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Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044
Gieseke, Kirchgaesser, Plätzer, Siodmok – arXiv:1808.06770
De Angelis, Forshaw, Plätzer – arXiv:181y.xxxxx
Forshaw, Holguin, Plätzer – arXiv:181y.xxxxx

More on soft gluons beyond leading colour

Simon Plätzer

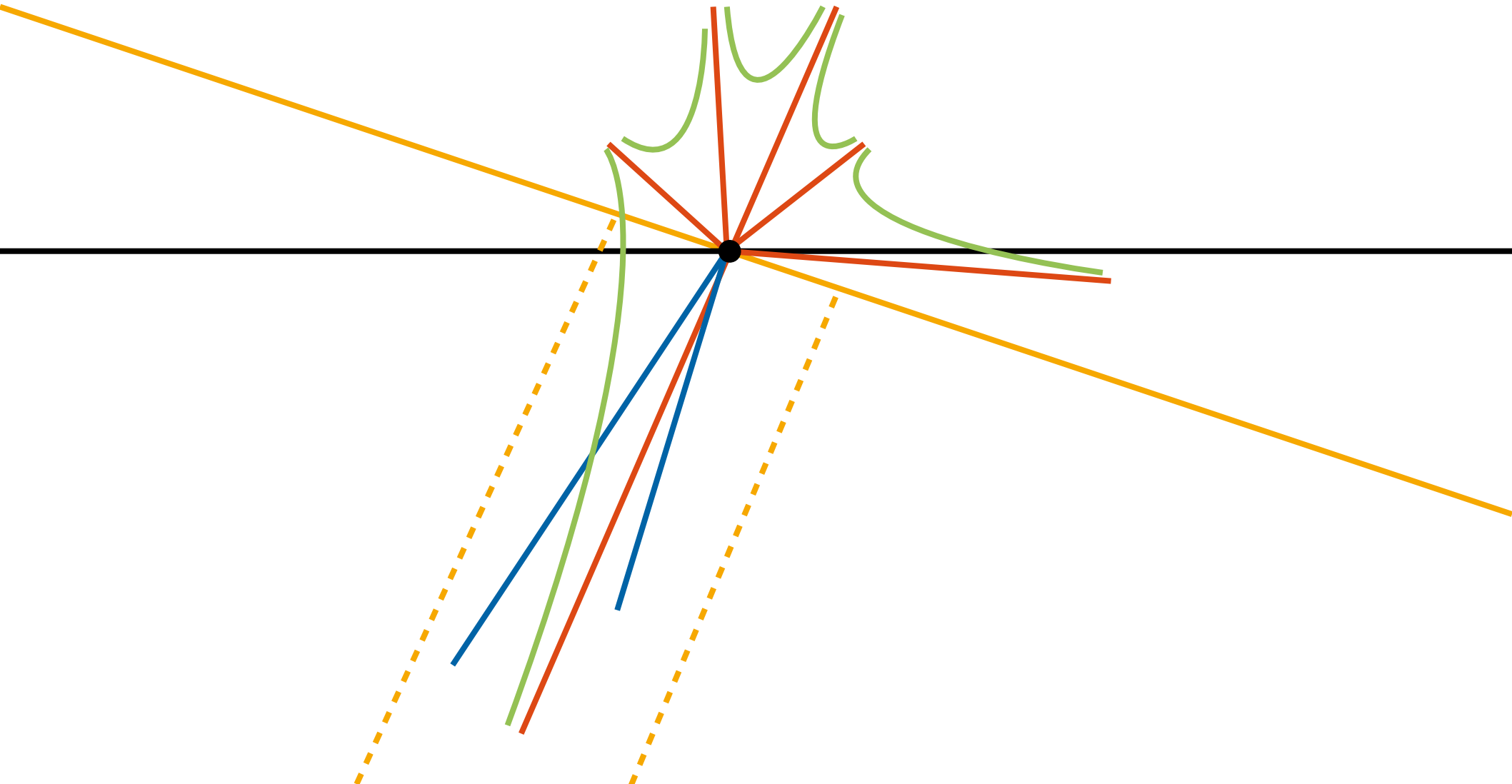
Particle Physics, University of Vienna

at the
HARPS Meeting
Genova | 13 September 2018



Resummation of non-global observables is
with dipole cascades in the large-N limit.

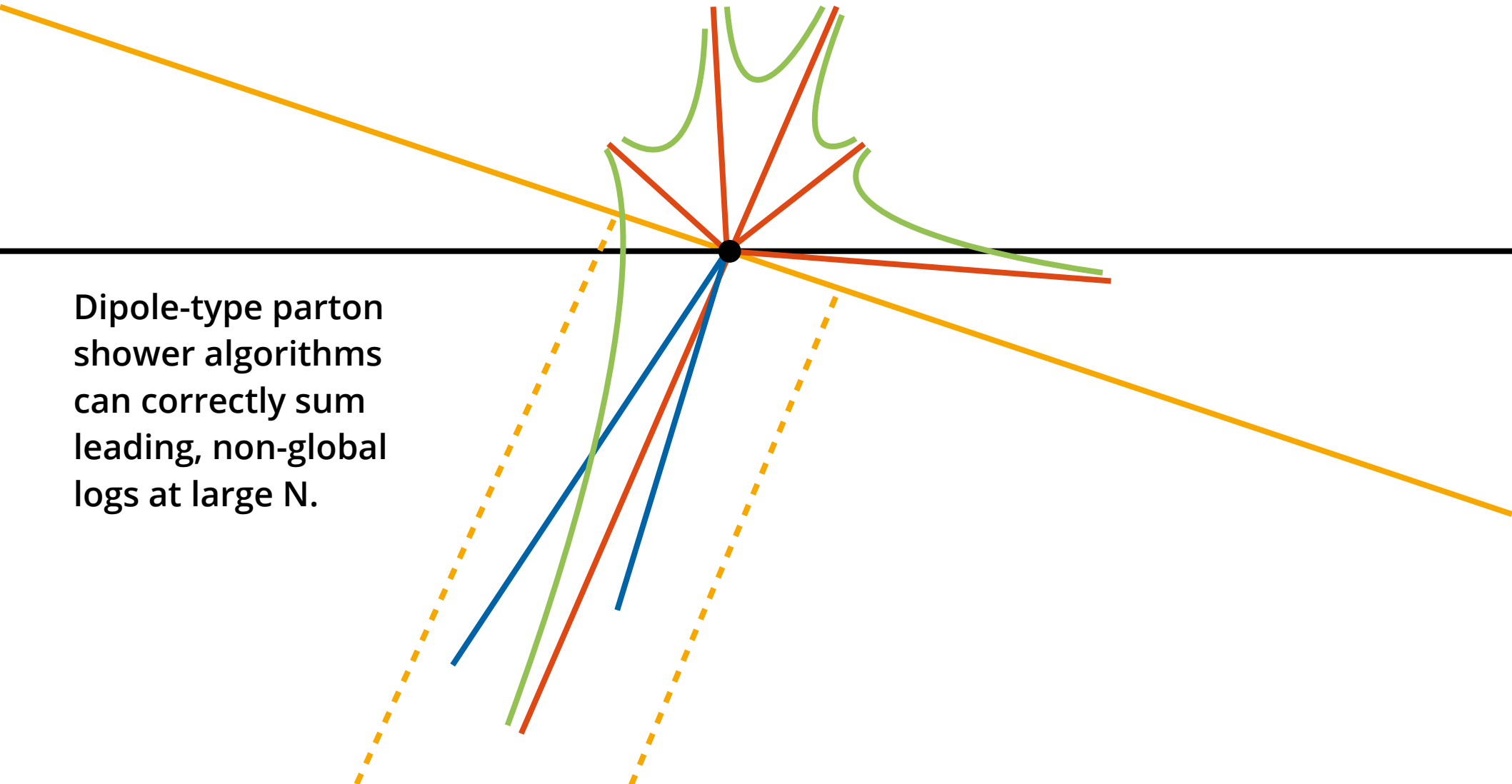
[Dasgupta, Salam – Phys.Lett. B512 (2001) 323]
[Balsinger, Becher, Shao – arXiv:1803.07045]



