





The role of plasma flow on quasi-helical states in reversed-field pinches

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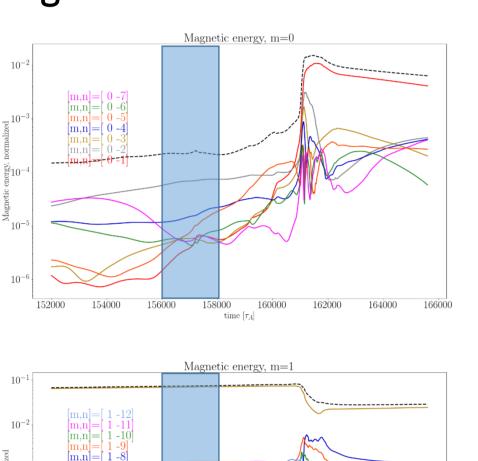
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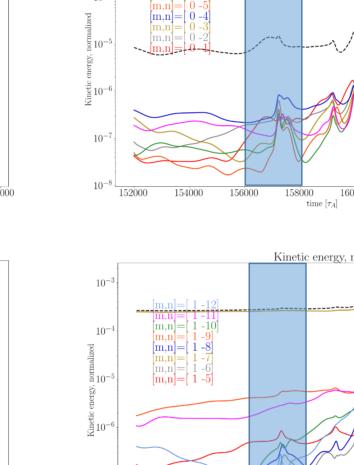
INTRODUCTION: Magnetic reconnection [1] involves changes in the connections of magnetic field lines, leading to modifications in the magnetic field topology. This phenomenon is common in both space and laboratory plasmas, such as solar flares [2], planetary magnetospheres, and laboratory experiments involving laser-plasma interactions. Magnetic reconnection is characterized by the release of magnetic energy, which is converted into kinetic and thermal energy within the plasma. It results in the acceleration of particles to non-thermal velocities [3] and the generation of waves and turbulence. Current sheets are often observed as a signature of reconnection.

Between the many configurations used to confine hot plasmas, both the RFP and the tokamaks exhibit a tendency to develop kink-like deformations, with RFPs forming a global helical configuration characterized by the quasi-single helicity (QSH) state [4, 5], interrupted by events which exhibit the signatures of magnetic reconnection [6]. These events have been referred to by various names in the RFP literature, including "dynamo relaxation events," "RFP sawtoothing," "discrete dynamo activity," and "localized reconnection events."

MAIN FEATURES of huge reconnection events in the reversed-field pinch

- release of magnetic energy, converted into kinetic and thermal internal energy of the plasma [1]
- presence of socalled current sheets [7, 8]
- acceleration of particles to nonthermal velocities
- generation of waves and turbulence [9]





AIM: answer the question: "what is the cause of the relaxation / reconnections events in the RFP?"

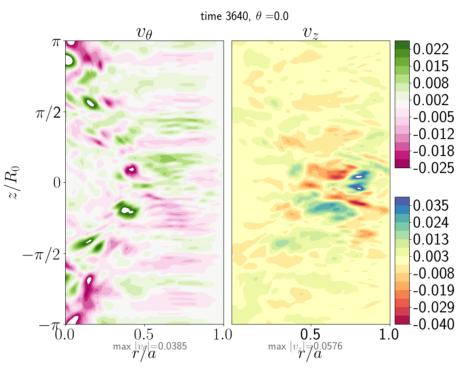
TOOLS: 3D nonlinear visco-resistive MHD modelling, and a wide set of diagnostics for the physical quantities like magnetic field, velocity field, current density field, electric field

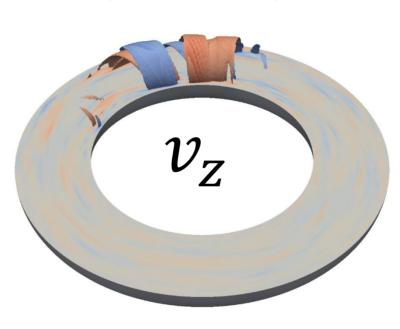
RESULTING INDICATION: The reconnection event is caused by the phenomenon of "collapse of the helixes", meaning that localized imperfections of the quasi-helical structure of the magnetic and velocity fields result in localized asymmetries in the plasma current density helical distribution. And given that co-currents tend to attract each others, a partial or total collapse of the structure is bound to happen. In this work we study both minor and major collapses of the RFP helical backbone. We focus on small/incomplete reconnection events, because the behaviour of the plasma is easier to observe than in the major events (which generate current sheets, high magnetic stochasticity).



