulsory. The background determines the statistical accuracy near the endpoint, i.e., in the region of the highest sensitivity to the ν mass. So: 1) R should be $\sim M_\nu$, 2) the smaller M_ν is, the smaller the background ($\sim M_\nu^2$) must be and the higher the statistics ($\sim M_\nu^{-3}$) must be. For example, suppose that for $M_\nu = 100$ eV we need resolution R , background Q , and statistics N . If $M_\nu = 30$ eV, to achieve the same $\Delta M/M$ they should be $R/3$, $Q/10$, and $N \times 30$, respectively.

The shorter the β -spectrum, the less it is spread due to R (as $R \sim \Delta p/p = \text{const.}$). A classical example is ^3H β -decay, which has 1) the smallest $E_0 \sim 18.6$ keV, 2) an allowed β -transition, simple nucleus, and simple theoretical interpretation, 3) highly reduced radioactivity. The first experiments with ^3H were by S. Curran et al. (1948) and G. Hanna, B. Pontecorvo (1949). Using ^3H gas in a proportional counter, they obtained $M_\nu \leq 1$ keV. Further progress required magnetic spectrometer development. This allowed the resolution to be improved considerably, and L. Langer and R. Moffat (1952) obtained $M_\nu \leq 250$ eV. The best value was obtained by K. Bergkvist (1972): $R \sim 50$ eV and $M_\nu \leq 55$ eV.

The ITEP spectrometer is of a new type: ironless, with toroidal magnetic field (E. Tretyakov, 1973). The principle of the toroidal magnetic field focusing systems was proposed by V. Vladimirovsky et al. (An example is a "Horn" of ν -beams.) It turns out that a rectilinear conductor (current) has a focusing ability for particles emitted perpendicular to the rotation axis. This system has infinite periodical focusing structure. The ITEP spectrometer is based on this principle.

Non-baryonic dark matter cosmologies



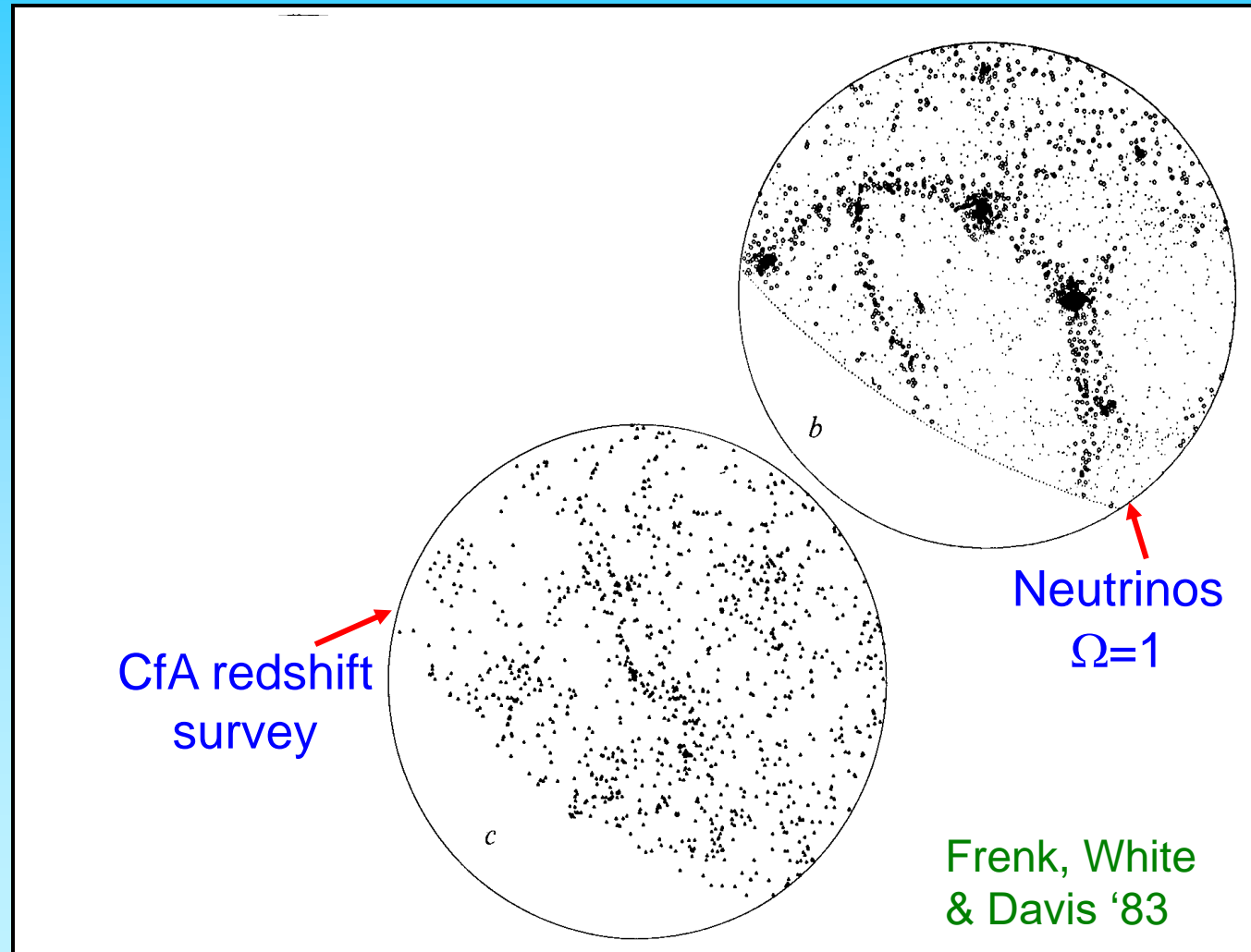
Neutrinos
 $\Omega=1$

Frenk, White
& Davis '83

Non-baryonic dark matter cosmologies

Neutrino DM \rightarrow
wrong clustering

Neutrinos cannot
make appreciable
contribution to Ω
 $\rightarrow m_\nu \ll 30$ eV



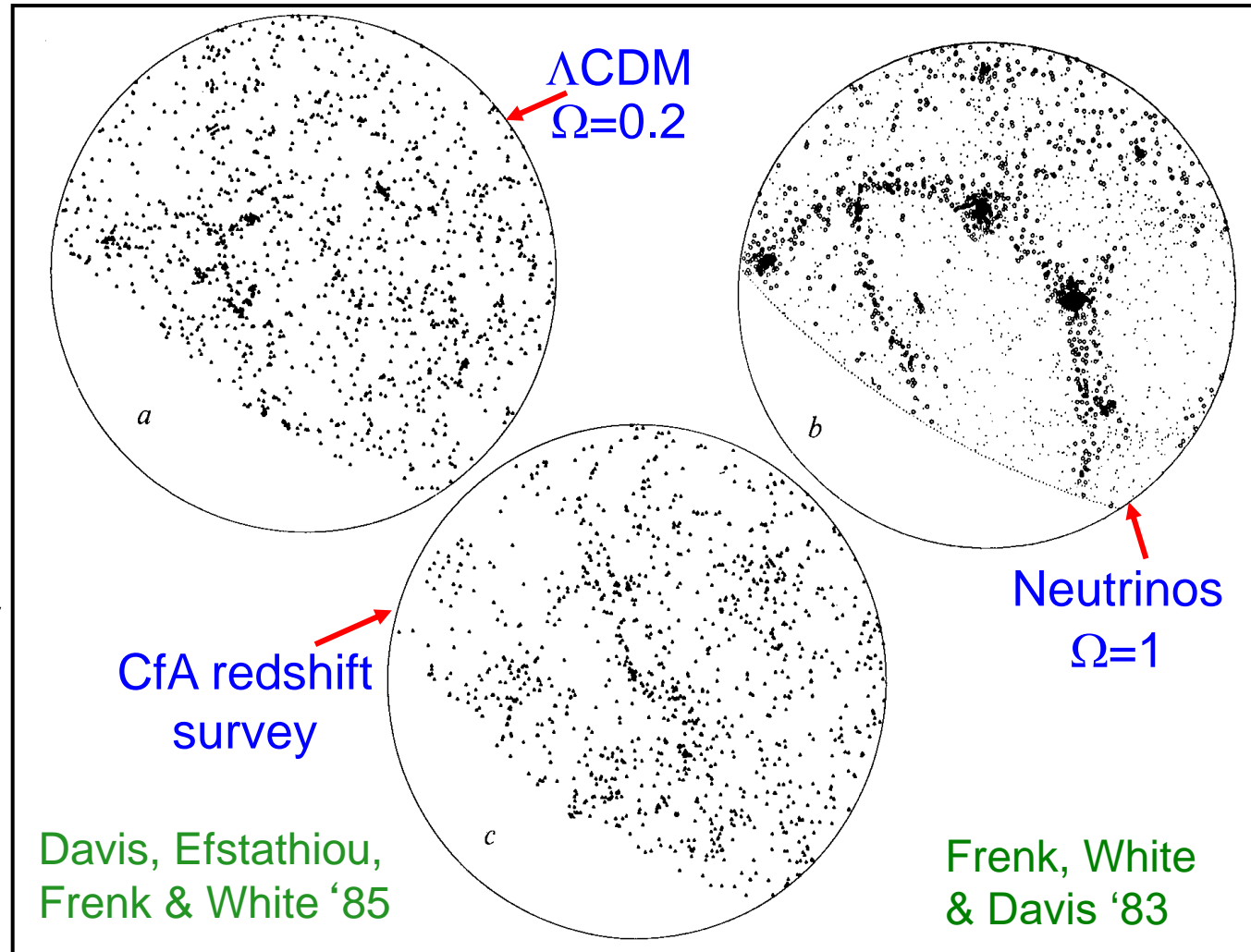
Non-baryonic dark matter cosmologies

Neutrino DM \rightarrow
wrong clustering

Neutrinos cannot
make appreciable
contribution to Ω
 $\rightarrow m_\nu \ll 30$ eV

Early CDM N-body
simulations gave
promising results

In CDM structure
forms hierarchically



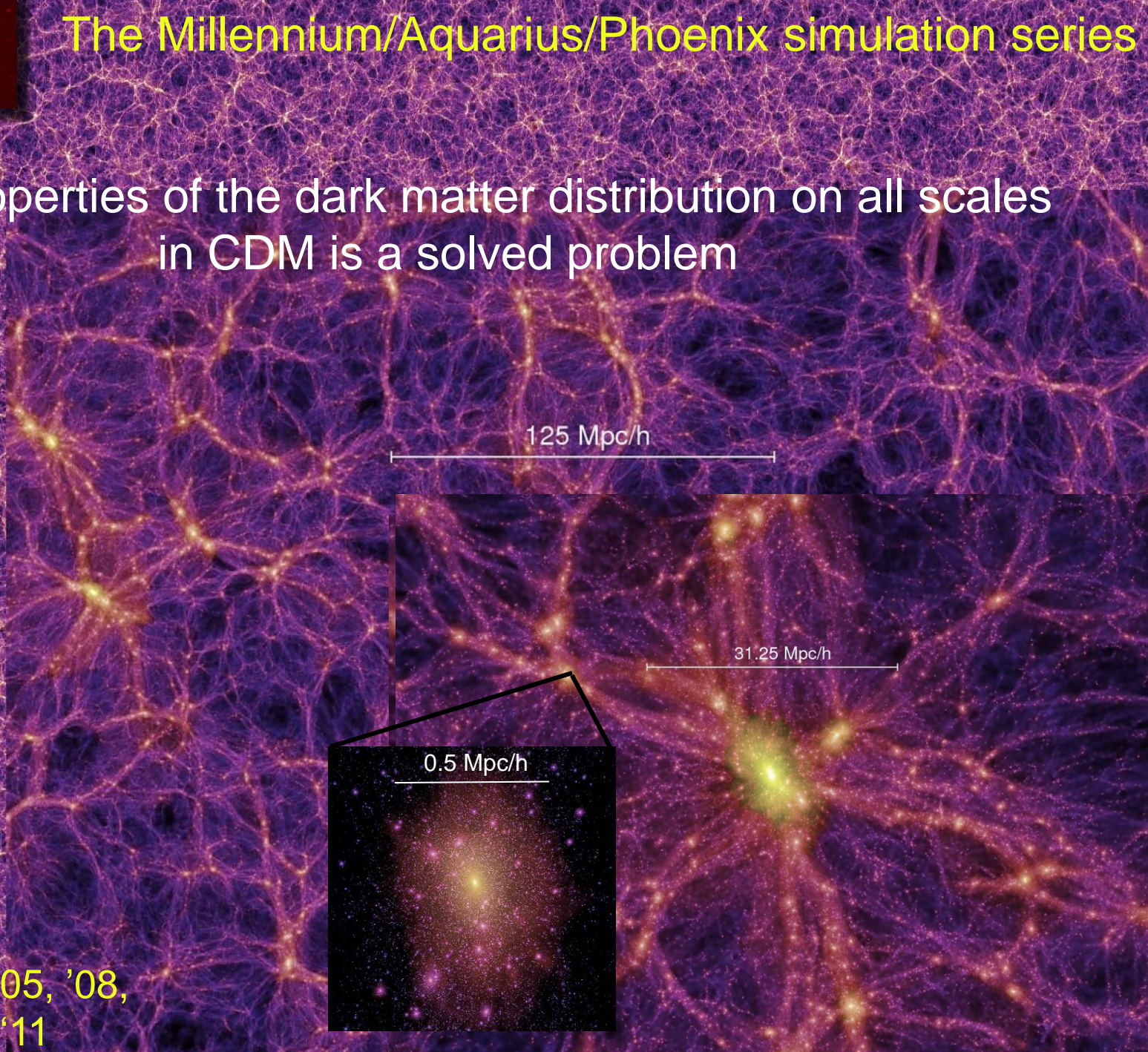


The properties of the dark matter distribution on all scales
in CDM is a solved problem

VIRGO

The Millennium/Aquarius/Phoenix simulation series

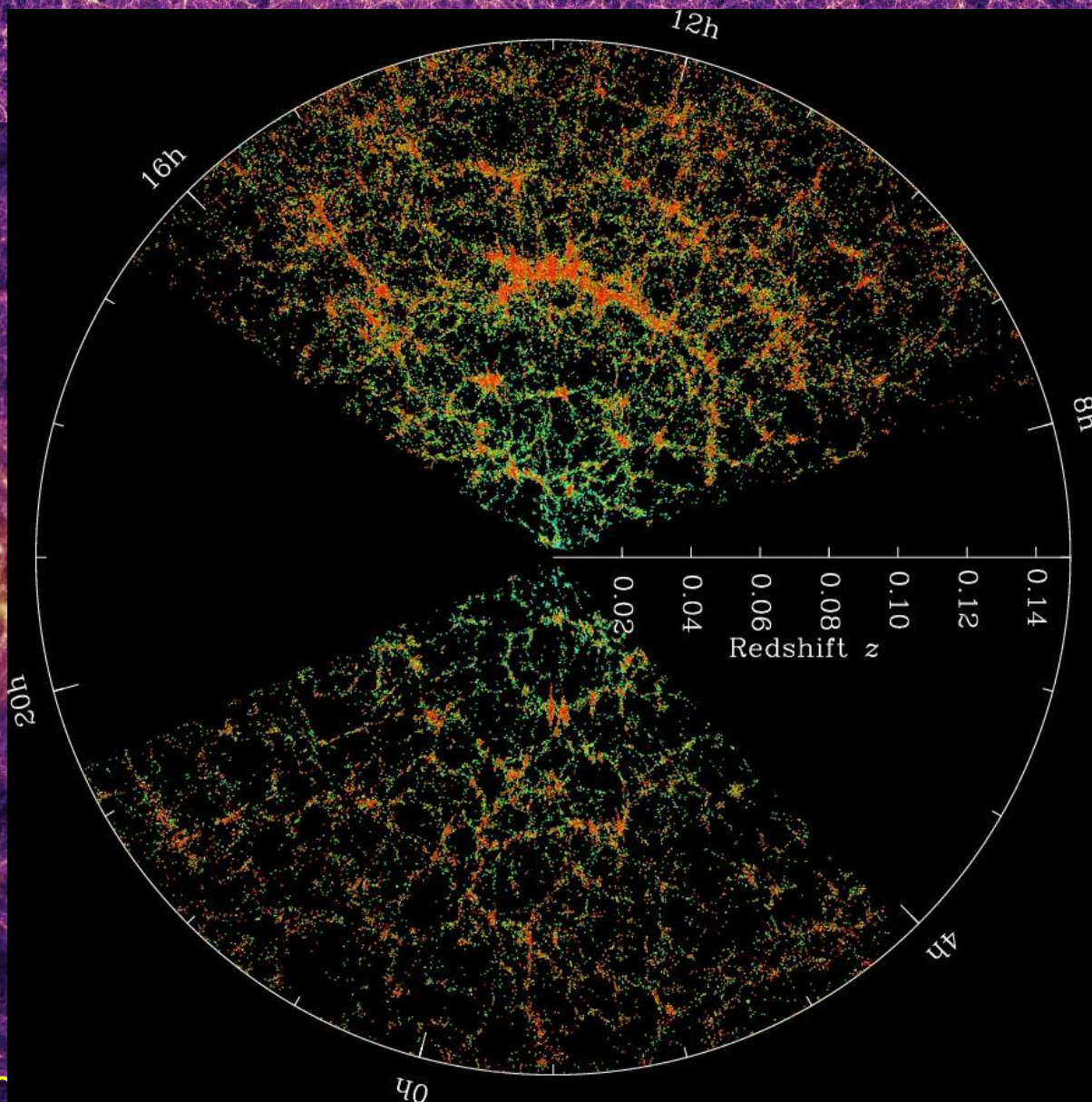
The properties of the dark matter distribution on all scales in CDM is a solved problem



Springel et al '05, '08,
Gao et al '11

VIRGO

The Millennium/Aquarius/Phoenix simulation series



Springel et al '05, '06,
Gao et al '11

Galaxy distribution encodes info about dark matter and dark energy

5 billion yrs

DESI already has > 10 million spectra

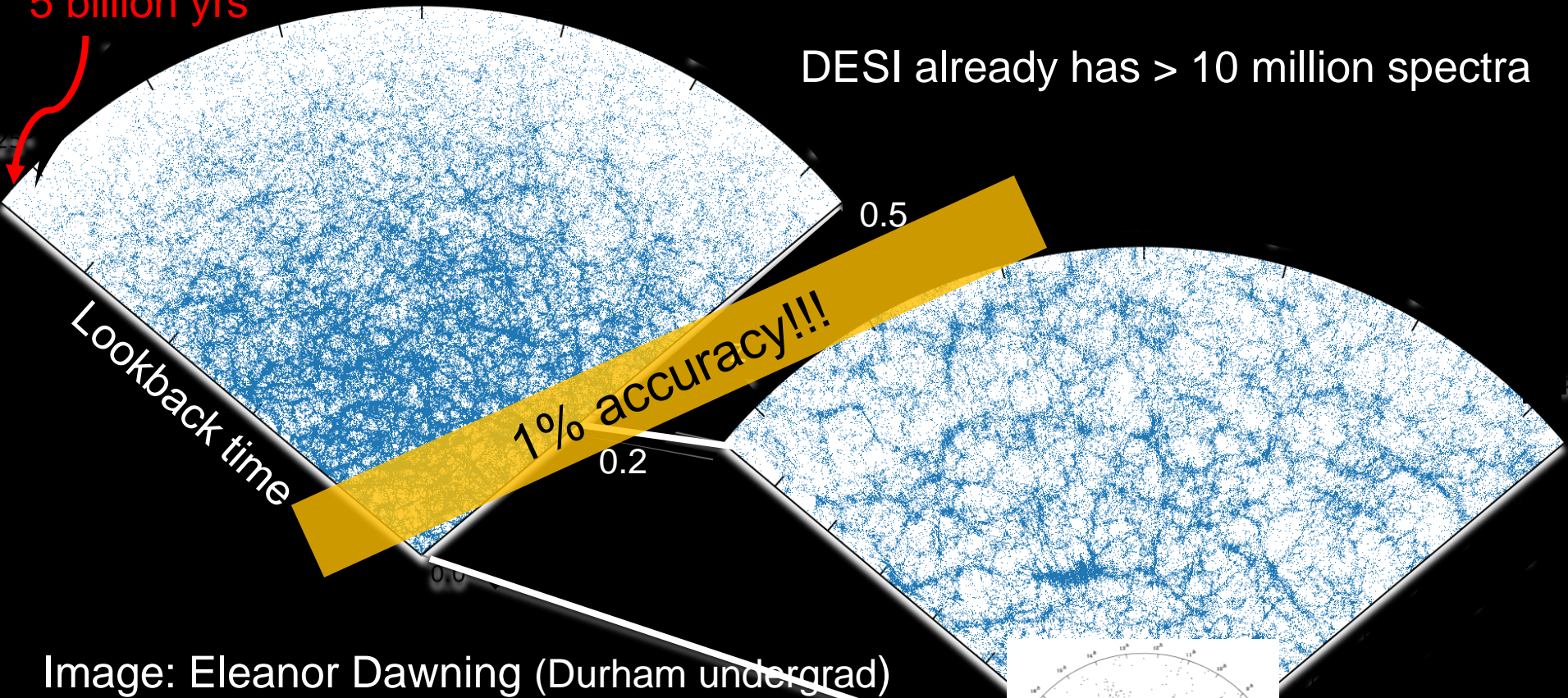
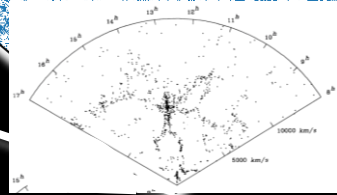
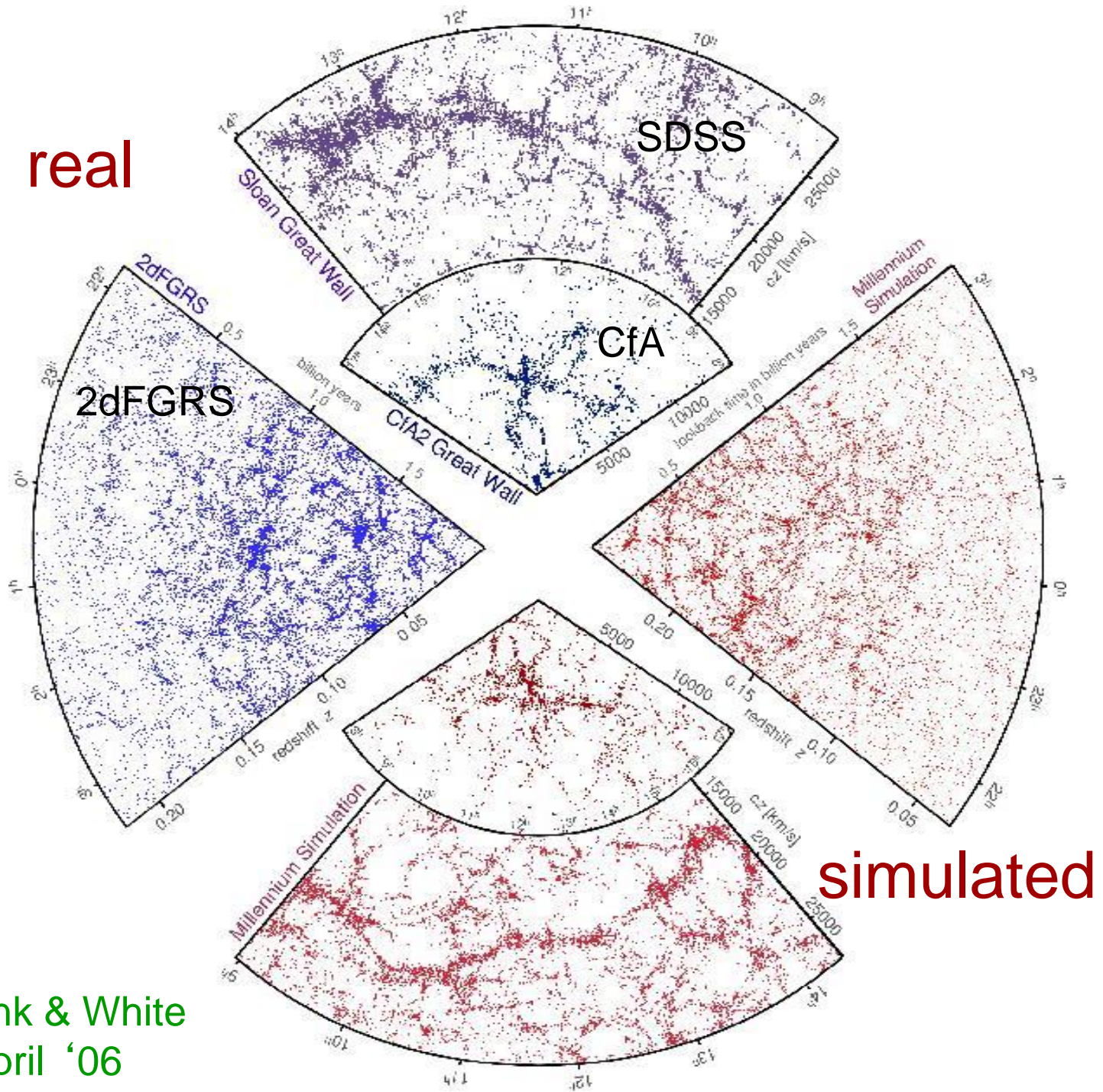


Image: Eleanor Dawning (Durham undergrad)