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Dipole-type parton
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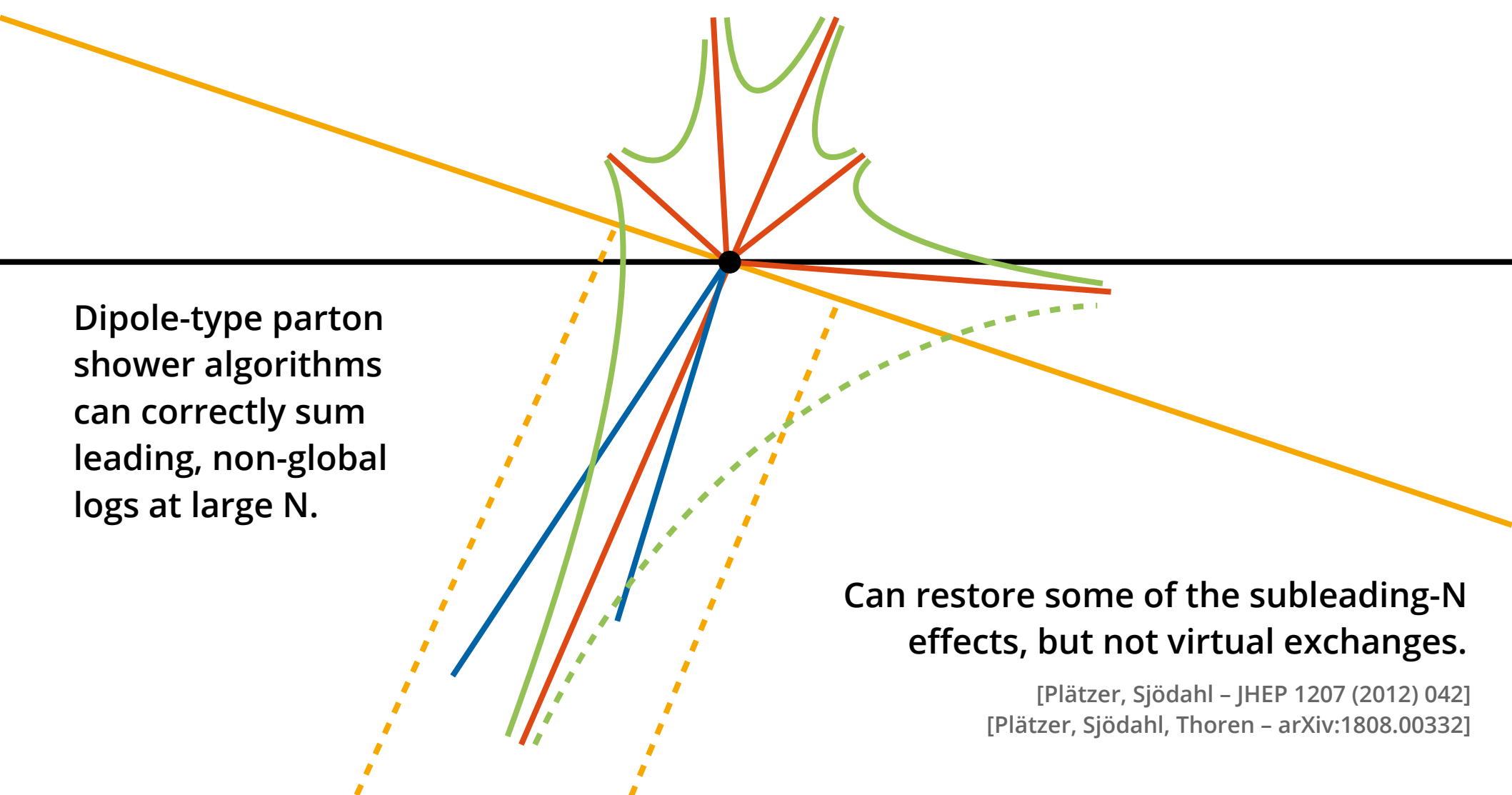
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Dipole-type parton shower algorithms can correctly sum leading, non-global logs at large N .

Can restore some of the subleading- N effects, but not virtual exchanges.

[Plätzer, Sjödal – JHEP 1207 (2012) 042]
[Plätzer, Sjödal, Thoren – arXiv:1808.00332]



Seek a framework to systematically address and improve **a new kind of parton shower algorithm**, not relying on ad-hoc constructions, treating colour exactly as far as possible.

