

# Strong coupling determination with grooming

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In Collaboration with:

Jeremy Baron and Simone Marzani [arXiv:1803.04719] and

Simone Marzani, Daniel Reichelt, Steffen Schumann, Gregory Soyez [in Progress]

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# Jet substructure

- Many jet substructure techniques developed; **Grooming** and Tagging
- Created with the purpose of distinguishing signal from background
- Removes soft wide-angle radiation
- Can also help reduce non-perturbative corrections

# mMDT & Soft drop

Main technique we will deal with is soft drop:

[Larkoski, Marzani, Soyez, Thaler; '14]

$$\frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_c \left( \frac{\Delta R_{12}}{R} \right)^\beta$$

or at  $e^+e^-$  colliders:

$$\frac{\min[E_i, E_j]}{E_i + E_j} > z_{\text{cut}} (1 - \cos \theta_{ij})^{\beta/2}$$

Makes use of Cambridge/Aachen clustering.

[Dokshitzer, Leder, Moretti, Webber; '97][Wobisch, Wengler; '99]

Reduces to modified Mass Drop Tagger (mMDT) for  $\beta = 0$

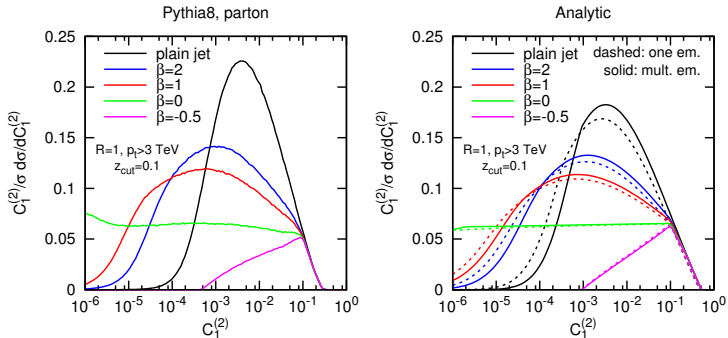
[Dasgupta, Fregoso, Marzani, Salam; '13]

# Need for resummation

For boosted jets, great separation of scales  $p_T \gg m$  leads to large logarithms:

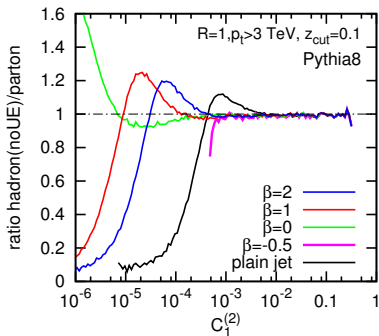
$$\log \left( \frac{m_J^2}{p_T^2 R^2} \right)$$

Large logarithms need to be resummed.



[Larkoski, Marzani, Soyez, Thaler; '14]

# Reduction in NP corrections



[Larkoski, Marzani, Soyez, Thaler; '14]

# Further work

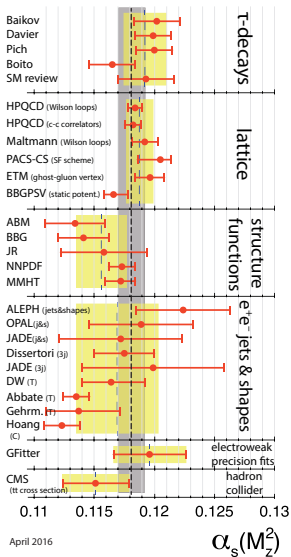
- Computation of soft drop using SCET at NNLL accuracy approximated for  $e_2^{(2)} \ll z_{cut}$  [Frye, Larkoski, Schwartz, Yan; '16]
- Calculation in dQCD including finite  $z_{cut}$  effects [Marzani, Schunk, Soyez; '17]
- Including jet radius resummation in SCET [Kang, Lee, Liu, Ringer; '18]
- Good agreement to experiments [CMS;'17] [ATLAS;'17]
- Application to top quark mass measurements [Hoang, Mantry, Pathak, Stewart;'17]
- And  $\alpha_s$  measurements at LHC [Les Houches;'18] and  $e^+e^-$  [Baron, Marzani, VT; '18]

