

Astro@TS

# Covariance matrix estimation for the statistics of galaxy clustering

Manuel Colavincenzo

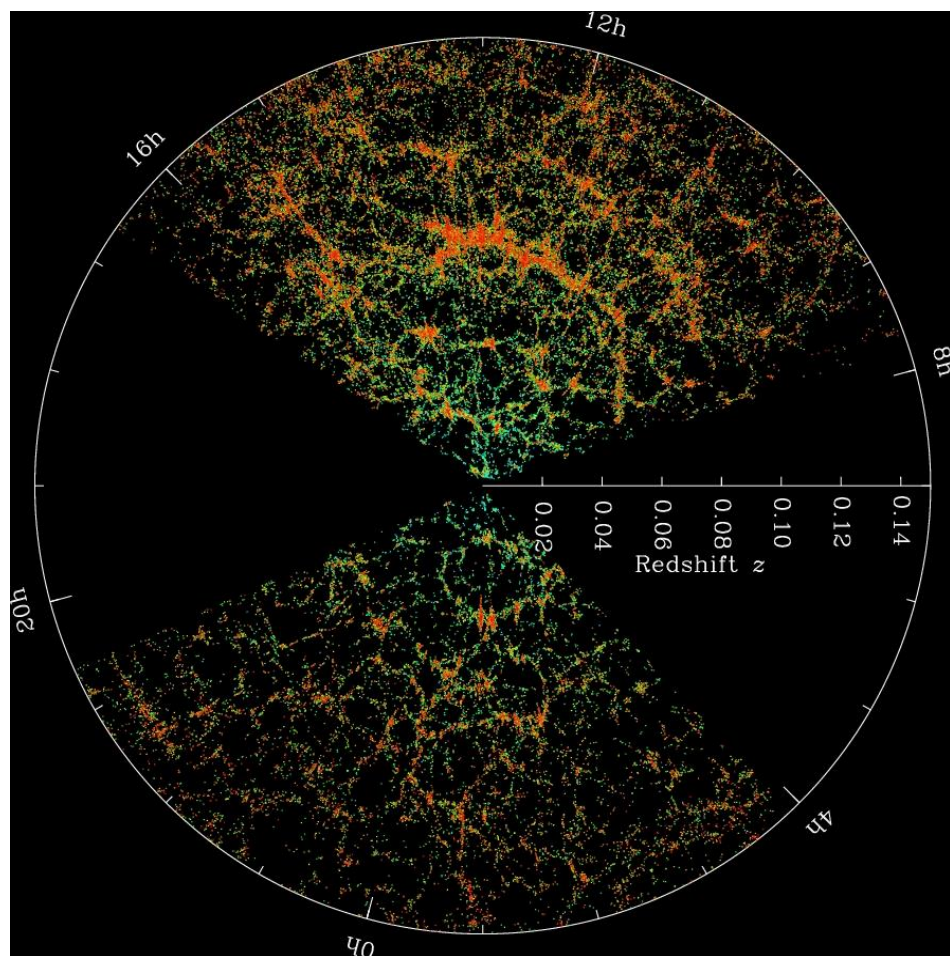
Supervisor: Pierluigi Monaco

Co-Supervisors: Stefano Borgani, Emiliano Sefusatti

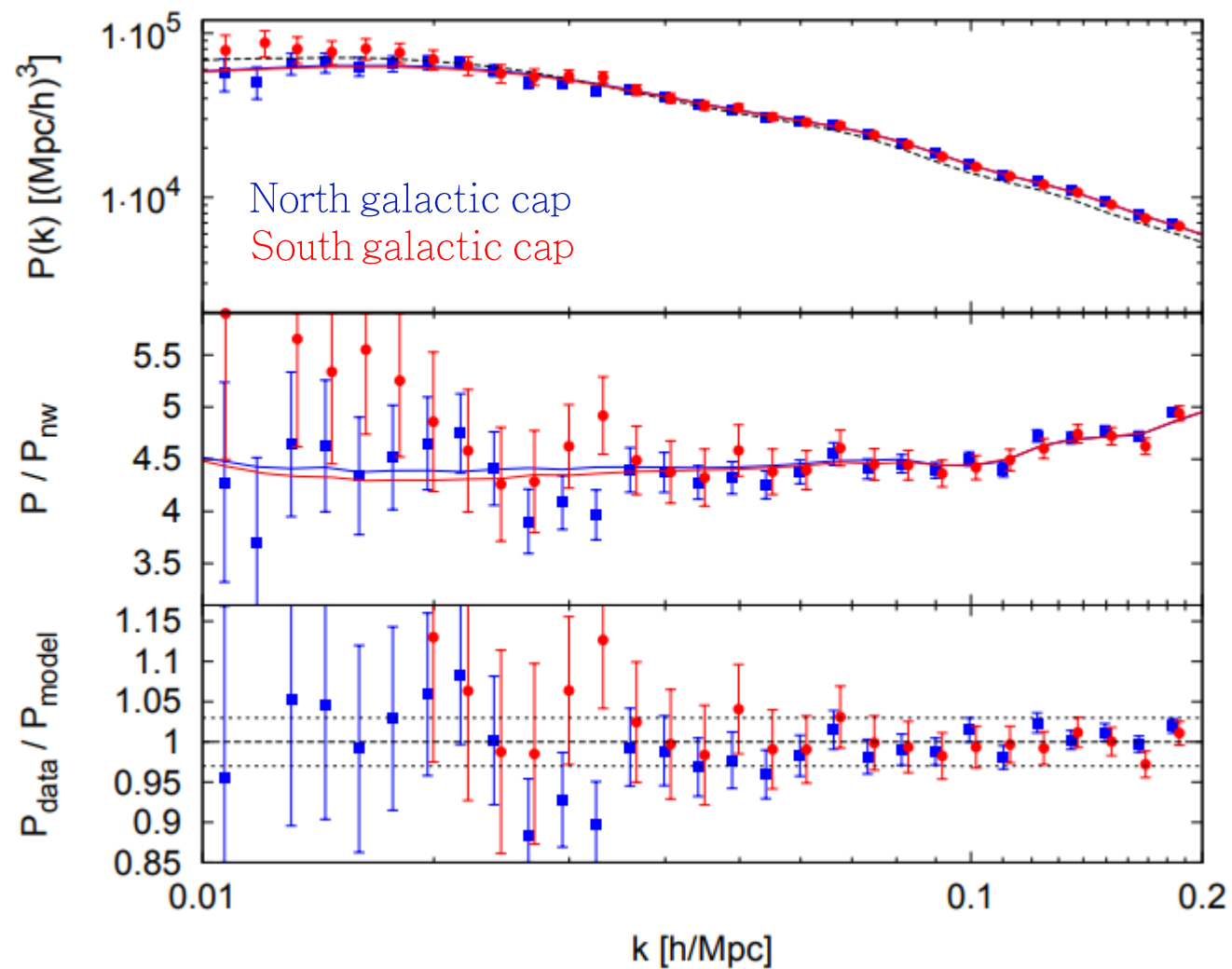


# General introduction

Gil-Marín et al. 2015  
(SDSS-BOSS paper)

