

10 years of Higgs boson measurements

lessons learned and plans for the future

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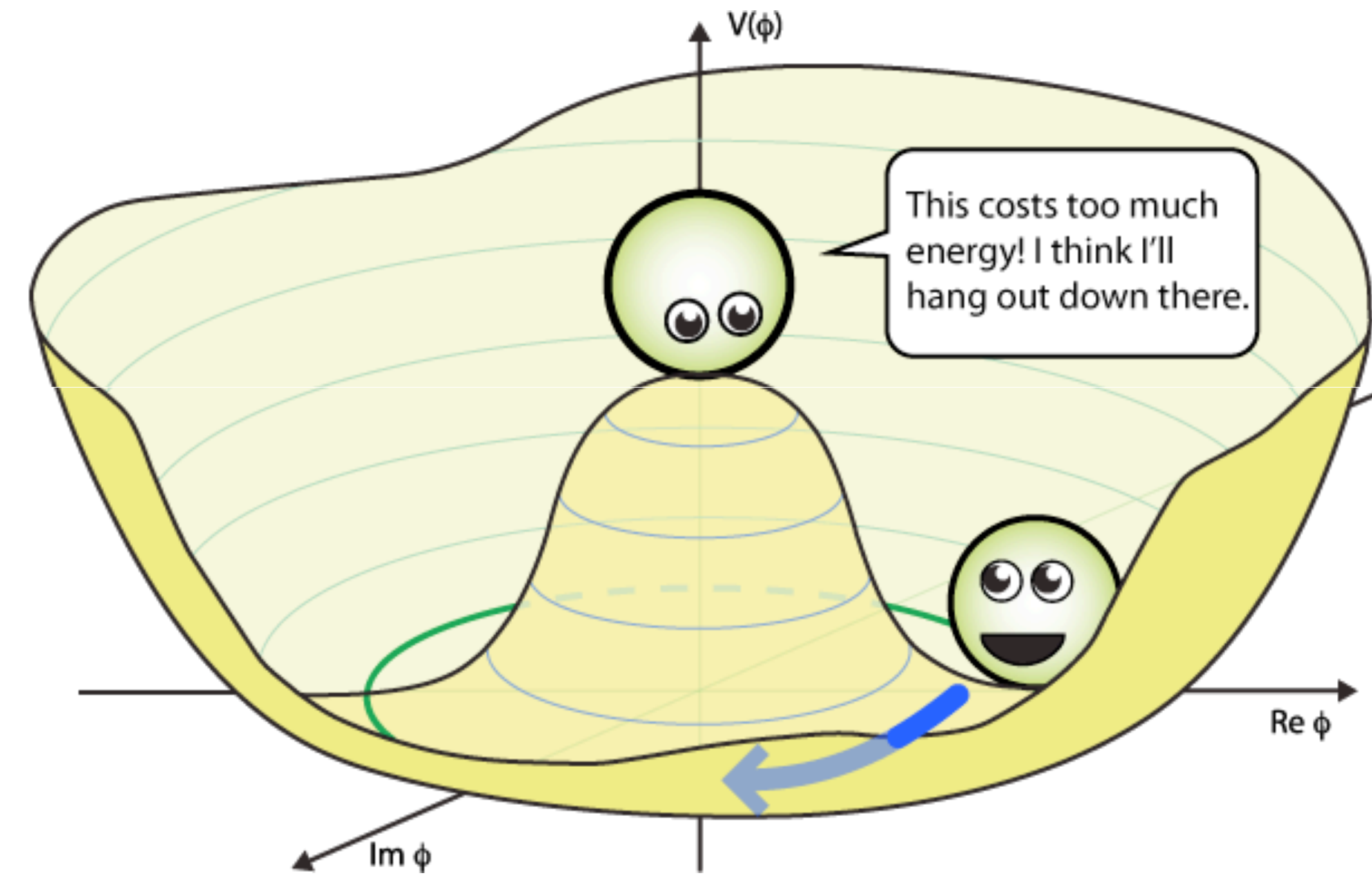


The Brout-Englert-Higgs (BEH) mechanism

"In one slide"

The economical¹⁾ way to endow fundamental particles with mass while keeping the theory gauge invariant and predictive

The field is responsible for the spontaneous breaking of electroweak symmetry



PRL 13, 321-323 (1964) Englert and Brout
PRL 13, 508-509 (1964) Higgs
PRL 13, 585-587 (1964) Guralnik, Hagen, Kibble

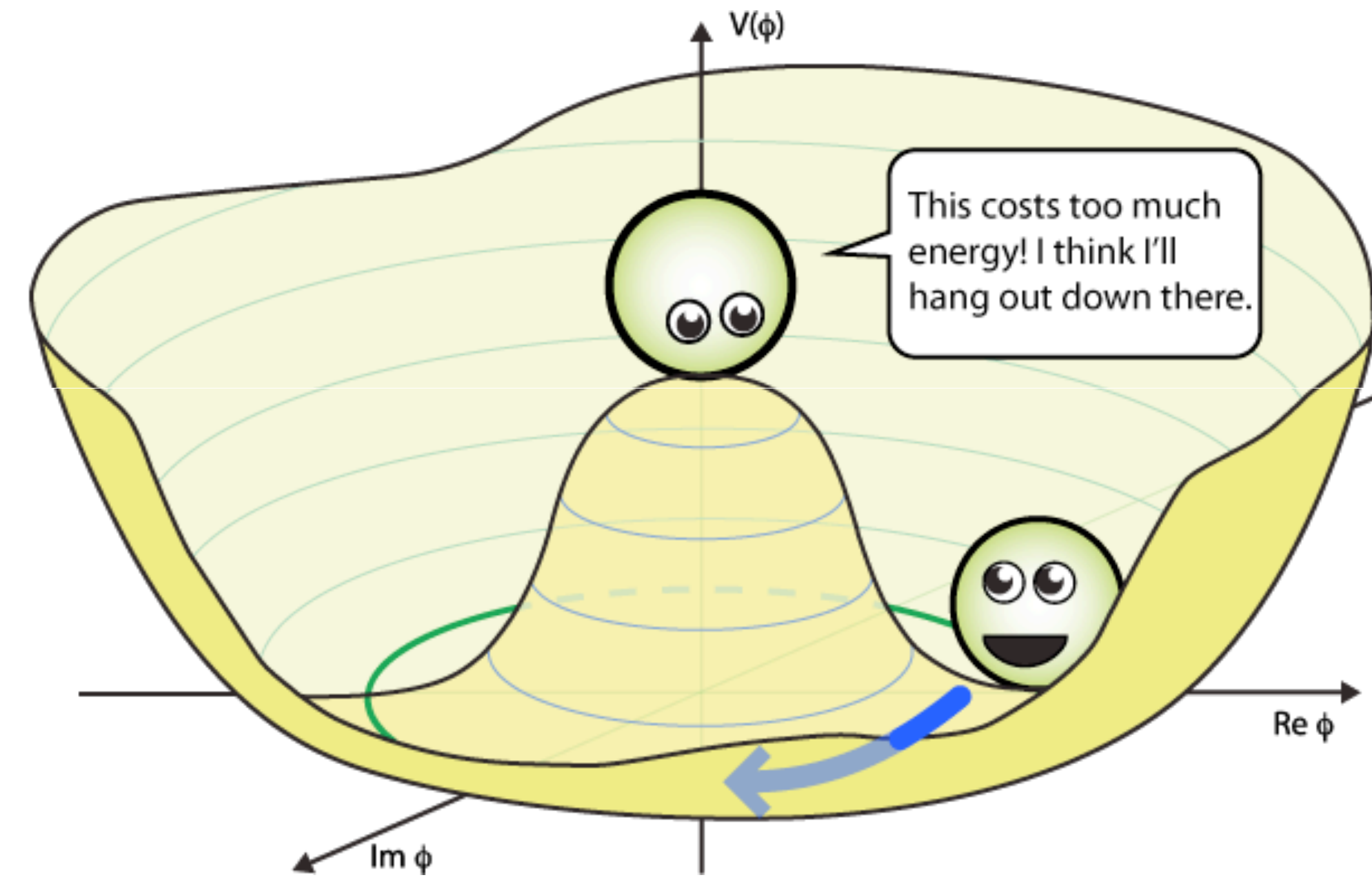
1) less economical (Higgs doublets, families of Higgses, ...)
or more complicated (Higgs-less solutions, Technicolor, ...) routes exist

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"Only"

1 new particle	the Higgs boson (H)
1 unknown	the Higgs boson mass (m_H)

1) less economical (Higgs doublets, families of Higgses, ...) or more complicated (Higgs-less solutions, Technicolor, ...) routes exist

The Higgs boson is special

It is a fundamental scalar particle (spin 0) and its theory is unlike anything else has been seen in Nature!

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi + h.c. \\ & + \bar{\psi}_i y_{ij} \psi_j \phi + h.c. \\ & + |D_\mu \phi|^2 - V(\phi) \end{aligned}$$