[also see Nagy, Soper]

Non-global observables are a unique playground: At large-N they provide a clean way of deriving a dipole-type parton shower, but the origin of the method used is much more general.

$$\sigma = \sum_n \int \text{Tr} [\mathbf{A}_n(\mu)] u(p_1, \dots, p_n) d\phi_n$$

$$\mathbf{A}_n(\mu) = |\mathcal{M}_n(\mu)\rangle\langle\mathcal{M}_n(\mu)|$$

Evolved `density operator`

Observable

Phase space

[Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044]

Successive virtual evolution/emission combinations down to an infrared cutoff, which will need to be removed at the end. Observable value itself can act as a cutoff, if fully inclusive for gluon energies below this scale.

$$\mathbf{A}_n(E) = \mathbf{V}(E, E_n) \mathbf{D}_n \mathbf{A}_{n-1}(E_n) \mathbf{D}_n^\dagger \mathbf{V}^\dagger(E, E_n) \theta(E - E_n)$$

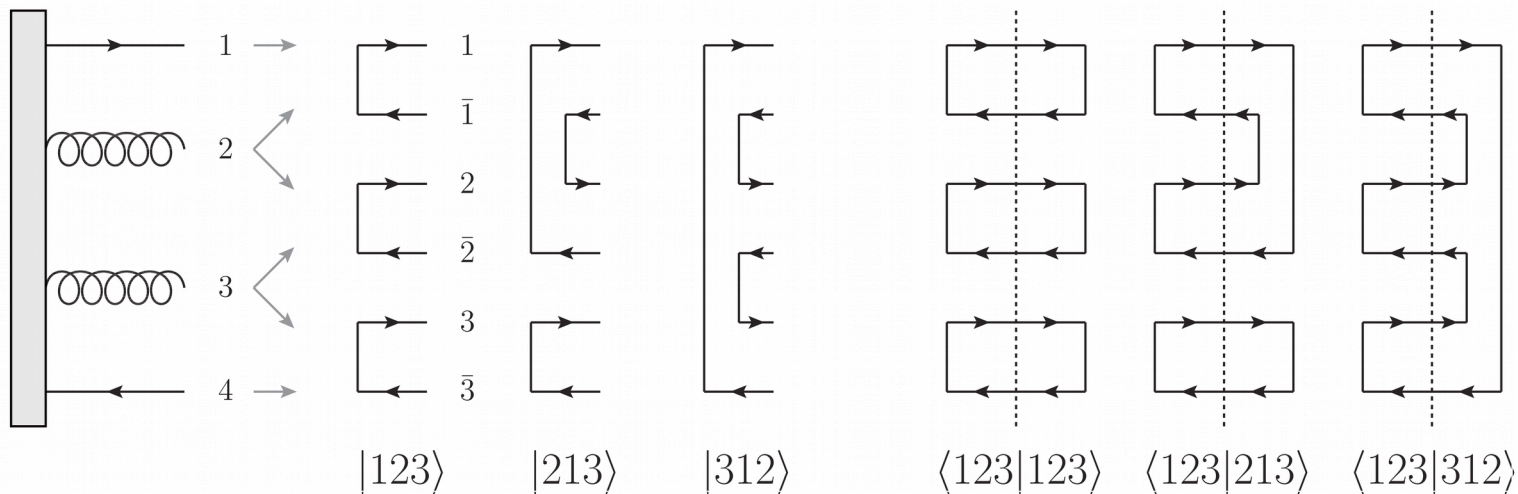
$$\mathbf{V}_n(E, Q) = \text{P exp} \left(- \int_E^Q \frac{dq}{q} \mathbf{\Gamma}_n(q) \right) \quad \mathbf{D}_n = \sum_{i=1}^{n-1} \frac{p_i \cdot \epsilon^*(p_n)}{p_i \cdot p_n} \mathbf{T}_i$$

Non-probabilistic evolution at the amplitude level, keeping full colour structure, virtual corrections encoded in anomalous dimension

$$\mathbf{\Gamma}_n = \frac{\alpha_s}{\pi} \sum_{i < j} \int d\Omega \frac{n_i \cdot n_j}{n_i \cdot n \ n \cdot n_j} (-\mathbf{T}_i \cdot \mathbf{T}_j)$$

Express amplitudes in combinations of fundamental/anti-fundamental indices:

$$|\mathcal{M}\rangle = \sum_{\sigma} \mathcal{M}_{\sigma} |\sigma\rangle \quad |\sigma\rangle = \left| \begin{array}{ccc} 1 & \cdots & n \\ \sigma(1) & \cdots & \sigma(n) \end{array} \right\rangle = \delta_{\alpha_{\sigma(1)}}^{\alpha_1} \cdots \delta_{\alpha_{\sigma(n)}}^{\alpha_n}$$



Non-orthogonal basis:

$$1 = \sum_{\sigma} |\sigma\rangle \langle \sigma| \quad [\sigma|\tau\rangle = \langle \sigma|\tau\rangle = \delta_{\tau\sigma} \quad \text{Tr}[\mathbf{A}] = \sum_{\tau, \sigma} [\tau|\mathbf{A}|\sigma\rangle \langle \sigma|\tau\rangle$$

Also overcomplete ... **but computationally very handy: It's all about permutations.**

Evolution operator in colour flow basis:

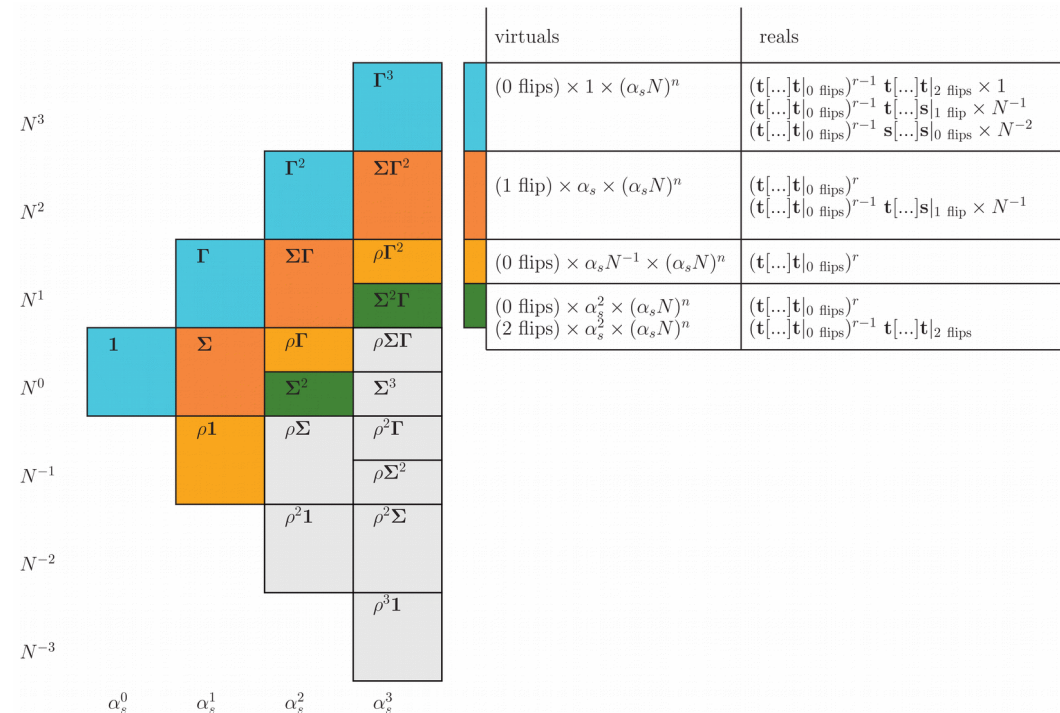
$$\begin{aligned}
 [\tau|e^{\Gamma}|\sigma\rangle &= \sum_{k=0}^{\infty} \frac{(-1)^k}{N^k} \sum_{l=k}^{\infty} \frac{(-\rho)^l}{l!} \sum_{\sigma_0, \dots, \sigma_{l-k}} \delta_{\tau\sigma_0} \delta_{\sigma_{l-k}\sigma} \prod_{\alpha=0}^{l-k-1} \Sigma_{\sigma_{\alpha}\sigma_{\alpha+1}} R(\{\sigma_0, \dots, \sigma_{l-k}\}) \\
 &= \delta_{\tau\sigma} \left(e^{-N\Gamma_{\sigma}} + e^{-N\Gamma_{\sigma}} \frac{\rho}{N} \right) - \frac{1}{N} \frac{e^{-N\Gamma_{\tau}} - e^{-N\Gamma_{\sigma}}}{\Gamma_{\tau} - \Gamma_{\sigma}} \Sigma_{\tau\sigma} + \text{NNLC}
 \end{aligned}$$

[Plätzer - EPJ C 74 (2014) 2907]

Sum terms enhanced by $\alpha_s N$ to all orders, insert perturbations in $1/N$.

Take into account real emission contributions and the final suppression by the scalar product matrix element.

$$\text{Tr}[\mathbf{A}] = \sum_{\tau\sigma} [\tau|\mathbf{A}|\sigma\rangle\langle\sigma|\tau\rangle \quad \langle\sigma|\tau\rangle = N^{n-\#(\tau,\sigma)}$$