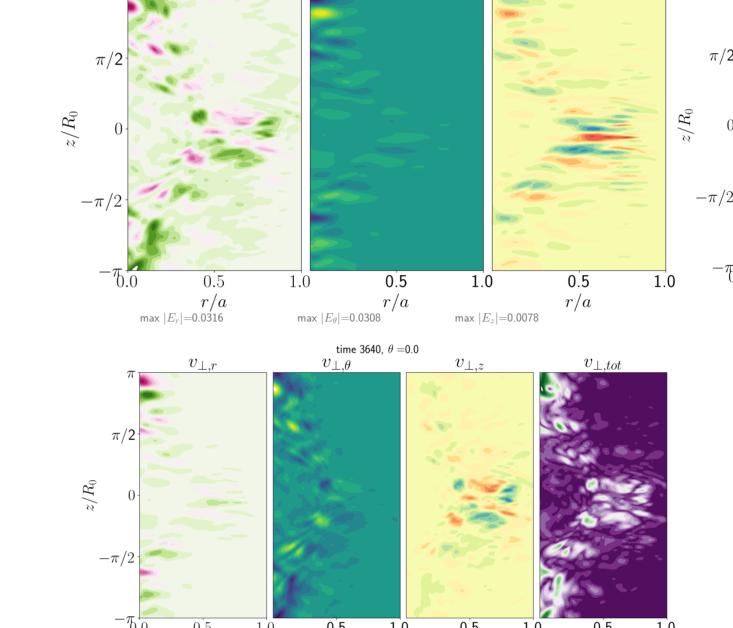


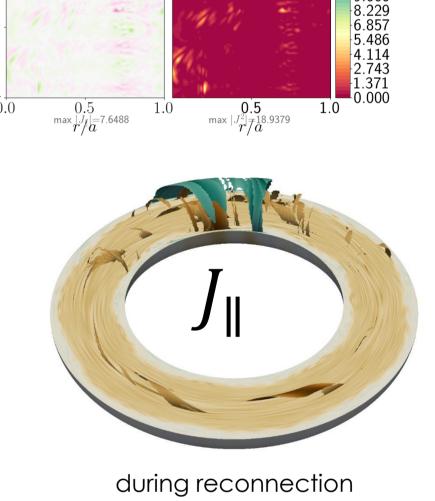






during reconnection





why study minor reconnection events

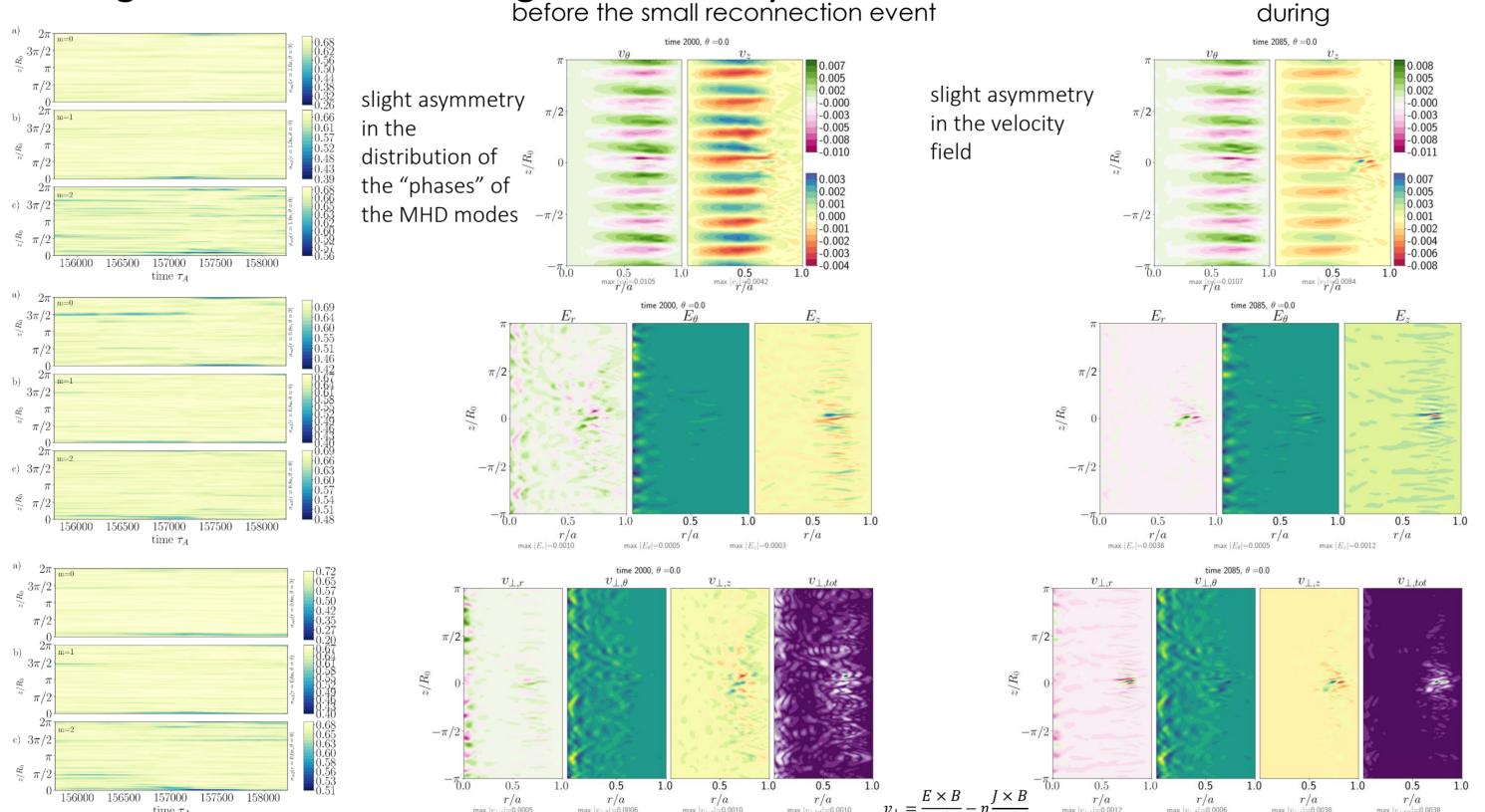
so-called "minor crash", i.e. a crash that does not destroy the helical state.

This allows avoiding the outburst of MHD modes and velocity associated with the main relaxation events, that obscures some of the features of the process.

In extreme synthesis:

- around the position of strong mode locking one can observe irregularities in the spatial distribution of the velocity field with respect to a quasi-helical structure.
- The locking provokes an outburst of velocity that aims at decreasing the locking itself