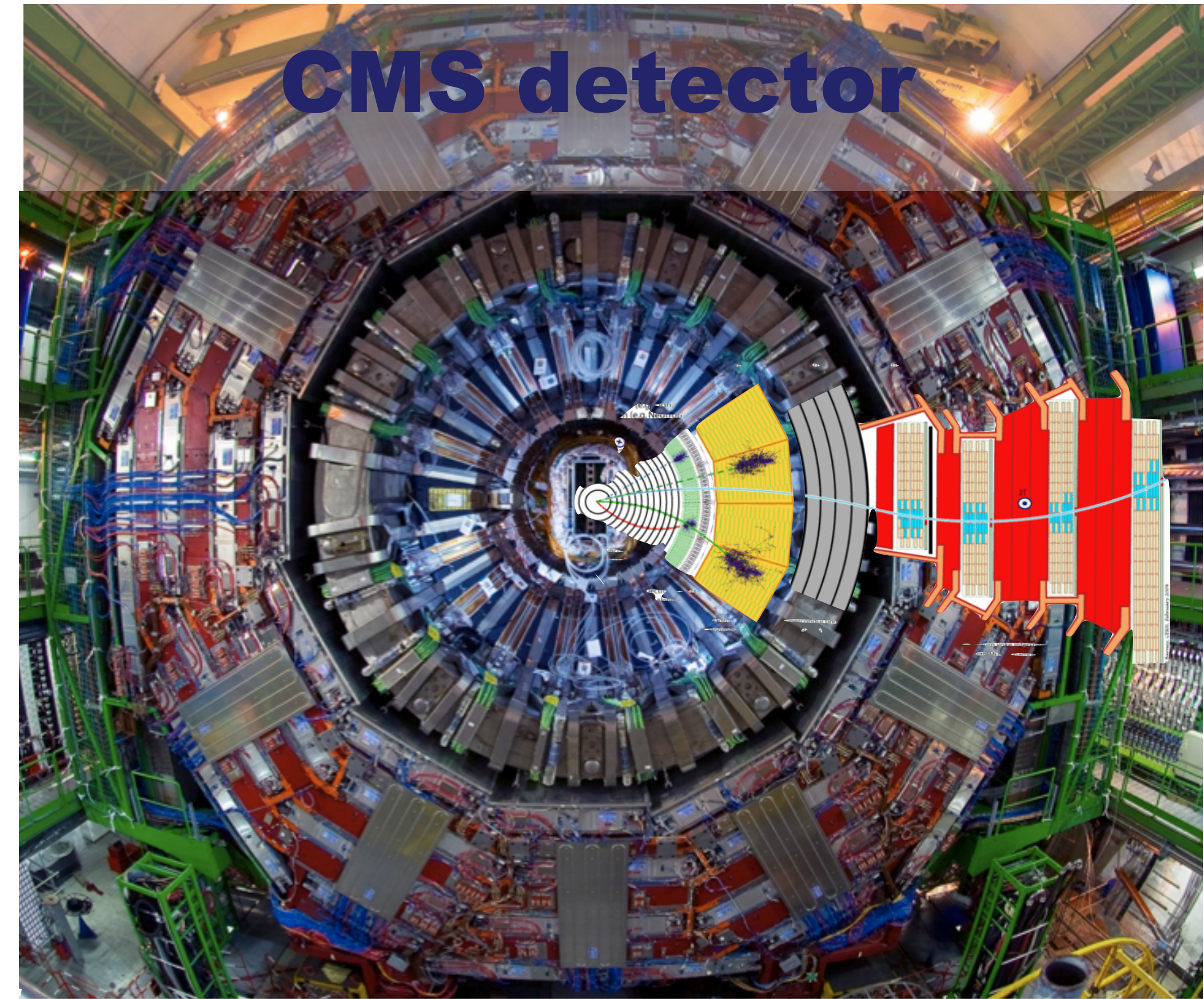
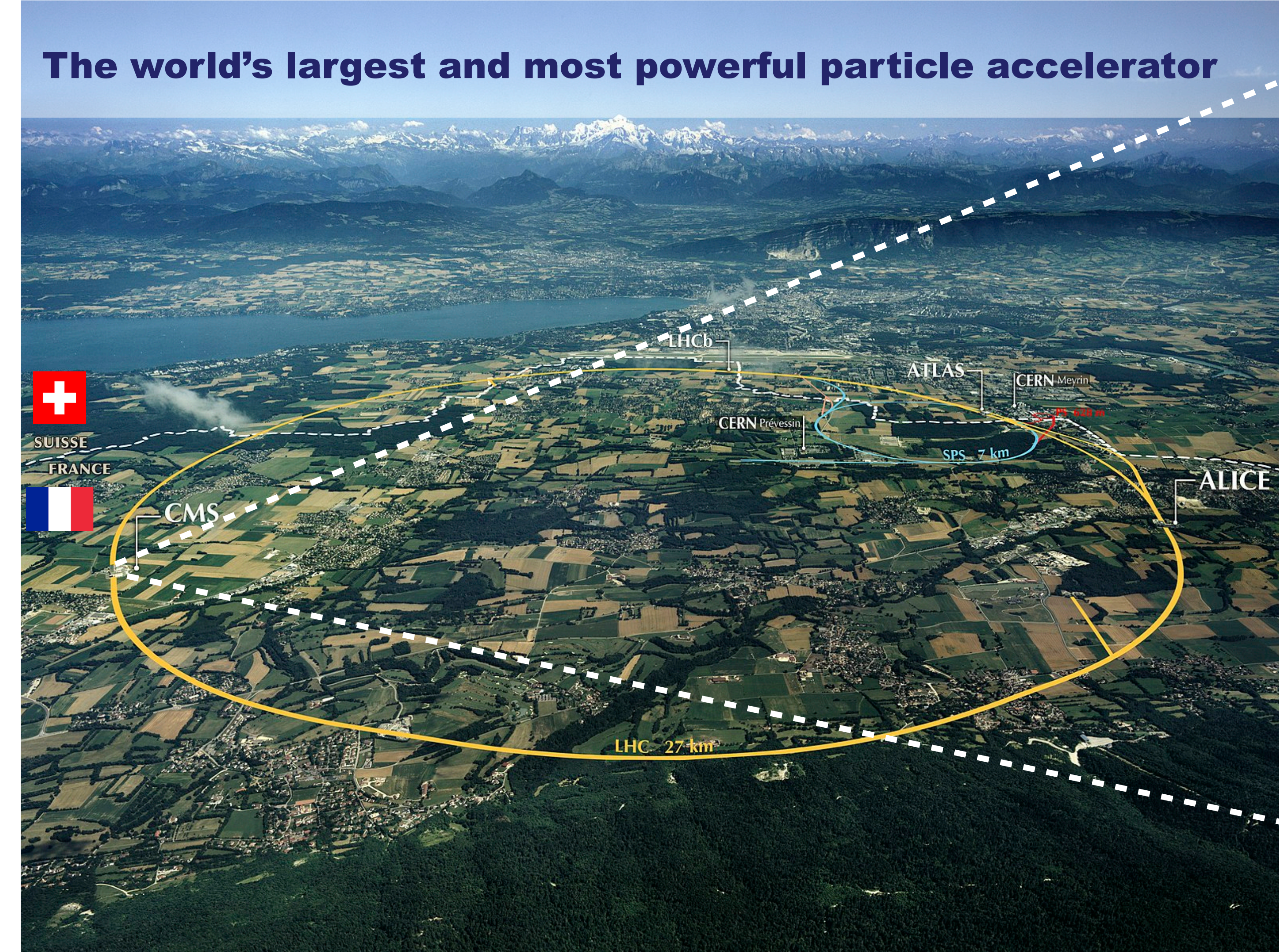
A gauge interaction with vector bosons

A Yukawa interaction with the fermions

A potential $V(\phi) \sim -\mu^2(\phi\phi^\dagger) + \lambda(\phi\phi^\dagger)^2$
the keystone of the BEH mechanism and SM

LHC : a new dimension in particle physics

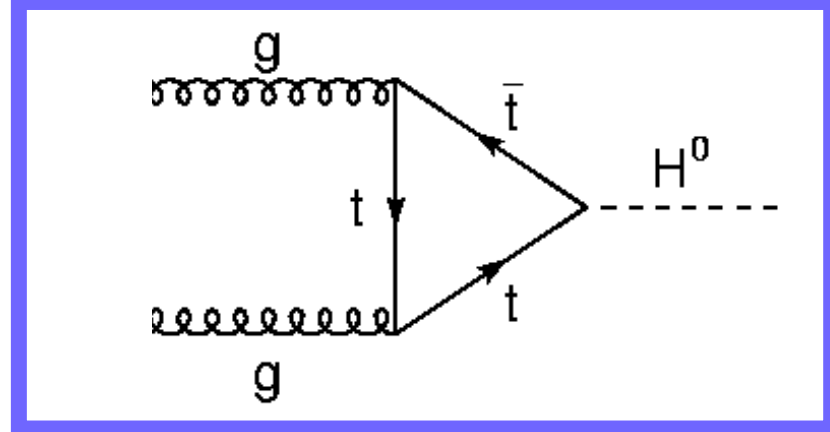
The world's largest and most powerful particle accelerator



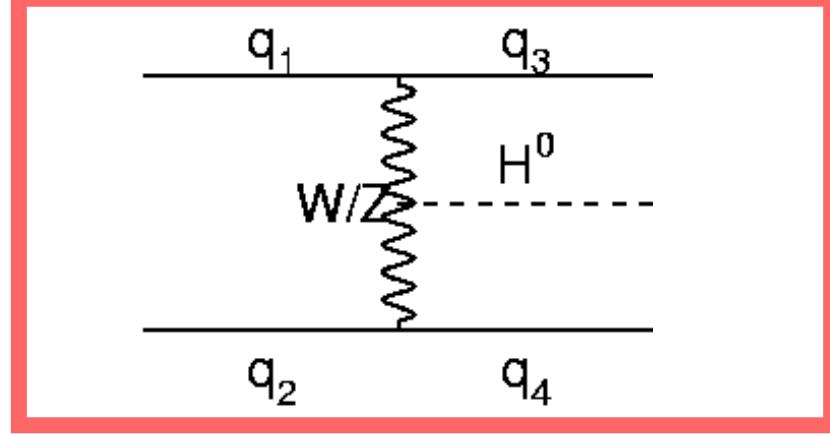
14000-tonne weight
21 metres long, 15 metres wide and 15 metres high
4 Tesla field ($\sim 10^6$ times the magnetic field of the Earth)

The Higgs boson production and decay "just a reminder"

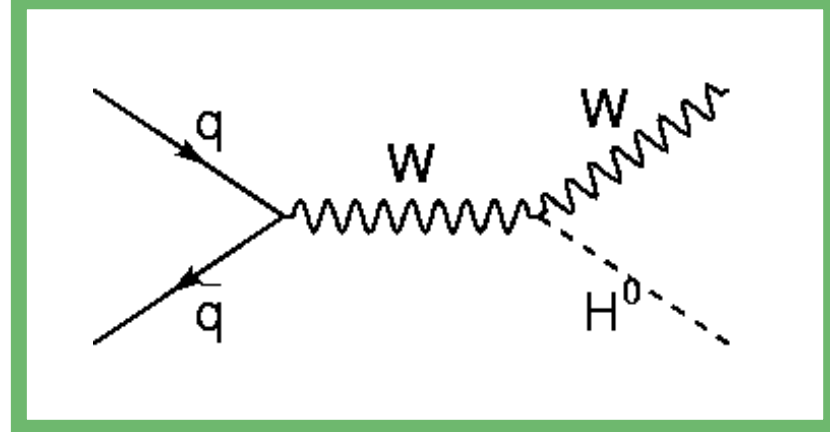
$\sigma=49 \text{ pb} / 6.9\text{M Higgs in } 140\text{fb}^{-1}$



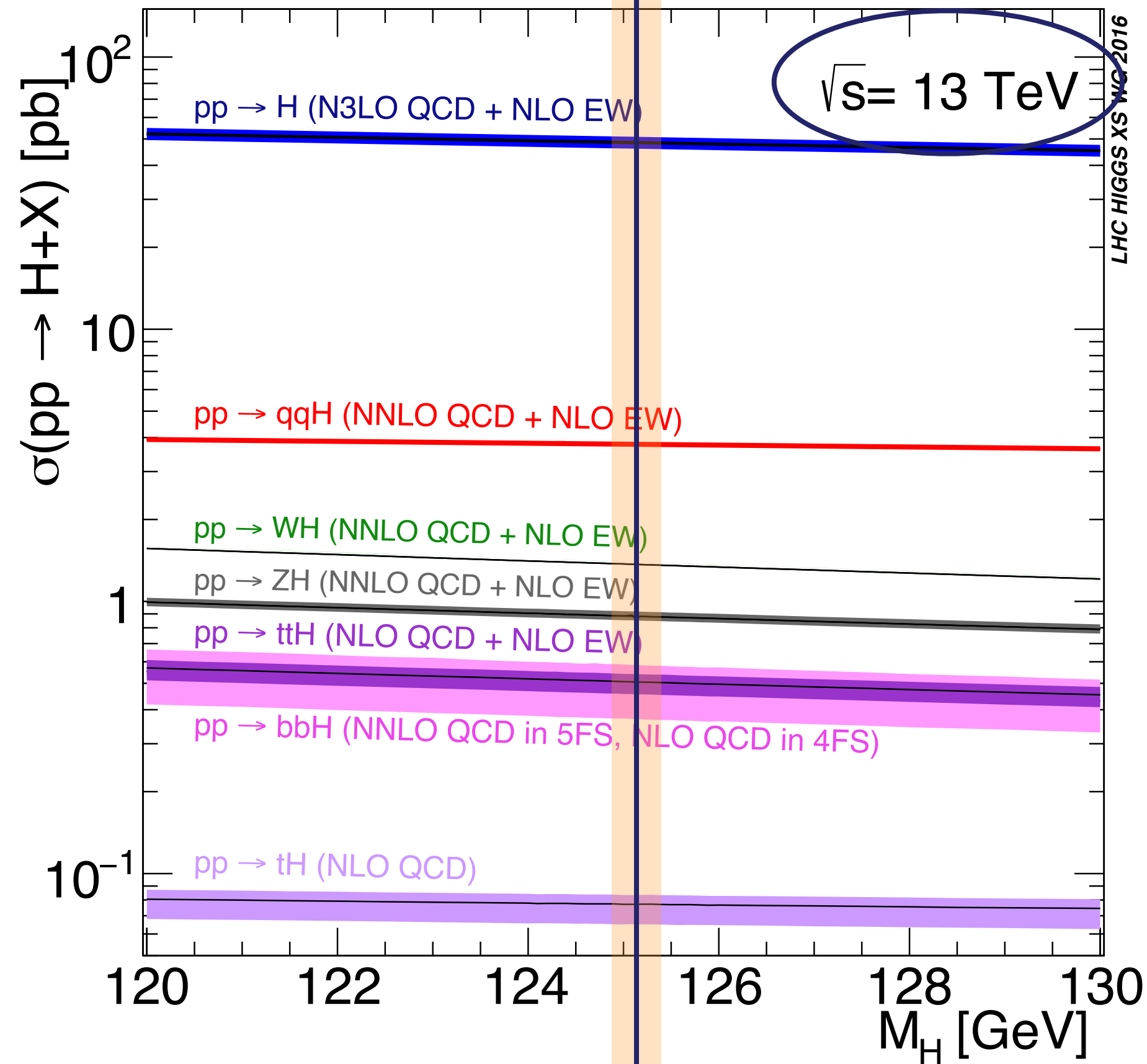
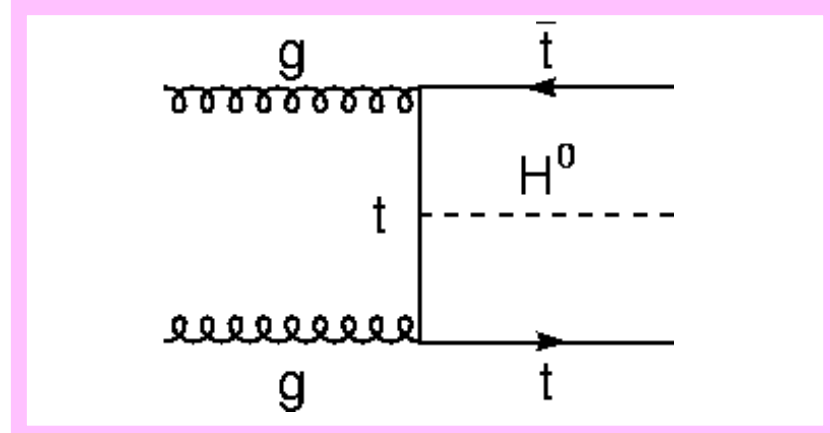
$\sigma=3.8 \text{ pb} / 520\text{k Higgs in } 140\text{fb}^{-1}$



