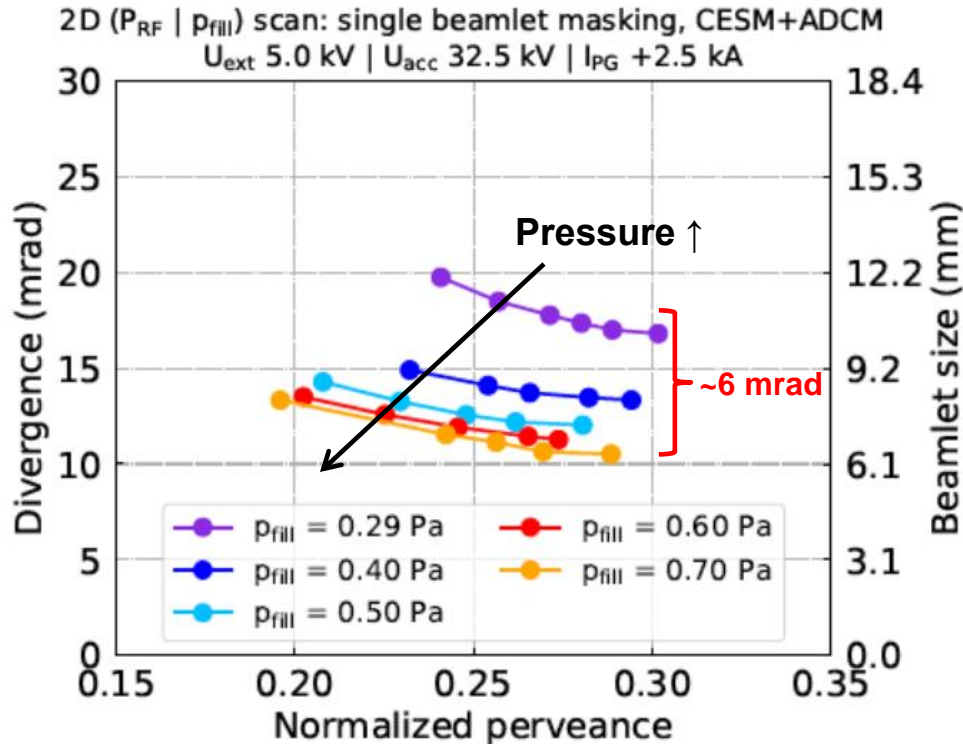


# Key parameters for the H<sup>-</sup> velocity distribution at the plasma meniscus of a caesiated negative ion source

A. Pimazzoni, E. Sartori, G. Serianni and P. Veltri

8<sup>th</sup> International Symposium on Negative Ions, Beams and Sources  
Padova, Italy, October 2-7, 2022

- **Motivation**
  
- ICARO: a test particle code w/ Monte Carlo collisions
  - Description and comparison with literature
  - Investigation of the pressure dependence
  
- Coupling of ICARO to a ray tracing code

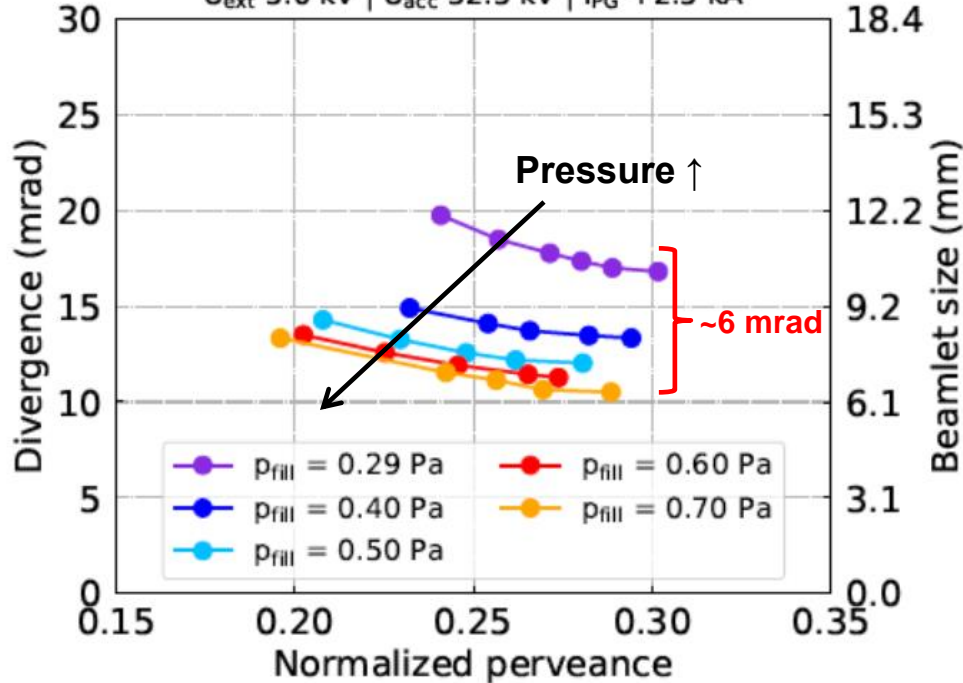


## Experimental evidence

Beam divergence from RF sources under development for the ITER HNB/DNB is highly sensitive to the source filling pressure

$$Div = Div(I_{ex}/U_{ex}^{1.5}, E_{beam}, T_H)$$

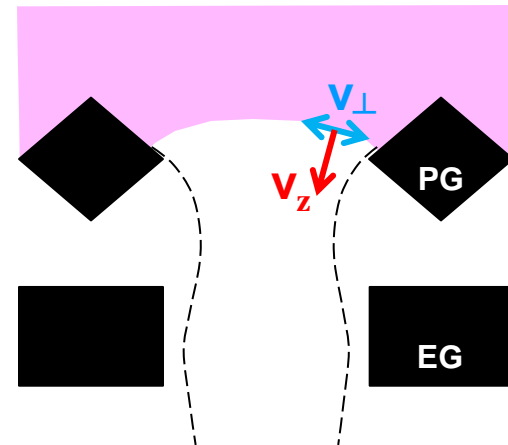
2D ( $P_{RF} | p_{fill}$ ) scan: single beamlet masking, CESM+ADCM  
 $U_{ext}$  5.0 kV |  $U_{acc}$  32.5 kV |  $I_{PG}$  +2.5 kA

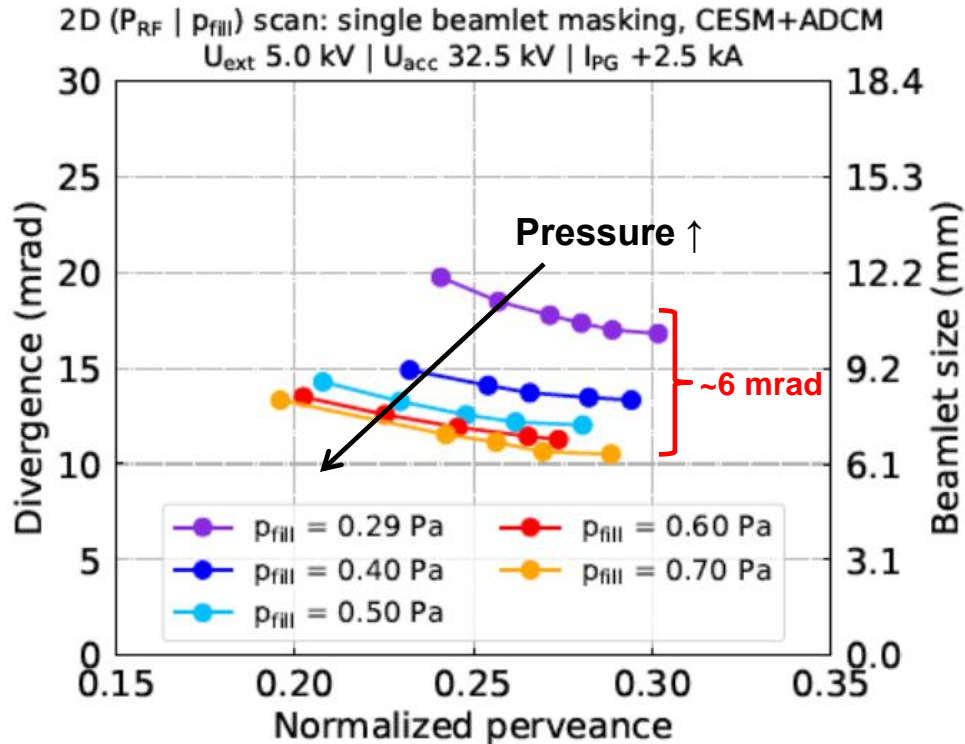


## Experimental evidence

Beam divergence from RF sources under development for the ITER HNB/DNB is highly sensitive to the source filling pressure

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## Experimental evidence

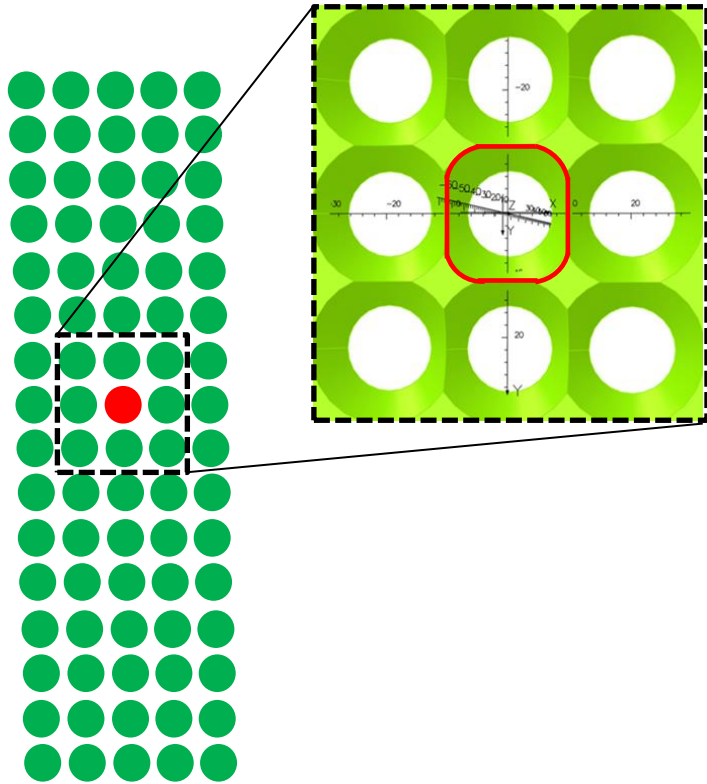
Beam divergence from RF sources under development for the ITER HNB/DNB is highly sensitive to the source filling pressure

$$Div = Div(I_{ex}/U_{ex}^{1.5}, E_{beam}, T_H)$$

## Hypothesis

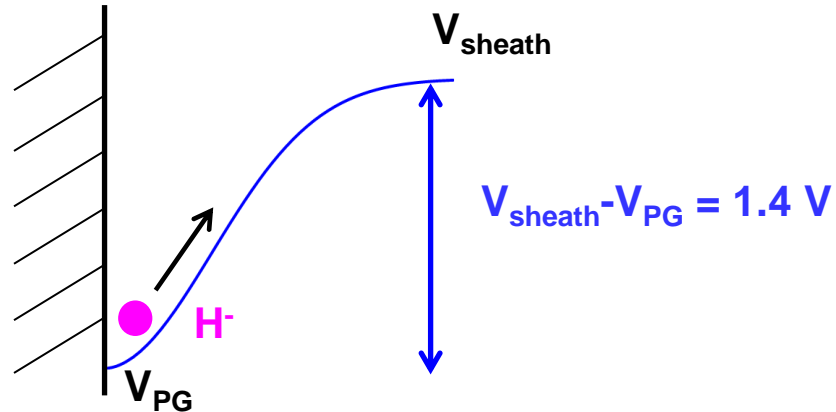
Being perveance and beam energy the same, we can conclude that pressure affects the negative ion velocity distribution at the plasma meniscus  
*But how?*

- Motivation
  
- **ICARO: a test particle code w/ Monte Carlo collisions**
  - **Description and comparison with literature**
  - Investigation of the pressure dependence
  
