

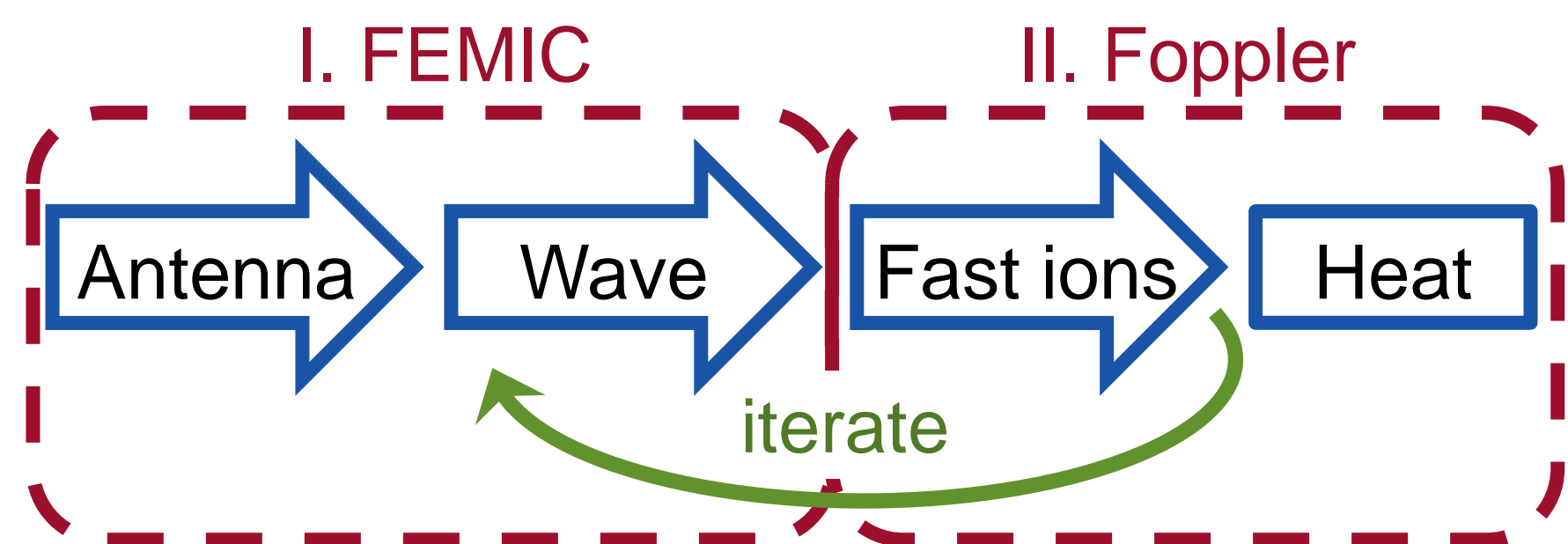
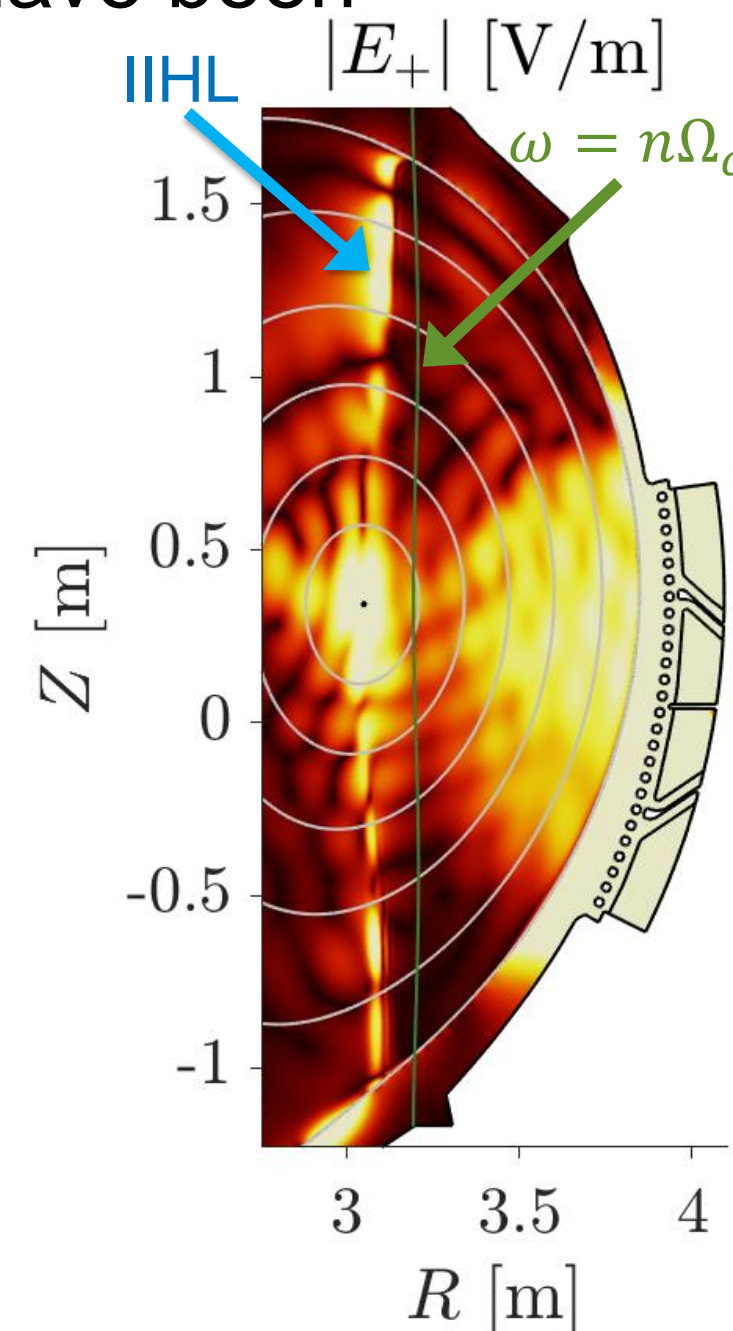
Pitch Angle Averaged Quasi-Linear Operator for ICRH