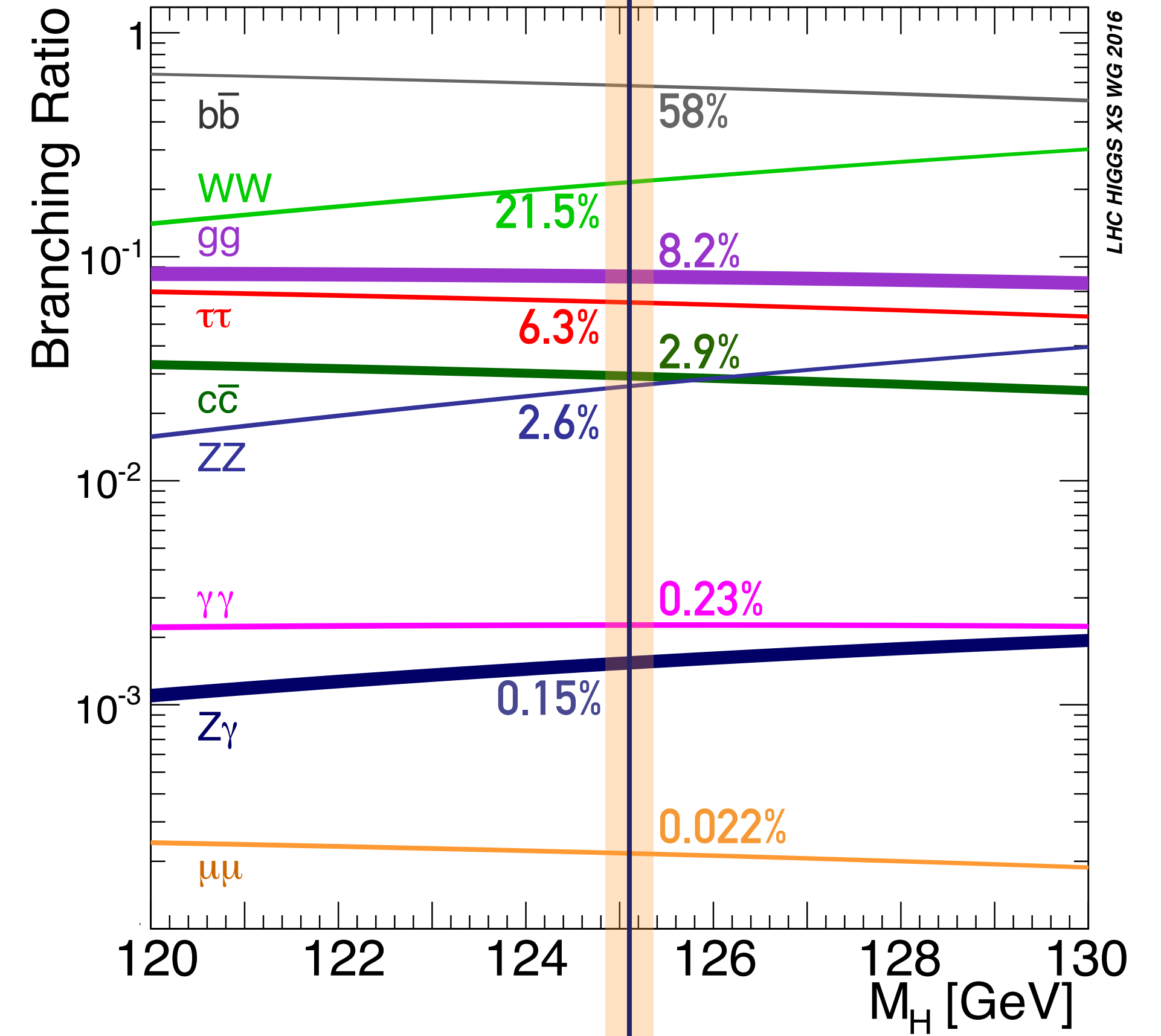
$\sigma=2.3 \text{ pb} / 320\text{k Higgs in } 140\text{fb}^{-1}$



$\sigma=0.5 \text{ pb} / 70\text{k Higgs in } 140\text{fb}^{-1}$



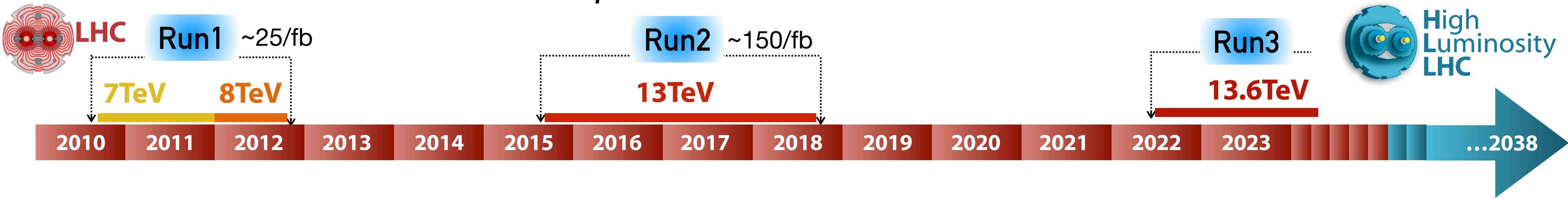
$125.09 \pm 0.24 \text{ GeV}$
LHC Run1 measurement



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LHC Run1 measurement

The Higgs boson timeline at LHC

Years of unprecedented moments in HEP



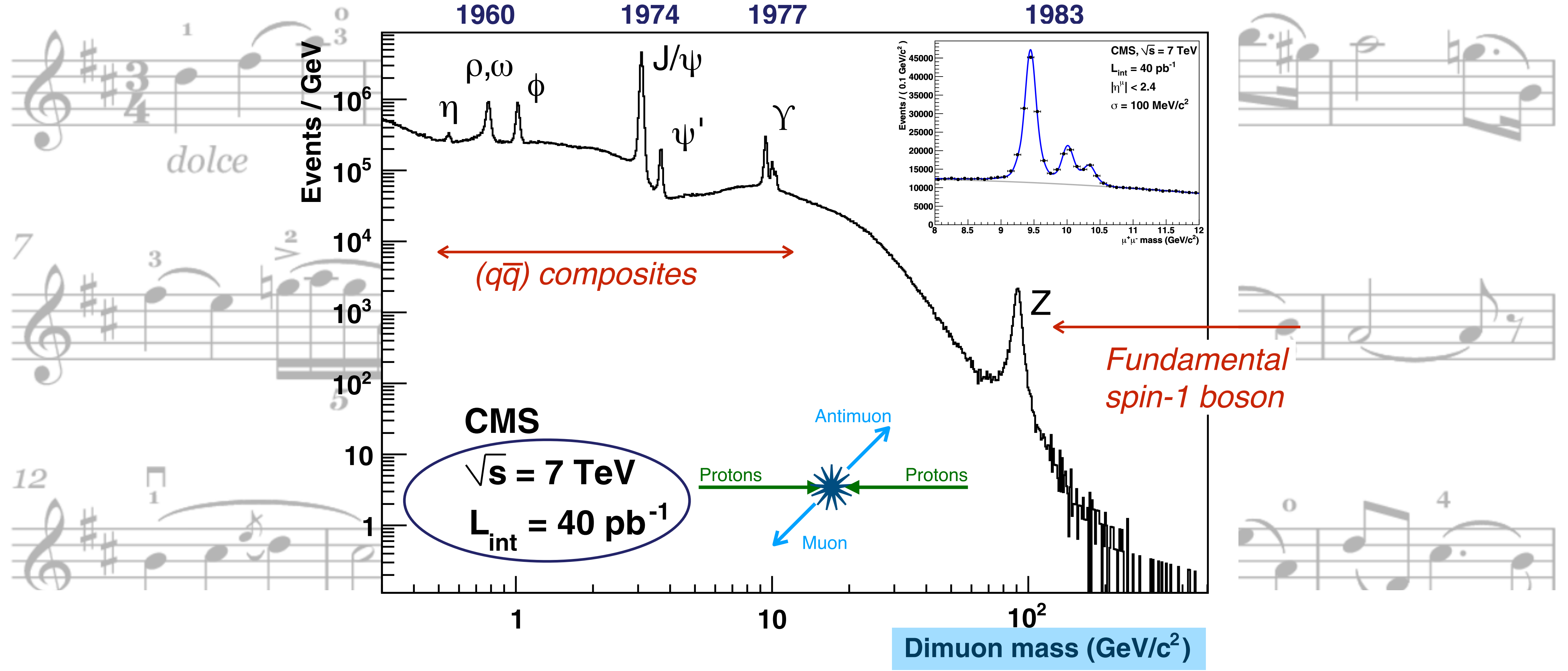
“Intermezzo”

In 2010 LHC started to deliver proton-proton collisions

The image displays a musical score for a piece titled "Intermezzo". The score is written in G major (one sharp) and 3/4 time. It consists of three staves of music, with measure numbers 1, 7, and 12 indicated at the beginning of each line. The first staff starts with the instruction "dolce". The second staff includes a "pp" (pianissimo) dynamic marking. The score is annotated with guitar-specific notation, including fingering numbers (0-4) and vibrato marks (v) above notes. A double bar line with a repeat sign is present between measures 6 and 7. The music features a mix of eighth and quarter notes, often beamed together, and includes some slurs and accents.

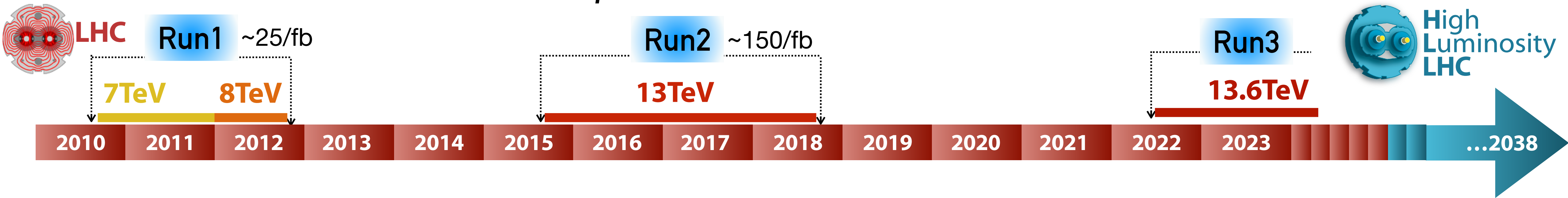
“Intermezzo”

In 2010 LHC started to deliver proton-proton collisions
 50 years of particle physics ... in few weeks of data taking