# The importance of $\alpha_s$

- Jet physics of great importance to the LHC
- Higher order perturbative corrections shown to be important scale with higher powers of  $\alpha_s$
- Higgs boson production scales as  $\alpha_s^2$

An accurate measurement of  $\alpha_s$  is necessary for precision LHC measurements

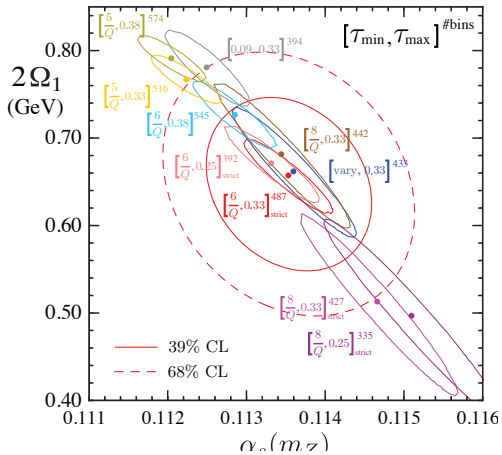
# $\alpha_S$ Measurement



[Particle Data Group; 16]



# NP contributions



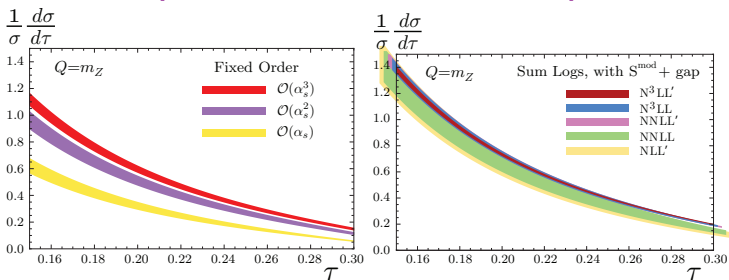
[Abbate, Fickinger, Hoang, Mateu, Stewart; '10]

# Thrust

$$\tau = 1 - T = \min_{\vec{n}} \left( 1 - \frac{\sum_i |\vec{n} \cdot \vec{p}_i|}{\sum_i |\vec{p}_i|} \right)$$

Minimize for thrust axis  $\vec{n}$

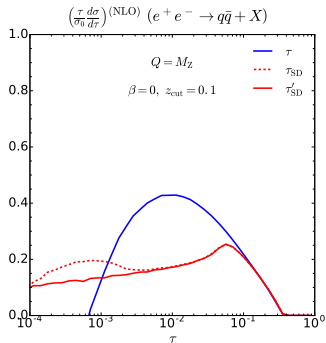
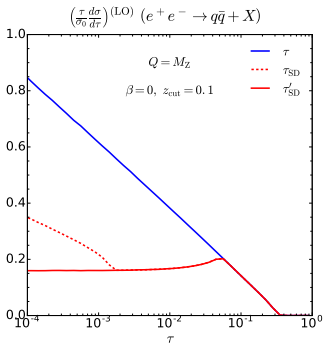
[Abbate, Fickinger, Hoang, Mateu, Stewart; '10]



## SD Distribution

Hemisphere jets at an  $e^+e^-$  collider  $\rightarrow$  Different soft drop condition:

$$\frac{\min[E_i, E_j]}{E_i + E_j} > z_{\text{cut}} (1 - \cos \theta_{ij})^{\beta/2}$$



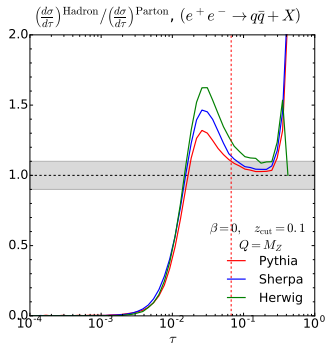
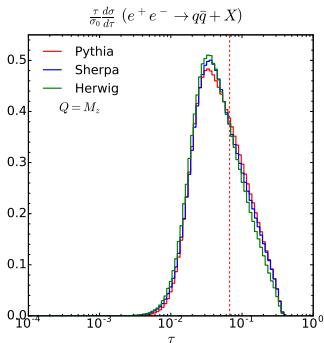
# Alternative definition