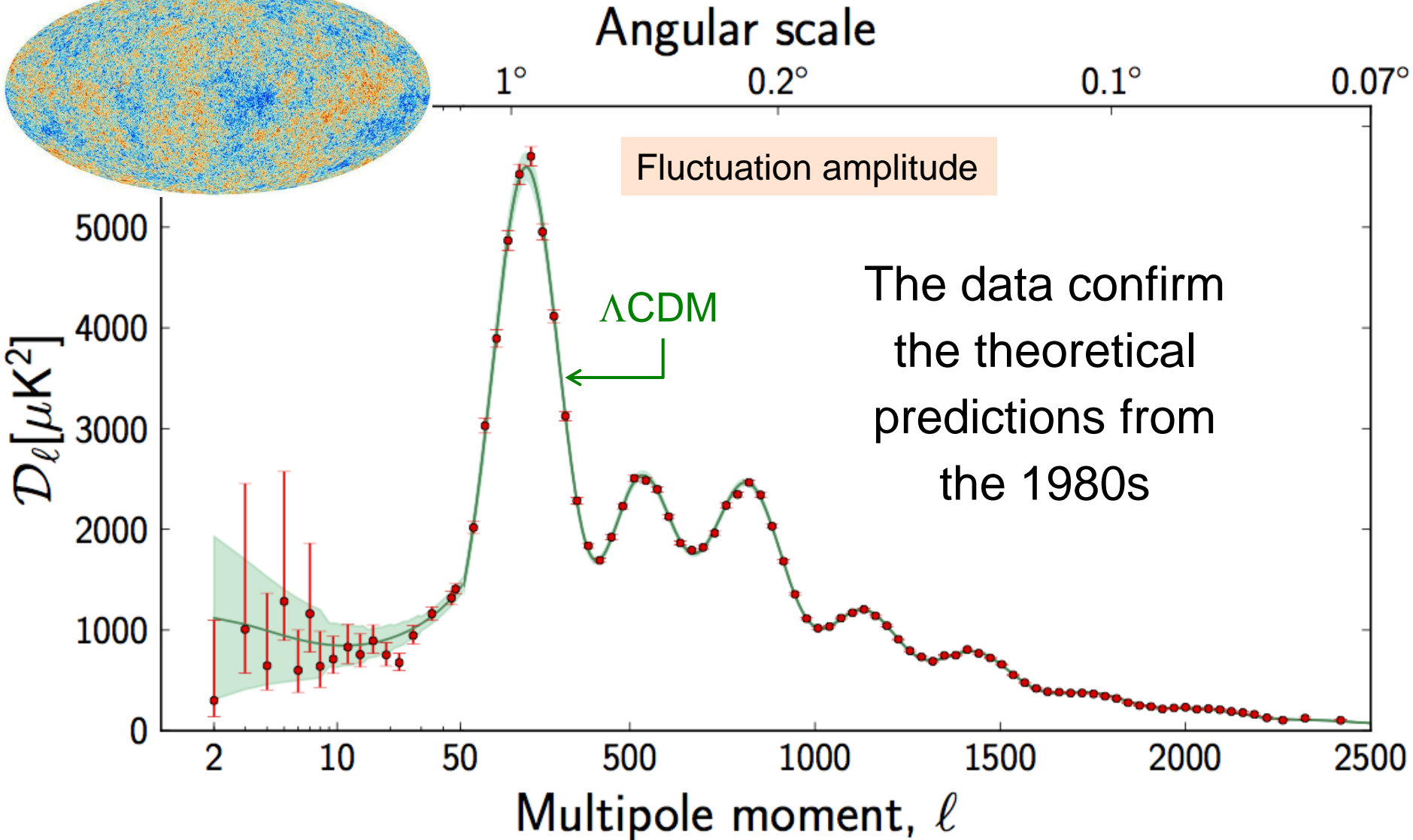
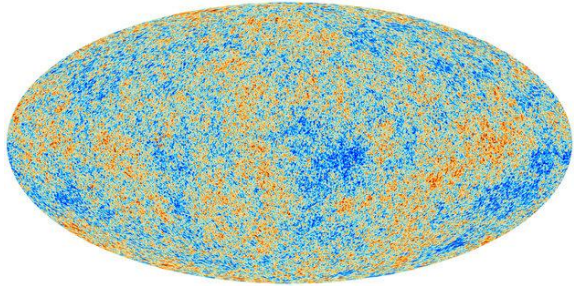
real



simulated

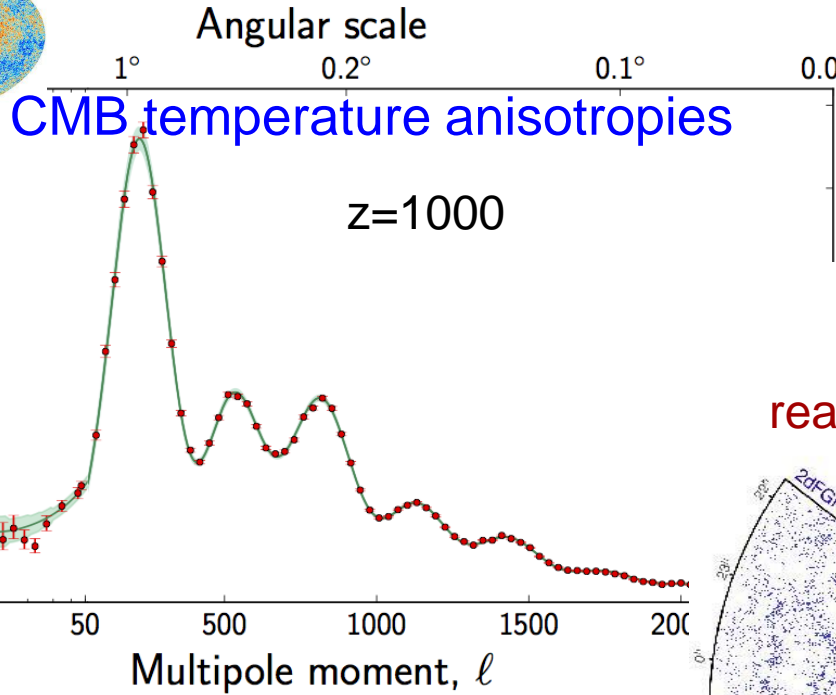
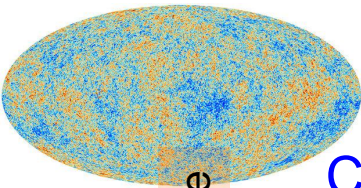
Springel, Frenk & White
Nature, April '06

Planck: CMB temperature anisotropies



The Λ CDM model of cosmogony

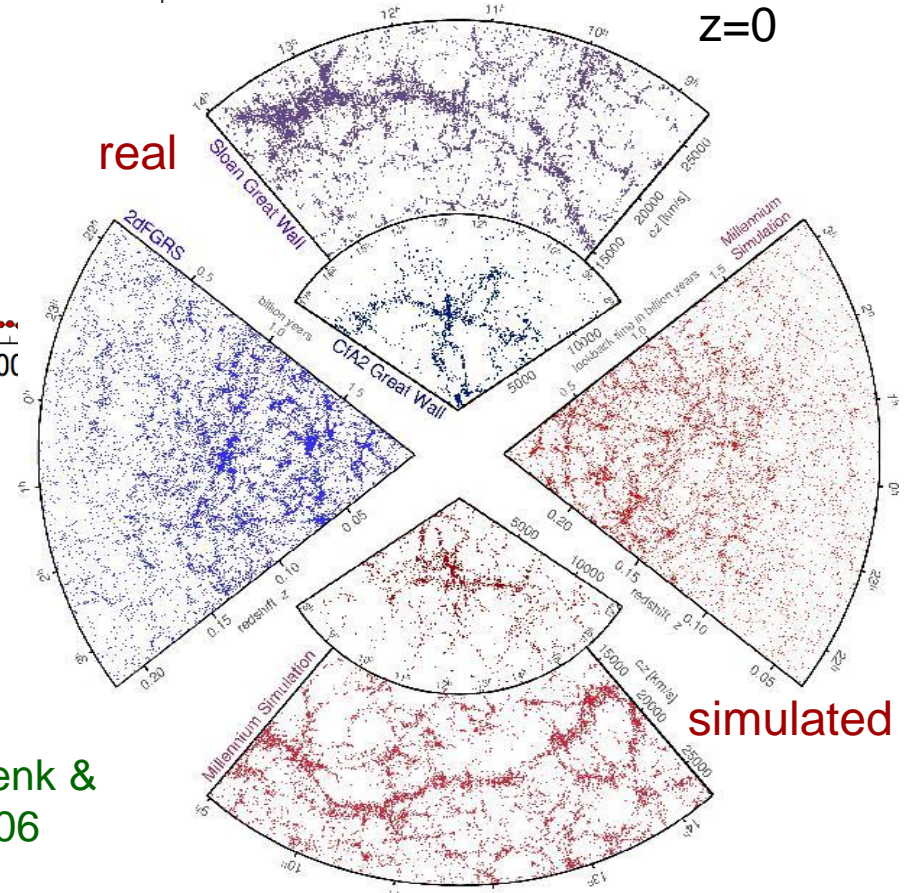
Proposed in 1980s; now empirically supported by:



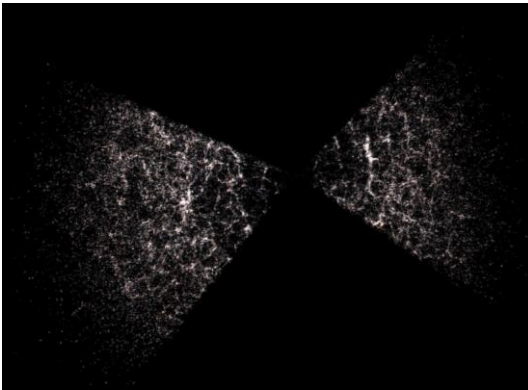
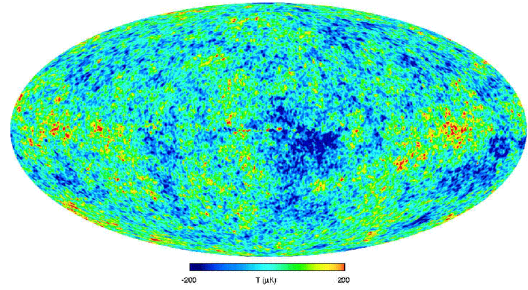
Planck coll. 2015

Springel, Frenk & White 2006

Galaxy clustering



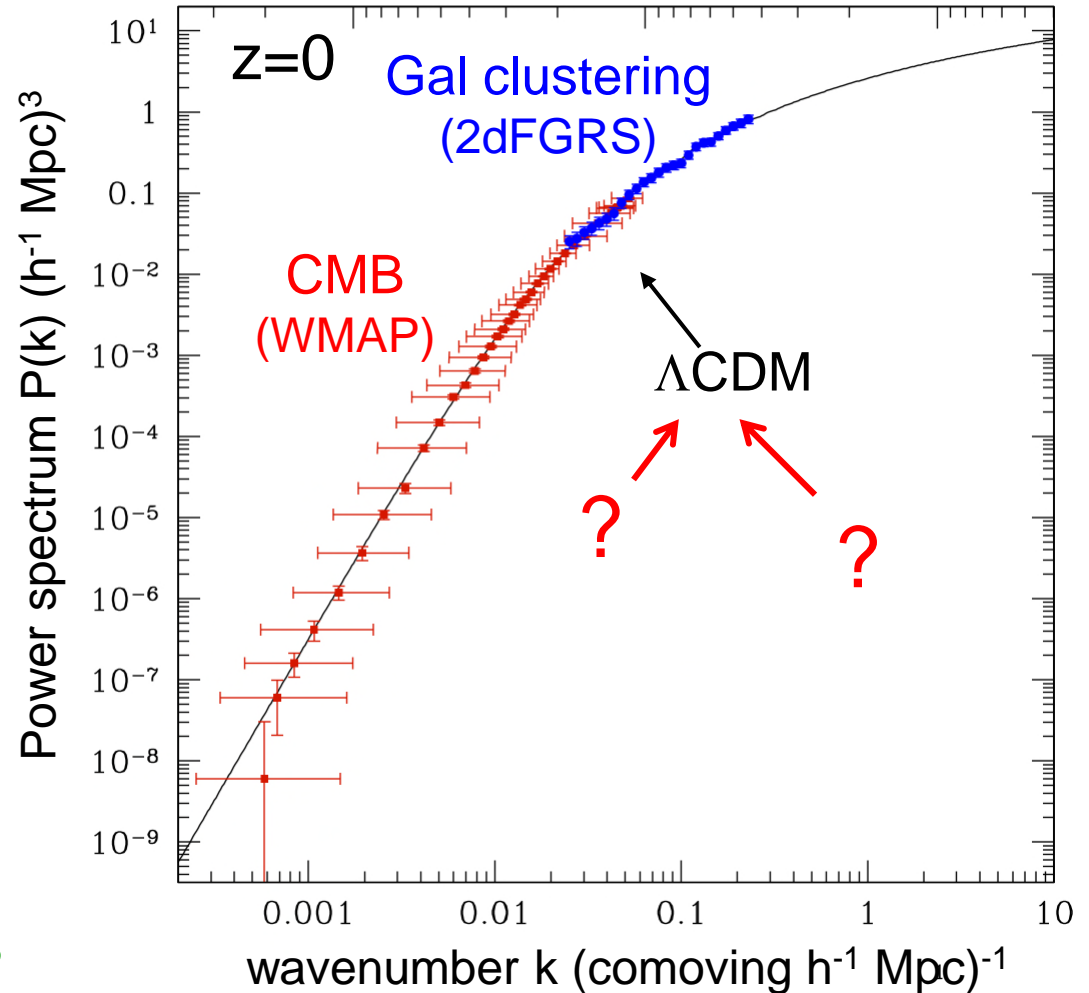
The cosmic power spectrum: from the CMB to the 2dFGRS



z~1000

Log $k^3 P(k)$

wavelength k^{-1} (comoving h^{-1} Mpc)



⇒ Λ CDM provides an excellent description of mass power spectrum from 10-1000 Mpc

Sanchez et al 06

The cosmic power spectrum: from the CMB to the 2dFGRS

Free streaming \rightarrow

$$\lambda_{\text{cut}} \propto m_x^{-1}$$

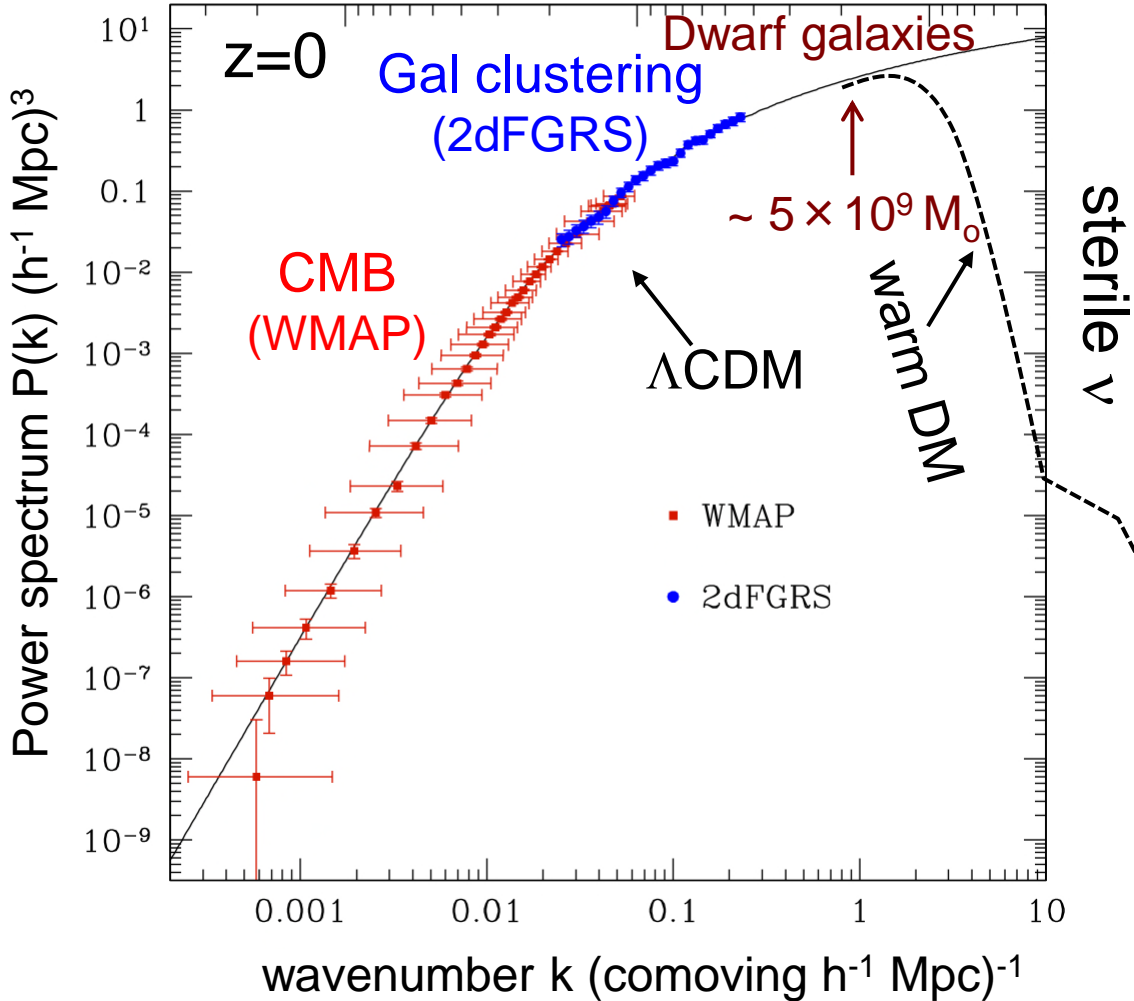
for thermal relic

$m_{\text{CDM}} \sim 100 \text{ GeV}$
 susy; $M_{\text{cut}} \sim 10^{-6} M_{\odot}$

$m_{\text{WDM}} \sim \text{few keV}$
 sterile ν ; $M_{\text{cut}} \sim 10^9 M_{\odot}$

Log $k^3 P(k)$ wavelength k^{-1} (comoving $h^{-1} \text{ Mpc}$)

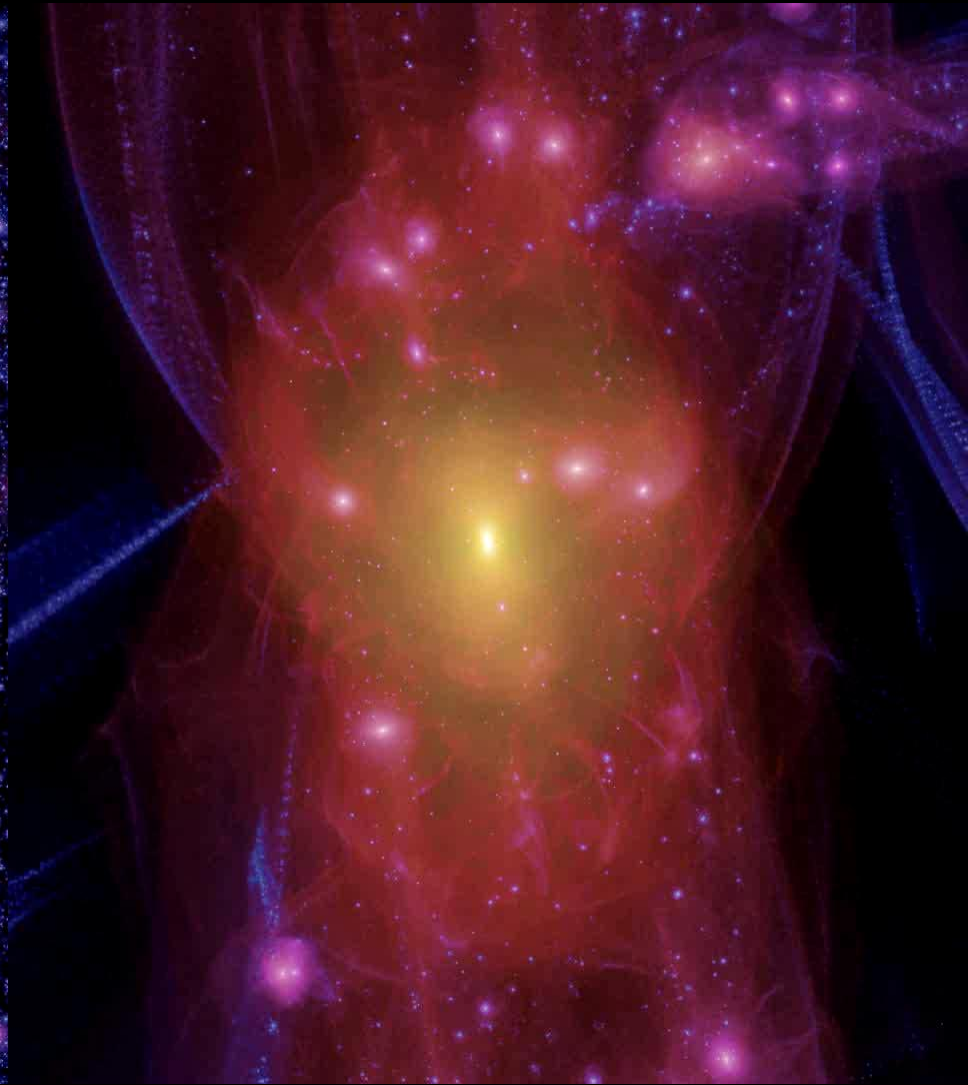
1 000 100 10



cold dark matter



warm dark matter



Lovell, Eke, Frenk, Gao, Jenkins, Wang, White, Theuns,
Boyarski & Ruchayskiy '12

The properties of the dark matter distribution on all scales in CDM is a solved problem

We now know:

- halo mass function down to the cutoff mass
- the internal structure of halos of all masses
- the spatial distribution of halos & diffuse DM

125 Mpc/h

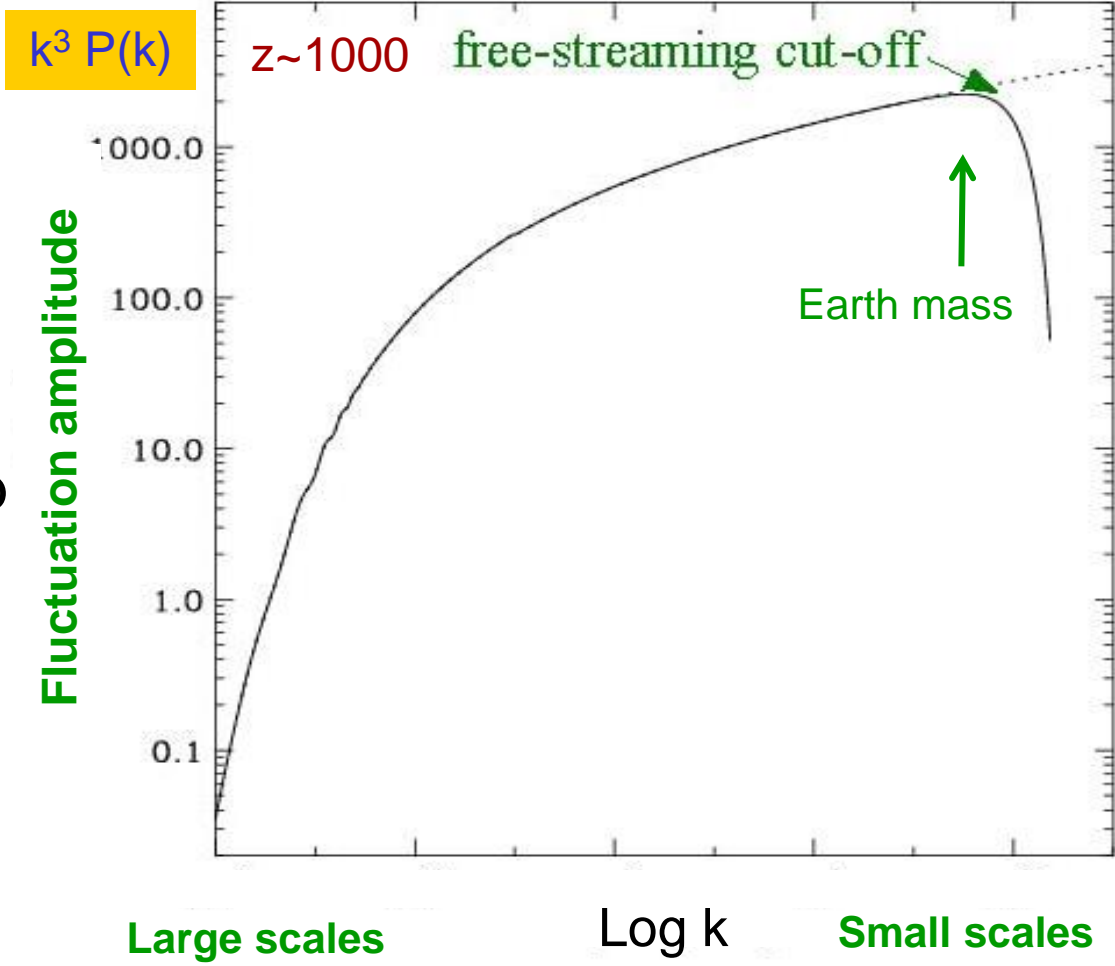
0.5 Mpc/h

31.2 Mpc/h

The cold dark matter power spectrum

The linear power spectrum
("power per octave")

Assumes a 100GeV wimp
Green et al '04



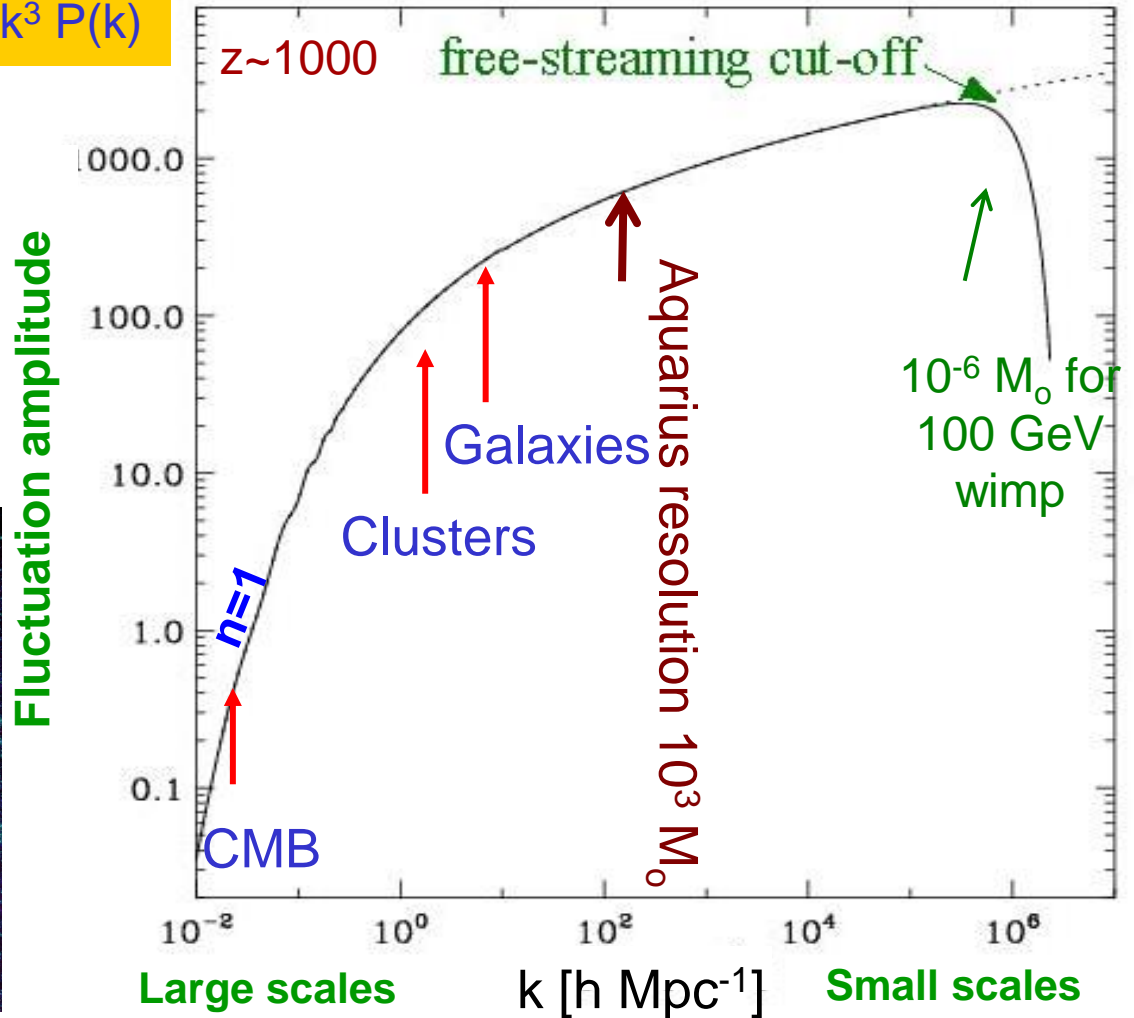
The cold dark matter linear power spectrum