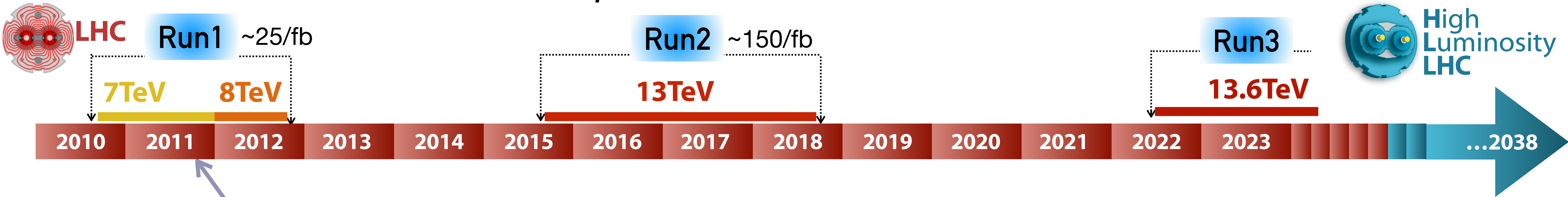
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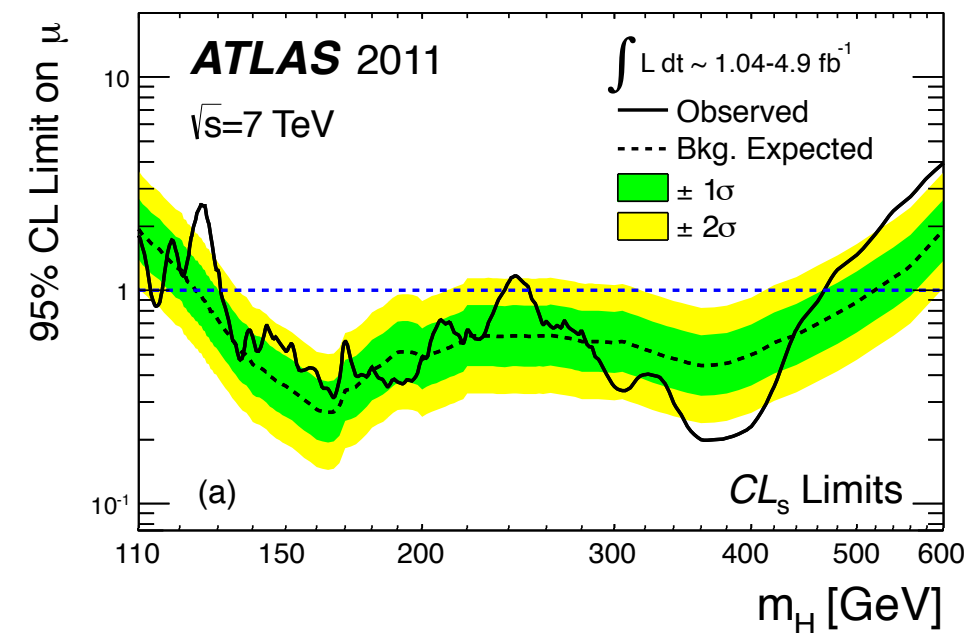
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Higgs boson searches

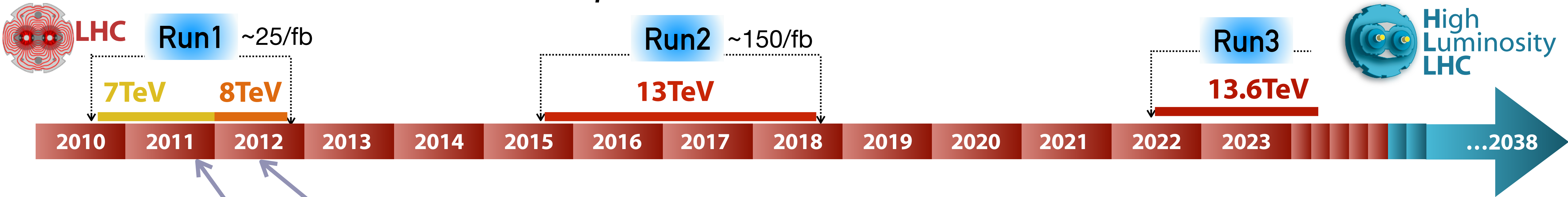
"if the SM Higgs boson exists, is most likely to have a mass constrained to 115-130 GeV "



Few highlights

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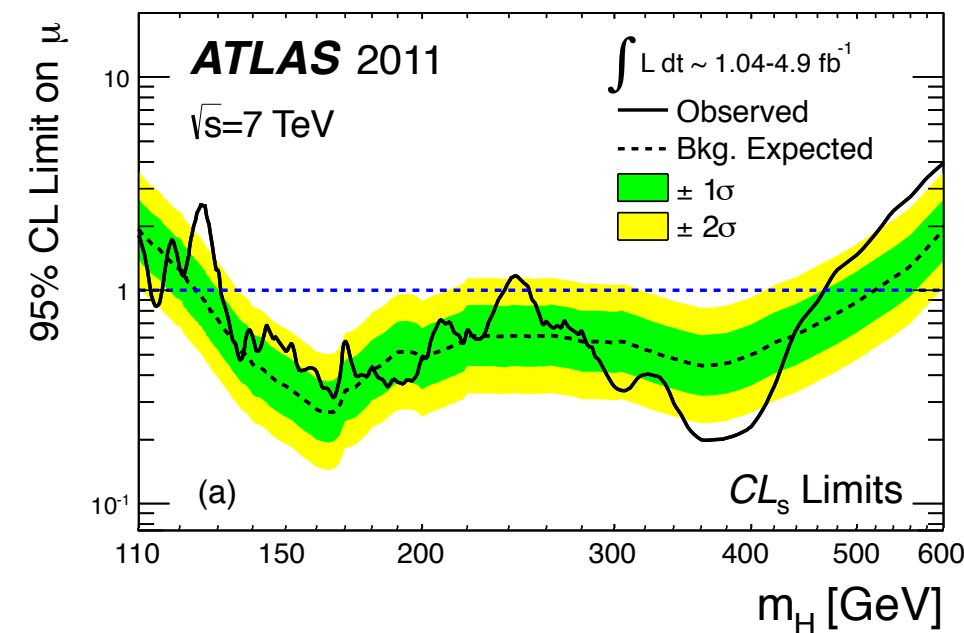
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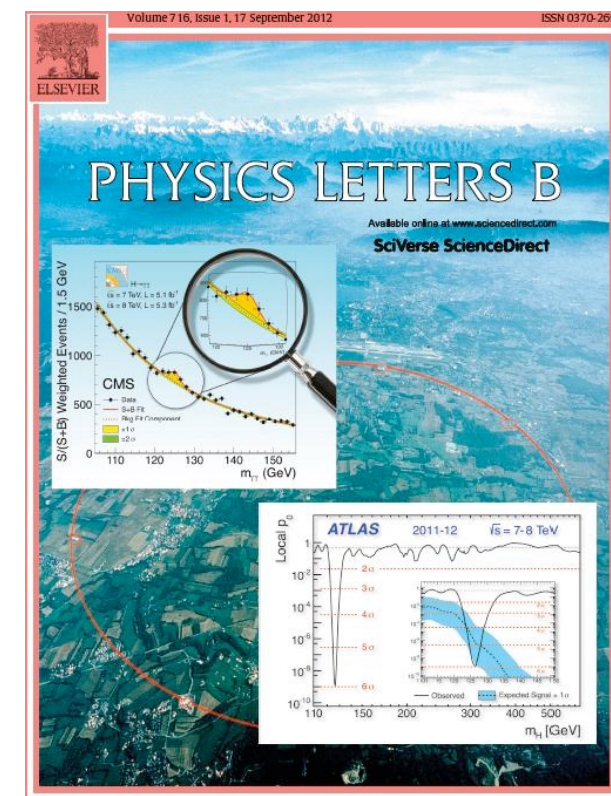
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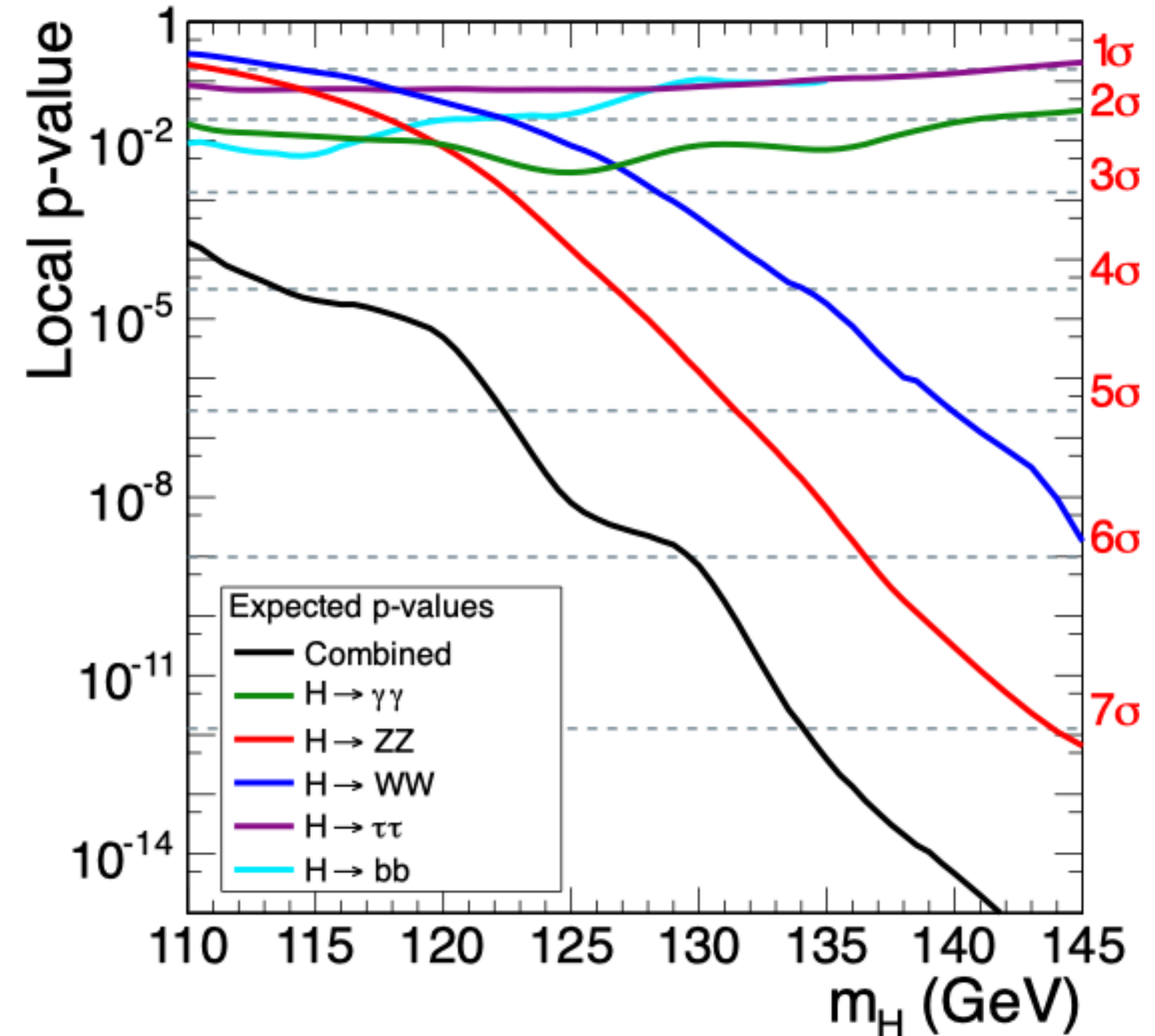
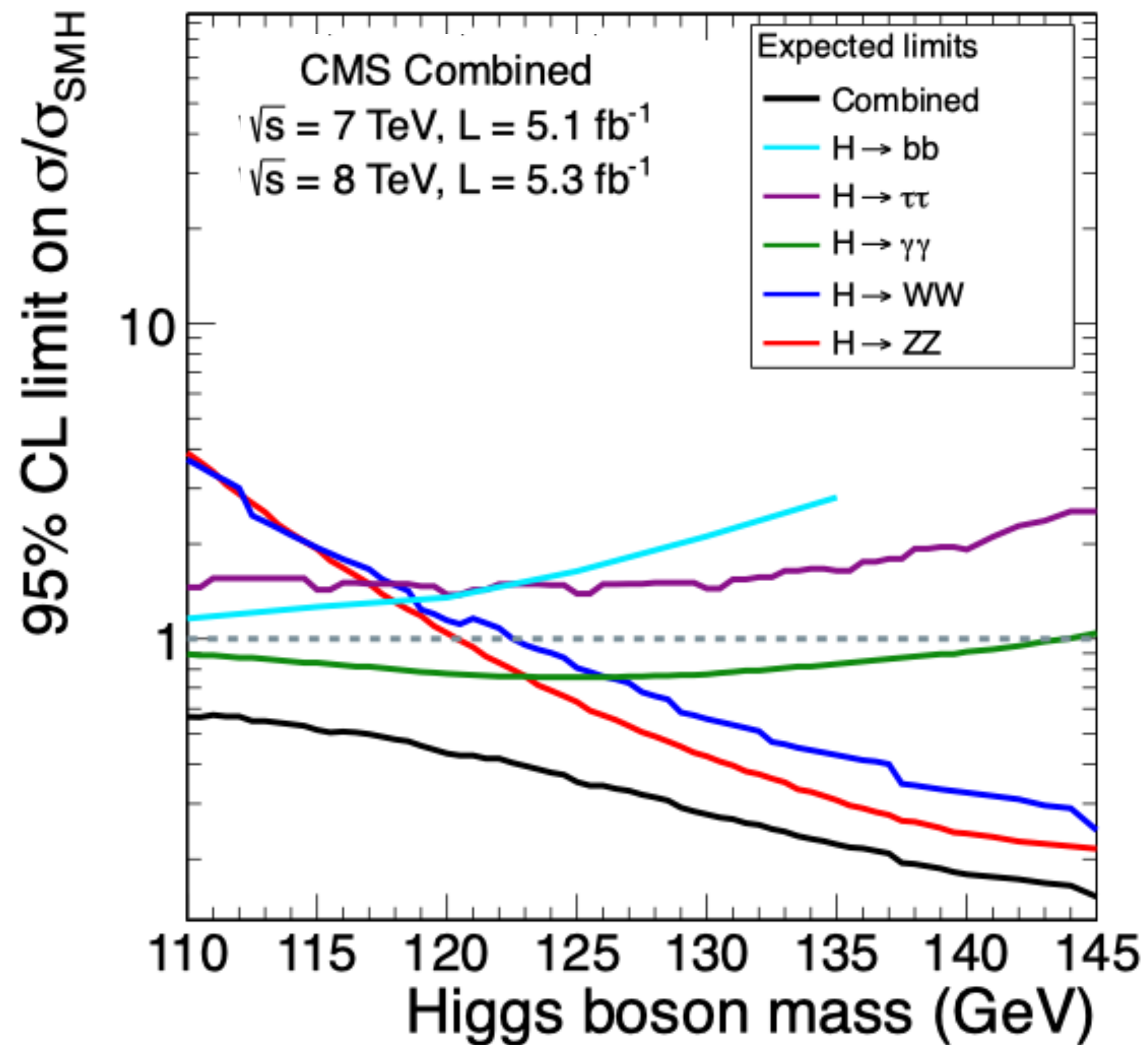