- Separation into two jets at the hand of thrust axis pre-softdrop
- After softdrop each hemisphere will have its own axis
- Each thrust axis is the jet axis

$$T'_{\text{SD}} = \frac{\sum_{i \in \mathcal{H}_{\text{SD}}^L} |\vec{n}_L \cdot \vec{p}_i|}{\sum_{i \in \mathcal{E}_{\text{SD}}} |\vec{p}_i|} + \frac{\sum_{i \in \mathcal{H}_{\text{SD}}^R} |\vec{n}_R \cdot \vec{p}_i|}{\sum_{i \in \mathcal{E}_{\text{SD}}} |\vec{p}_i|}$$

# MC studies

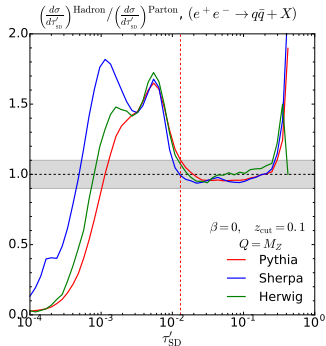
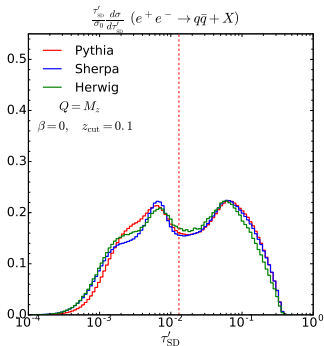
[Baron, Marzani, VT; '18]



Non-perturbative corrections above 10% around  $\tau \simeq 0.07$

# MC studies with soft drop

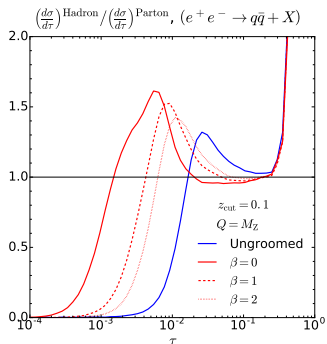
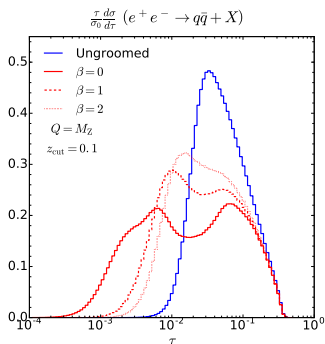
[Baron, Marzani, VT; '18]



Non-perturbative corrections above 10% around  $\tau \simeq 0.001$   
 Reduction in non perturbative corrections.

# $z_{\text{cut}}$ & $\beta$ values

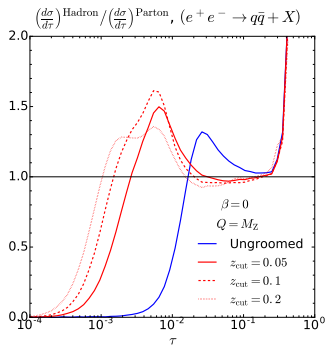
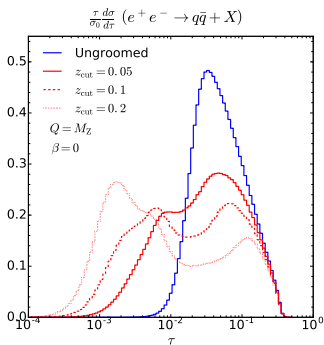
[Baron, Marzani, VT; '18]



Different values of  $\beta$  do not offer improvement

# $z_{\text{cut}}$ & $\beta$ values

[Baron, Marzani, VT; '18]



Smaller values of  $z_{\text{cut}}$  offer more data in the relevant region with only a slight increase in non-perturbative corrections.