The linear power spectrum
 (“power per octave”)

$$\lambda_{\text{cut}} \propto m_x^{-1}$$

Assumes a 100GeV wimp
 Green et al '04

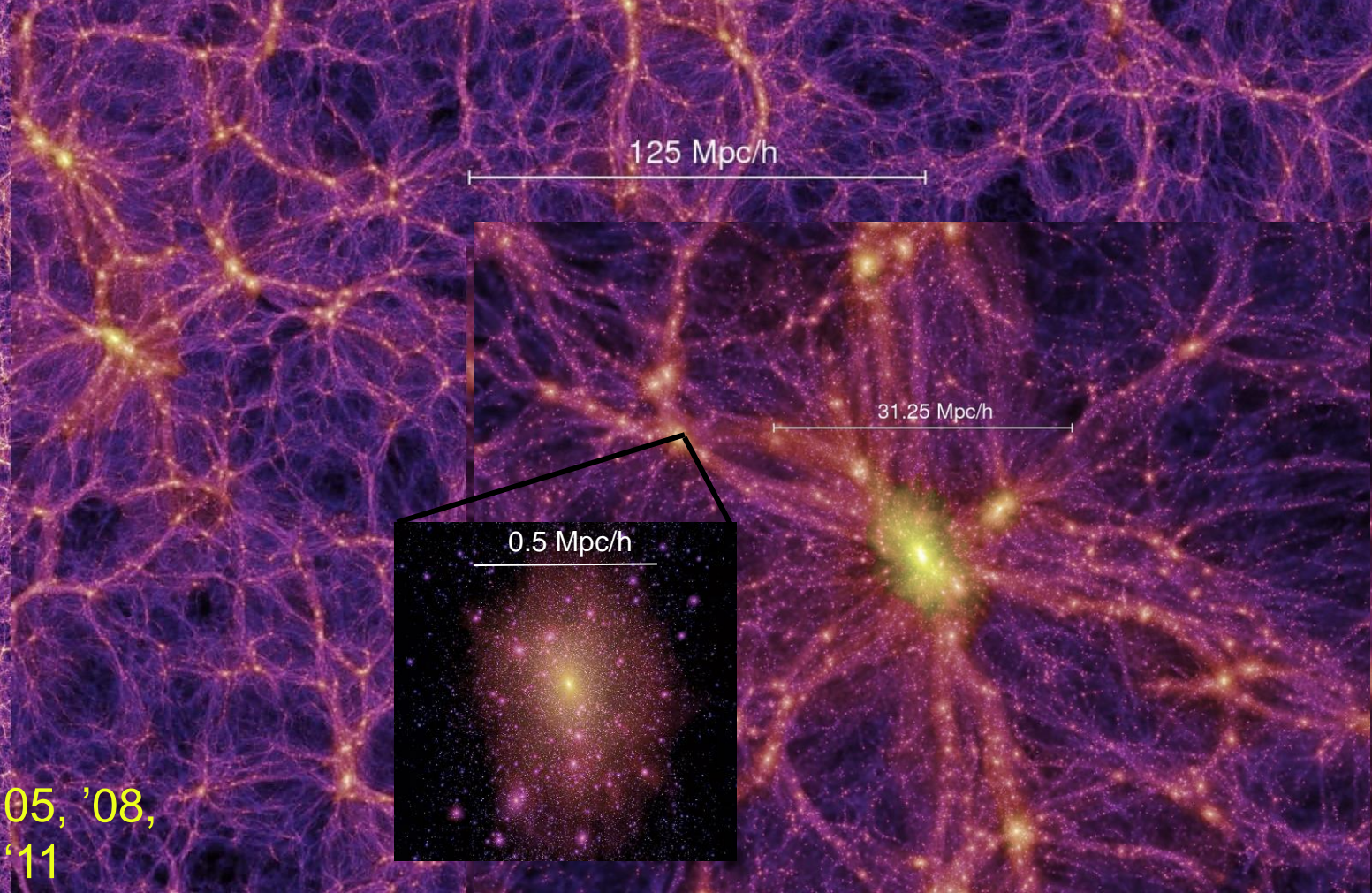
$k^3 P(k)$



VIRGO

The Millennium/Aquarius/Phoenix simulation series

To resolve Earth-mass halos in a cosmological simulation would require 10^{27} particles \rightarrow impossible