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Introduction

- Rapid ICRH modelling is needed for integrated modelling, scenario development and analysis, machine design etc.
- 1D Fokker-Planck models have been successful [1,2]
- Doppler physics is missing in 1D model:
 - Important for interaction with transport and MHD
 - Doppler shifted absorption at the ion-ion hybrid layer (IIHL) enhances the absorption, see Fig.
- Foppler**, new 1D-FP code including Doppler physics
- Combines the Stix-solution [3], with rabbit ears from Dendy [4].
- Uses waves from FEMIC [5]. Self-consistent coupling with FEMIC is under development.



Our goals are:

- Present 1D Fokker-Planck model Foppler
- Ad-hoc pitch angle distribution, from Dendy model, enables pitch angle averaging.
- Clarify impact of Doppler shift on quasilinear diffusion and distribution fcn.
- Quantify importance of Doppler shifts in JET D-T plasmas

Fokker-Planck model

- Fokker-Planck (FP) eq. predicts generation of fast ions and heating of thermal plasma
- Quasilinear time scales can be described by bounce averaged FP, for $f(v, \lambda, p_\phi, \sigma)$, where $\lambda = B_0 \mu / E$

$$C(f) + Q(f) = 0$$
 - $C(f)$ is the Coulomb collision operator
 - $Q(f)$ is the quasilinear diffusion operator
- Simplifications required for rapid simulations
 - Thin orbit, $p_\phi \rightarrow \psi$
 - Ansatz $f = f_s(v) f_D(v, \lambda)$, $\langle f_D \rangle_\lambda = 1$
 - Solve pitch angle averaged FP equation for f_s
 - Use ad-hoc model for f_D