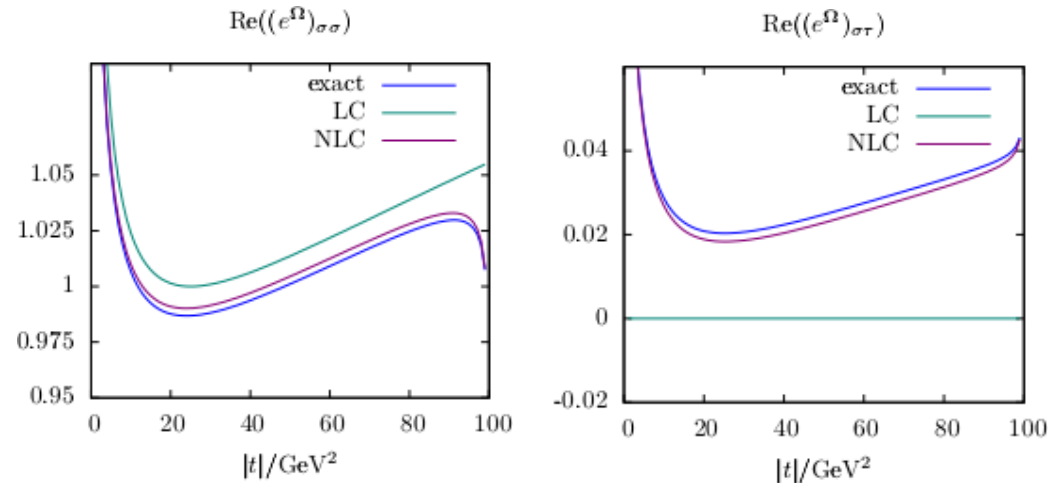
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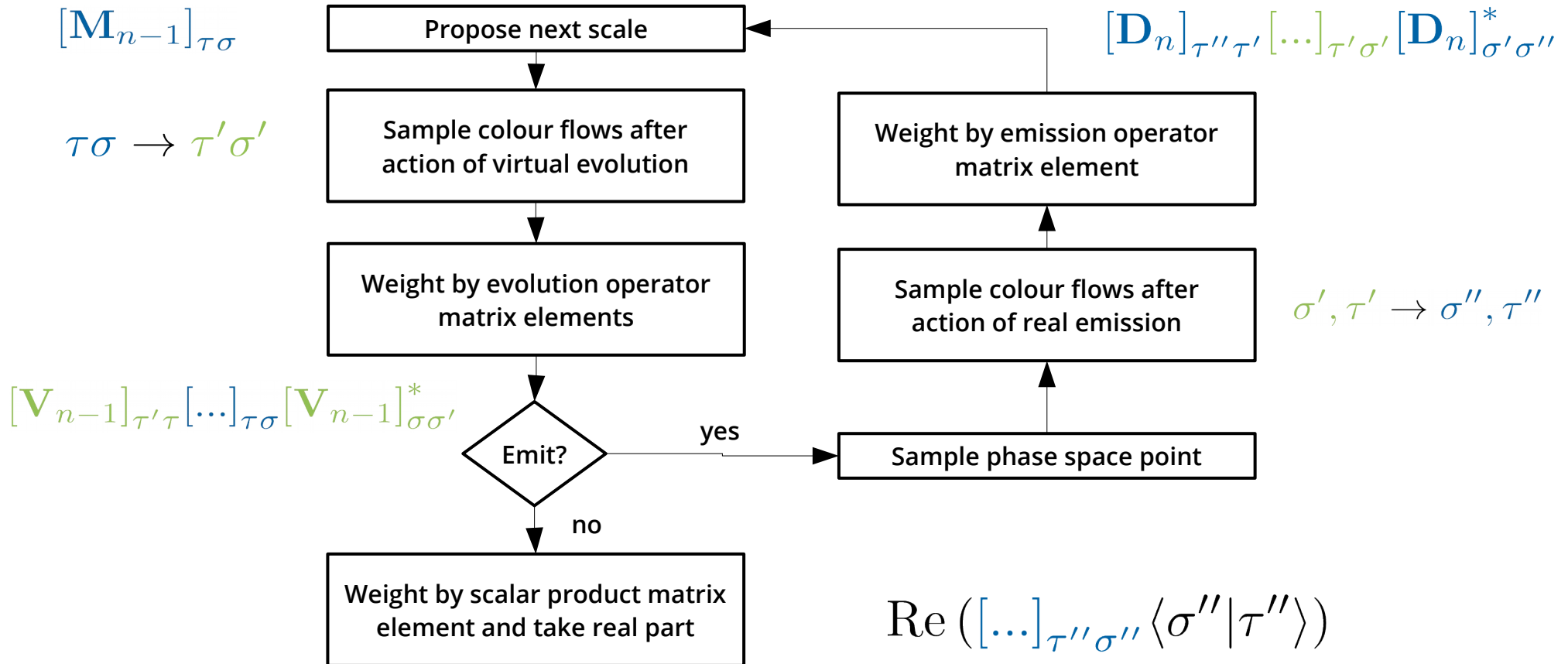
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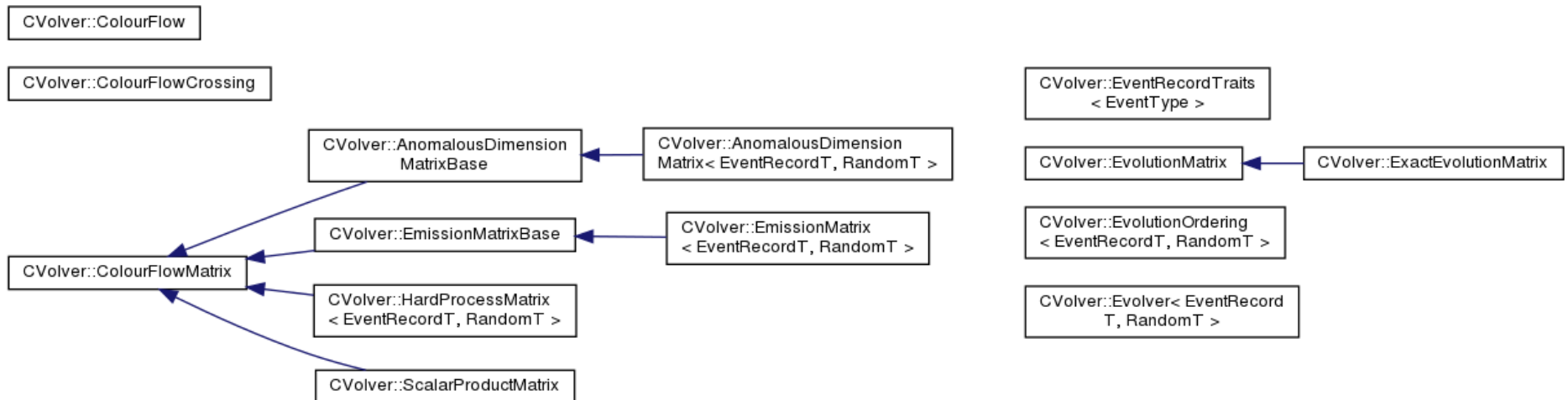
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A framework to solve multi-differential evolution equations in colour space.
Concise, simple, and light-weight code structure.

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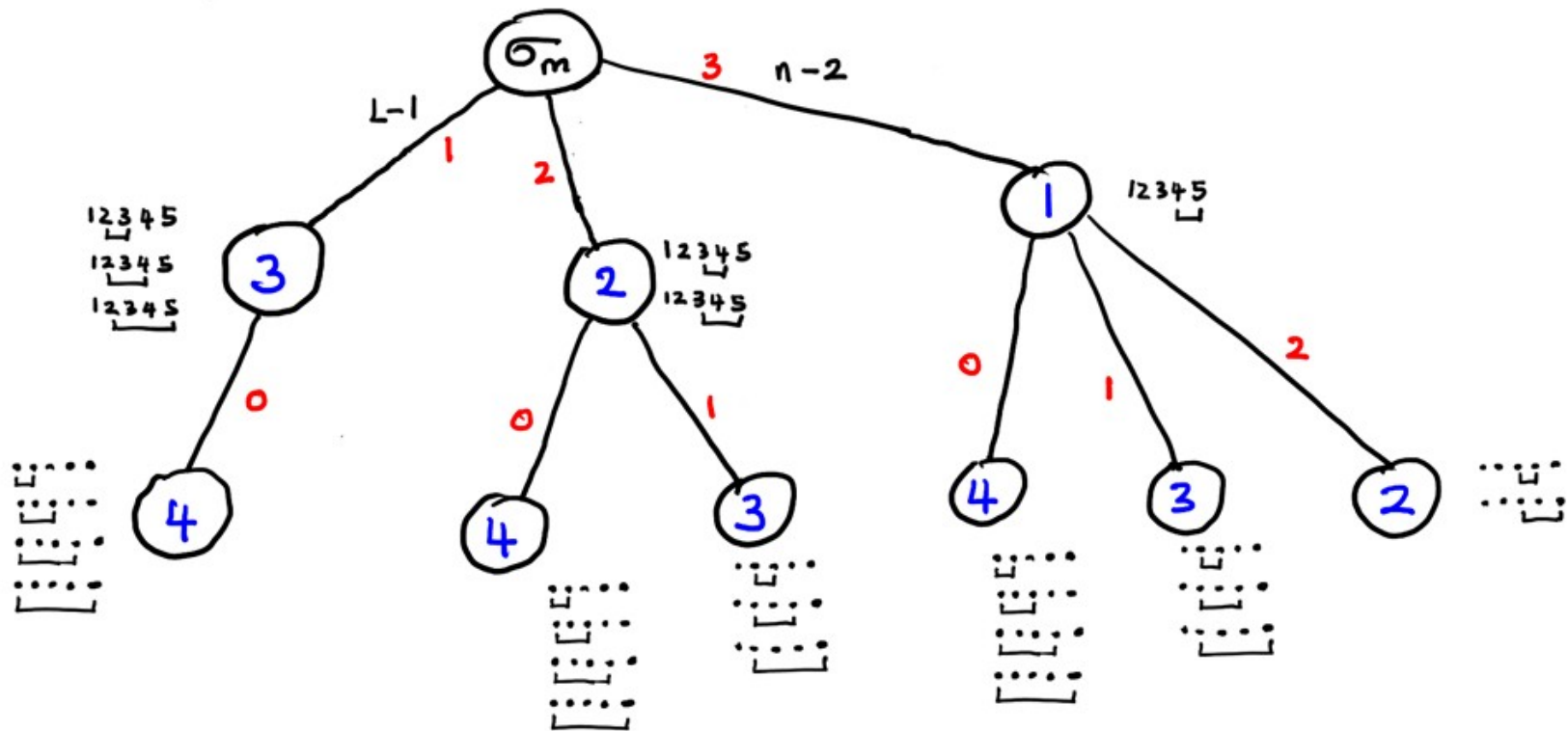