Springel et al '05, '08,
Gao et al '11

The VVV simulation

Planck cosmology

Dark matter only

Dynamic range of
30 orders of
magnitude in mass

$$M_{\text{char}} = 10^{14} M_{\odot}$$

Base Level

L0

150 Mpc

The VVV simulation

Planck cosmology

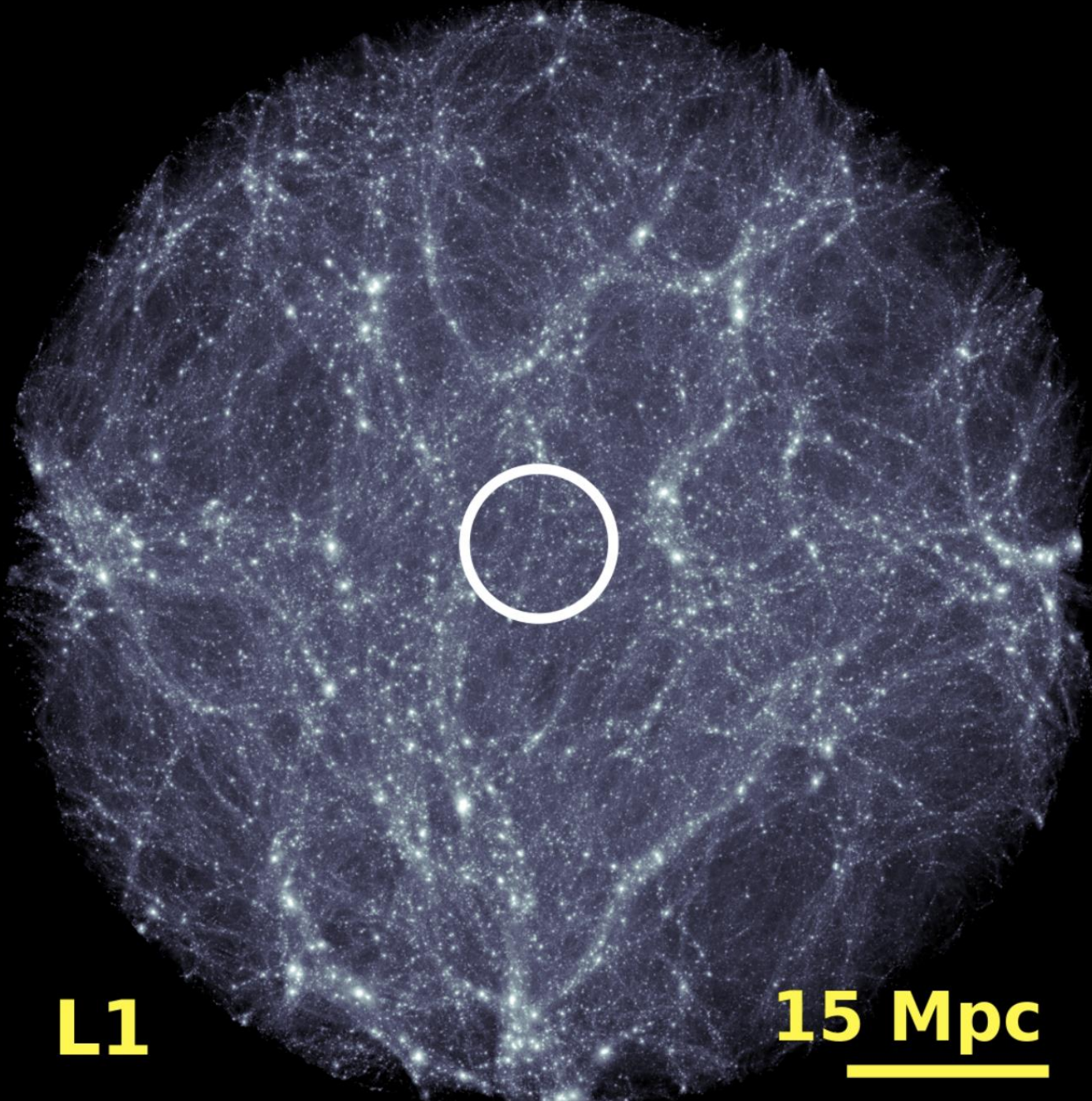
Dark matter only

Dynamic range of
30 orders of
magnitude in mass

$$M_{\text{char}} = 10^{12} M_{\odot}$$

Zoom Level 1

Wang, Bose et al 2020



L1

15 Mpc

The VVV simulation

Planck cosmology

Dark matter only

Dynamic range of
30 orders of
magnitude in mass

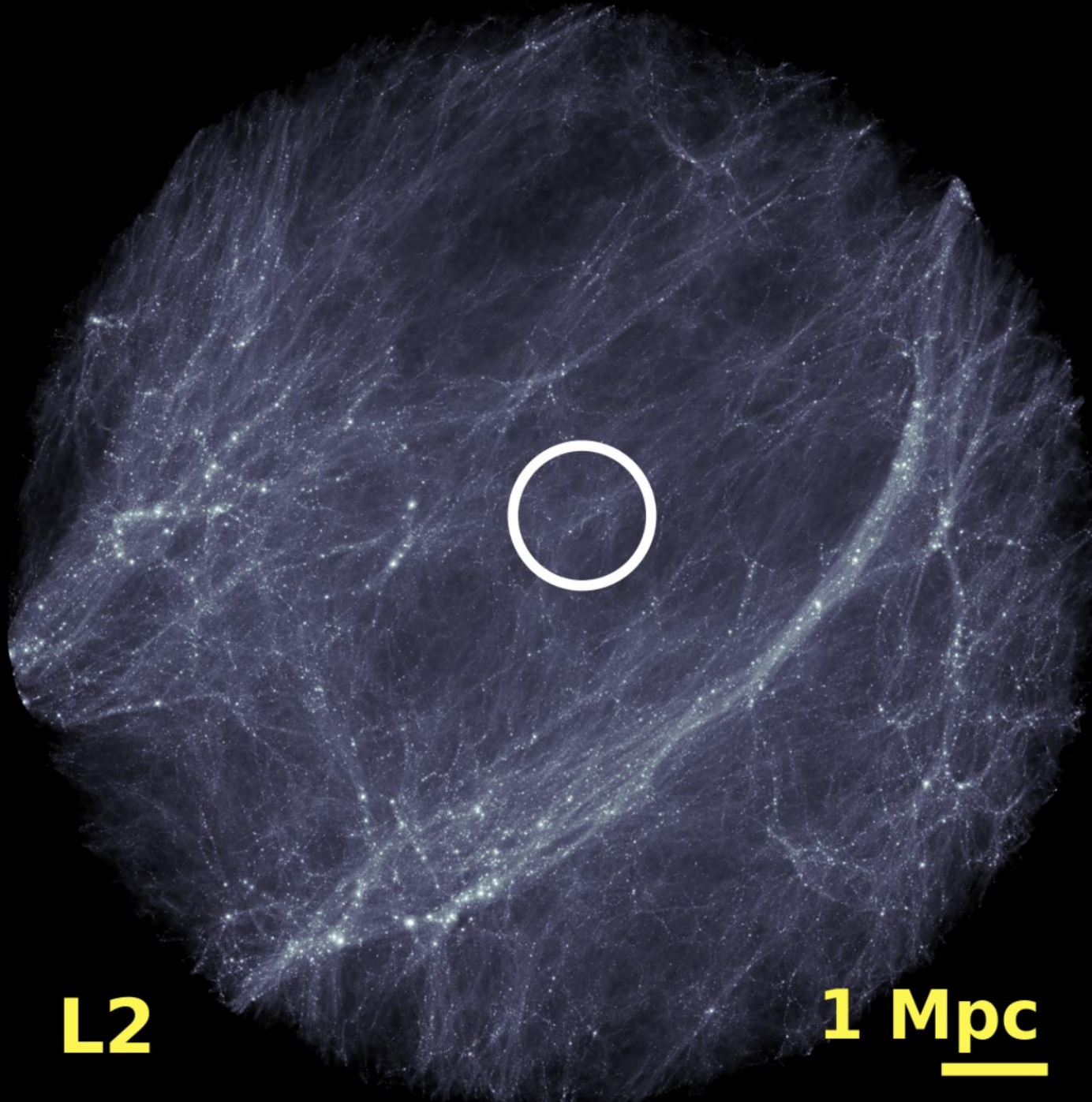
$$M_{\text{char}} = 10^9 M_{\odot}$$

Zoom Level 2

Wang, Bose et al 2020

L2

1 Mpc



The VVV simulation

Planck cosmology

Dark matter only

Dynamic range of
30 orders of
magnitude in mass

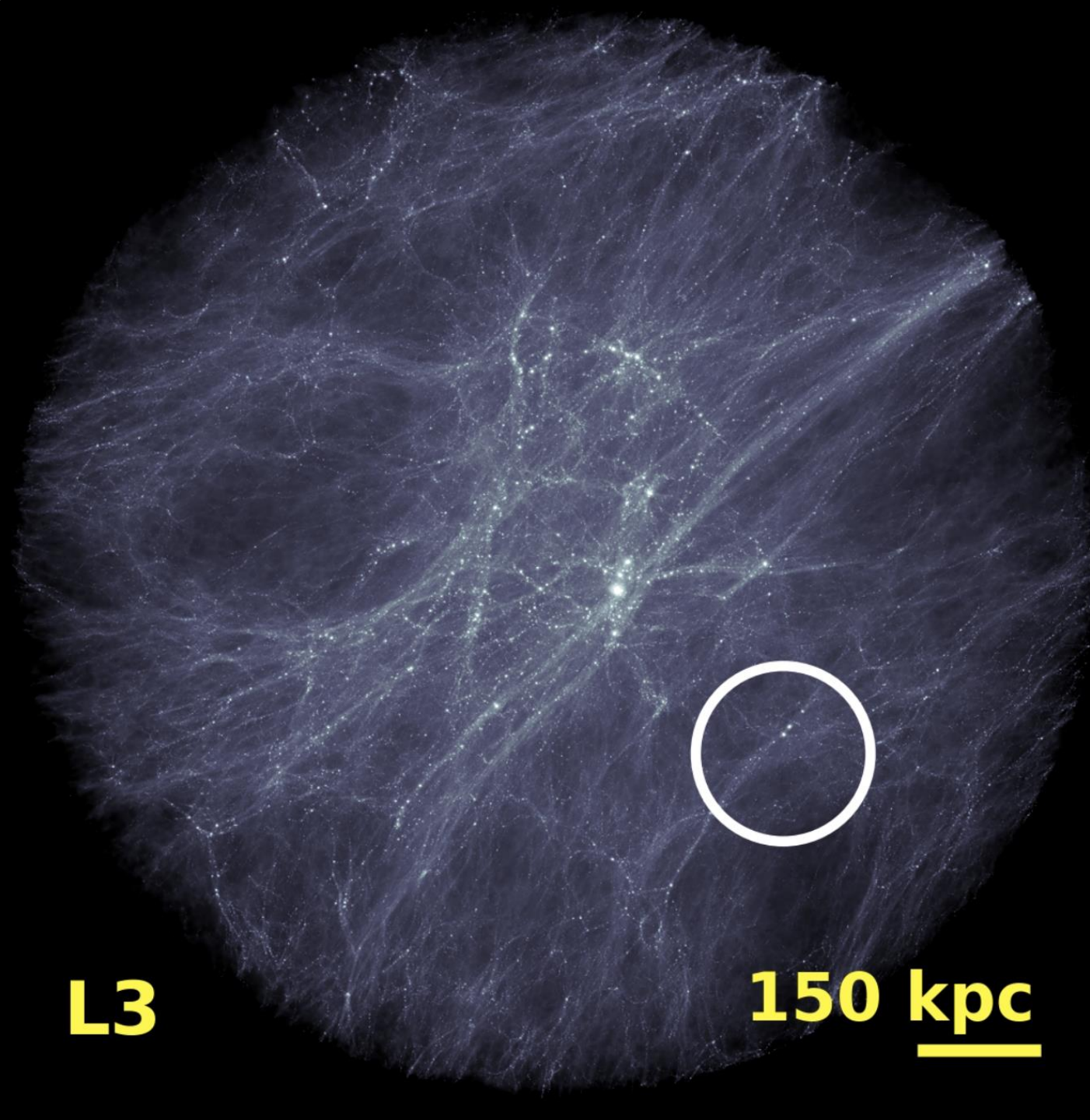
$$M_{\text{char}} = 10^6 M_{\odot}$$

Zoom Level 3

Wang, Bose et al 2020

L3

150 kpc



The VVV simulation

Planck cosmology

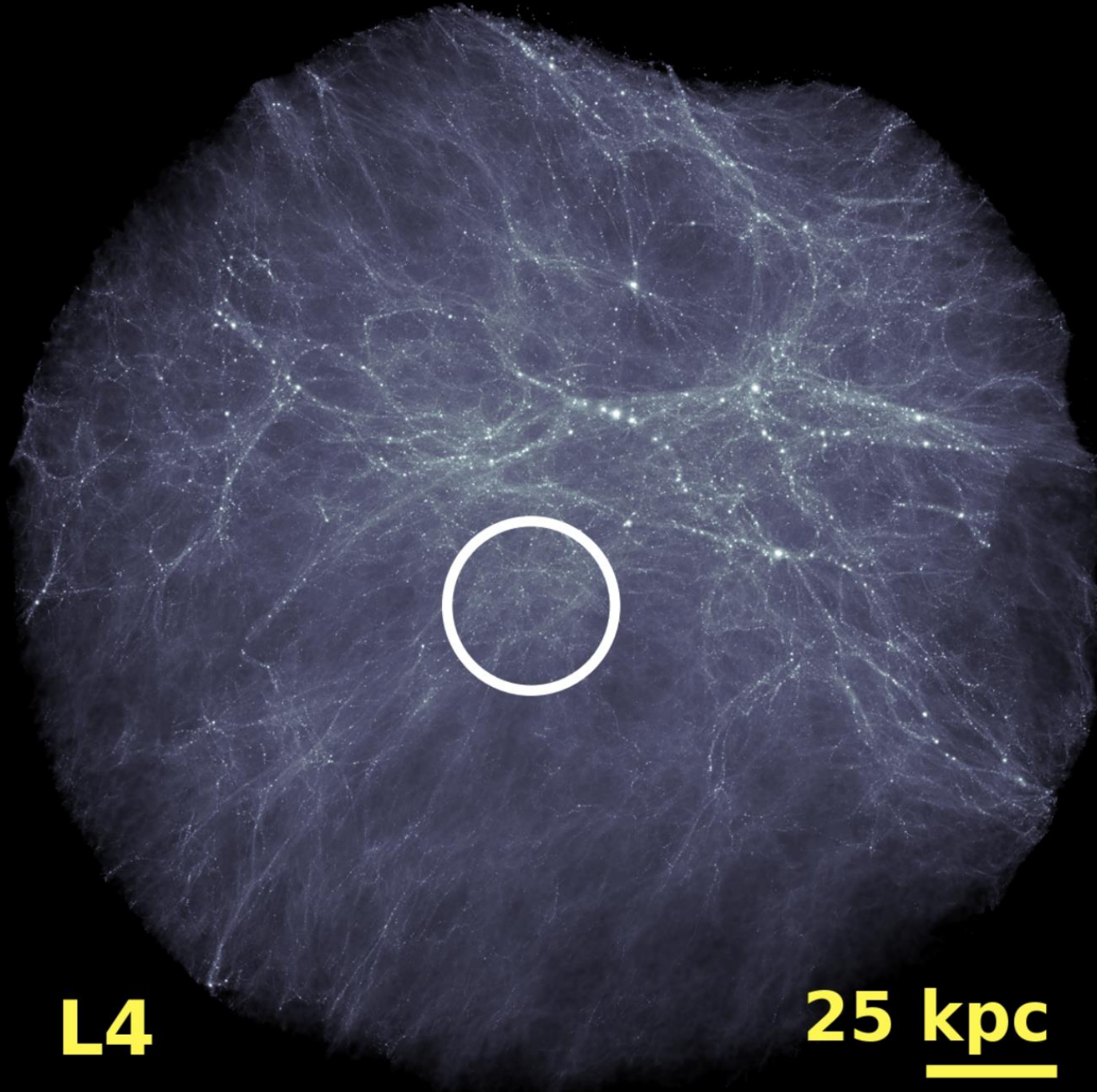
Dark matter only

Dynamic range of
30 orders of
magnitude in mass

$$M_{\text{char}} = 10^3 M_{\odot}$$

Zoom Level 4

Wang, Bose et al 2020



L4

25 kpc

The VVV simulation

Planck cosmology

Dark matter only

Dynamic range of
30 orders of
magnitude in mass

$$M_{\text{char}} = 10 M_{\odot}$$

Zoom Level 5



L5

5 kpc

Wang, Bose et al 2020

The VVV simulation

Planck cosmology

Dark matter only

Dynamic range of
30 orders of
magnitude in mass

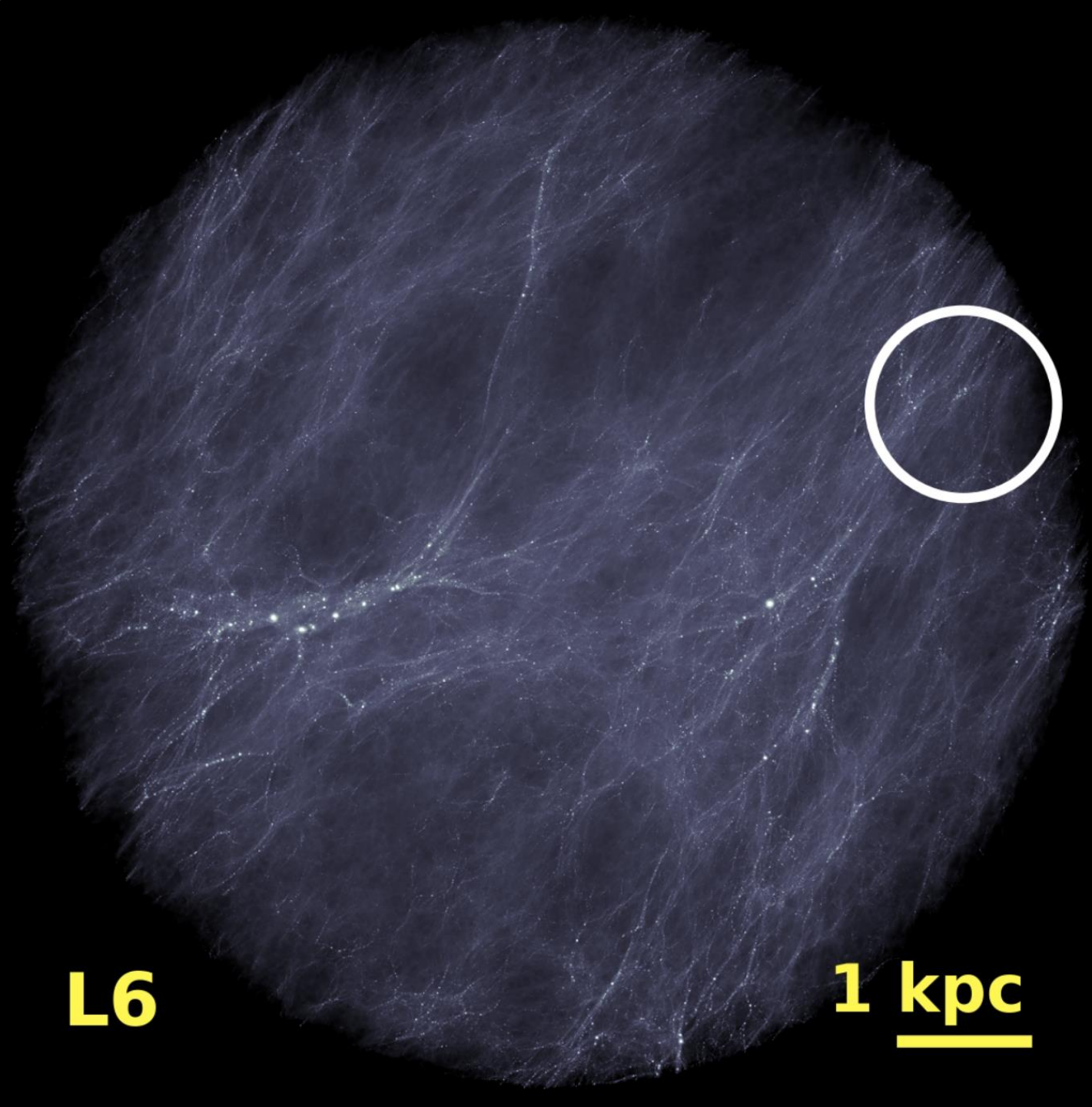
$$M_{\text{char}} = 10^{-1} M_{\odot}$$

Zoom Level 6

L6

1 kpc

Wang, Bose et al 2020



The VVV simulation

Planck cosmology

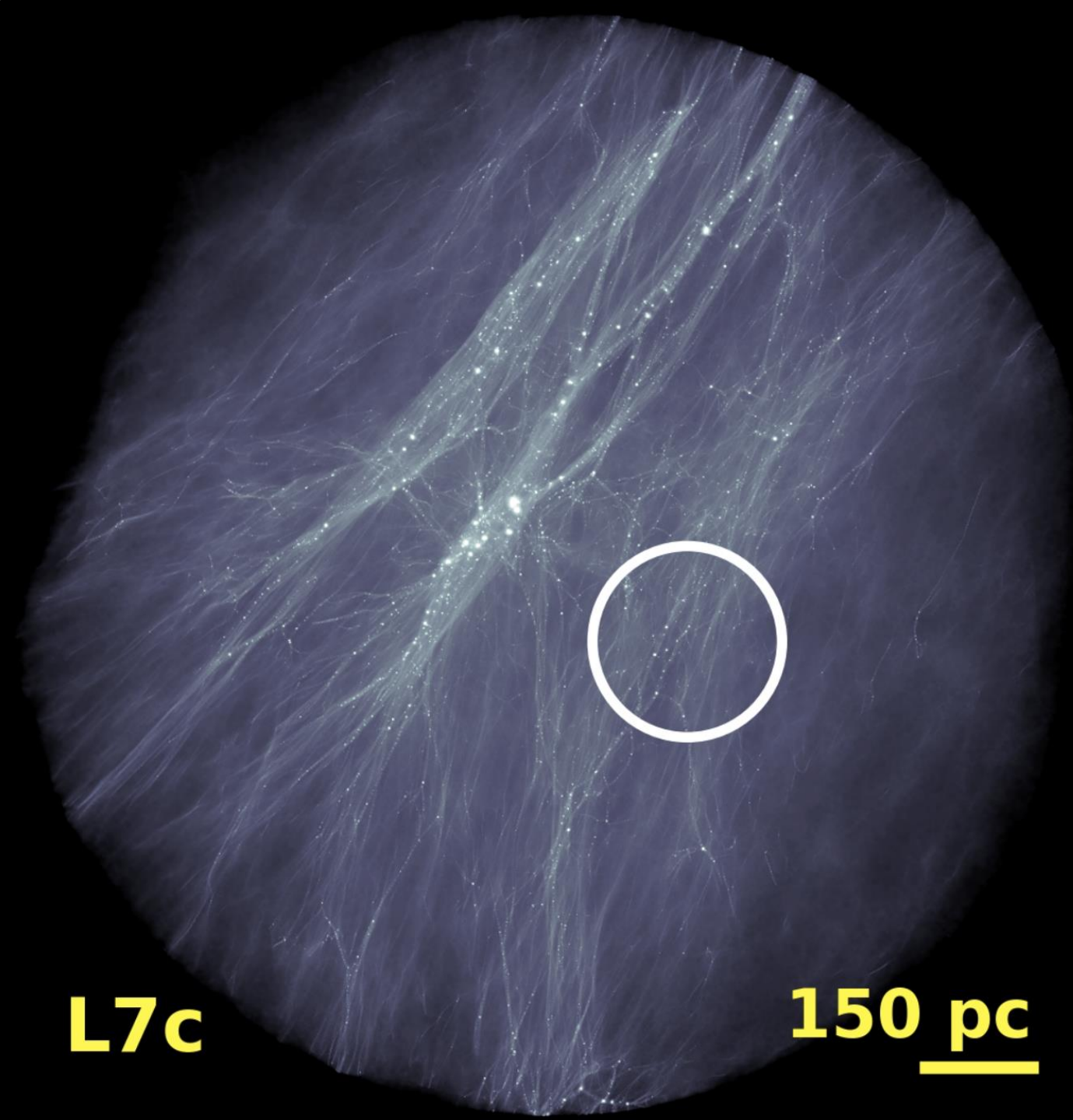
Dark matter only

Dynamic range of
30 orders of
magnitude in mass

$$M_{\text{char}} = 10^{-4} M_{\odot}$$

Zoom Level 7

Wang, Bose et al 2020



L7c

150 pc

The VVV simulation

Planck cosmology

Dark matter only

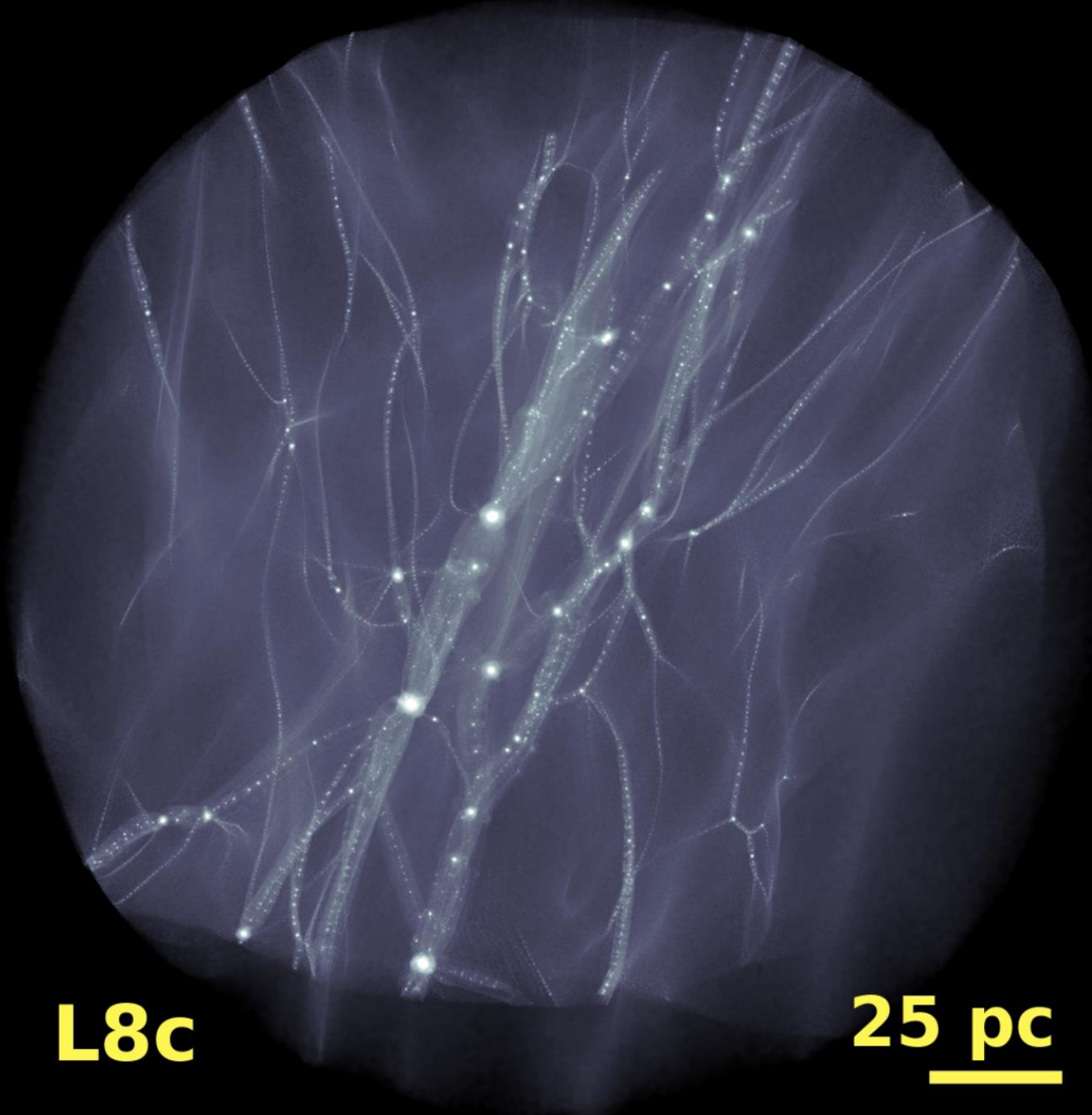
Dynamic range of
30 orders of
magnitude in mass

$$M_{\text{char}} = 10^{-6} M_{\odot}$$

Zoom Level 8

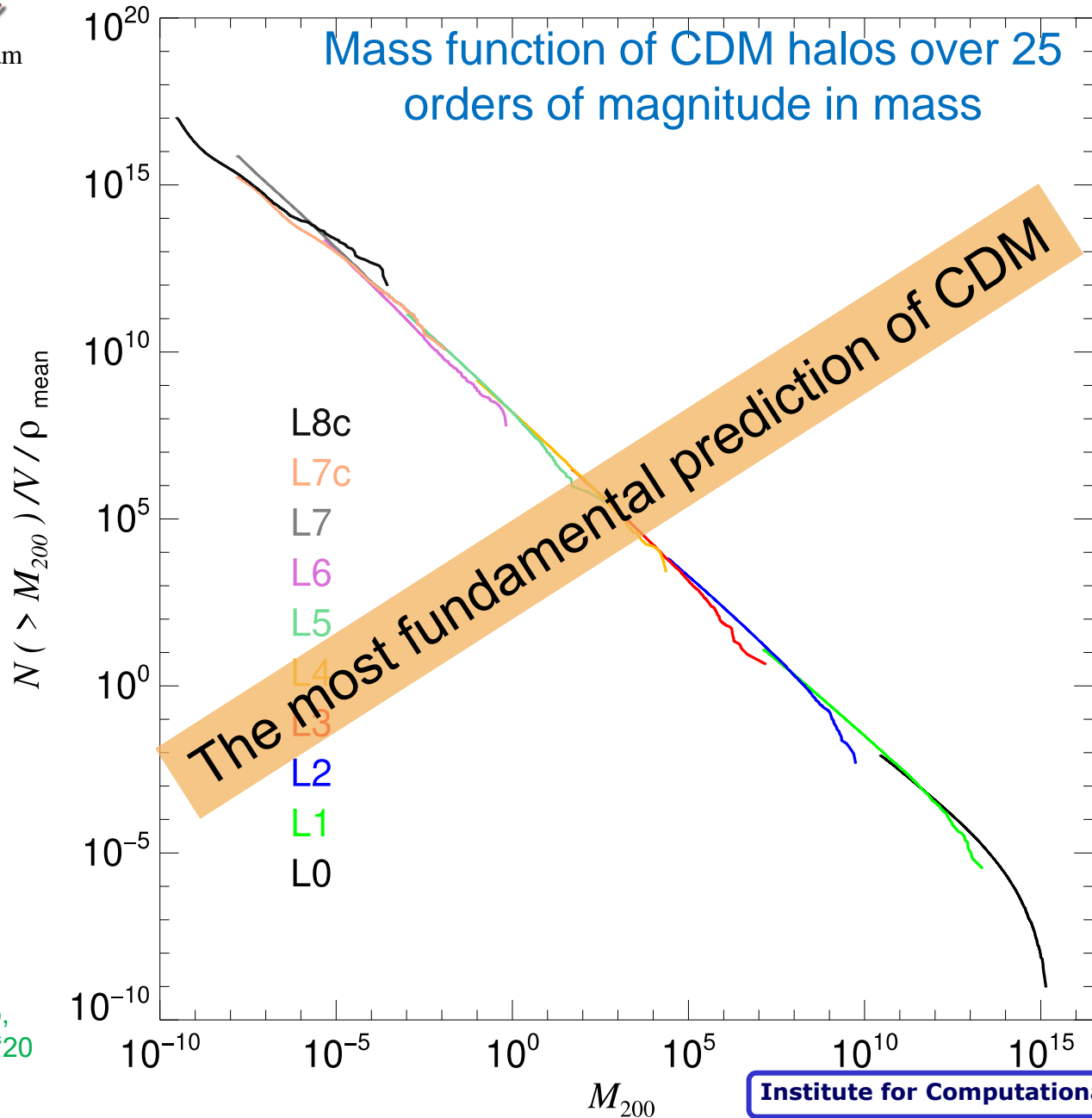
The density of
this region is
only $\sim 3\%$ of the
cosmic mean

Wang, Bose et al 2020



L8c

25 pc

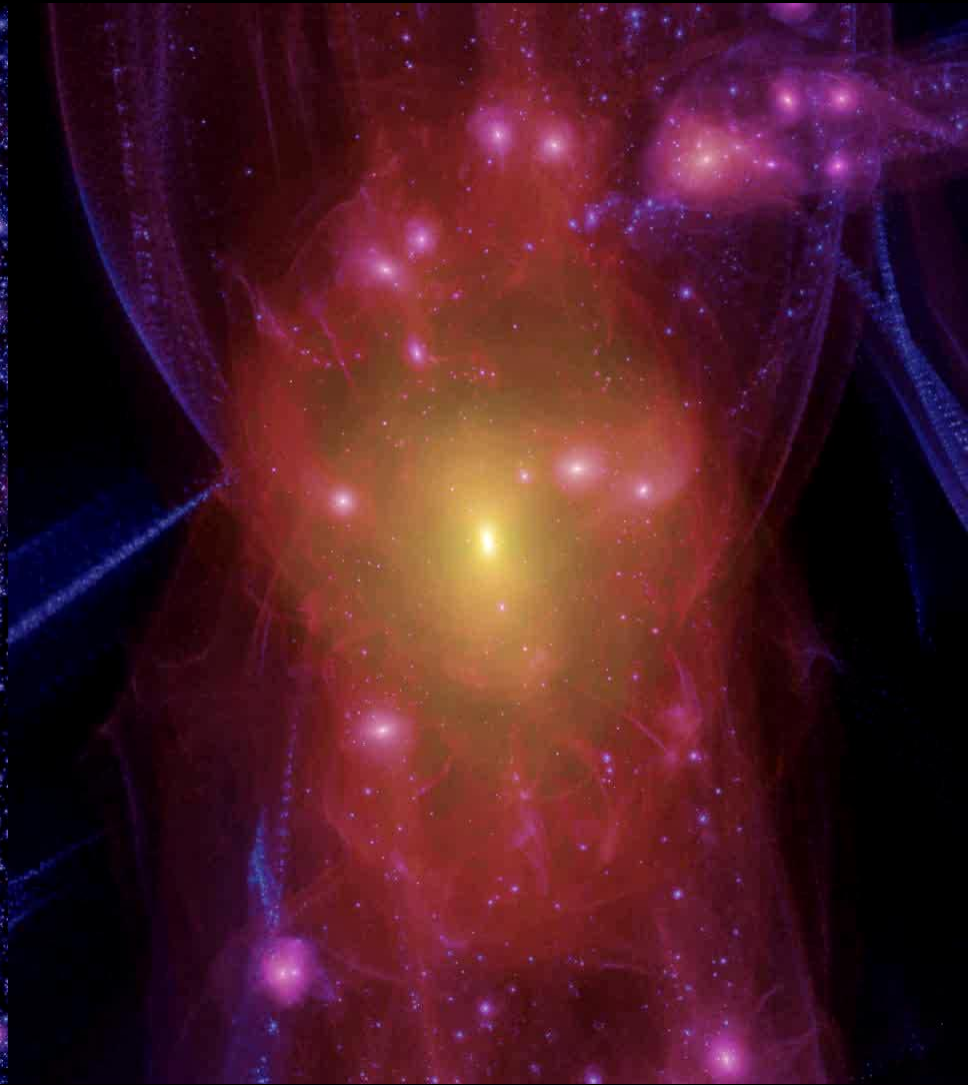


Wang, Bose, CSF, Gao,
Jenkins, Springel, White '20

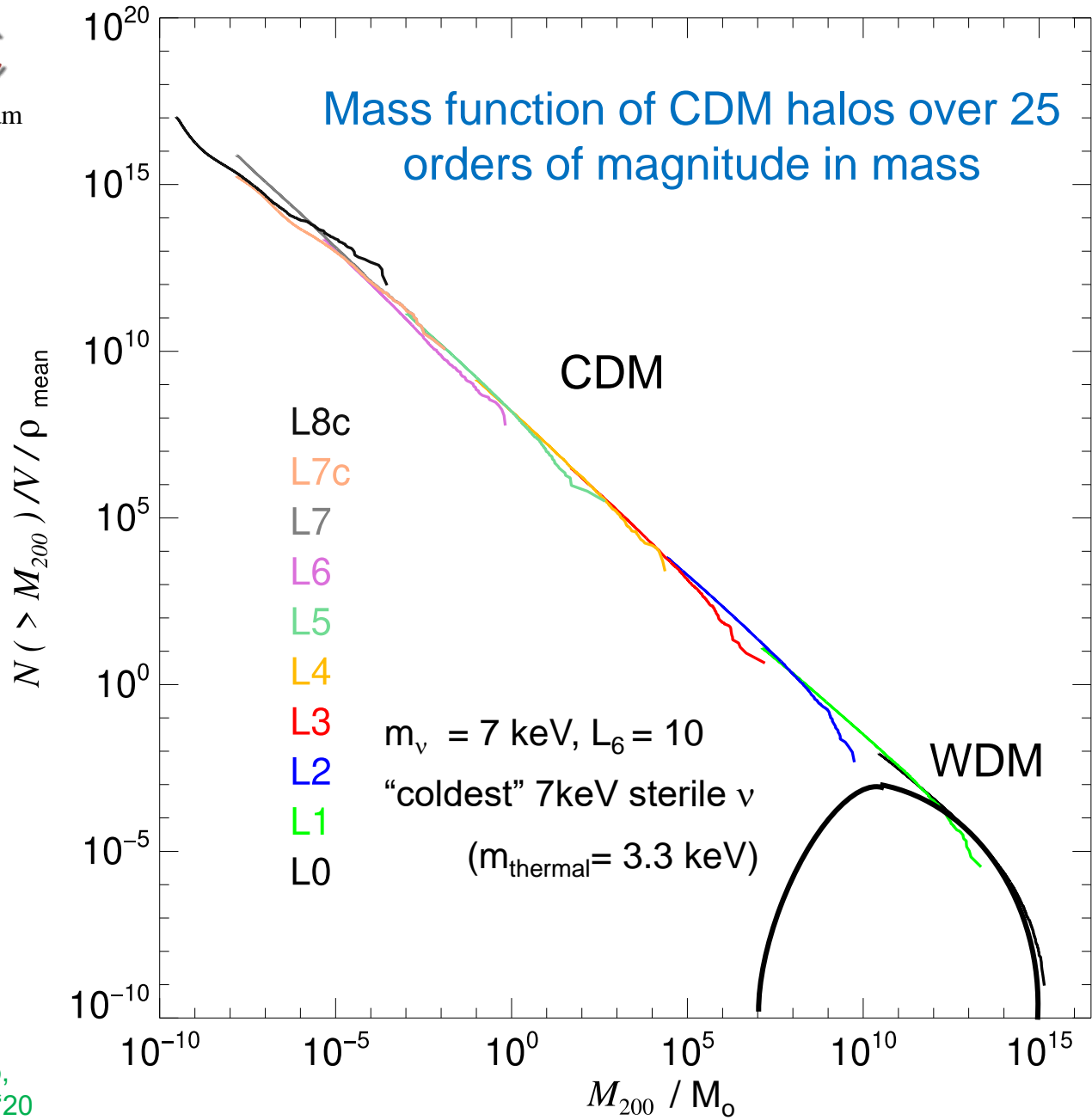
cold dark matter



warm dark matter



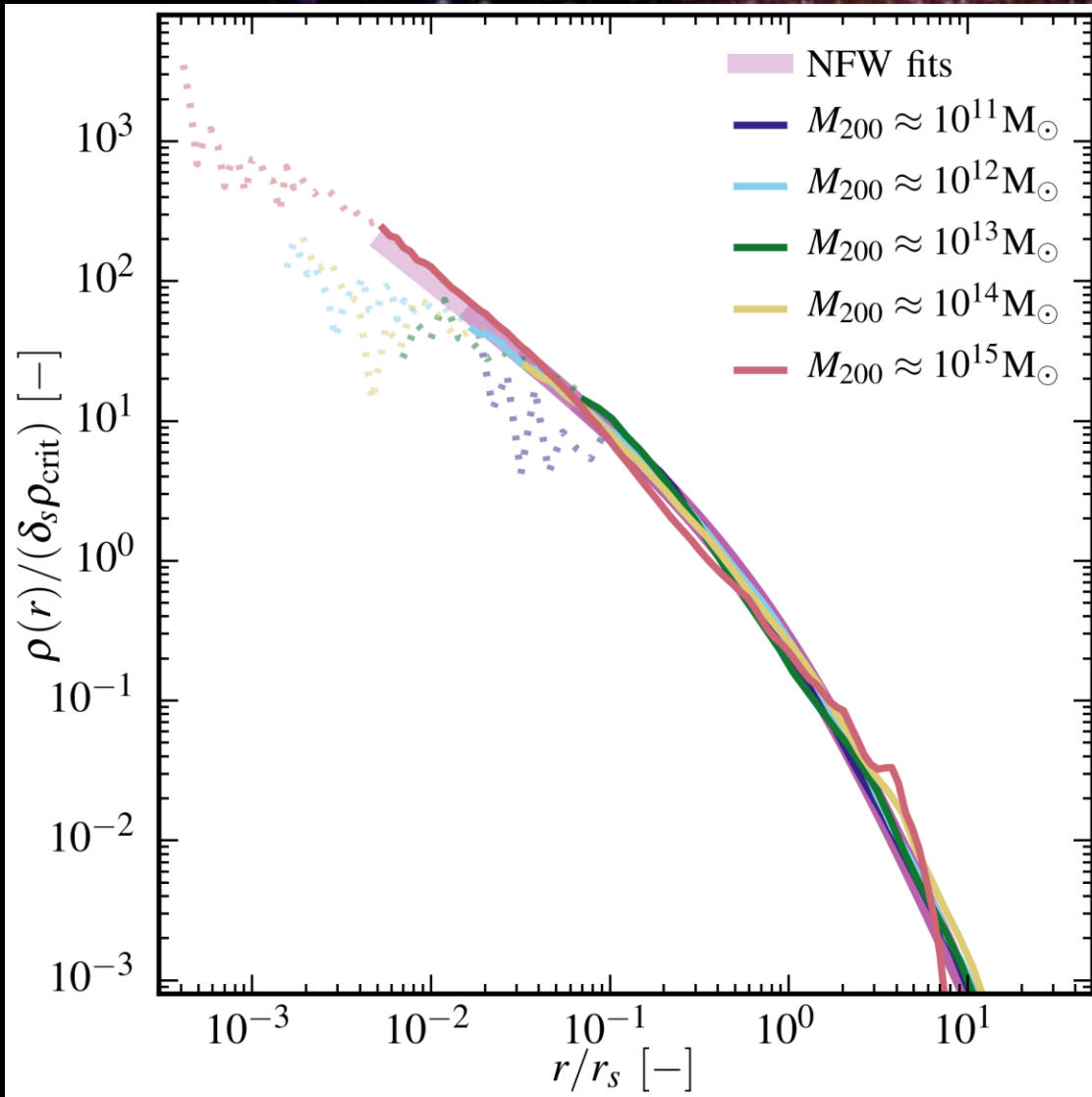
Lovell, Eke, Frenk, Gao, Jenkins, Wang, White, Theuns,
Boyarski & Ruchayskiy '12





The structure of dark matter
halos of all masses

The Density Profile of Cold Dark Matter Halos



Shape of halo profiles
~independent of halo mass &
cosmological parameters

Density profiles are “cuspy” -
no ‘core’ near the centre

Fitted by simple formula:

$$\frac{\rho(r)}{\rho_{\text{crit}}} = \frac{\delta_c}{(r/r_s)(1+r/r_s)^2}$$

(Navarro, Frenk & White '97)

More massive halos and
halos that form earlier have
higher densities (bigger δ)

Universal halo density profiles

$$\frac{\rho(r)}{\rho_{crit}} = \frac{\delta_c}{(r/r_s)(1+r/r_s)^2}$$

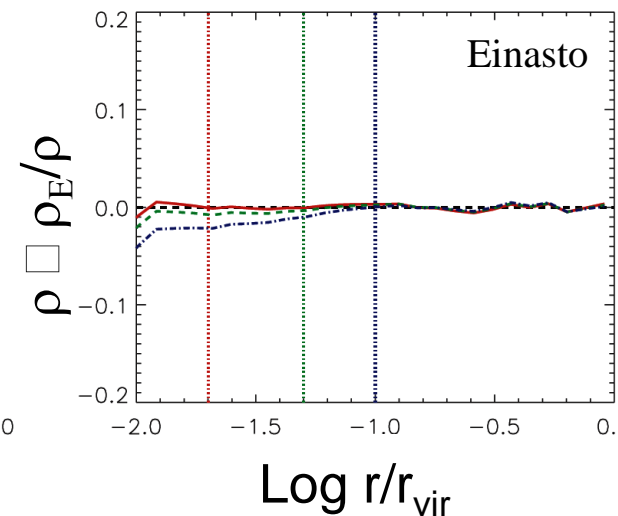
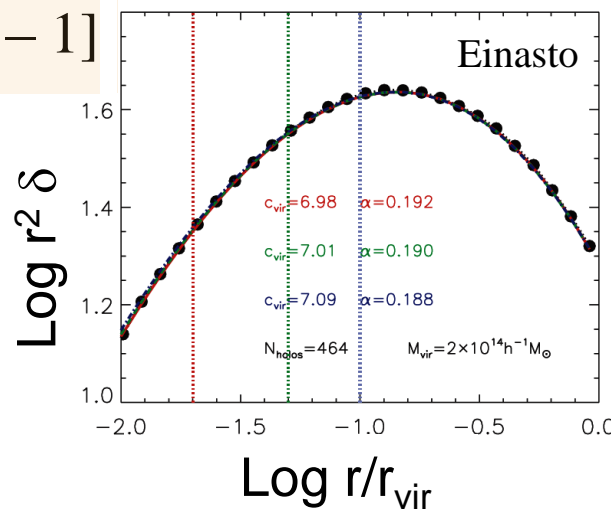
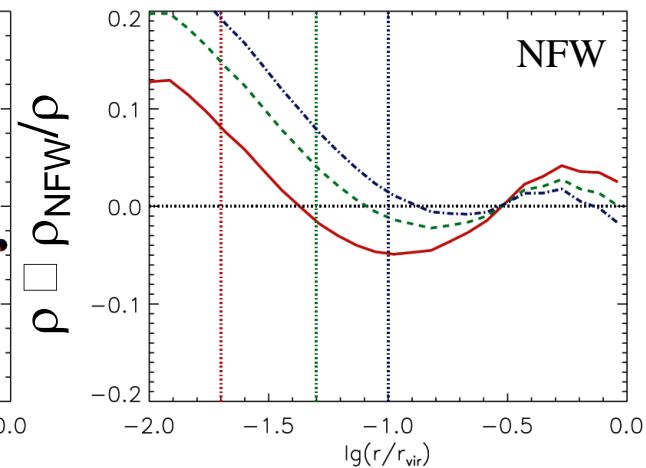
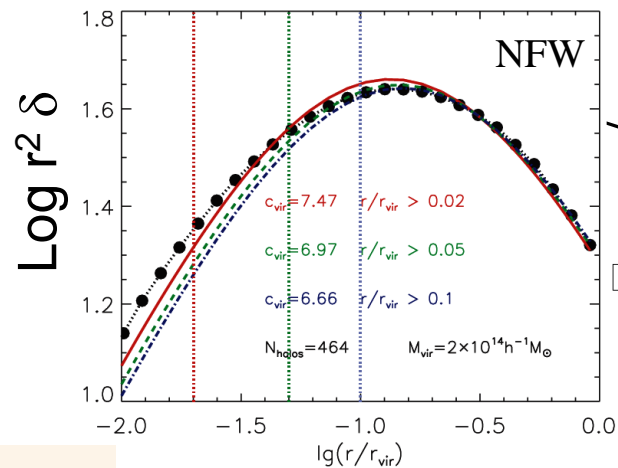
The “Einasto” formula

$$\ln(\rho(r)/\rho_{-2}) = (-2/\alpha) [(r/r_{-2})^\alpha - 1]$$

Fits mean profiles even better

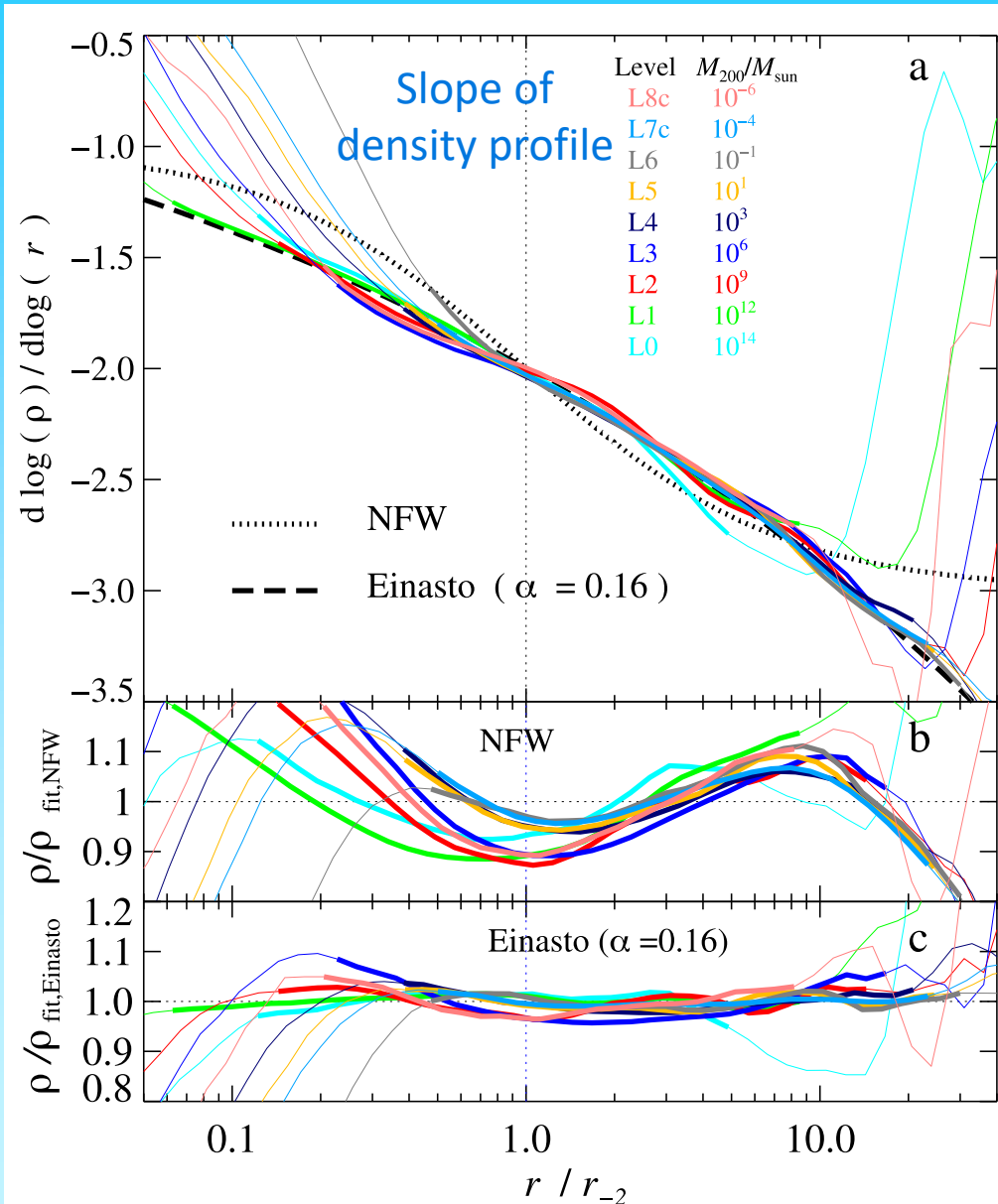
Gao, N, F, W + 2008

Averaged cluster mass halos fit with NFW and Einasto



Density profile shapes

Over **20 orders** of magnitude in halo **mass** and 4 orders of magnitude in density, the mean density **profiles** of halos are **fit** by **NFW** to within **20%** and by **Einasto** ($\alpha = 0.16$) to within **7%**

