

Ultra long turbulent eddies, magnetic topology, and the triggering of internal transport barriers in tokamaks

EFTC 2023

Arnas Volčokas, Justin Ball, Stephan Brunner

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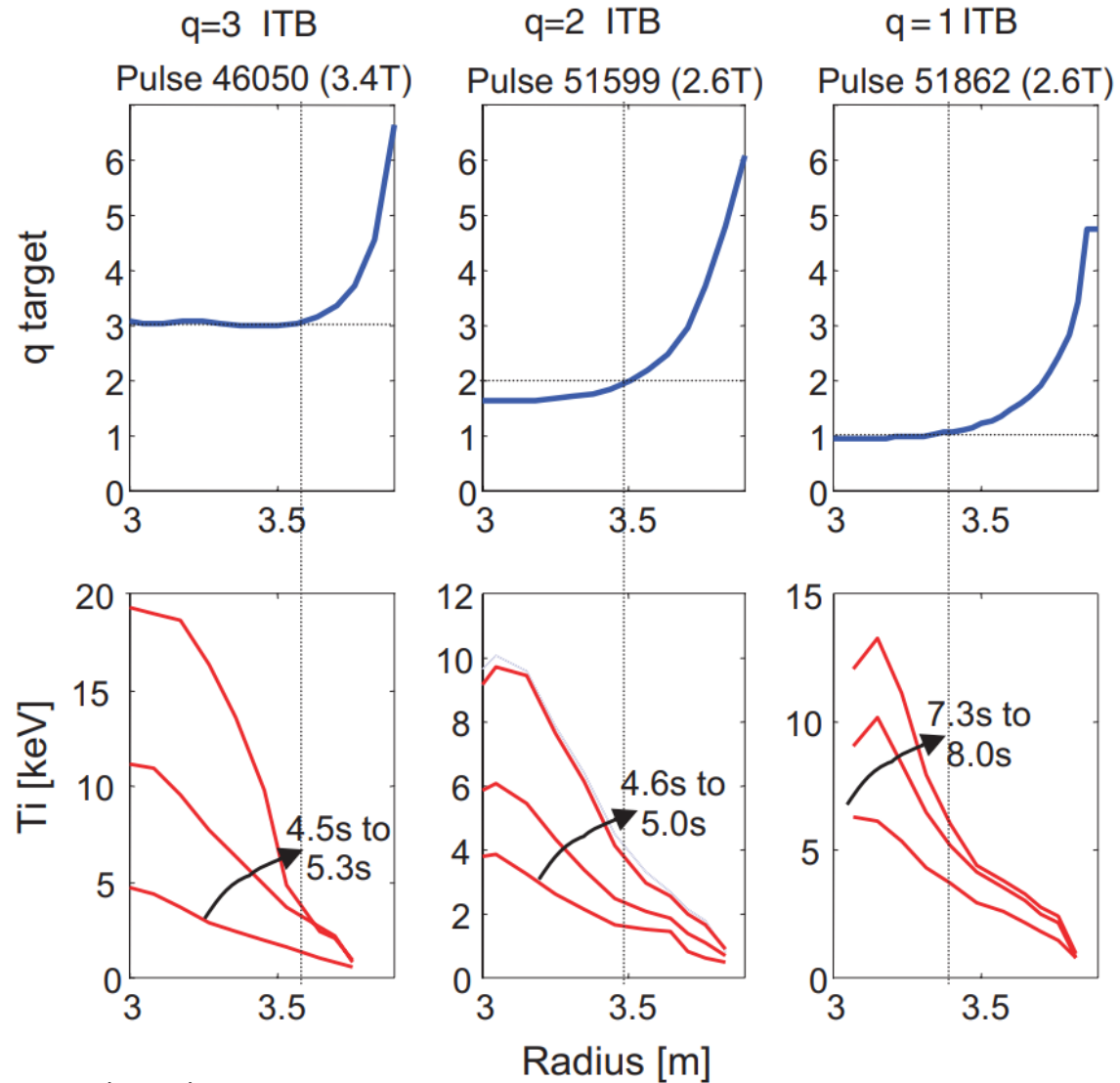
EPFL Outline

- Motivation and background
- Methods
- Numerical results:
 - Low magnetic shear simulations
 - Ultra-long turbulent eddies $\hat{s} = 0$
 - Non-uniform magnetic shear simulations
- Conclusions

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Internal Transport Barriers (ITBs)

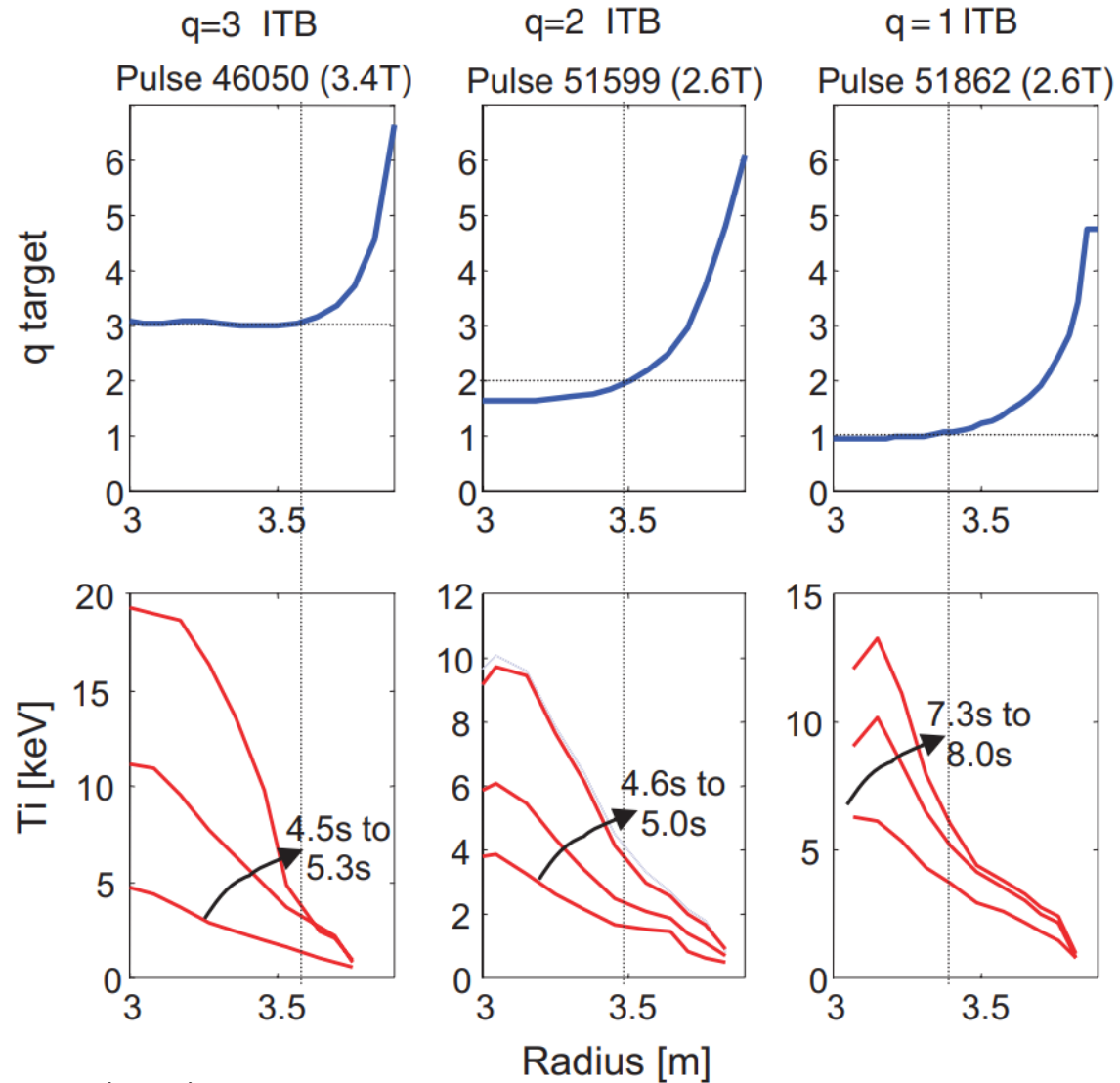


Examples
from JET

EPFL Internal Transport Barriers (ITBs)

ITBs are formed when:

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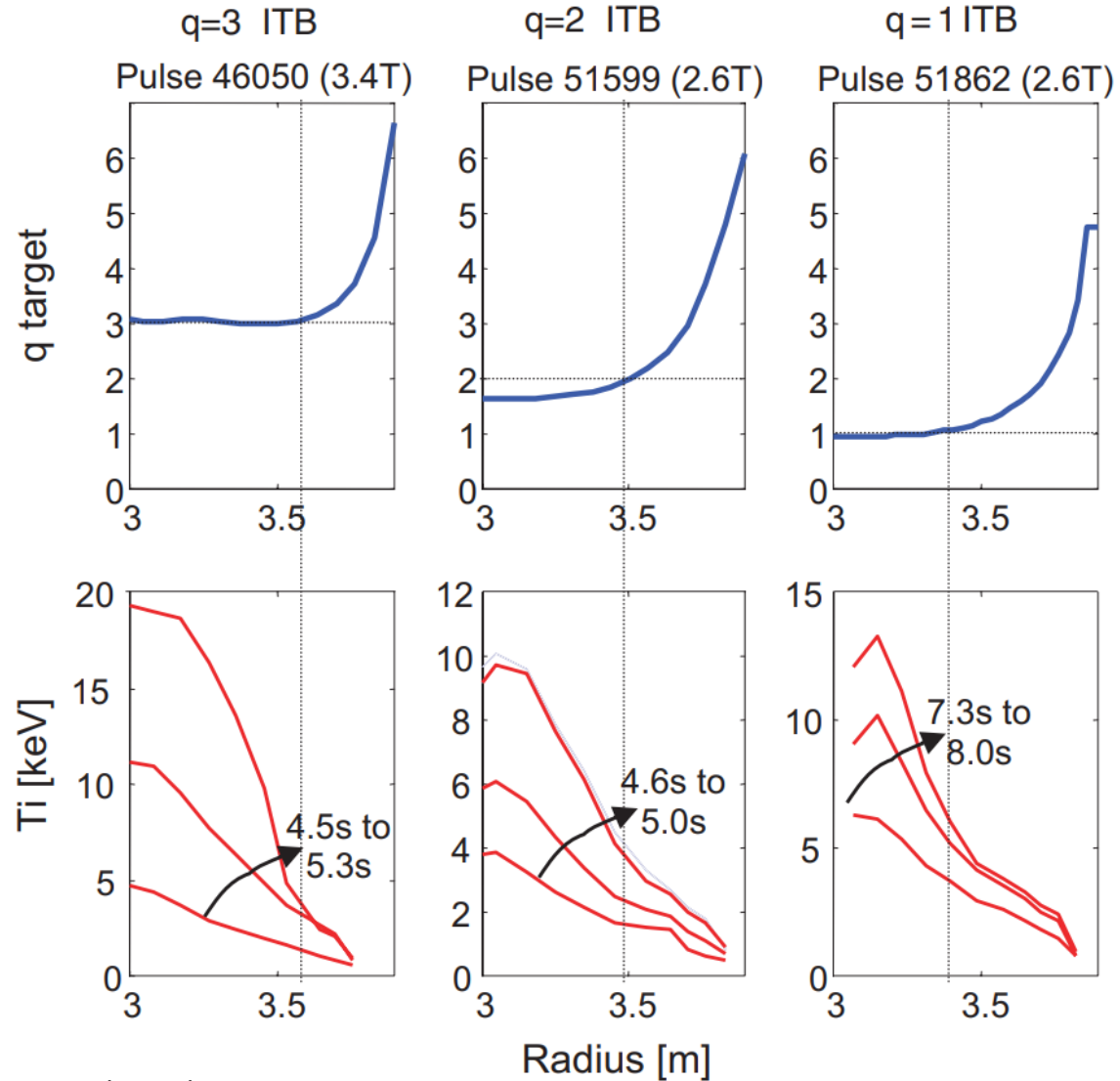


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EPFL Internal Transport Barriers (ITBs)

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- Facilitated by integer or low order rational $q = M/N$ with $\hat{s} \approx 0$

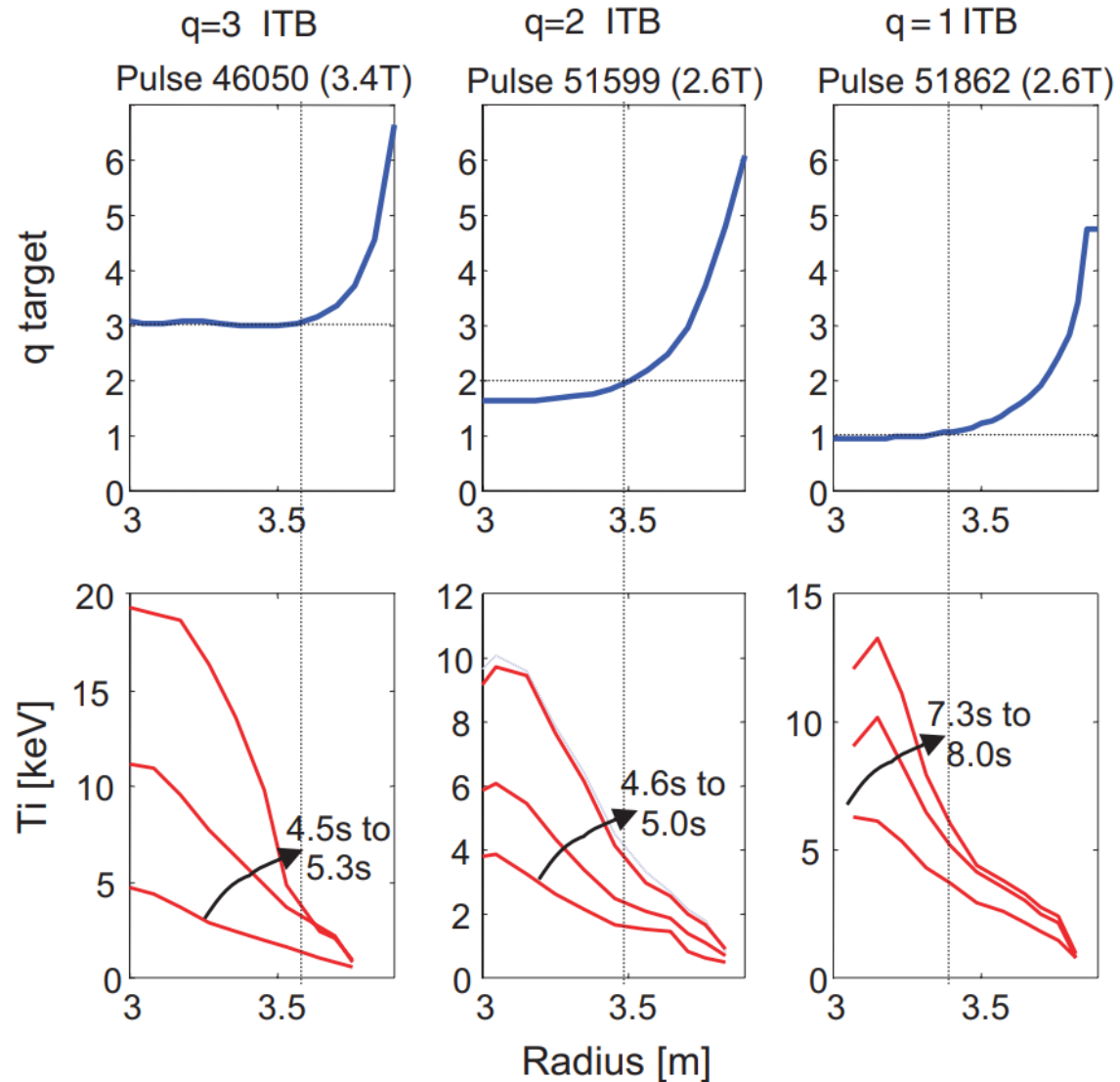


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- Other conditions...

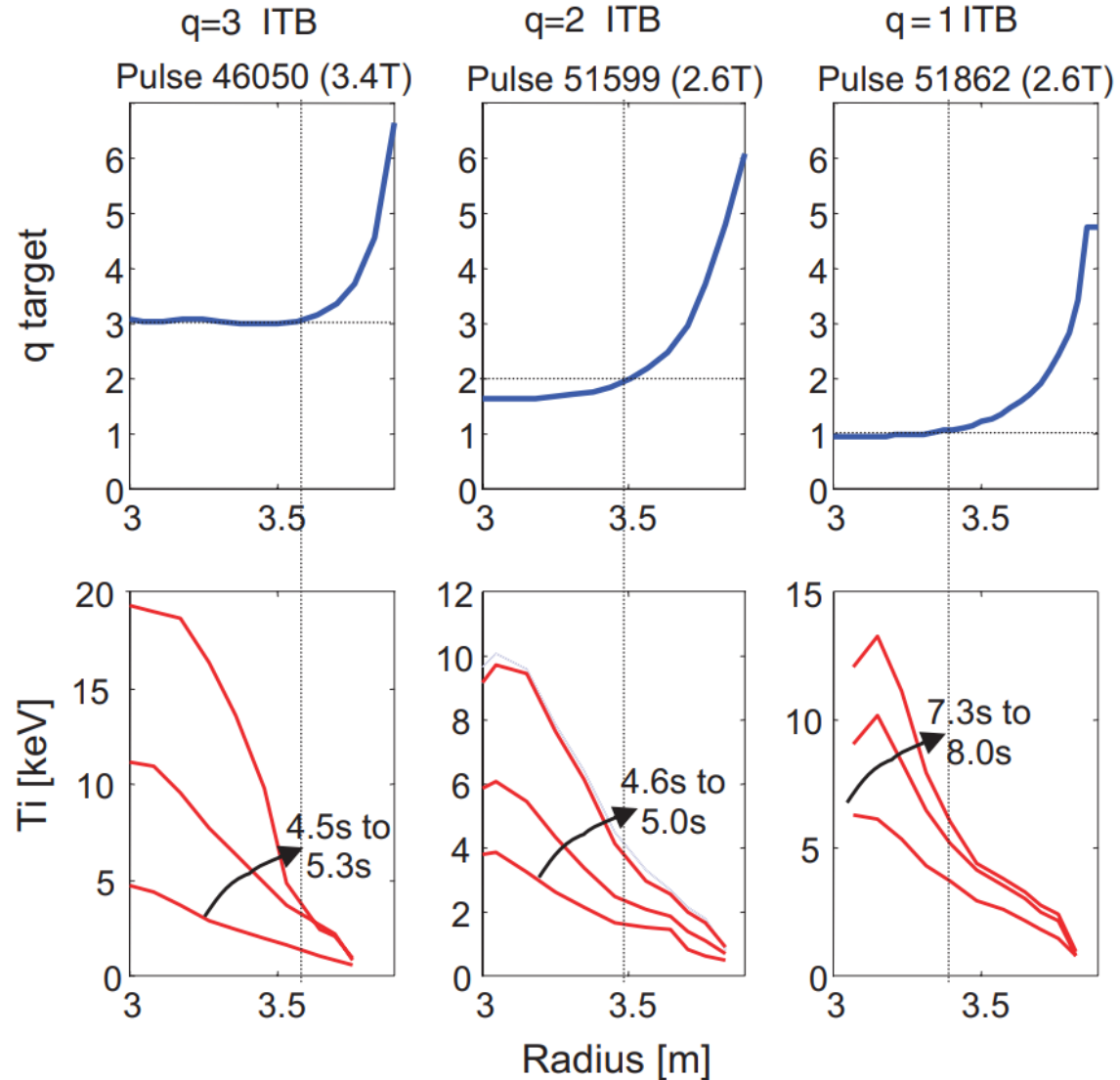


Examples
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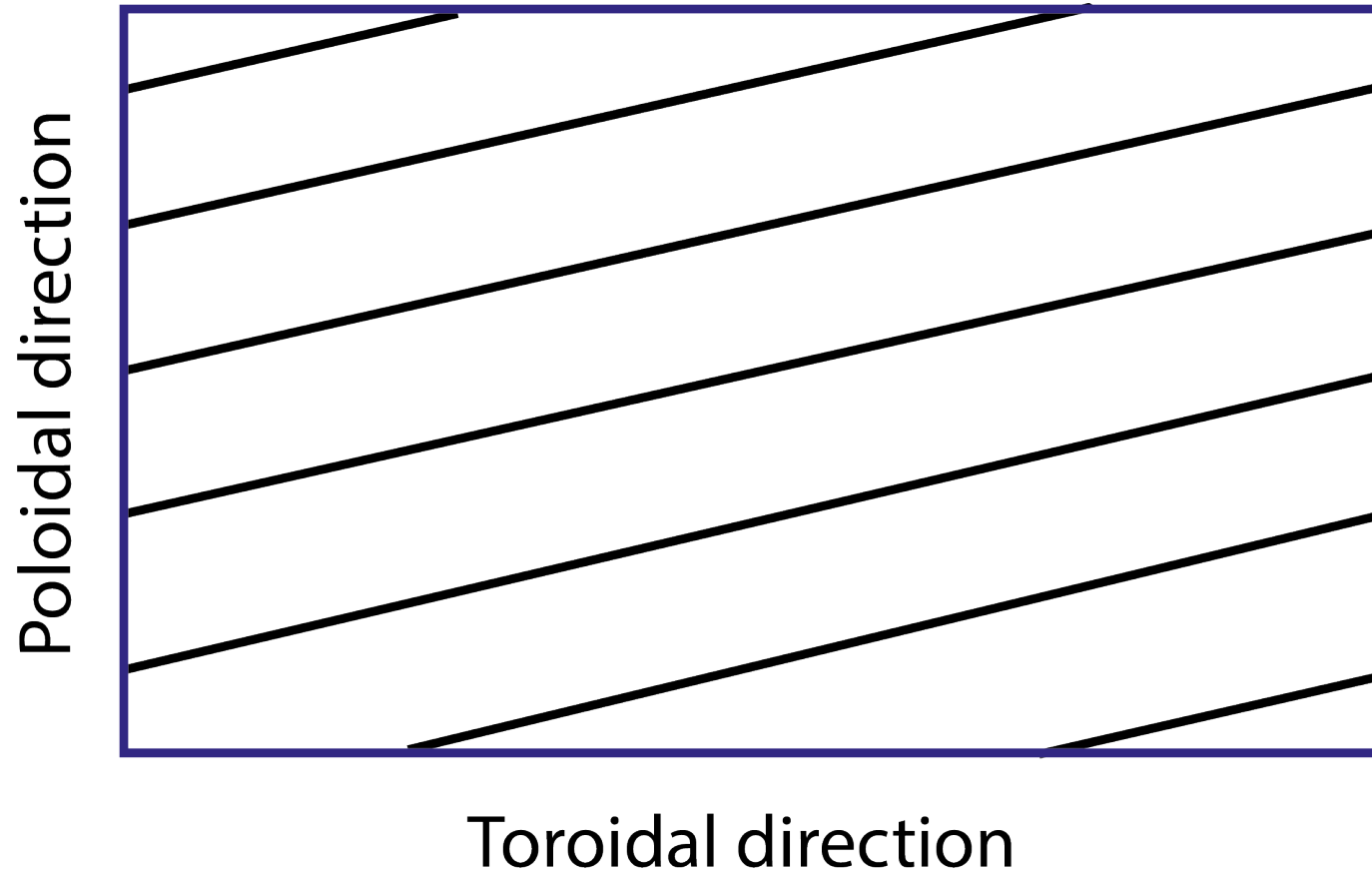
Examples from JET

What is the role of turbulent self-interaction in ITB formation?