References

- [1] L.-G. Eriksson et al., Nucl. Fusion 33, 1037 (1993)
- [2] D. Van Eester, Journal of Plasma Physics 87, 855870202 (2021)
- [3] T. H. Stix, Waves in Plasmas (New York: American Institute of Physics, 1992)
- [4] R. O. Dendy et al., Phys. Plasmas 2, 1623 (1995)
- [5] P. Vallejos et al., Nuclear Fusion 59, 076022 (2019)
- [6] L.-G. Eriksson and P. Helander, Physics of plasmas 1, 308–314 (1994)
- [7] M.J. Mantinen et al., to be submitted to Nuclear Fusion

Pitch angle averaged FP

- Pitch angle average, $\langle X \rangle_\lambda = I(X)/I(1)$, where

$$I(X) = \int_\psi^{\psi+d\psi} d^3r \int_v^{v+dv} d^3v X \propto \int \tau_b X d\lambda$$
 - $\langle C \rangle$ includes energy diffusion, β , and slowing down, κ

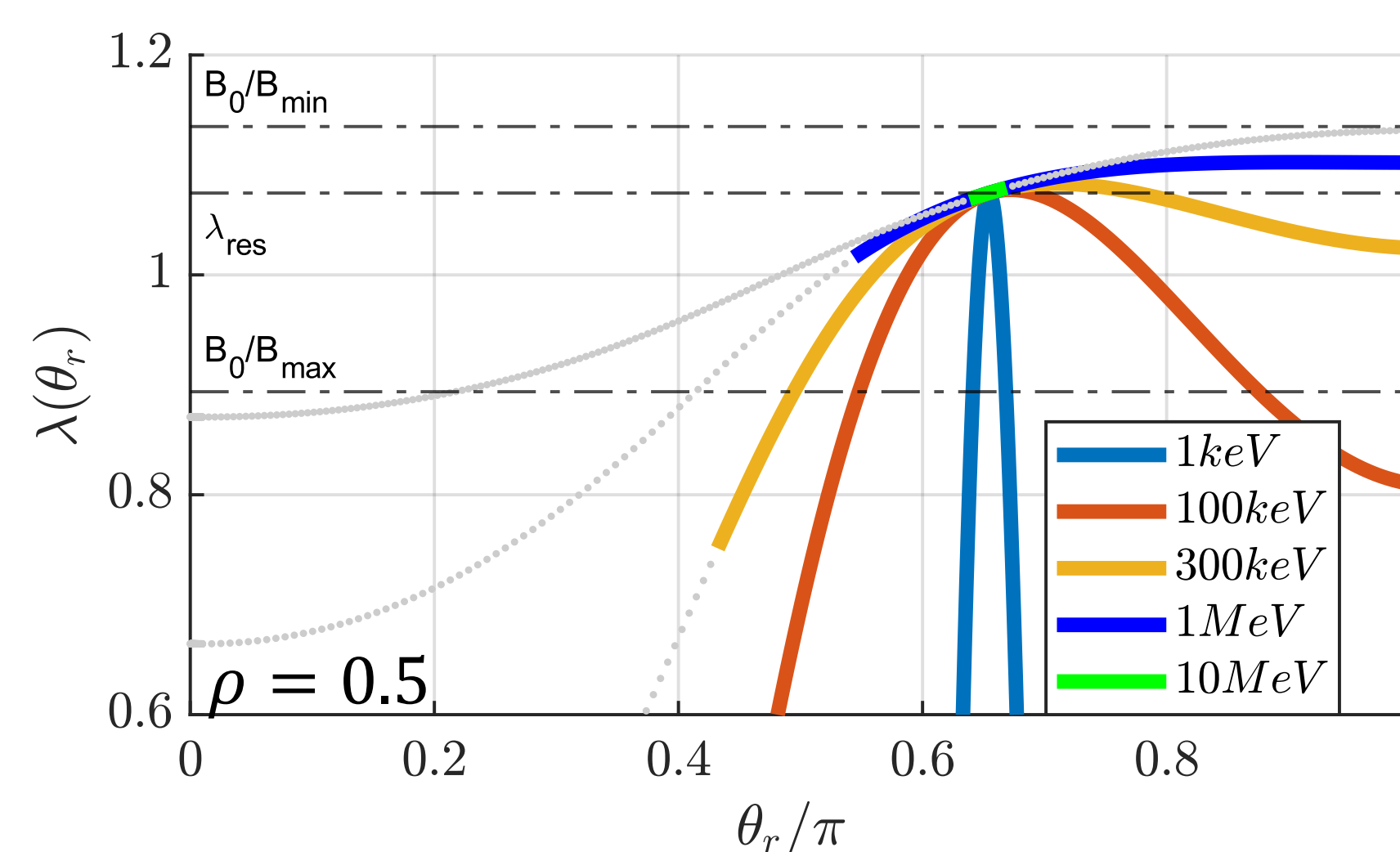
$$\langle C \rangle = \frac{1}{v^2} \frac{\partial}{\partial v} v^2 \left(-\kappa f_s + \frac{\beta}{2} \frac{\partial f_s}{\partial v} \right)$$
 - Quasilinear operator