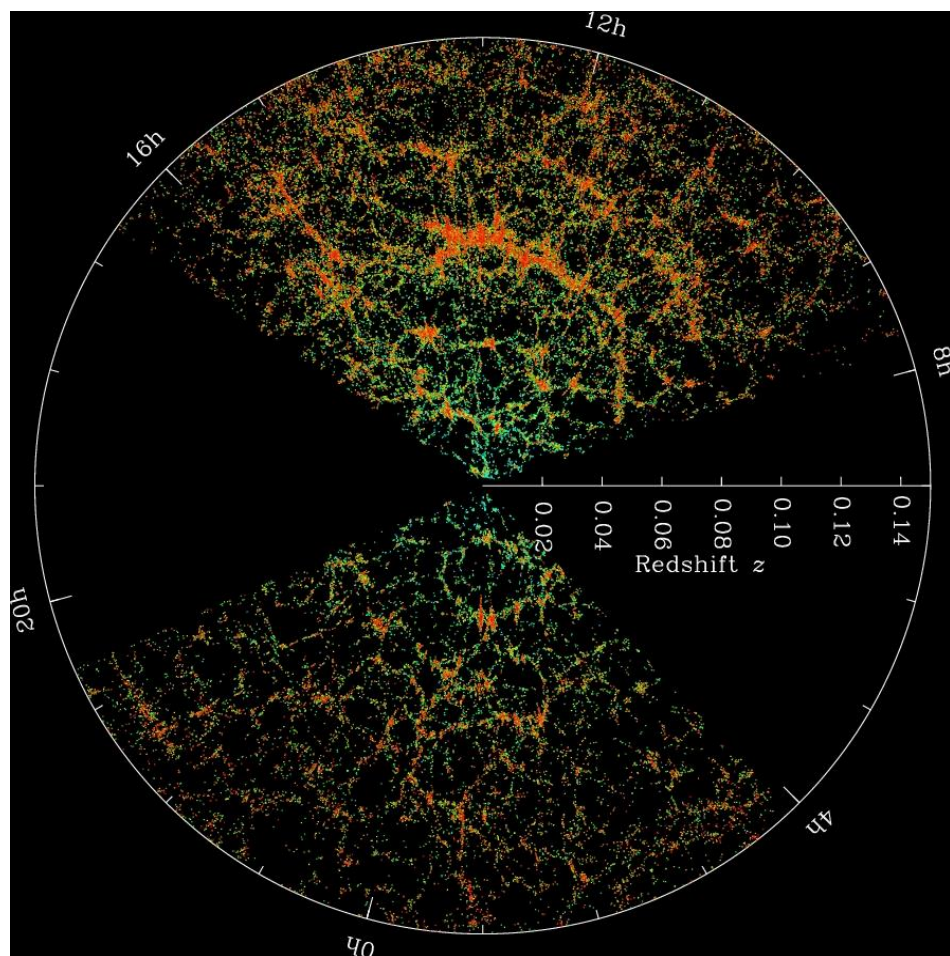


Credits: SDSS

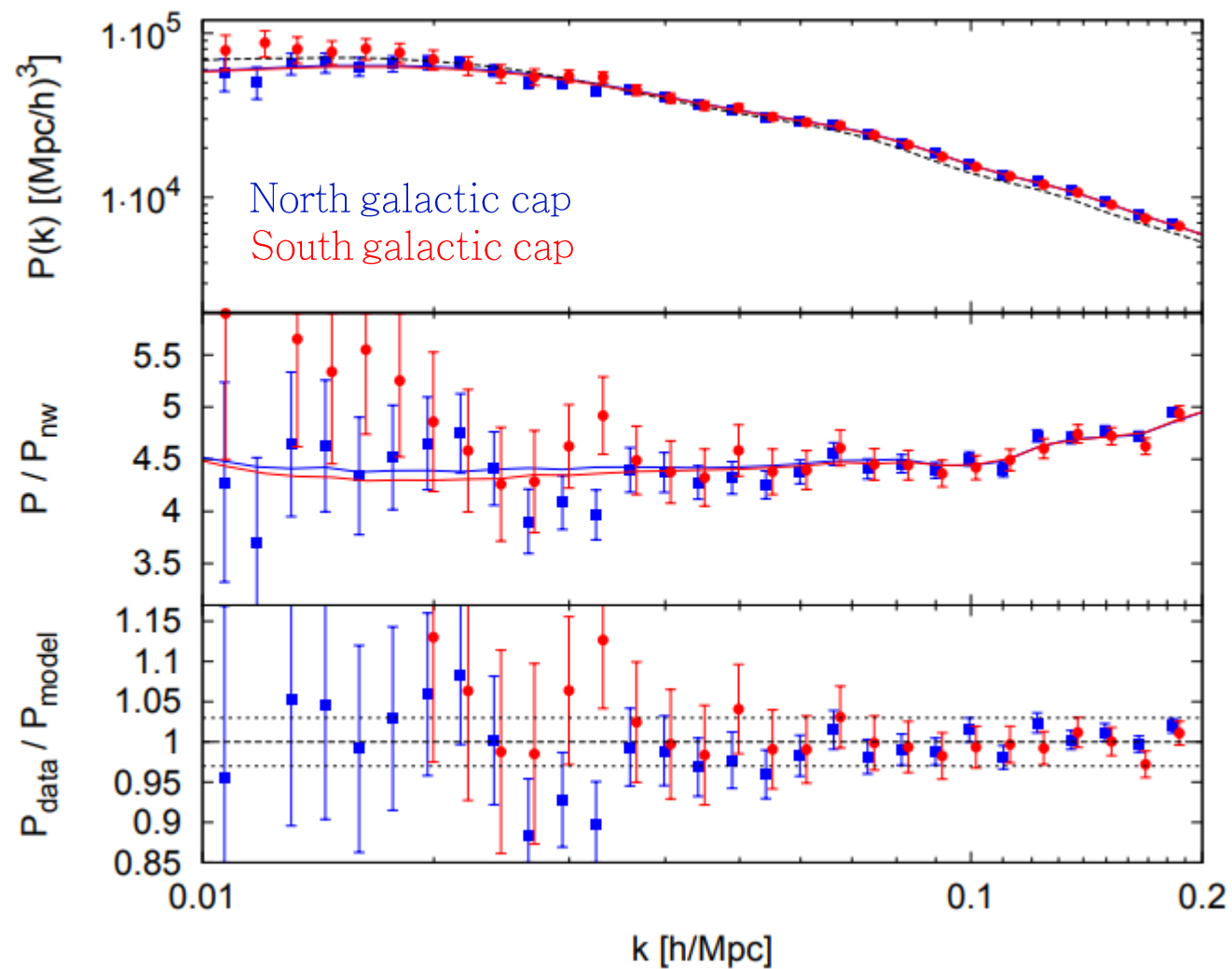


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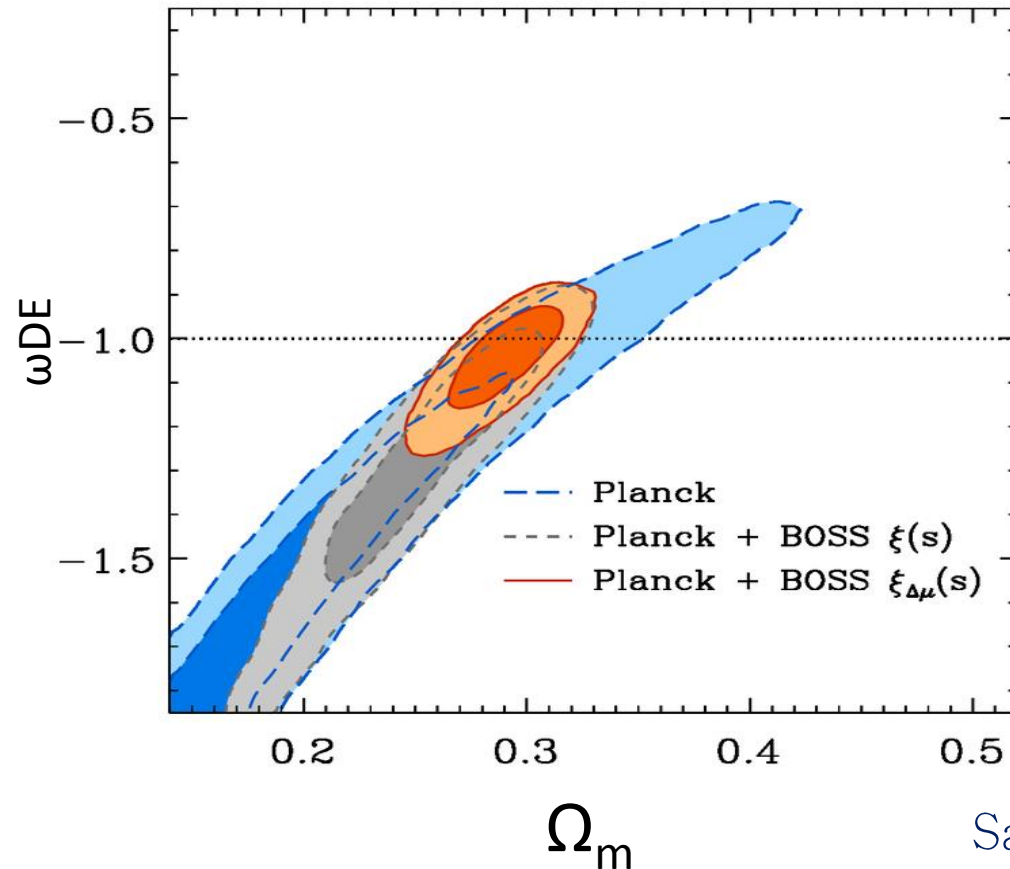


Error bars: dispersion among 2048 simulated galaxy catalogues

# General introduction

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An accurate modeling of the power spectrum covariance matrix is needed to constrain **cosmological parameters**



# Uncertainty in the visibility mask of a survey and its effects on the clustering of biased tracers

M. Colavincenzo,<sup>a</sup> P. Monaco,<sup>a,b,d</sup> E. Sefusatti,<sup>c,e</sup> S. Borgani<sup>a,b,d</sup>

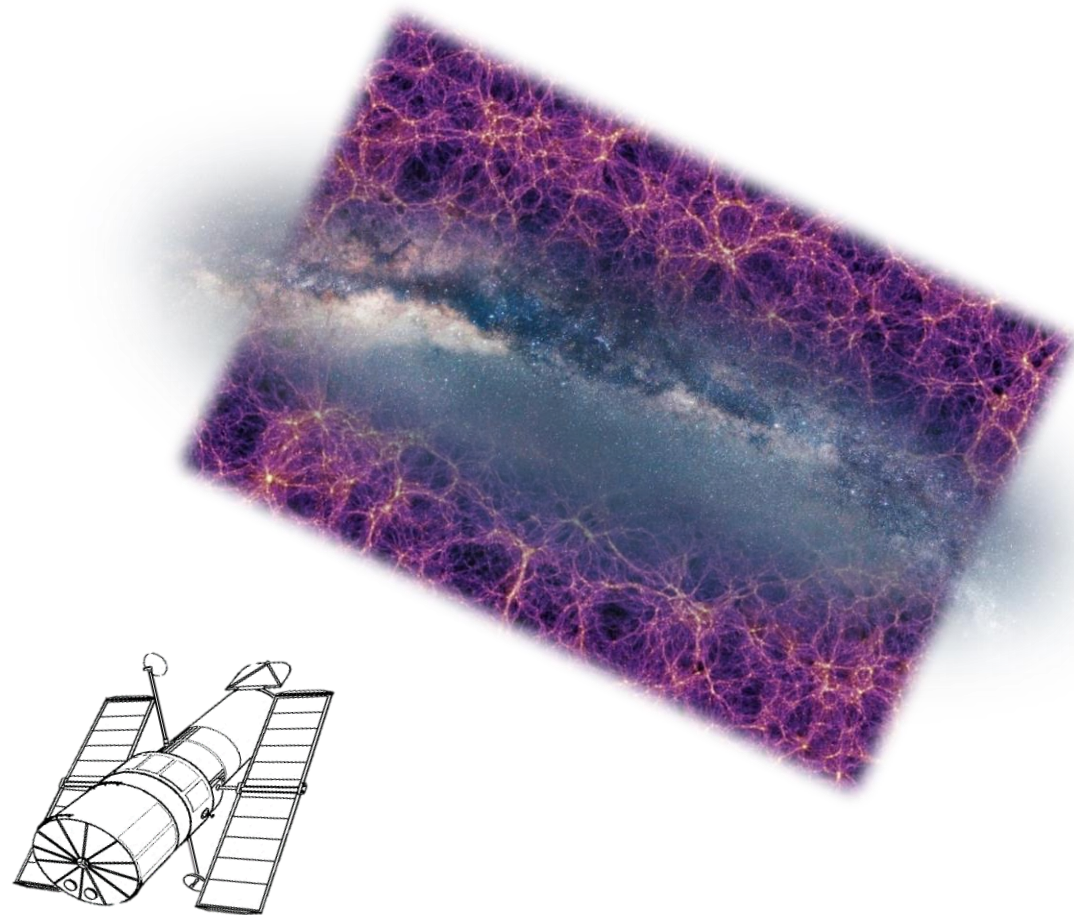
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[Journal of Cosmology and Astroparticle Physics](#), [Volume 2017](#), [March 2017](#)