- Coupling of ICARO to a ray tracing code

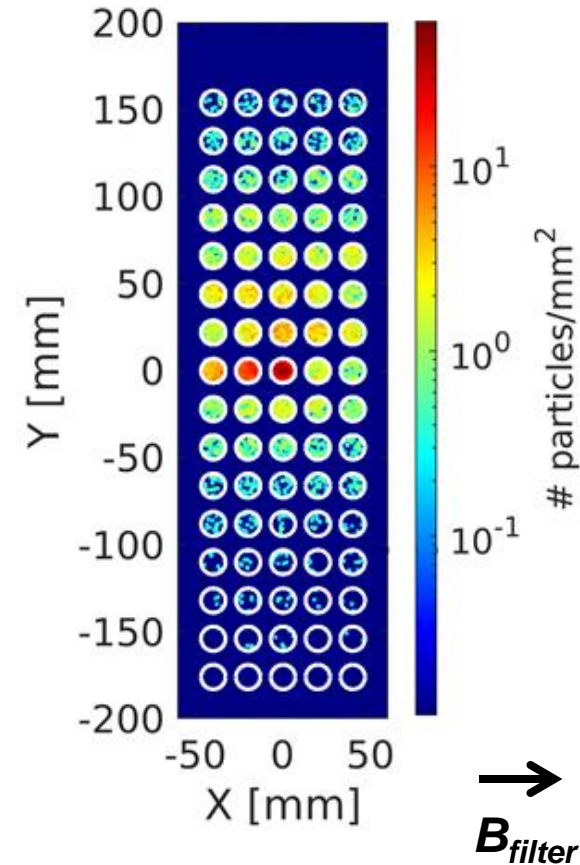
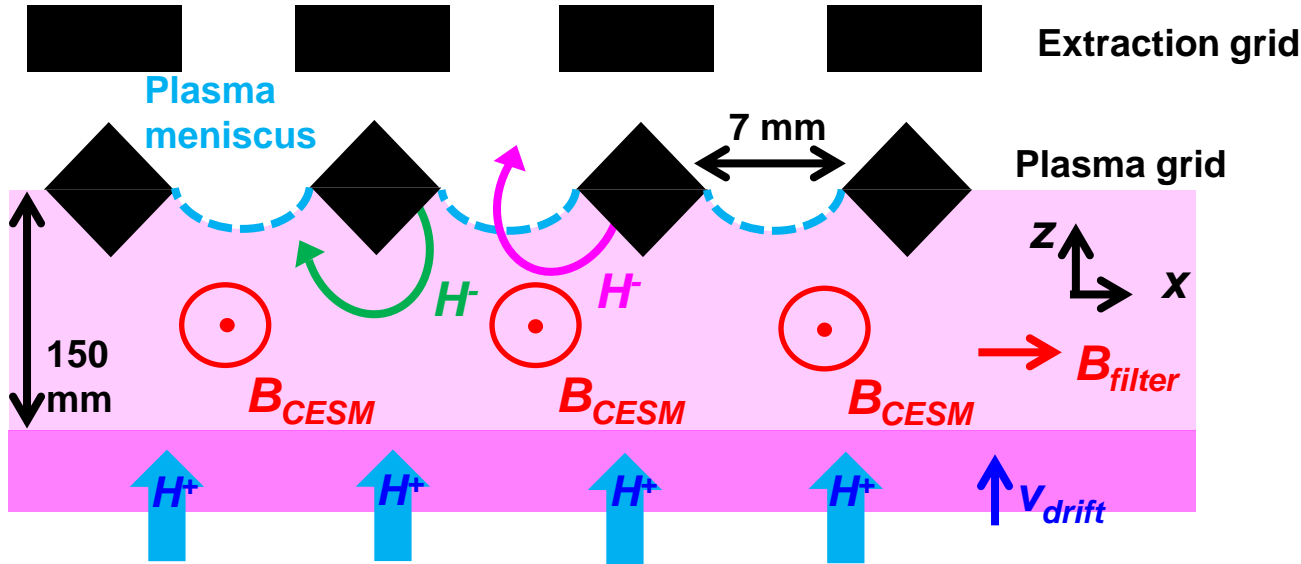


SPIDER beamlet group:  
5x16 apertures

- Emission with a cosine distribution from the upstream chamfered part of the PG (**only from the aperture centered in [0,0] mm**)
- $H^-$  are emitted from the PG with an initial energy that already considers the potential difference  $V_{\text{sheath}} - V_{\text{PG}}$ 
  - Virtual cathode is not considered yet



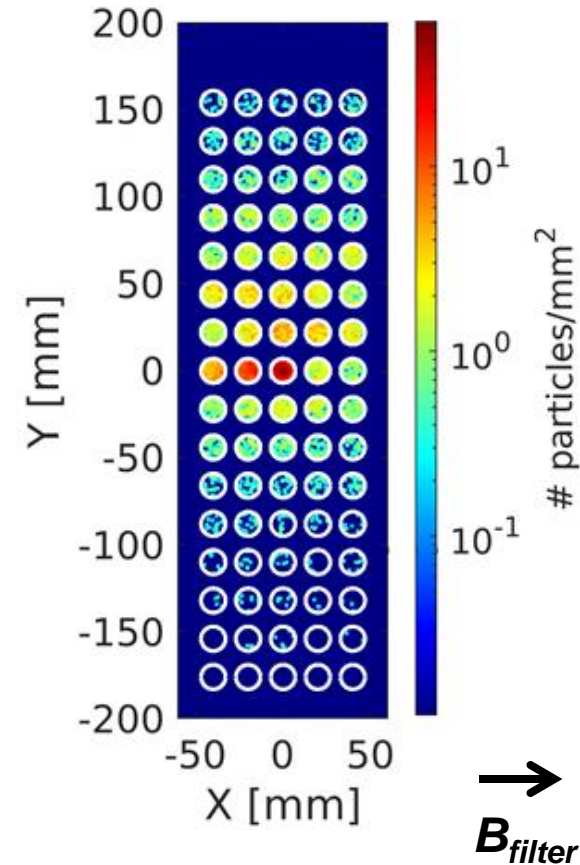
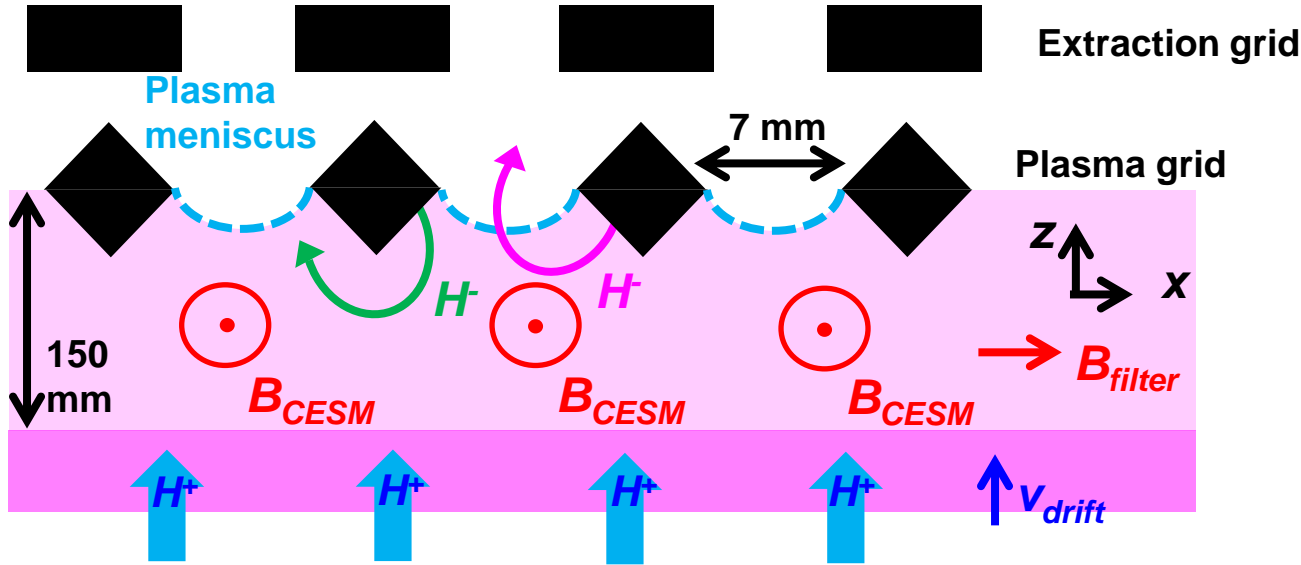
# ICARO: Plasma domain and particle motion



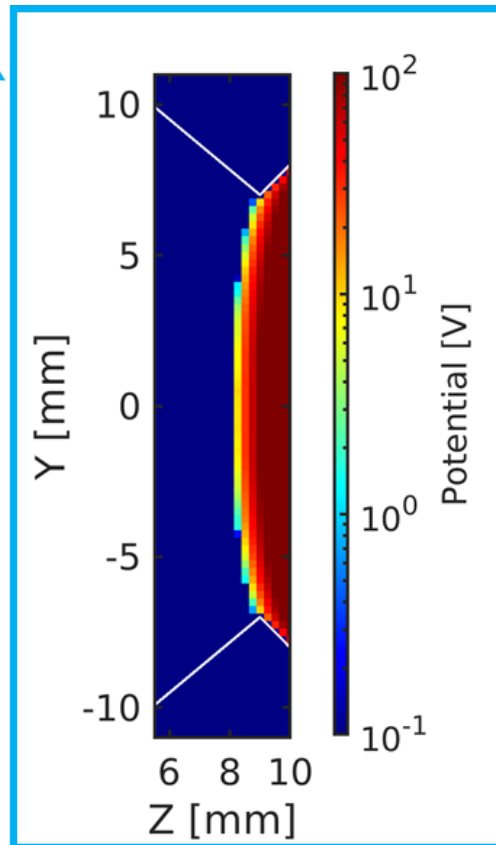
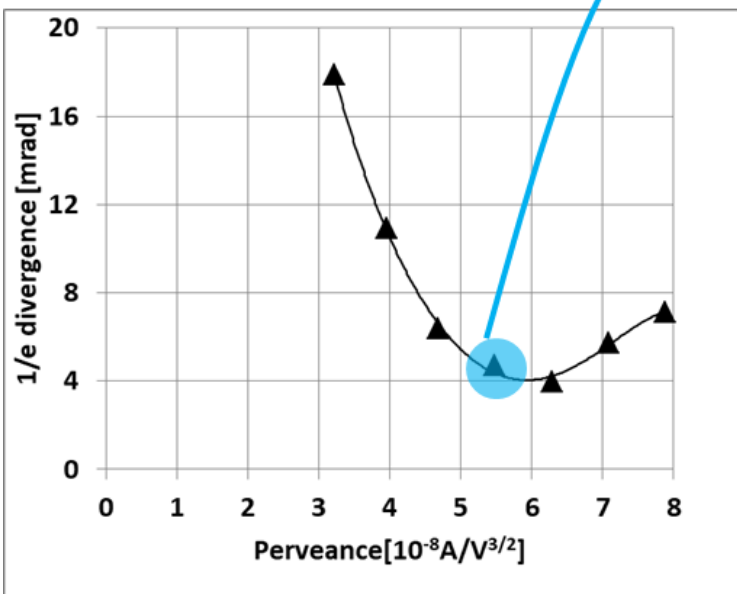
- Particles are moved in the magnetic field (CESM=magnets in the EG; Filter field  $B_{\text{filter}}=B_x=2\text{mT}$ )
- $dt$  is adjusted so to keep  $ds=0.5\text{ mm}$  at each step
- The domain length is set to 150 mm



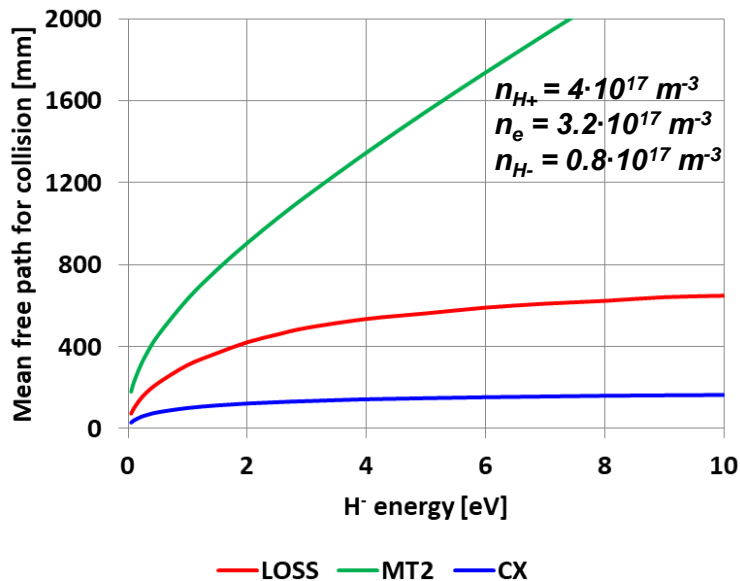
# ICARO: Plasma domain and particle motion



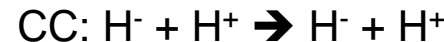
- Plasma parameters ( $n_{H^+}, n_{H^-}, n_e, T_e, T_{H^+}, v_{drift}$ ) and gas parameters ( $n_{H_0}, n_{H_2}, T_{H_0}, T_{H_2}$ ) assumed homogeneous
- Only  $H^+$  are considered as positive ions
- Collisions are implemented with Monte Carlo method



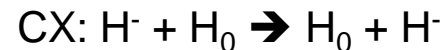
- The meniscus surface is loaded from OPERA3D
- **Perveance match** case ( $\Pi=5.5-6.3 \cdot 10^{-8} \text{ A/V}^{3/2}$ ) are considered; this meniscus is very flat, so that the direct extraction is expected to be very little