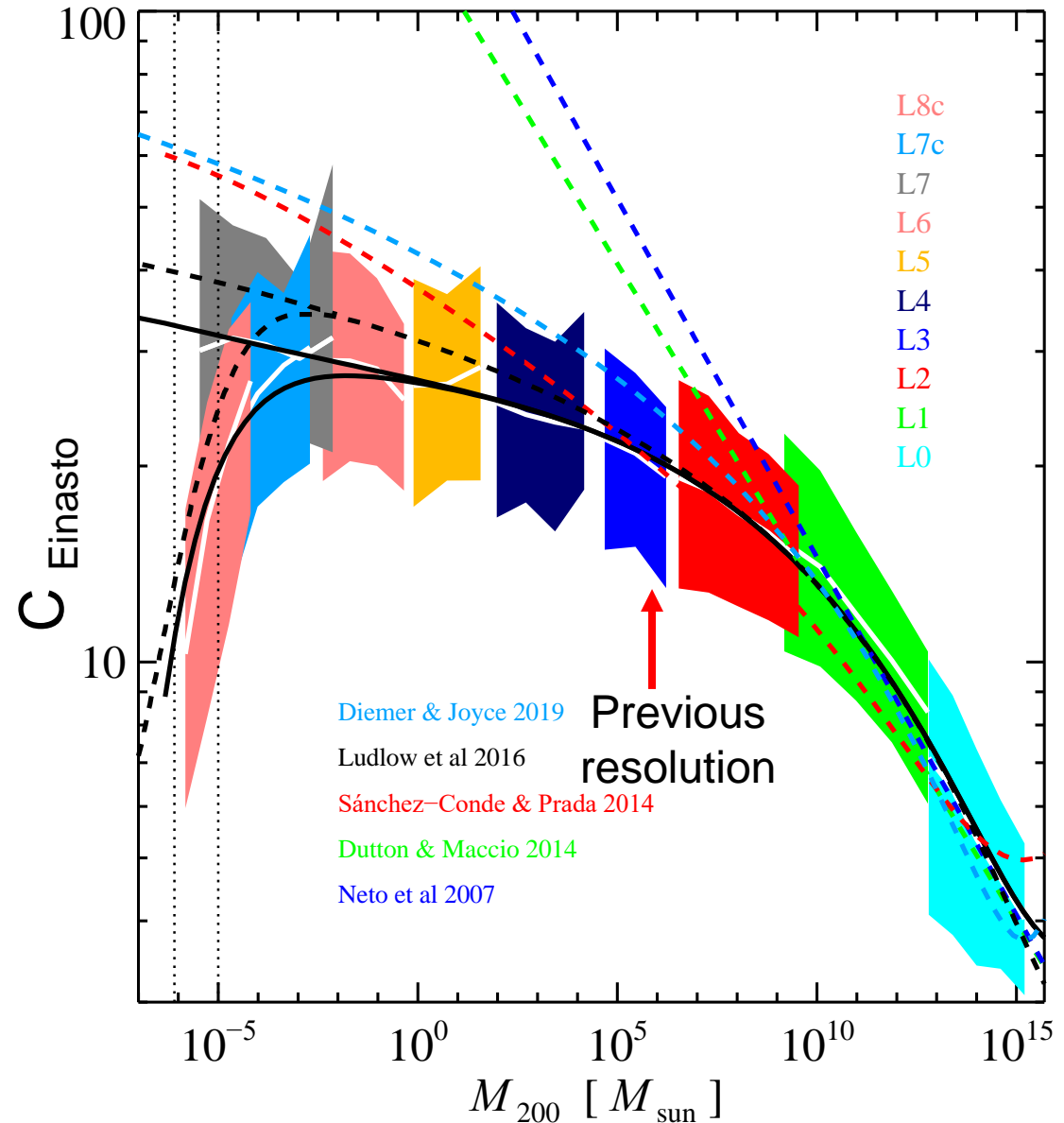


Concentration-mass relation

Concentrations at small mass are **lower** than all previous extrapolations by up to factors of tens.

A **turndown** at 10^3 Earth masses is due to the **free-streaming limit**.

The **scatter** depends only weakly on halo mass



A digression: the cosmic neutrino background

Neutrinos are the only non-baryonic form of **dark matter** known!

They make a small contribution, $\sim 1\%$, to the dark matter

Cosmic neutrinos were produced a few seconds after the Big Bang and produce a **cosmic background** today

May be detected by Ptolomy experiment

The cosmic neutrino background

Willem Elbers

Elbers, CSF, Jenkins, Li, Pascoli, Lavaux, Jasche, Springel '23

Constrained realization simulations

Simulations from CDM initial conditions, with phases adjusted to reproduce the local observed galaxy clustering

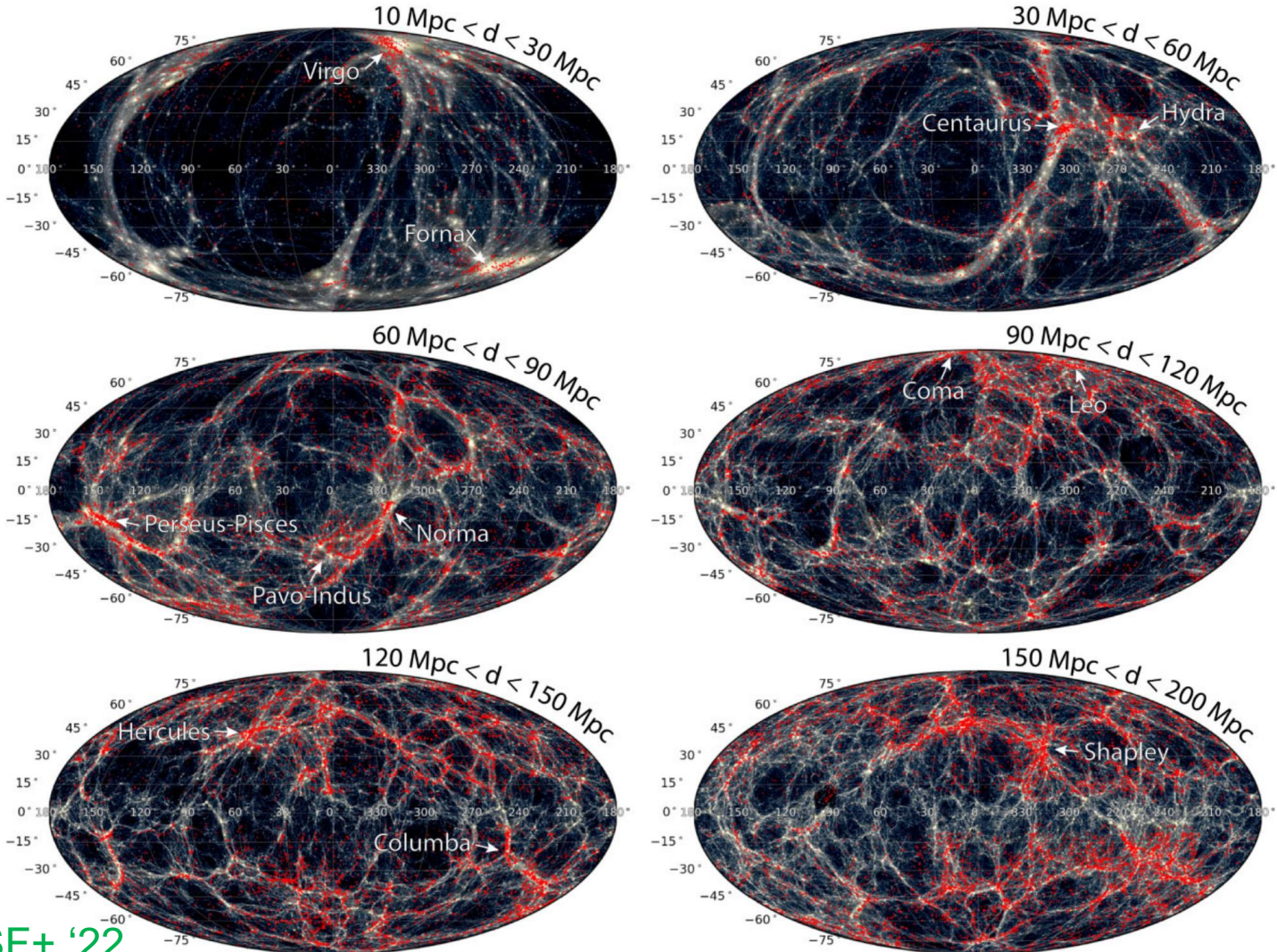
Constrained by 2MASS++ survey

Constrained simulations

SIBELIUS DARK

Grey: dark matter
in SIBELIUS DARK

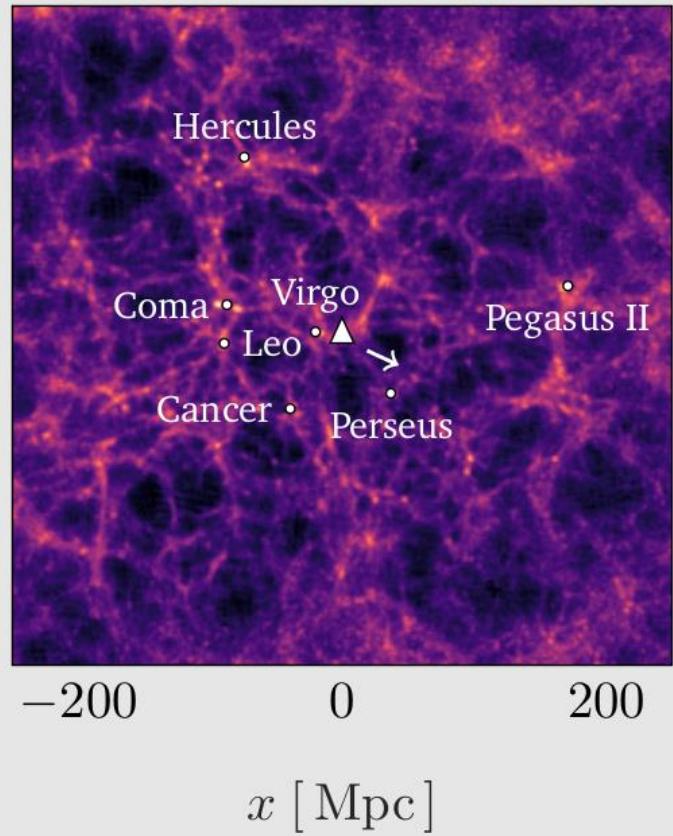
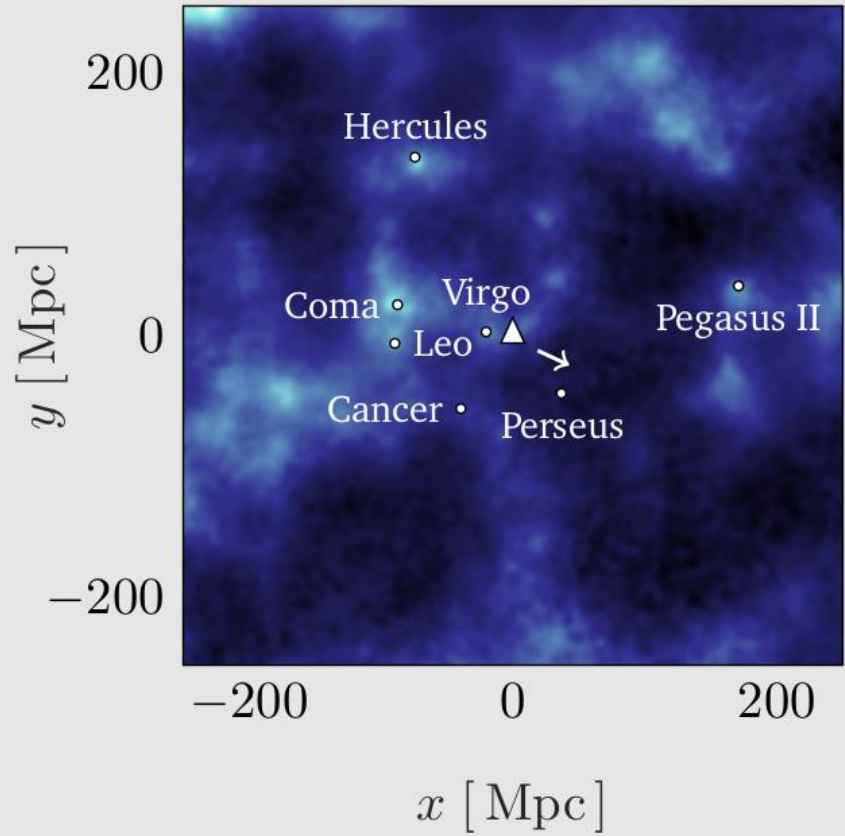
Red: galaxies in
2M++ survey,
used for
reconstruction



Constrained simulations with ν_s

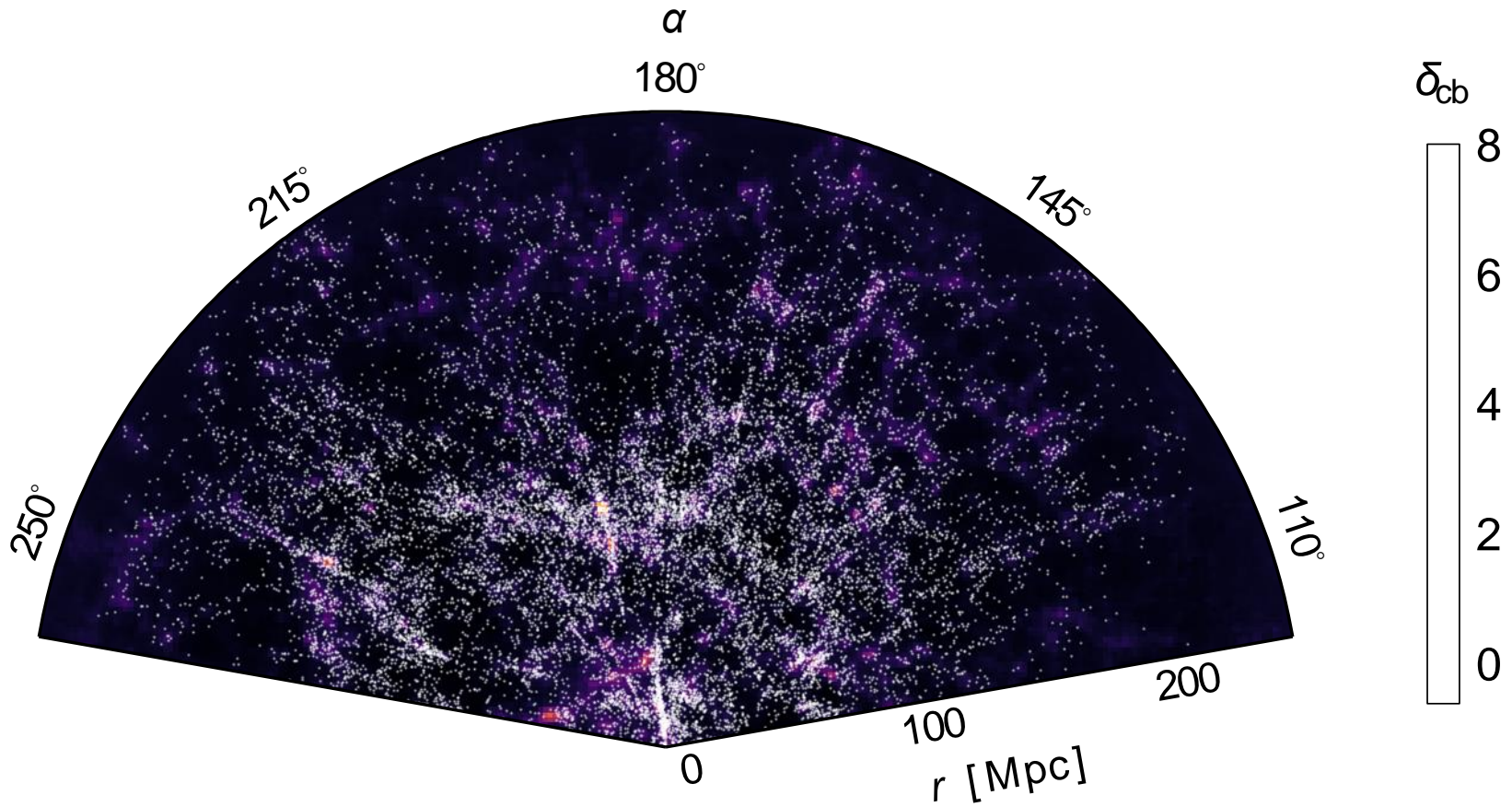
$m_\nu = 0.05 \text{ eV}$

CDM

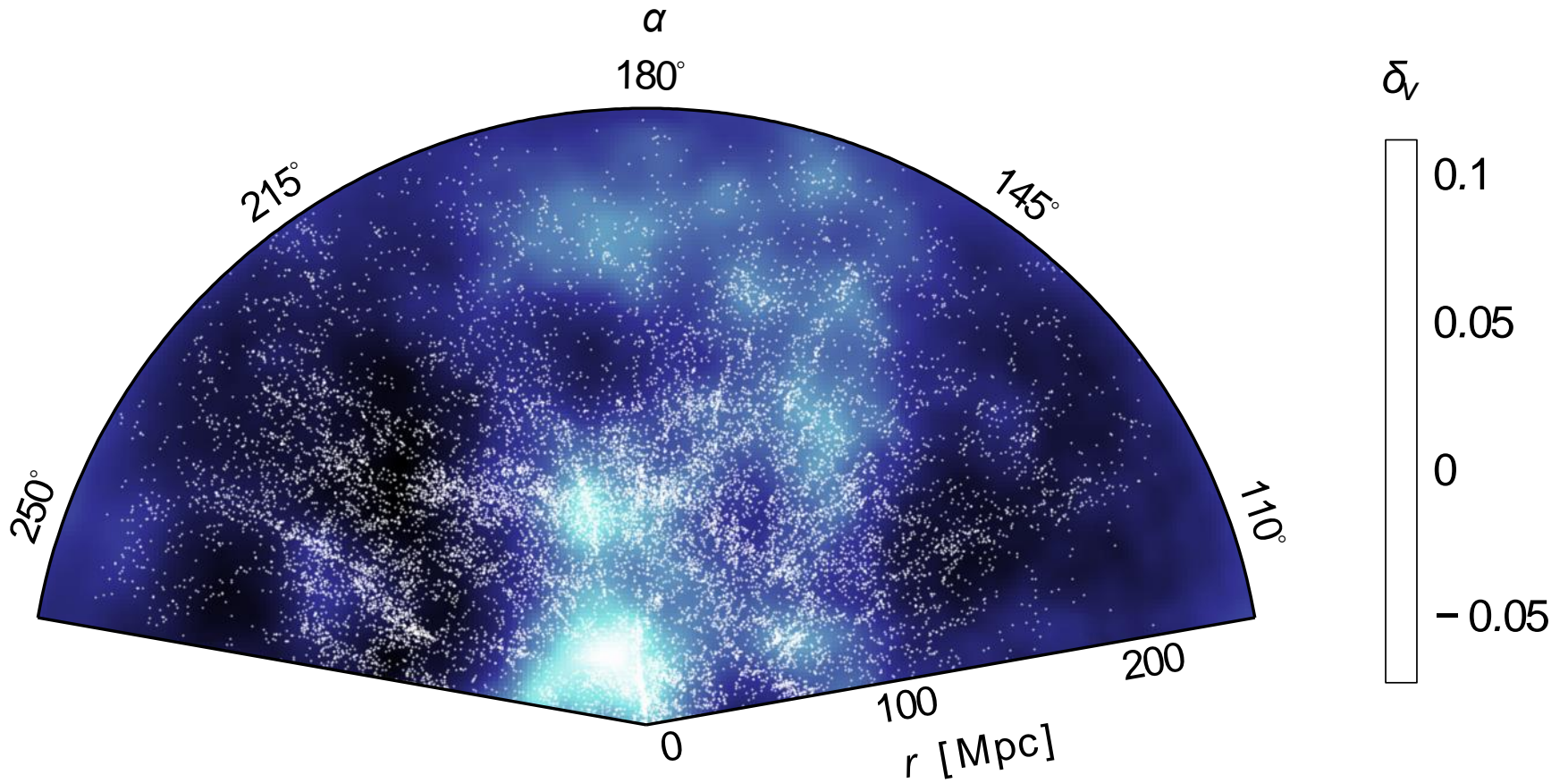


The position of the Milky Way is indicated by a white triangle

Local CDM distribution in z-space

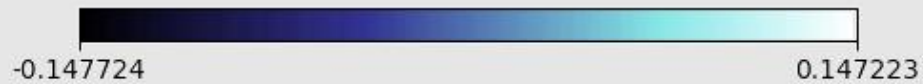
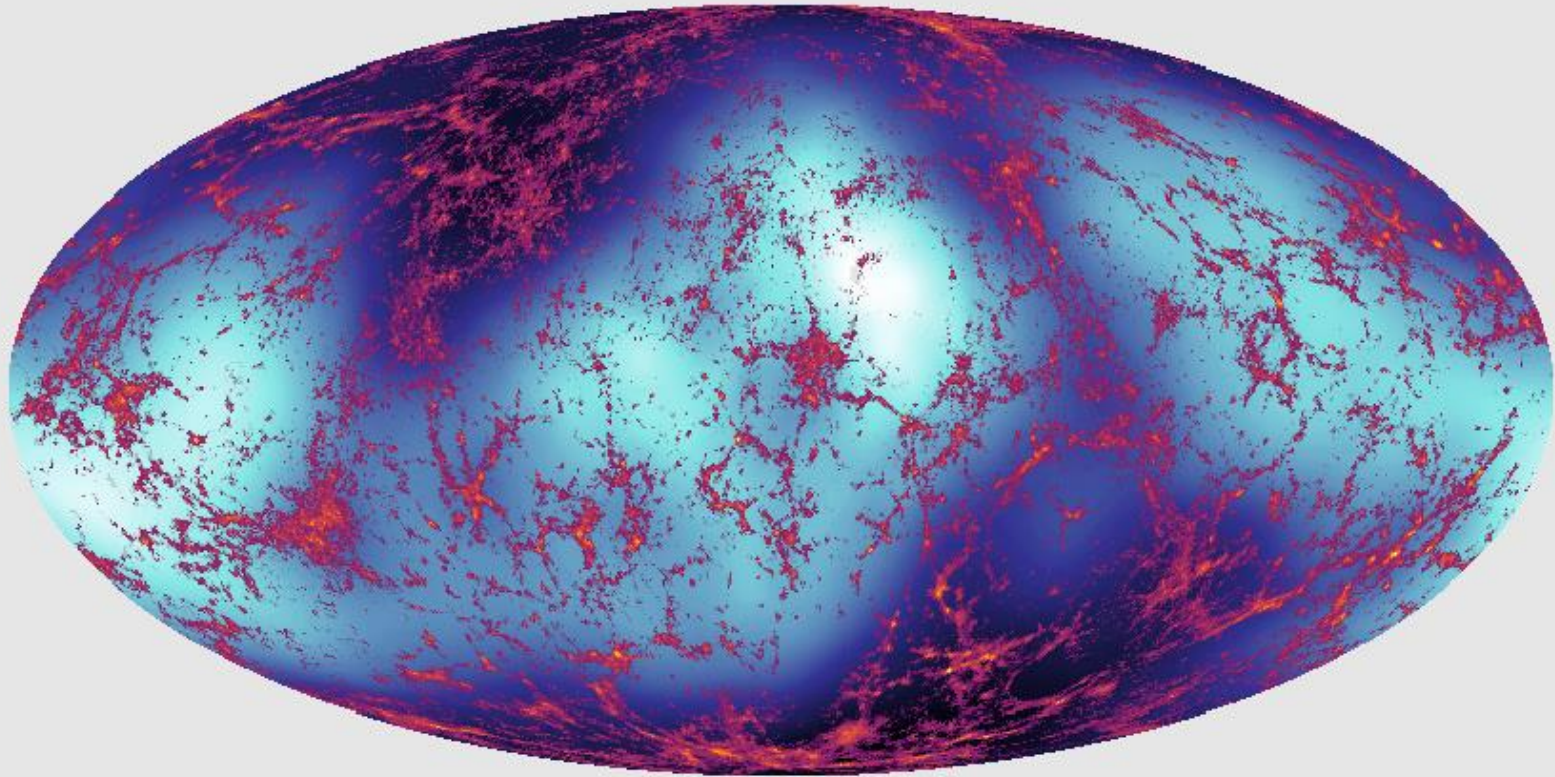