Self-interaction – the nonlinear interaction of turbulent eddy with itself

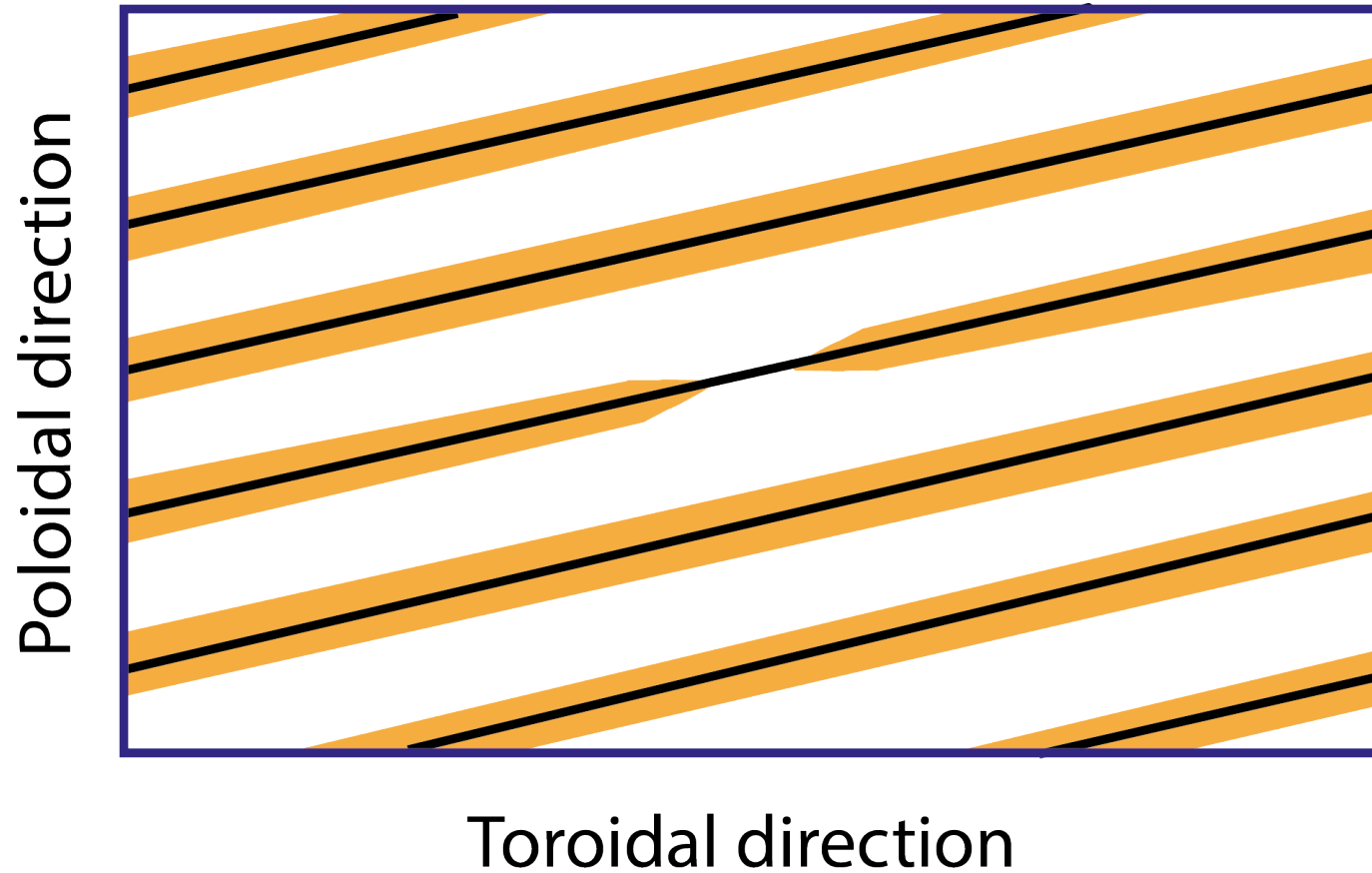
- J. Ball *et al.* 2020 *Journal of Plasma Physics* **86(2)**, 905860207
- Ajay CJ, Studying the effect of non-adiabatic passing electron dynamics on microturbulence self-interaction in fusion plasmas using gyrokinetic simulations, Thesis EPFL Lausanne, 2020
- J. Dominski *et al.* 2015 *Physics of Plasmas* **22**, 062303

EPFL No self-interaction



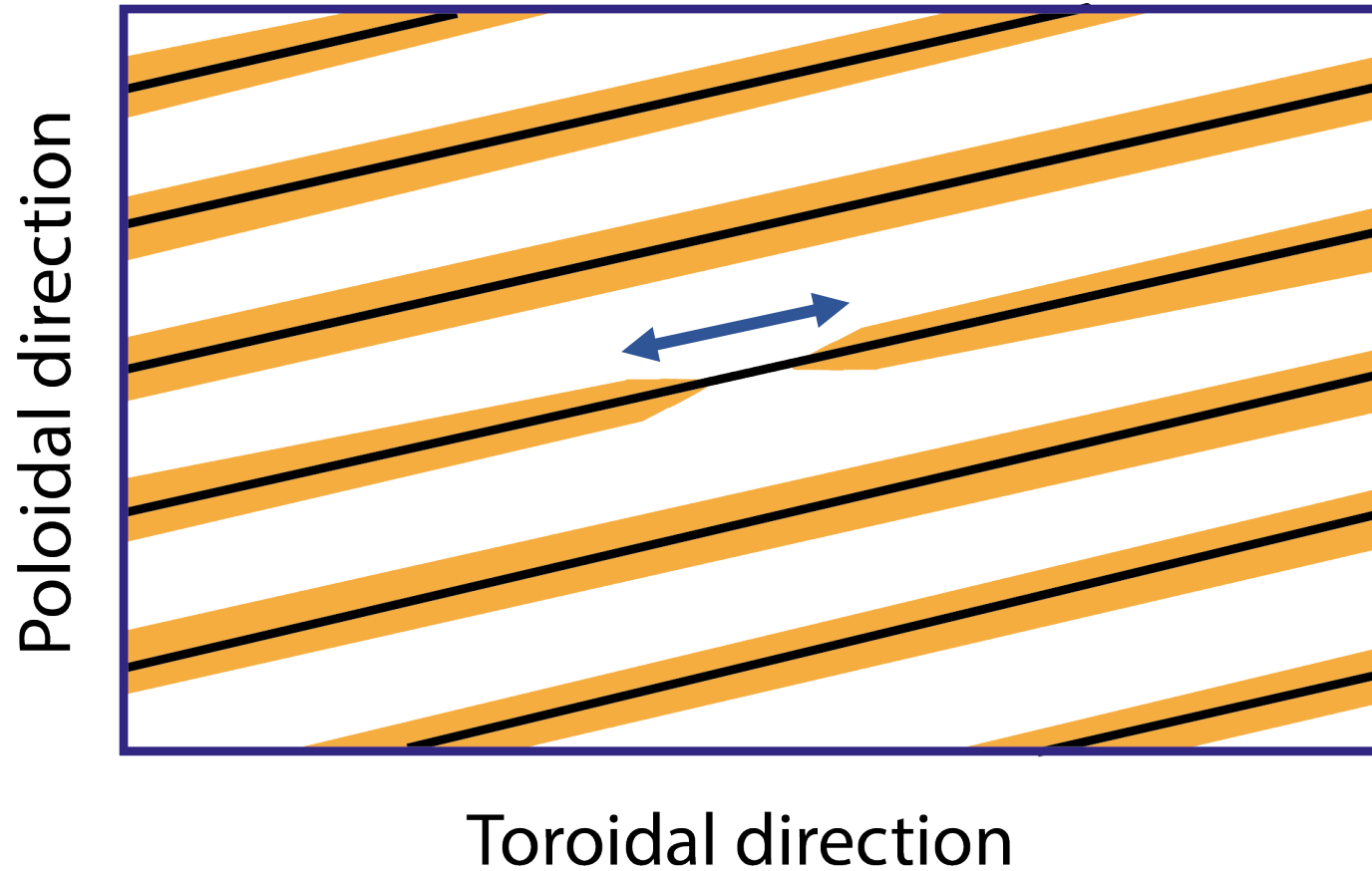
$$q = 2.5$$

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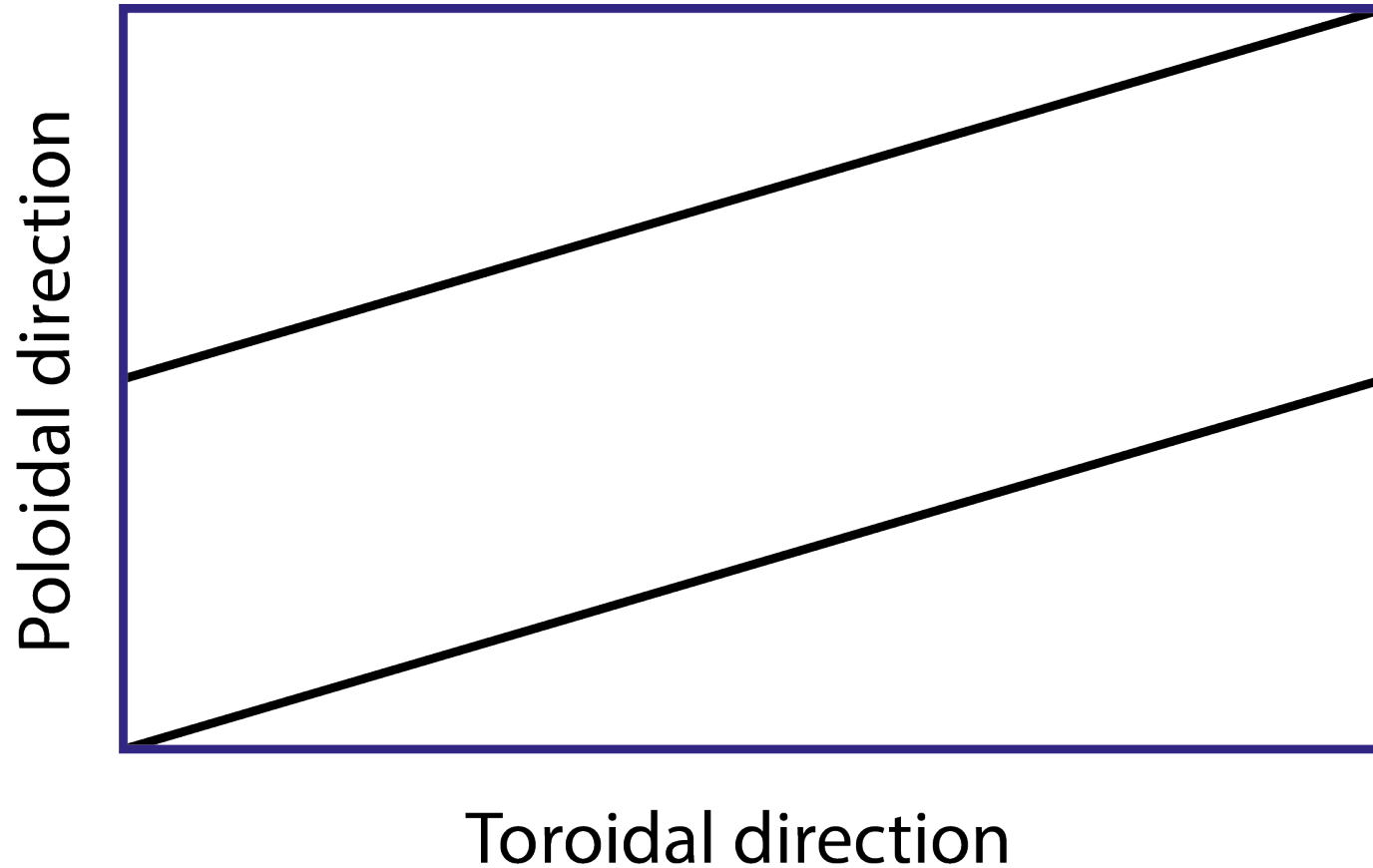
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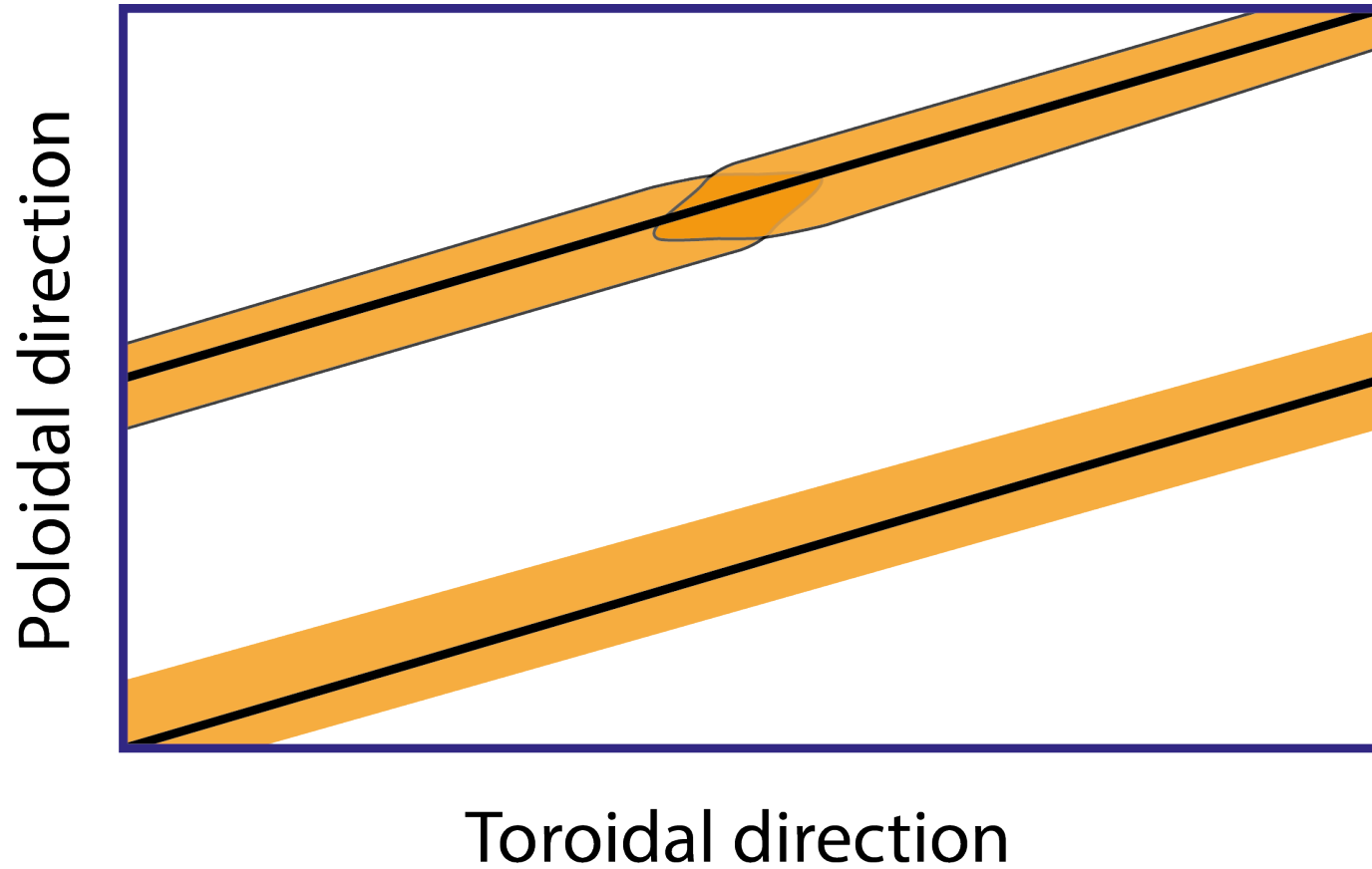
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EPFL Self-interaction



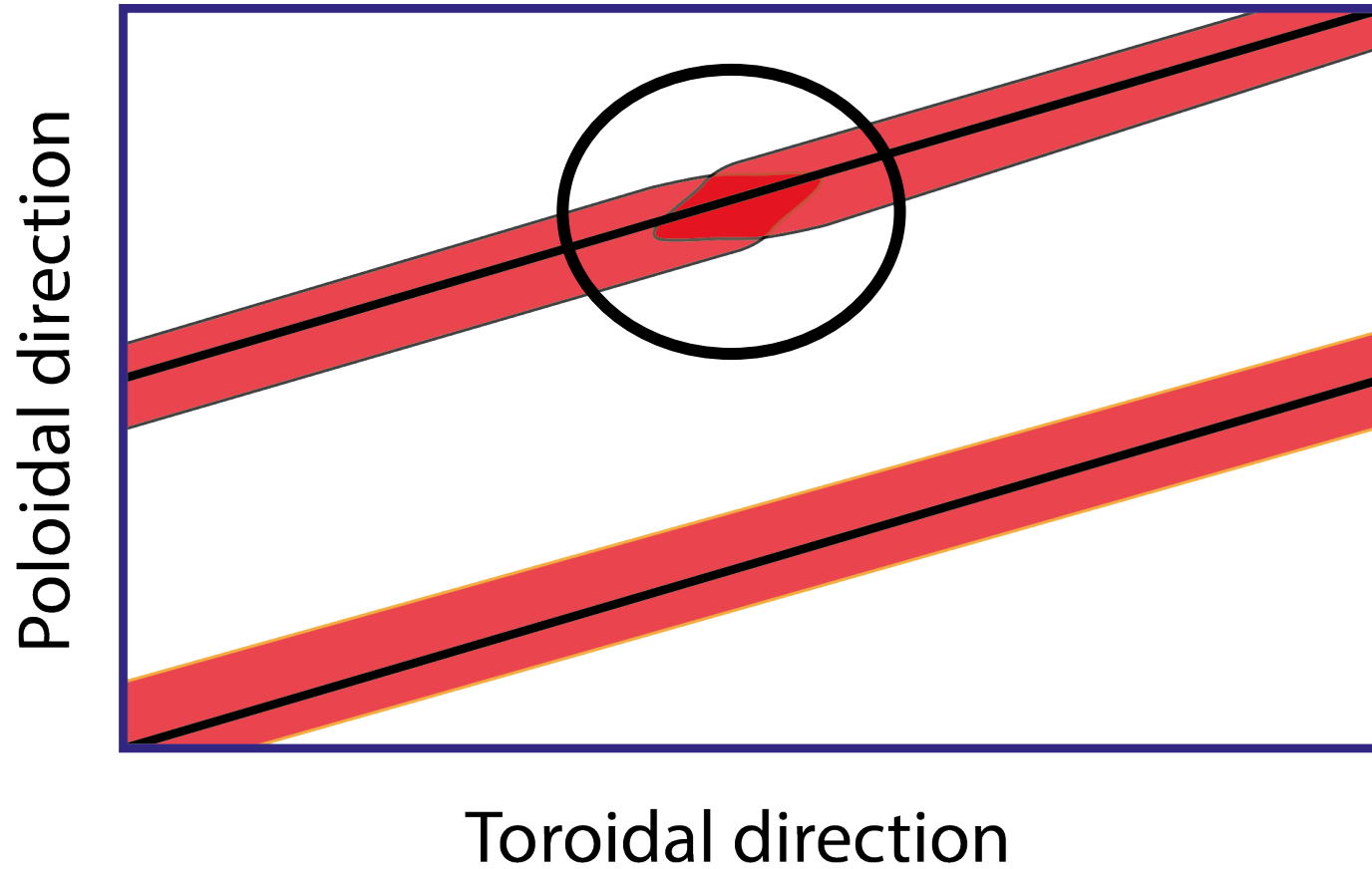
$$q = 2$$

EPFL Self-interaction



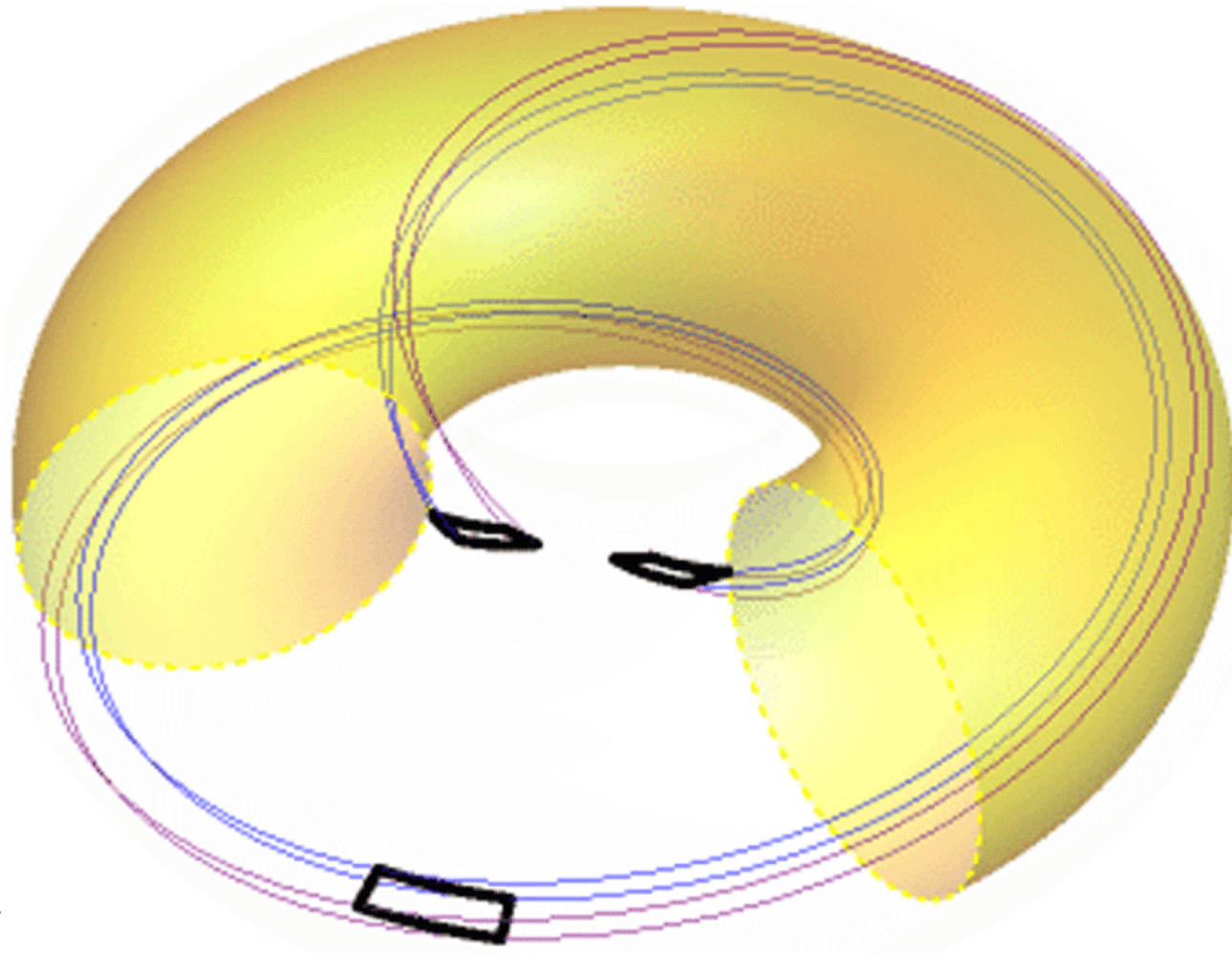
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EPFL Self-interaction



EPFL Outline

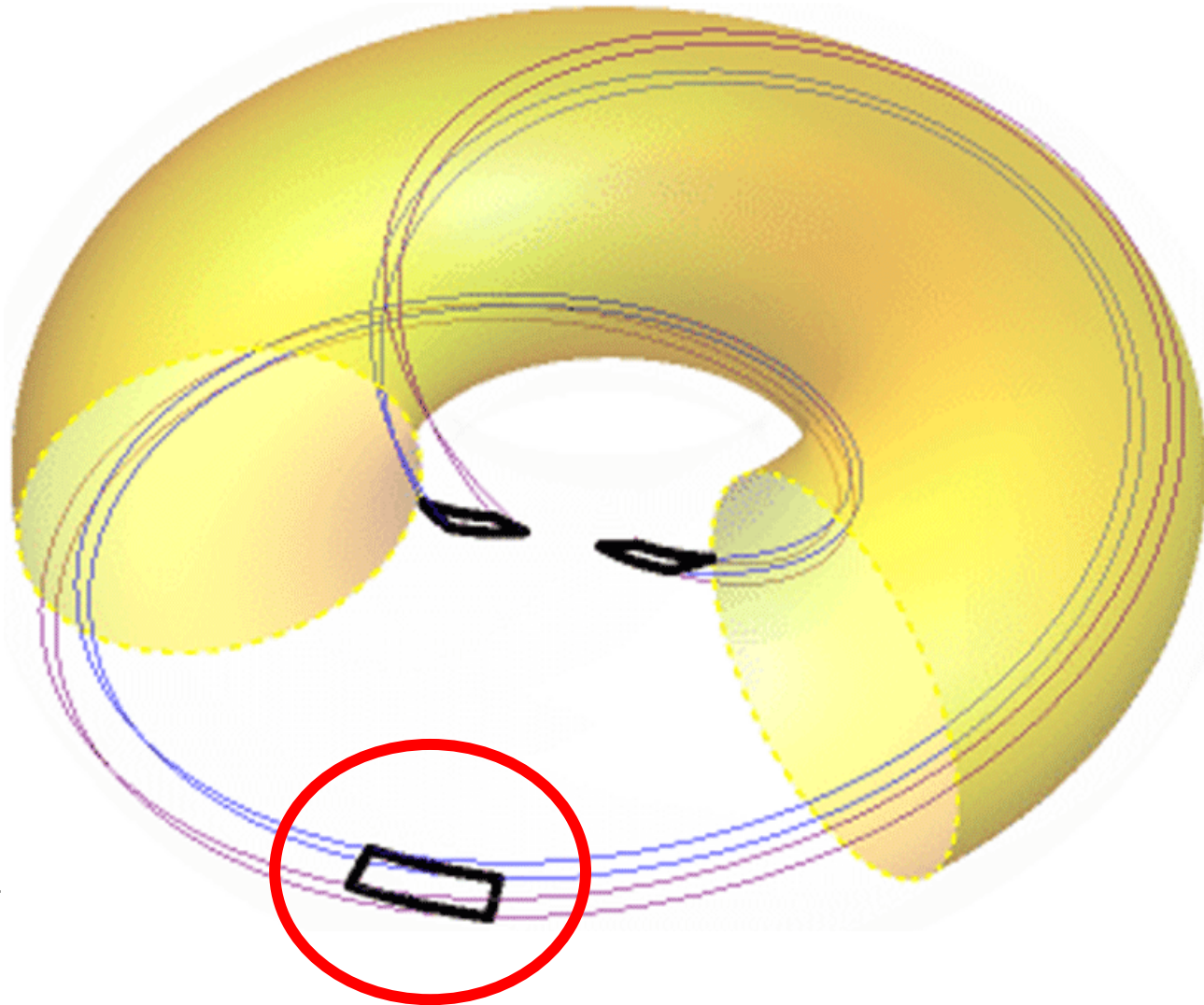
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$$N_{pol} = 1$$