$$\langle Q \rangle = \frac{1}{v^2} \frac{\partial}{\partial v} v^2 \langle f_D D \rangle \frac{\partial f_s}{\partial v}$$
- Here $D(v, \lambda)$ is bounce averaged diffusion coefficient [6]
- Integral $\int \tau_b f_D D d\lambda$ has singularity at tangential resonance that is removed when integrating in poloidal angle θ_r , where

$$\omega - k_\parallel(\theta_r) v_\parallel(\theta_r) - n \Omega_c(\theta_r) = 0$$

$$\langle f_D D \rangle = \sum_{n, \lambda_\phi, \omega} \frac{Z^2 e^2 L}{8 v m^2 \ell} \int_0^{2\pi} f_D \lambda \frac{|E_+ J_{n-1} + E_- J_{n+1}|^2}{k_\parallel} d\theta_r$$

where $\ell = v \int \tau_b d\lambda$ and L is circumference of flux surface.



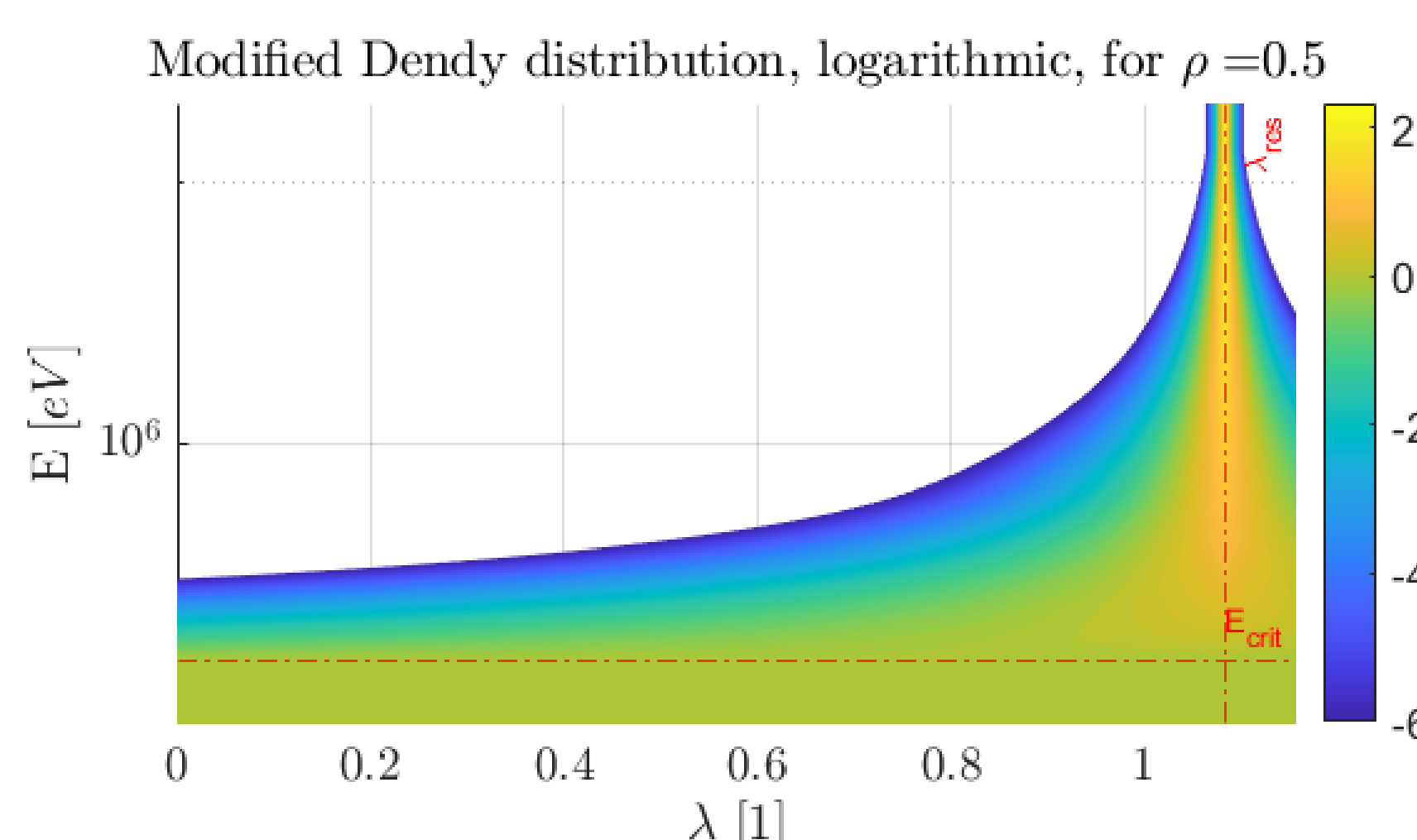
Pitch angle distribution

- The solution to the 1D Fokker-Planck equation as proposed by Stix [3]:

$$f_s(v) = f(0) \exp \left[- \int_0^v \frac{2\kappa}{\beta + \langle D \rangle} dv' \right]$$
- We propose a modified Dendy model [4]