# The foreground problem

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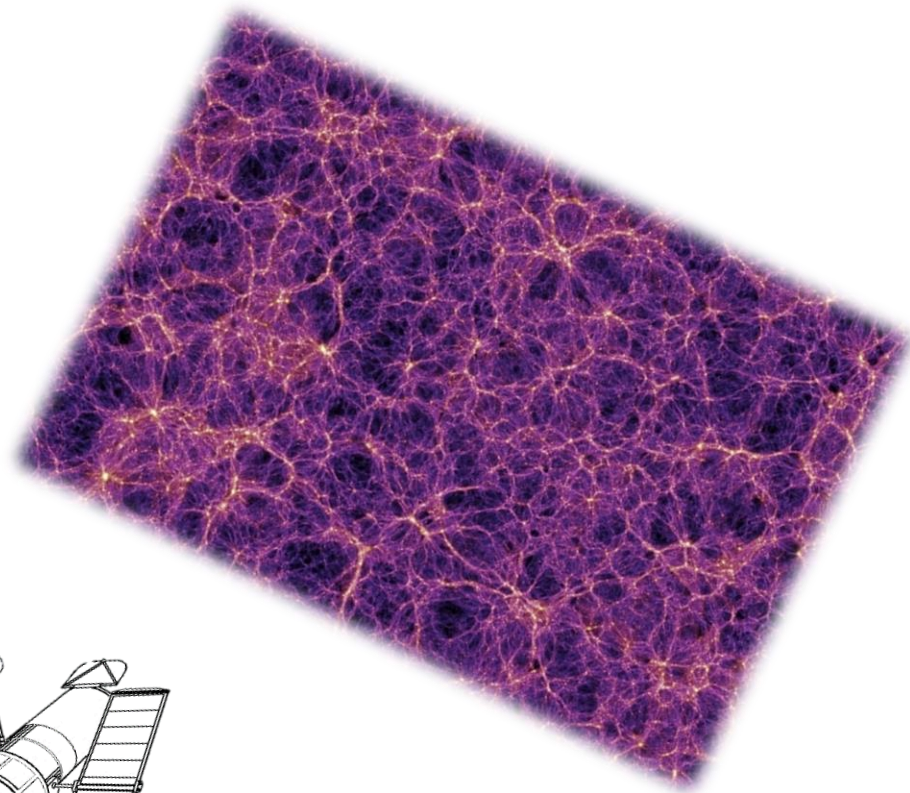
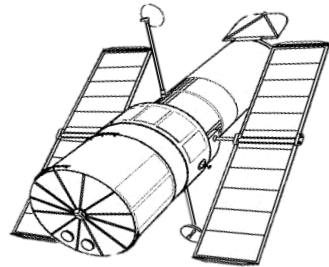
The limiting magnitude of a galaxy sample is modulated by foregrounds



# The foreground problem

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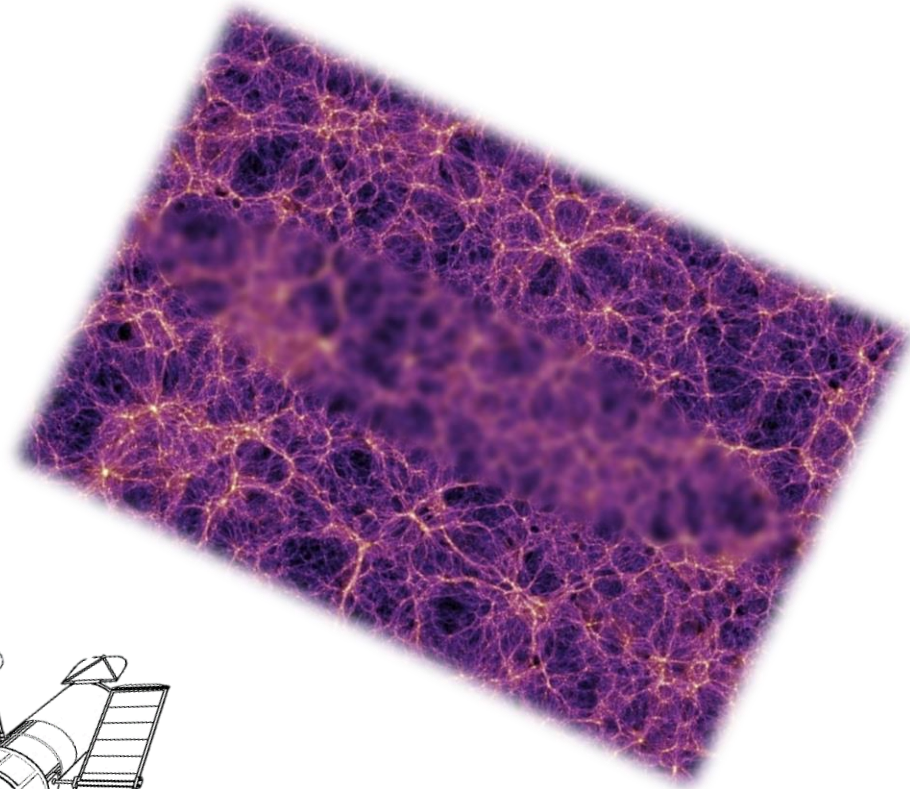
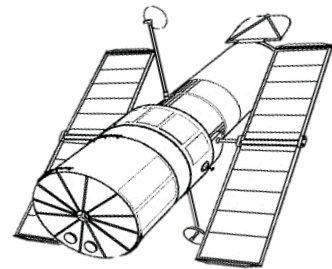
Cosmological information can be extracted after foreground subtraction



# The foreground problem

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This subtraction is subject to an uncertainty and the residual foreground will be correlated on the sky and will create structure on large scales





# The foreground problem

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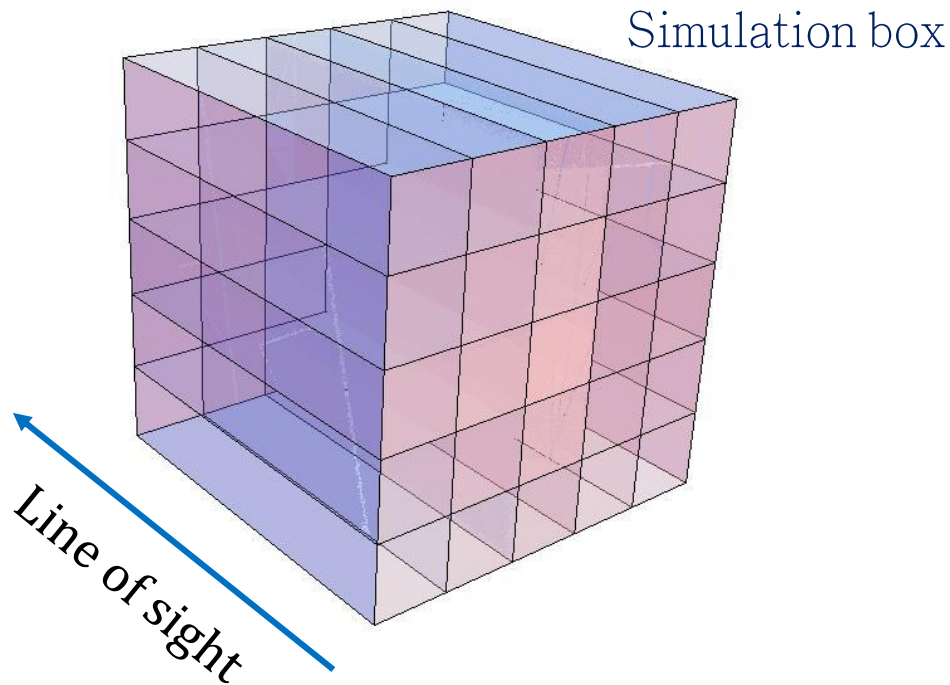
Why is it fundamental the foreground analysis?