Gyrokinetic
Electromagnetic
Numerical
Experiment-
Eulerian δf code

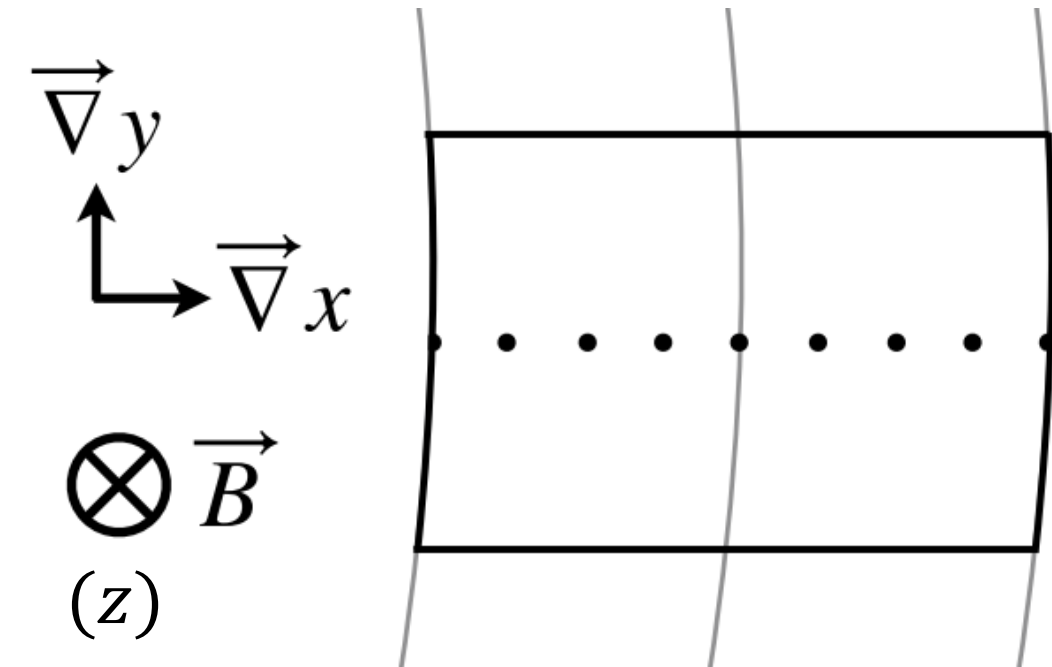
N_{pol} - number of
poloidal turns



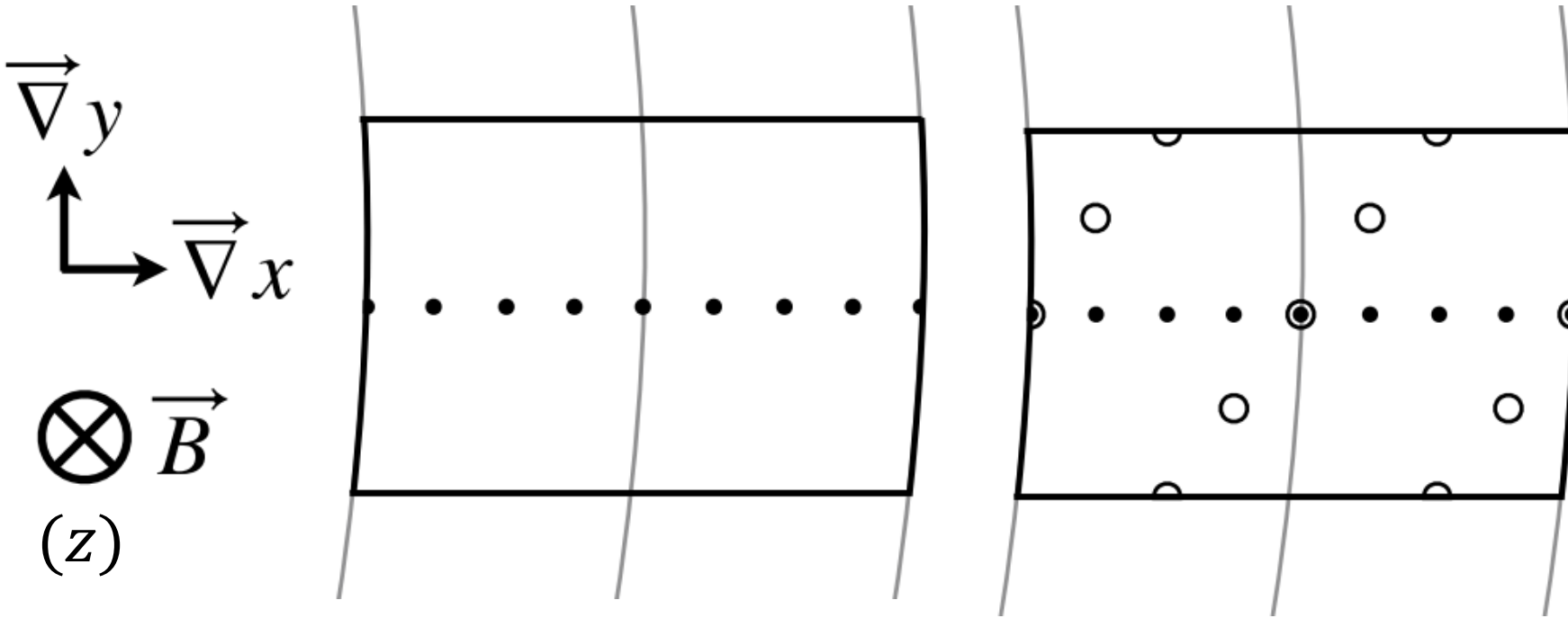
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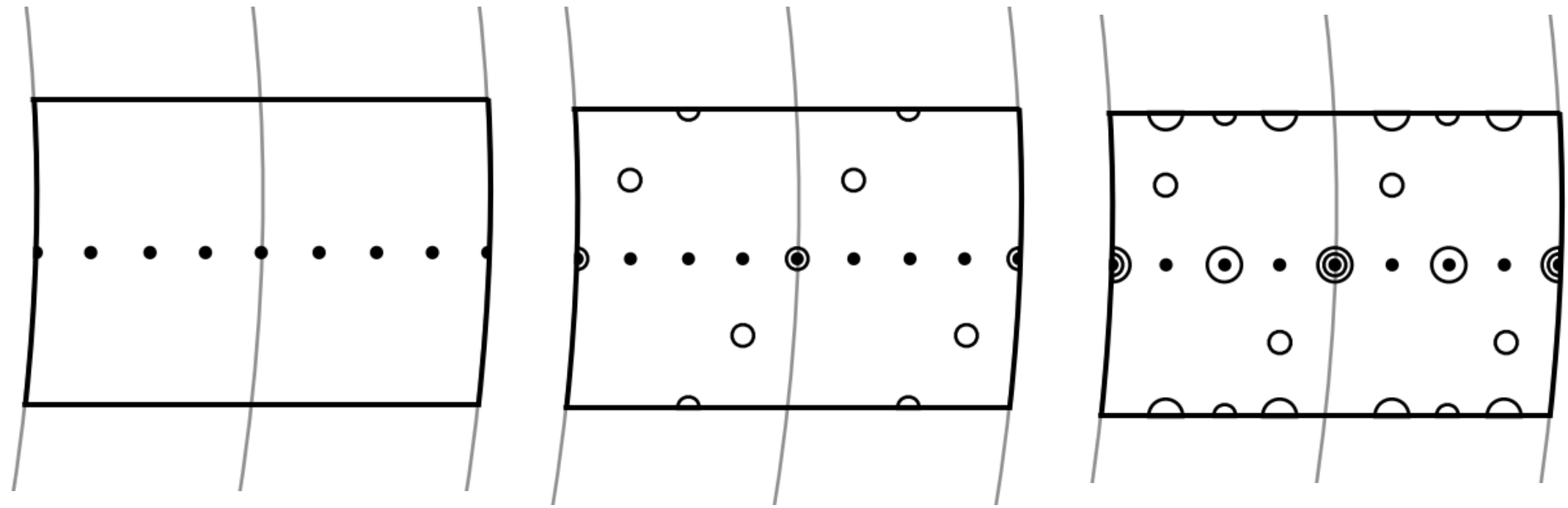
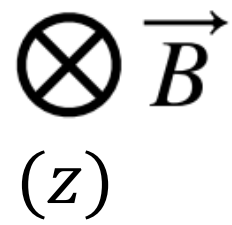
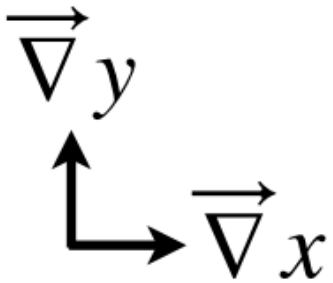
Twist-and-shift $\hat{s} \neq 0$ 

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1st pass through
parallel boundary

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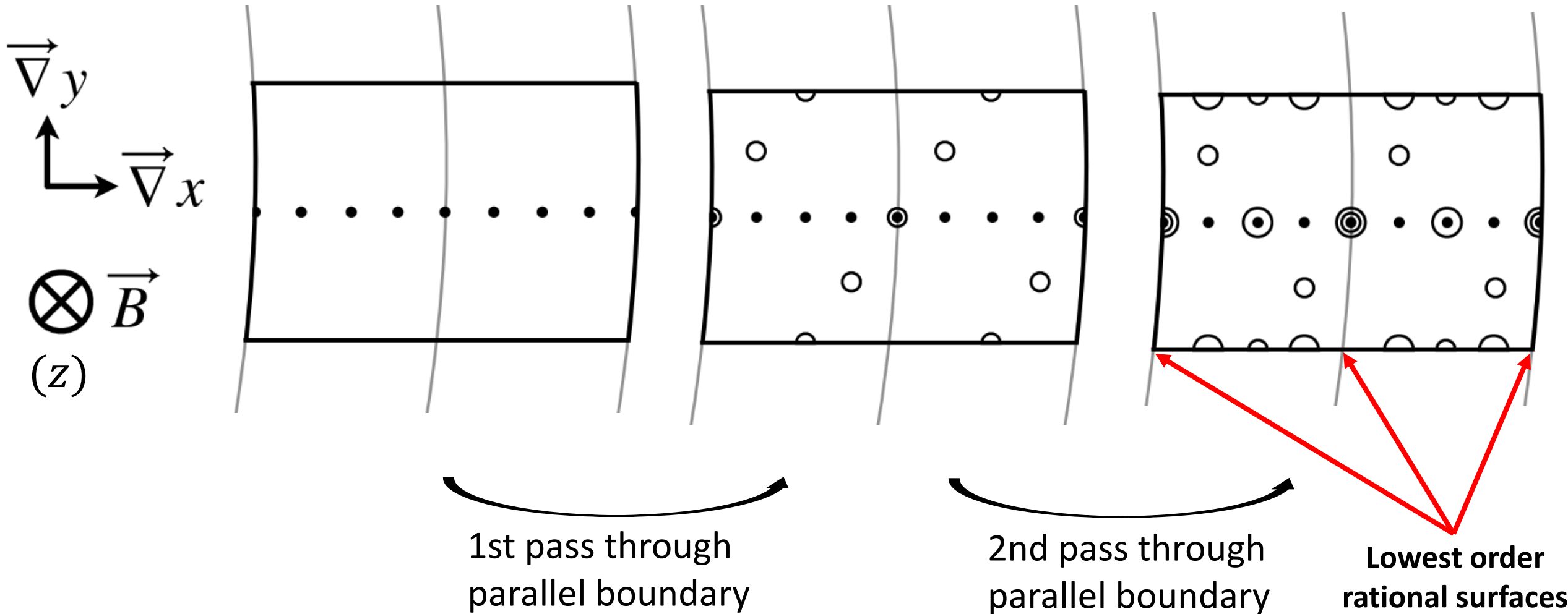


1st pass through parallel boundary



2nd pass through parallel boundary

Twist-and-shift $\hat{s} \neq 0$

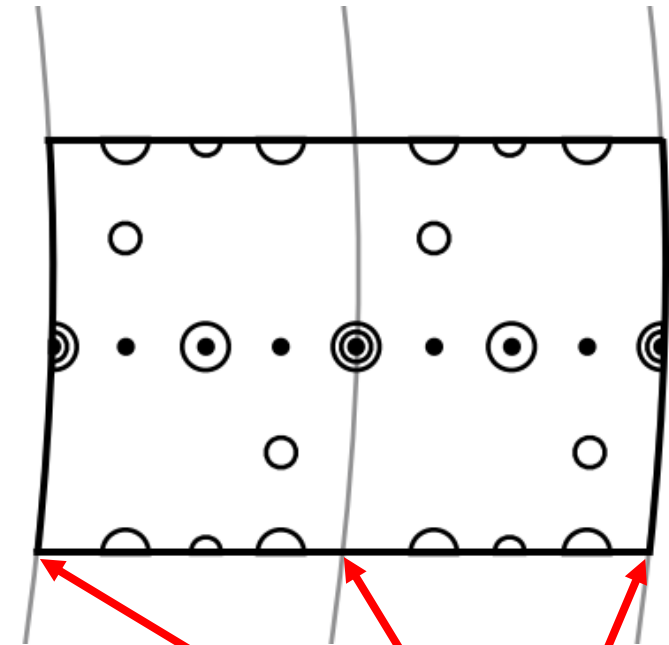


EPFL N_{pol} - adjusted to safety factor

$$q = 2 = \frac{2}{1} \rightarrow N_{pol} = 1$$

$$q = 2.5 = \frac{5}{2} \rightarrow N_{pol} = 2$$

$$q = 2.125 = \frac{17}{8} \rightarrow N_{pol} = 8$$



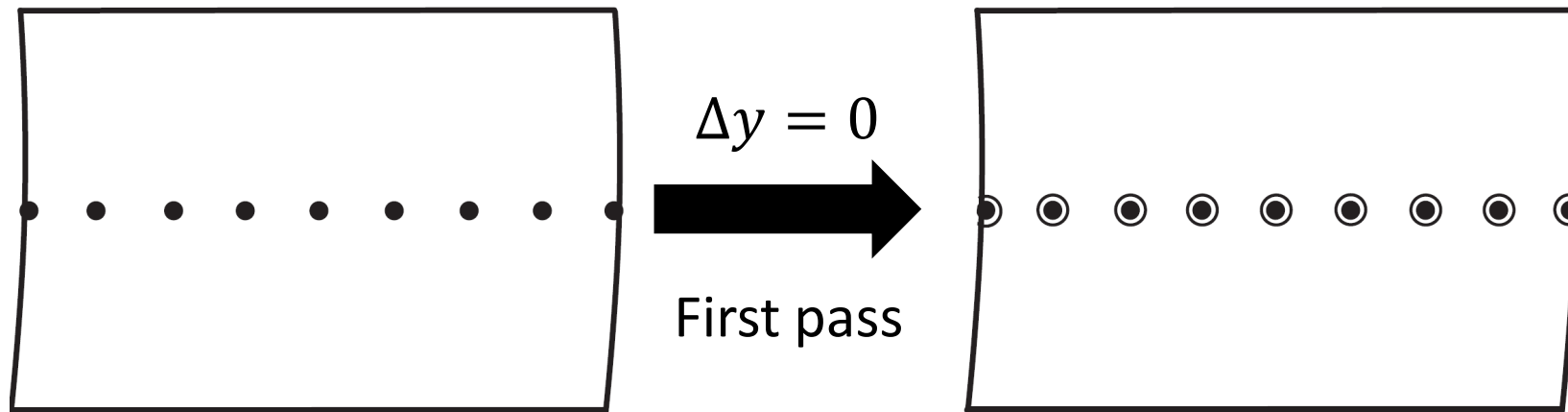
Lowest order
rational surfaces

EPFL Parallel boundary $\hat{s} = 0$

Boundary condition: $\hat{A}_{k_x, k_y}(z) = C(\Delta y) \hat{A}_{k_x, k_y}(z + 2\pi)$

Phase factor: $C(\Delta y) \neq 1$

Shift: Δy

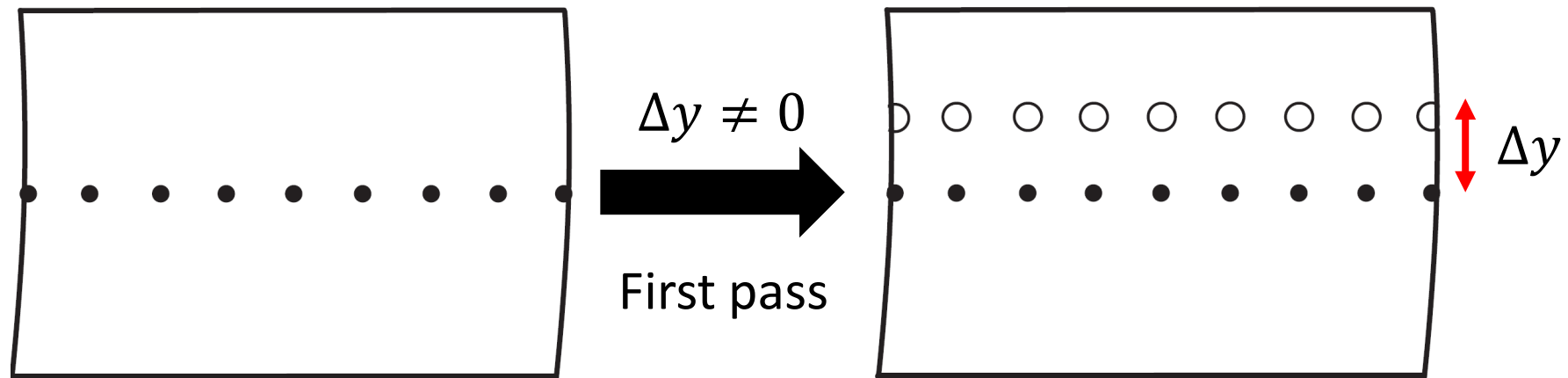


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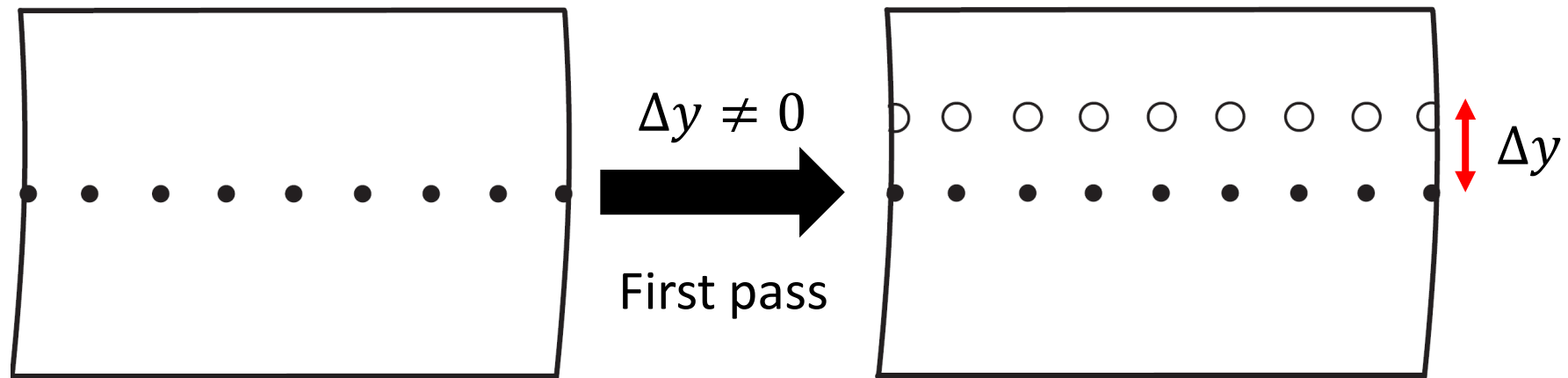


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Shift: Δy



$$q \Leftrightarrow N_{pol} \text{ and } \Delta y$$

Usual simulation parameters

- $\beta = 10^{-4}$
- Kinetic electrons (!)
- Two cases – Cyclone Base Case (CBC) or pure ITG (pITG) drive

- $T_e = T_i$

- $R/L_T = 6.96$

- $R/L_n = 2.22$

- $q = 1.4$

- $T_e = T_i$

- $R/L_{T,i} = 6.96$

- $R/L_{T,e} = 0$

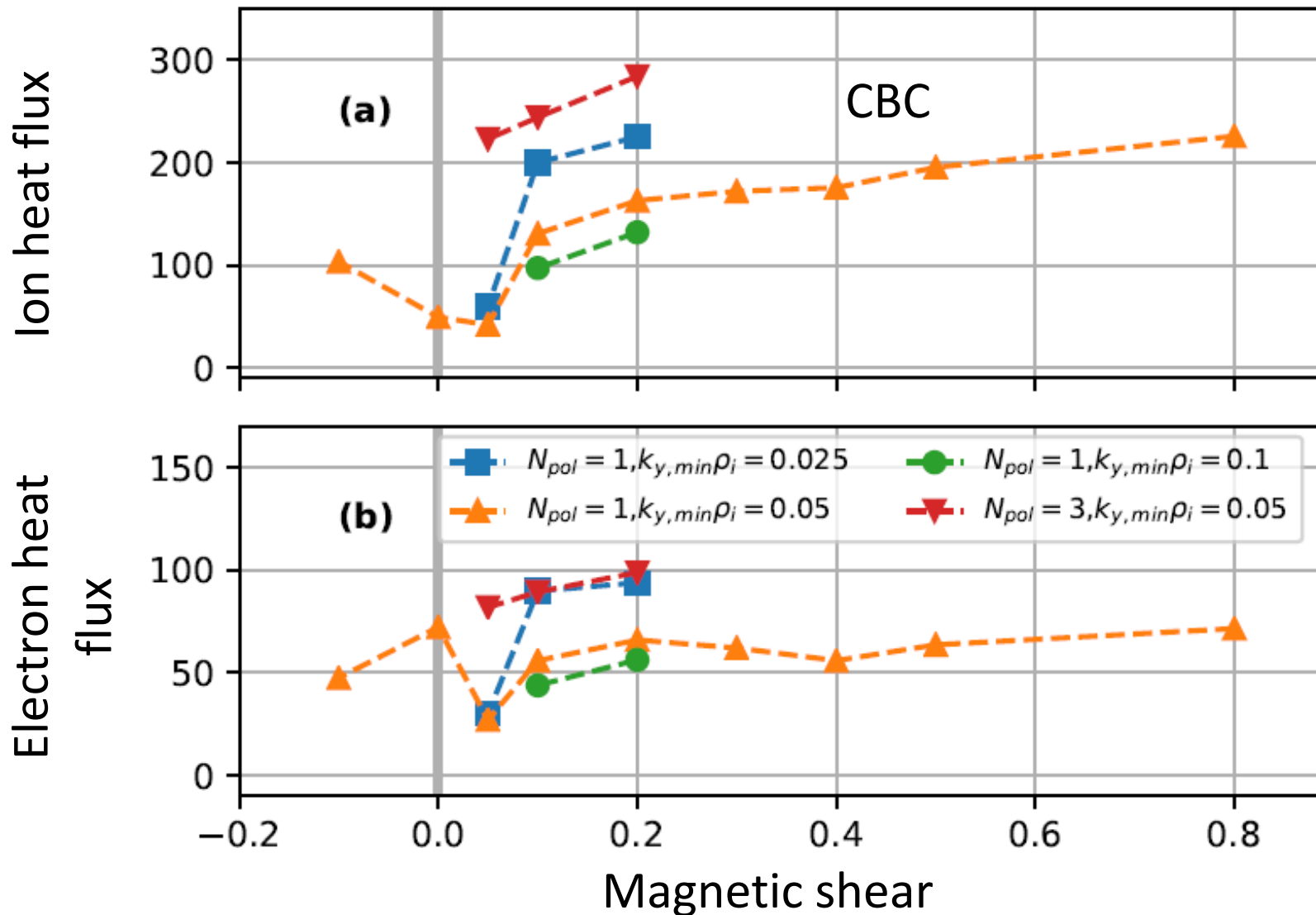
- $R/L_n = 0$

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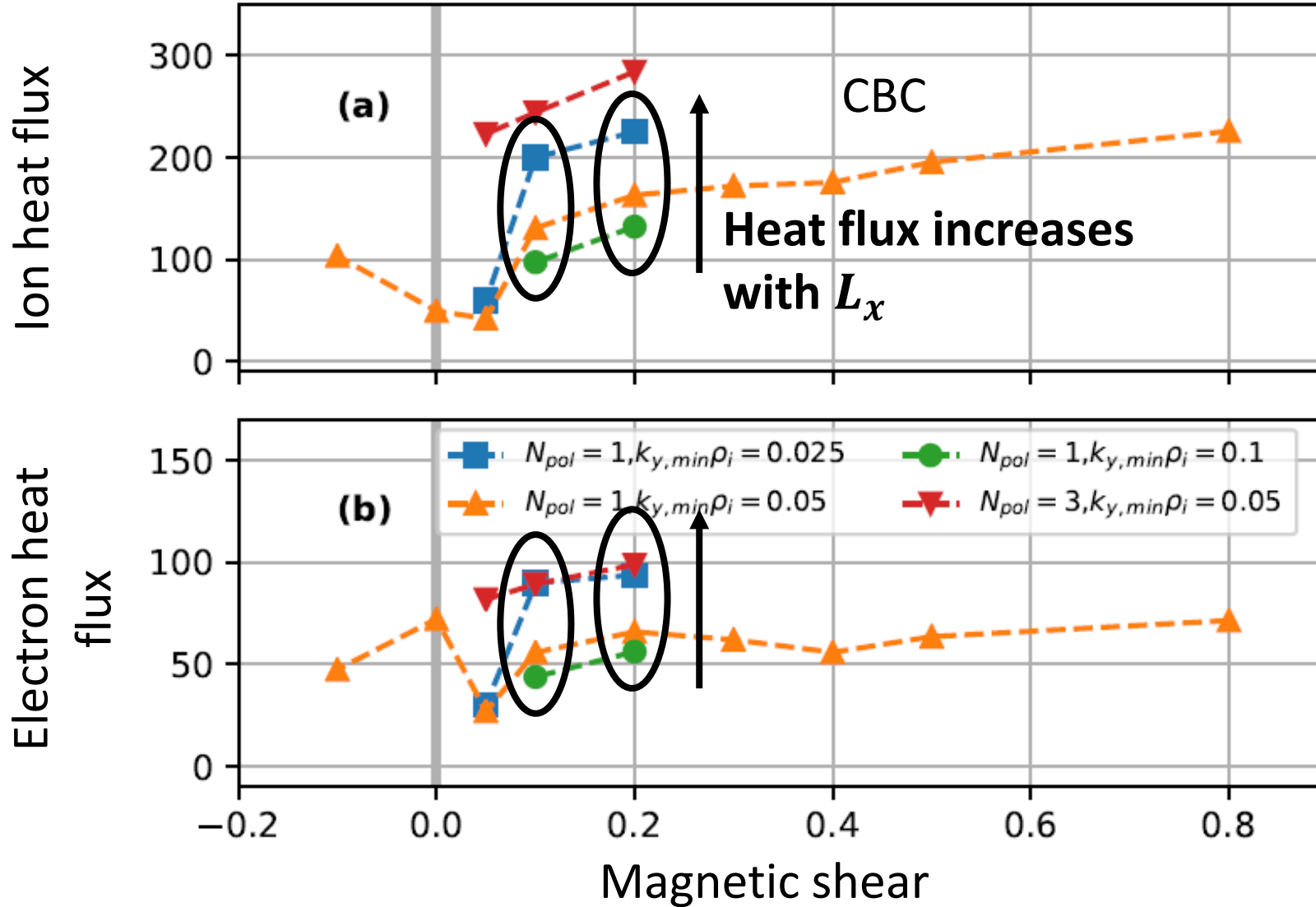
Heat flux at finite shear



Self-interaction trends at low shear agree with previous study at $\hat{s} = 0.8$

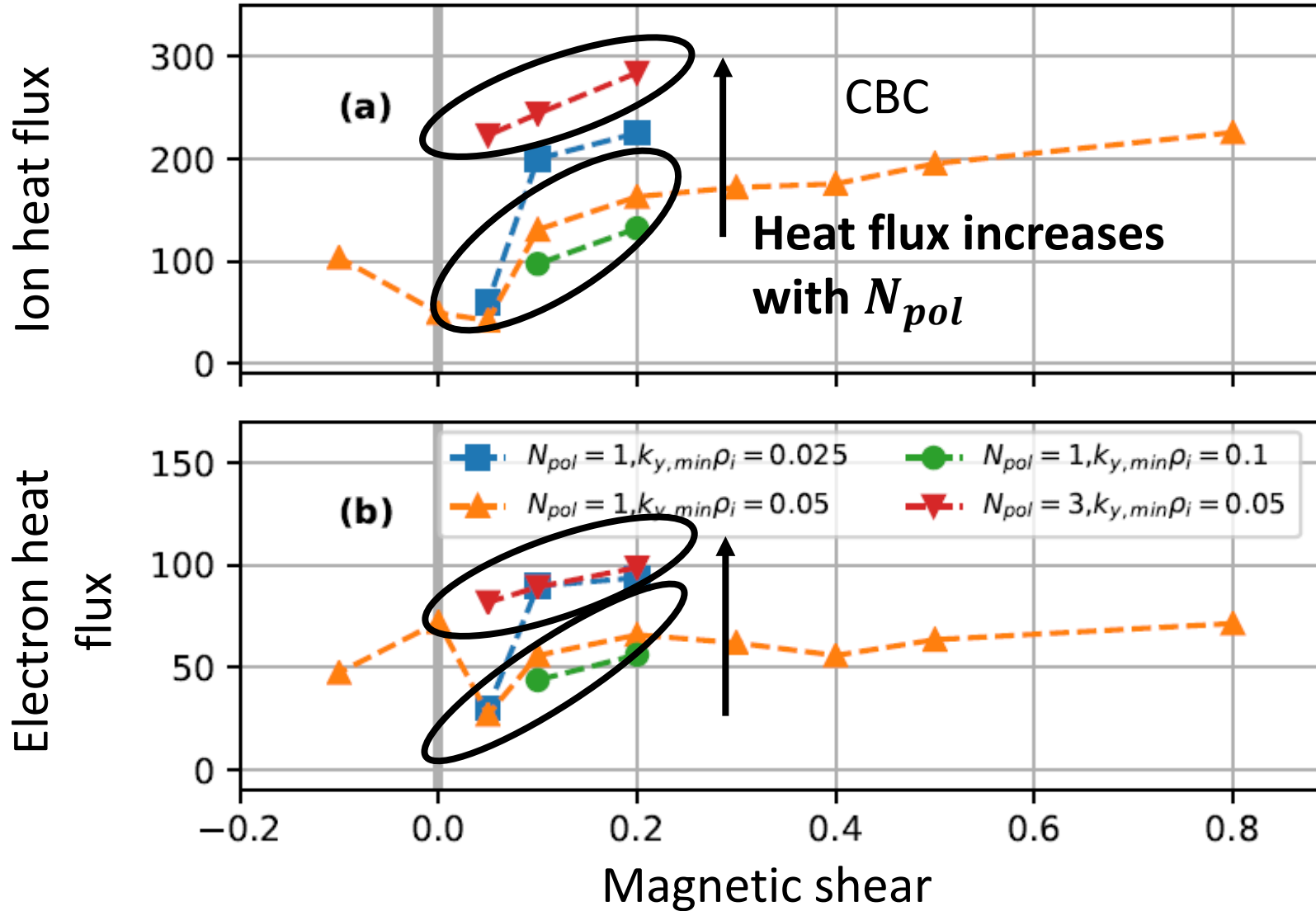
• J. Ball et al. 2020 *Journal of Plasma Physics* **86**(2), 905860207