For more observables [talk Jeremy Baron]



# Factorization

Factorization for  $\tau \ll z_{\text{cut}} \ll 1$  [Frye, Larkoski, Schwartz, Yan; '16]:

$$\frac{d\sigma}{d\tau} = H(Q) S_G(z_{\text{cut}}, \beta) [S_C(\tau, z_{\text{cut}}, \beta) \otimes J(\tau)]^2$$

Computed in Laplace space and inverted leading to:

$$\Sigma(\tau) = \left[ 1 + \left( \frac{\alpha_s}{\pi} \right) C^{(1)} + \dots \right] \exp \left[ \frac{1}{\alpha_s} g_1(-\lambda_\tau, \lambda_{z_{\text{cut}}}) + g_2(-\lambda_\tau, \lambda_{z_{\text{cut}}}) + \dots \right]$$

for  $\lambda_x = \alpha_s b_0 \log x$  and confirmed using dQCD.

With matching:

$$\tau \frac{d\sigma^{\text{LO+NLL}'}}{d\tau} = \tau \frac{d\sigma^{\text{LO}}}{d\tau} + \left[ \tau \frac{d\sigma^{\text{NLL}'}}{d\tau} - \tau \frac{d\sigma^{\text{NLL}'|_{\text{LO}}}}{d\tau} \right]$$

# Analytic computation

Additional calculation for contributions where  $\tau \sim z_{\text{cut}}$  at NLL' accuracy:

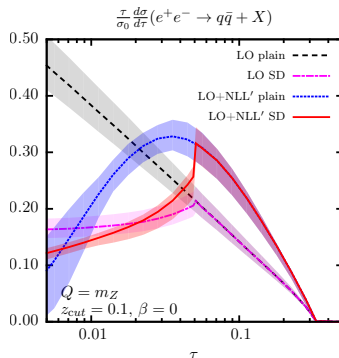
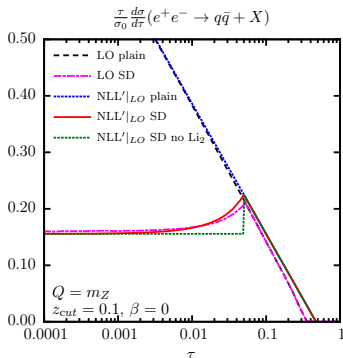
$$\frac{\alpha_s}{\pi} C_F (\beta + 2) \text{Li}_2 \left[ \frac{1}{2} \left( \frac{2\tau}{z_{\text{cut}}} \right)^{\frac{2}{\beta+2}} \right]$$

Can be neglected for  $\tau \ll z_{\text{cut}}$ , but offers a constant contribution near the transition point  $\tau = z_{\text{cut}}/2$ .

Additional corrections for the end-point of the resummation and expansion.

# Resummation results

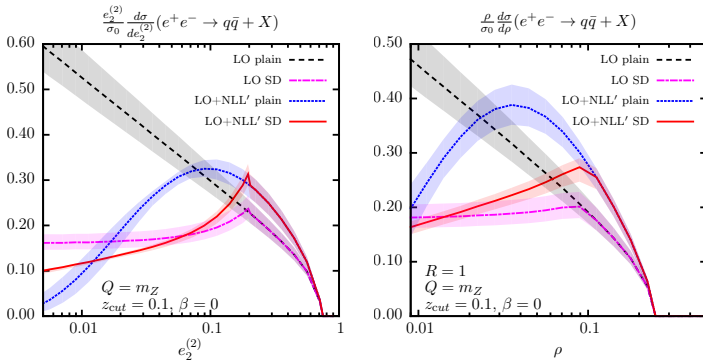
[Baron, Marzani, VT; '18]



- Expansion offers a good approximation for fixed order
- Transition corrections are important for thrust

# Alternative observables

[Baron, Marzani, VT; '18]



Other observables allow for a reduction in transition point effects.