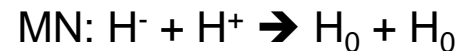
## Momentum transfer w/ plasma



## Momentum transfer w/ gas



## Neutralization by the plasma



## Neutralization by the gas



### Cross-sections

- C.F. Barnett, *Atomic Data for Fusion*, Volume 1, (ORNL-6086), 1990.
- T. Tabata et al., *Atomic Data and Nuclear Data Tables* **76**, 1–25 (2000)

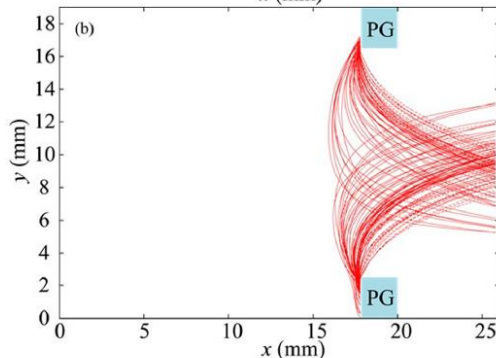
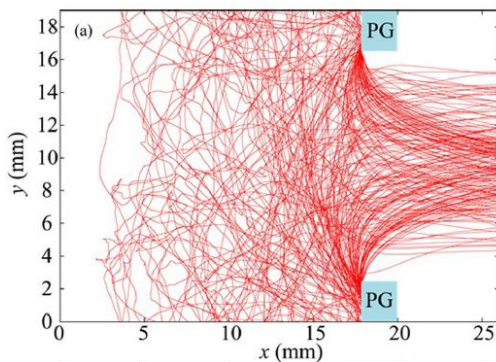
### Gas parameters

- U. Fantz et al., *Front. Phys.* **9**:709651 (2021)

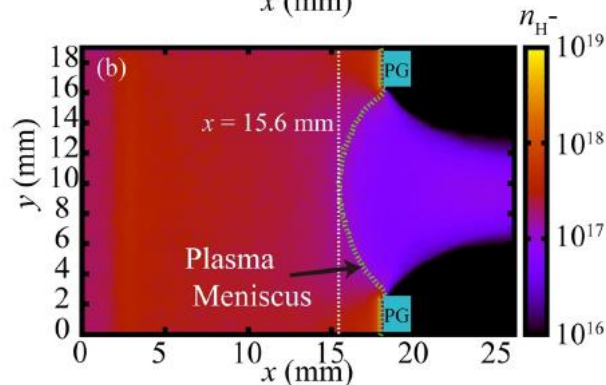
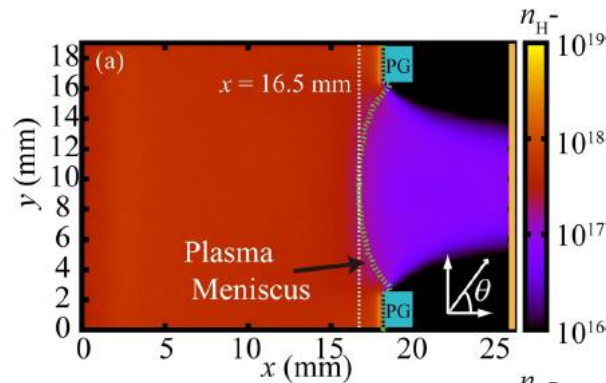
CC are treated as in T. Takizuka, H. Abe *J. Comput. Phys.*, **25** (1977)

# The role of Coulomb collisions (CC)

w/ CC  
( $H^+ - H^-$ )



w/o CC



Coulomb collisions strongly affects the extraction probability and the meniscus.

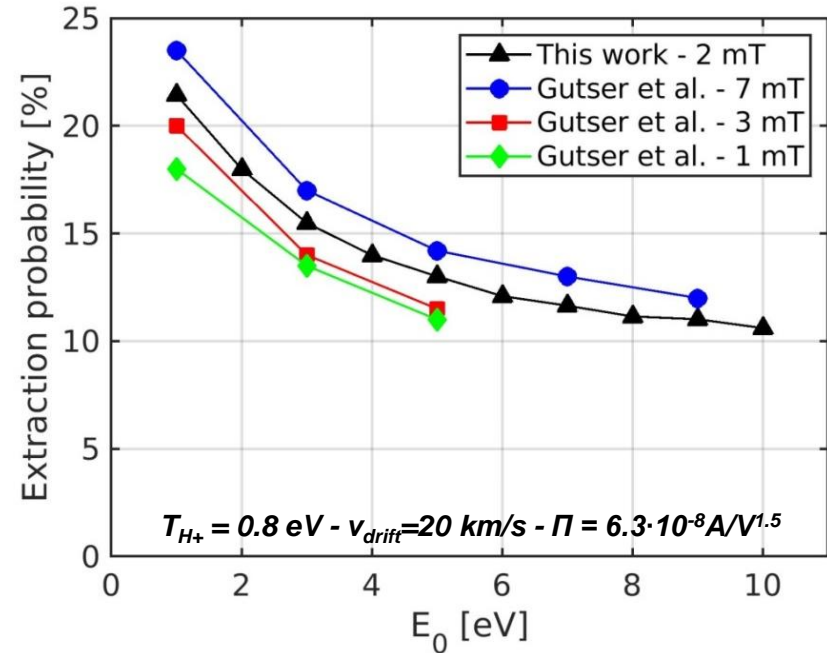
W/o CC, particles are extracted only through direct extraction → Poor optics

$$V_{CC} \sim |v_{H^+} - v_{H^-}|^{-3}$$

S. Nishioka et al. *J Appl. Phys.* **123**, 063302 (2018)

	R. Gutser et al. *	This work
$T_{H2,1}$ [K]	1200	630 **
$T_{H2,2}$ [K]		4600 **
$T_{H0,1}$ [eV]	0.8	0.19 **
$T_{H0,2}$ [eV]		2.5 **
$n_{H2,1}$ [m <sup>-3</sup> ]	$4.0 \cdot 10^{19}$	$1.0 \cdot 10^{19}$
$n_{H2,1}$ [m <sup>-3</sup> ]		$0.5 \cdot 10^{19}$
$n_{H0,1}$ [m <sup>-3</sup> ]	$10 \cdot 10^{18}$	$2.2 \cdot 10^{18}$
$n_{H0,2}$ [m <sup>-3</sup> ]		$2.2 \cdot 10^{18}$
$T_{H+}$ [eV]	0.8	[0.4,1.6]
$v_{drift}$ [km/s]	0	[10,20]
$T_e$ [eV]	2.0	2.0
$n_e$ [m <sup>-3</sup> ]	$5.0 \cdot 10^{17}$	$3.2 \cdot 10^{17}$
$n_{H+}$ [m <sup>-3</sup> ]	$5.5 \cdot 10^{17}$	$4.0 \cdot 10^{17}$
$n_{H-}$ [m <sup>-3</sup> ]	$0.5 \cdot 10^{17}$	$0.8 \cdot 10^{17}$

Extraction probability  $p_{extr} = \frac{N_{extr}}{N_{tracked}}$



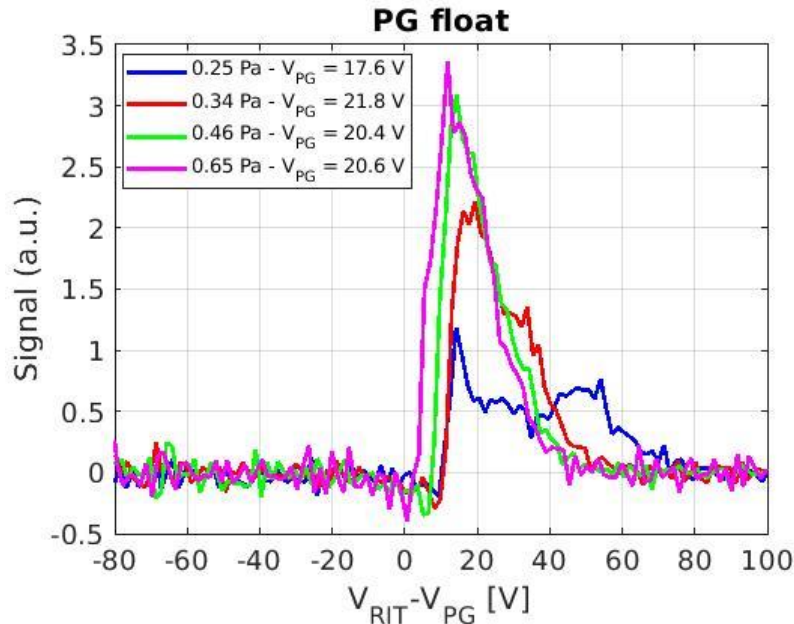
\* R. Gutser et al. PPCF **52** (2010) 045017

\*\* U. Fantz et al., Front. Phys. 9:709651 (2021)

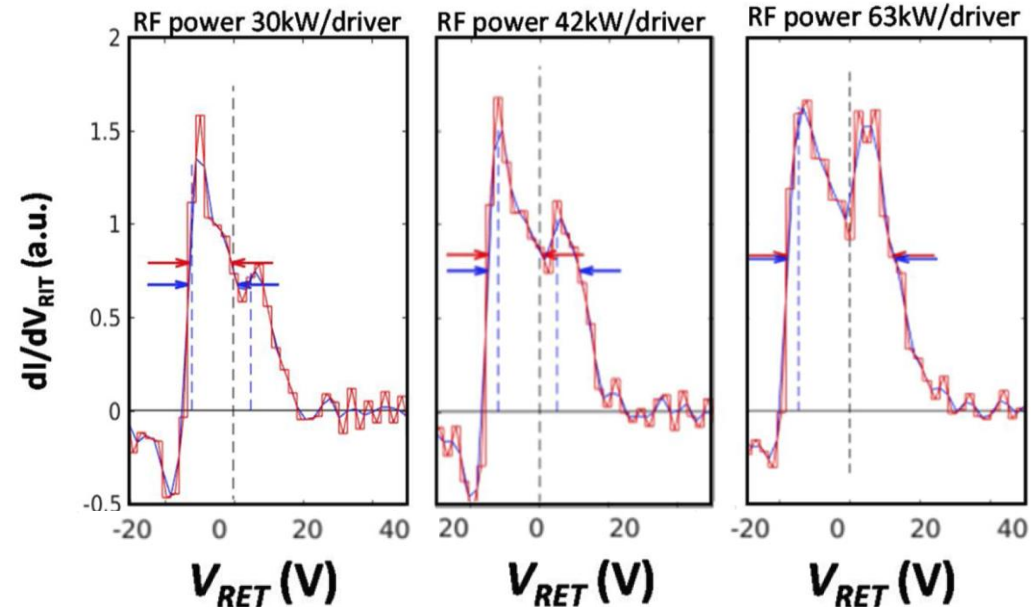
# Energy distribution of positive ions

The dependence of the positive ion energy distribution on the source parameters was investigated in SPIDER with a retarding field energy analyzer (RFEA)

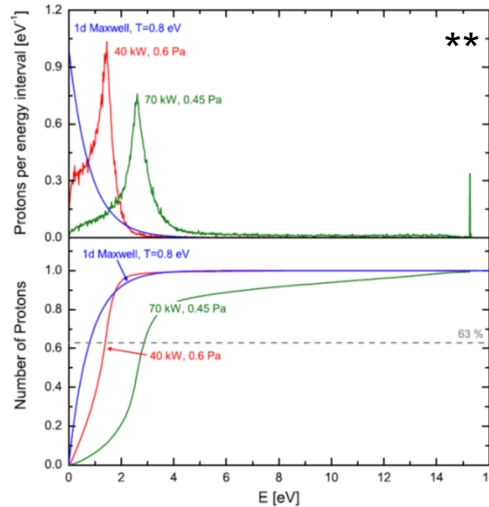
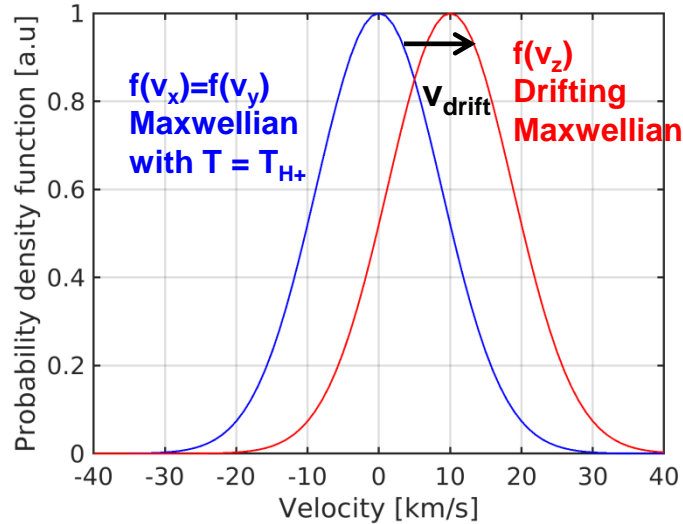
Pressure dependence



RF power dependence



# ICARO: Positive ion velocity distribution



**@ BATMAN (IPP, Garching)**

Mach probe:

$v_{drift} \sim 10 \text{ km/s}$  for  $p_{fill} = 0.7 \text{ Pa}^*$

Probes+ MC models:

$T_{H+} \sim 0.8 \text{ eV}$  @ PG for

$P_{RF} = 40 \text{ kW/dr} - p_{fill} = 0.6 \text{ Pa}^{**}$

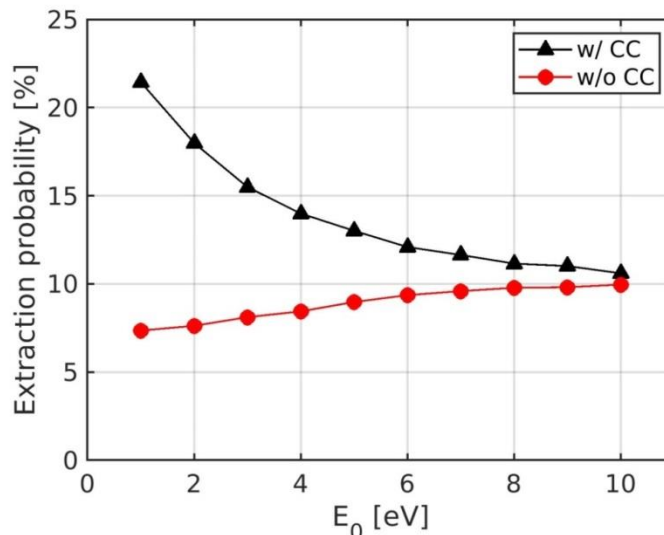
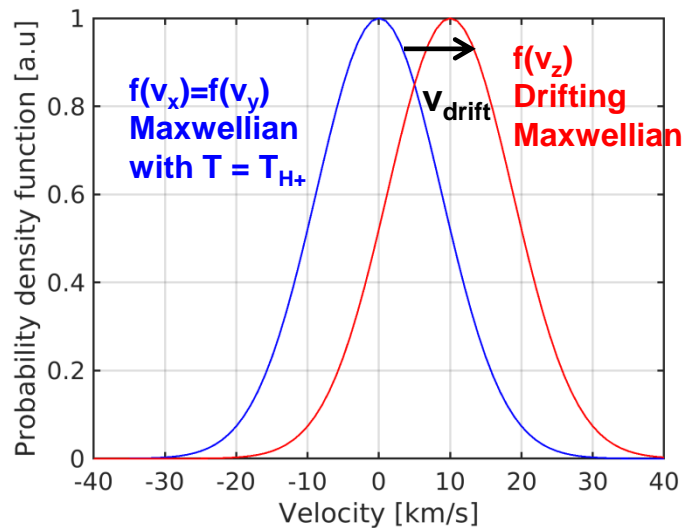
The efficacy of CC depends on:

- $E_{H-}$
- $n_{H+}$
- $f(E_{H+}) \rightarrow T_{H+}, v_{drift}$

\* M. Bandyopadhyay et al. *A.J. Appl. Phys.* **96**, 4107 (2004)

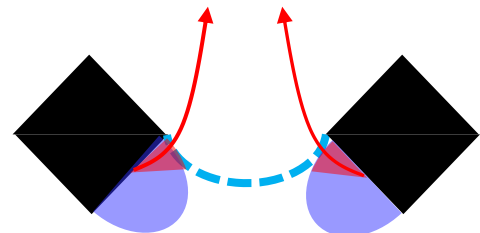
\*\* D. Wuenderlich et al., *PPCF54* (2012) 125002





( $E_0$ : emission energy of  $H^+$  before plasma sheath)

$$p_{extr} = \frac{N_{extr}}{N_{tracked}}$$



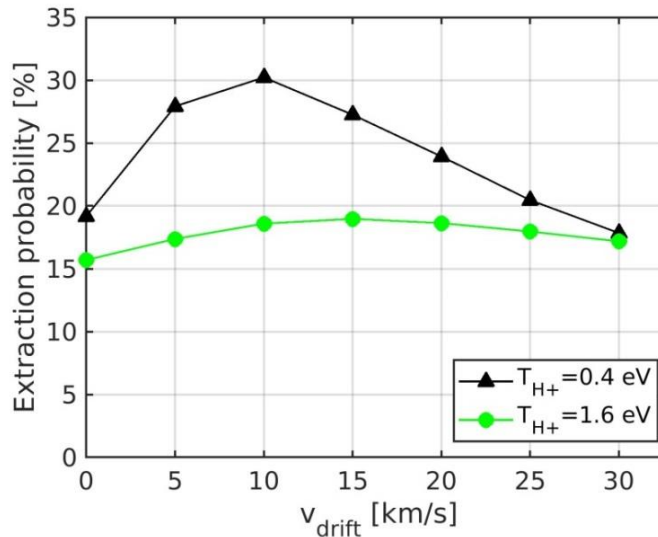
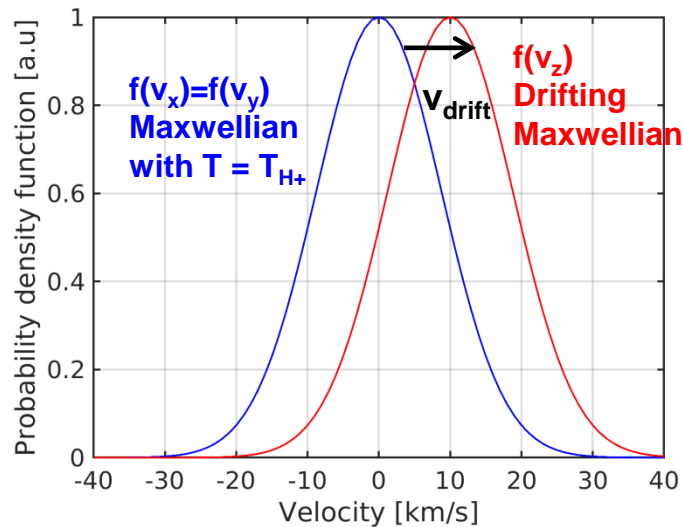
**Favourable emission angle**

The efficacy of CC depends on:

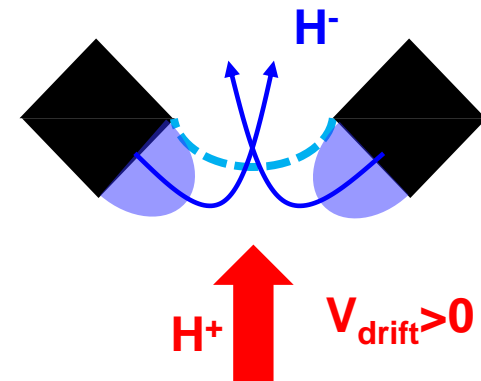
- $E_{H^+}$
- $n_{H^+}$
- $f(E_{H^+}) \rightarrow T_{H^+}, v_{drift}$

Since  $v_{CC} \sim |\mathbf{v}_{H^+} - \mathbf{v}_{H^-}|^{-3}$  CC are more effective for low energy  $H^+$





$$p_{extr} = \frac{N_{extr}}{N_{tracked}}$$

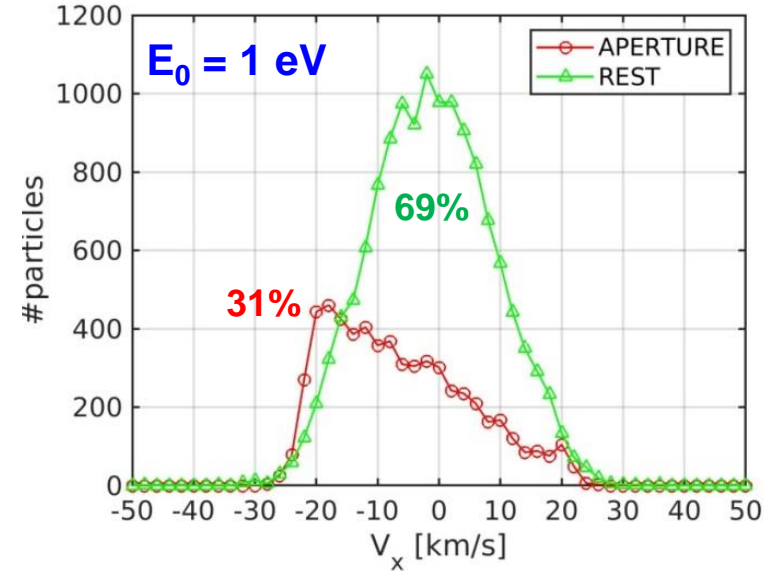
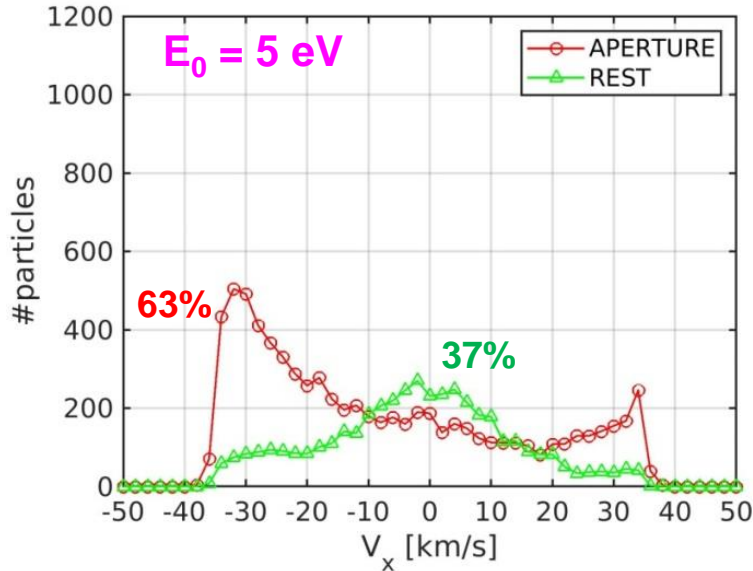
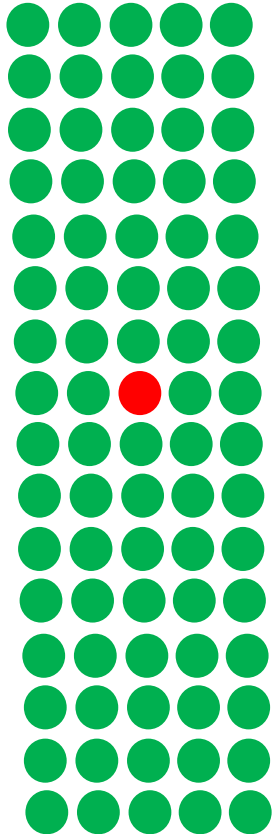


The efficacy of CC depends on:

- $E_{H-}$
- $n_{H+}$
- $f(E_{H+}) \rightarrow T_{H+}, v_{drift}$

Since  $v_{CC} \sim |\mathbf{v}_{H+} - \mathbf{v}_{H-}|^{-3}$  there is an optimum drift velocity for  $H^+$  ions. Beyond this velocity, CC lose efficacy

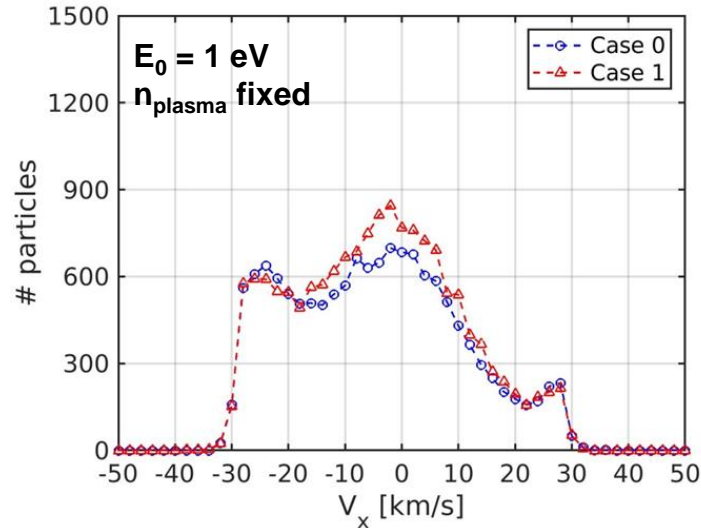
# Results: velocity distribution at meniscus



- $H^-$  emitted with large energies are unlikely to be brought back to the apertures by collisions and magnetic fields.
- For  $H^-$  emitted with low energy, CC are effective in reversing the  $H^-$  velocity

- Motivation
  
- ICARO: a test particle code w/ Monte Carlo collisions
  - Description and comparison with literature
  - **Investigation of the pressure dependence**
  