By spacing out integer surfaces



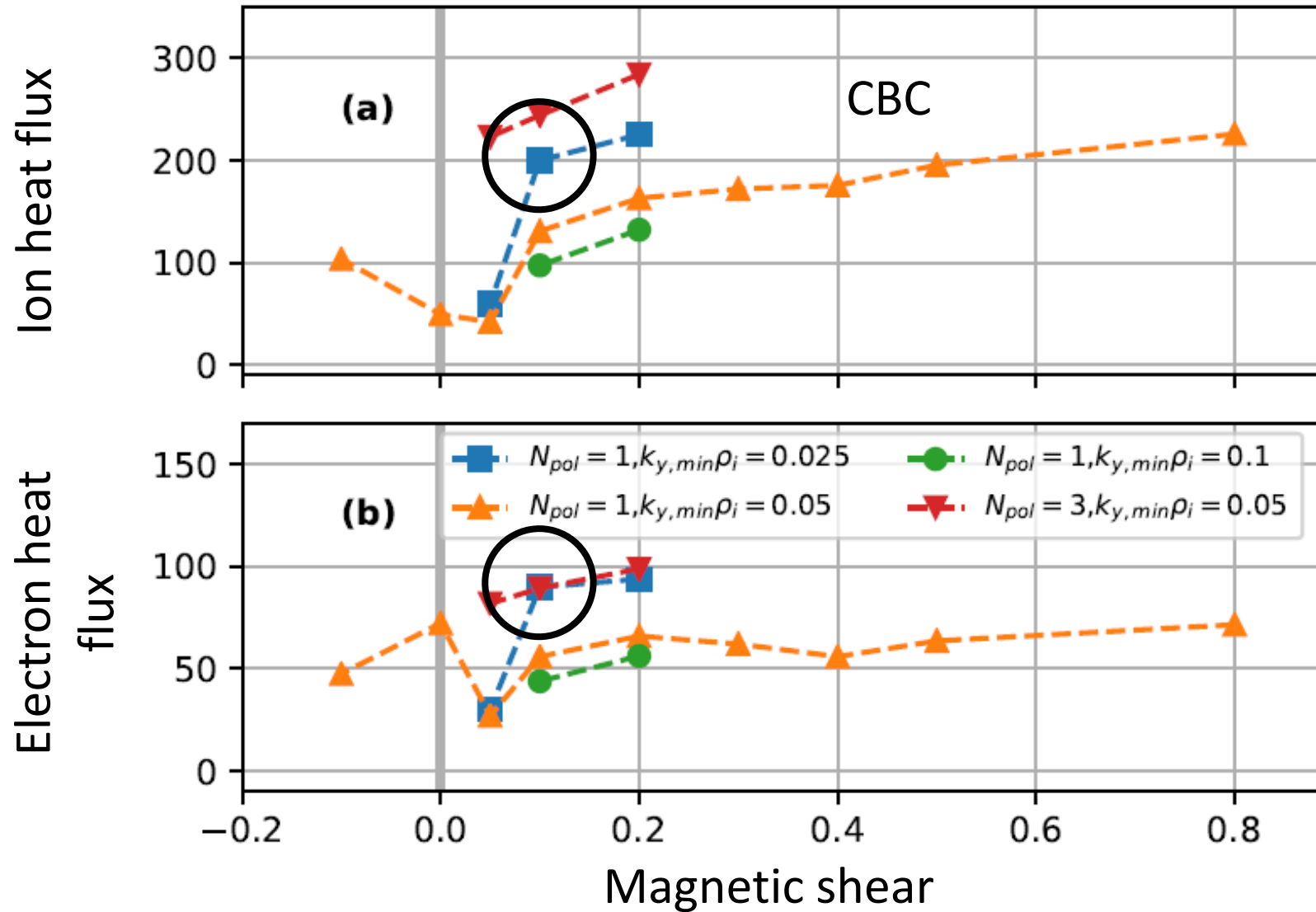
$$L_x \propto (k_{y,min}\hat{S})^{-1}$$

By increasing domain length



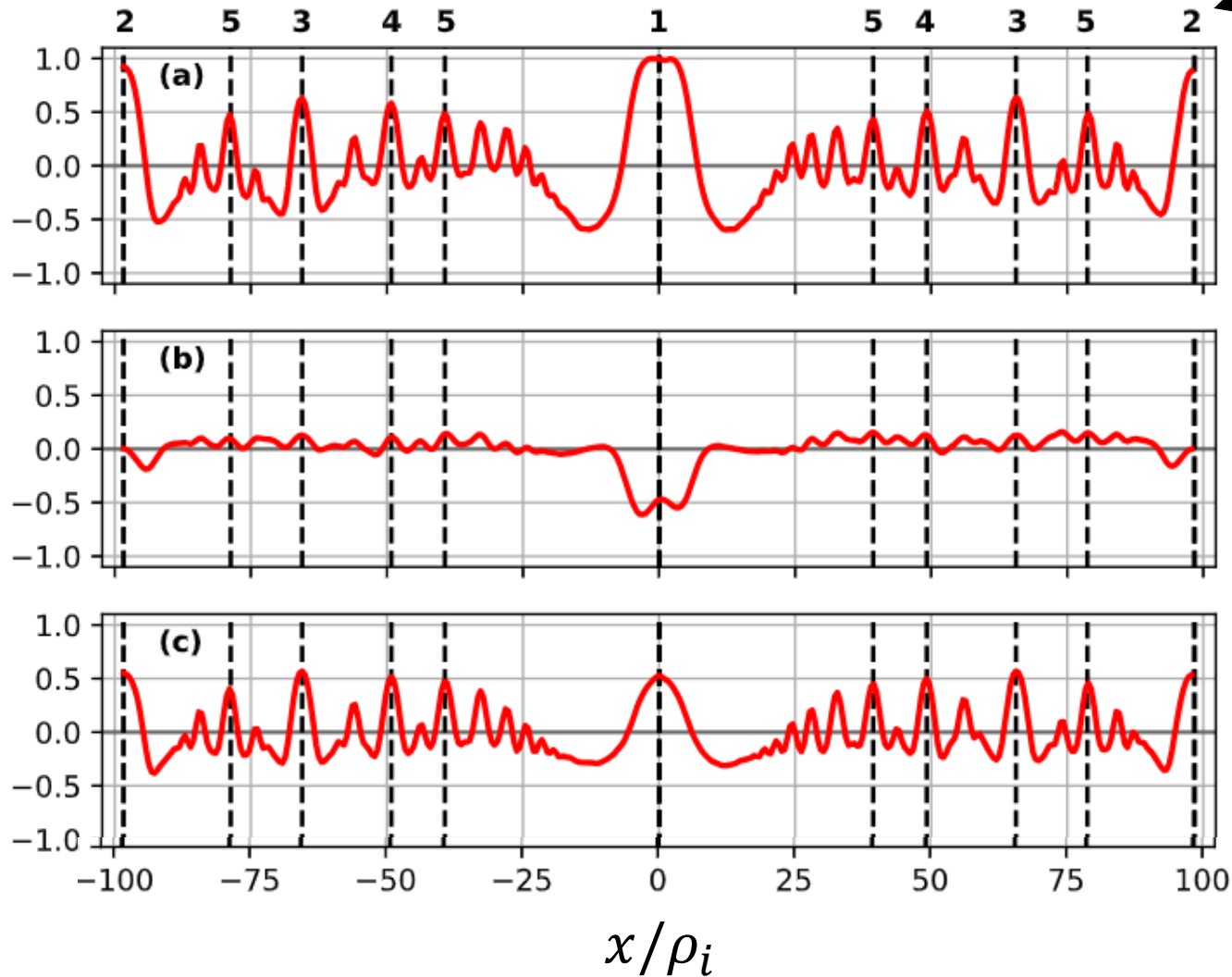
Due to
finite eddy
length

Radial plasma profiles at $\hat{s} = 0.1$



EPFL Profile corrugations

Order of rational surface



Perturbed density gradient

Perturbed ion temperature gradient

Perturbed electron temperature gradient

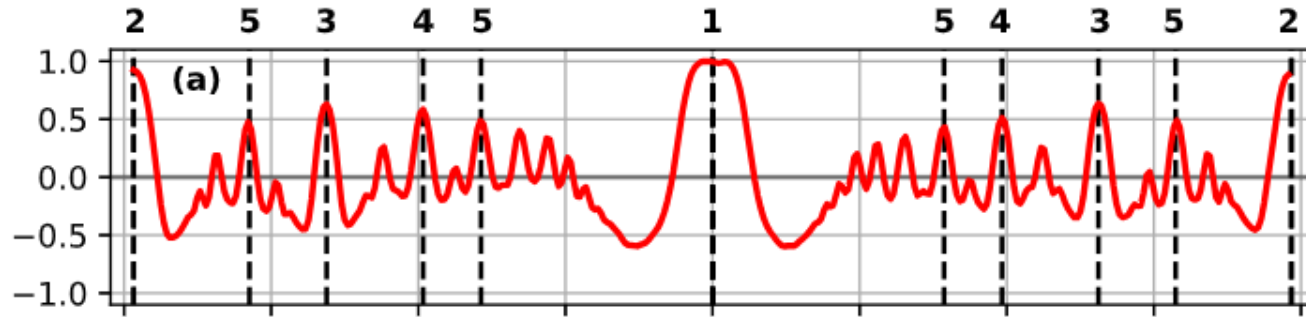
$$\propto \partial_x \langle \delta n \rangle_{y,z,t}$$

$$\hat{s} = 0.1$$

$$N_{pol} = 1$$

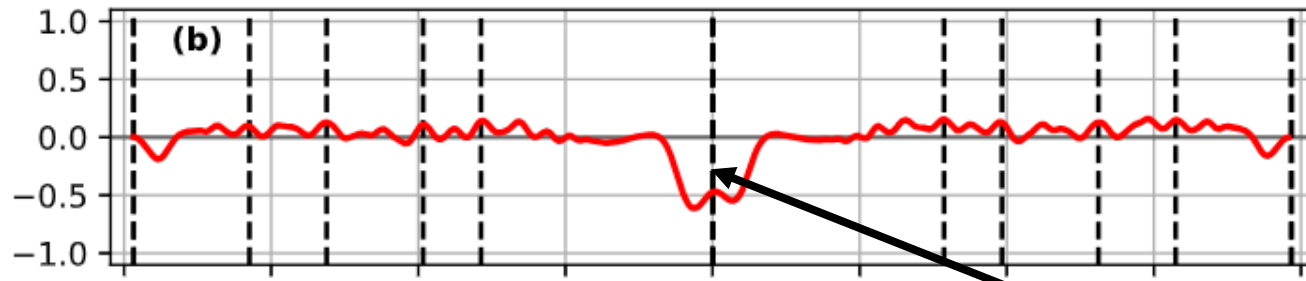
EPFL Profile corrugations

Perturbed
density gradient



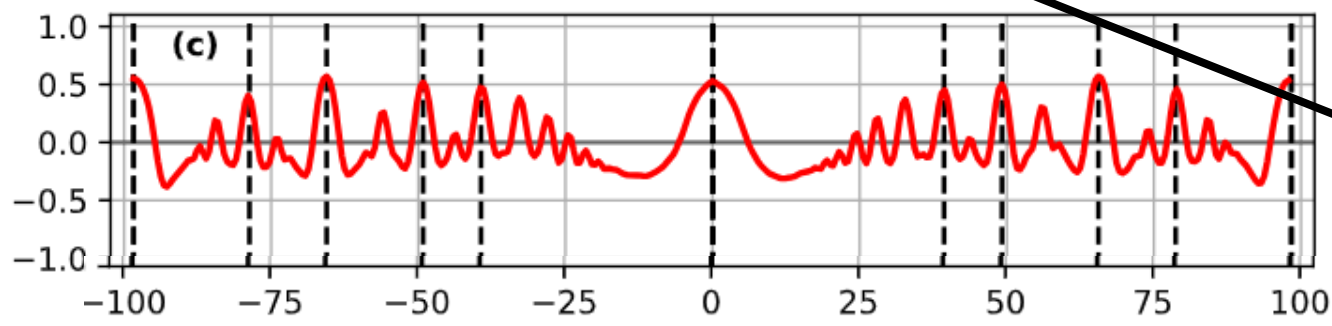
Positive → flattening
Negative → steepening

Perturbed
ion temperature
gradient



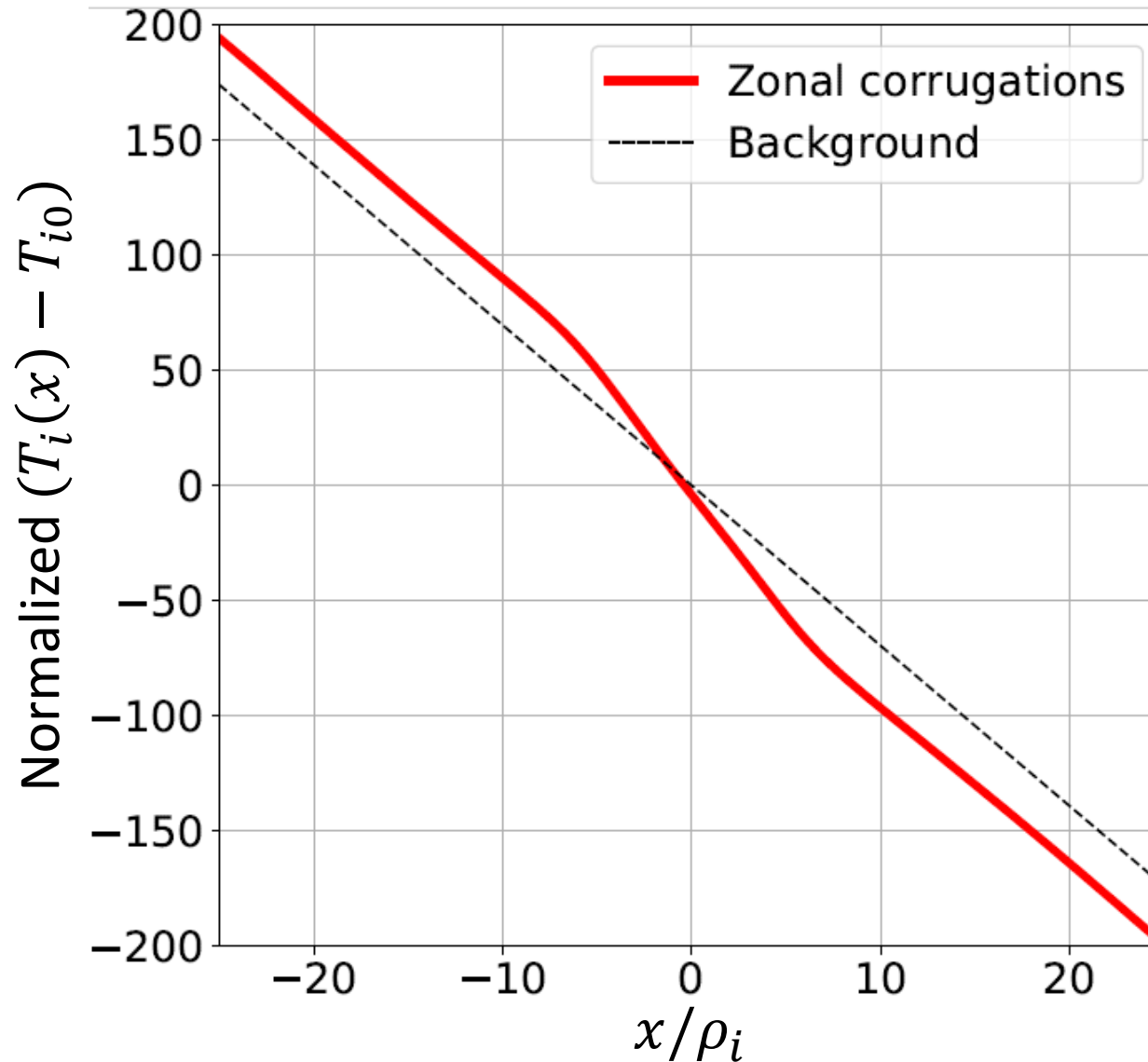
$\hat{s} = 0.1$
 $N_{pol} = 1$

Perturbed
electron
temperature
gradient



Profile
steepening
by 50 %

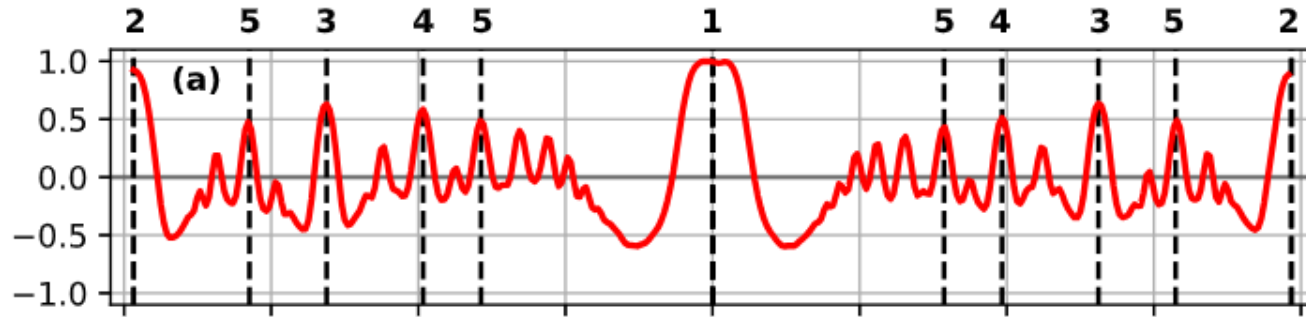
x/ρ_i



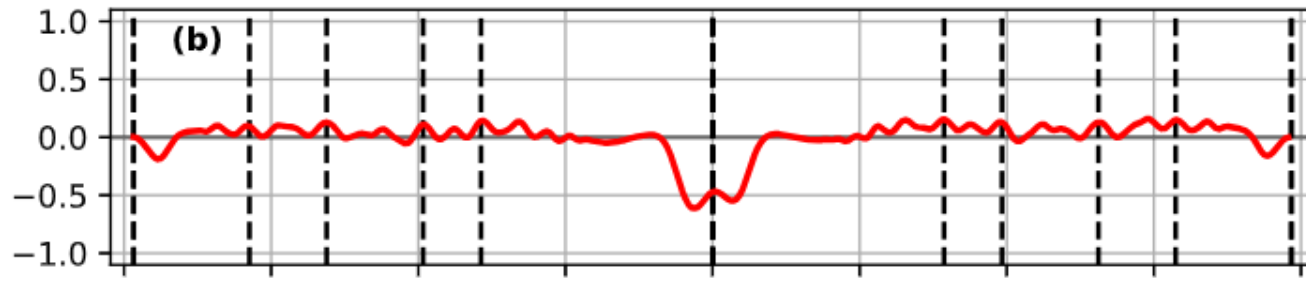
$$\hat{s} = 0.1$$
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Estimate of eddy length with \hat{s}

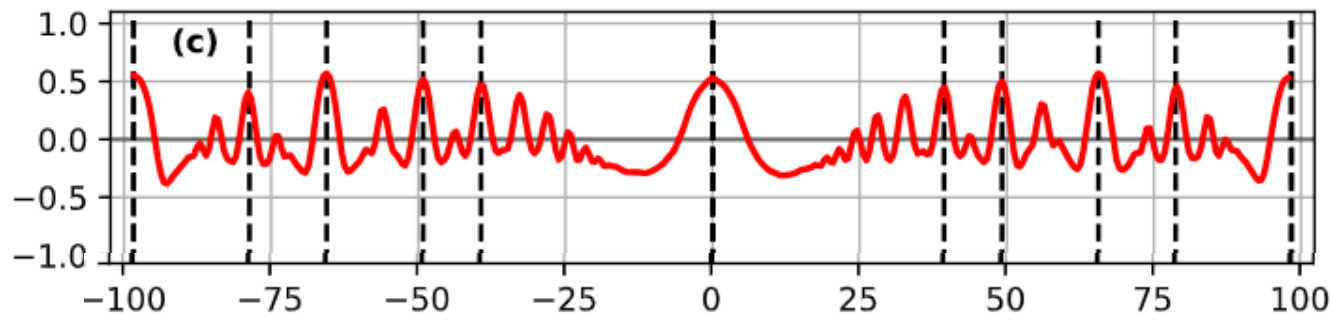
Perturbed
density gradient



Perturbed
ion temperature
gradient



Perturbed
electron
temperature
gradient

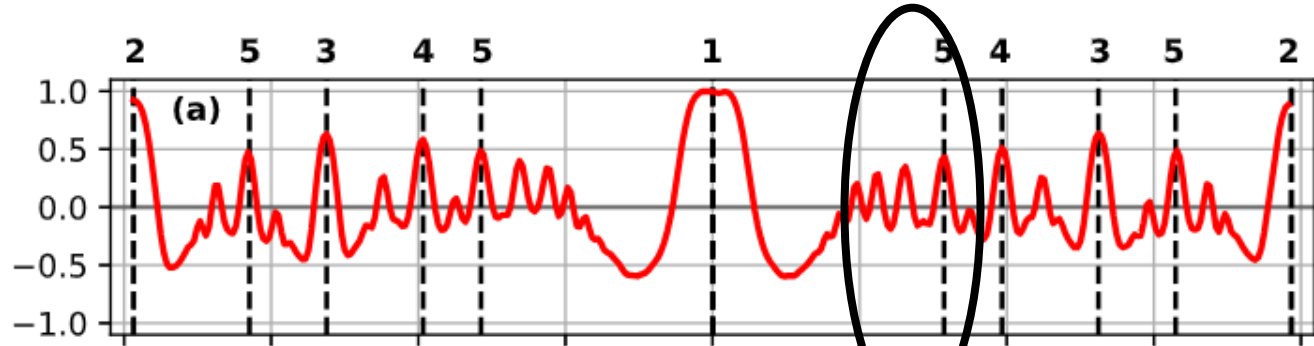


$\hat{s} = 0.1$
 $N_{pol} = 1$

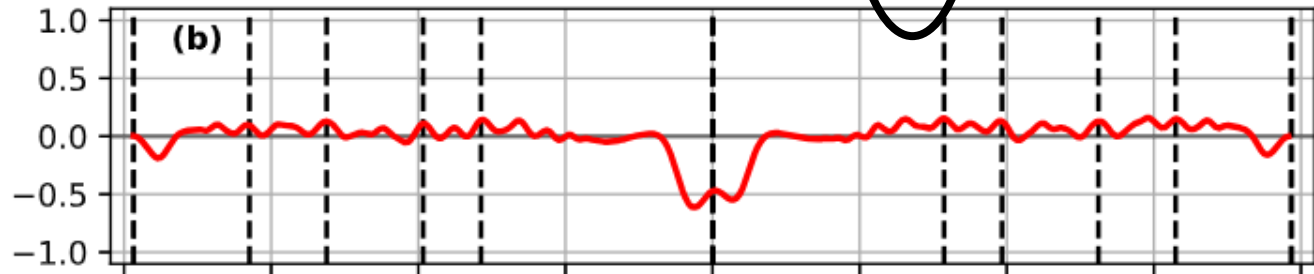
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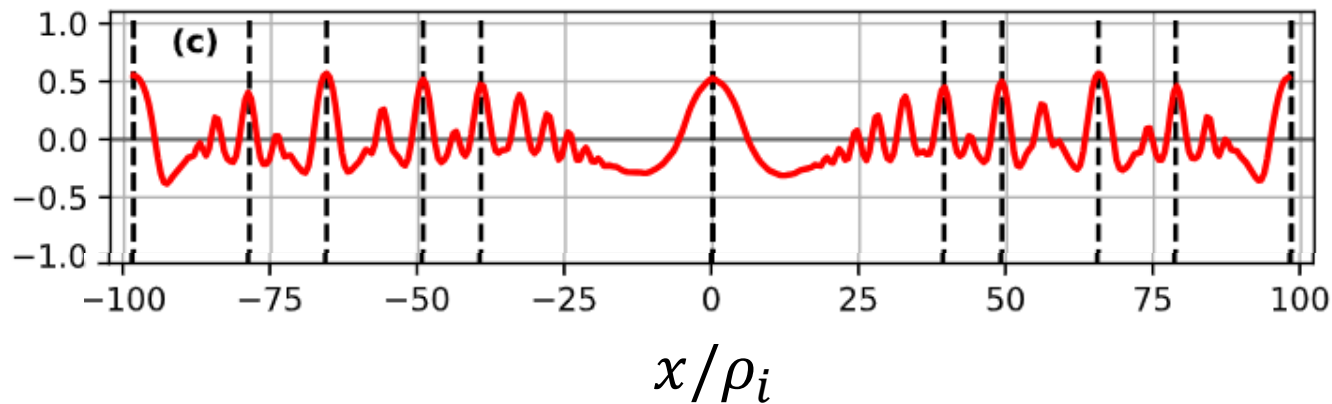
Perturbed
density gradient