Dedicated Monte Carlo algorithms to sample colour structures.

Plugin approach can accommodate anything from (N)GLs to full parton showers.

Importance sampling in colour space rules: $\#(\tau, \tau') \sim 1/N^{\#(\tau, \tau')}$

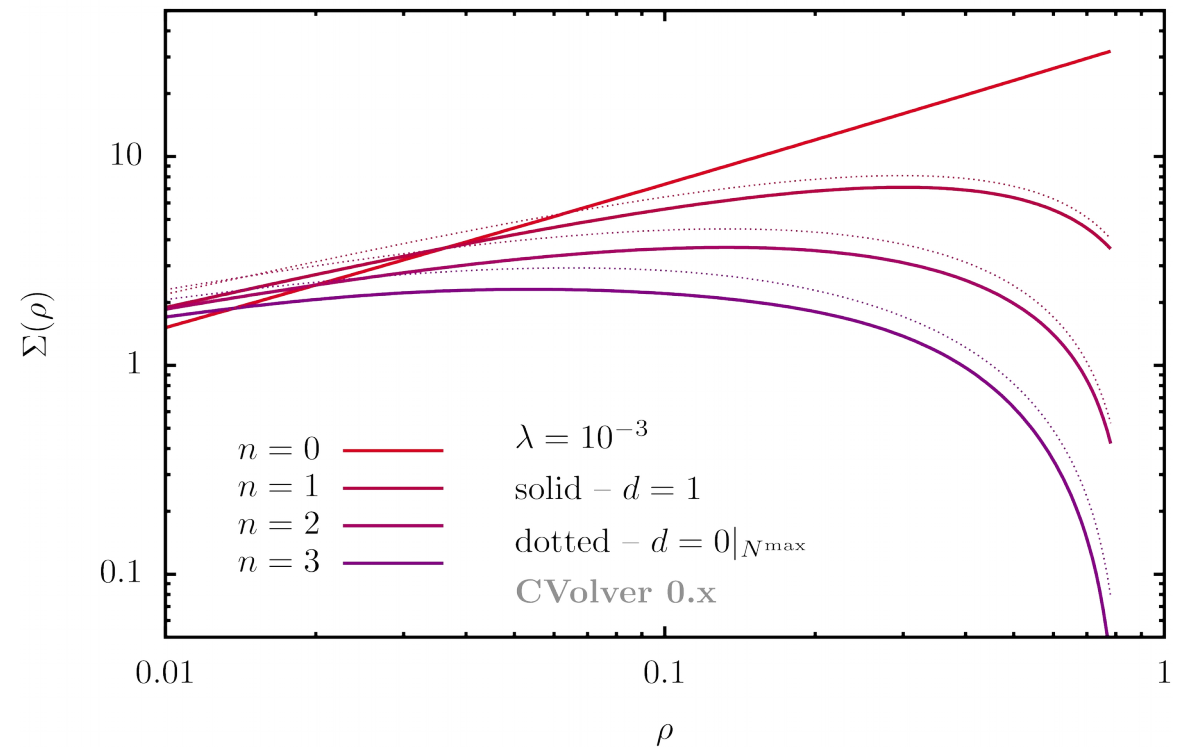
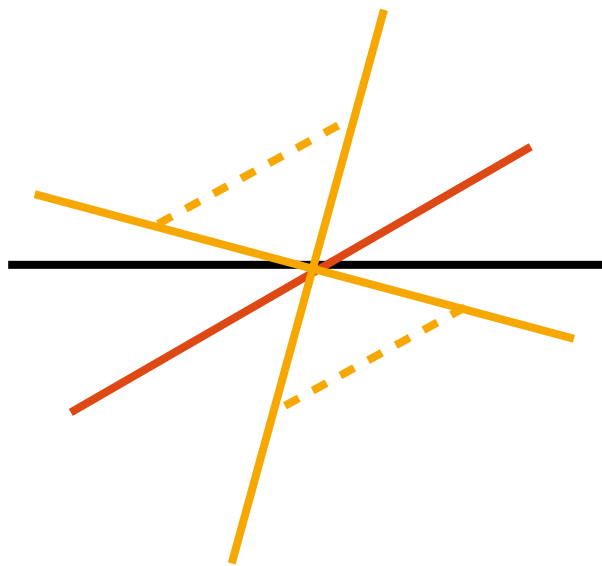
Enumerate and address permutations with fixed cycle length:



[De Angelis, Forshaw, Plätzer – arXiv:181y.xxxxx]

Code is differential for a large class of (non-global) observables.
Example: Cone-dijet veto cross section.