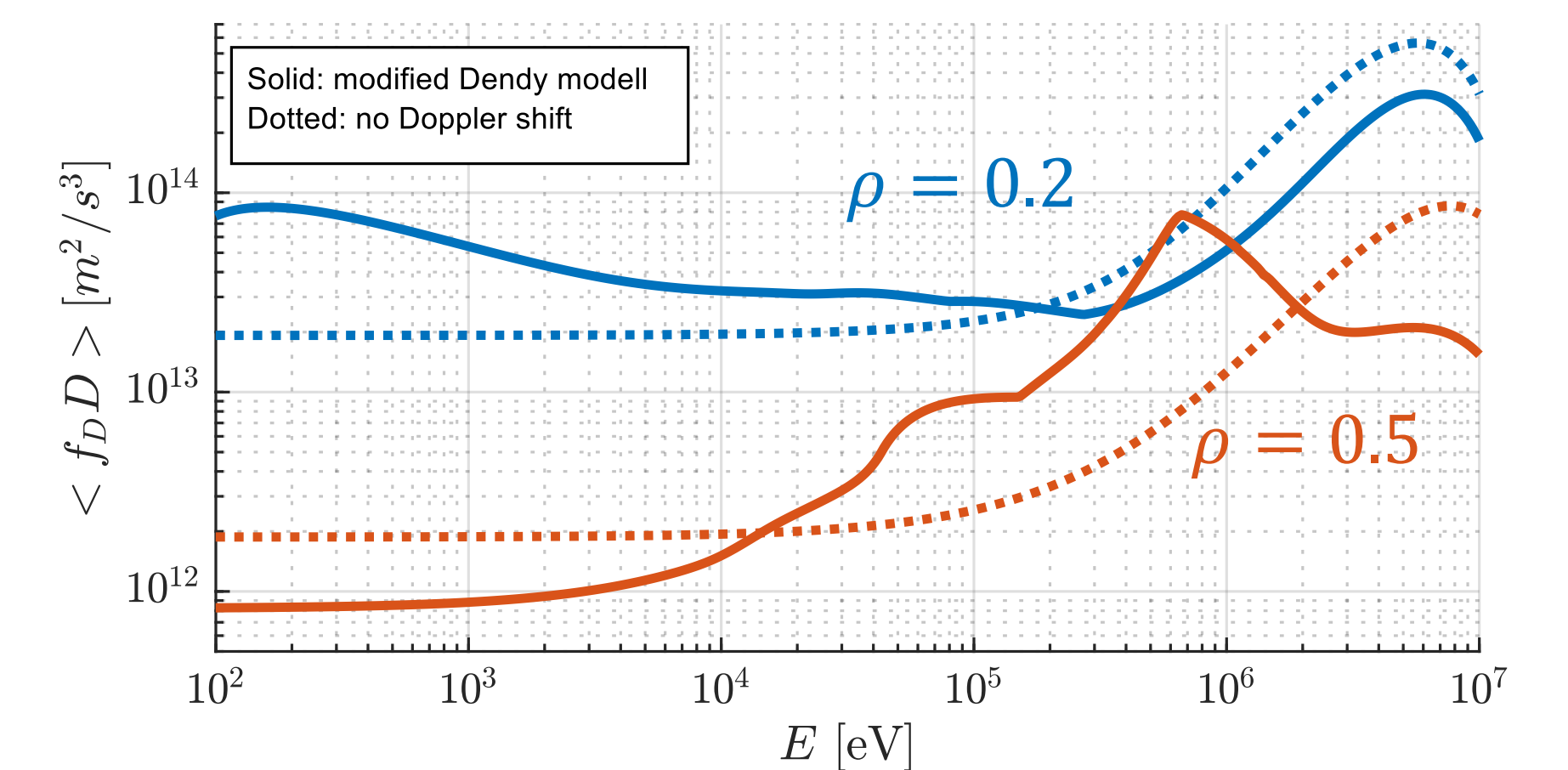
$$f_D(v, \lambda) \propto \exp \left[- \frac{E_\perp}{T_\perp} - \frac{|E_\parallel|}{T_\parallel} - \frac{E_\parallel - |E_\parallel|}{T_\perp} \right]$$

where $E_\perp = B_r \mu$ and $E_\parallel = E - E_\perp$ and $B_r = m\omega / (nZe)$, $T_\perp = -[\partial_E \ln(f_s)]^{-1}$, $T_\parallel = \max(T_\perp, E_\gamma)$, and E_γ is the critical energy.