- Coupling of ICARO to a ray tracing code

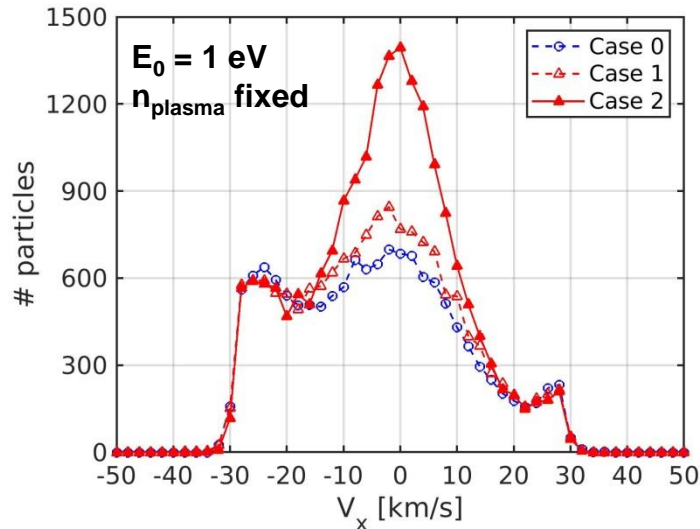
	$p_{\text{fill}}$ [Pa]	Density of gas species	$T_e$ [eV]	$T_{H^+}$ [eV]	$v_{\text{drift}}$ [km/s]
<b>Case 0</b>	0.3	$n_{\text{standard}}$	2	0.8	20
<b>Case 1</b>	0.6	$2n_{\text{standard}}$	1	0.8	20



The optics dependence on pressure:

- **Cannot be explained** by the increased gas density in the extraction region (Case 0 vs Case 1)

	$p_{\text{fill}}$ [Pa]	Density of gas species	$T_e$ [eV]	$T_{H^+}$ [eV]	$v_{\text{drift}}$ [km/s]
<b>Case 0</b>	0.3	$n_{\text{standard}}$	2	0.8	20
<b>Case 1</b>	0.6	$2n_{\text{standard}}$	1	0.8	20
<b>Case 2</b>	0.6	$2n_{\text{standard}}$	1	0.4	10 *



The optics dependence on pressure:

- **Cannot be explained** by the increased gas density in the extraction region (Case 0 vs Case 1)
- **Can be explained** by reduced energy of the positive ions (Case 0 vs Case 2)

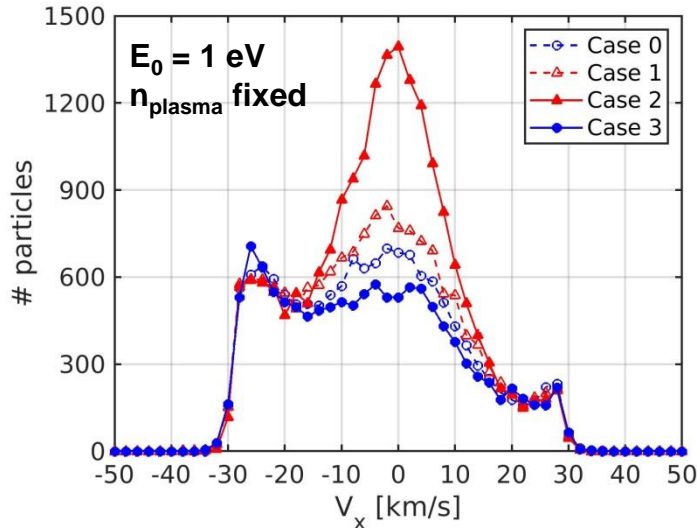
\* M. Bandyopadhyay et al. *A.J. Appl. Phys.* **96**, 4107 (2004)

	$p_{\text{fill}}$ [Pa]	Density of gas species	$T_e$ [eV]	$T_{H^+}$ [eV]	$v_{\text{drift}}$ [km/s]
<b>Case 0</b>	0.3	$n_{\text{standard}}$	2	0.8	20
<b>Case 1</b>	0.6	$2n_{\text{standard}}$	1	0.8	20
<b>Case 2</b>	0.6	$2n_{\text{standard}}$	1	0.4	10 *
<b>Case 3</b>	0.3	$n_{\text{standard}}$	2	1.6	20

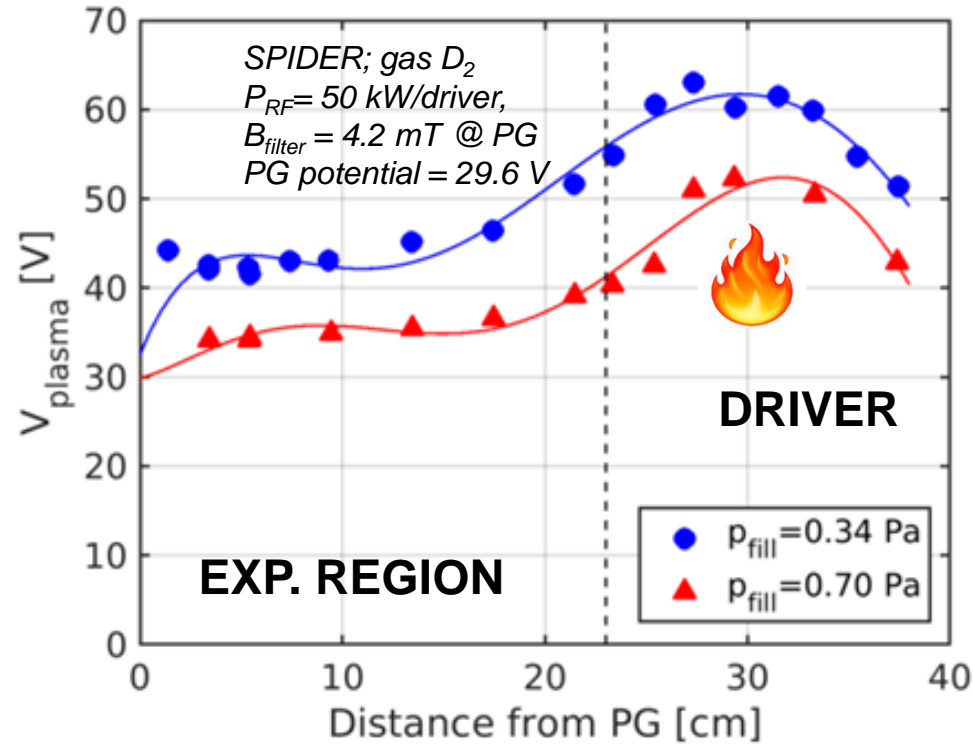


*The effect of pressure is not local, due to MT with neutrals, but global, as pressure affects the whole discharge*

*Pressure influences the  $H^+$  extraction via the positive ion energy distribution (which controls CC efficacy)*

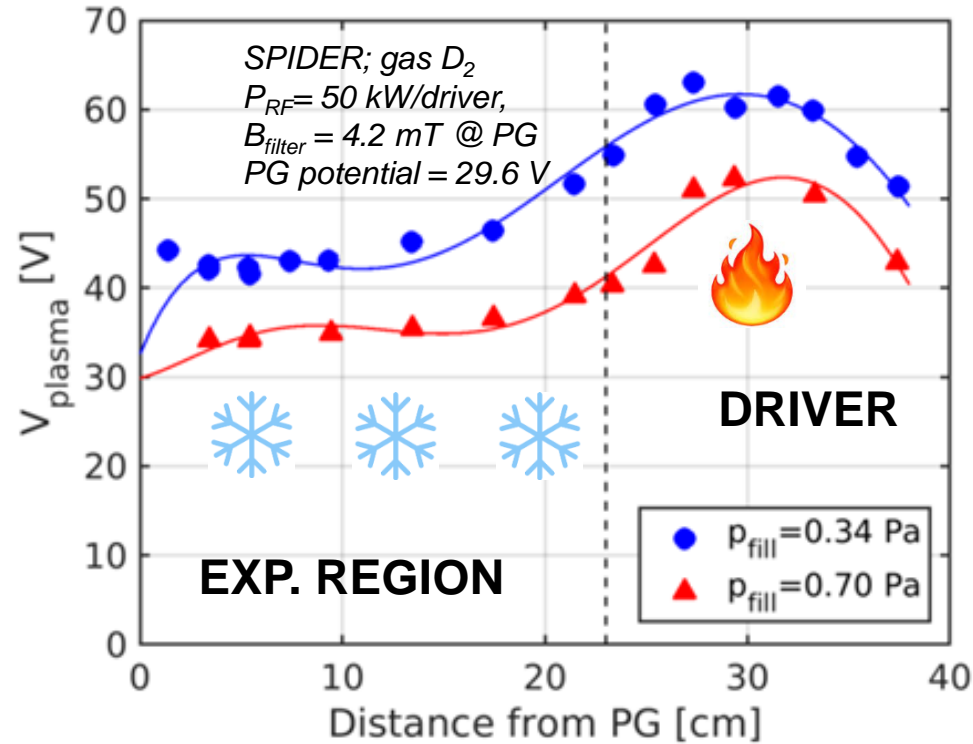


# Energy balance for positive ions



**Energy source term:**  
potential slide between RF  
drivers and plasma grid

# Energy balance for positive ions



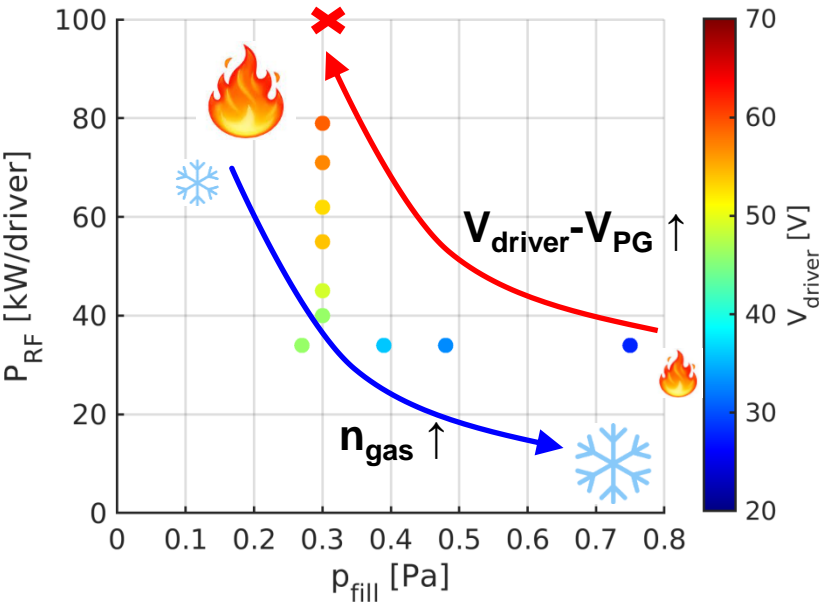
**Energy source term:**  
potential slide between RF  
drivers and plasma grid



**Energy loss term:**  
Collisions with neutrals ( $H^0, H_2$ )  
in the expansion region



Expected working point (100 kW/dr – 0.3 Pa)



Data are from:  
P. McNeely et al., *J Appl. Phys.* **18** (2009) 014011

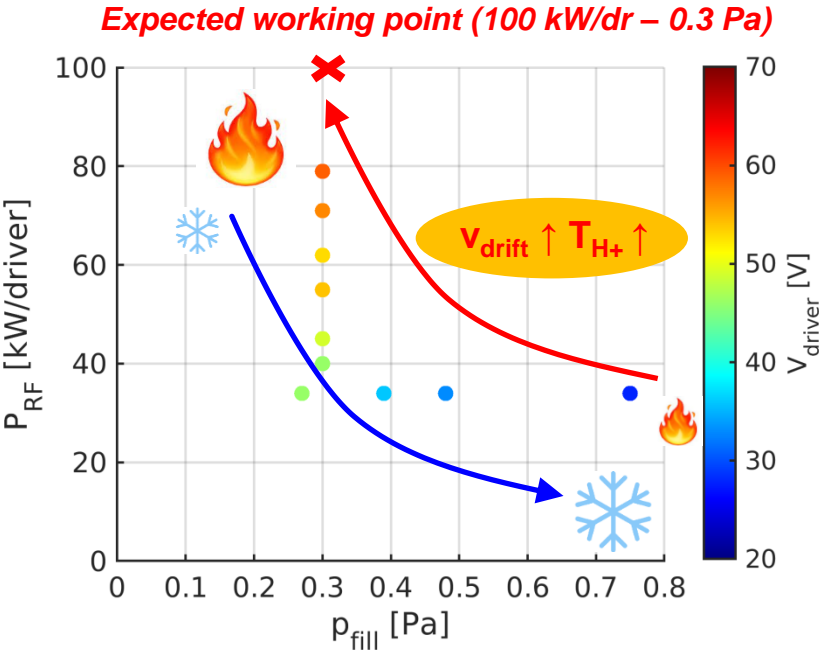
$p_{fill} \leq 0.3$  Pa is needed to limit stripping losses  
**We cannot increase the energy loss term for positive ions**

$P_{RF} \sim 100$  kW/driver is needed to assure the target current density ( $330-355$  A/m<sup>2</sup>) with  $p_{fill} = 0.3$  Pa. By increasing  $P_{RF}$ :