Future galaxy surveys will provide very high statistics -----> the error budget will be dominated by **systematics**

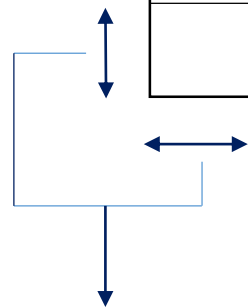
{ Milky way extinction  
Instrumental and survey features  
...

The error on foreground subtraction must be properly propagated to correctly assess the error on cosmological parameters estimation.

# The model for the “mask”



$A_1$	$A_2$	$A_3$	$A_4$	
				$A_N$



Correlation scales

$$\langle A \rangle = \frac{1}{N} \sum_{i=1}^N A_i = 0$$

$$\langle A^2 \rangle = 0.01, 0.05, 0.1, 0.2$$

$$A(\boldsymbol{\theta}) \equiv \frac{\delta M(\boldsymbol{\theta})}{M_0} \longrightarrow \delta_{\text{mask}}(\mathbf{x}) = -\frac{M_0 \bar{\Phi}(M_0)}{\bar{n}} A(\boldsymbol{\theta}) + \mathcal{O}(A^2)$$

$$\sigma_{\text{mask}}^2 \simeq \frac{M_0^2 \bar{\Phi}^2(M_0)}{\bar{n}^2(M_0)} \sigma_A^2$$

# Prediction for the PS covariance matrix

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$$C_{ij} \equiv \text{cov}[\hat{P}(k_i), \hat{P}(k_j)] = \langle \delta \hat{P}(k_i) \delta \hat{P}(k_j) \rangle$$

$$C_{ij}^{\text{obs}} = C_{ij}^{\text{cosm}} + C_{ij}^{\text{mask}} + C_{ij}^{\text{mixed}}$$

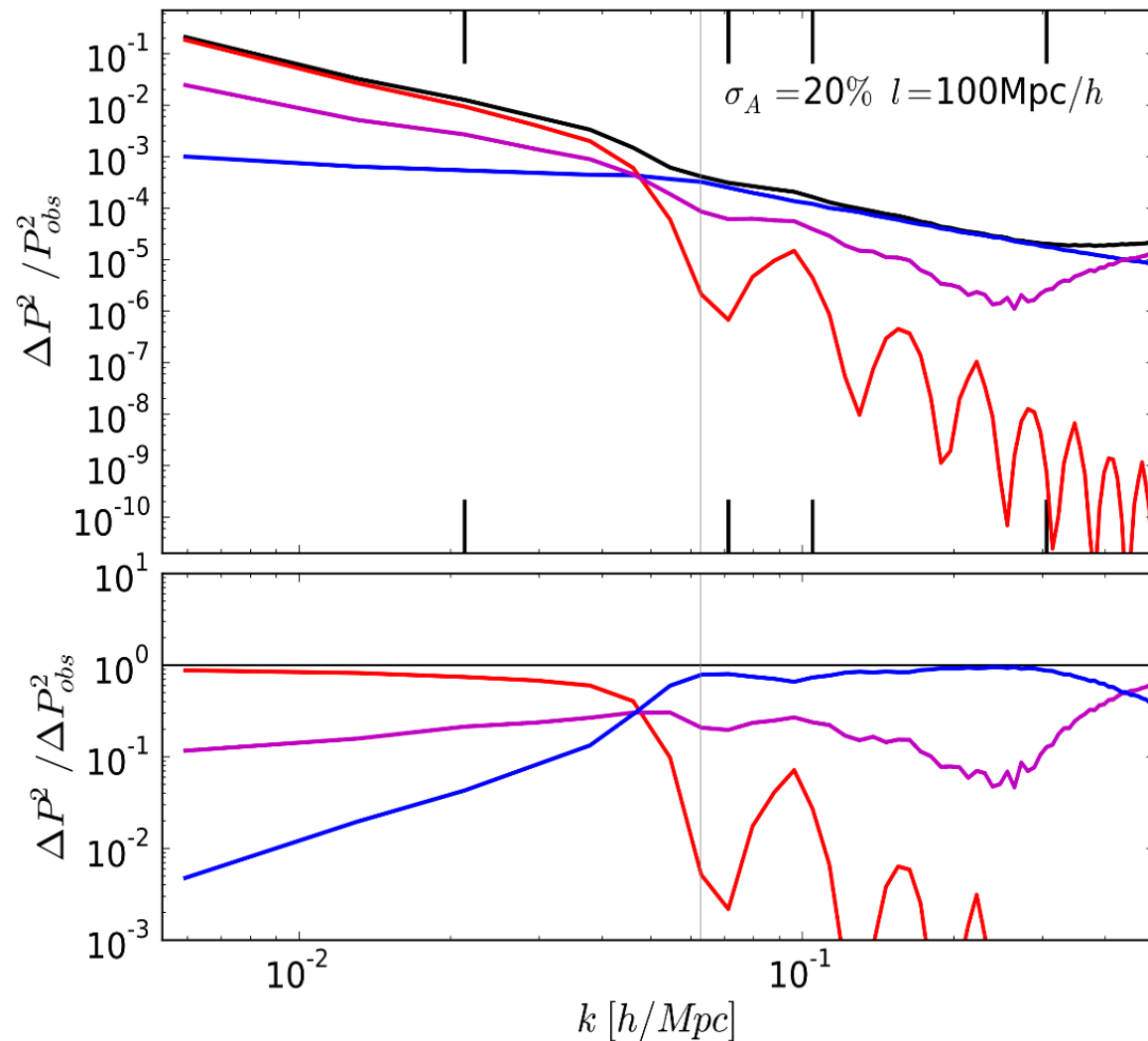
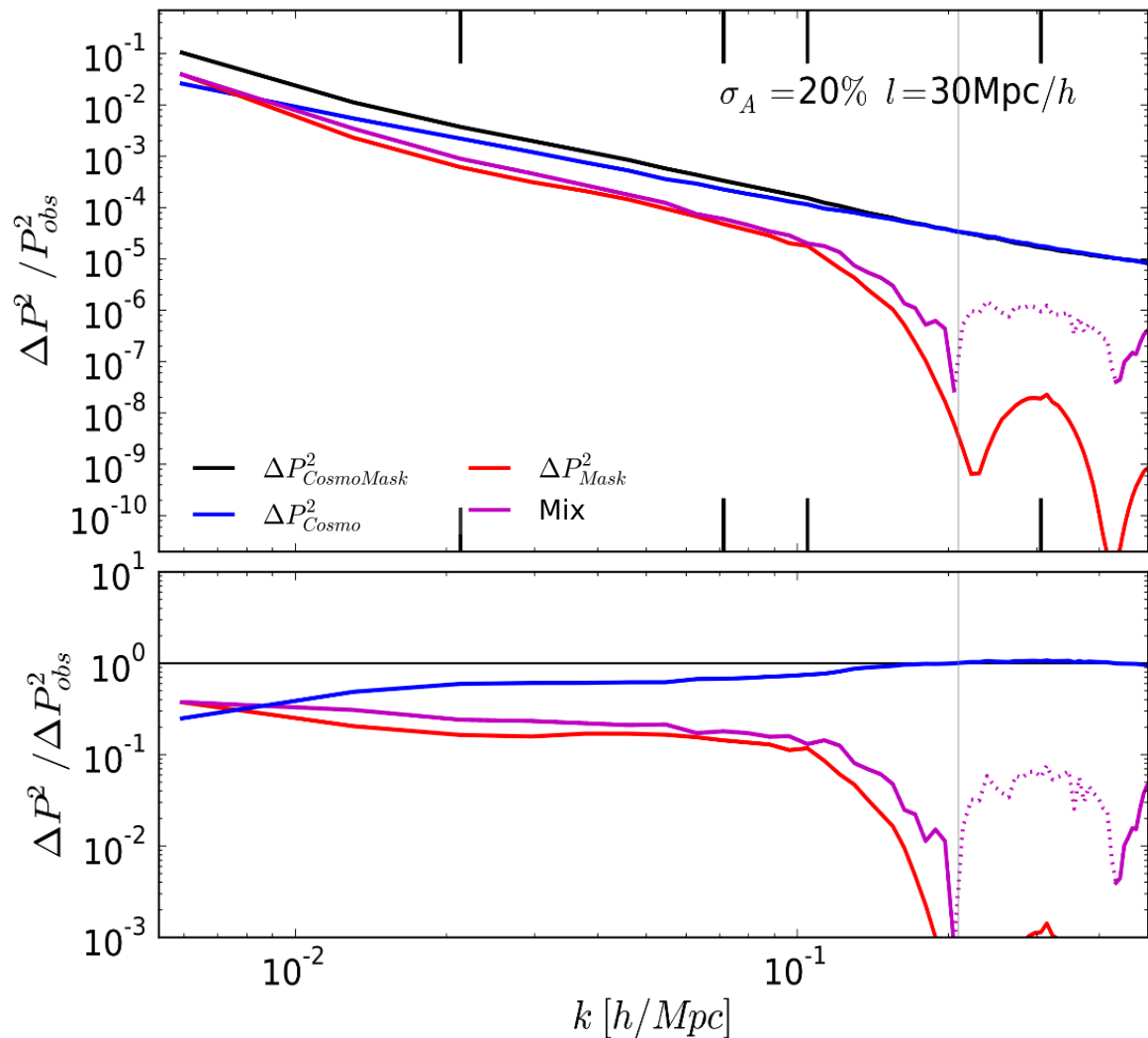
$$\begin{aligned} C_{ij}^{\text{mixed}} = & \langle \hat{P}_{\text{conv}}(k_i) \hat{P}_{\text{conv}}(k_j) \rangle - \langle \hat{P}_{\text{conv}}(k_i) \rangle \langle \hat{P}_{\text{conv}}(k_j) \rangle + \\ & \langle \hat{P}_{\text{cosmo}}(k_i) \hat{P}_{\text{conv}}(k_j) \rangle - \langle \hat{P}_{\text{cosmo}}(k_i) \rangle \langle \hat{P}_{\text{conv}}(k_j) \rangle + \\ & \langle \hat{P}_{\text{mask}}(k_i) \hat{P}_{\text{conv}}(k_j) \rangle - \langle \hat{P}_{\text{mask}}(k_i) \rangle \langle \hat{P}_{\text{conv}}(k_j) \rangle + \\ & \langle \hat{P}_{\text{cosmo}}(k_i) \hat{G}(k_j) \rangle + \langle \hat{P}_{\text{mask}}(k_i) \hat{G}(k_j) \rangle + \\ & \langle \hat{P}_{\text{conv}}(k_i) \hat{G}(k_j) \rangle + \langle \hat{G}(k_i) \hat{G}(k_j) \rangle \end{aligned}$$