Local v distribution in z -space

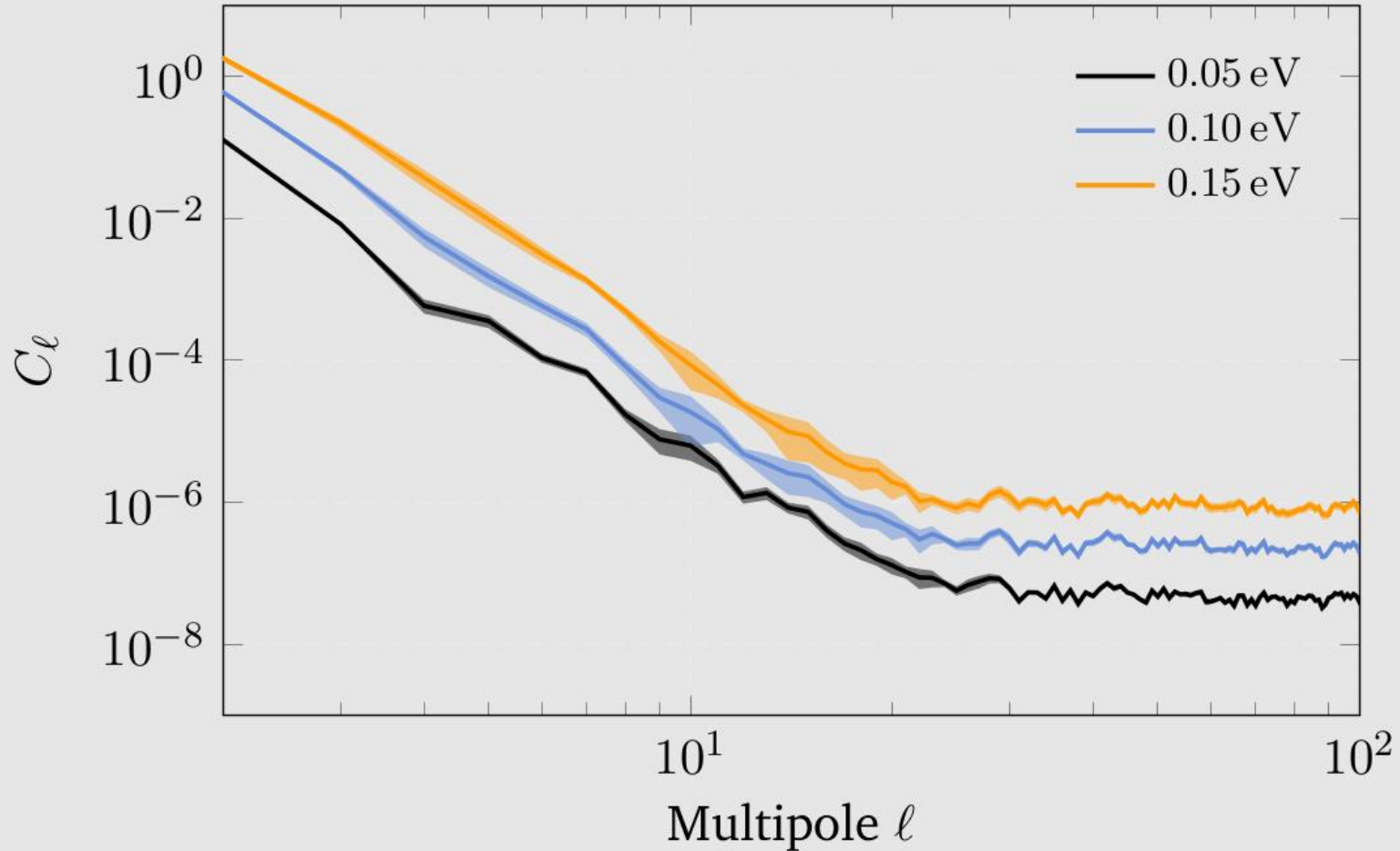


Angular anisotropies

Local neutrino density perturbations without dipole:



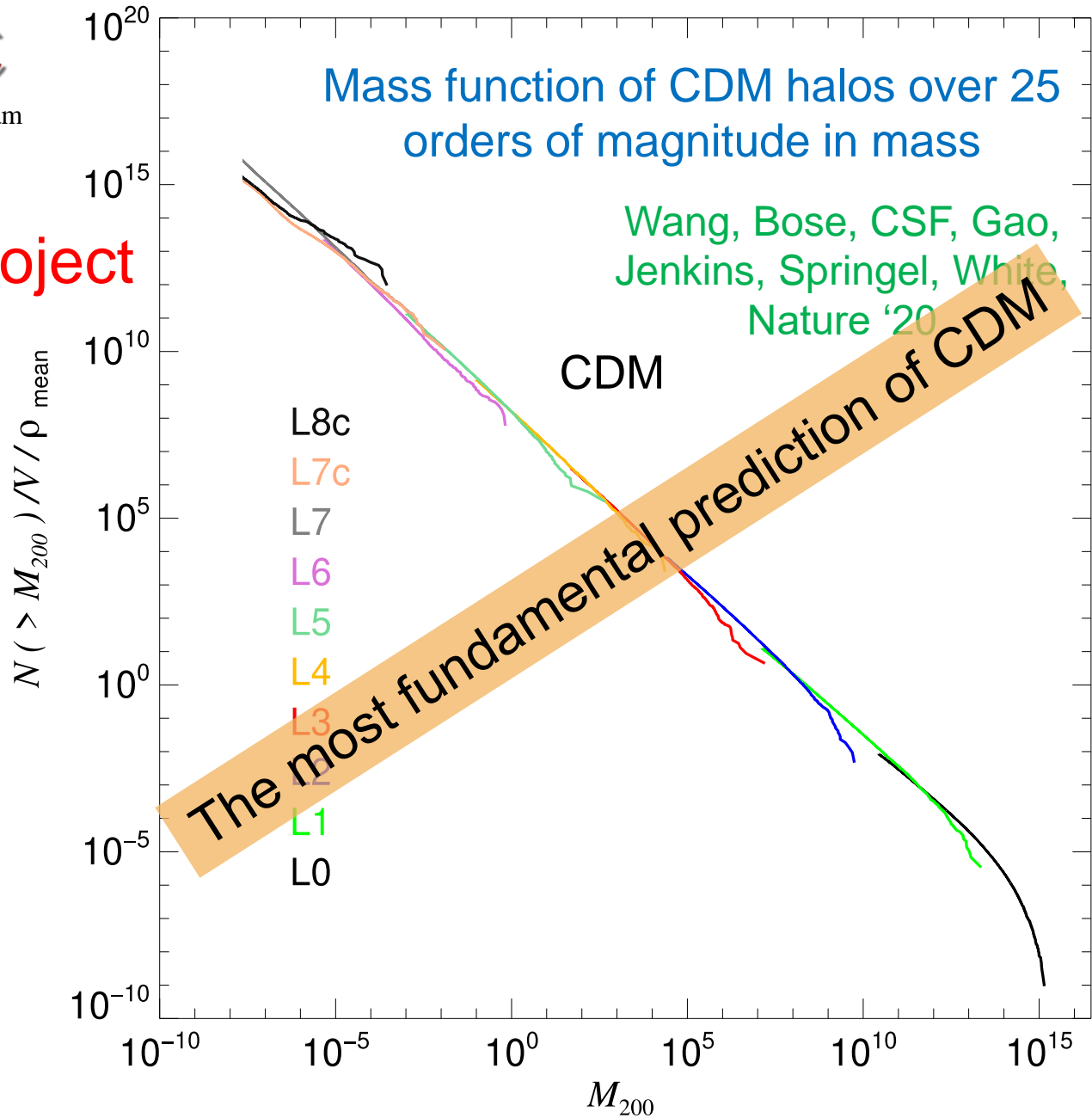
Angular power spectrum of ν perturbations





A conclusive test of CDM

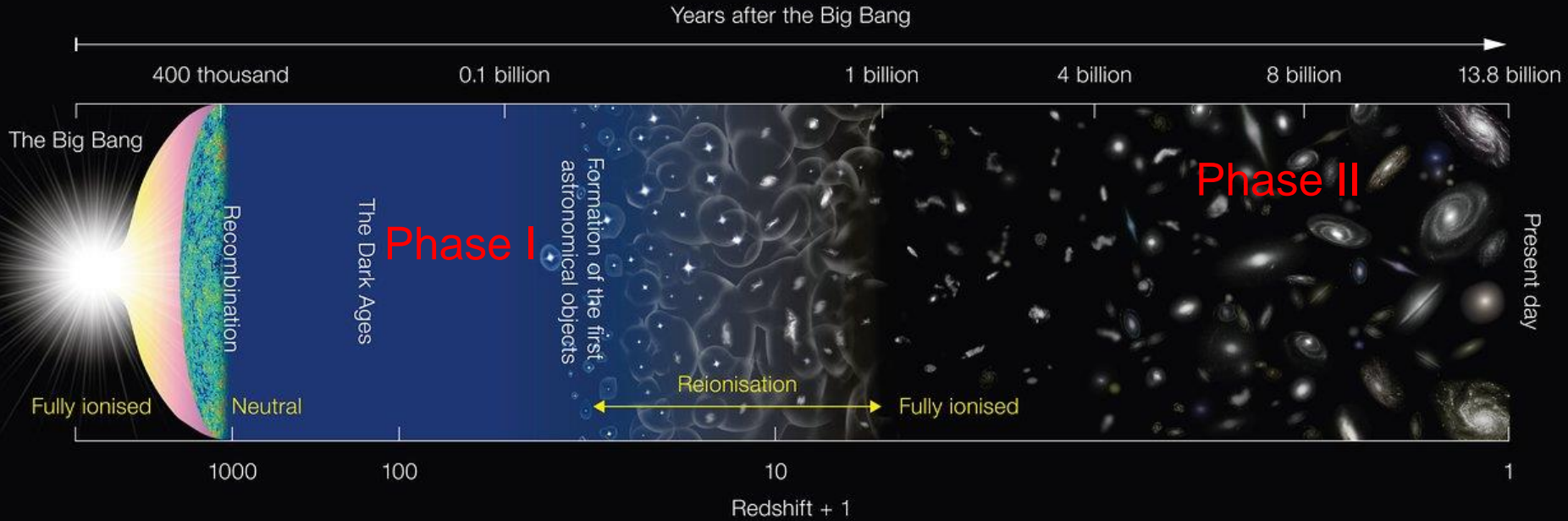
The VVV project



CDM

Most subhalos never make a galaxy!

The two phases of galaxy formation



Phase I: H gas is neutral \rightarrow can only cool in halos $m > m_{thr,1}$

First stars reionize H and heat it up to 10^4K

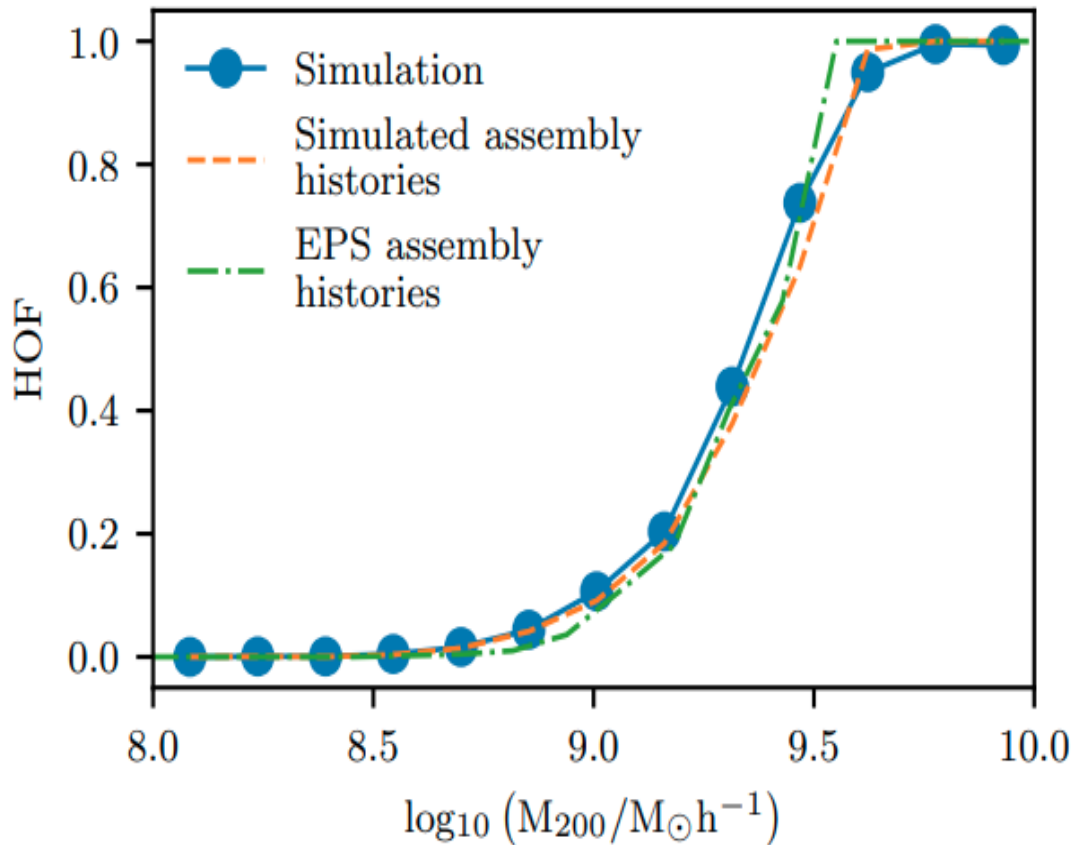
Phase II: H Gas is ionized ($T_{vir} > 10^4$)

\rightarrow can only cool in halos $m > m_{thr,2}$

A galaxy formation primer

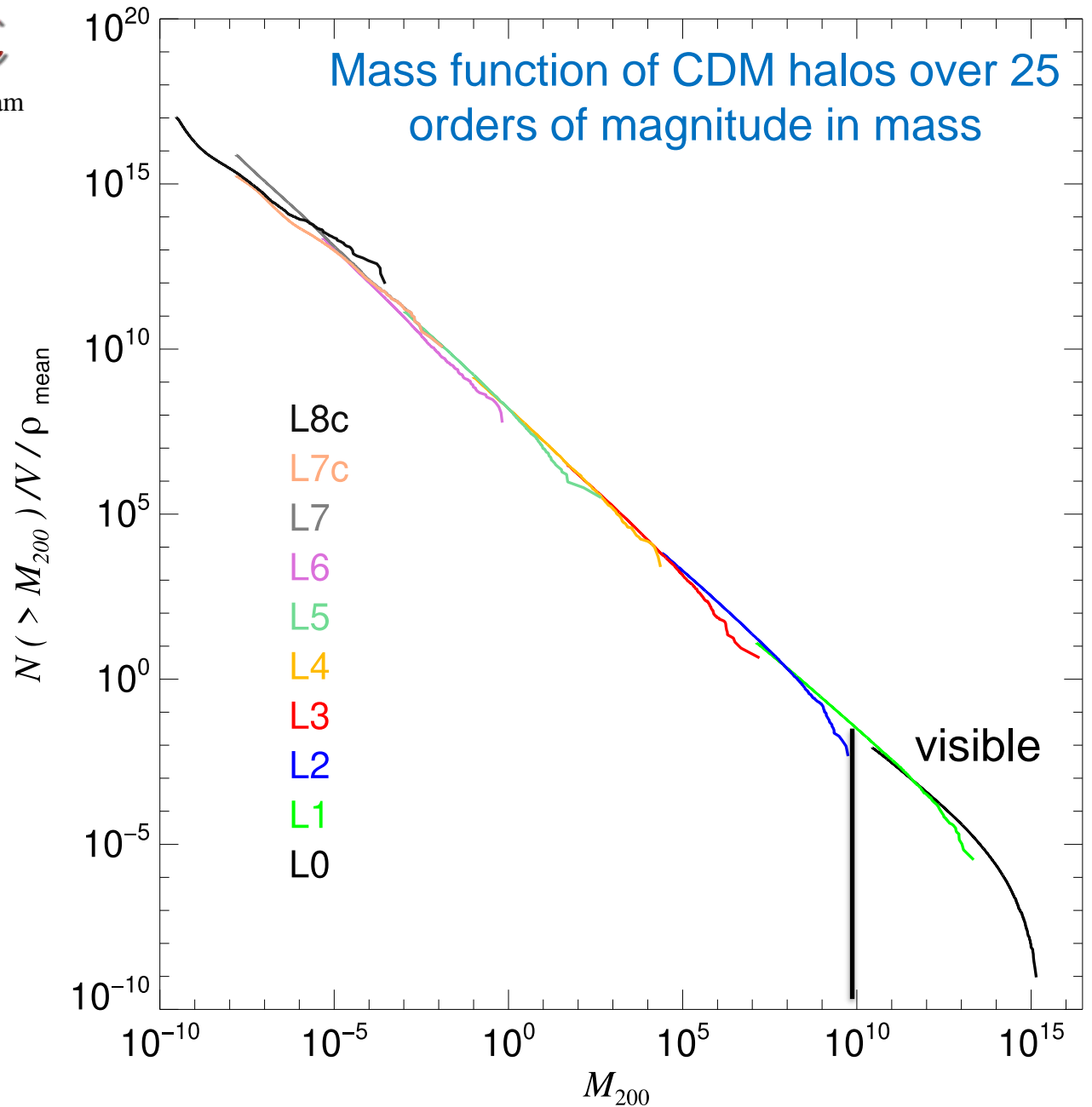
Halo Occupation Fraction (HOF): fraction of halos of a given mass today that host a galaxy

fraction of halos that host a galaxy



$M < 3 \times 10^8 M_{\odot}$
 → dark

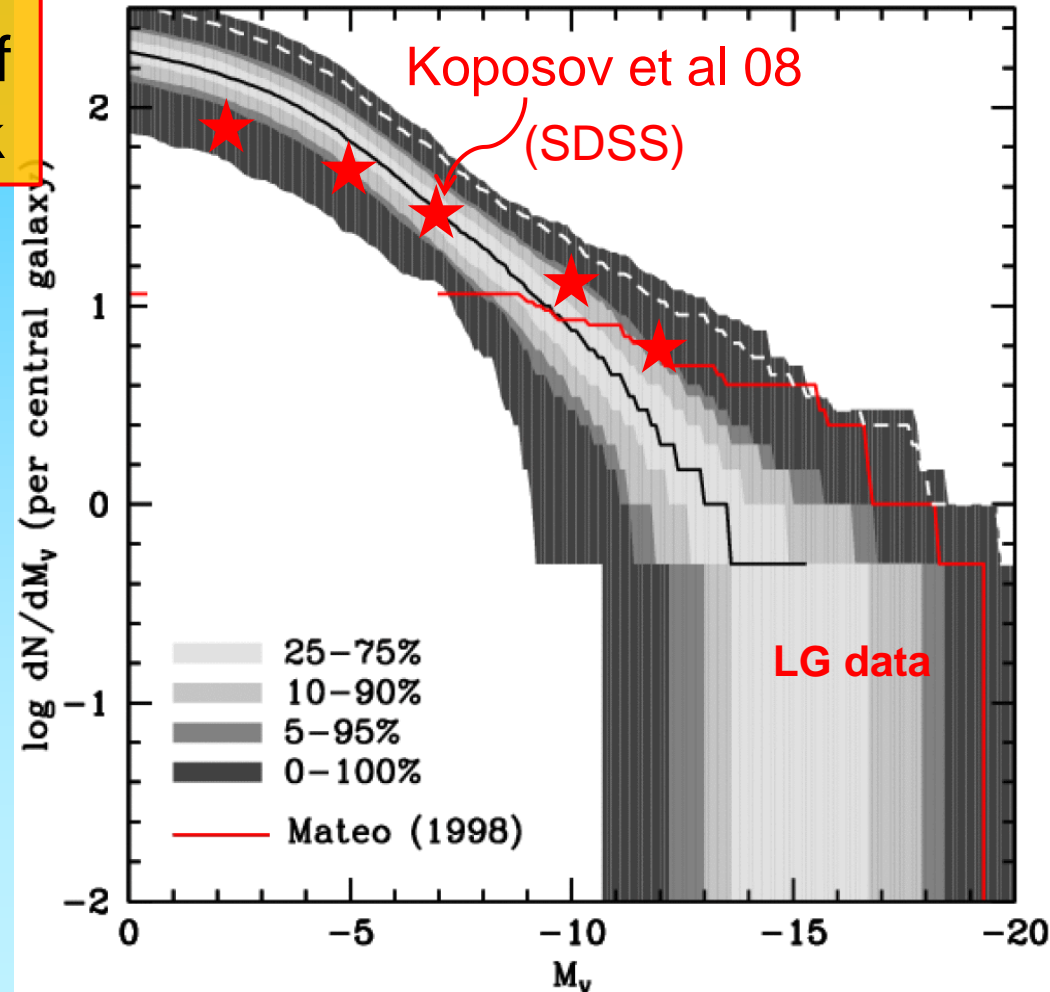
$M > 3 \times 10^9 M_{\odot}$
 → visible



Luminosity Function of Local Group Satellites

Semi-analytic model of galaxy formation including effects of reionization and SN feedback

- Median model \rightarrow correct abundance of sats brighter than $M_V = -9$ ($V_{\text{cir}} > 12$ km/s)
- Model predicts many, as yet undiscovered, faint satellites



Benson, Frenk, Lacey, Baugh & Cole '02
(see also Kauffman+ '93, Bullock+ '00, Somerville '02)

CDM

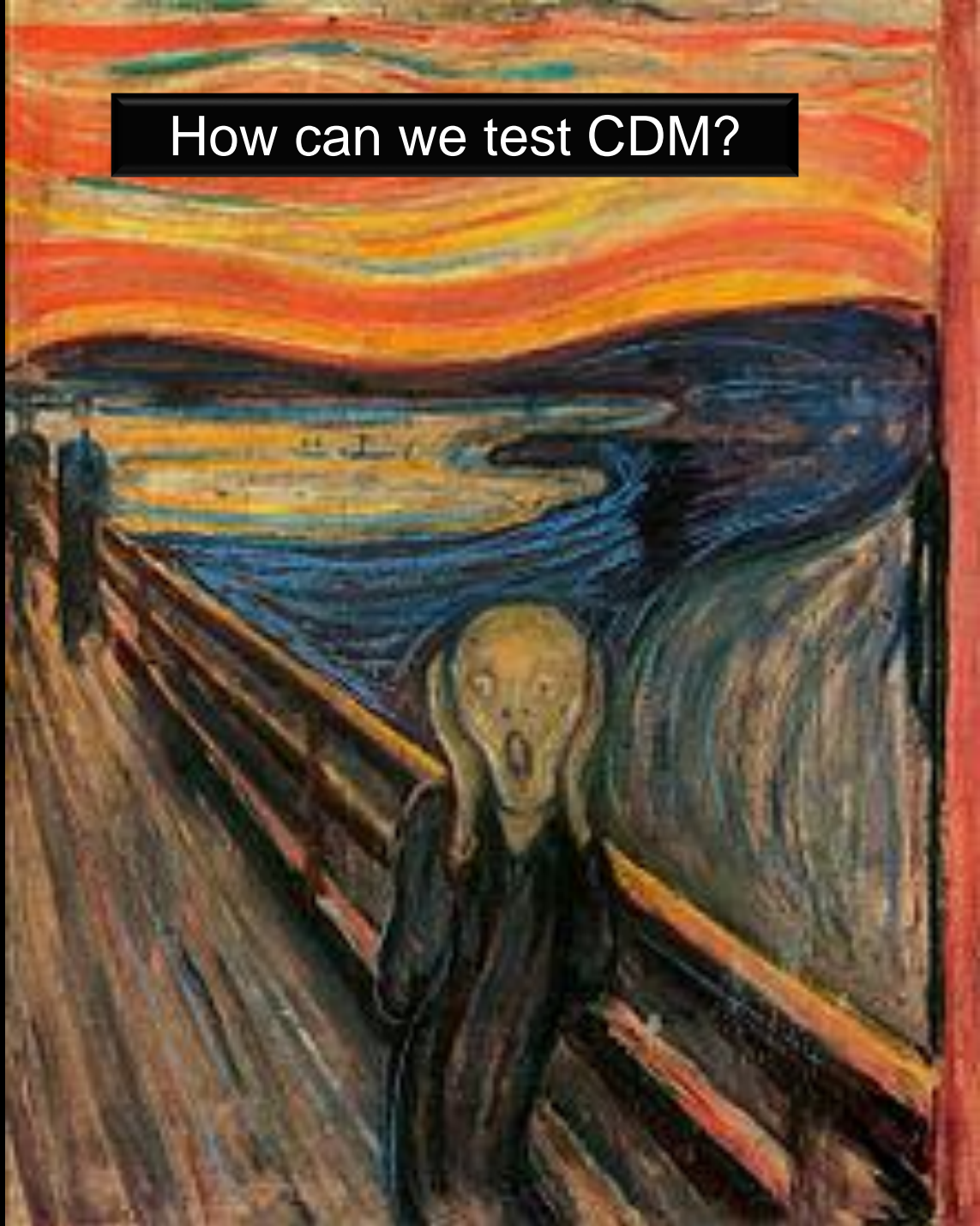
Most subhalos never make a galaxy!

→
CDM predicts the observed
abundance of satellites

→
There is **no** such thing as a “**missing
satellite problem**” in CDM!



How can we test CDM?





... and distinguish CDM/WDM?