**The energy source term for positive ions would increase**

**The energy loss term would decrease (neutral depletion)**

**$T_{H^+}$  and  $v_{drift}$  are expected to increase**



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P. McNeely et al., *J Appl. Phys.* **18** (2009) 014011

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**$T_{H+}$  and  $v_{drift}$  are expected to increase**

ICARO simulating  $H^-$  emission  
from PG with

$$T_{H^-} \sim R_E T_{H0} = 1.75 \text{ eV}$$

$$T_{H0} = 2.5 \text{ eV}^*$$

$$R_E = 0.7^{**}$$

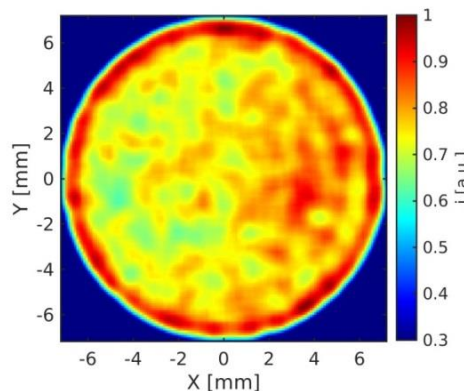
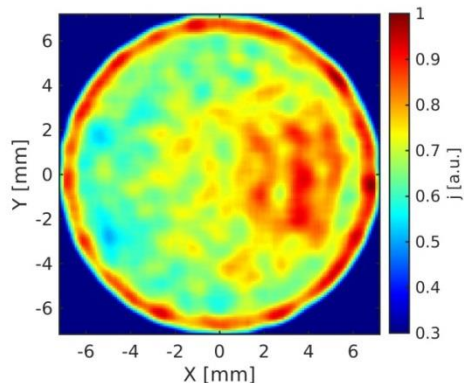
Only  $H^-$  generation from atoms

$$0.3 \text{ Pa} - T_e = 2 \text{ eV}$$

$$T_{H^+} = 1.6 \text{ eV}$$

$$v_{drift} = 20 \text{ km/s}$$

**Current density at meniscus  
(normalized)**



$$0.6 \text{ Pa} - T_e = 1 \text{ eV}$$

$$T_{H^+} = 0.4 \text{ eV}$$

$$v_{drift} = 10 \text{ km/s}$$

# ICARO: emittance at the meniscus

ICARO simulating H<sup>+</sup> emission from PG with

$$T_{H^+} \sim R_E T_{H0} = 1.75 \text{ eV}$$

$$T_{H0} = 2.5 \text{ eV}^*$$

$$R_E = 0.7^{**}$$

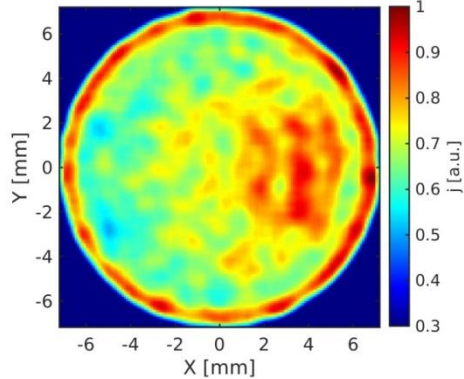
Only H<sup>+</sup> generation from atoms

$$0.3 \text{ Pa} - T_e = 2 \text{ eV}$$

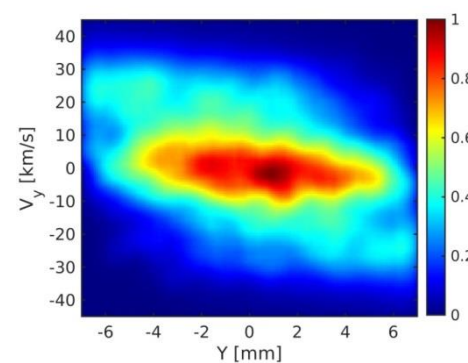
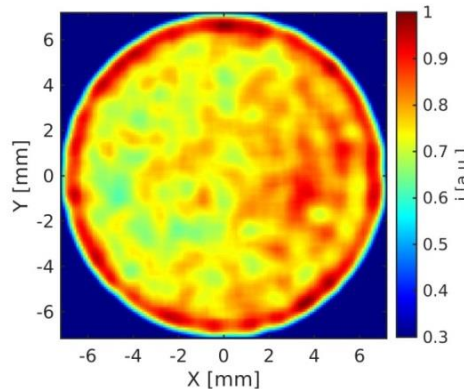
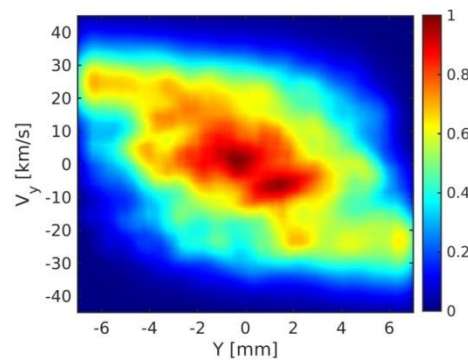
$$T_{H^+} = 1.6 \text{ eV}$$

$$v_{drift} = 20 \text{ km/s}$$

**Current density at meniscus (normalized)**



**Emittance at meniscus (normalized) \***



$$0.6 \text{ Pa} - T_e = 1 \text{ eV}$$

$$T_{H^+} = 0.4 \text{ eV}$$

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# ICARO: emittance at the meniscus

ICARO simulating H<sup>+</sup> emission from PG with

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$$T_{H0} = 2.5 \text{ eV}^*$$

$$R_E = 0.7^{**}$$

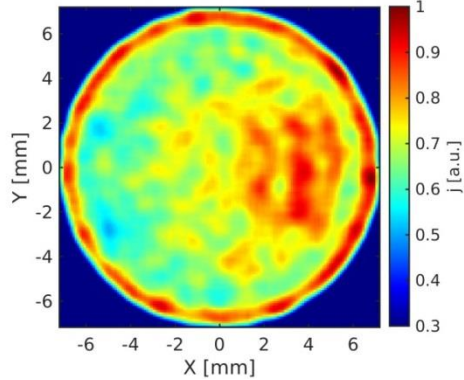
Only H<sup>+</sup> generation from atoms

$$0.3 \text{ Pa} - T_e = 2 \text{ eV}$$

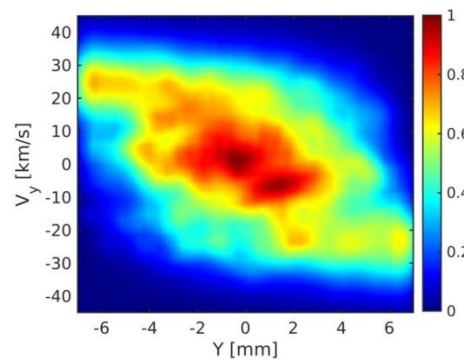
$$T_{H^+} = 1.6 \text{ eV}$$

$$v_{drift} = 20 \text{ km/s}$$

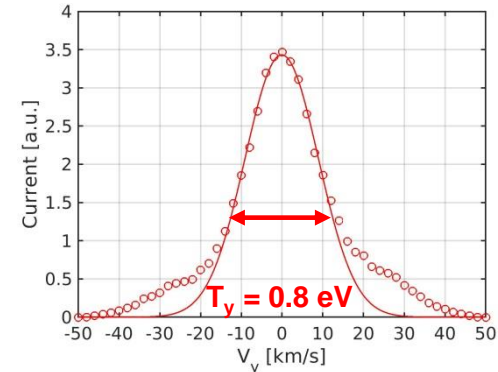
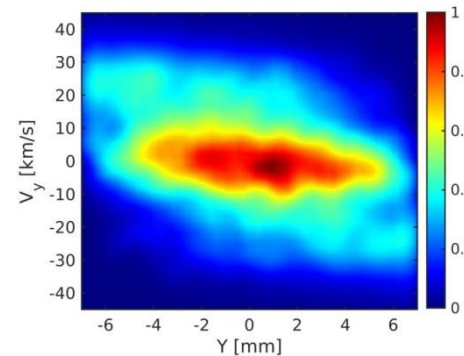
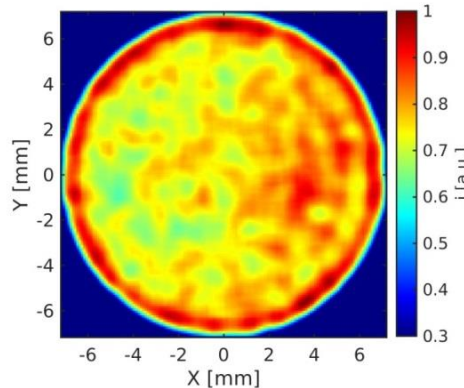
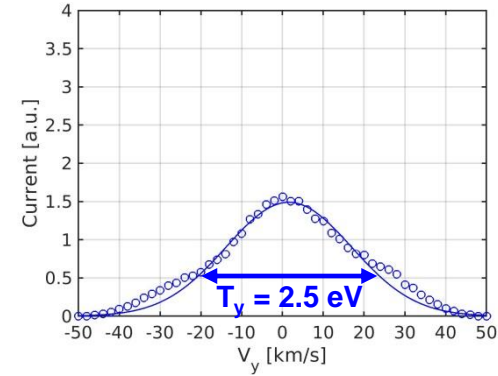
**Current density at meniscus (normalized)**



**Emittance at meniscus (normalized) \***



**Velocity distribution at meniscus**



$$0.6 \text{ Pa} - T_e = 1 \text{ eV}$$

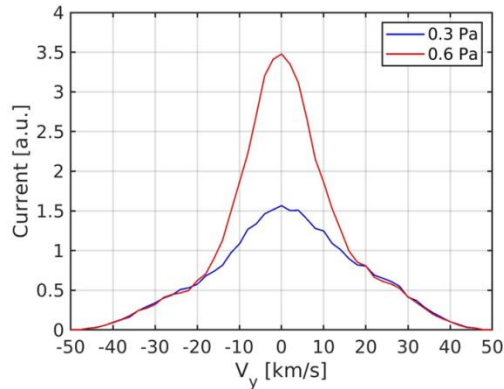
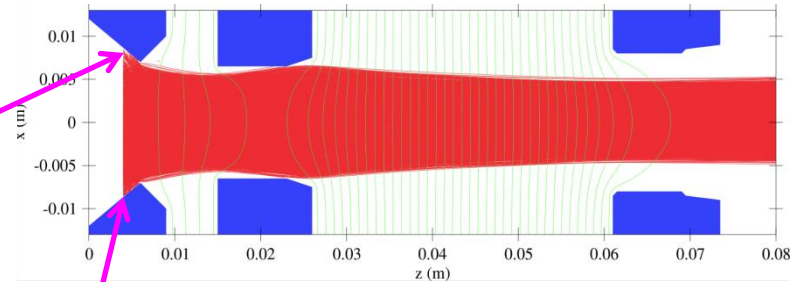
$$T_{H^+} = 0.4 \text{ eV}$$

$$v_{drift} = 10 \text{ km/s}$$

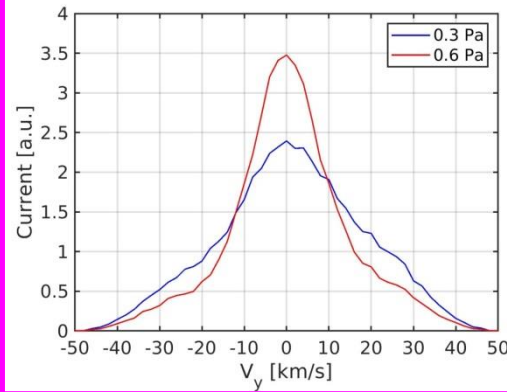
- Motivation
  
- ICARO: a test particle code w/ Monte Carlo collisions
  - Description and comparison with literature
  - Investigation of the pressure dependence
  
- **Coupling of ICARO to a ray tracing code**

- Current is scaled so to have the same perveance and thus the same meniscus shape
- We set same  $I_{ex}$ ,  $U_{ex}$  and  $U_{acc}$  to assess only the role of the emittance at the meniscus
- Settings:  $j_{ext,avg} = 126 \text{ A/m}^2$   
 $U_{ex} = 5 \text{ kV} - U_{acc} = 47.5 \text{ kV}$

\*T. Kalvas et al., *Rev. Sci. Instrum.* **81**, 02B703 (2010)



At meniscus

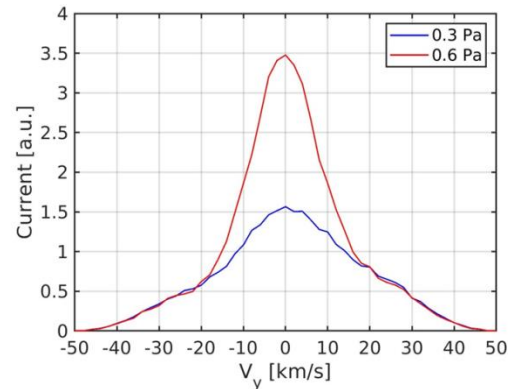
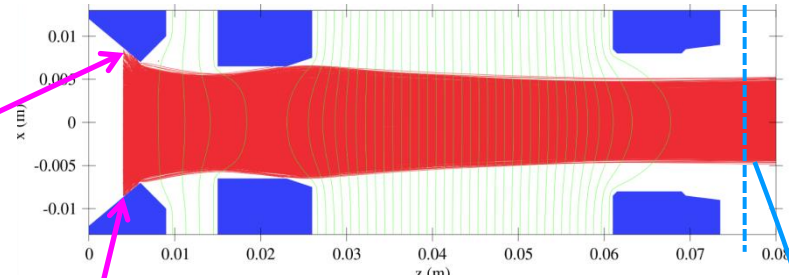