dijet veto, $\delta = \pi/4$

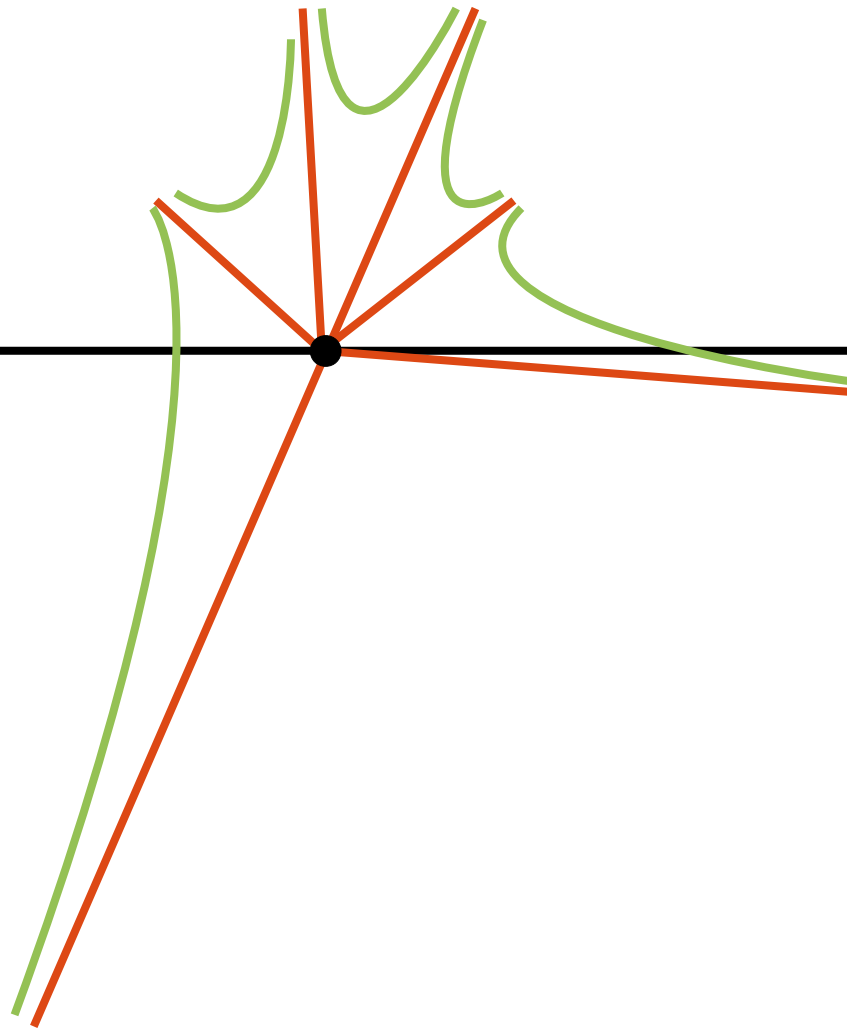


1/N breakdowns possible, scales up to several 10s of emissions for $d=2$.

The cluster model is based on a single colour flow after the shower has stopped.

Essentially no evolution is considered, just decays.

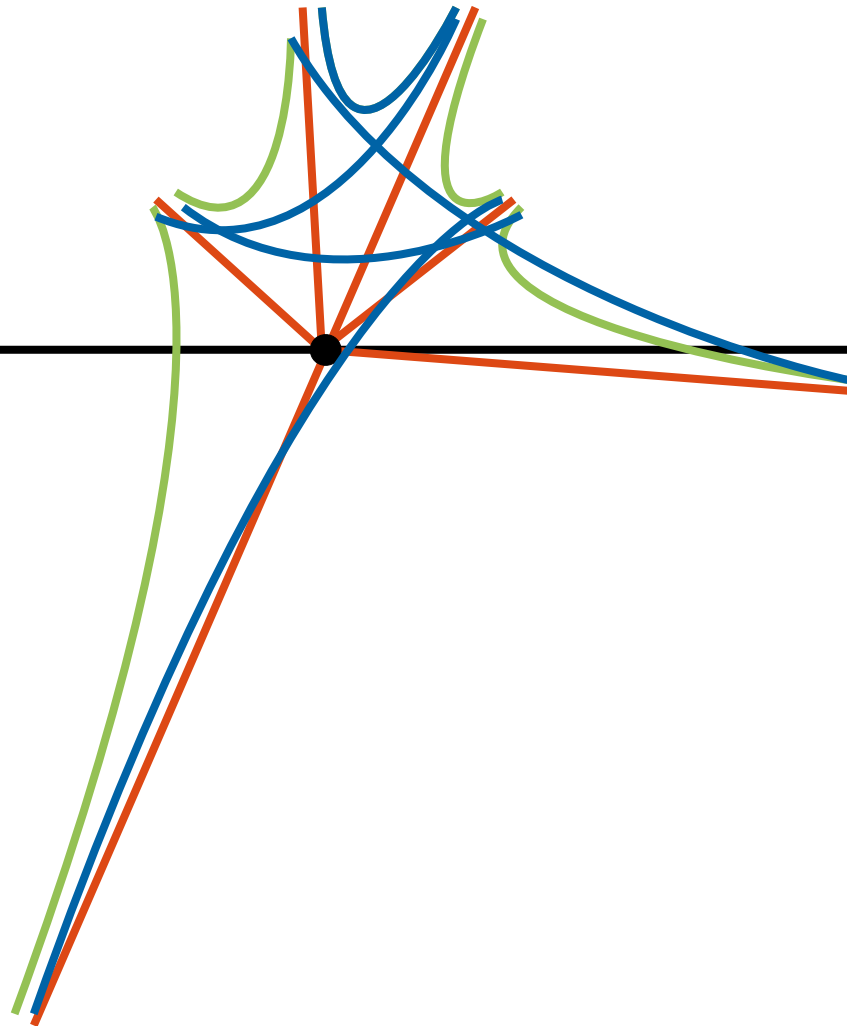
[Gieseke, Kirchgaesser, Plätzer, Siodmok – arXiv:1808.06770]