Perturbed
ion temperature
gradient

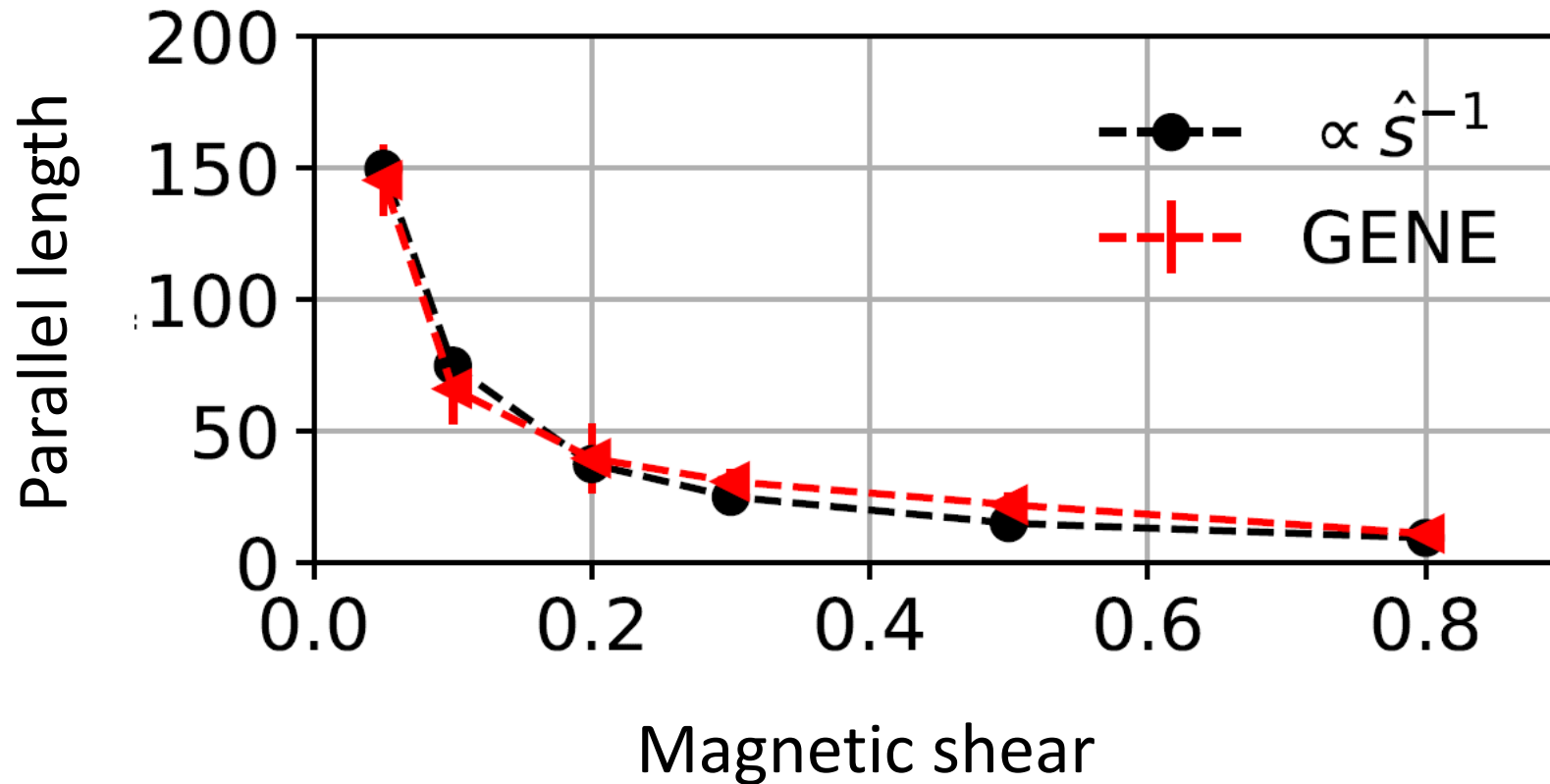


Perturbed
electron
temperature
gradient



Based on the highest
order rational
surfaces with
significant zonal
structures

Estimate of eddy length with \hat{S}



Based on the highest order rational surfaces with significant zonal structures

Due to **magnetic drifts** and **FLR effects** parallel eddy length scales like

$$\hat{S}^{-1}$$

EPFL Summary: Finite shear study

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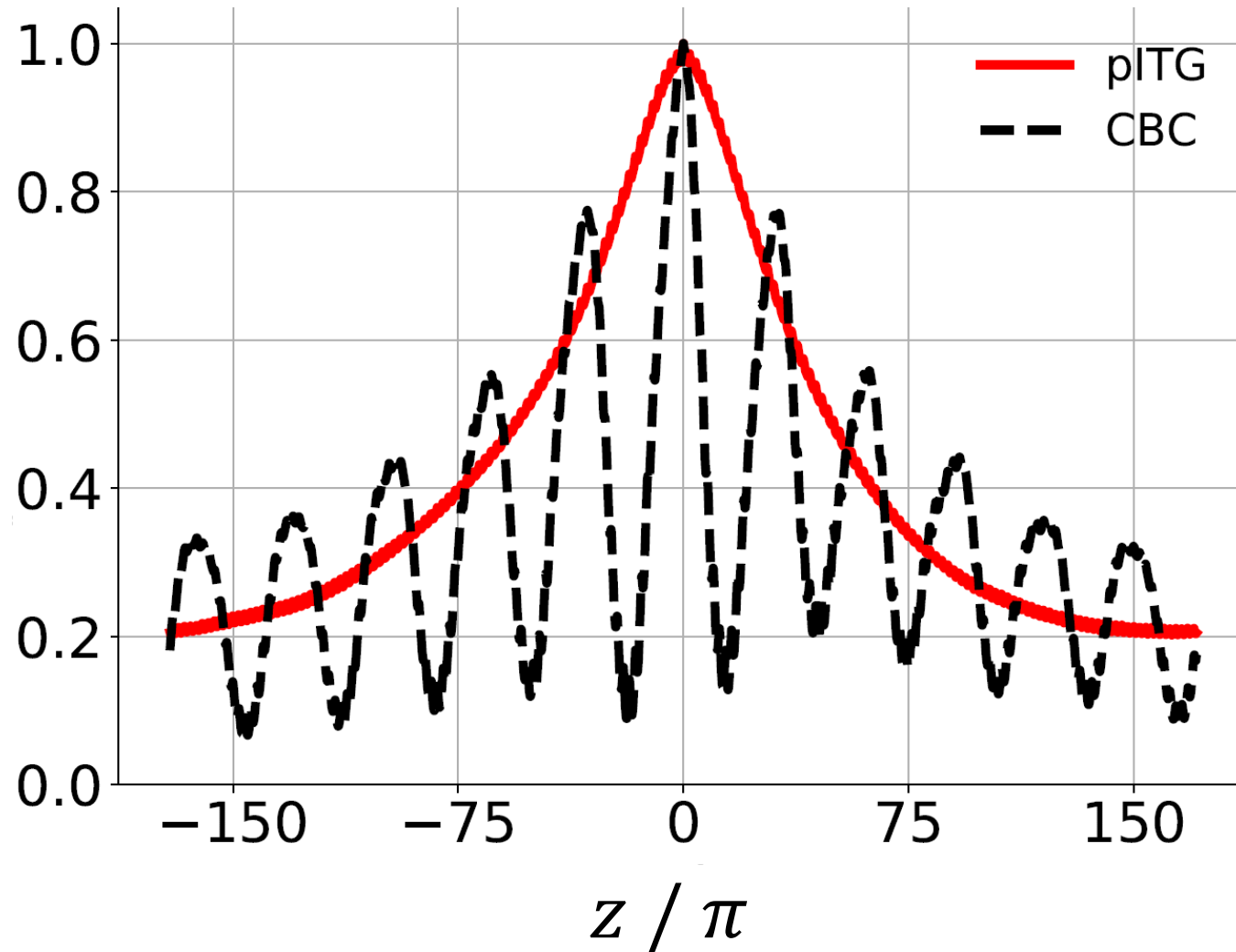
Summary: Finite shear study

- Self-interaction follows the same paradigm as established in previous work
- Eddies extend in the parallel direction when shear is reduced as \hat{s}^{-1}
- Profile corrugations appear at multiple rational surfaces
- Simulations become computationally expensive as shear is reduced [$L_x \propto (k_y \hat{s})^{-1}$]

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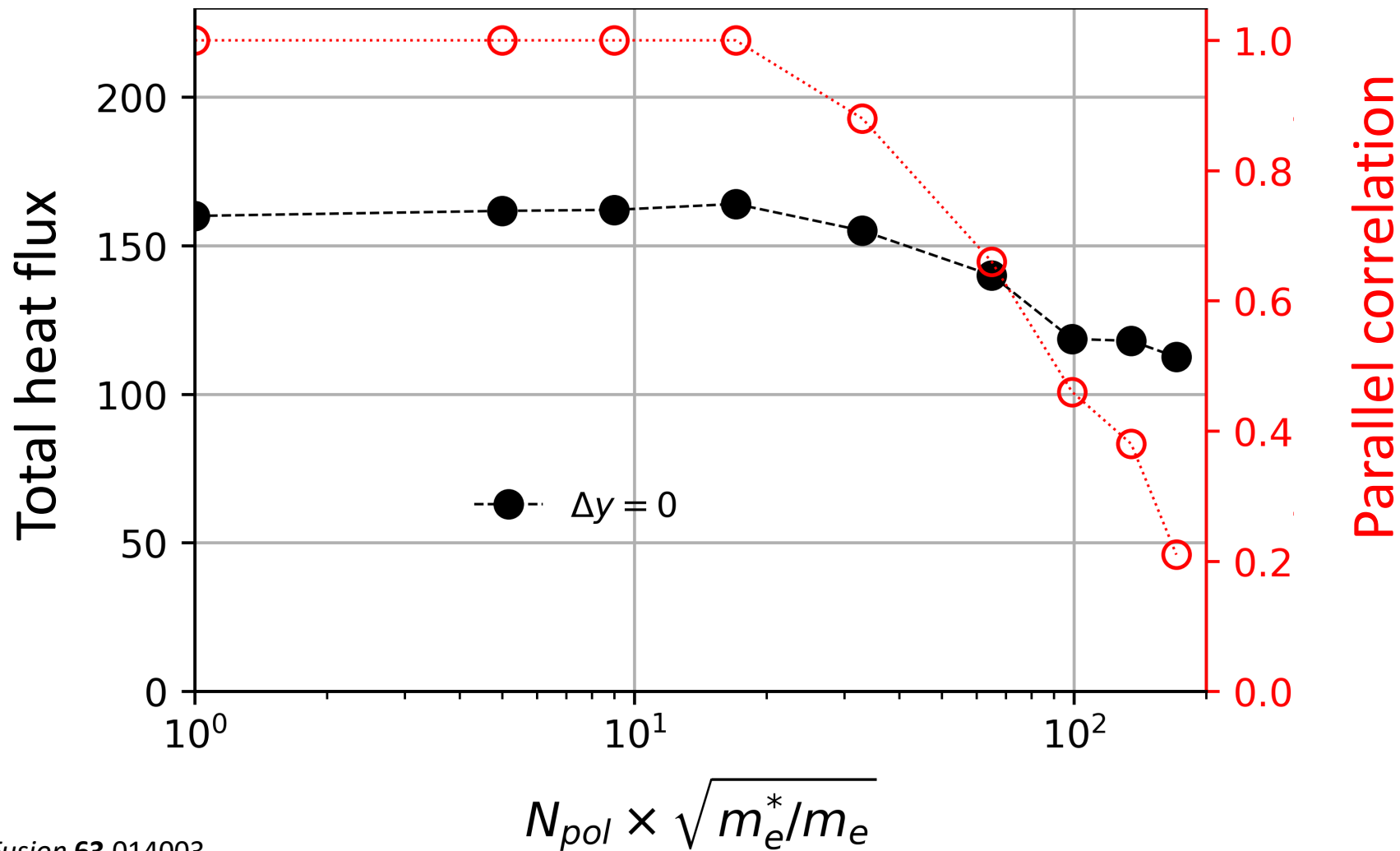
Parallel correlation
along field lines



$$N_{pol} = 171$$

pITG heat flux and correlation

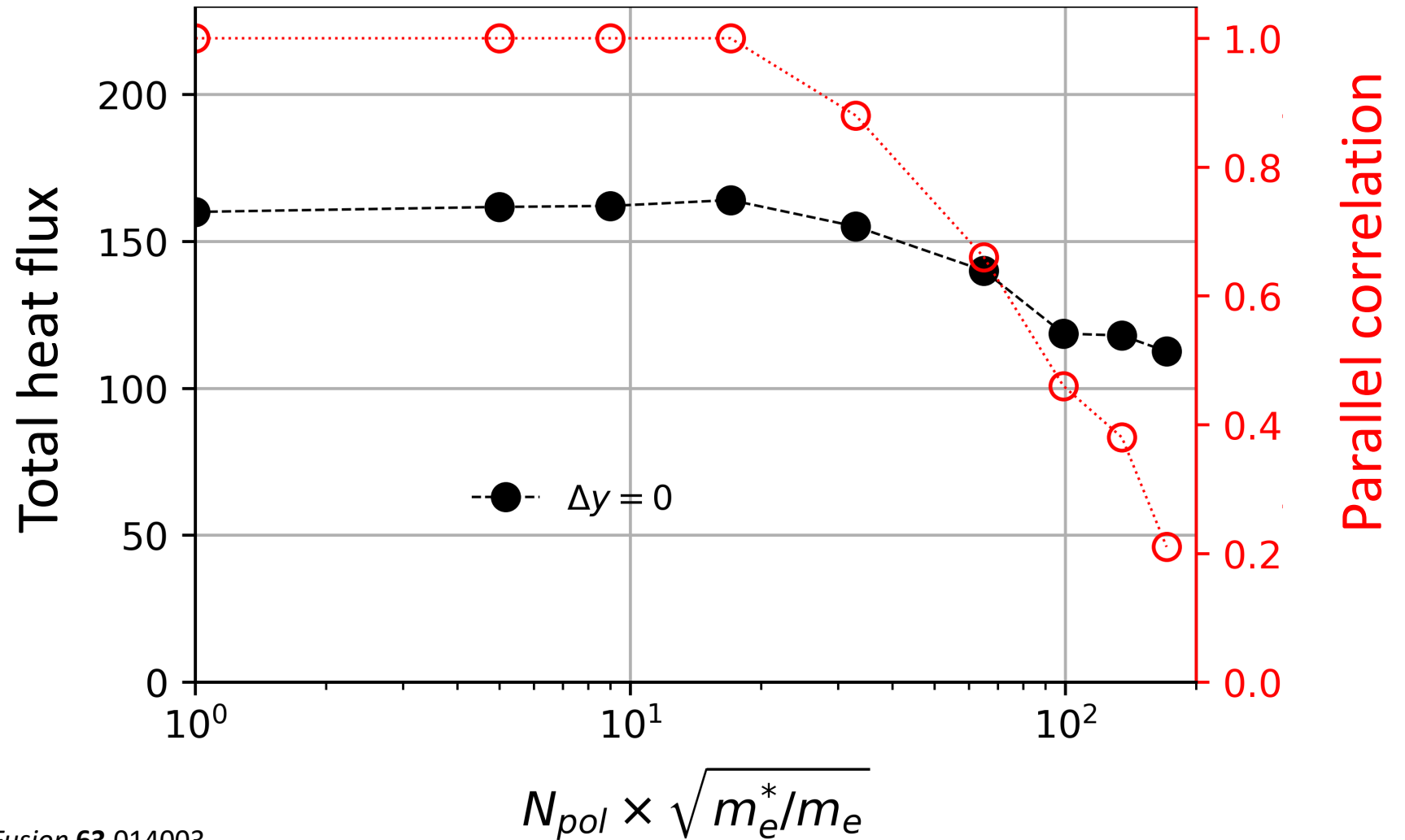
Periodic parallel
boundary condition
($\Delta y = 0$)



EPFL pITG heat flux and correlation

**Kinetic electrons
are key**

Periodic parallel
boundary condition
($\Delta y = 0$)



EPFL Electrons set eddy length

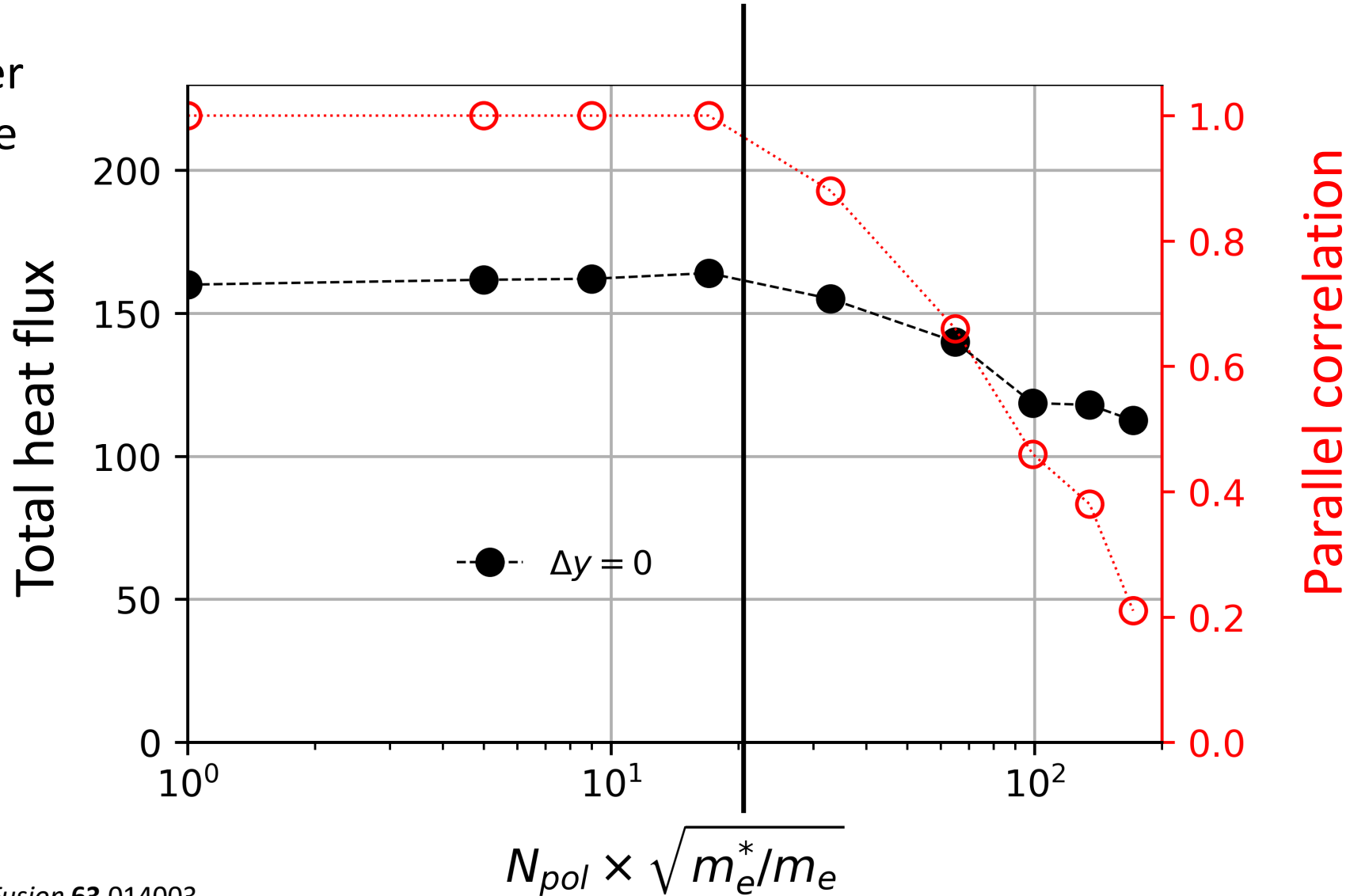
How far **electrons** travel over turbulent time scale sets the parallel eddy size.

$$\sqrt{m_i/m_e} \approx 60$$

$$\tau_{turb} \approx 3 R/c_i$$

$$L_{\parallel} \approx v_{th,e} \tau_{turb}$$

$$\frac{L_{\parallel}}{L_{N_{pol}=1}} \approx 20$$



EPFL Check with heavier electrons

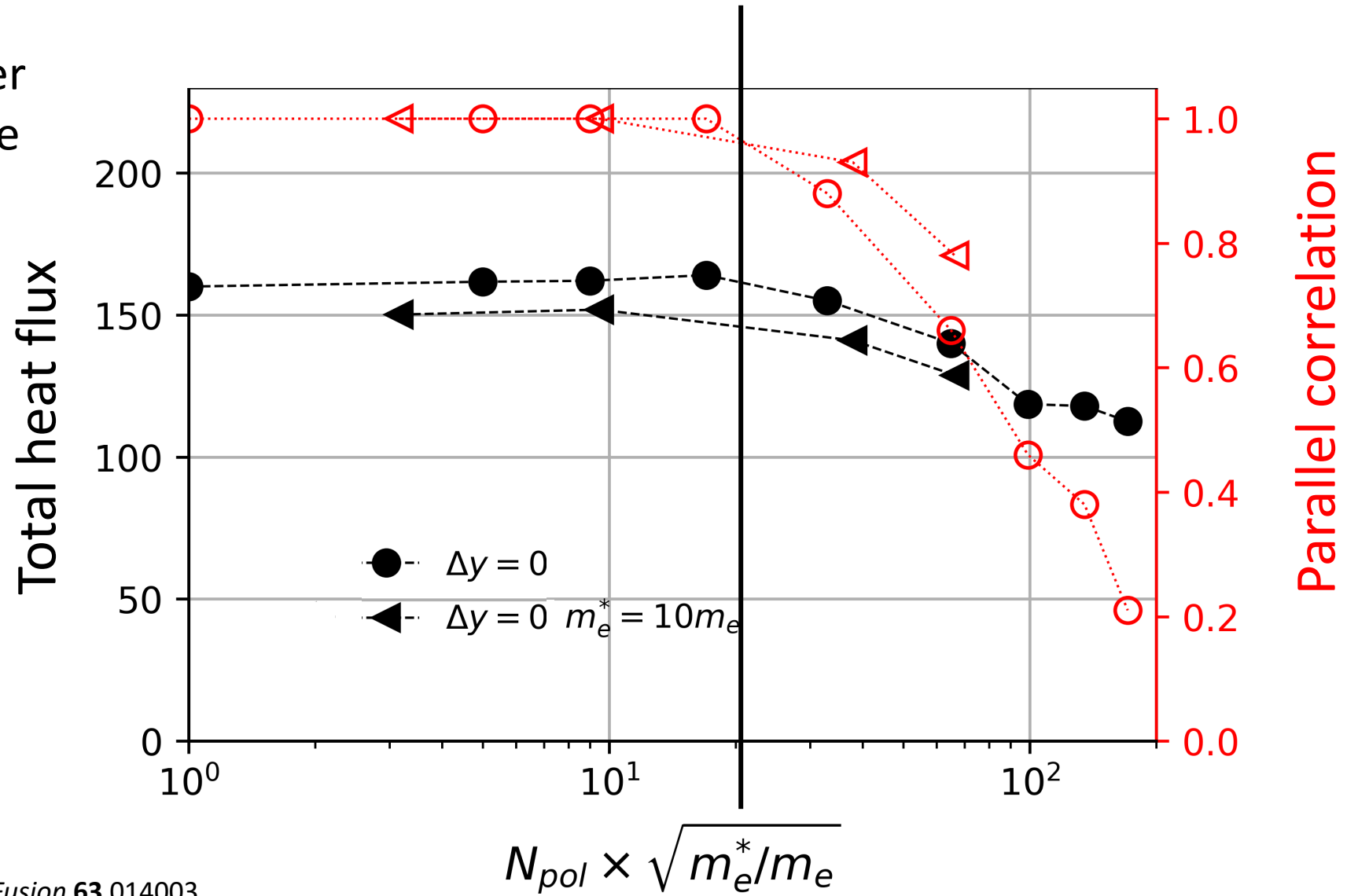
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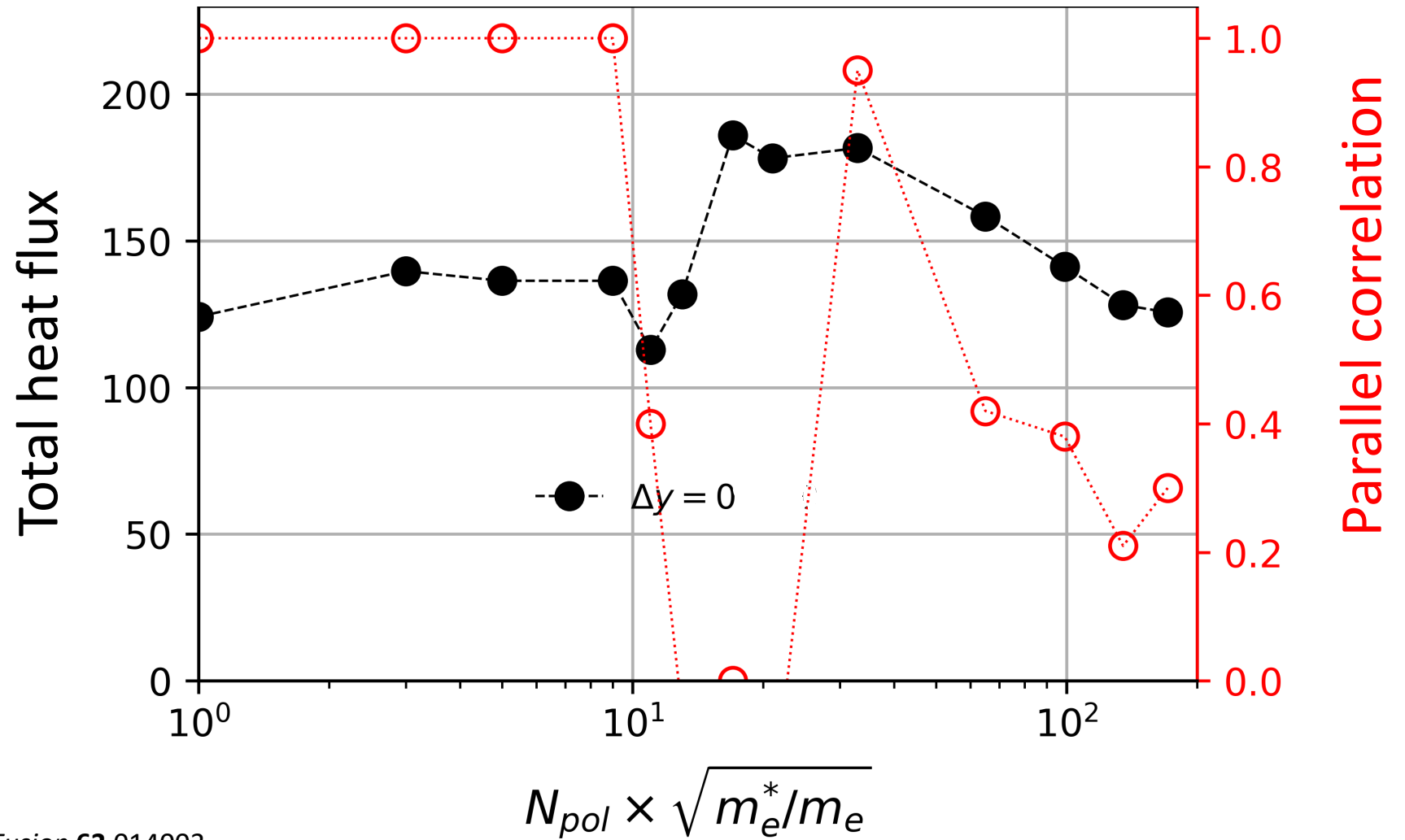
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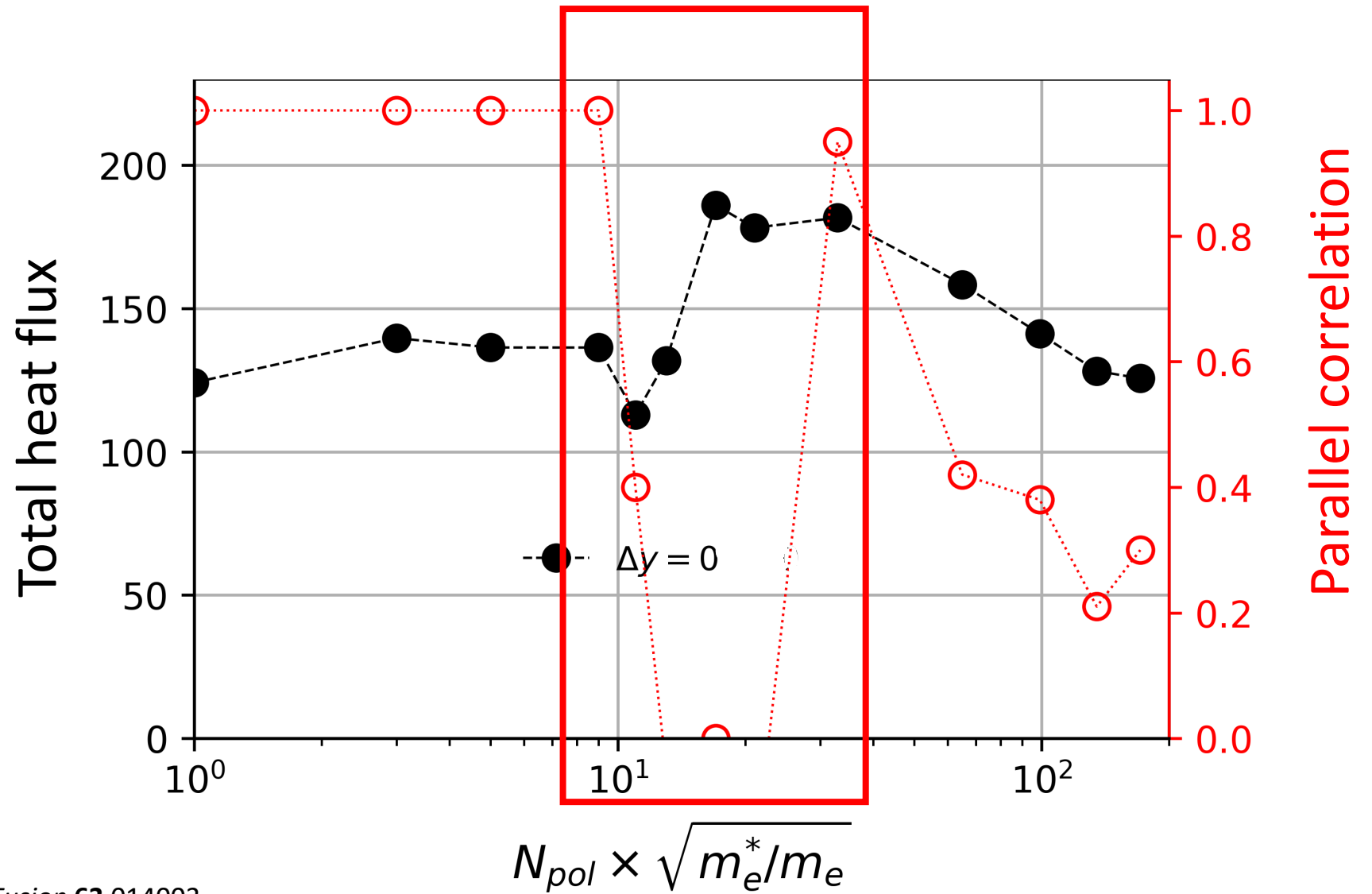