$$\hat{G} = 2\delta_{\mathbf{q}}\delta_{\text{mask},\mathbf{q}} - \delta_{\mathbf{q}}\delta_{\text{conv},\mathbf{q}} + \delta_{\text{mask},\mathbf{q}}\delta_{\text{conv},\mathbf{q}}$$

# Results: Variance

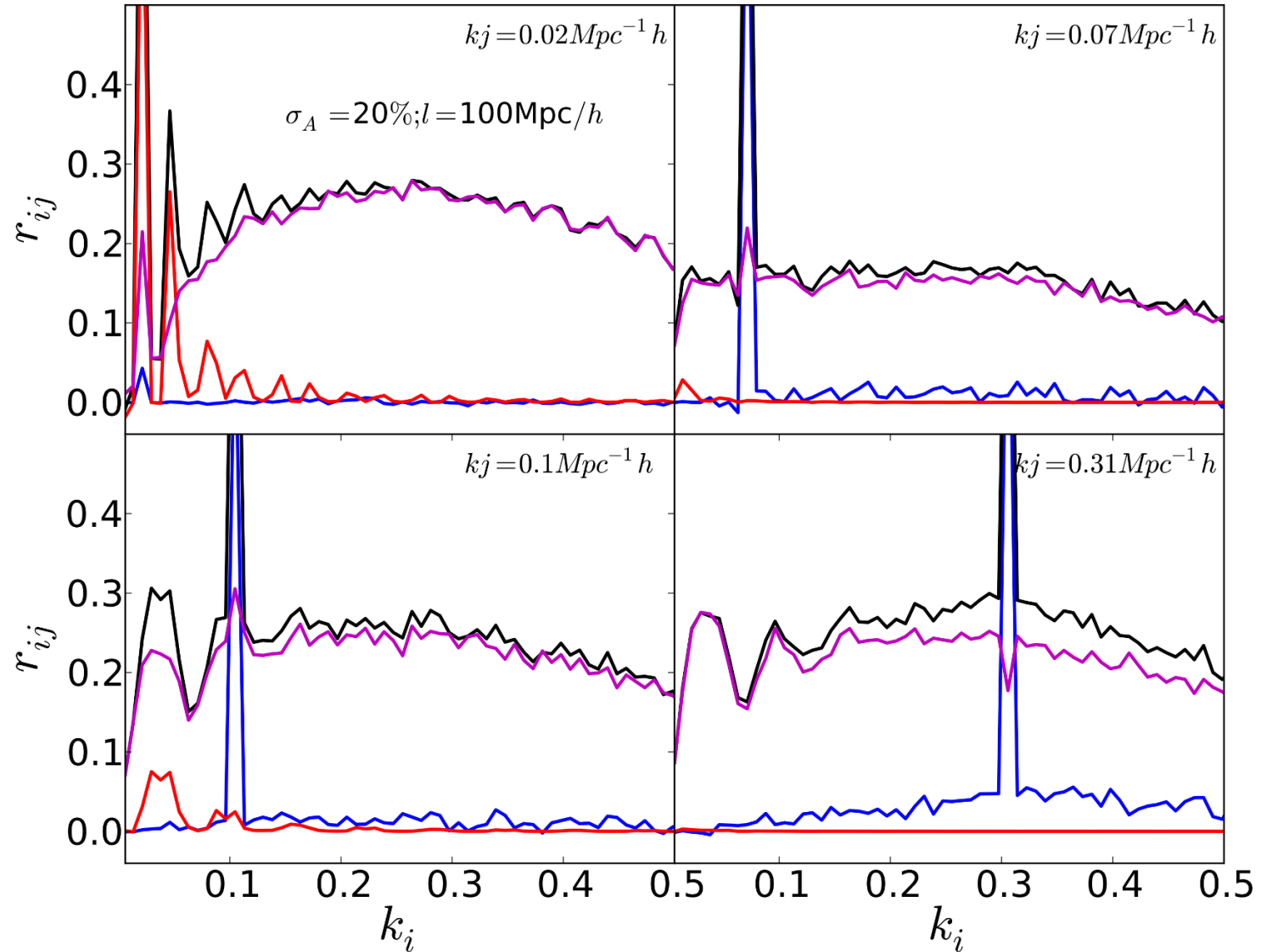
10 000 simulated galaxy catalogs



# Results: Covariance

10 000 simulated galaxy catalogs

$$\frac{C_{ij}}{\sqrt{C_{ii}C_{jj}}} = \frac{\text{cov}_{ij}}{\sqrt{\text{var}_i \text{var}_j}}$$



# Results

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The error in the subtraction of the foreground is relevant and may be not negligible.

The convolution term, in the power spectrum, enters in the covariance mixing the modes and transferring power from large to small scales.

Non negligible off-diagonal terms could be enhanced by the coupling of the mask and the cosmological power spectrum

# Results

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The convolution term, in the power spectrum, enters in the covariance mixing the modes and transferring power from large to small scales.

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## Next steps

Study the effects using more realistic foreground (extinction maps, zodiacal light maps, ...)