Higgs boson observation



The big five

Expected results



High mass resolution decay modes:

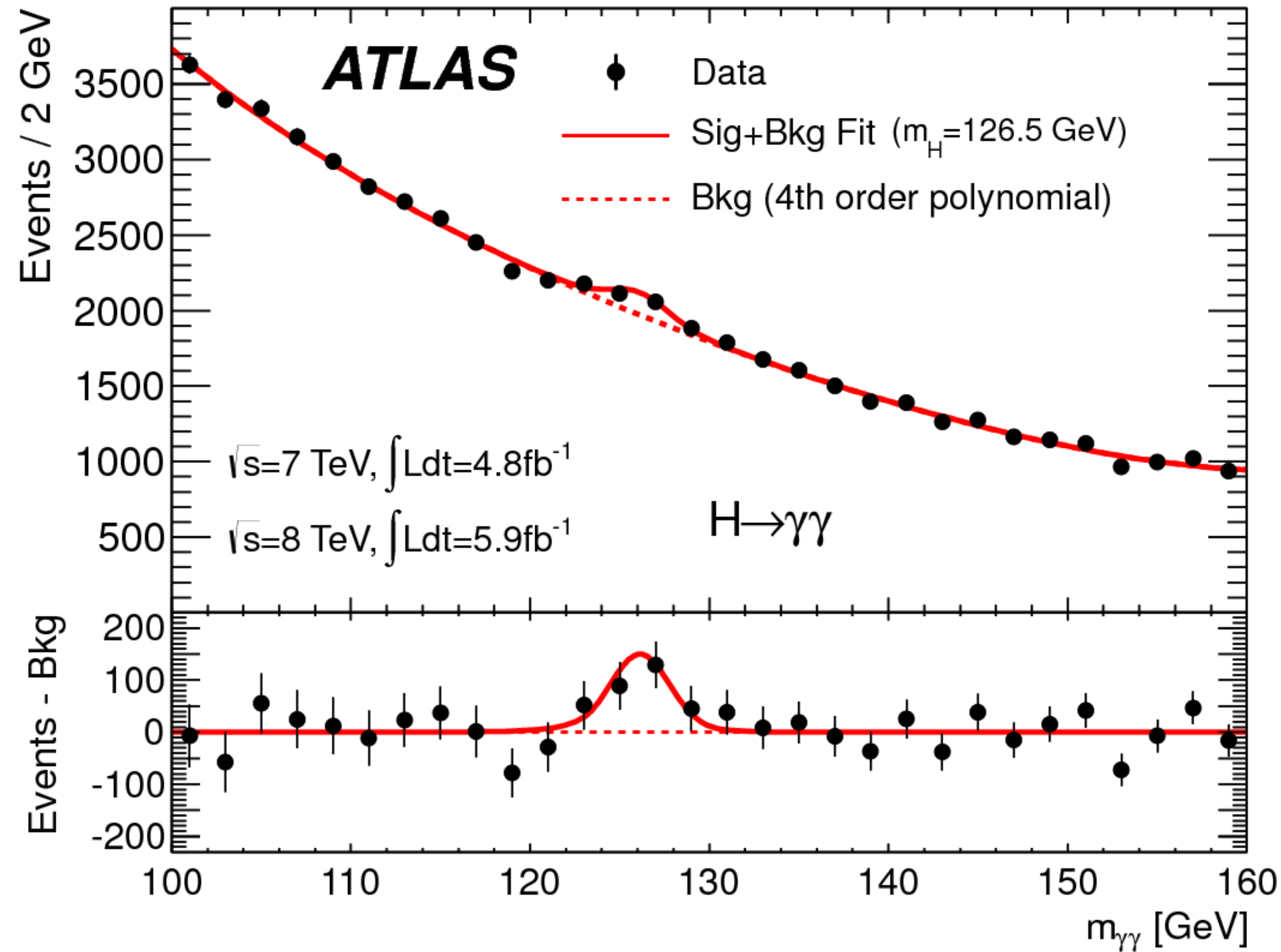
$H \rightarrow ZZ$ / $H \rightarrow \gamma\gamma$

Lower mass resolution decay modes:

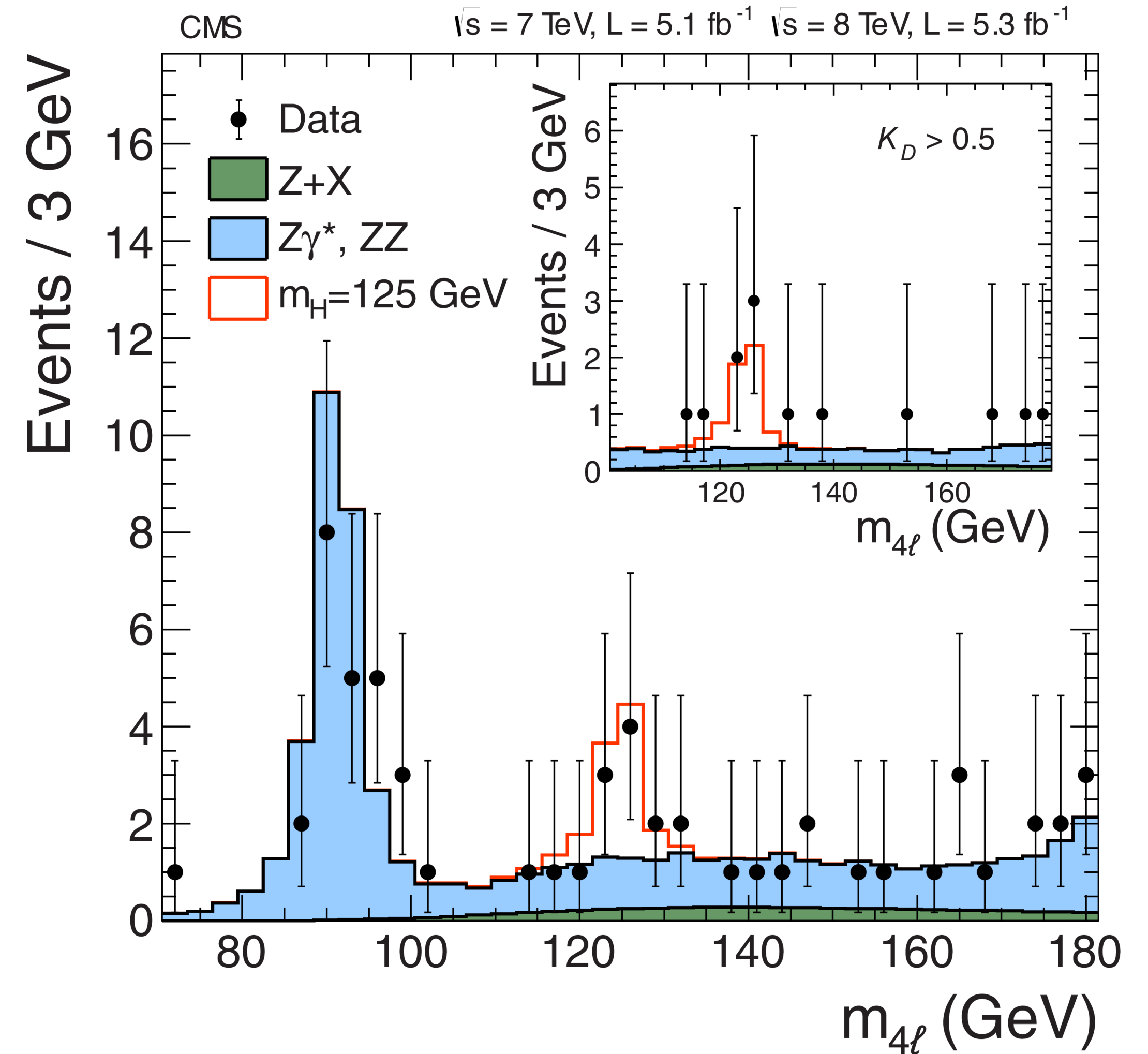
$H \rightarrow WW$ / $H \rightarrow bb$ / $H \rightarrow \tau\tau$