... and distinguish CDM/WDM?

cold dark matter

warm dark matter

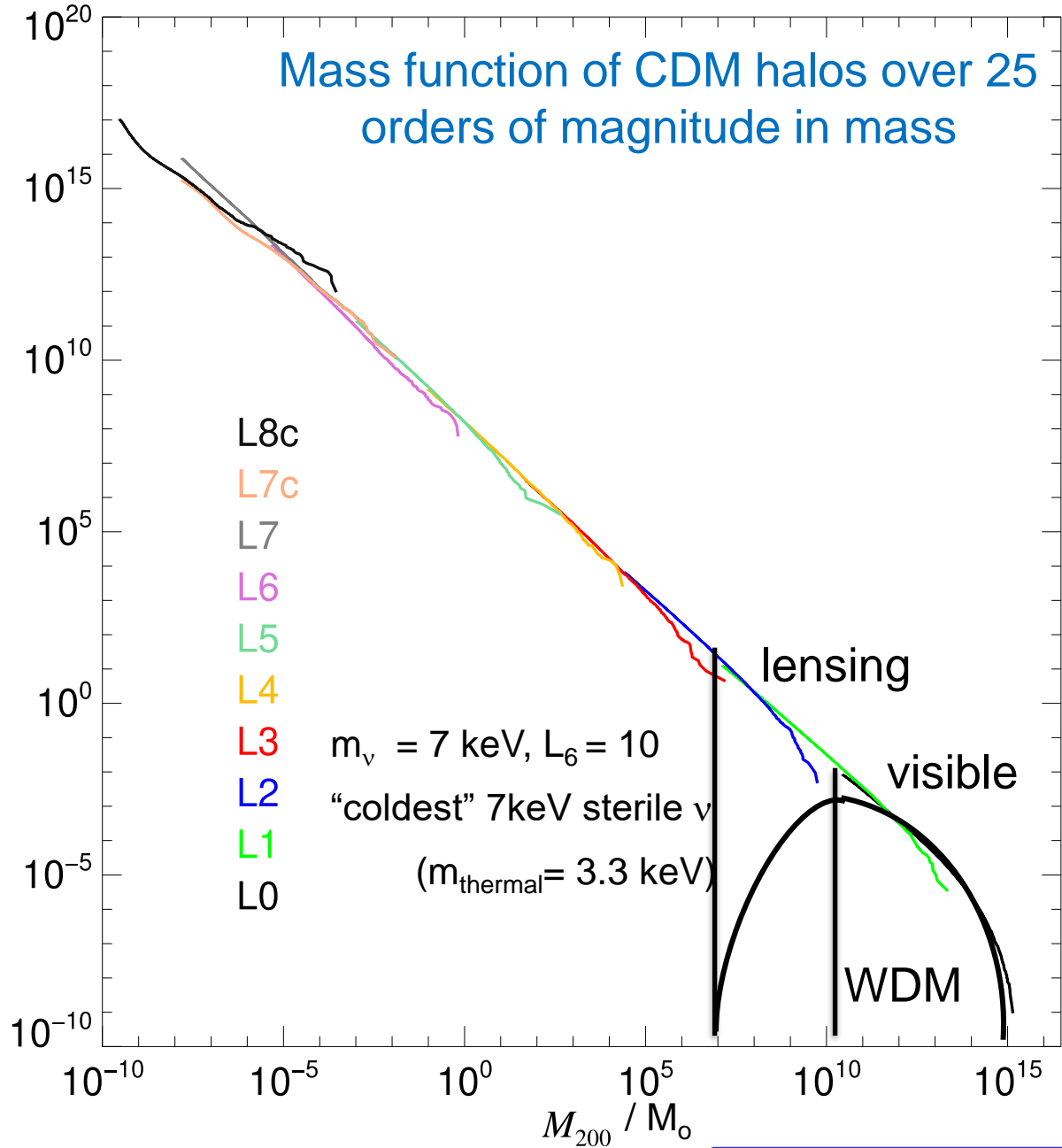
Rather than counting faint galaxies,
count the number of starless dark halos

Can we count dark haloes?

cold dark matter

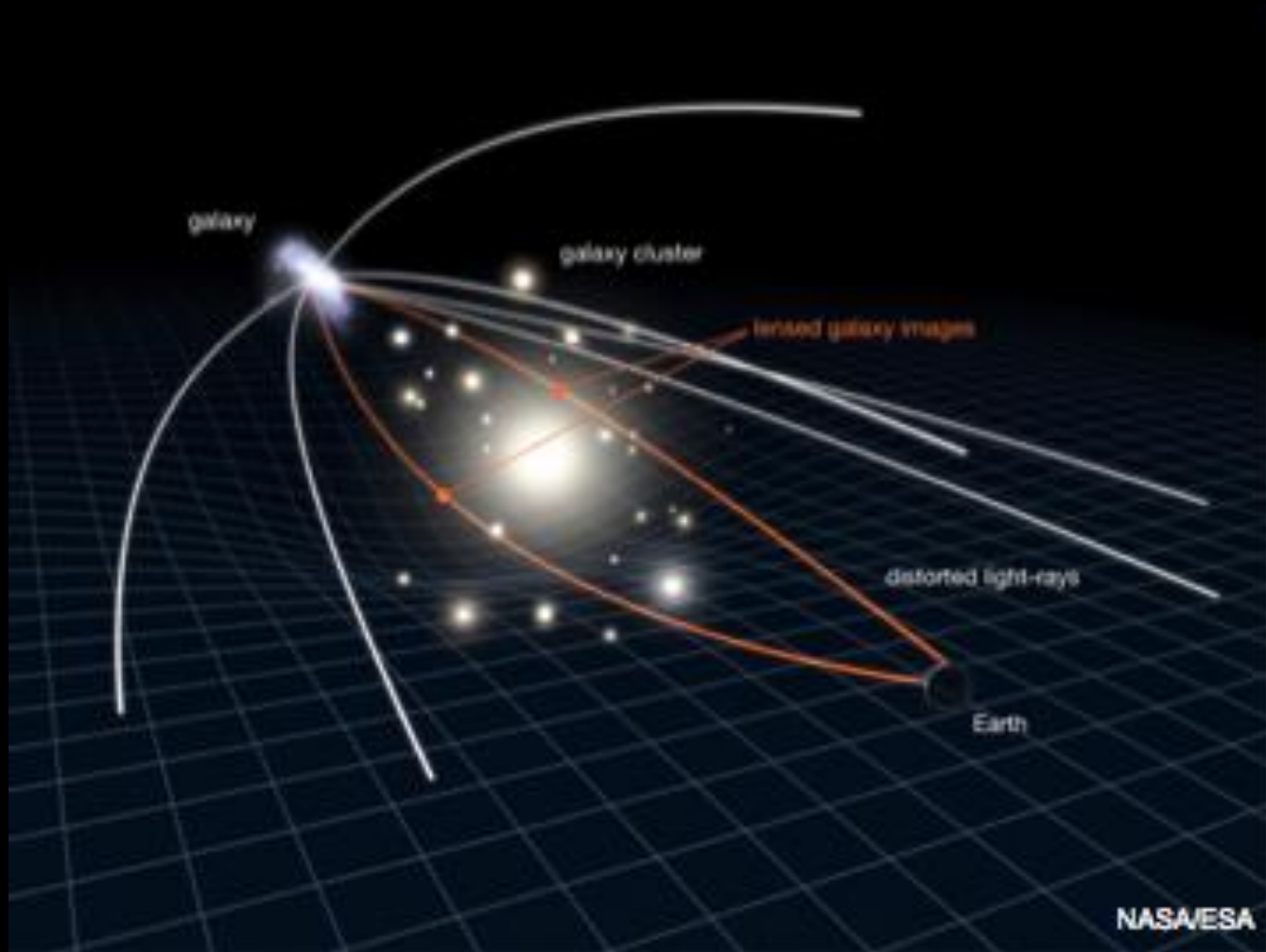
warm dark matter

→ Gravitational lensing



Wang et al '20

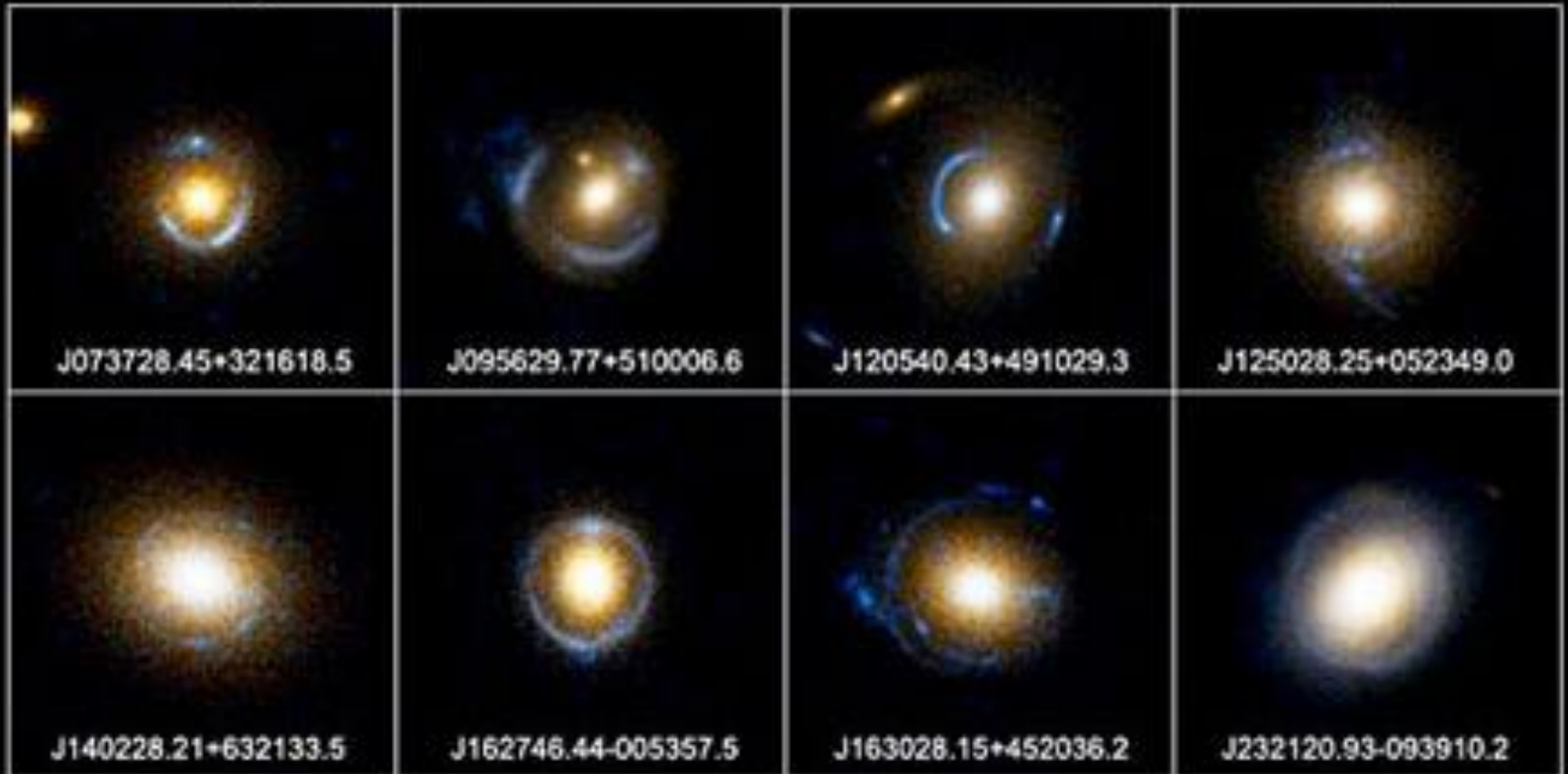
Gravitational lensing: Einstein rings



When the source and the lens are well aligned → strong arc or an Einstein ring

Einstein Ring Gravitational Lenses

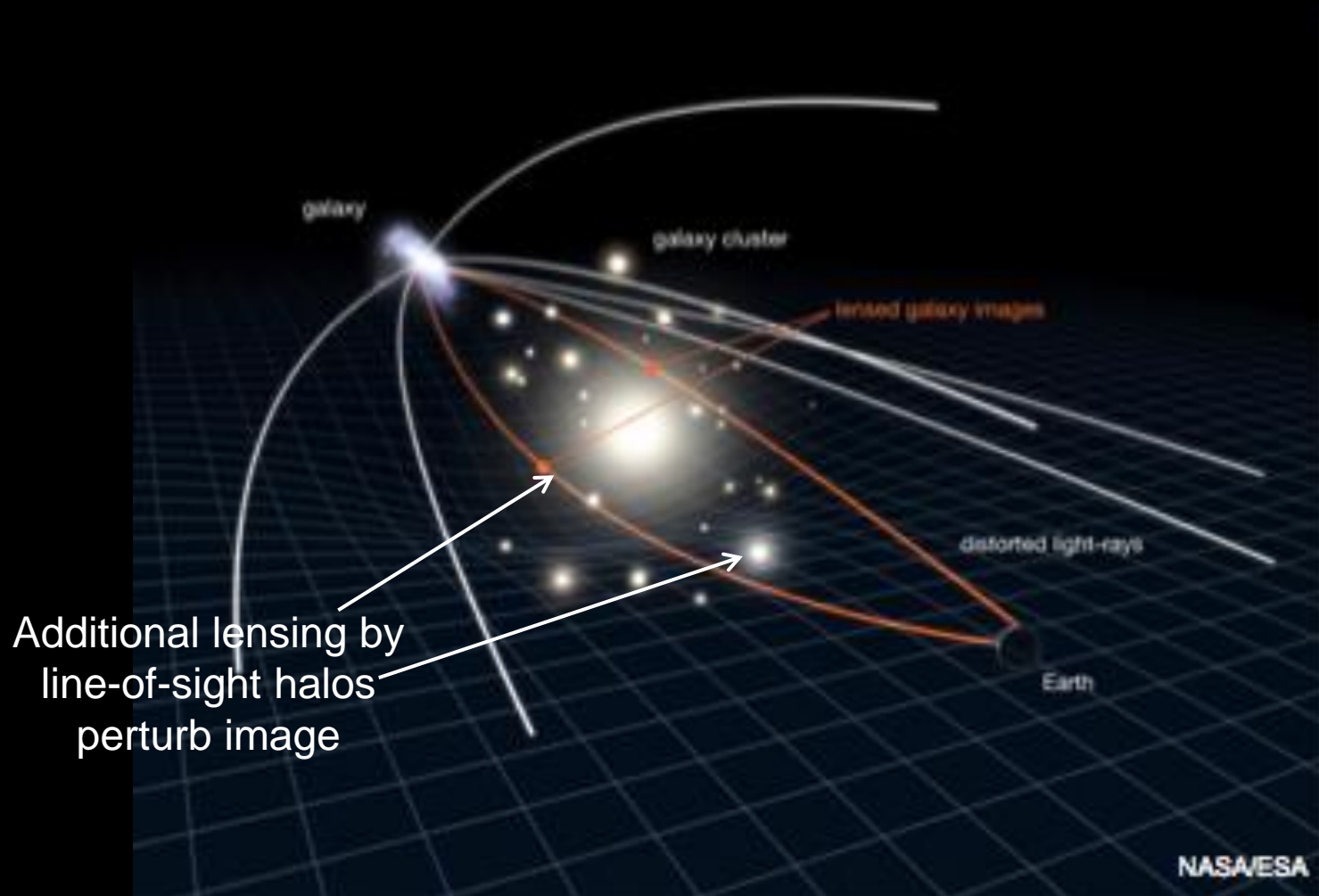
Hubble Space Telescope • ACS



NASA, ESA, A. Bolton (Harvard-Smithsonian CfA), and the SLACS Team

STScI-PRC05-32

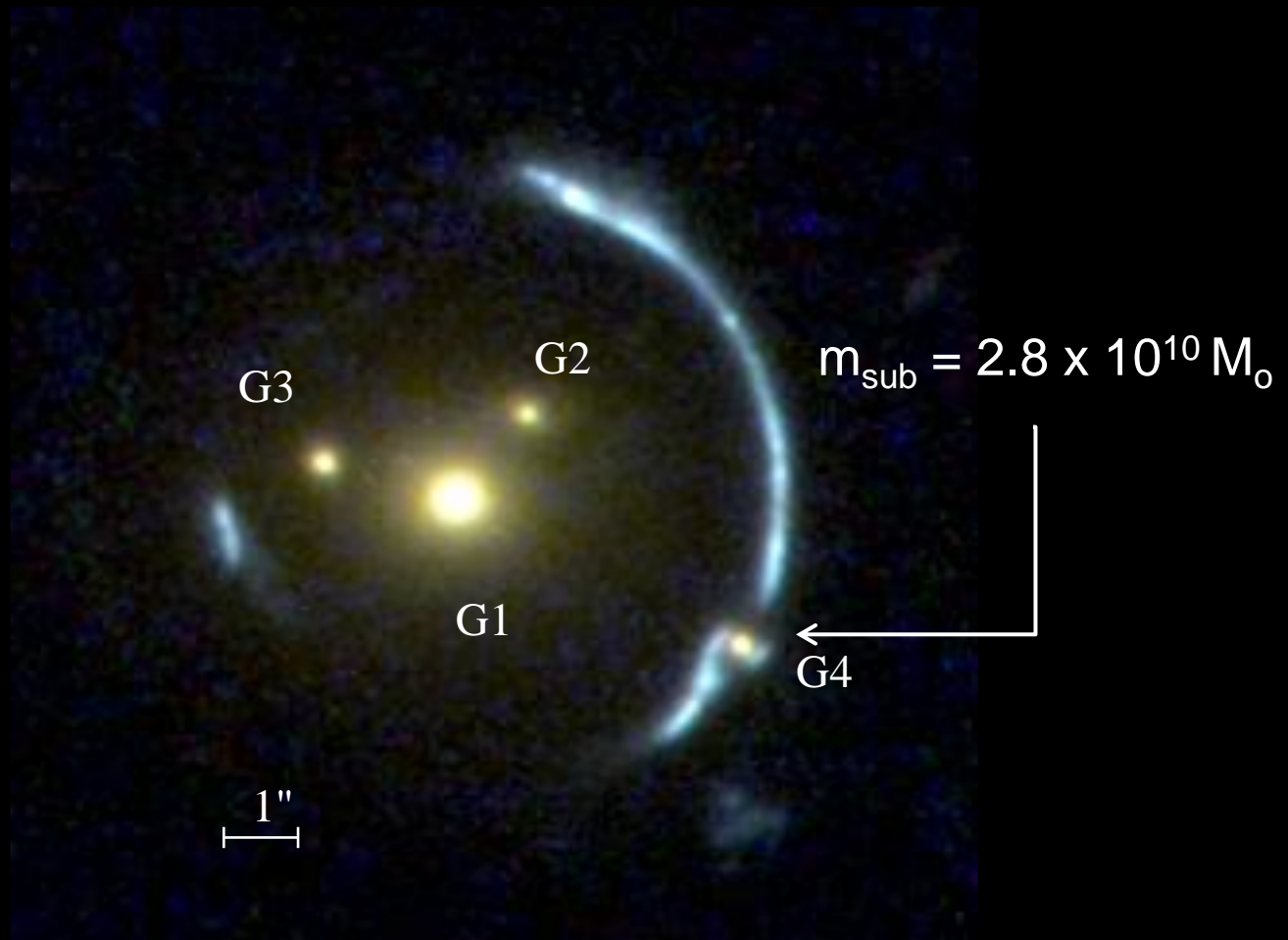
Gravitational lensing: Einstein rings



When the source and the lens are well aligned → strong arc or an Einstein ring

Gravitational lensing: Einstein rings

Halos projected onto an Einstein ring distort the image

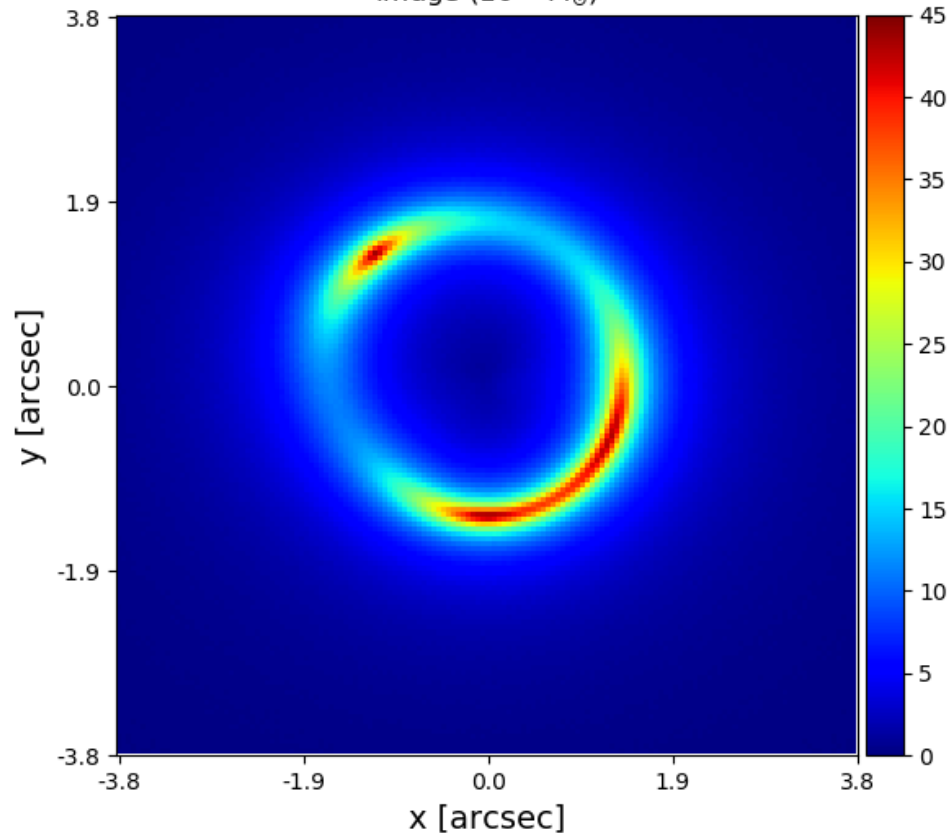


HST “data”: $z_{\text{source}}=1$; $z_{\text{lens}}=0.2$

$10^{10}M_{\odot}$ halo – **easy to spot**

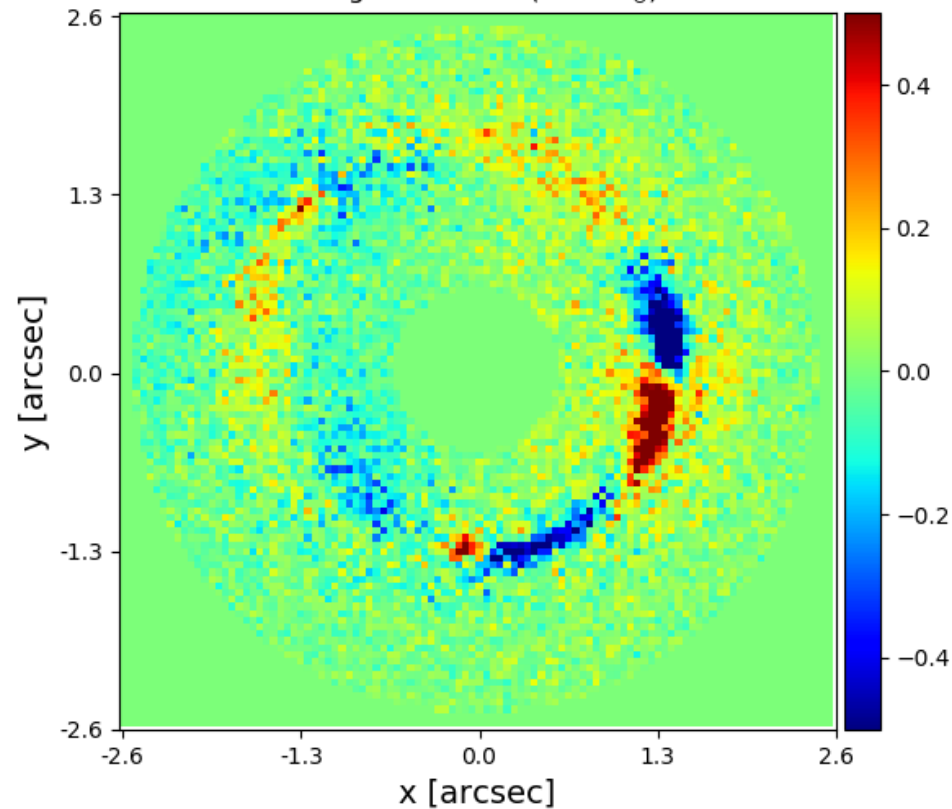
Image

Image ($10^{10} M_{\odot}$)



Residuals

Image Residuals ($10^{10} M_{\odot}$)