# Higher accuracy for fit

Some changes in the setup:

- Matching to NLO
- End point corrections from plain thrust applied over full range
- Additional transition point effect from multiple emissions

$$\Sigma_{\text{res}}|_{\tau > z_{\text{cut}}/2} = C \frac{e^{R(\tau) + \gamma_E R'(\tau)}}{\Gamma(1 - R'(\tau))} \exp \left[ R'(\tau) \frac{z_{\text{cut}}}{2\tau} {}_3F_2 \left( 1, 1, 1 + R'(\tau); 2, 2; \frac{z_{\text{cut}}}{2\tau} \right) \right] \\ \times \exp \left[ \log \left( \frac{2\tau}{z_{\text{cut}}} \right) \{ R''(\tau) - R''(\tau = z_{\text{cut}}/2, z_{\text{cut}}) \} \right]$$

# End point corrections

[Catani, Trentadue, Turnock, Webber; '93] [Jones, Ford, Salam, Stenzel, Wicke; '03]

Modification of the logarithm:

$$\log(x_L \tau) \rightarrow -\frac{1}{p} \log \left( \frac{1}{(x_L \tau)^p} - \frac{1}{(x_L \tau_{\max})^p} + 1 \right)$$

And Resummation:

$$\Sigma = C \exp \left[ \tilde{R}(\tau) - \frac{\tau}{\tau_{\max}} \tilde{R}'(\tau \rightarrow \tau_{\max}) \log \bar{\tau} \right]$$

Ensures that  $\frac{d\sigma}{d\tau}(\tau_{\max}) = 0$  for resummation and expansion

# Fitting setup

- Fit to ALEPH data for thrust *[ALEPH;'04]*
- Fit to MEPS@NLO 2-5j Sherpa result for soft drop *[Schumann, Krauss;'07] [Gehrmann, Hoche, Krauss, Schonherr, Siegert;'12]*  
Assume ALEPH uncertainty with  $\sqrt{\sigma}$  scaling
- Minimization of chi-squared
- Neglect any correlation effects in chi-squared
- Using range  $\tau \in [0.06, 0.25]$
- Experimental uncertainty at the hand of  $\Delta\chi^2 = 1$
- Theoretical uncertainty based on 500 random variations:
  - Variation of  $\mu_R$  and  $x_L$  with  $1/2 < \mu_R x_L / Q < 2$
  - Variation of end-point  $p$  variable 1 or 2
  - Switch between matching scheme: Multiplicative, additive and Multiplicative expanded out

Setup is similar to a combination of methods from

*[Abbate, Fickinger, Hoang, Mateu, Stewart;'12]* and *[Gehrmann, Luisoni, Monni;'12]*

# Non-perturbative corrections

For the non-perturbative model use a shift in  $\tau$  and  $z_{\text{cut}}$  from

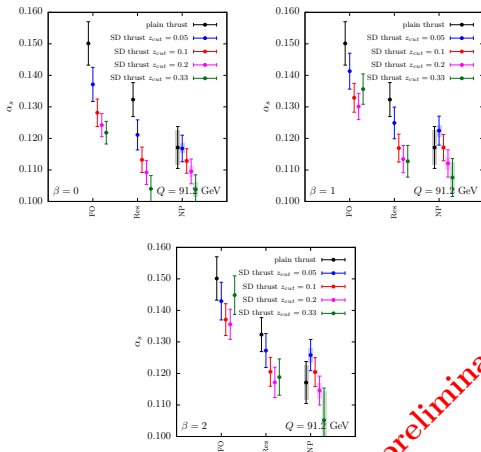
*[Dasgupta, Fregoso, Marzani, Salam;'13] [Marzani, Schunk, Soyez;'17]*

- Shift in  $\tau$  from non-perturbative emission within a cone defined by thrust
- Shift in  $z_{\text{cut}}$  from non-perturbative emission that reduces energy leading to it being groomed
- Both are computed in a  $2 \rightarrow 2 + \text{NP}$  with small  $\tau$  limit configuration
- More detailed calculation can be performed as a better approximation



# Variation of $z_{\text{cut}}$ and $\beta$

[Marzani, Reichelt, Schumann, Soyez, VT; In Preparation]