Discovery of a new particle

$H \rightarrow \gamma\gamma$

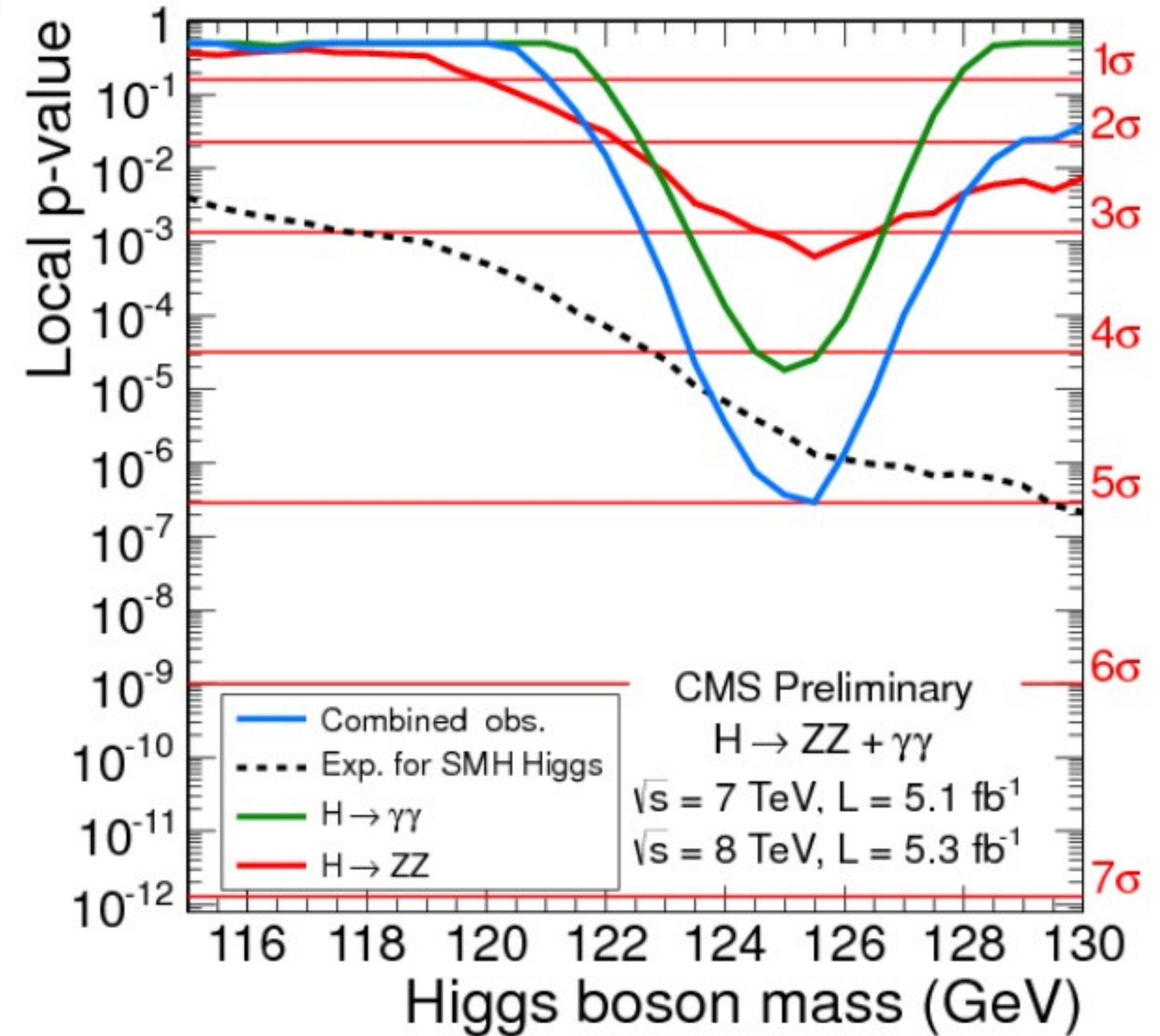
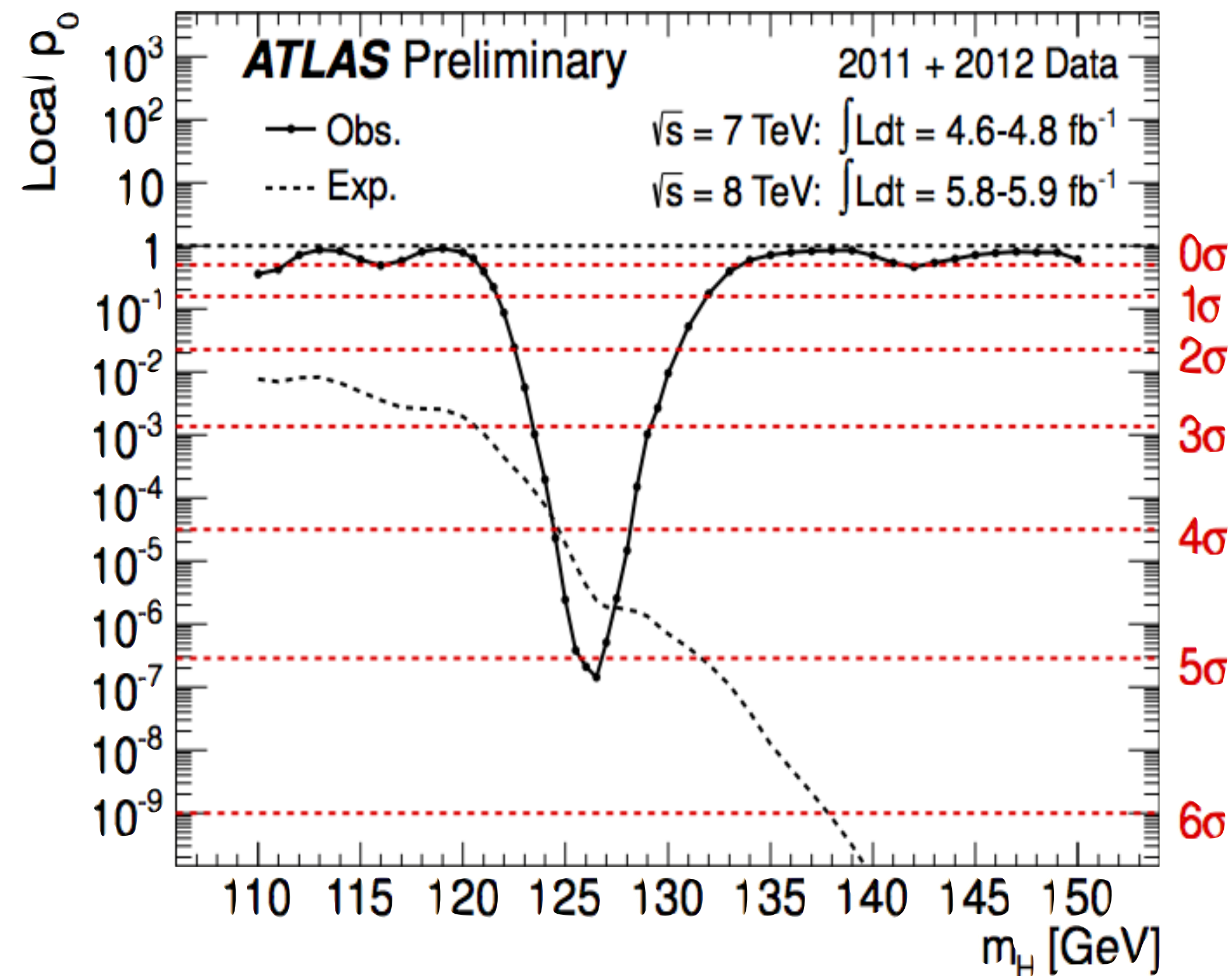


$H \rightarrow ZZ$



Discovery of a new particle

The open boxes on July 4th 2012



Discovery of a new particle



Discovery of New Particle Could Redefine Physical World

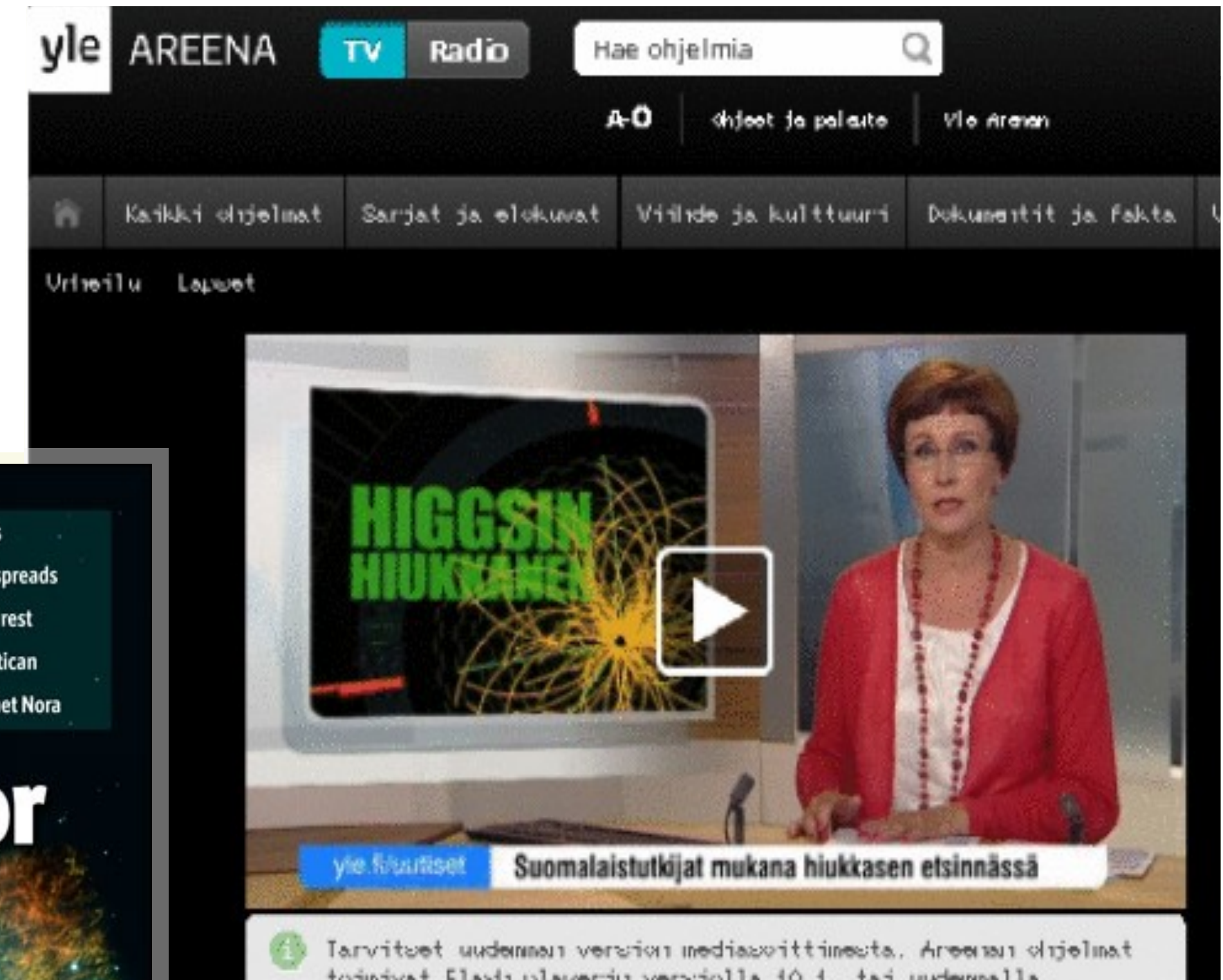
By DENNIS OVERBYE
21 minutes ago

The discovery by physicists at CERN's Large Hadron Collider, if confirmed to be the Higgs boson particle, could lead to a new understanding of how the universe began.

• The Lede Blog: What in the World Is a Higgs Boson?
4:16 AM ET



Fabrice Coffrini/Agence France-Presse — Getty Images



More than a billion people saw the news on TV, on five thousand broadcasts on a thousand TV stations

Ten thousand news articles in a hundred countries

Discovery of a new particle

July 2013

European Physics Society HEPP Prize

“**Discovery** of a Higgs boson, a scalar particle whose associated field breaks electroweak symmetry and generates mass”
awarded to the **ATLAS and CMS Collaborations**

October 2013

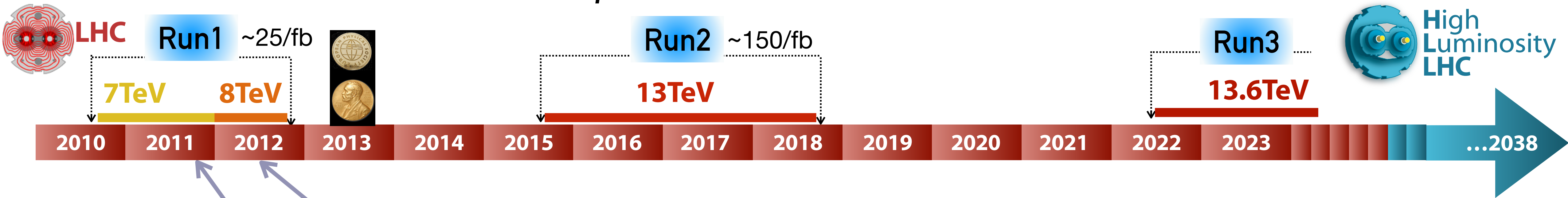
Nobel Prize

“For the **theoretical discovery** of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, **by the ATLAS and CMS experiments at CERN's LHC**”
awarded jointly to **François Englert and Peter W. Higgs**