Searched for substructure in 55 lenses with good HST imaging

→ 2 detections: G2
G3

SLACS0946+1006 → $\text{Log } M_{\text{sub}} = 11.59^{+0.18 - 0.34}$

BELLS1226+5457 → $\text{Log } M_{\text{sub}} = 11.80^{+0.16 - 0.30}$

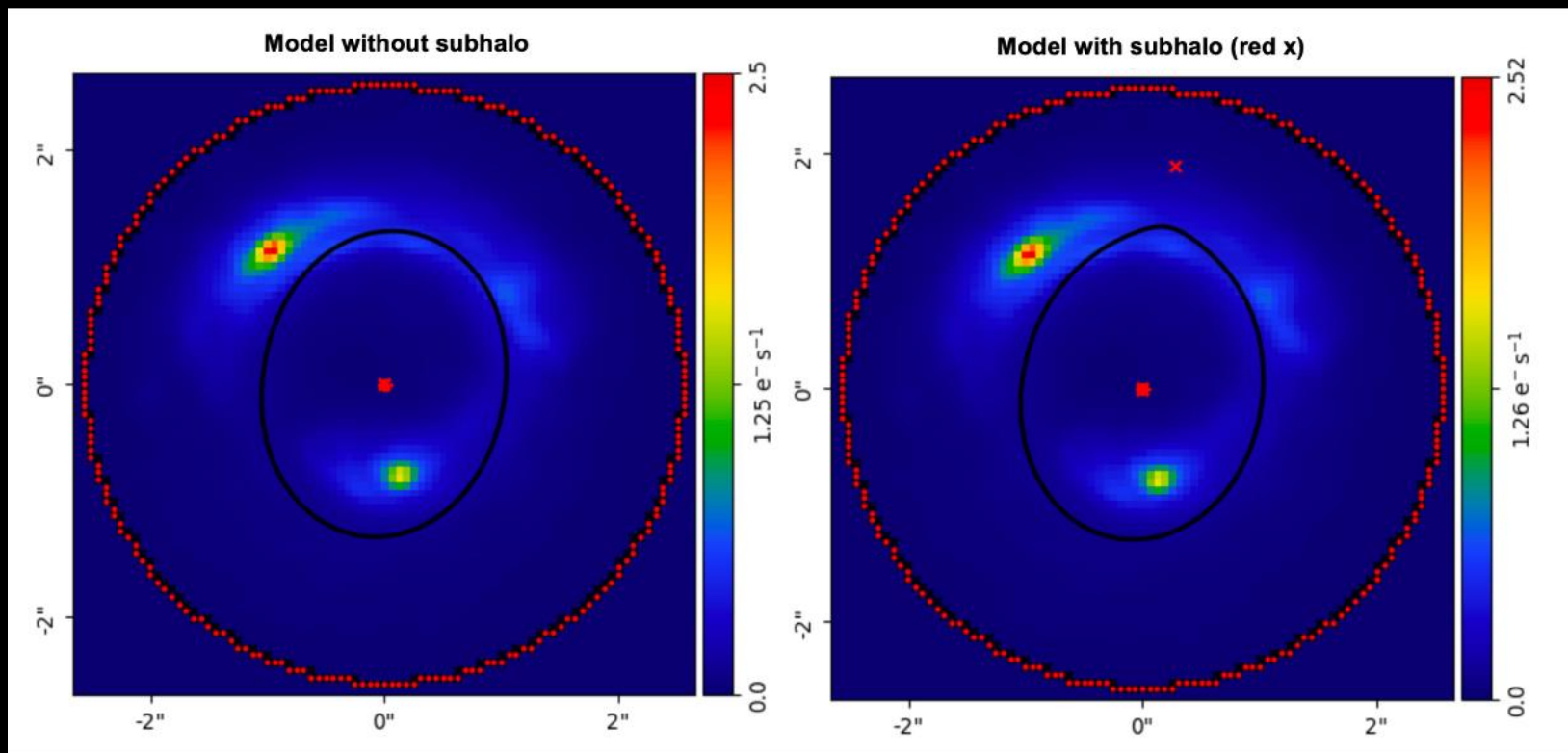
G1

Nightingale + '22

G4

1"
┌───┐

JWST



And another one in JWST data:

$$\rightarrow \text{Log } M_{\text{sub}} = 11.59^{+0.18}_{-0.34}$$

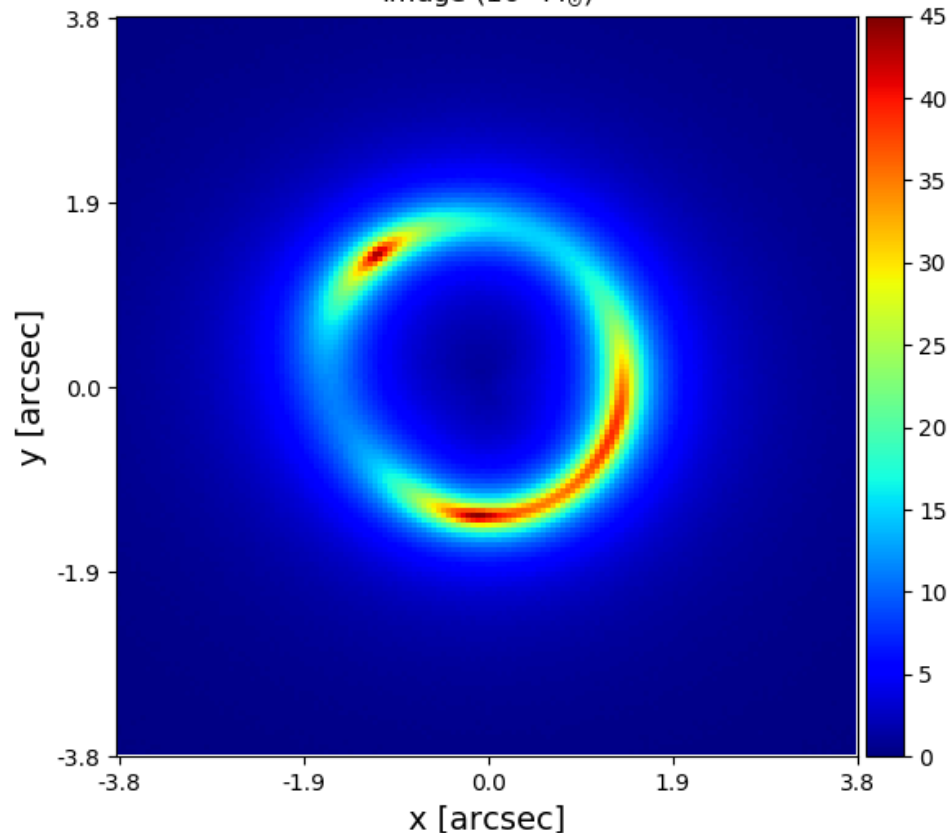
Lange, Nightingale, CSF+ '23

Strong lensing: detecting small halos

HST “data”: $z_{\text{source}}=1$; $z_{\text{lens}}=0.2$ $10^7 M_{\odot}$ halo – **NOT** so easy to spot

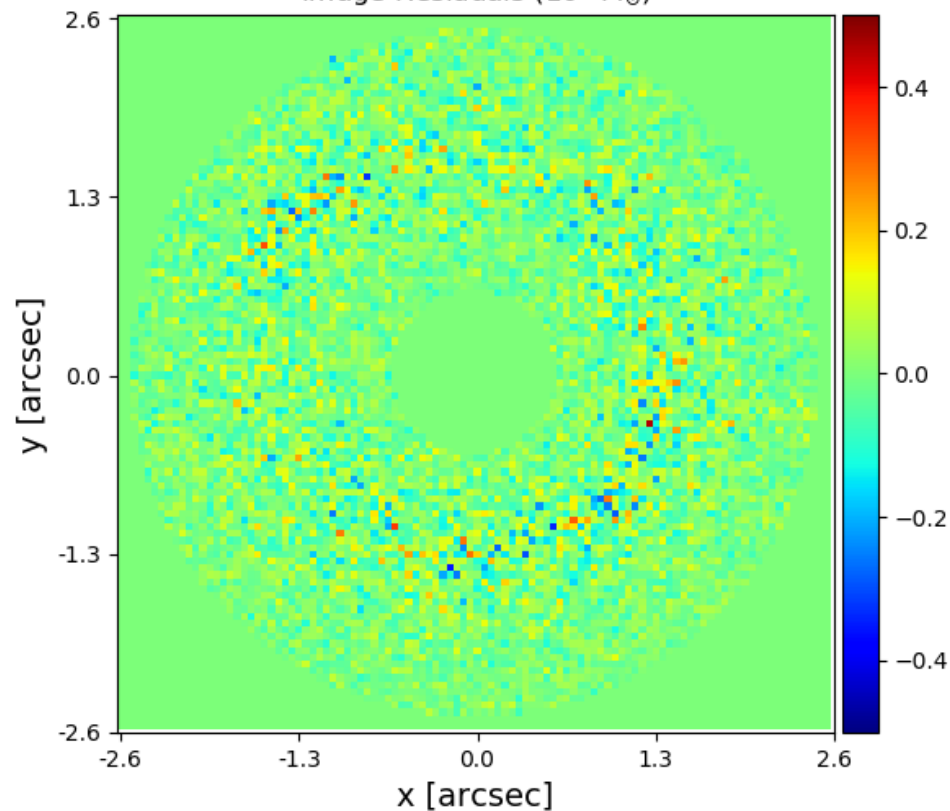
Image

Image ($10^7 M_{\odot}$)



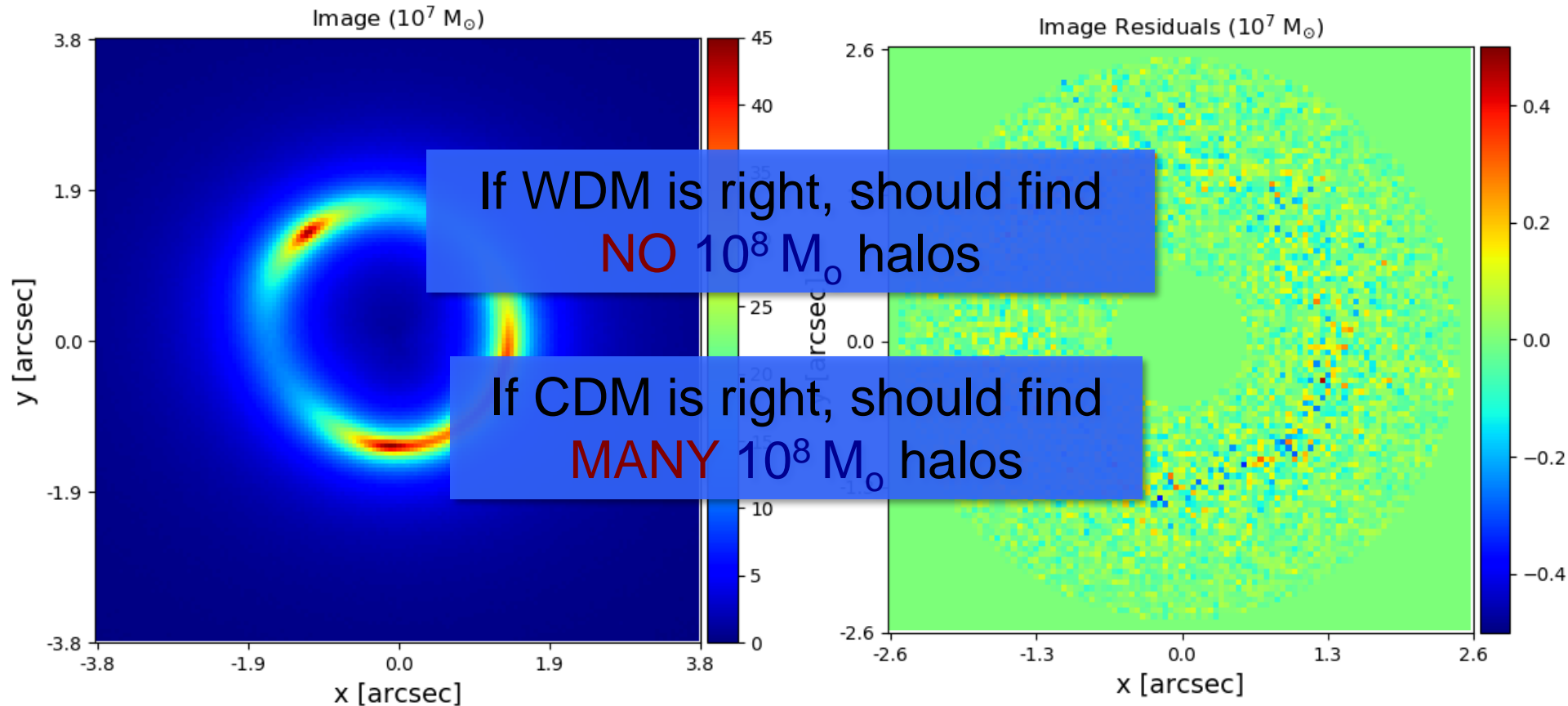
Residuals (image – smooth model)

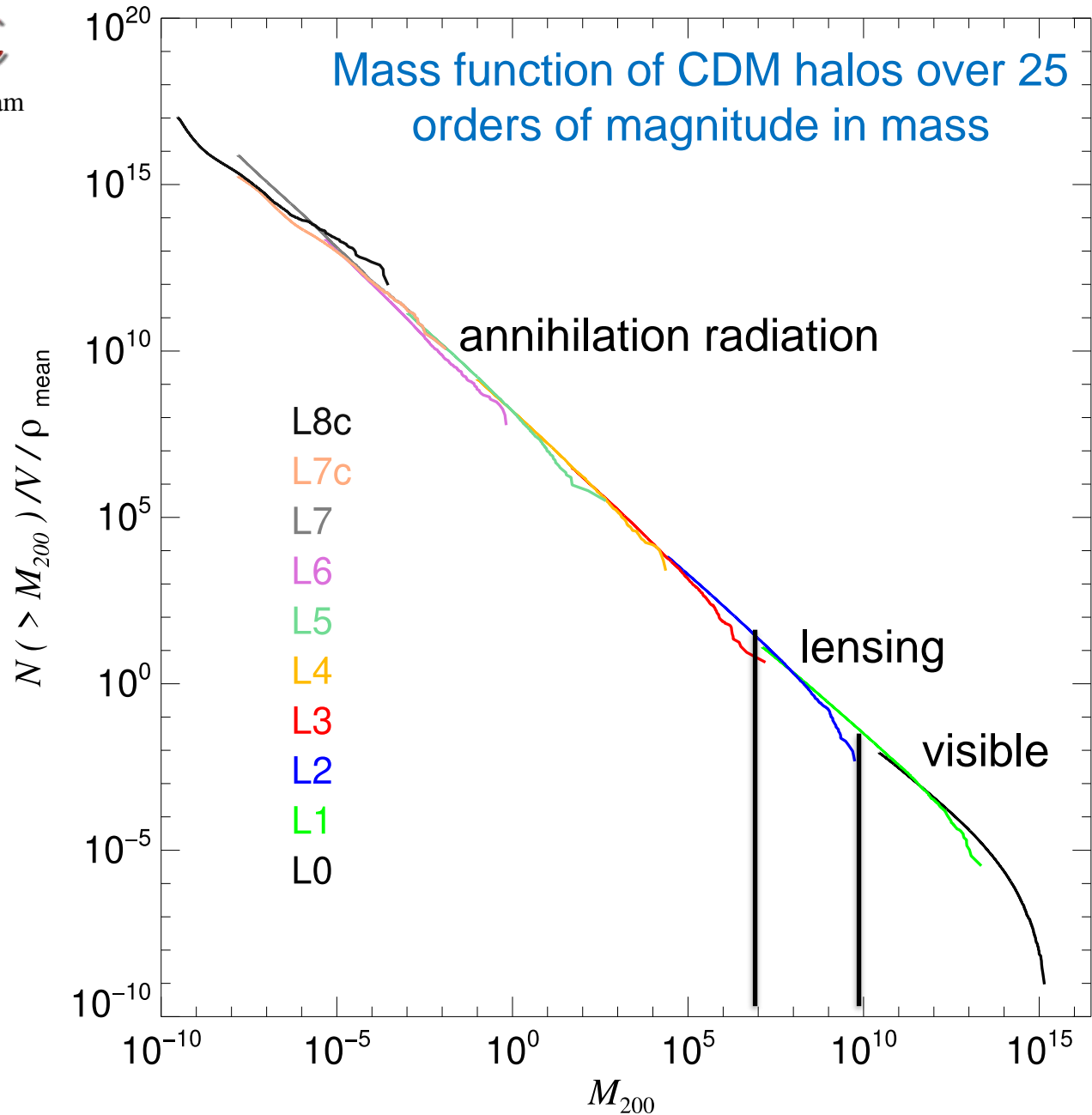
Image Residuals ($10^7 M_{\odot}$)



Detecting halos w. strong lensing

Can detect halos as small as $10^7 M_\odot$





Indirect CDM detection through annihilation radiation

Supersymmetric particles are Majorana particles → **annihilate** into Standard Model particles (including **γ -rays**)

Intensity of annihilation radiation at x is:

$$I(x) = \frac{1}{8\pi} \sum_f \frac{dN_f}{dE} \langle \sigma_f v \rangle \int_{los} \left(\frac{\rho_x}{M_x} \right)^2 dl$$

↑ cross-section (particle physics)
↓ halo density at x (astrophysics)

$\langle \sigma v \rangle = 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$ → relic abundance in simple SUSY models

⇒ Theoretical expectation requires knowing $\rho(x)$

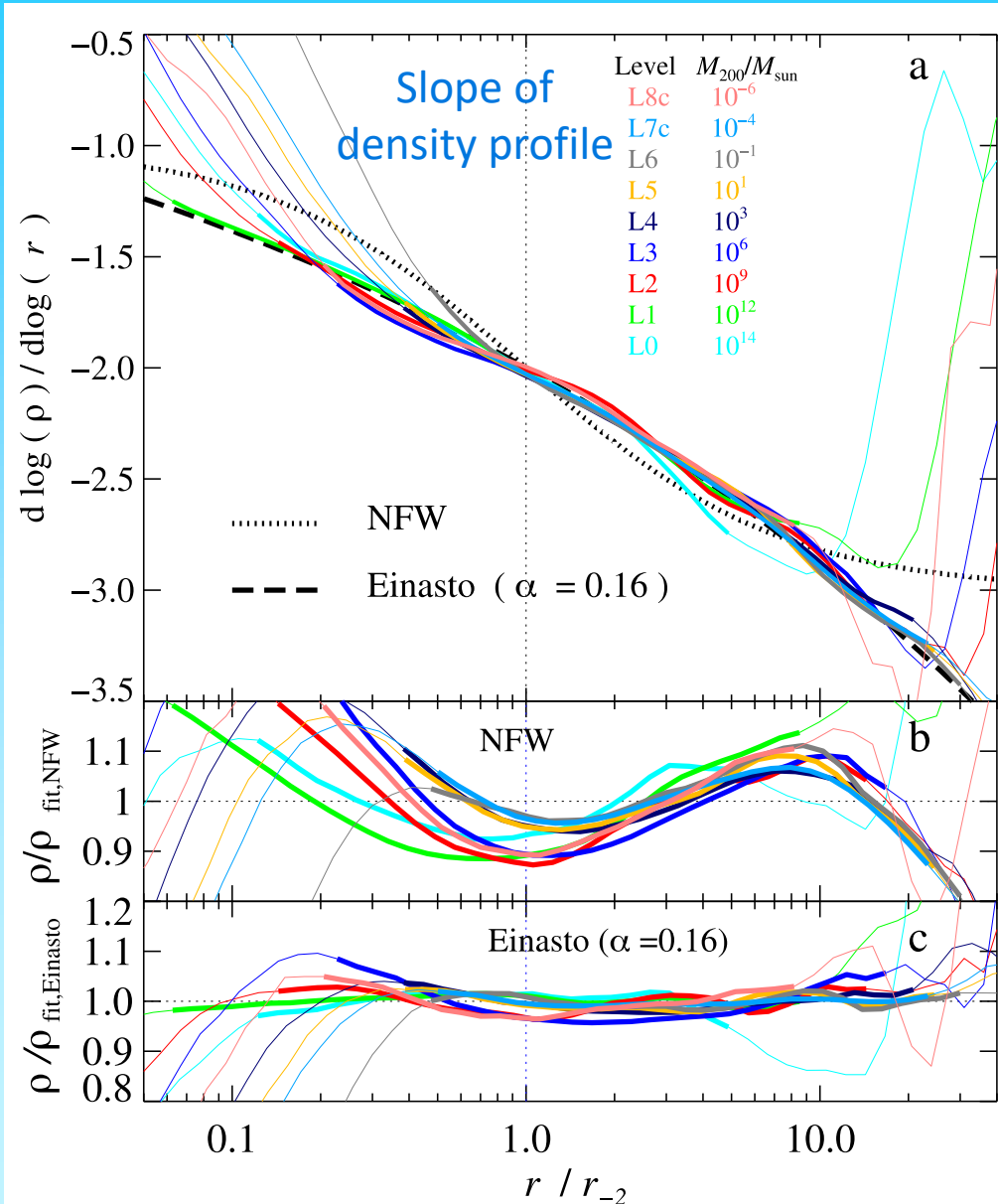
⇒ Accurate high resolution **N-body** simulations of **halo** formation from **CDM initial conditions**

Density profile shapes

Over **20 orders** of magnitude in halo **mass** and 4 orders of magnitude in density, the mean density **profiles** of halos are **fit** by **NFW** to within **20%** and by **Einasto** ($\alpha = 0.16$) to within **7%**

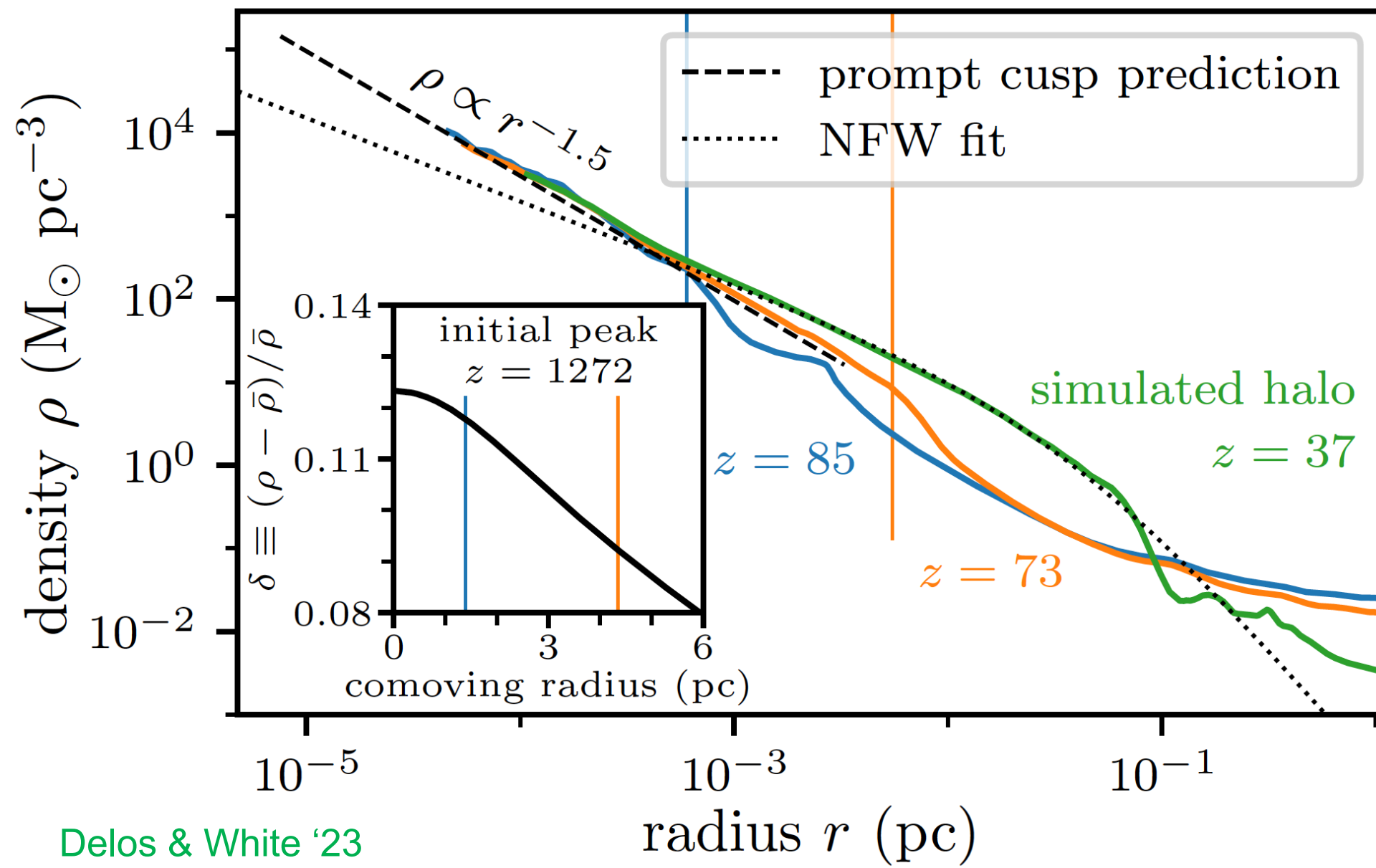
Earth-mass halo

Wang, Bose, CSF + '20



Earth-mass halo

Prompt cusp and subsequent halo growth for a peak with $z_{\text{coll}} = 87$



Prompt cusps

Delos & White '22

- For standard WIMPs, they make up $\sim 1\%$ of all the dark matter
- In MW, disrupted by tides & stellar encounters within ~ 20 Kpc
- Dominate DM annihilation signal from outer halo of MW and all extragalactic objects, leading to $L \propto \rho_{\text{DM}}$, not ρ_{DM}^2

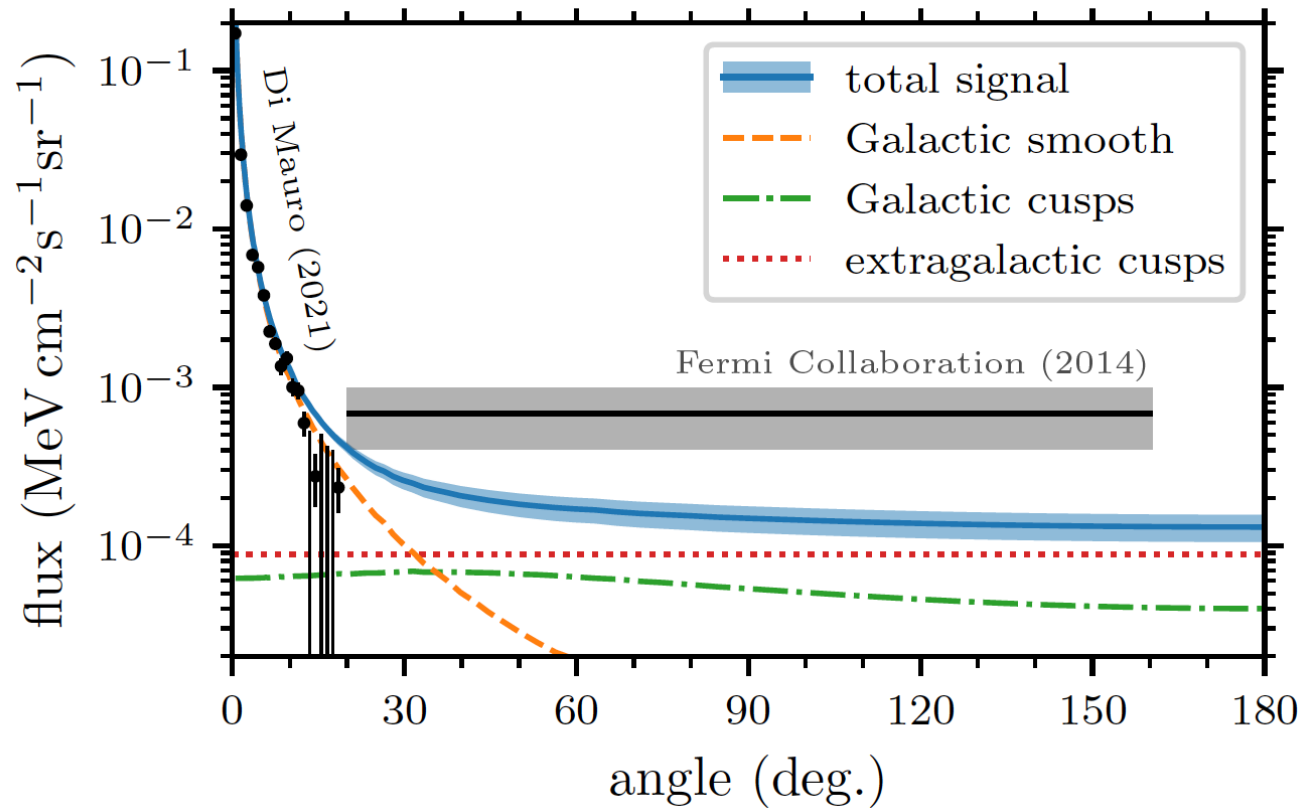
The Galactic Centre excess

Cusp emission dominates at $>20^\circ$ from centre

Surface brightness from MW halo comparable to that from halos of external galaxies

Prompt cusps dominate the 1 – 10 GeV background

Delos & White '23



Only just consistent with γ -ray background

Conclusions

- The properties CDM (no baryons) on all scales - solved problem
- The **abundance & structure** of CDM halos (down **to Earth's mass**) and of WDM (no halos of mass $<10^8 M_{\odot}$) is now known
- CDM (and WDM) **halos** of all masses have **NFW** density profiles (except in **very inner** parts for halos near the **cutoff**)
- There is NO “small-scale” **crisis** in CDM
- Local large-scale structure is reflected in the **cosmic neutrino background**. MW dipole & angular PS depend in the ν mass
- Distortions of strong **gravitational lenses** offer a **clean test** of CDM vs WDM \rightarrow and can potentially rule out CDM!
- Prompt **cusps** of Earth mass **dominate** the annihilation radiation from **outer halo** of MW and **extragalactic objects**