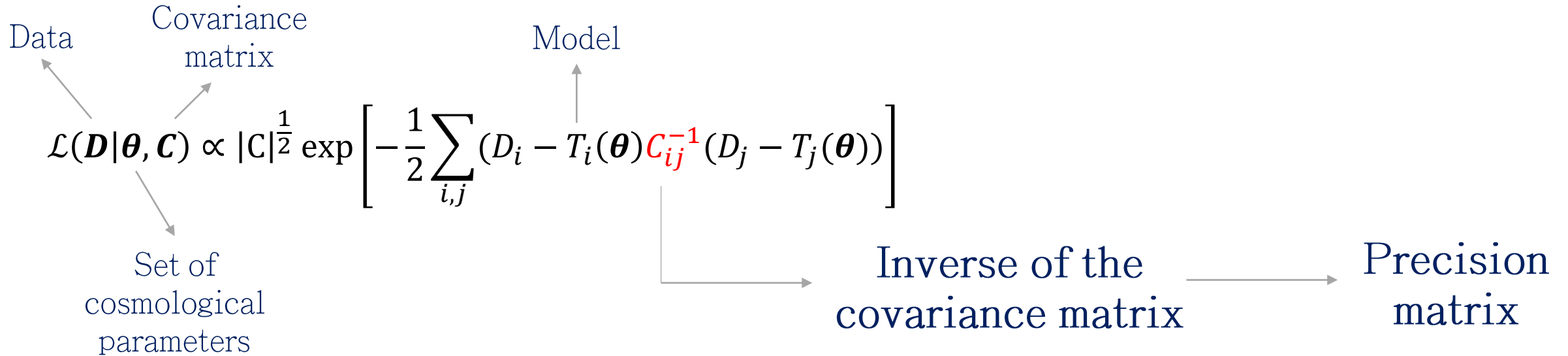


# Covariance matrix comparison

M. Colavincenzo & L. Blot, M. Crocce, P. Monaco, A. Sanchez, E. Sefusatti, et al. (in preparation)



# The approximate methods

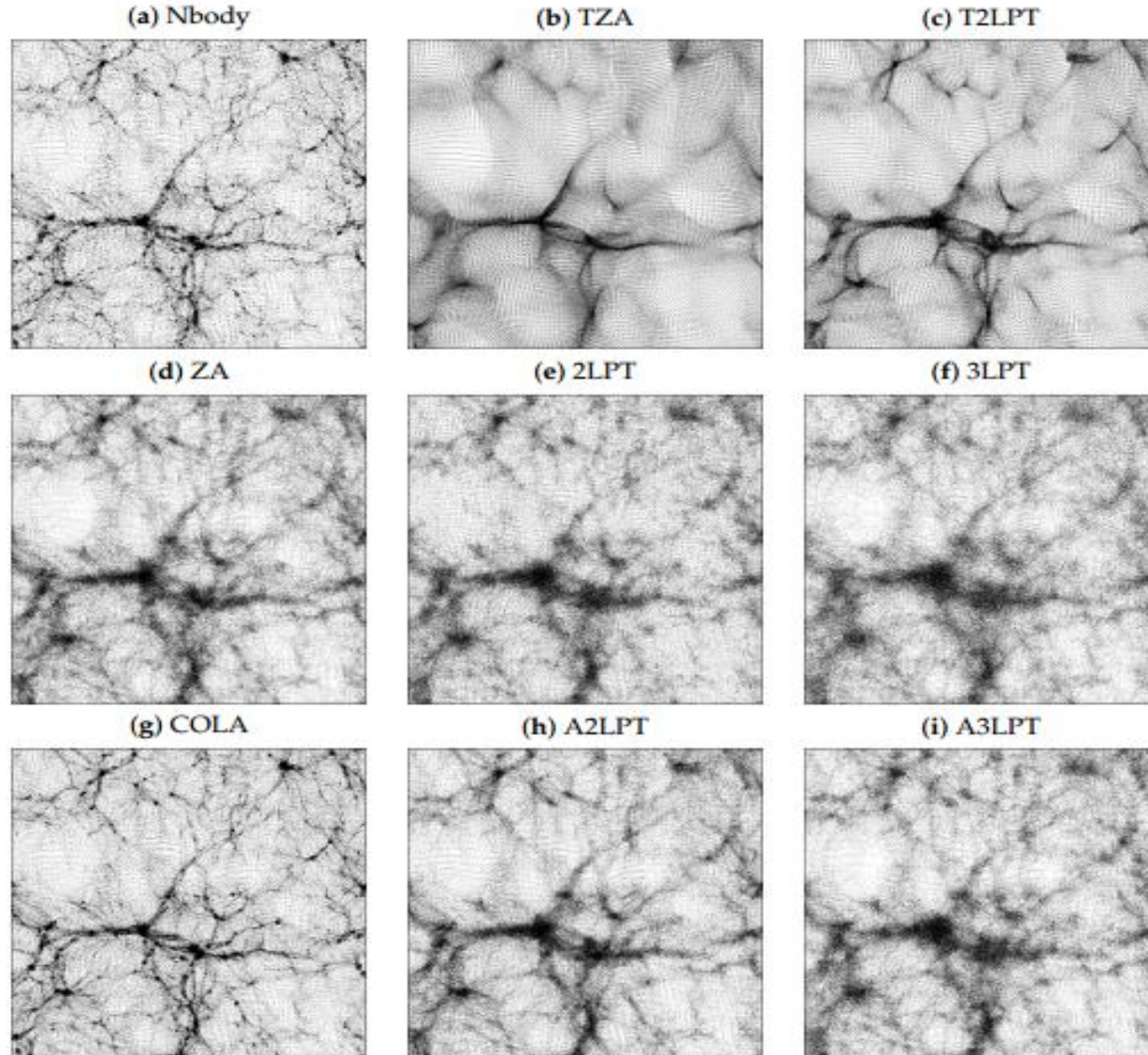


In most analyses of clustering measurements, the **precision matrix** is estimated from an ensemble of simulated catalogues

All estimates of the precision matrix based on a finite number of catalogues are affected by noise.

A large number of simulated catalogues might be necessary in order to keep the uncertainty under control

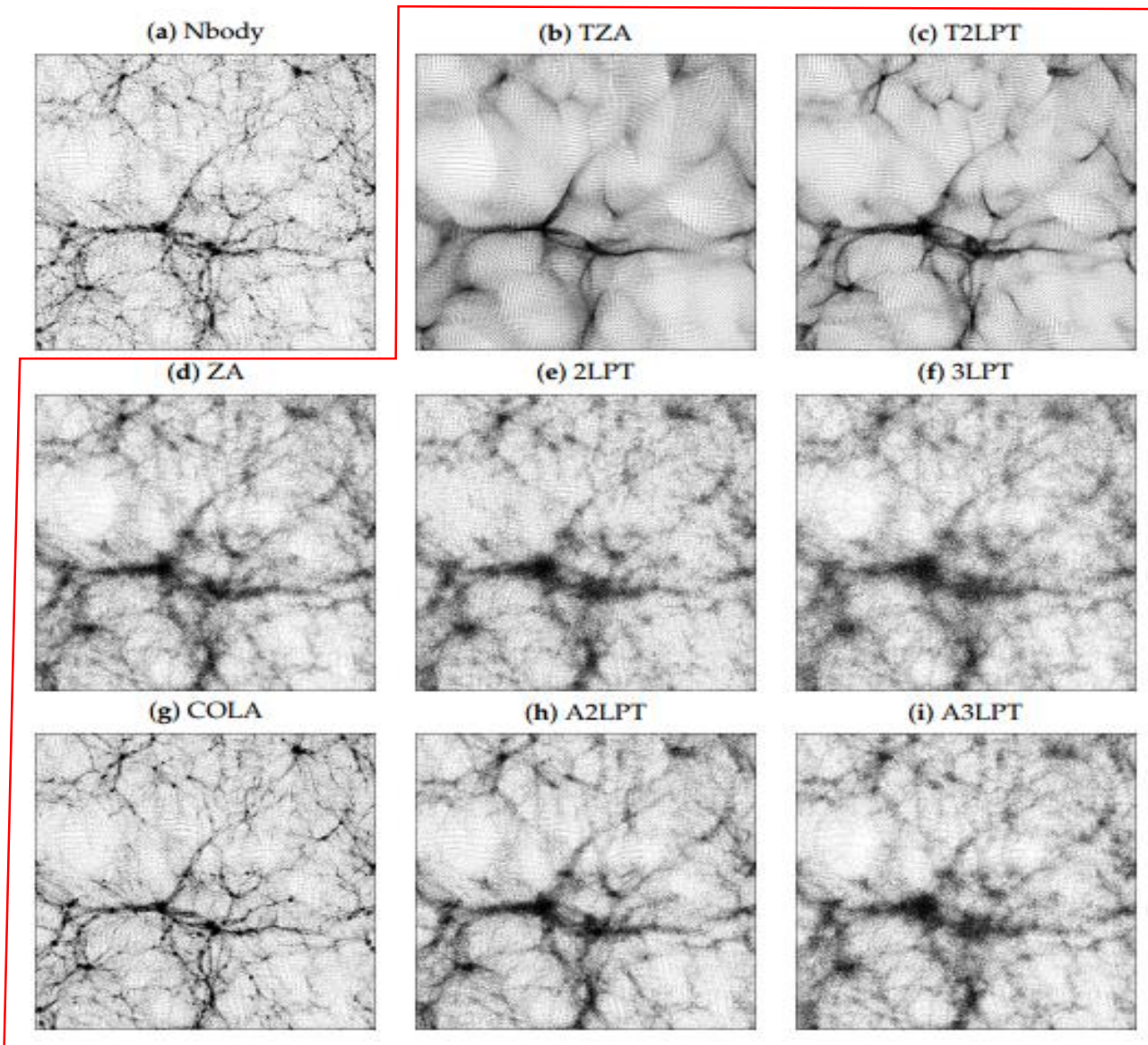
# The approximate methods



N-Body simulations: calculate the **non-linear growth** of structures in the Universe by following the trajectories of  $N$  particles interacting between each others through gravity

Approximate methods: take advantage of **analytic approximations** to produce a large number of galaxy catalogues reducing the computing time of a factor  $\sim 1000$ , paying a loss in accuracy of the small scales ( $k \sim 0.5 h^{-1} \text{Mpc}$ )

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# The approximate methods

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## Main questions:

- Can we use approximate methods to estimate the covariance matrices of clustering measures?
- Do approximate methods reproduce the error on cosmological parameters?

## Tools we will use to answer:

- Halo power spectrum (L. Blot et al.)
- Bispectrum (M. Colavincenzo et al.)
- Two-point correlation function (A. Sanchez et al.)

## Approximate methods:

- Pinocchio
- Cola
- Halogen
- PeakPatch
- Patchy,...