The Higgs boson timeline at LHC

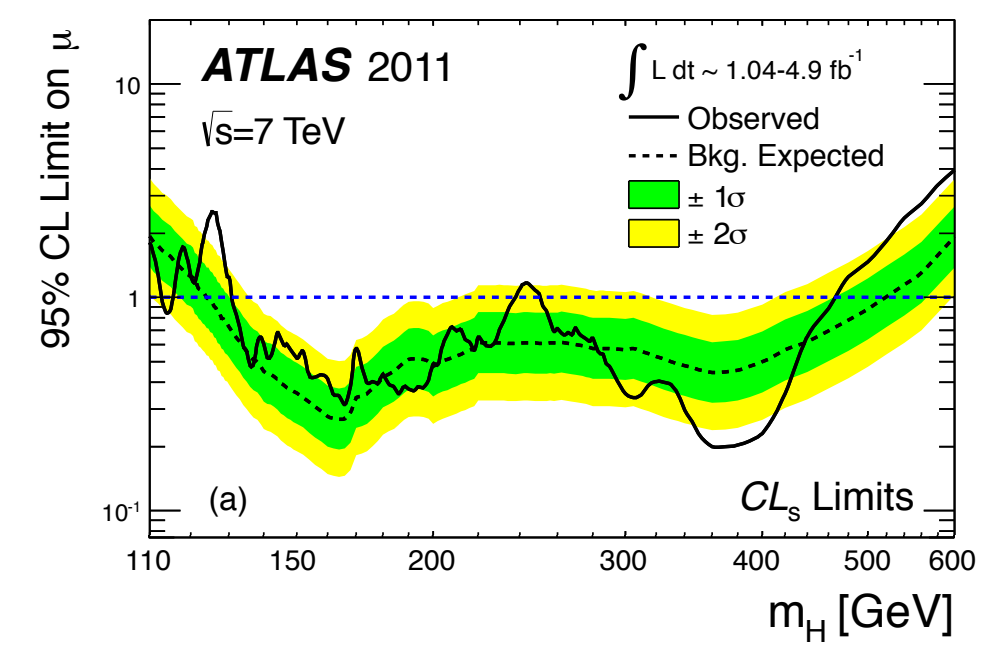
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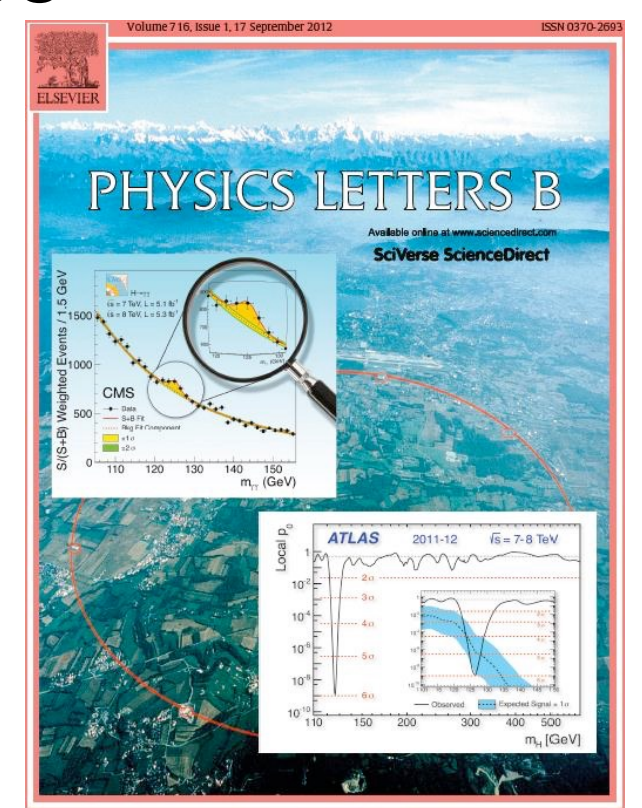
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"if the SM Higgs boson exists, is most likely to have a mass constrained to 115-130 GeV"

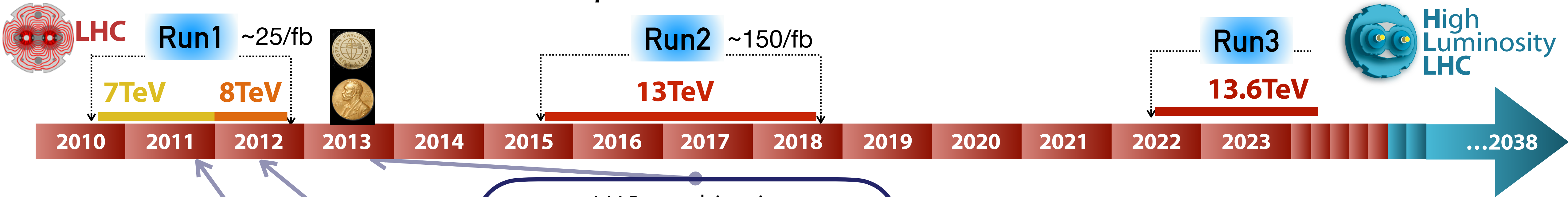


Higgs boson observation



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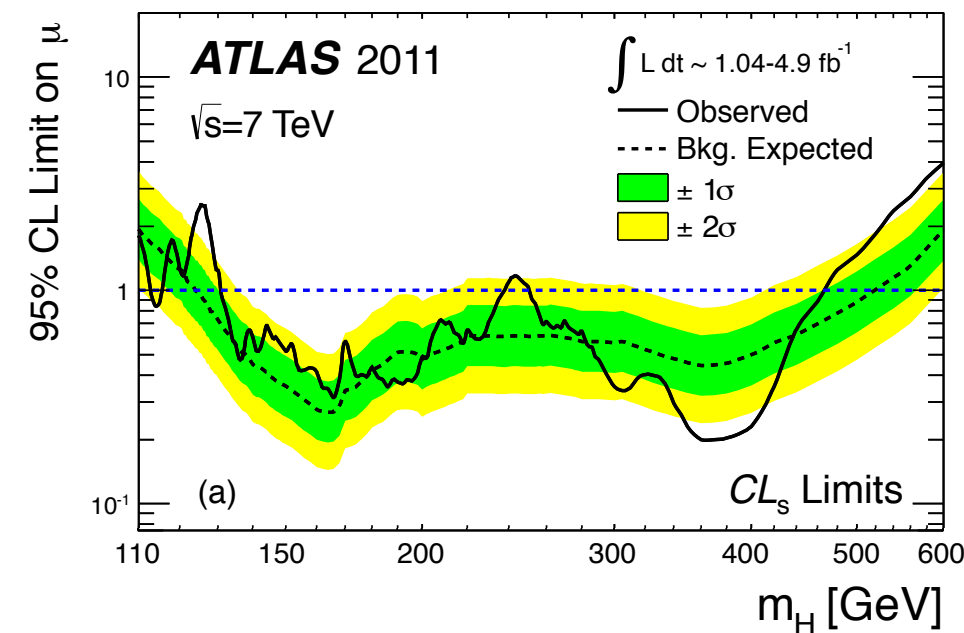
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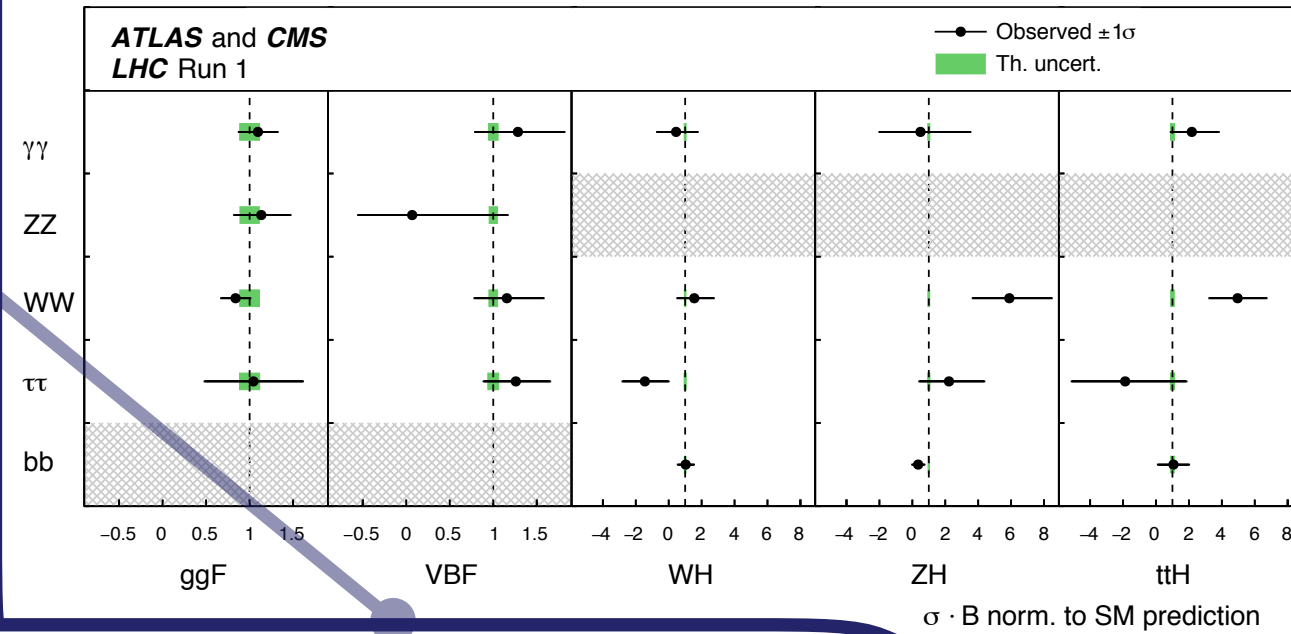
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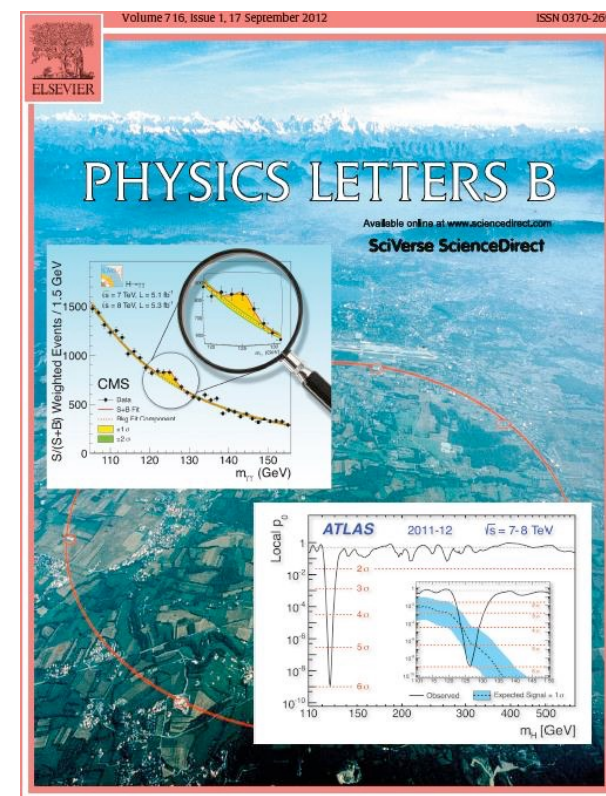


LHC combination

In one word "SM like"