Similar picture to pITG



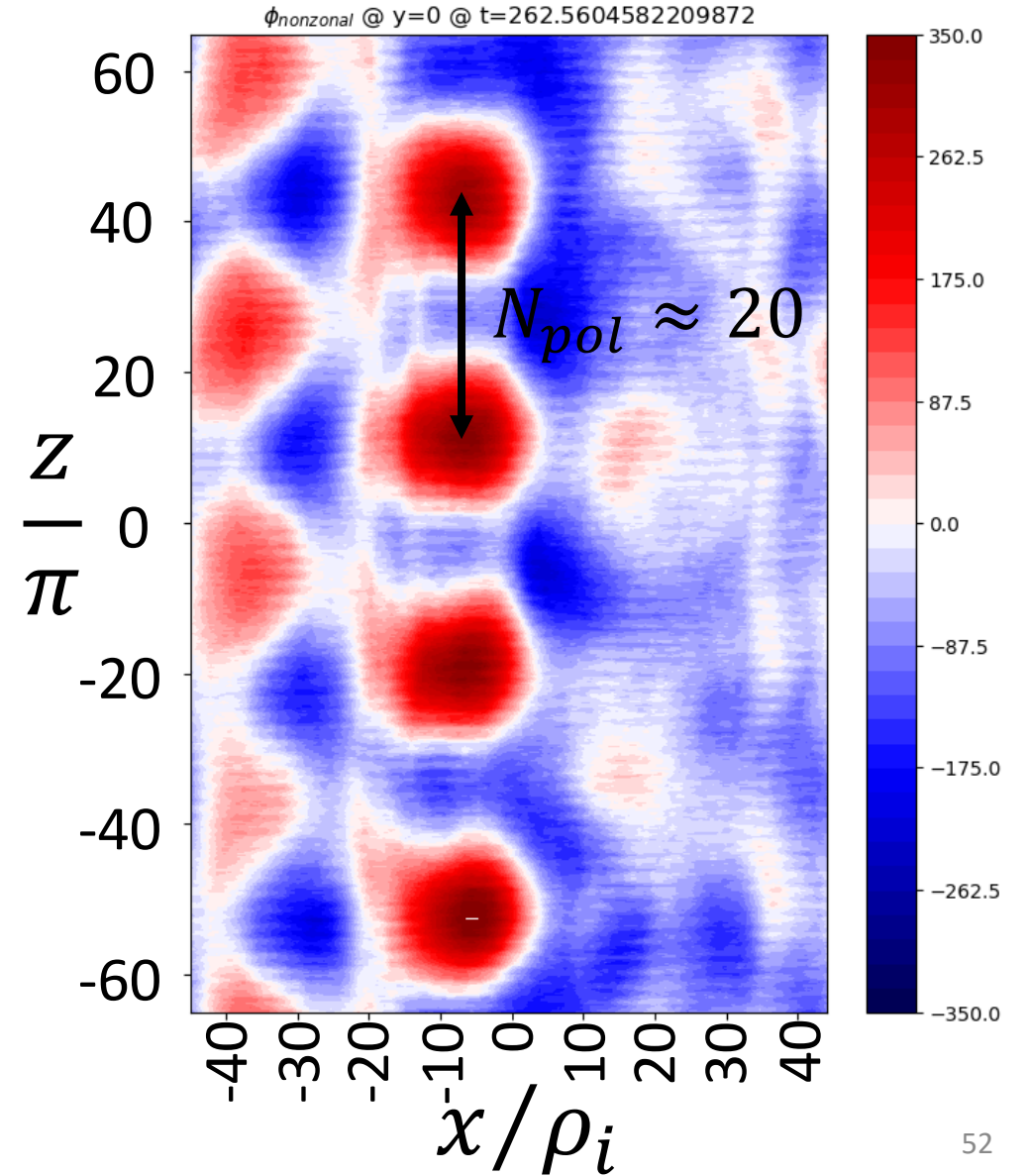
EPFL Parallel waves

Long parallel waves
 $\lambda \approx 20 N_{pol}$ appear



Long parallel wave-like structures

$$N_{pol} = 65$$



Robustness of ultra-long turbulent eddies

- **Collisions** do not reduce ultra-long eddy length but eliminate parallel waves
 - No **plasma shaping** (elongation or triangularity) effects on ultra-long eddies
- Ultra-long eddies seem to **be a robust plasma feature** at low magnetic shear $\hat{s} \ll 1$.

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- Simulations with kinetic electrons at zero magnetic shear **require hundreds of poloidal turns** to achieve convergence

Summary: Eddy parallel length study

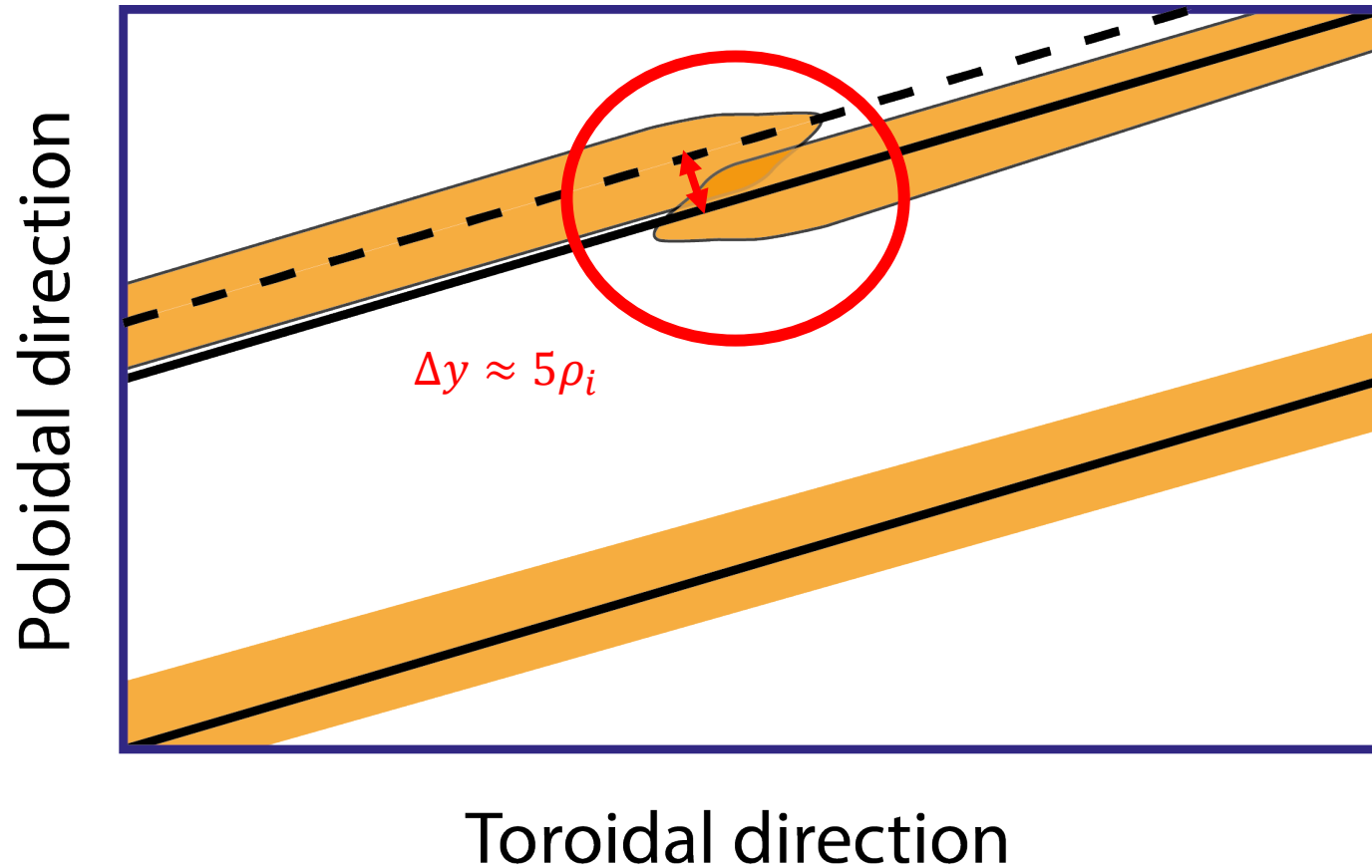
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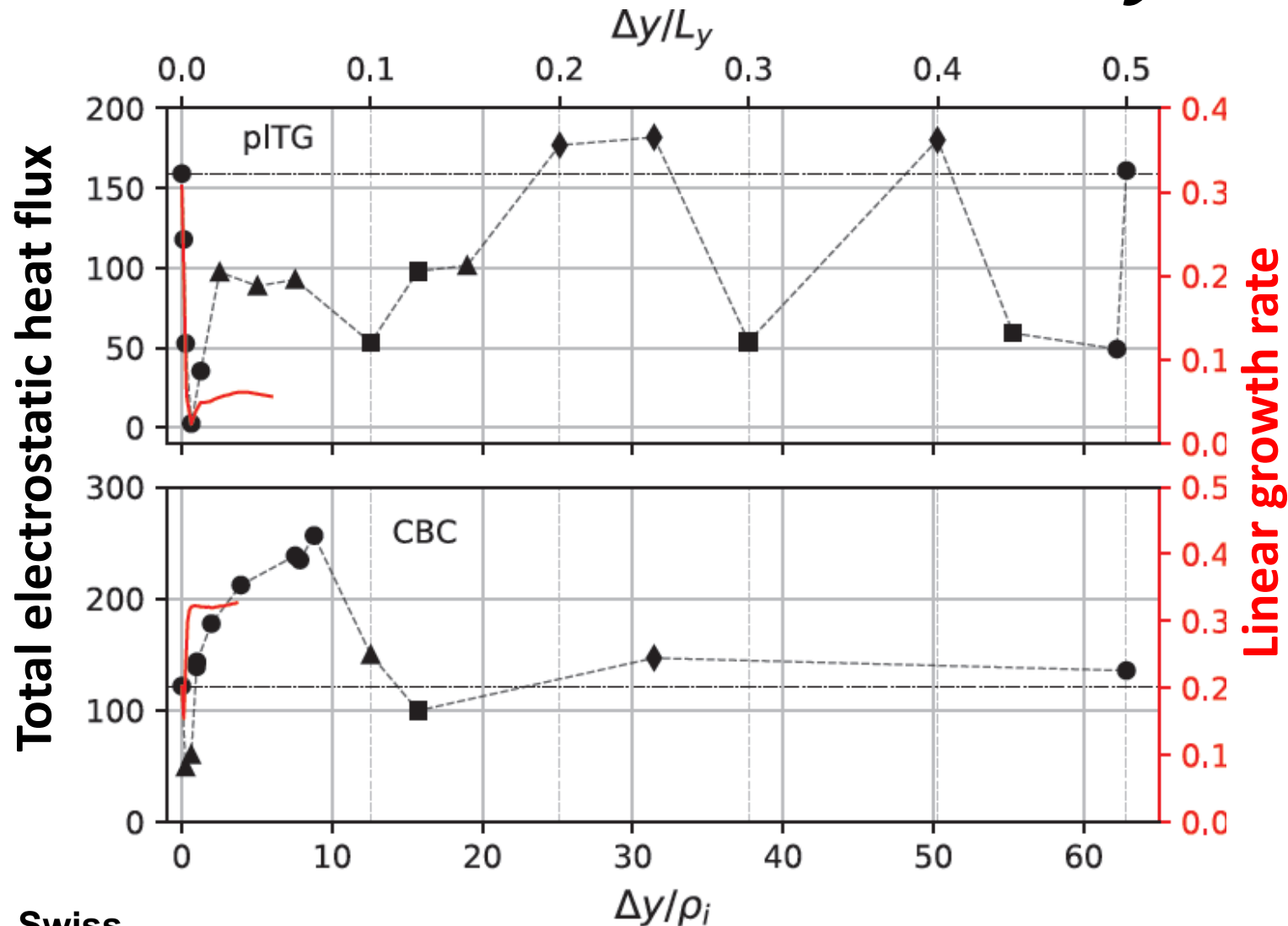
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- Heat flux is reduced when eddies are unable to self-interact

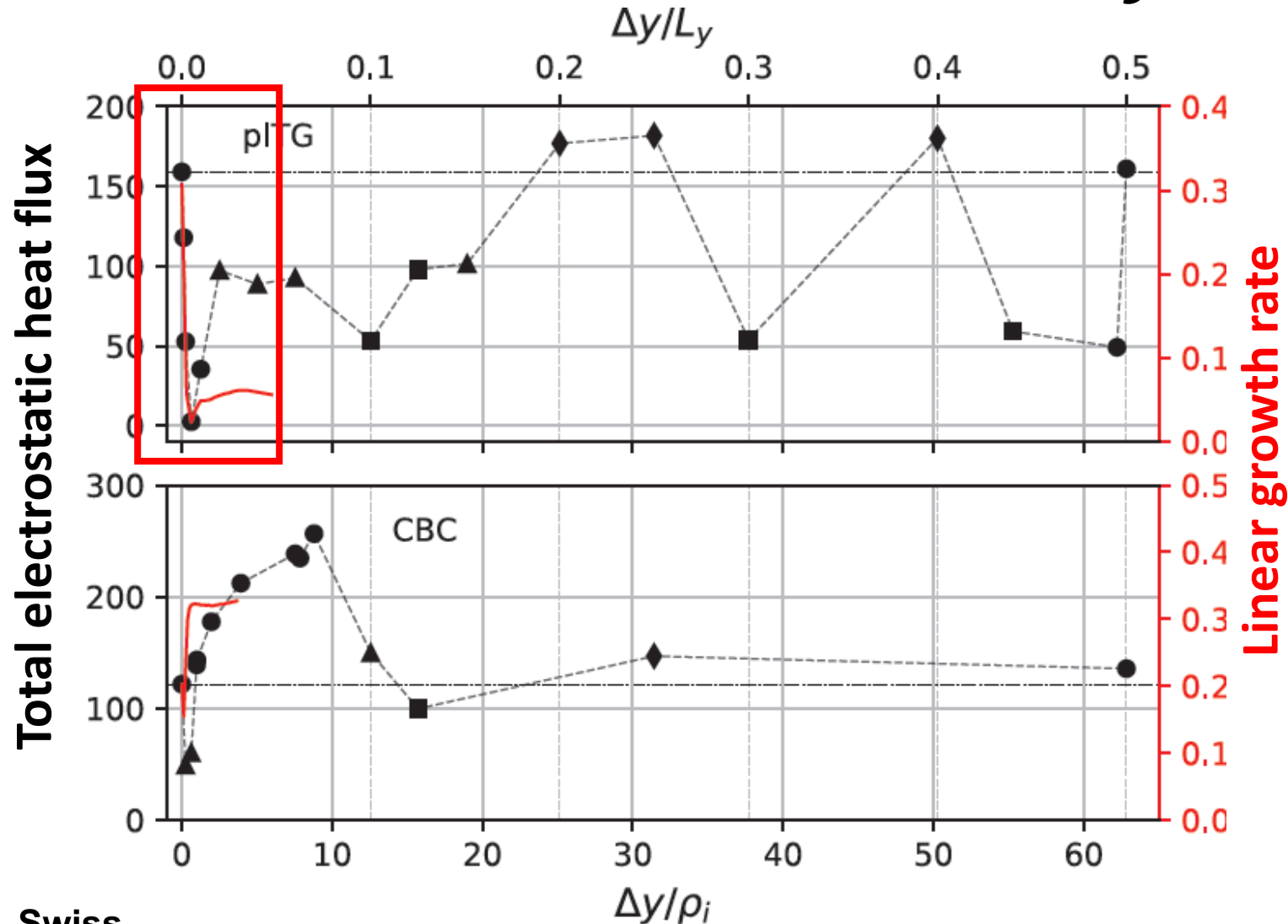
Summary: Eddy parallel length study

- Simulations with kinetic electrons at zero magnetic shear **require hundreds of poloidal turns** to achieve convergence
- **Kinetic electrons** set the parallel turbulence length scale
- Heat flux is reduced when eddies are unable to self-interact
- In simulations with electron temperature gradient long parallel waves emerge



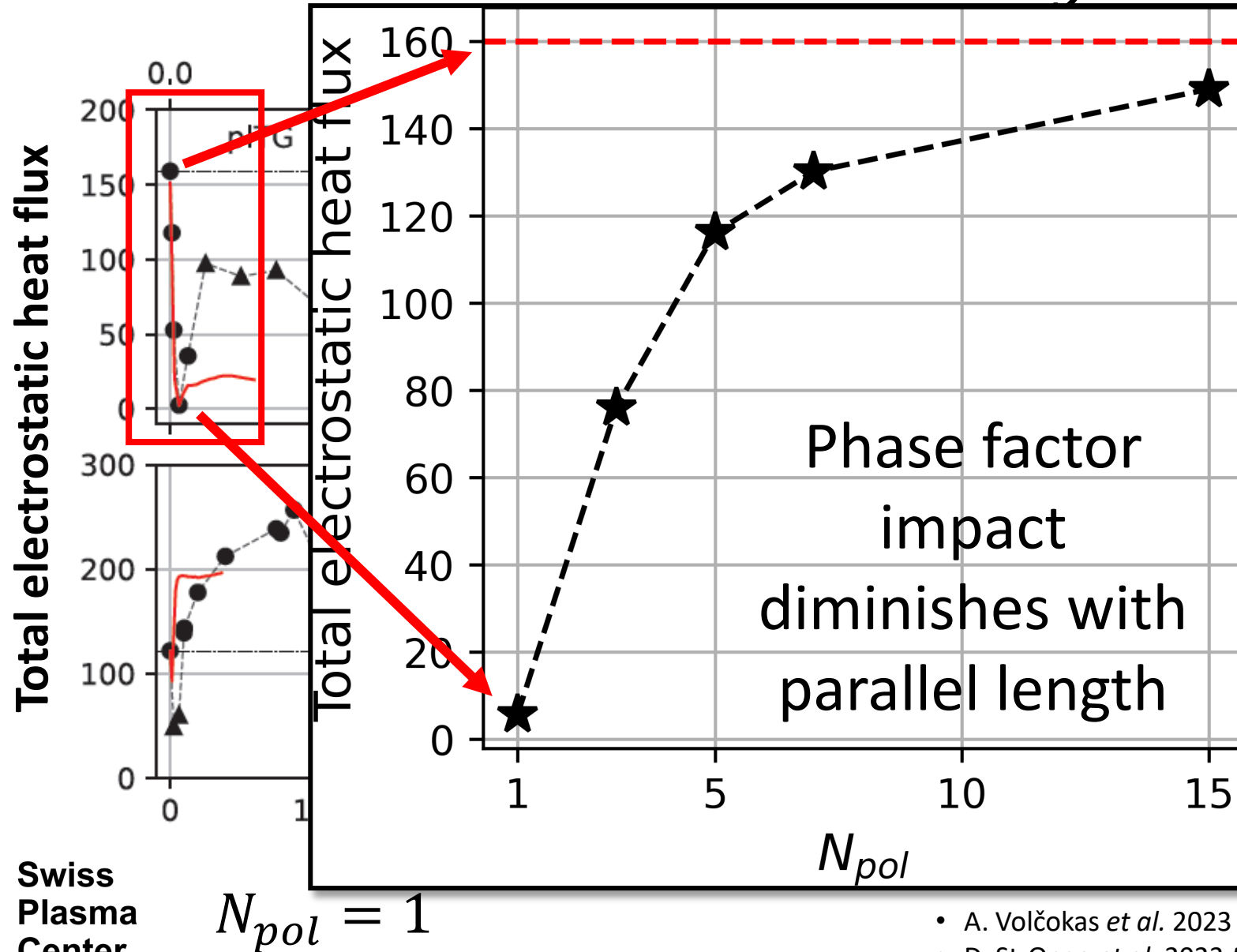
Binormal shift Δy when $\hat{s} = 0$ 

- Turbulence can be strongly affected by the magnetic field topology, leading to **complete stabilization** in some cases.