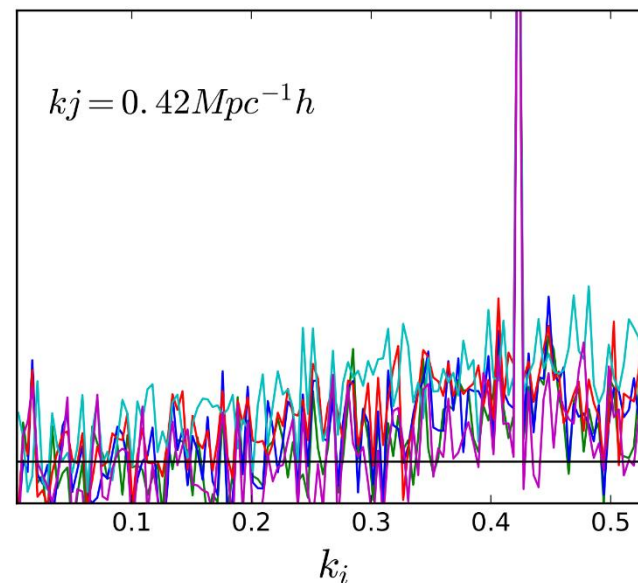
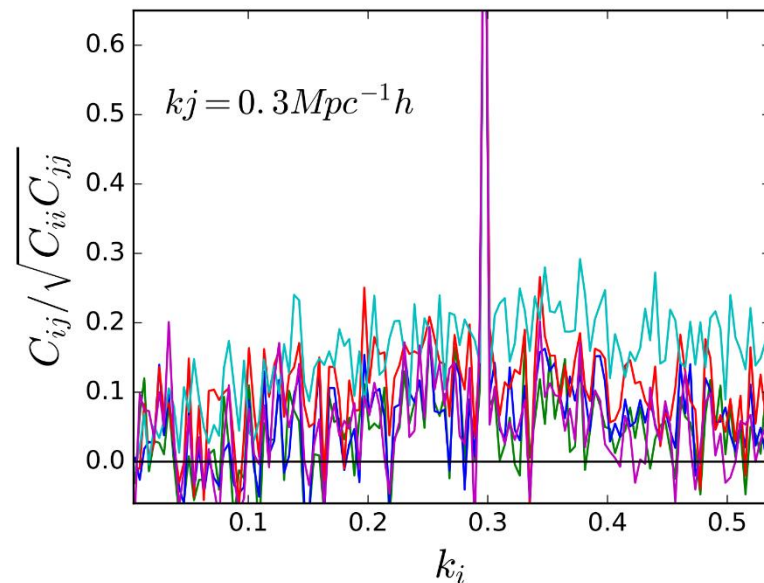
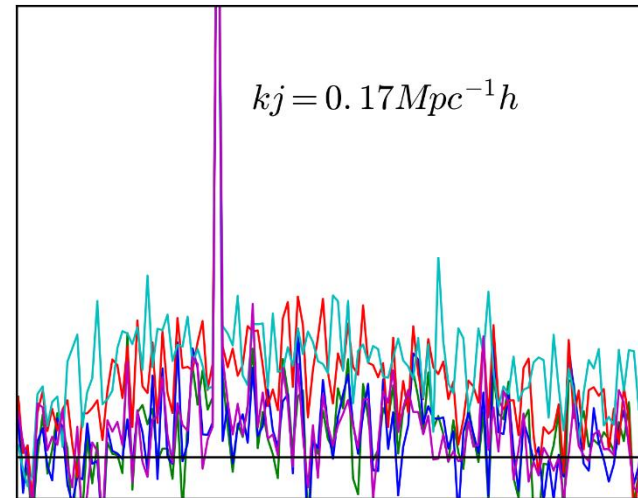
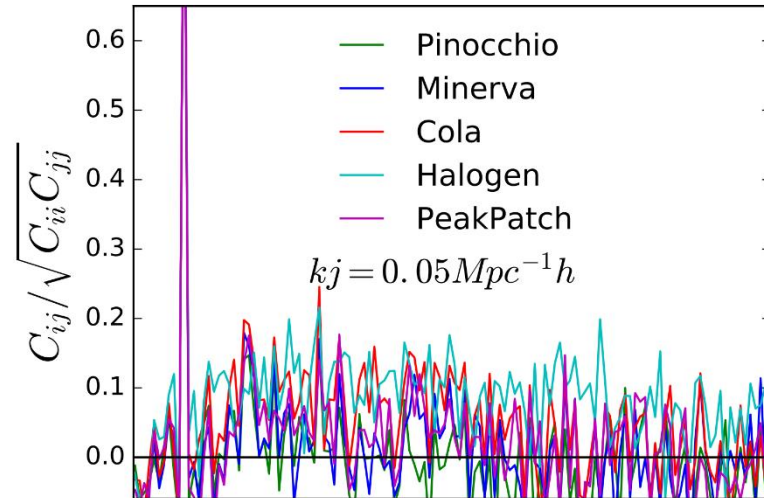
## Reference simulations:

- N-body (Minerva)  
[Gadget-2]

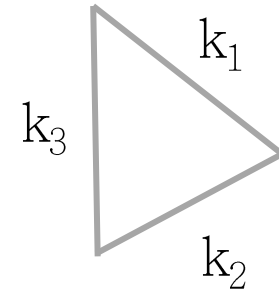
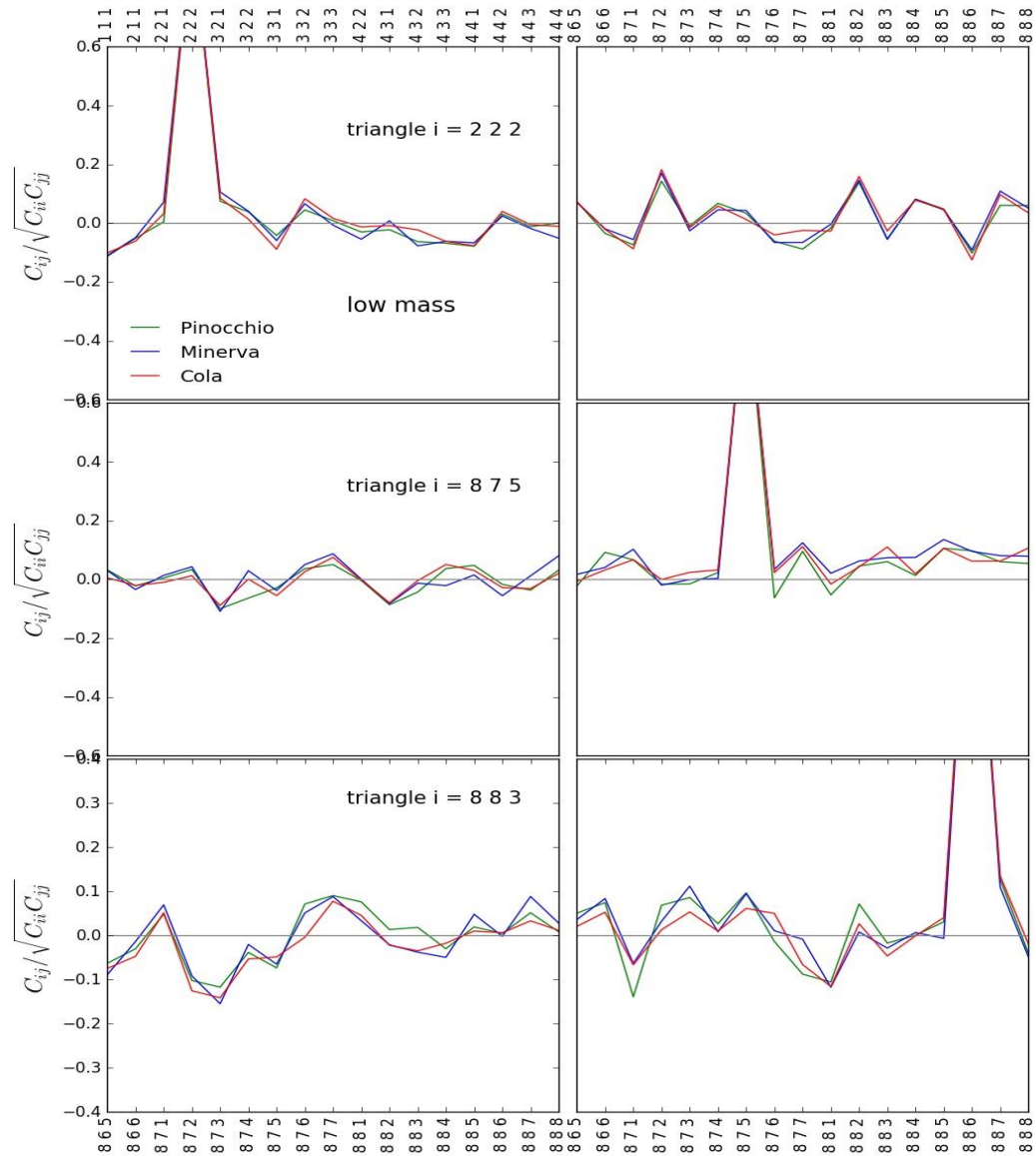
## Simulation infos:

- Cubic box: 1500 Mpc/h
- $z = 1$
- Real and Redshift space
- 300 realizations

# Power spectrum analysis



# Bispectrum analysis



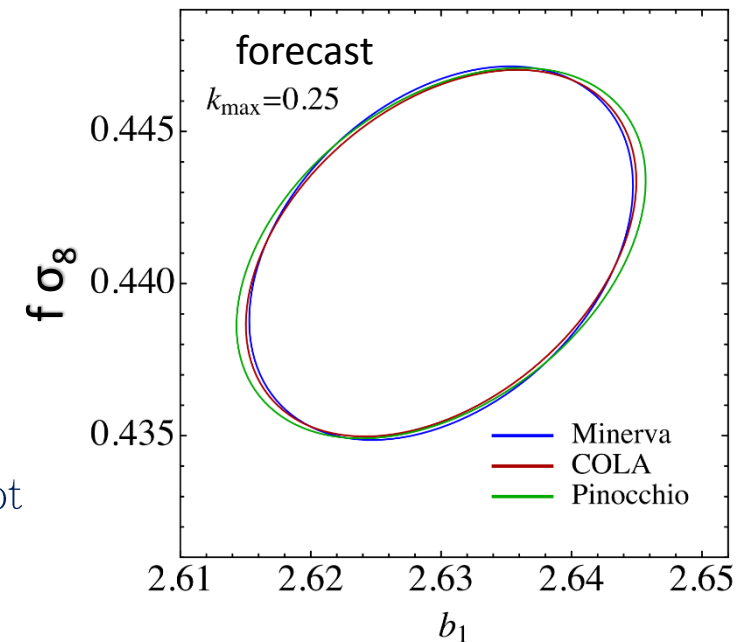
# Results

- Different approximate methods allow us to obtain the same correlated noise in the power spectrum and bispectrum covariance matrix
- We expect the physical signal is the same