At meniscus (normalized)

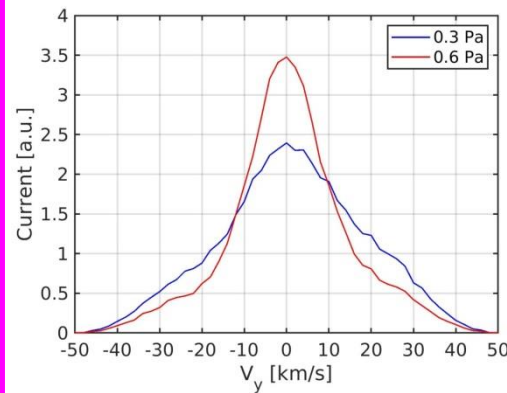


\*T. Kalvas et al., *Rev. Sci. Instrum.* **81**, 02B703 (2010)

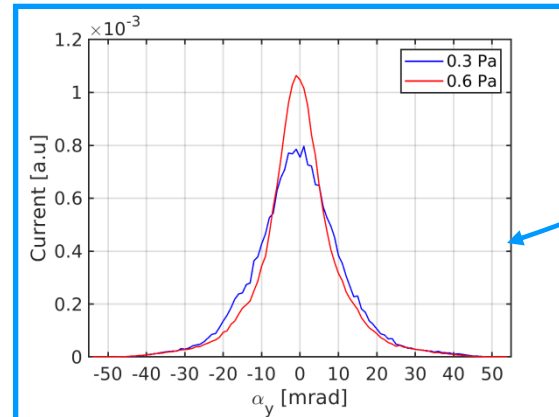
- Current is scaled so to have the same perveance and thus the same meniscus shape
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 $U_{ex} = 5 \text{ kV} - U_{acc} = 47.5 \text{ kV}$



At meniscus



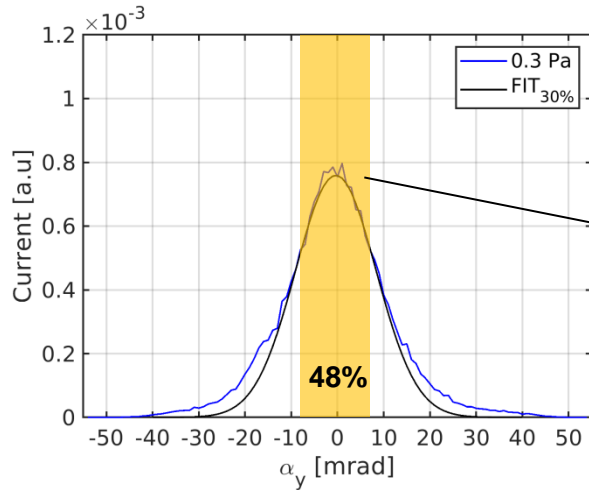
At meniscus (normalized)



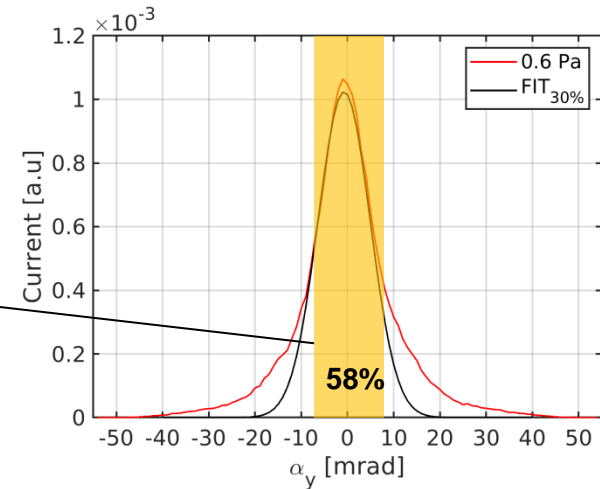
After GG



$P_{\text{fill}}$ [Pa]	$n_{\text{H}0,1}$ [ $10^{18}\text{m}^{-3}$ ]	$n_{\text{H}0,2}$ [ $10^{18}\text{m}^{-3}$ ]	$n_{\text{H}2,1}$ [ $10^{18}\text{m}^{-3}$ ]	$n_{\text{H}2,2}$ [ $10^{18}\text{m}^{-3}$ ]	$T_e$ [eV]	$T_{\text{H}^+}$ [eV]	$V_{\text{drift}}$ [km/s]	div <sub>30%</sub> [mrad]	Fraction in ITER angular acceptance (-7,7) mrad [%]
0.3	2.2	2.2	10.0	5.0	2	1.6	20	12.3	48
0.6	4.5	4.5	19.9	9.9	1	0.4	10	8.0	58



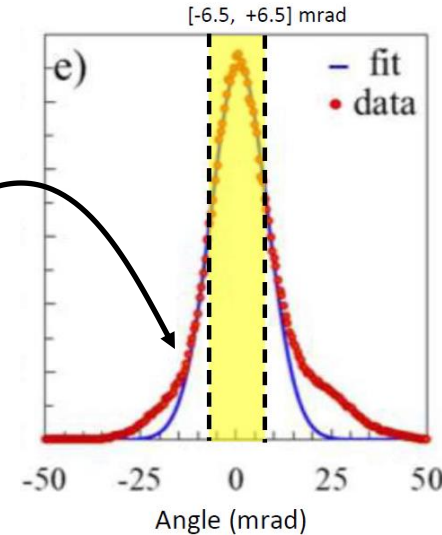
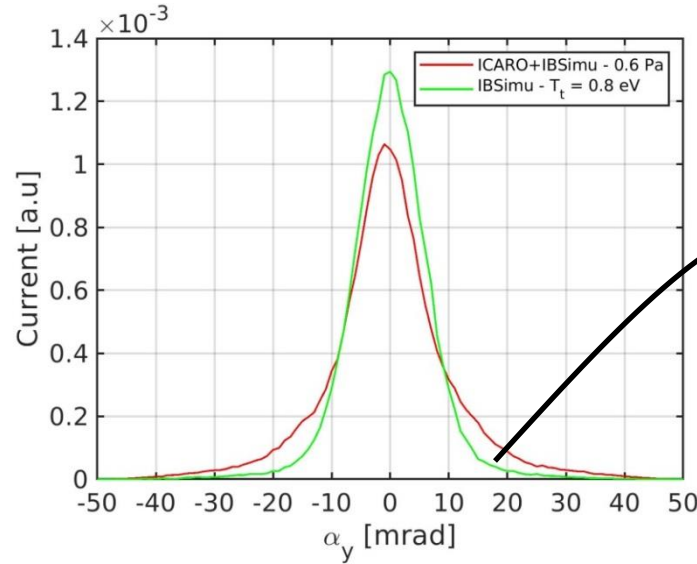
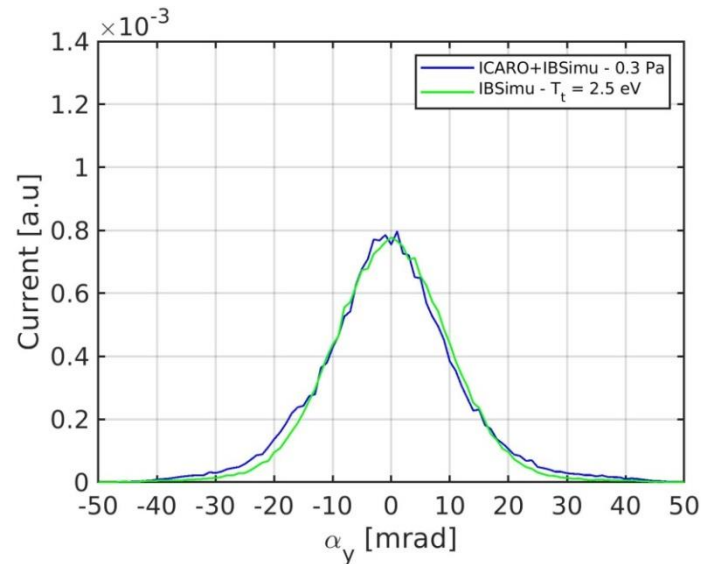
ITER HNB  
«transmission  
window»



**0.3 Pa -  $T_e = 2$  eV -  $T_{\text{H}^+} = 1.6$  eV -  $v_{\text{drift}} = 20$  km/s**

**0.6 Pa -  $T_e = 1$  eV -  $T_{\text{H}^+} = 0.4$  eV -  $v_{\text{drift}} = 10$  km/s**

We compare the results by ICARO+IBSimu with the ones of a **standard IBSimu run**. In the IBsimu standard simulation, we set  $T_t$  equal to the estimate by ICARO



**0.3 Pa -  $T_e = 2$  eV -  $T_{H^+} = 1.6$  eV -  $v_{drift} = 20$  km/s**

**0.6 Pa -  $T_e = 1$  eV -  $T_{H^+} = 0.4$  eV -  $v_{drift} = 10$  km/s**

Using ICARO the tails seen by the beam diagnostics are reproduced

Talk by P. Veltri  
Mon O5

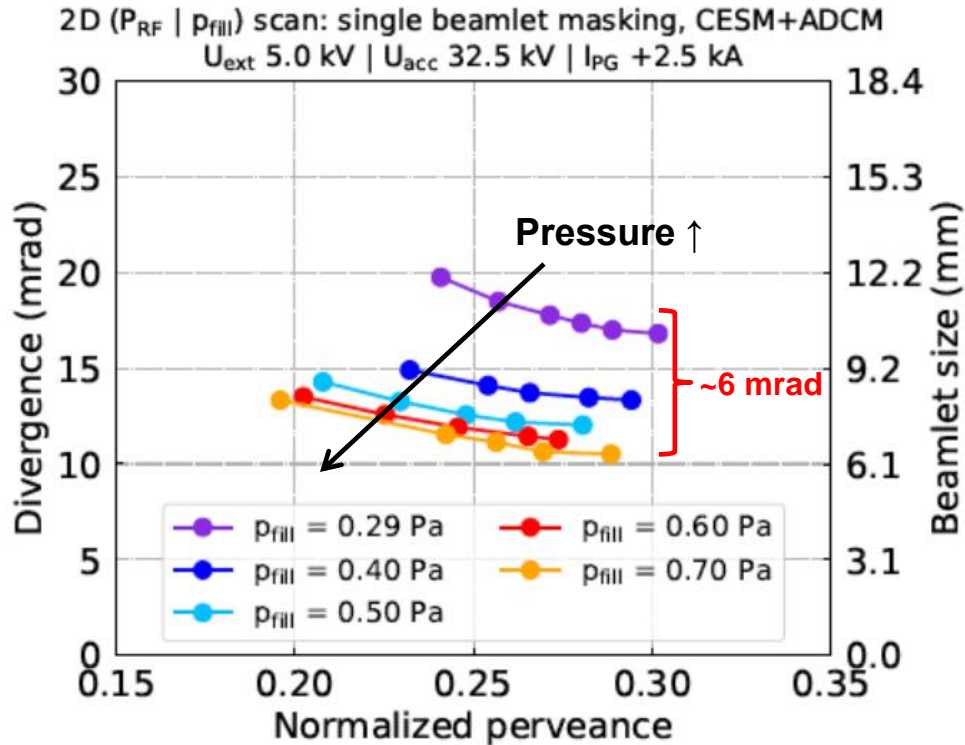
- A particle tracing code with Monte-Carlo collisions, named ICARO, was developed; similar to the code TrajAn by Gutser et al.
  - Trends of previous literature reproduced
  - Highlighted the effect of Coulomb collisions on the H<sup>-</sup> velocity distribution
- Starting from recent measurements of the gas parameters, the code:
  - Provides a possible explanation for the optics dependence on pressure, but only postulating non-local effects ( $v_{\text{drift}}$ ,  $T_{\text{H}^+}$  from RF drivers)
  - Justifies the existence of a broad component (halo) for the beamlet
  - Simulated velocity distribution at GG exit seems more similar to experiments than those from standard beam acceleration models (e.g. IBsimu  $T_{\text{t}}$ )