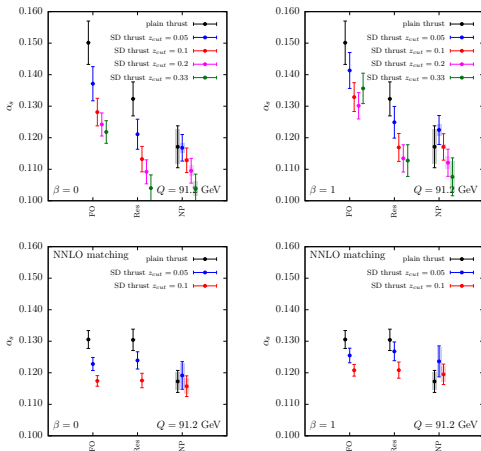


preliminary

Significant reduction in hadronization uncertainty and shift

# NNLO corrections

[Marzani, Reichelt, Schumann, Soyez, VT; In Preparation]



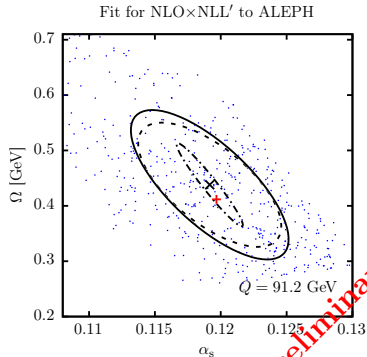
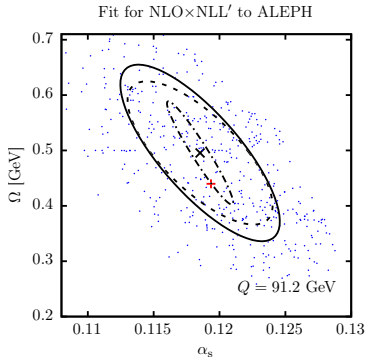
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Using NNLO corrections [Kardos, Somogyi, Trócsányi; '18]

Reduction of uncertainty and stabilization for NLL corrections

# Effect of smearing NP corrections

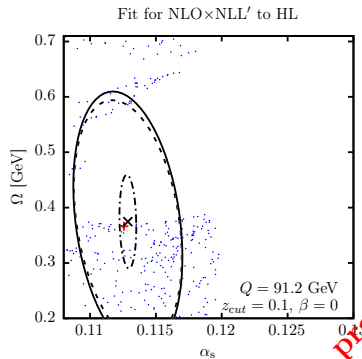
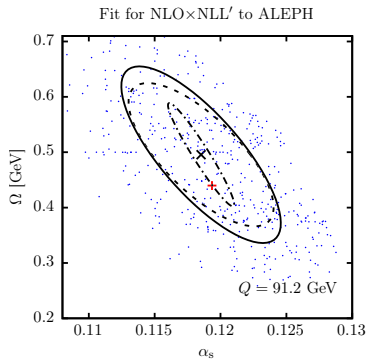
*[Marzani, Reichelt, Schumann, Soyez, VT; In Preparation]*



preliminary

Smearing does not impact fit significantly at this accuracy, helps reduce uncertainty

## Reduction of NP uncertainty

*[Marzani, Reichelt, Schumann, Soyez, VT; In Preparation]*

preliminary

Significant reduction in correlation between  $\alpha_s$  and  $\Omega$

# Fit using MC for hadronization effects

If Parton level agrees well with resummation [*talk Daniel Reichelt*], can make use of hadronization model of Monte Carlo:

- Sherpa MEPS@NLO 2-3j results  
[*Schumann, Krauss;'07*] [*Gehrmann, Hoche, Krauss, Schonherr, Siegert;'12*]
- Sherpa's Cluster model [*Winter, Krauss, Soff;'03*]
- Pythia's Lund string model in Sherpa [*Sjostrand, Mrenna, Skands;'06*]
- Uncertainty given by difference and value is average

For soft drop thrust with  $z_{\text{cut}} = 0.1$  and  $\beta = 0$ :

$$\alpha_s = 0.1163 \pm 0.0008(\text{exp}) \pm 0.0030(\text{had}) \pm 0.0042(\text{th})$$

Compared to:

$$\alpha_s = 0.1128 \pm 0.0007(\text{exp}) \pm 0.0006(\text{had}) \pm 0.0038(\text{th})$$