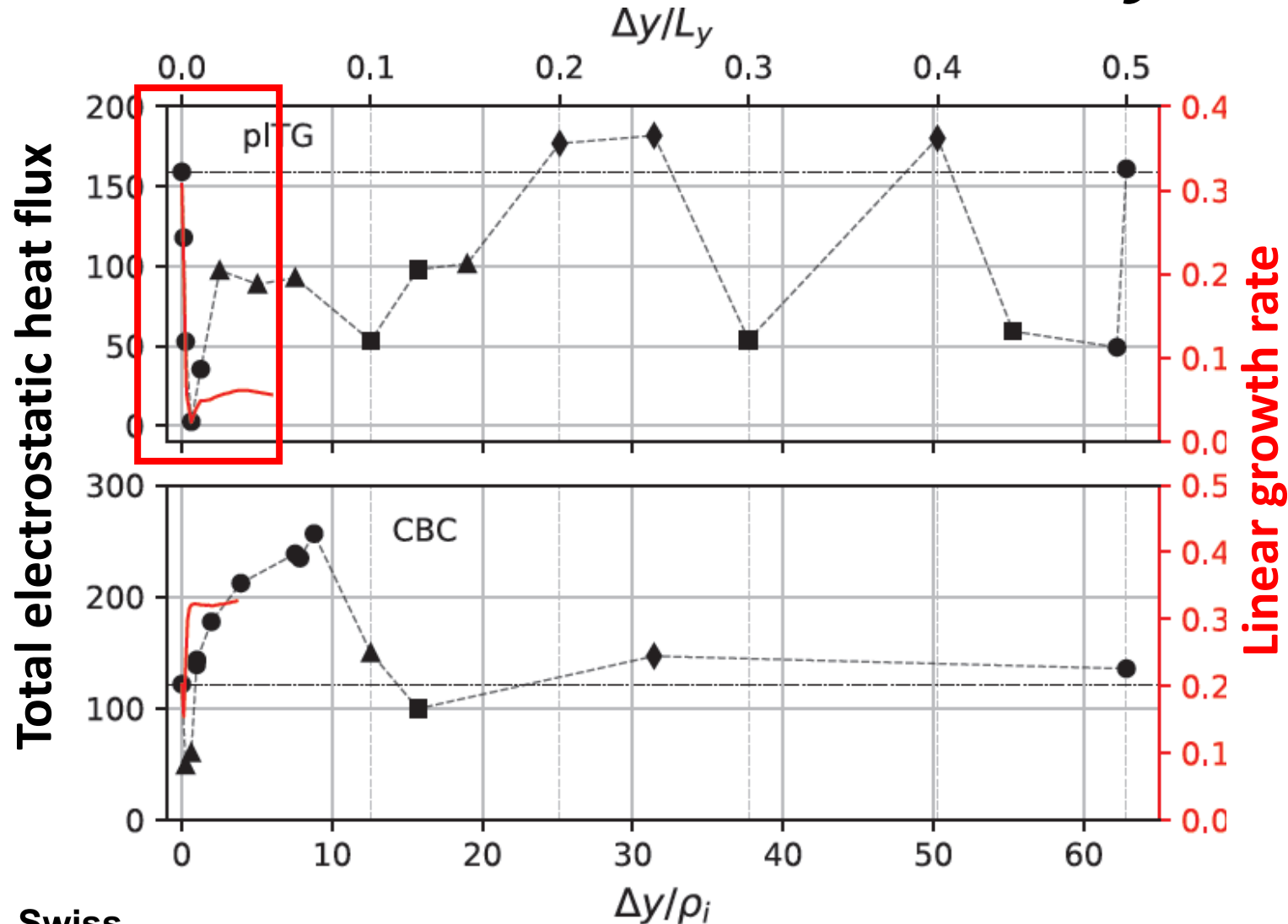
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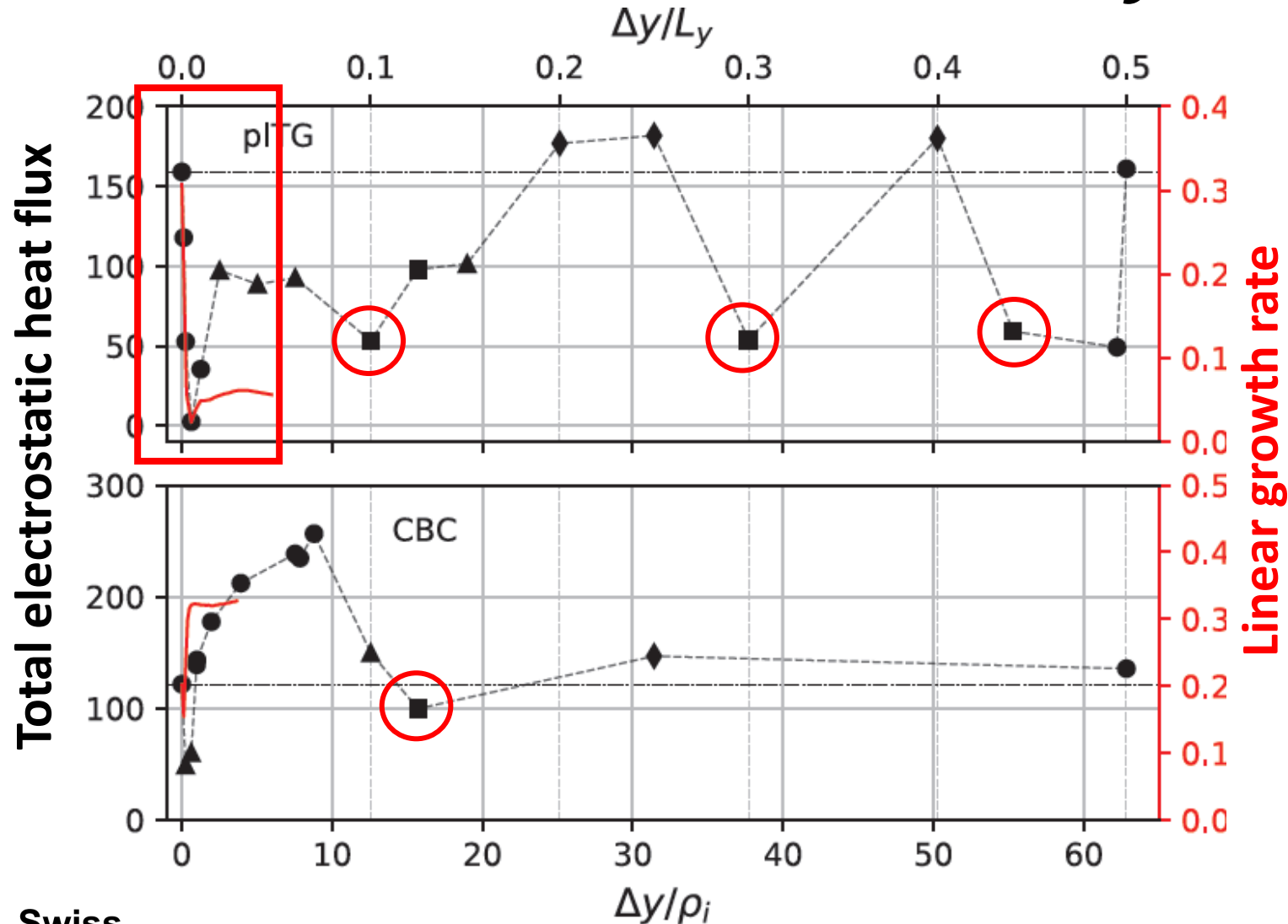
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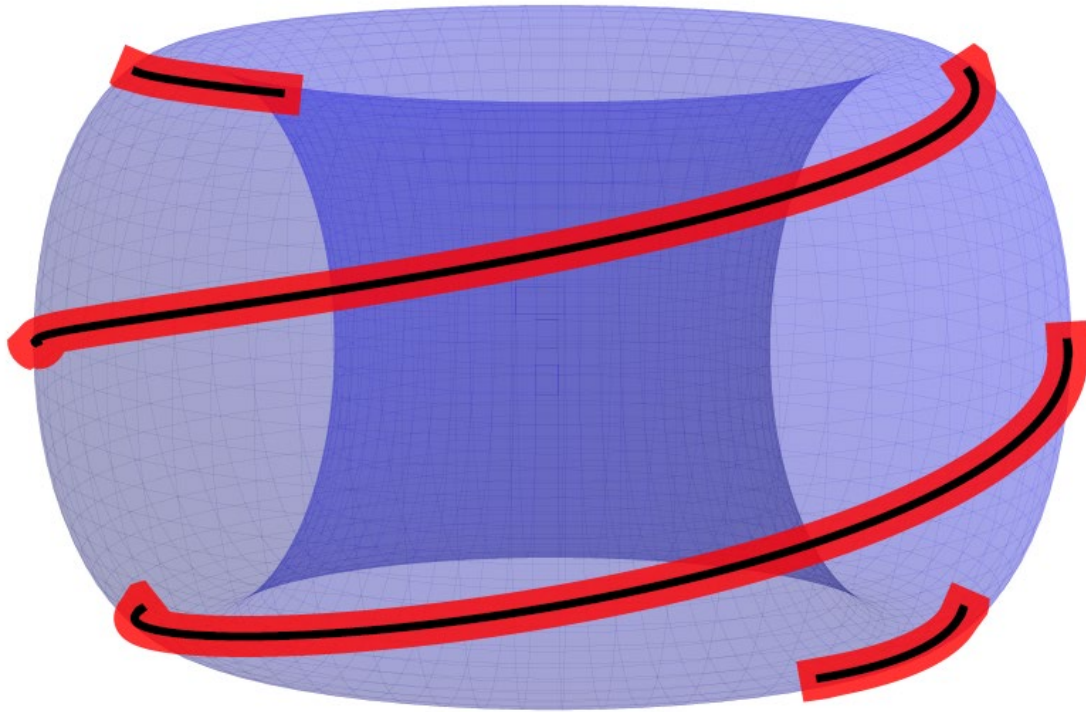
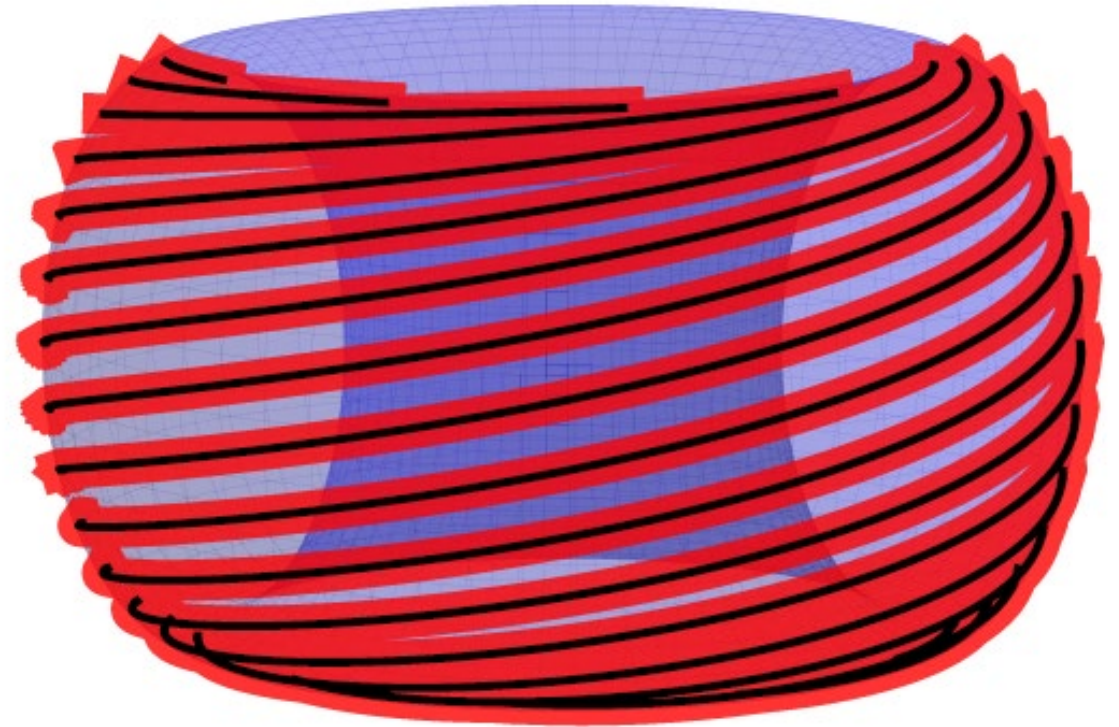
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Binormal shift Δy when $\hat{s} = 0$ 

- Turbulence can be strongly affected by the magnetic field topology, leading to **complete stabilization** in some cases.
- A single eddy can cover an entire flux surface leading to reduced flux

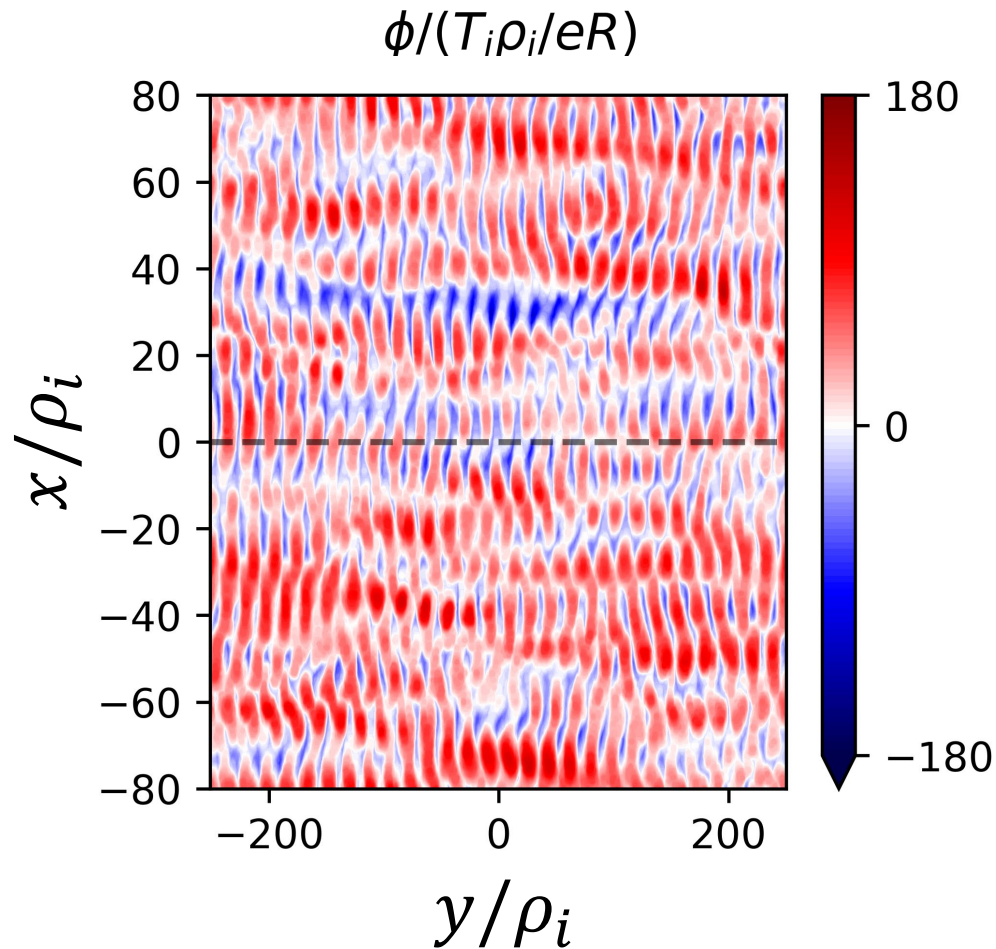
“Eddy squeezing”

Perpendicular self-interaction

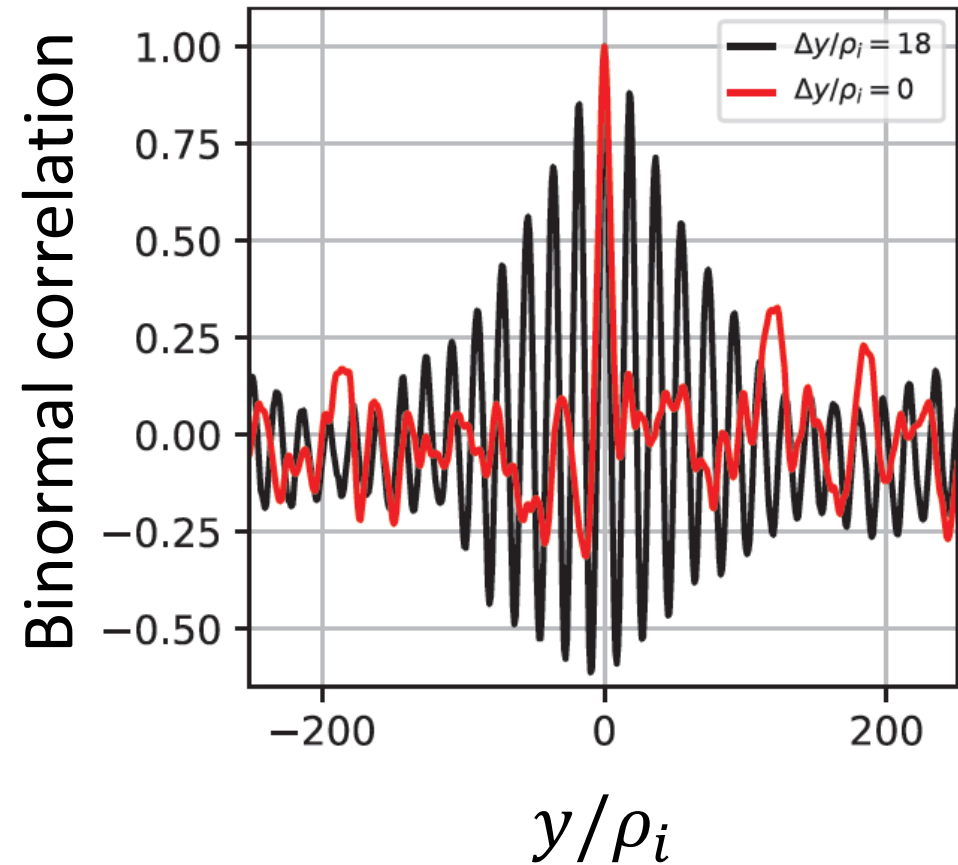
 $q=2.5$  $q=2.7$

$$\hat{s} = 0$$

- Increased electron mass
- $k_{y,min}\rho_i = 0.0125$
- **Two-fold change in the heat flux**



@ $t = 200 R/c_i$



- A. Volčokas *et al.* 2023 *Nucl. Fusion* **63** 014003

EPFL Summary: Binormal shift study $\hat{s} = 0$

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- Allows to study self-interaction at irrational q values and **in particular close to rational- q**
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- Provides experimentally testable predictions regarding turbulence self-organization in tokamak core

EPFL Outline

- Motivation and background
- Methods
- Numerical results:
 - Low magnetic shear simulations
 - Ultra-long turbulent eddies $\hat{s} = 0$
 - **Non-uniform magnetic shear simulations**
- Conclusions

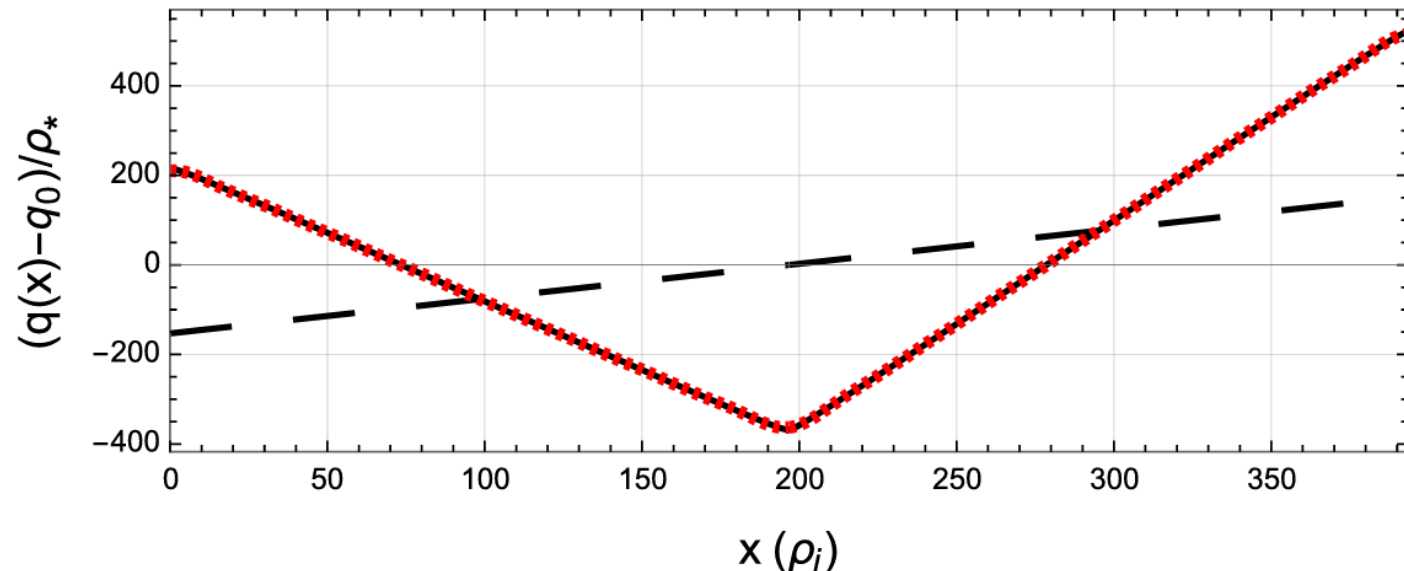


Key idea:

Create a safety factor profile that varies on the gyroradius-scale, which is rigorously derived from a current drive source inspired by ECCD

- The simulations are no longer “local” but still performed in a flux tube domain

$$\tilde{q}(r) = \sum_{n=1}^{\infty} \left[\tilde{q}_n^C \cos\left(\frac{2\pi n}{L_r}(r-r_0)\right) + \tilde{q}_n^S \sin\left(\frac{2\pi n}{L_r}(r-r_0)\right) \right].$$



EPFL Goals

- Increase realism
- Investigate electromagnetic effects
- Investigate extremely low magnetic shear
- Include safety factor curvature effects

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Simulate the minimum in the safety factor profile

EPFL Full safety factor profile

Safety factor profile



$$\frac{\bar{q} - q_0}{\rho_*} = \frac{r}{q_0} \hat{s} x$$



Linear shear
part

EPFL Full safety factor profile

Safety factor profile



$$\frac{\bar{q} - q_0}{\rho_*} = \frac{r}{q_0} \hat{s}x + \tilde{q}(x)$$

Non-uniform shear
part



Linear shear
part



EPFL Full safety factor profile

Safety factor profile



$$\frac{\bar{q} - q_0}{\rho_*} = \frac{r}{q_0} \hat{s}x + \tilde{q}(x) + \tilde{q}_{\delta A_{\parallel}}(x)$$

Non-uniform shear

part



$$\delta \mathbf{B} = \nabla \times \delta A$$



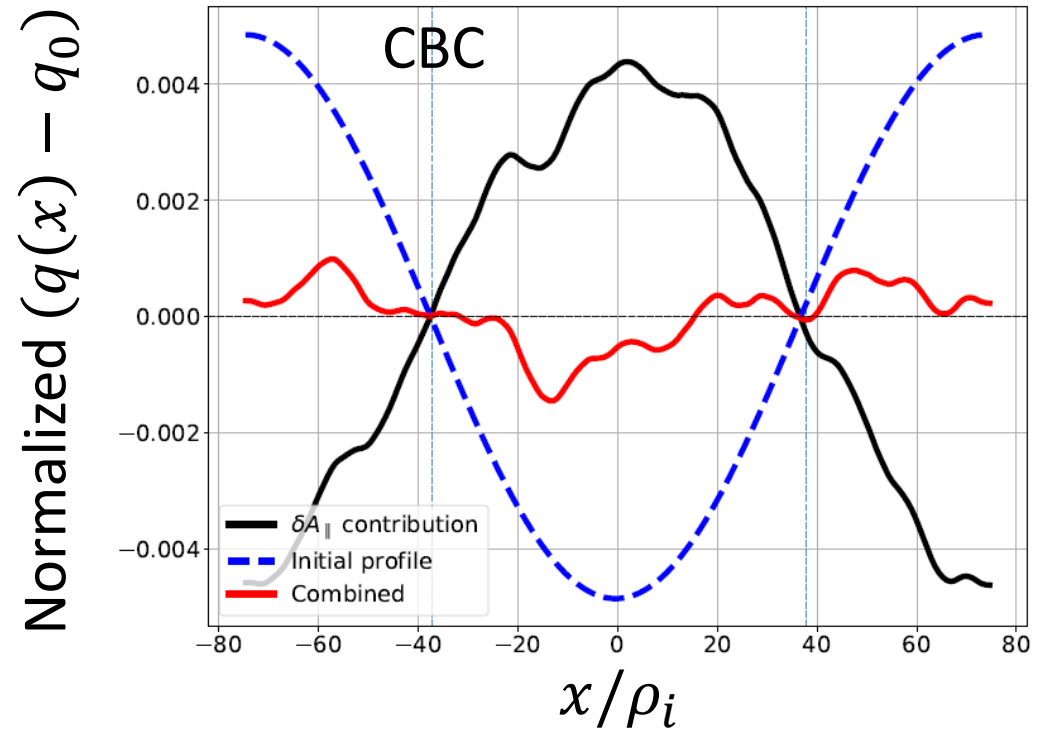
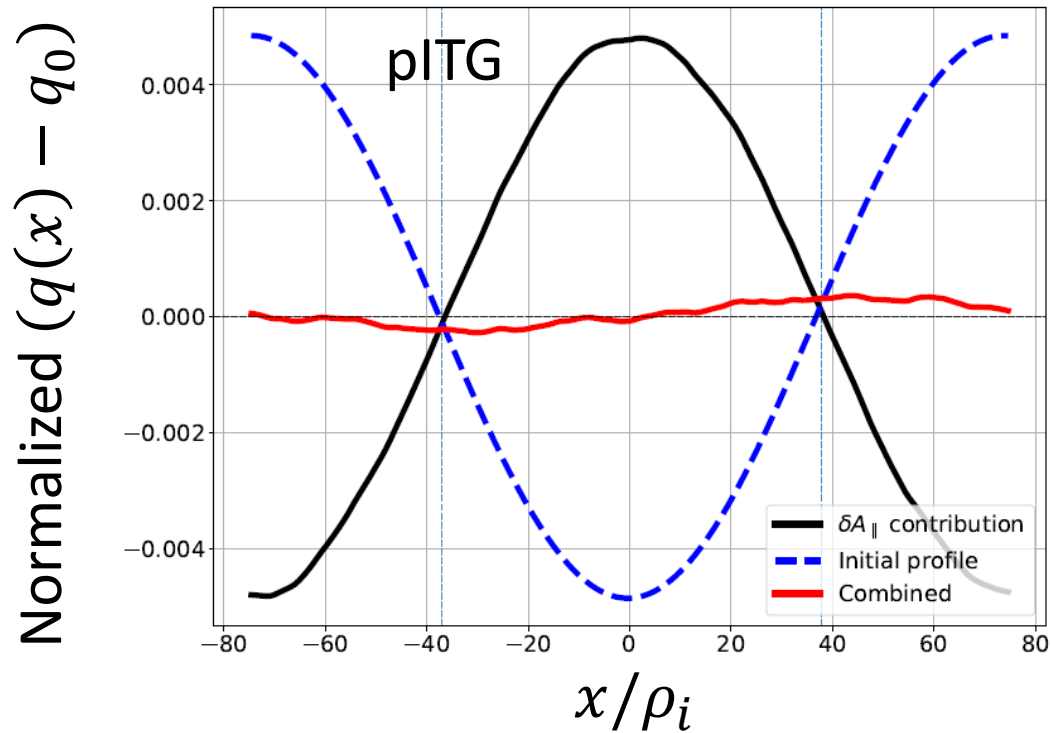
Linear shear
part



Correction due to
zonal part of the
magnetic potential

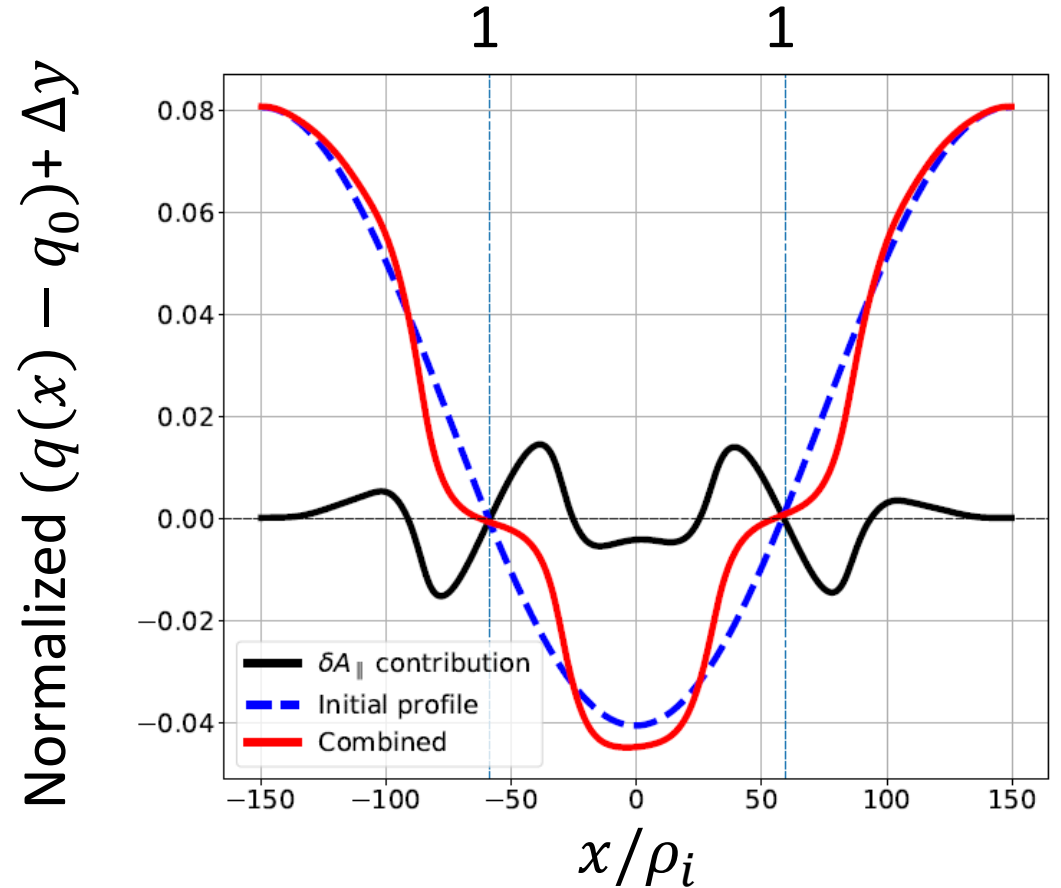
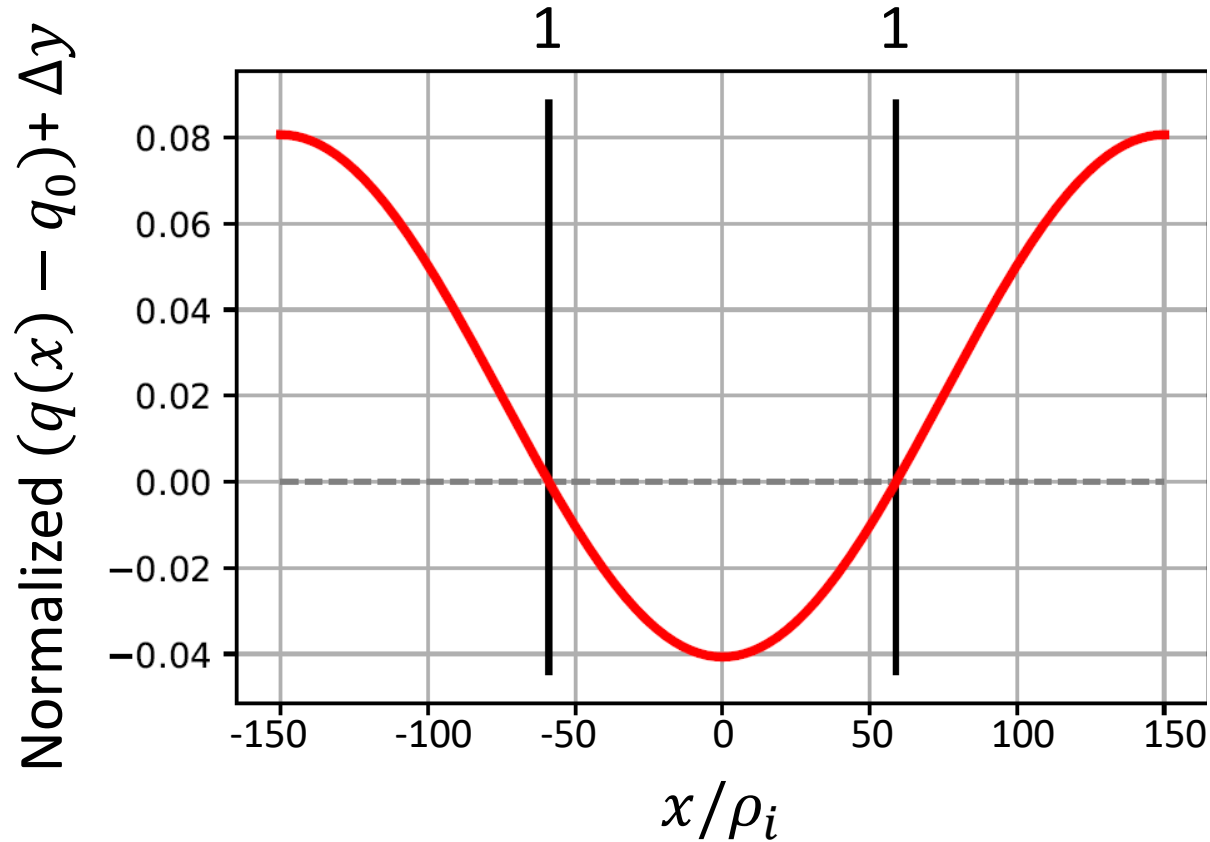