- generation of localized electric field, perpendicular velocity, intense plasma current spikes around the position of high locking.

from magnetic field mode locking to the velocity and electric field before the small reconnection event

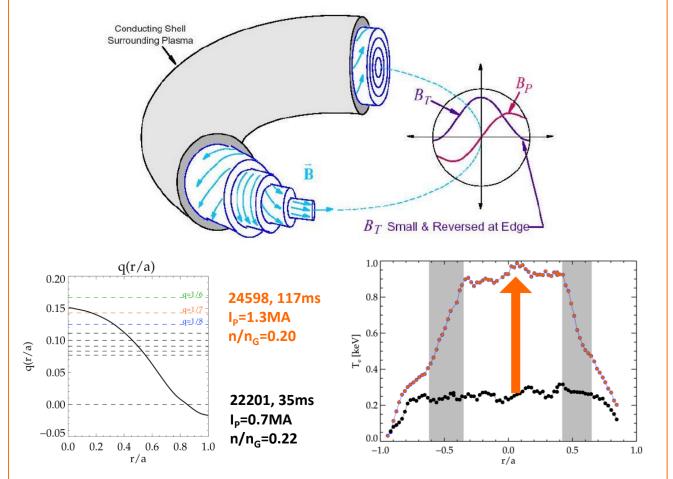


plasma current density features before the small reconnection event during

THE REVERSED-FIELD PINCH: similar to the tokamak but 10x stronger plasma current (with same B_{ϕ})

consequences:

- i) the current produces most of the confining magnetic field;
- toroidal field coils need to produce just the small reversed field -> simplicity
- iii) high current can heat the plasma