**preliminary**

# Other observables

This approach can also be applied to heavy hemisphere mass.  
 There is tension for thrust and heavy hemisphere mass without grooming  
*[Chien, Schwartz;'10]*

For soft drop heavy hemisphere mass (for  $\rho_+ \in [0.08, 0.18]$ ):

$$\alpha_s = 0.1159 \pm 0.0060(\text{exp}) \pm 0.0011(\text{had}) \pm 0.0036(\text{th})$$

Compared to thrust:

$$\alpha_s = 0.1163 \pm 0.0008(\text{exp}) \pm 0.0030(\text{had}) \pm 0.0042(\text{th})$$

Leading to a combined fit:

$$\alpha_s = 0.1160 \pm 0.0005(\text{exp}) \pm 0.0019(\text{had}) \pm 0.0039(\text{th})$$

No tension in the two fits

**preliminary**

# Summary

## Conclusions

- Soft drop can help reduce dependence on non-perturbative corrections for thrust
- Can help break degeneracy between non-perturbative contributions and  $\alpha_s$  in fit leading to significant reduction in hadronization uncertainty
- Can help reduce tension between different observables
- Higher order corrections can help further reduce uncertainty and increase stability

## Future work

- Heavy hemisphere mass with non-perturbative shift
- More precise calculation of shift
- Potential use of  $\Delta PS$  for other values of  $z_{\text{cut}}$  and  $\beta$

# Summary

## Further future work

- Explore an increased fitting range
- Move to NNLL accuracy including transition point effects
- Analyze different types of observables
- Potential to simultaneously study two or more independent observables
- Hopefully one day a full measurement



# Summary

## Further future work

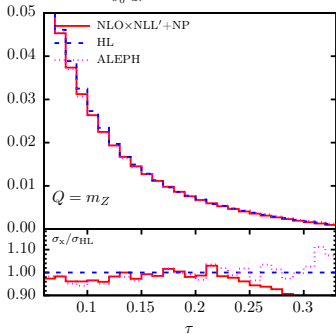
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**Thank you for your attention**

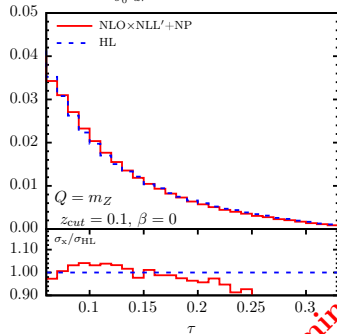
# Comparison HL to fit

[Marzani, Reichelt, Schumann, Soyez, VT; In Preparation]

$$\frac{\Delta\tau}{\sigma_0} \frac{d\sigma}{d\tau}(e^+e^- \rightarrow q\bar{q} + X)$$



$$\frac{\Delta\tau}{\sigma_0} \frac{d\sigma}{d\tau}(e^+e^- \rightarrow q\bar{q} + X)$$

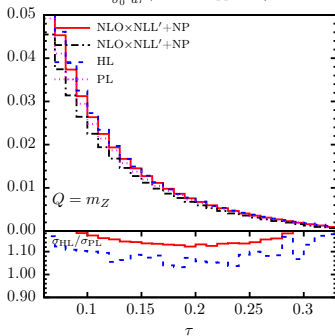


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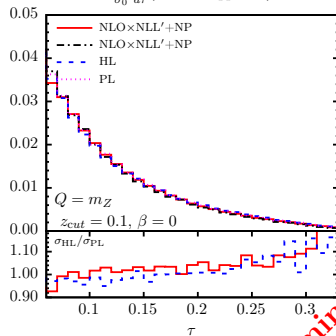
# Comparison HL/PL to fit

[Marzani, Reichelt, Schumann, Soyez, VT; In Preparation]

$$\frac{\Delta\tau}{\sigma_0} \frac{d\sigma}{d\tau}(e^+e^- \rightarrow q\bar{q} + X)$$



$$\frac{\Delta\tau}{\sigma_0} \frac{d\sigma}{d\tau}(e^+e^- \rightarrow q\bar{q} + X)$$



preliminary