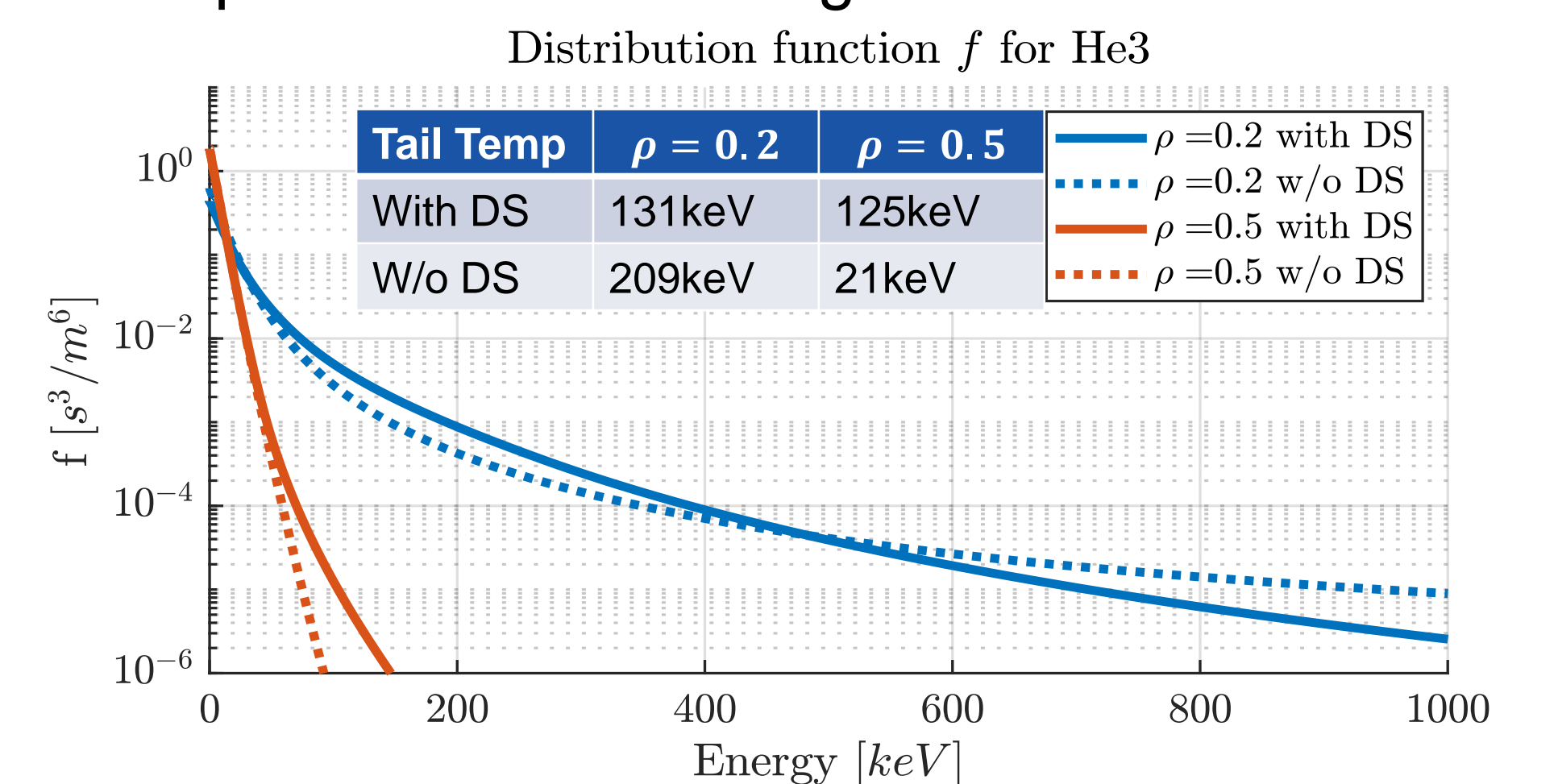
Diffusion coefficient $\langle f_D D(v) \rangle$

- Example, JET D-T plasmas with fast ³He (similar to #99629) [7]
- Doppler shift changes absorption by spreading out resonances to regions with different polarization and phase integral
- Pitch distribution narrows Doppler shifted resonance region for high energies