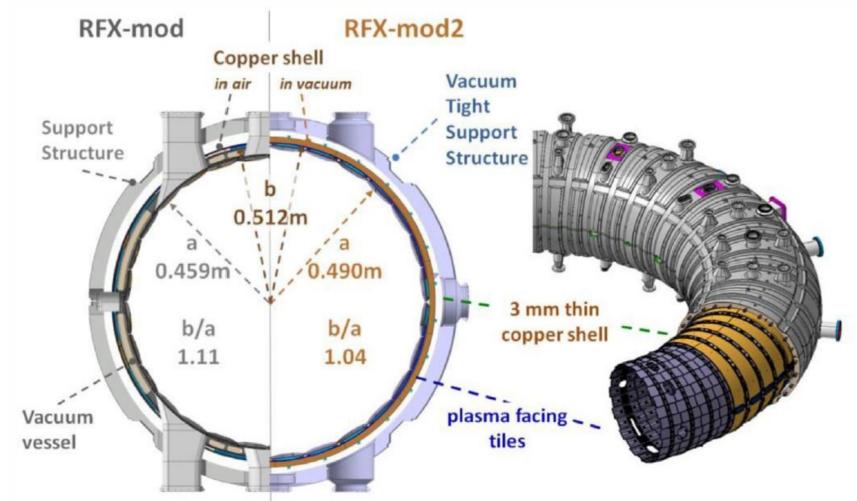
References:

[1] M. Yamada, R. Kulsrud, H. Ji Rev Mod Physics 82, 603 (2010) [2] D. Innes, B. Inhester et al, Nature 368, 811 (1997) [3] L. Chen, A. Bhattacharjee et al, Nature Physics 4, 19 (2008)

[4] R. Lorenzini et al, Nature Physics 5, 570 (2009) [5] D.F. Escande, P. Martin et al, PRL 85, 1662 (2000)

[6] B. Momo et al, Nucl. Fusion 60 056023 (2020)

RFX-MOD2 NEW MAGNETIC LAYOUT



The highly resistive vessel will be removed, graphite tiles will be attached to the copper stabilizing shell and the stainless steel support structure will be modified in order to be vacuum tight. In RFX-mod2, the shell-plasma proximity passes from b/a = 1.11to b/a = 1.04 and copper, instead of Inconel, will be the continuous conducting structure nearest to the plasma [10].

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- [8] M. Veranda et al, Rendiconti Lincei. Scienze Fisiche e Naturali 31 963 (2020) [9] A. Lazarian, G. Eyink et al Philos. Trans. R. Soc. A Math. Phys. Eng. Sci. 373, 2041, 20140144 (2015)
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MHD DESCRIPTION OF RECONNECTION EVENTS

3D nonlinear viscoresistive MHD modelling gives a description of mode locking similar to the one coming from the experimental measurements previously described.

Peculiarity: the phase of growth of the helical state is often disturbed by the growth of secondary perturbations not having the necessary strength to destroy the helical state. One can observe the quasi-periodical emergence of a QSH state, interrupted by violent eruption observed as a sharp increase of the kinetic energy.

 $\partial_t oldsymbol{v} + oldsymbol{v} \cdot
abla oldsymbol{v} = oldsymbol{J} imes oldsymbol{B} + rac{1}{M}
abla^2 oldsymbol{v} + rac{1}{O} oldsymbol{S}_M$ $\partial_t oldsymbol{B} =
abla imes (oldsymbol{v} imes oldsymbol{B} - rac{1}{S} oldsymbol{J})$ $abla imes oldsymbol{B} = oldsymbol{J}$ $\nabla \cdot \boldsymbol{B} = 0$ $\sigma_m = 1/N_{tot} \sum_{j=1}^{n_{max}-1} \sum_{k=i+1}^{n_{max}} \left| \sin(\frac{\phi_{m,j} - \phi_{m,k}}{2}) \right|$

- solved by SpeCyl [10] Physical approximations: $\rho = cost, p = 0;$
- some definitions: cylindrical geometry: periodical boundary conditions on θ and z;

Dissipation coefficients radial profile: $\eta(r,t) = \eta_0(1+20r^{10}); \ \nu(r,t) = \nu_0$

- resistivity: $\eta_0 = \frac{\tau_A}{\tau_B} = S^{-1}$ (inverse Lundquist number). viscosity $v_0 = \frac{\tau_A}{T} = M^{-1}$ (inverse viscous Lundquist number).
- $M=10^4$ $n_{MP}=7$ ideal wall

 - S_M represents an ad-hoc momentum source • Edge MP is imposed at r=a with amplitude given by the quantity

 - Edge MP is a single helical modulation of the radial component of the field, typically with m=1.