***Any input to experiments?***

***Yes: Reduce energy source term for positive ions***



# SPARE SLIDES

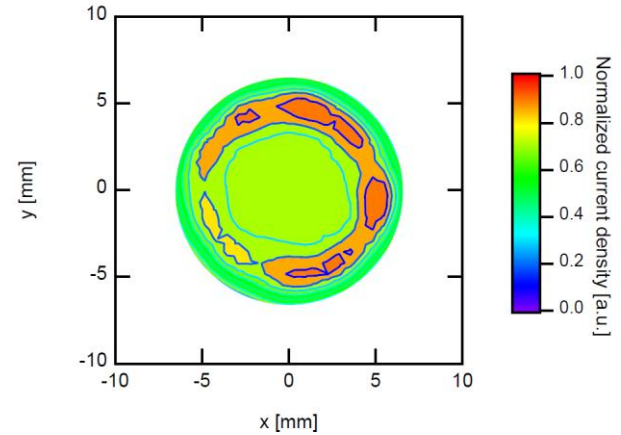
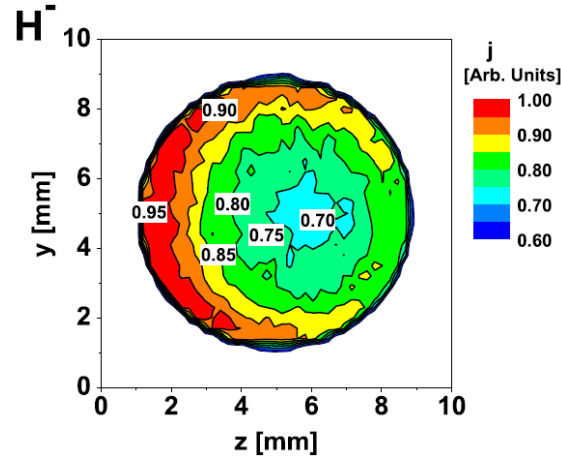
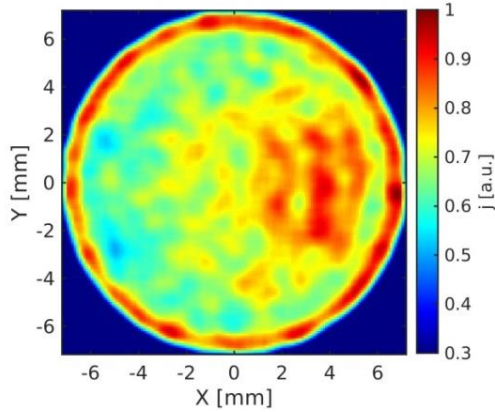


## Experimental evidence

Beam divergence from RF sources under development for the ITER HNB/DNB is highly sensitive to the source filling pressure

$$Div = Div(I_{ex}/U_{ex}^{1.5}, E_{beam}, T_H)$$

# ICARO: simulating a Maxwellian distribution



ICARO simulating  $H^-$  emission from PG with  $T_{H^-} \sim R_E T_{H0} = 1.75$  eV

$T_{H0} = 2.5$  eV\*

$R_E = 0.7$  \*\*

Only  $H^-$  generation from atoms

R. Gutser et al. *PPCF* **52** (2010) 045017

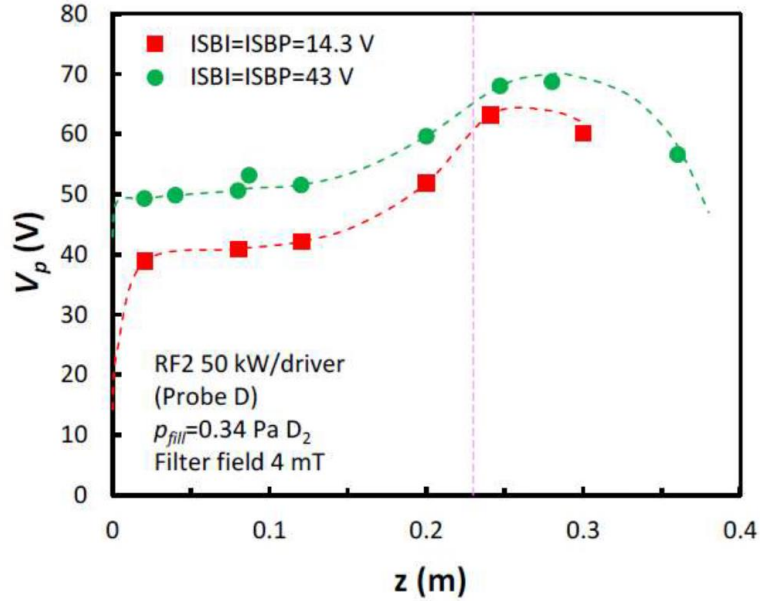
Similar code applied to the case of BATMAN (IPP, Garching)

M. Kisaki et al., *Phys.: Conf. Ser.* **2244** 012061 (2022);

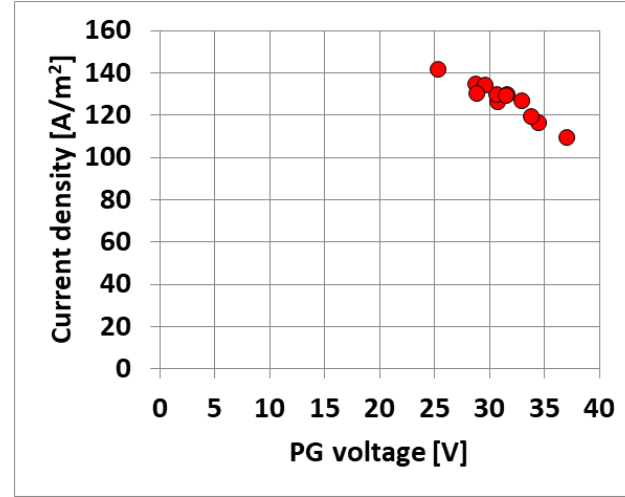
From backward tracing of particles starting from emittance measurements at the Megavolt Test Facility (QST, Japan)

\* U. Fantz et al., *Front. Phys.* **9**:709651 (2021)

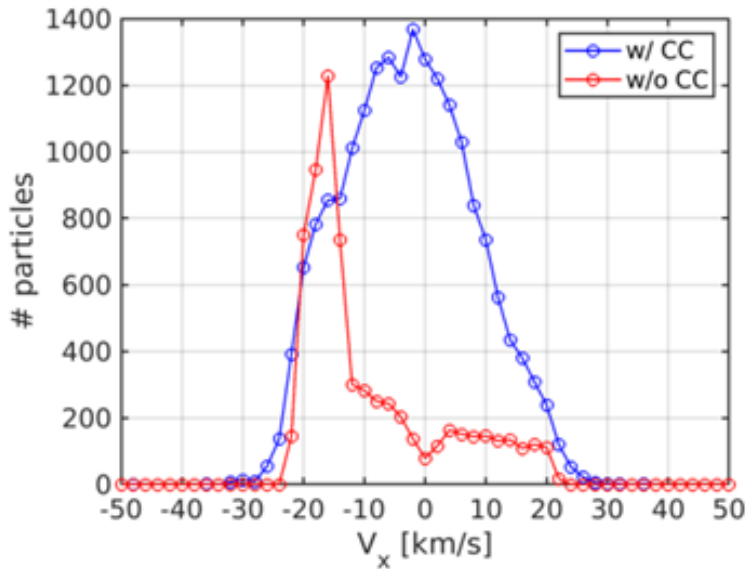
\*\* W. Eckstein and J. P. Biersack, *Appl. Phys. A* **38**, 123 (1985)



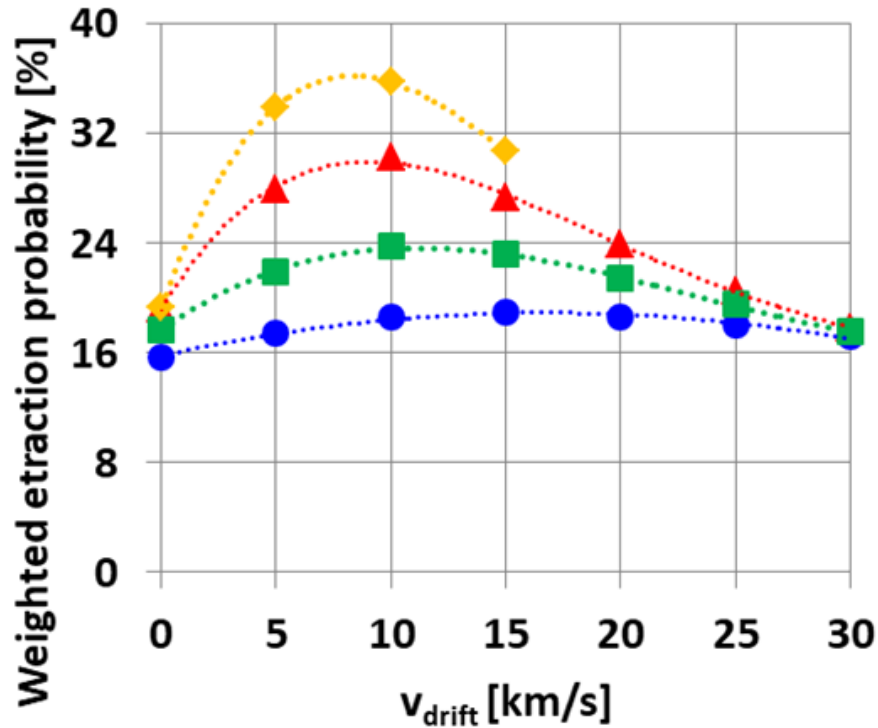
E. Sartori et al., Fusion Engineering and Design Volume 169, 2021)



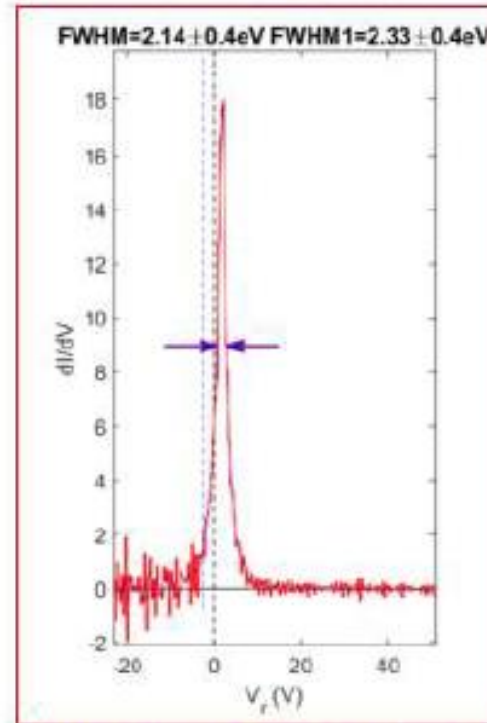
Hydrogen operation  
 $P_{RF} = 23 \text{ kW/driver}$   
 $p_{fill} = 0.36 \text{ Pa}$   
 $B_{filter} = 1.7 \text{ mT}$







●  $T=1.6$  eV    ▲  $T=0.4$  eV    ■  $T=0.8$  eV    ◆  $T=0.2$  eV

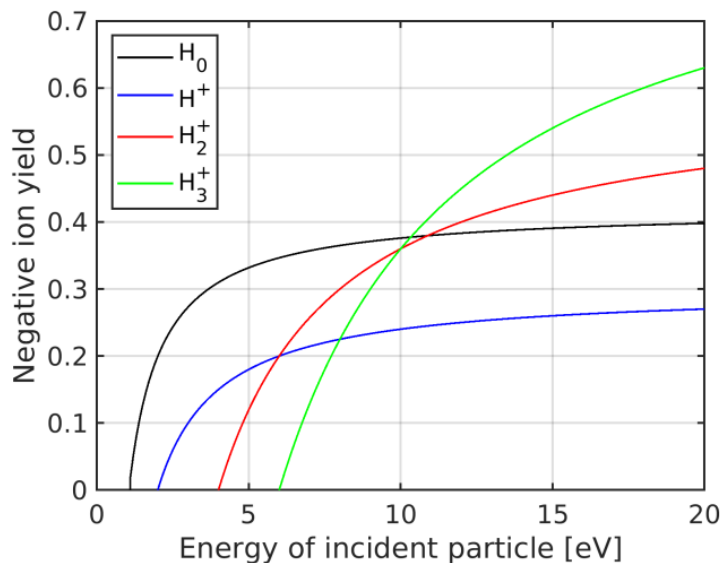


Width of Distribution

1-2 eV

Arc Source (R-NITS, @NIFS)

RFEA in arc source



$$Y(E_{in}) = R_N \eta_0 \left( 1 - \frac{E_{th}/R_E}{E_{in}} \right)$$

Parameters from:

M. Seidl et al. *Journal of Applied Physics* **79**,  
2896 (1996)

By scaling probe measurements at SPIDER towards 70 kW/0.3 Pa:

Typical positive ion flux onto the PG:  $3\text{-}4 \cdot 10^{21} \text{m}^{-3}$ .

Maxwellian flux of H<sup>0</sup>:  $0.25 n_{\text{H}0,1} (8kT_{\text{H}0}/3.14m_{\text{H}})^{0.5} = 1.4 \cdot 10^{22} \text{m}^{-3}$

The majority of the H<sup>-</sup> is obtained from H<sup>0</sup>.