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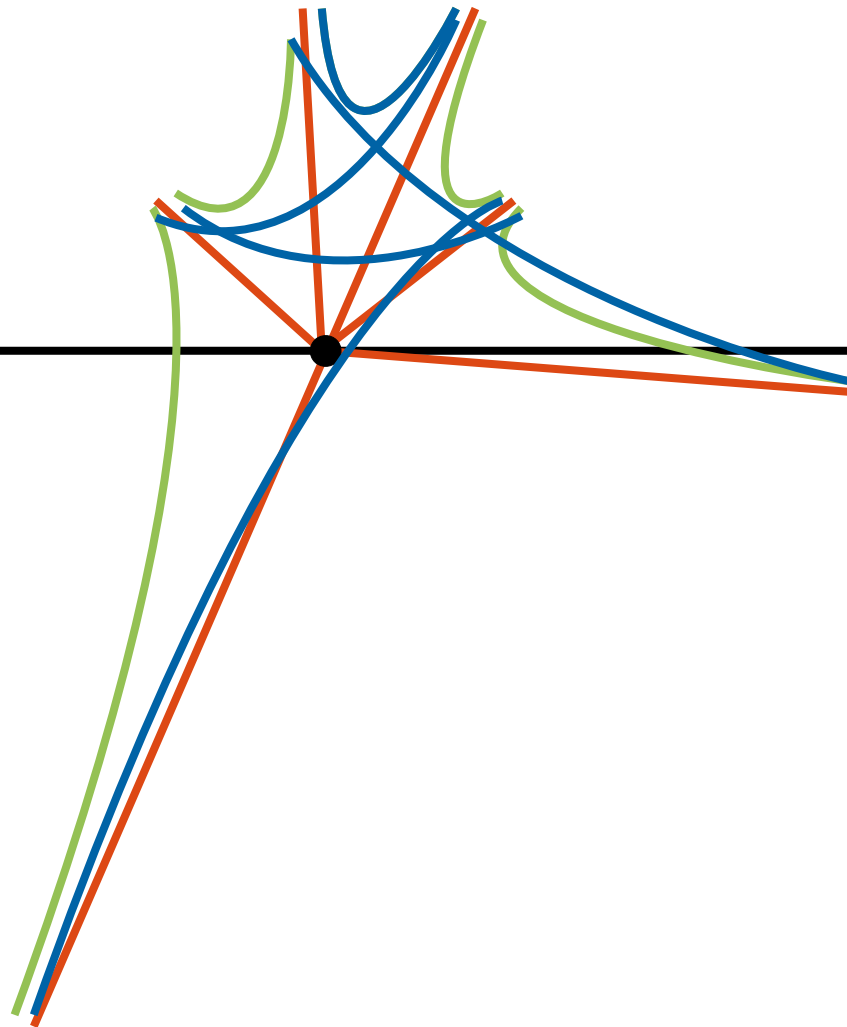
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View as an evolving amplitude, driven by a single initial colour flow:

$$|\mathcal{M}\rangle = e^{\Gamma} |\text{clusters}\rangle$$

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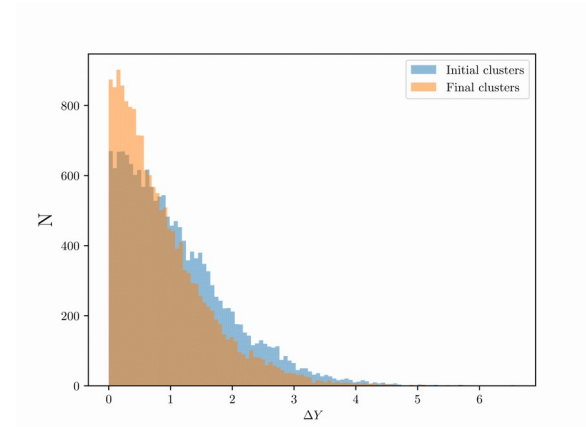
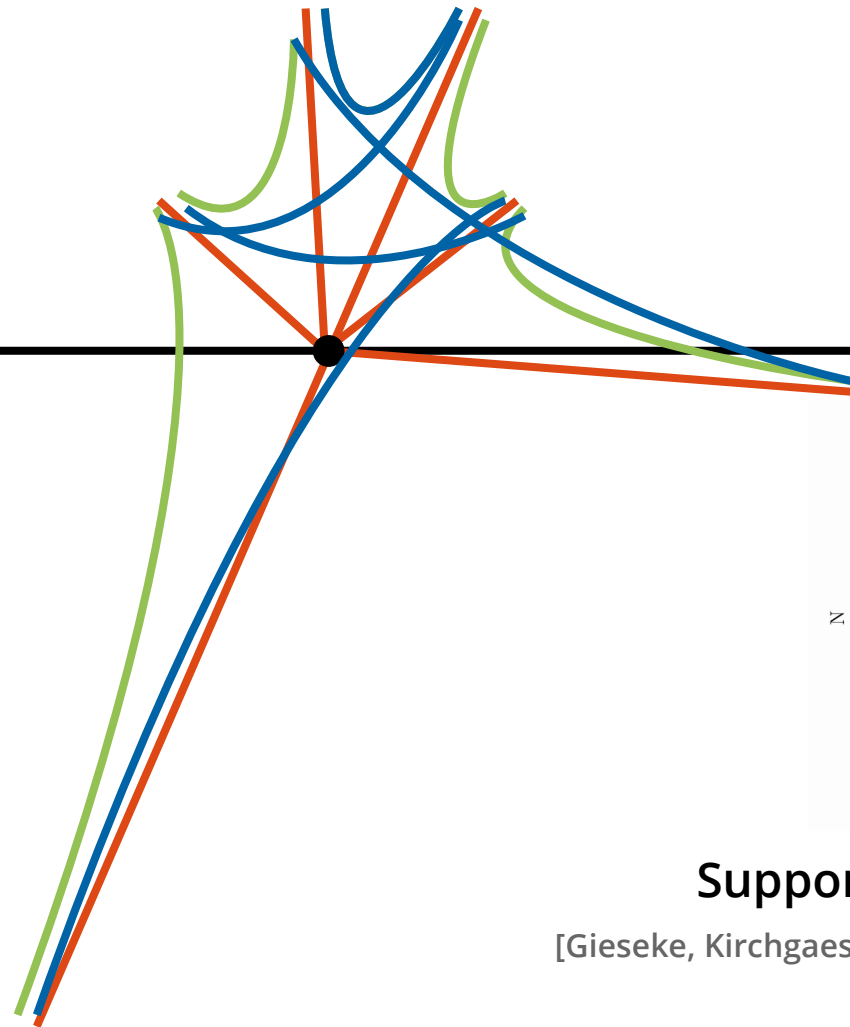
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Supports geometric models

[Gieseke, Kirchgaesser, Plätzer – EPJ C 78 (2018) 99]

A framework for amplitude-level evolution at full colour, including a Monte Carlo method to produce actual `events`.

Serves to host a new approach to parton shower algorithms, which are systematically improvable, and first steps in that direction are underway.

Resummation of non-global observables requires the very same methods, and so provide an ideal playground to quantify these effects in a controlled way.

Details of colour expansion under study.

Beyond the perturbative evolution this provides us with unique insight into colour reconnection models.