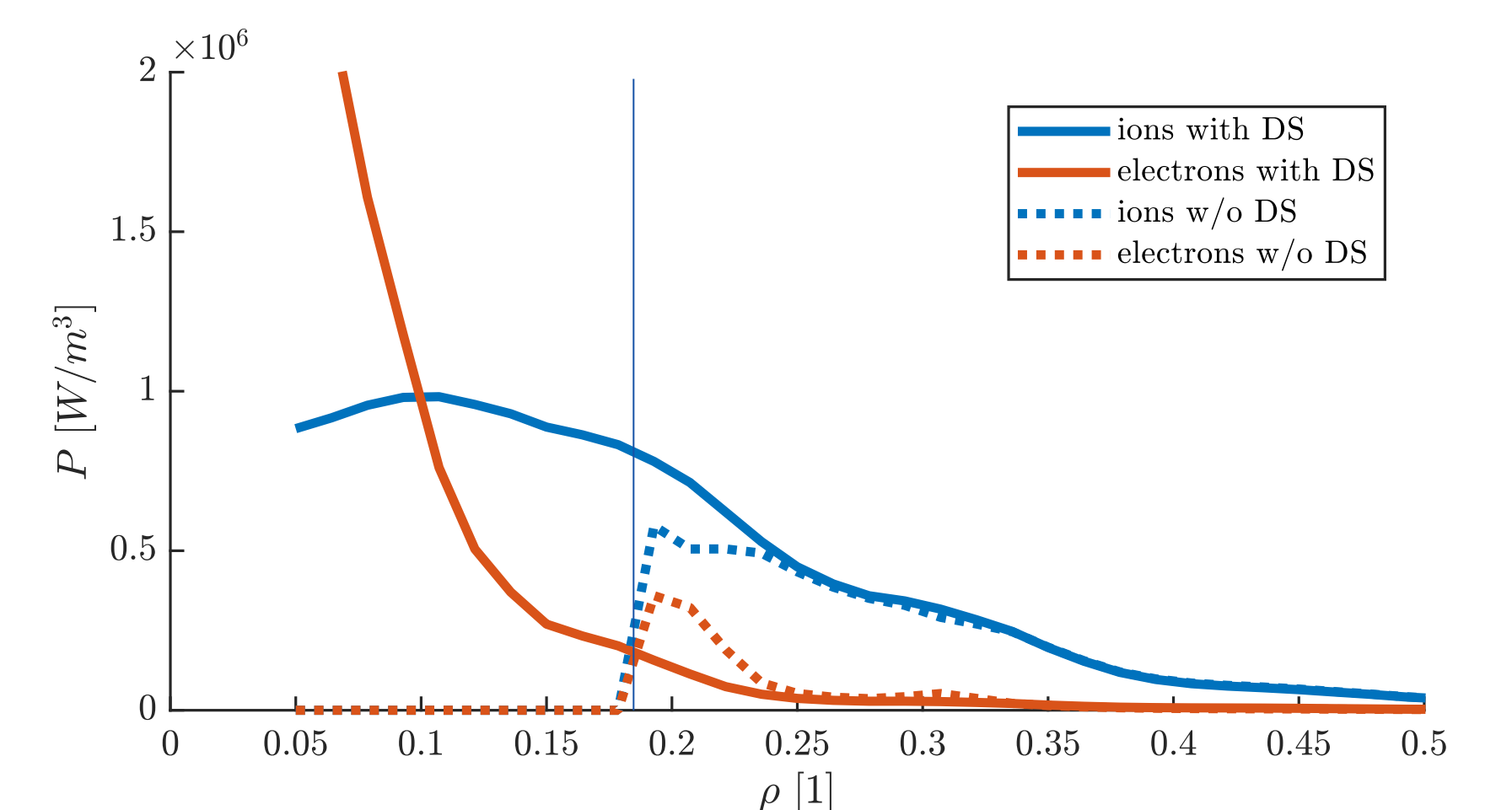
Distribution f(v)

- Including Doppler shift (DS) predicts increased or decreased fast ion temperatures in different regions of plasma
- Without DS localised high fast ion temperature around tangential resonance



Power Absorption

- $\rho < 0.2$: strong ³He absorption at IIHL
- $\rho < 0.2$: No resonances w/o Doppler shift
- $\rho \sim 0.2$: Increased electron absorption w/o Doppler shift
- $\rho \sim 0.2$: Doppler shift pushes tangential resonances away from equatorial plane



Discussion and Conclusions

- 1D Fokker-Planck model is advantageous in simplicity, transparency and rapidity
- Despite being a 1D model, we include Doppler shift and a pitch angle distribution
- Singularities in RF diffusion coefficient have been resolved
- Fast ion temperature modified by DS
- LFS resonance, strong damping at IIHL requires Doppler shifts to model.