Full flattening of q profile



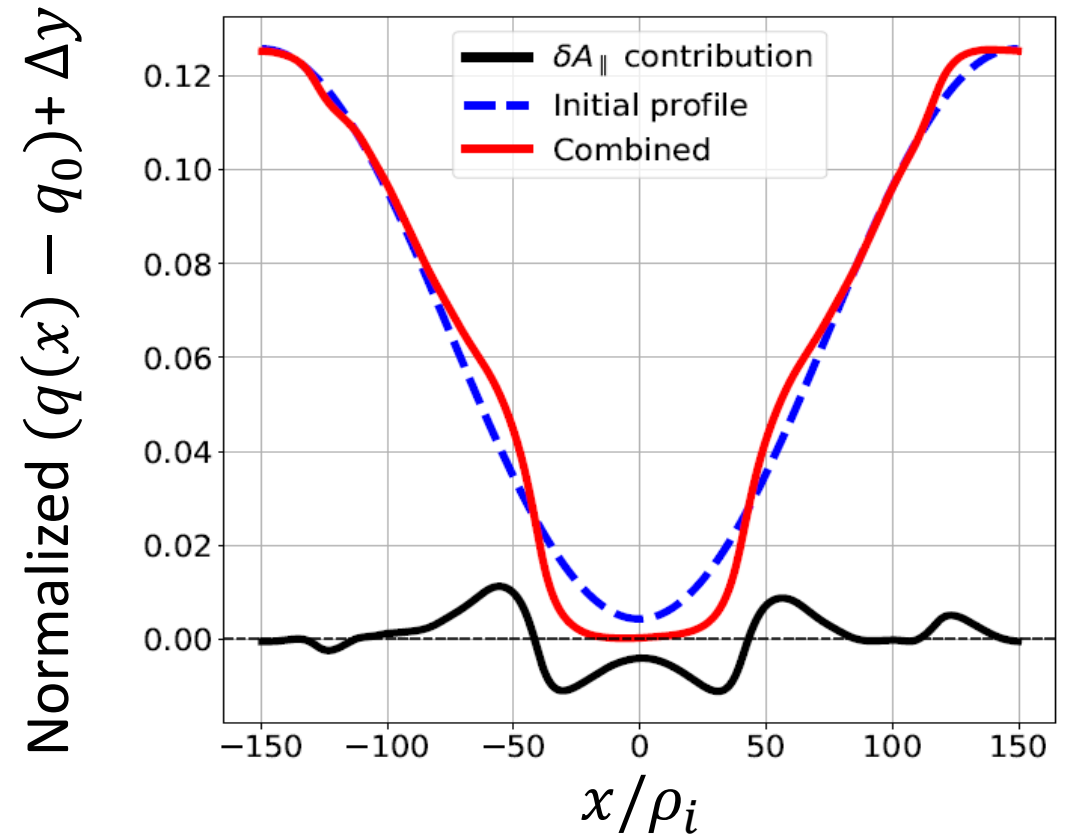
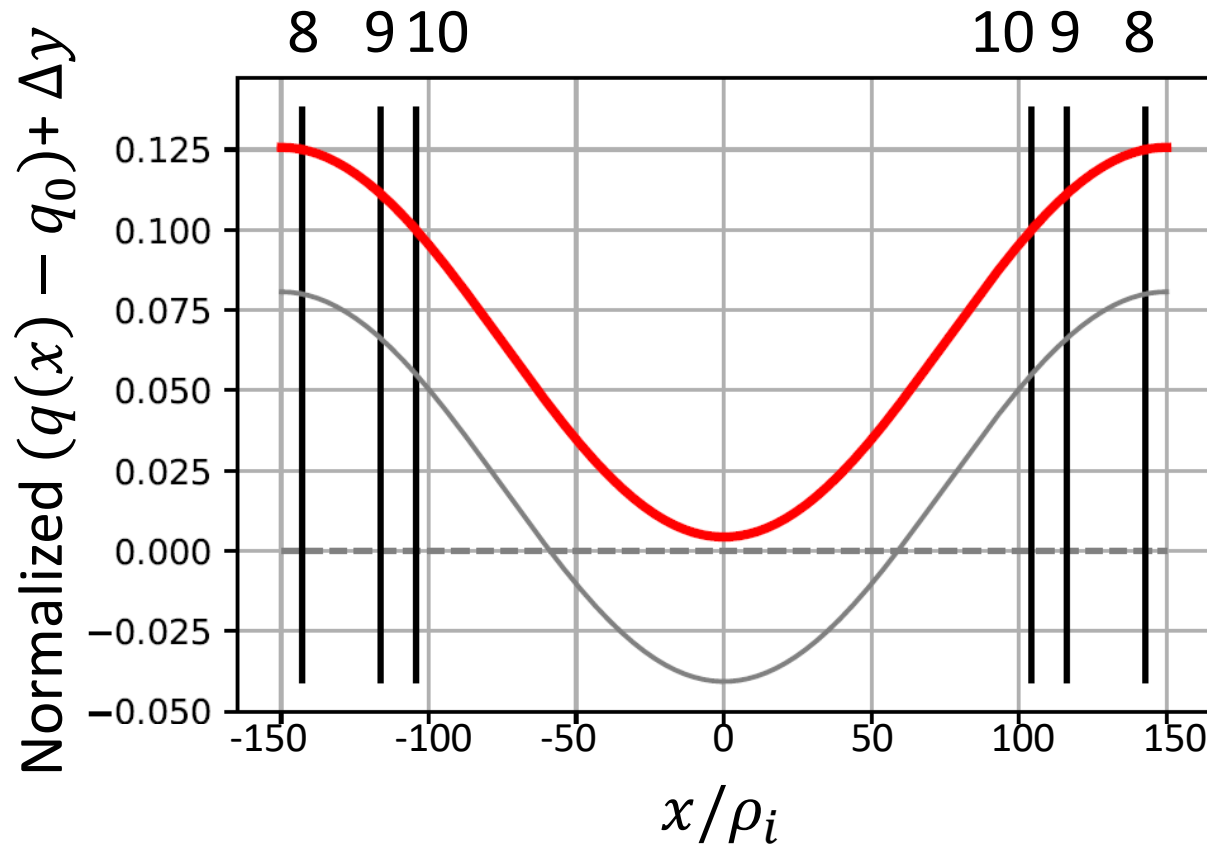
- $\beta = 0.001$
- $\hat{s}_S^1 = -0.004$

EPFL Two integer surfaces



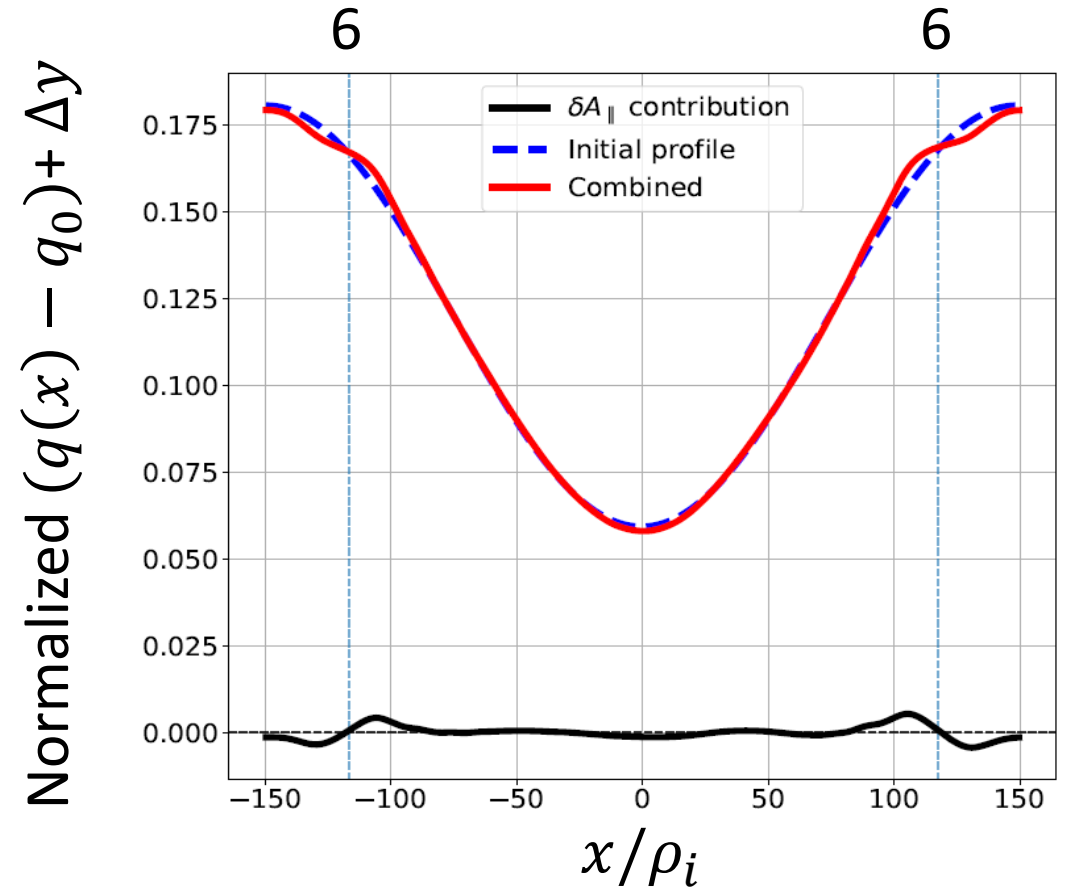
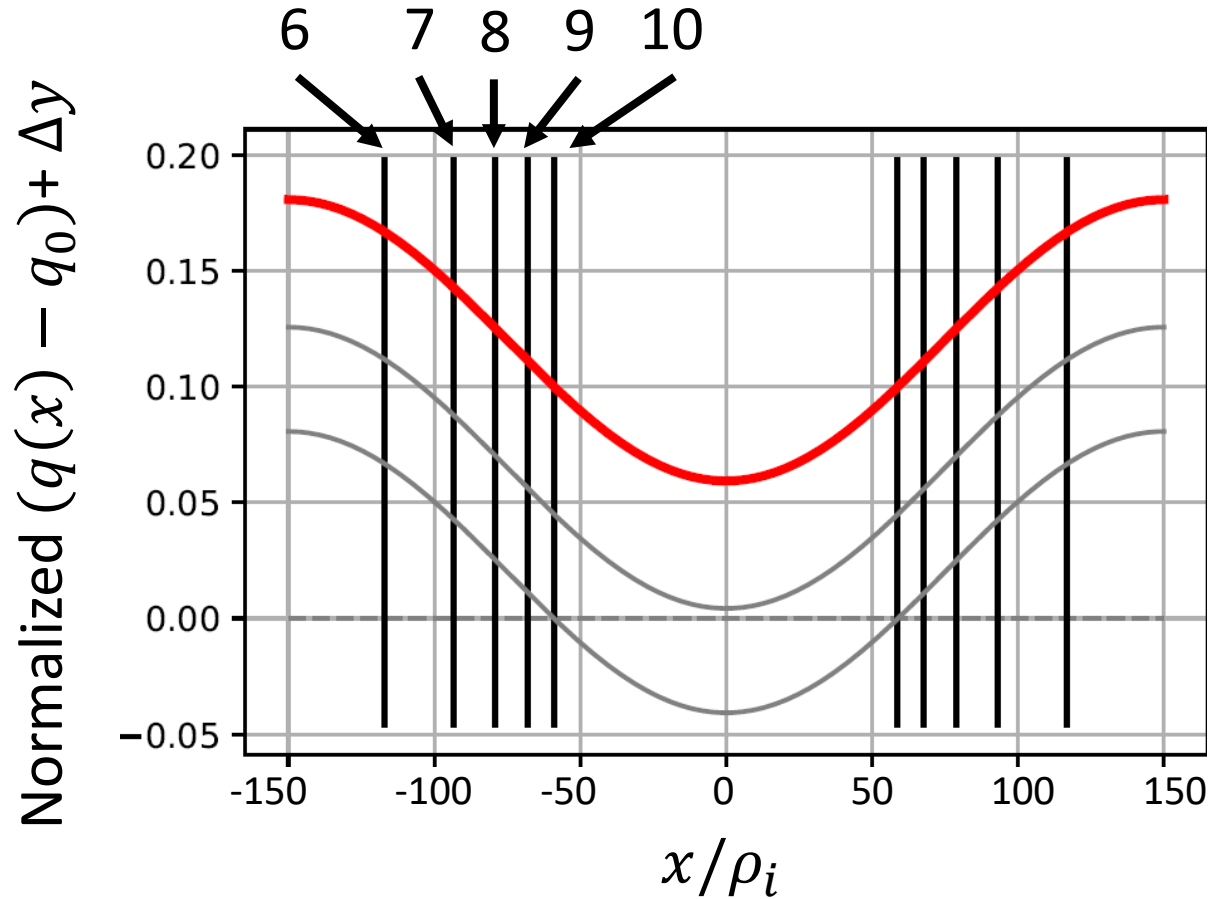
- $\beta = 0.0001$
- $\hat{s}_S^1 = -0.025$

Close to an integer surface



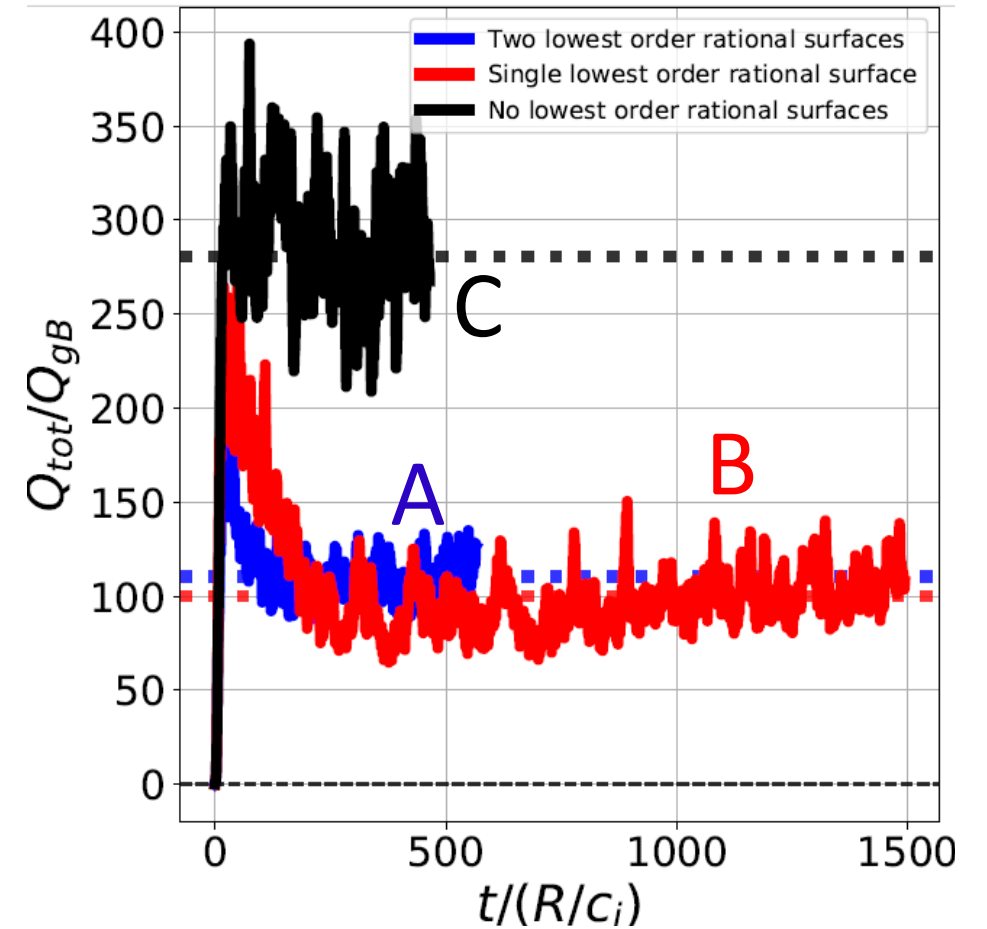
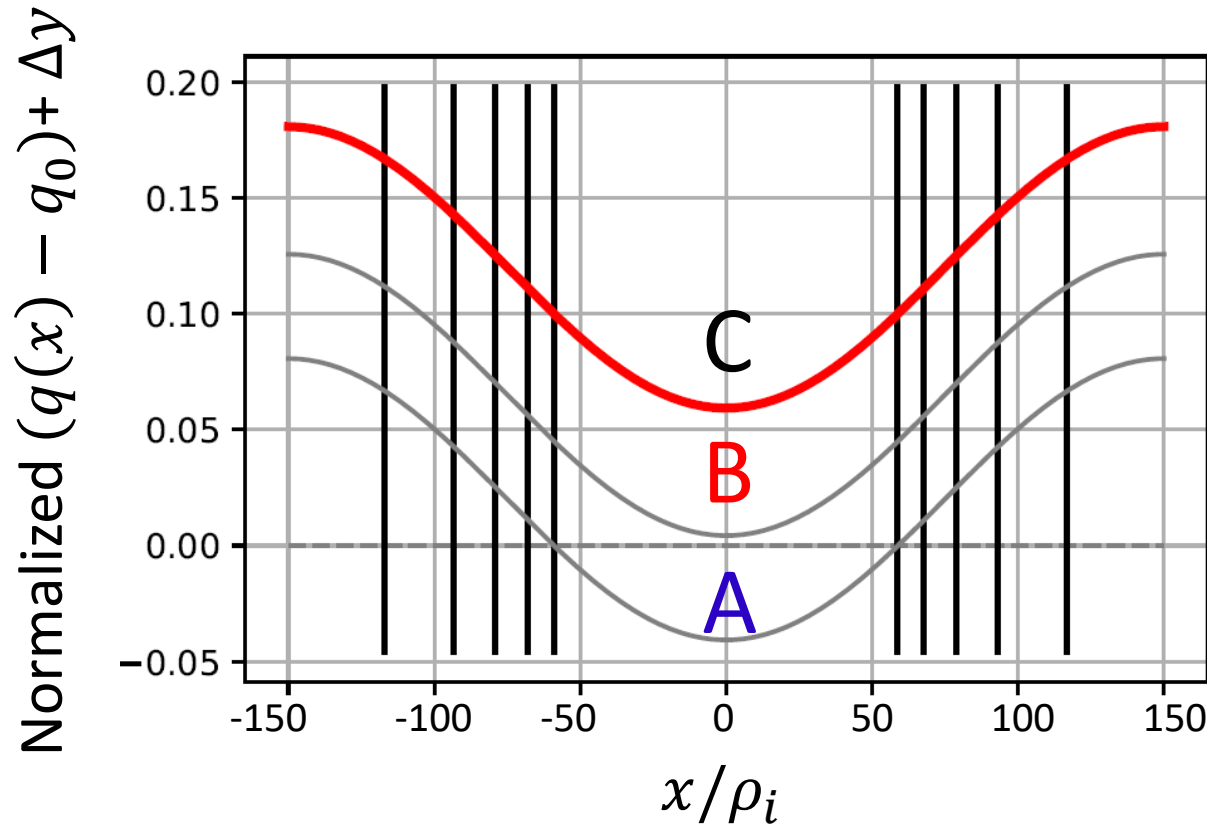
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Away from an integer surface



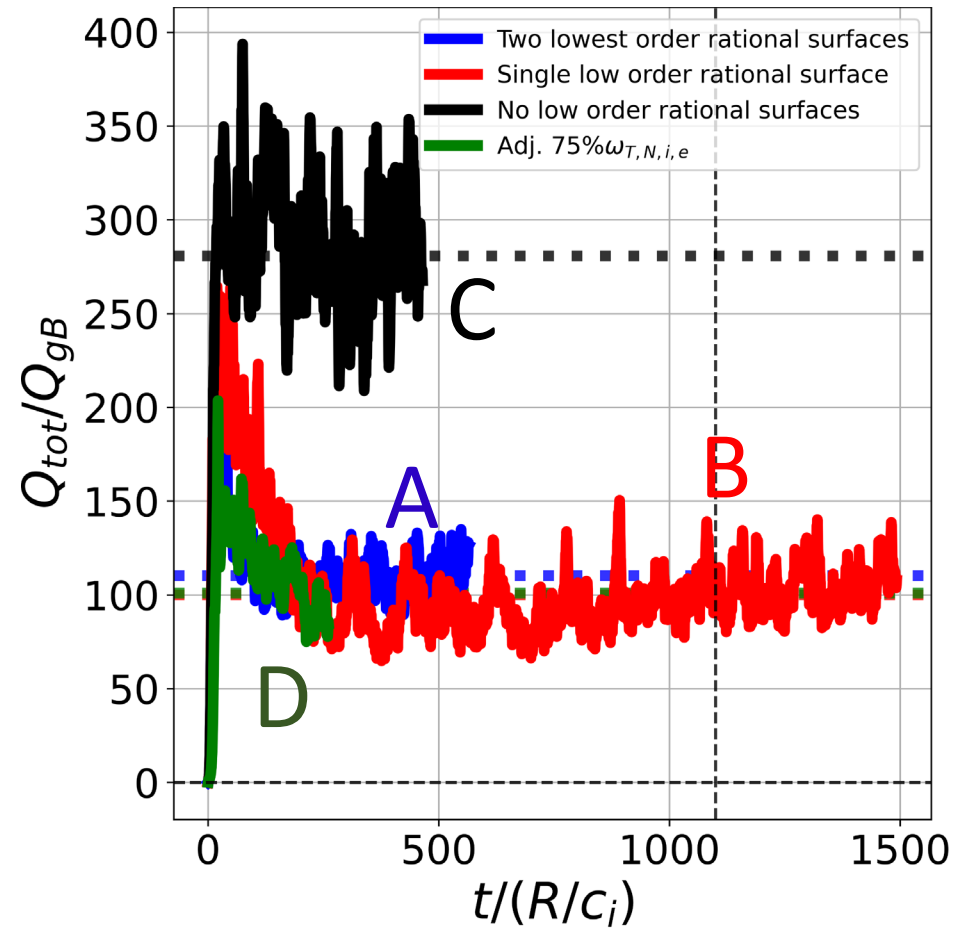
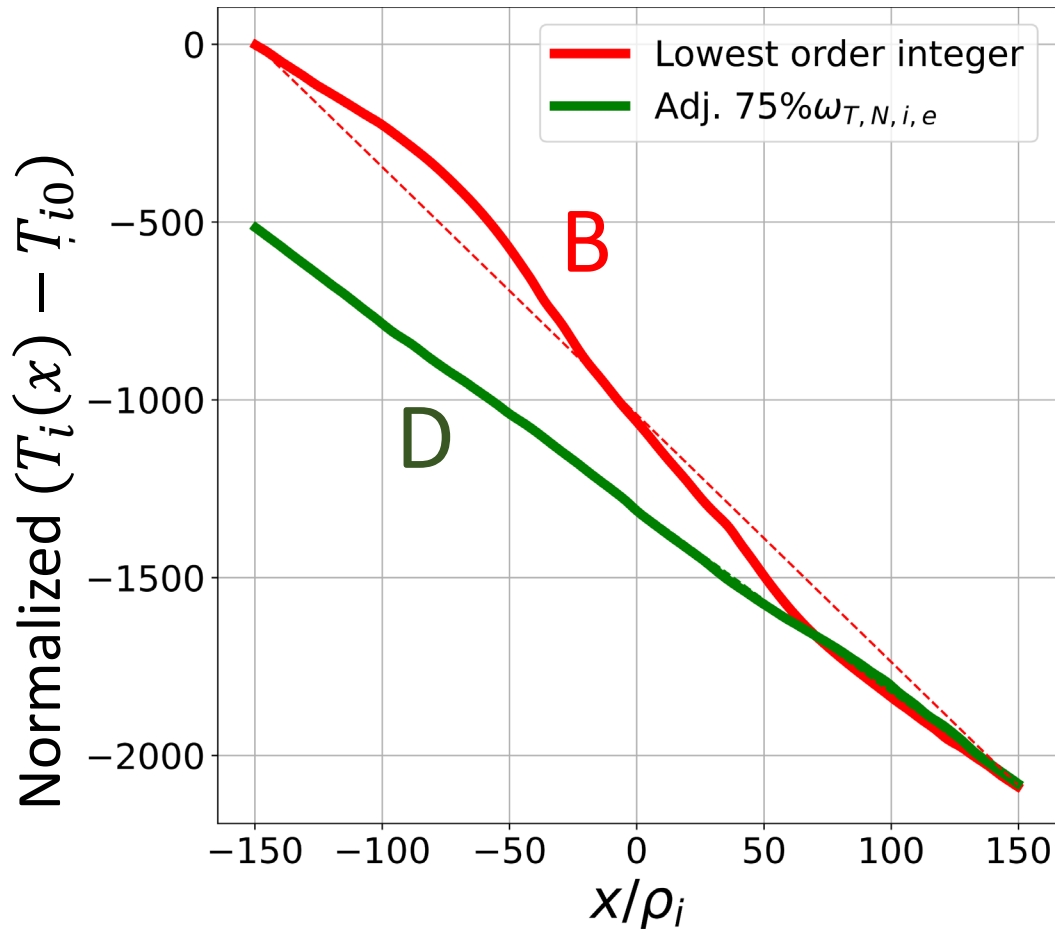
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q profile impact on heat flux



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Profile steepening



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Thank you for your attention!