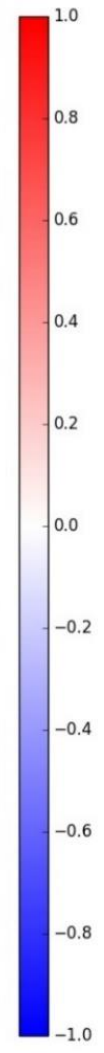
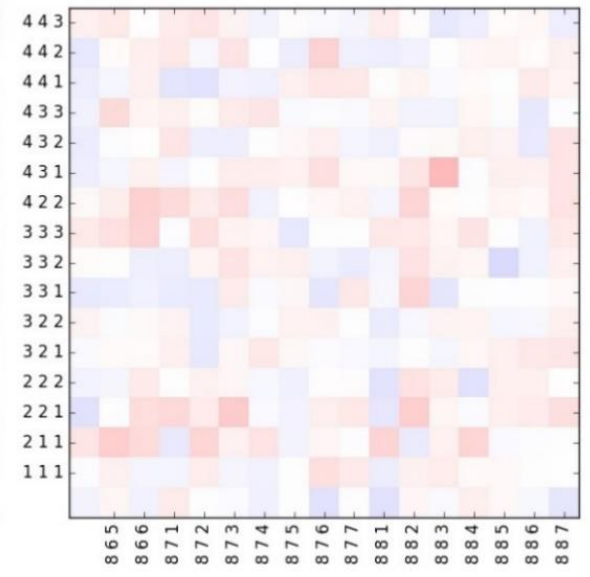
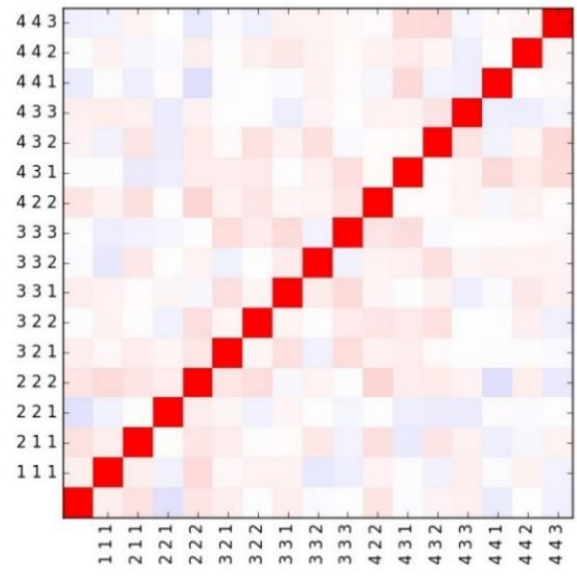
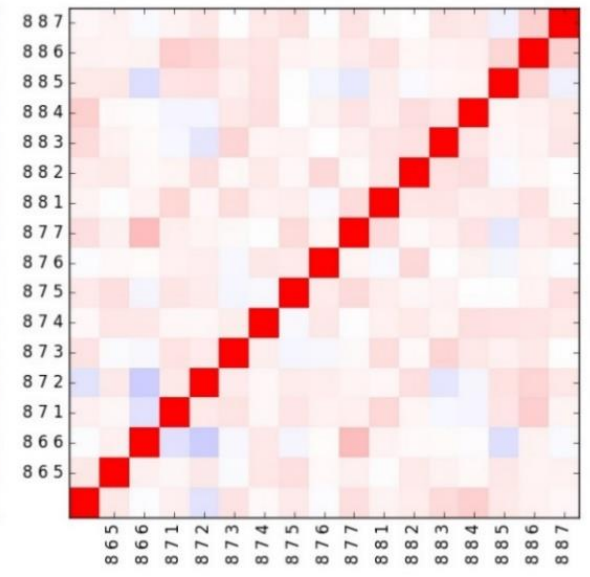
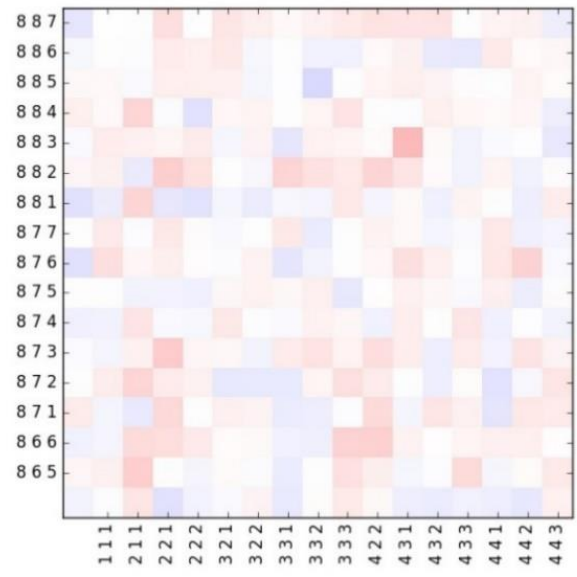
## Preliminary

MCMC results - 3 multipoles - $k_{\text{Max}} = 0.17 h \text{ Mpc}^{-1}$					
	$f \times \sigma_8$	$b_1$	$b_2$	$\gamma_3$	$a_{\text{vir}}$
Minerva	$0.444 \pm 0.022$	$2.625 \pm 0.015$	$2.669 \pm 0.4$	$1.83 \pm 0.134$	$1.14 \pm 0.76$
COLA	$0.446 \pm 0.023$	$2.627 \pm 0.013$	$2.467 \pm 0.348$	$1.86 \pm 0.097$	$0.80 \pm 0.27$
Pinocchio	$0.451 \pm 0.019$	$2.627 \pm 0.012$	$2.326 \pm 0.254$	$1.87 \pm 0.086$	$0.63 \pm 0.35$

Credits: Linda Blot



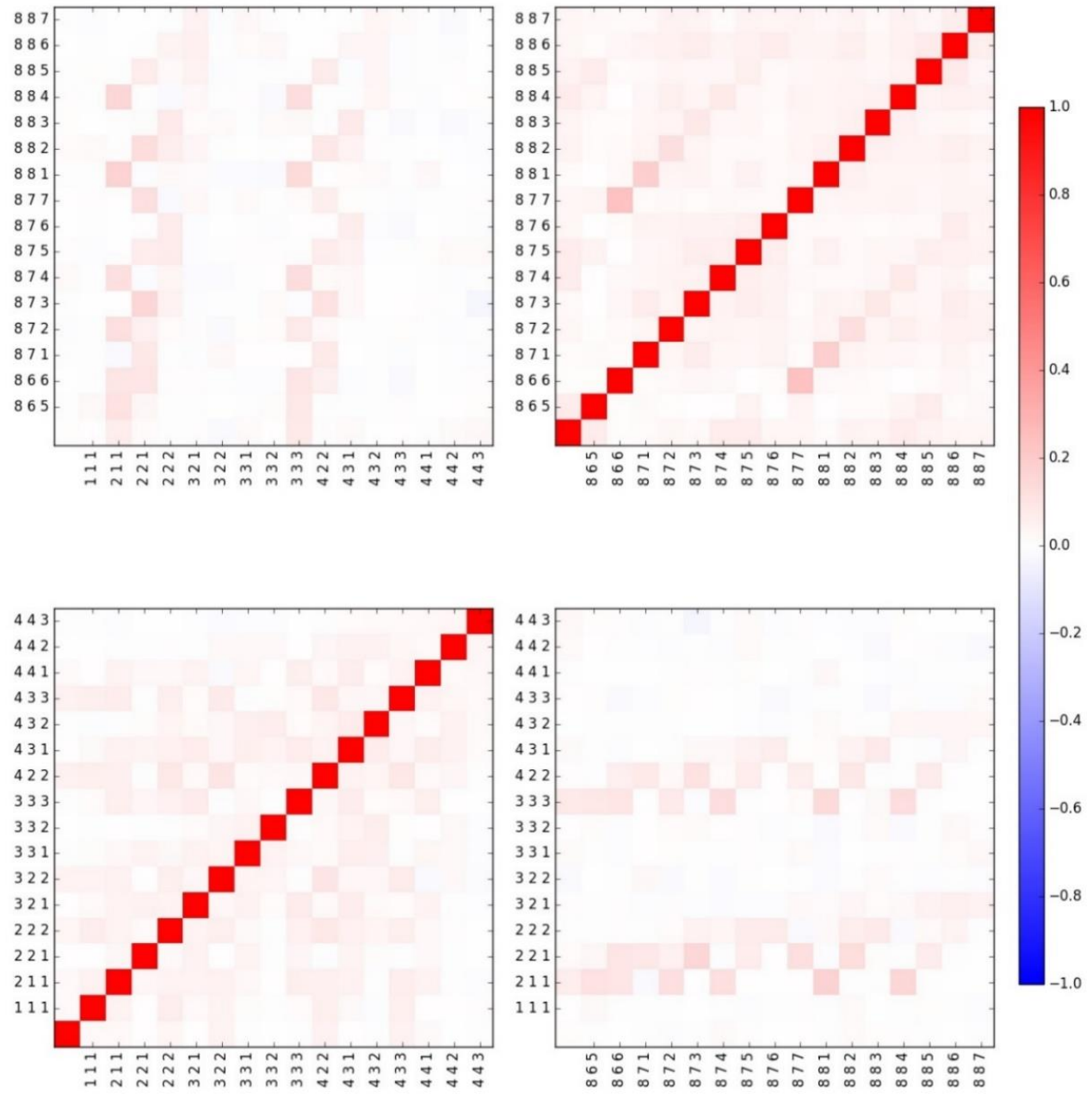
$$\frac{C_{ij}}{\sqrt{C_{ii}C_{jj}}} = \frac{\text{cov}_{ij}}{\sqrt{\text{var}_i \text{var}_j}}$$



300 N-body  
simulations



$$\frac{C_{ij}}{\sqrt{C_{ii}C_{jj}}} = \frac{\text{COV}_{ij}}{\sqrt{\text{var}_i \text{var}_j}}$$



10 000  
PINOCCHIO  
realizations