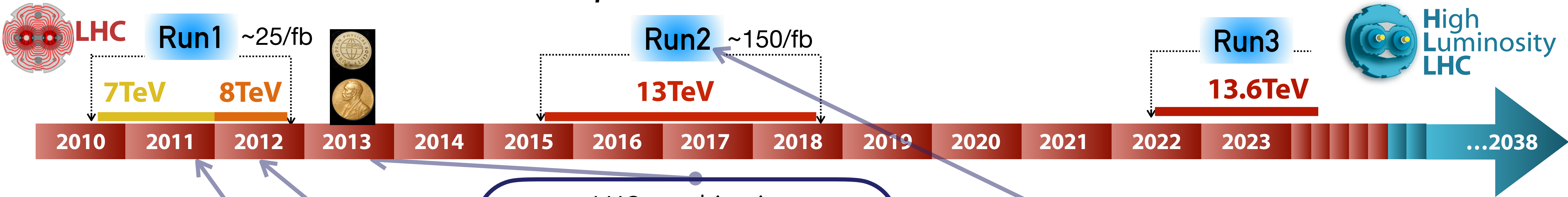


Higgs boson observation



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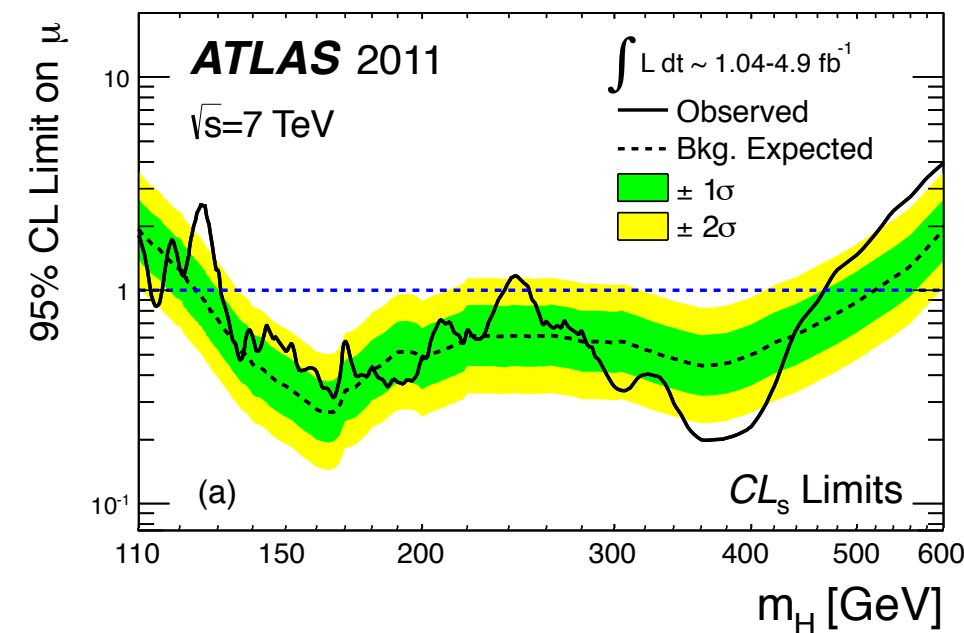
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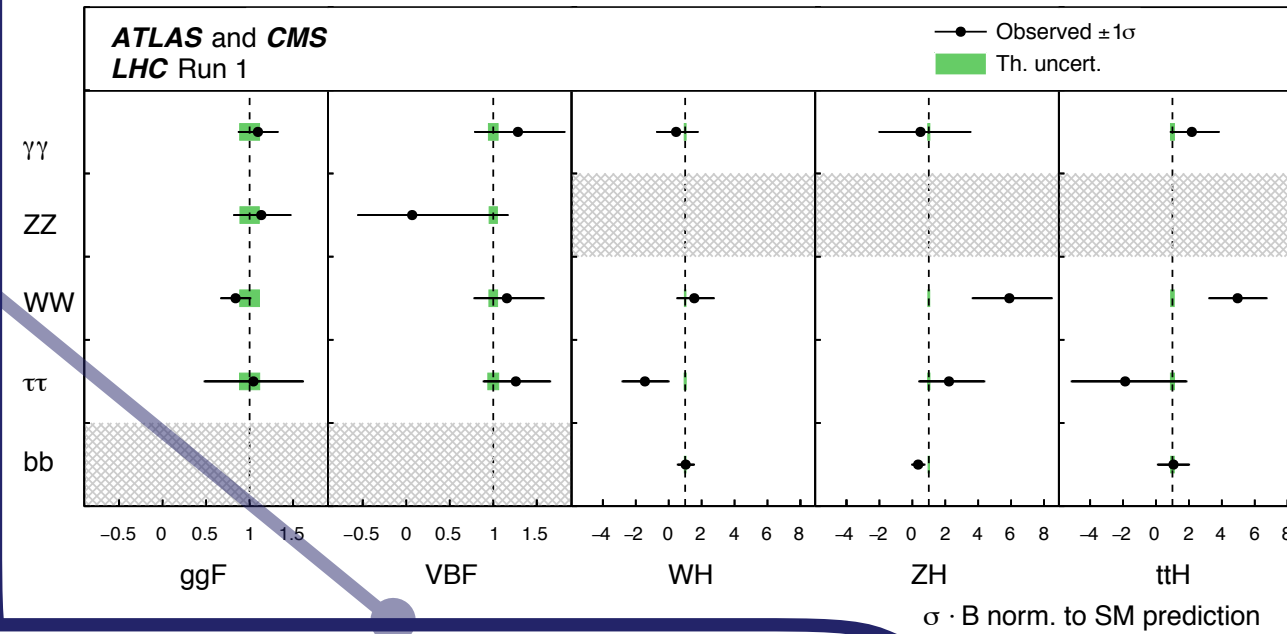
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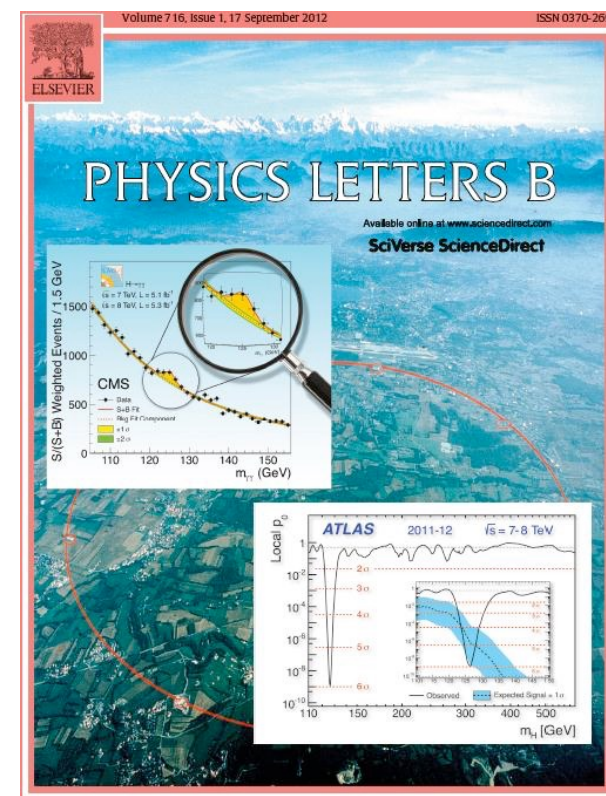


LHC combination

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Higgs boson observation



Towards the precision physics era

A huge landscape of possibilities opens



A huge landscape of possibilities opens



Complete characterisation of the Higgs boson profile

(exploiting both H-Z, H-W, and H- γ interactions
and its couplings with 3rd/2nd-family fermion)

A huge landscape of possibilities opens



Study of the Higgs boson self-coupling

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Study of the Higgs boson self-coupling

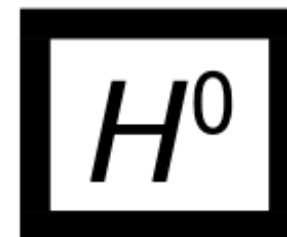


The Higgs boson as a tool to reveal the mysteries of Universe

(Dark matter, BSM,)

The experimental H profile

Citation: R.L. Workman *et al.* (Particle Data Group), Prog.Theor.Exp.Phys. **2022**, 083C01 (2022)



$$J = 0$$



In the following H^0 refers to the signal that has been discovered in the Higgs searches. Whereas the observed signal is labeled as a spin 0 particle and is called a Higgs Boson, the detailed properties of H^0 and its role in the context of electroweak symmetry breaking need to be further clarified. These issues are addressed by the measurements listed below.

Concerning mass limits and cross section limits that have been obtained in the searches for neutral and charged Higgs bosons, see the sections “Searches for Neutral Higgs Bosons” and “Searches for Charged Higgs Bosons (H^\pm and $H^{\pm\pm}$)”, respectively.

H^0 MASS

<u>VALUE (GeV)</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
125.25 ± 0.17 OUR AVERAGE	Error includes scale factor of 1.5. See the ideogram below.		
125.46 ± 0.16	¹ SIRUNYAN	20L CMS	$pp, 13 \text{ TeV}, \gamma\gamma, ZZ^* \rightarrow 4\ell$

The Higgs boson mass: “is not” vs. “is”

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- ✓ The Higgs boson mass **is not** a prediction of the theory but the Higgs boson mass **is** free input parameter of the theory

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- ✓ The Higgs boson mass measurement **is not** a test of the SM but the Higgs boson mass **is** an important¹⁾ ingredient in SM predictions of many “ m_H -dependent” SM observables:
 - **Higgs boson observables : couplings, branching ratios, width**
 - **Electroweak observables : mass of the W boson, mass of the top quark, effective weak mixing angle,**

1) I would argue that is the most important

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 - **Higgs boson observables : couplings, branching ratios, width**
 - **Electroweak observables : mass of the W boson, mass of the top quark, effective weak mixing angle,**
- ✓ The Higgs boson mass value **is** connected to the Fermi and the Planck scales and ultimately with the vacuum stability

1) I would argue that is the most important

Approximated parametrisation for the mass of W boson (m_W) :

$$M_W = M_W^0 - c_1 dH - c_2 dH^2 + c_3 dH^4 + c_4(dh - 1) - c_5 d\alpha + c_6 dt - c_7 dt^2 \\ - c_8 dH dt + c_9 dh dt - c_{10} d\alpha_s + c_{11} dZ,$$

where

$$\boxed{dH = \ln\left(\frac{M_H}{100 \text{ GeV}}\right), \quad dh = \left(\frac{M_H}{100 \text{ GeV}}\right)^2}$$
$$dt = \left(\frac{m_t}{174.3 \text{ GeV}}\right)^2 - 1,$$
$$dZ = \frac{M_Z}{91.1875 \text{ GeV}} - 1, \quad d\alpha = \frac{\Delta\alpha}{0.05907} - 1, \quad d\alpha_s = \frac{\alpha_s(M_Z)}{0.119} - 1,$$

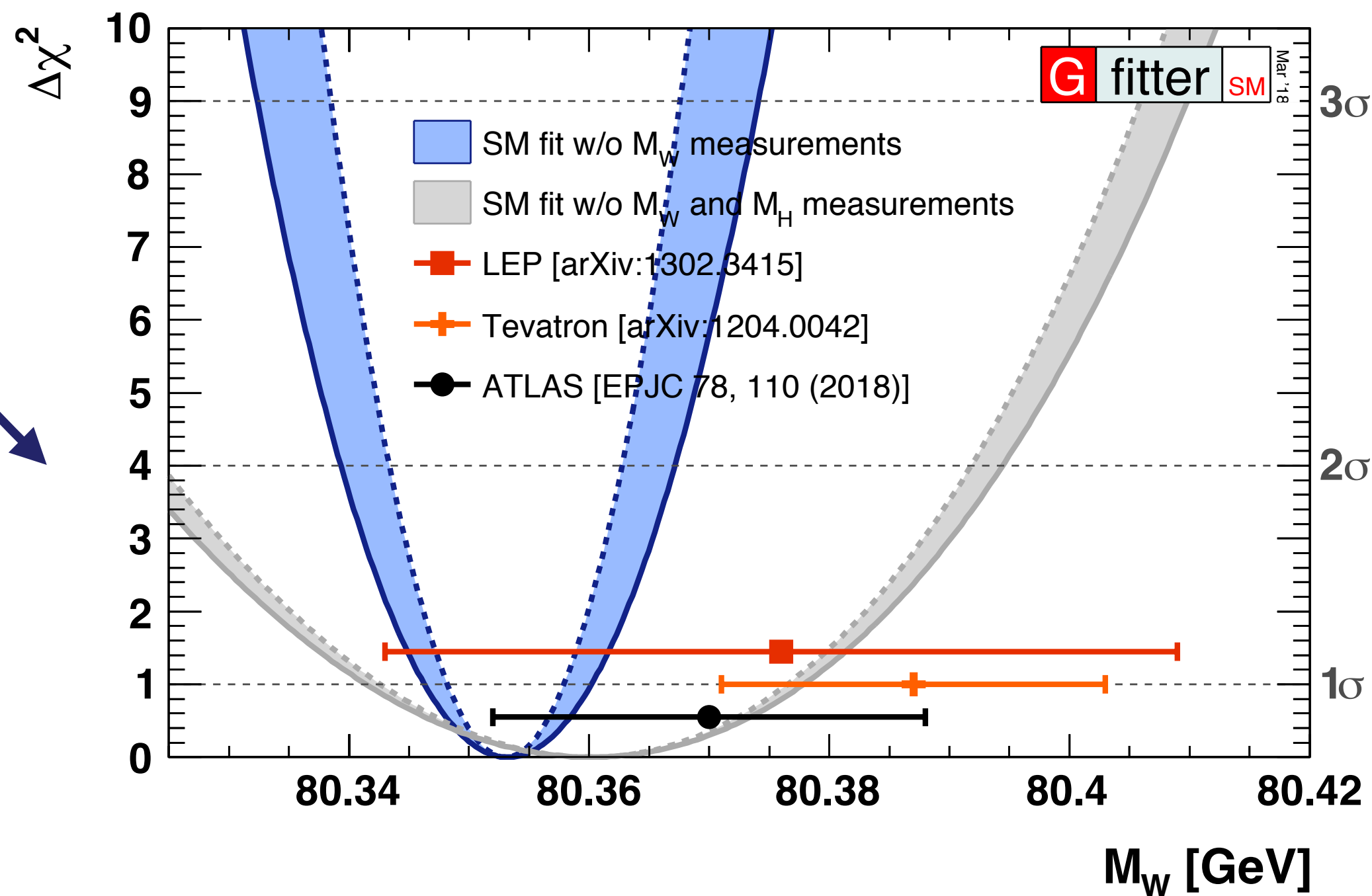
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Significant one loop corrections growing like the logarithm of m_H .

The m_H measurement

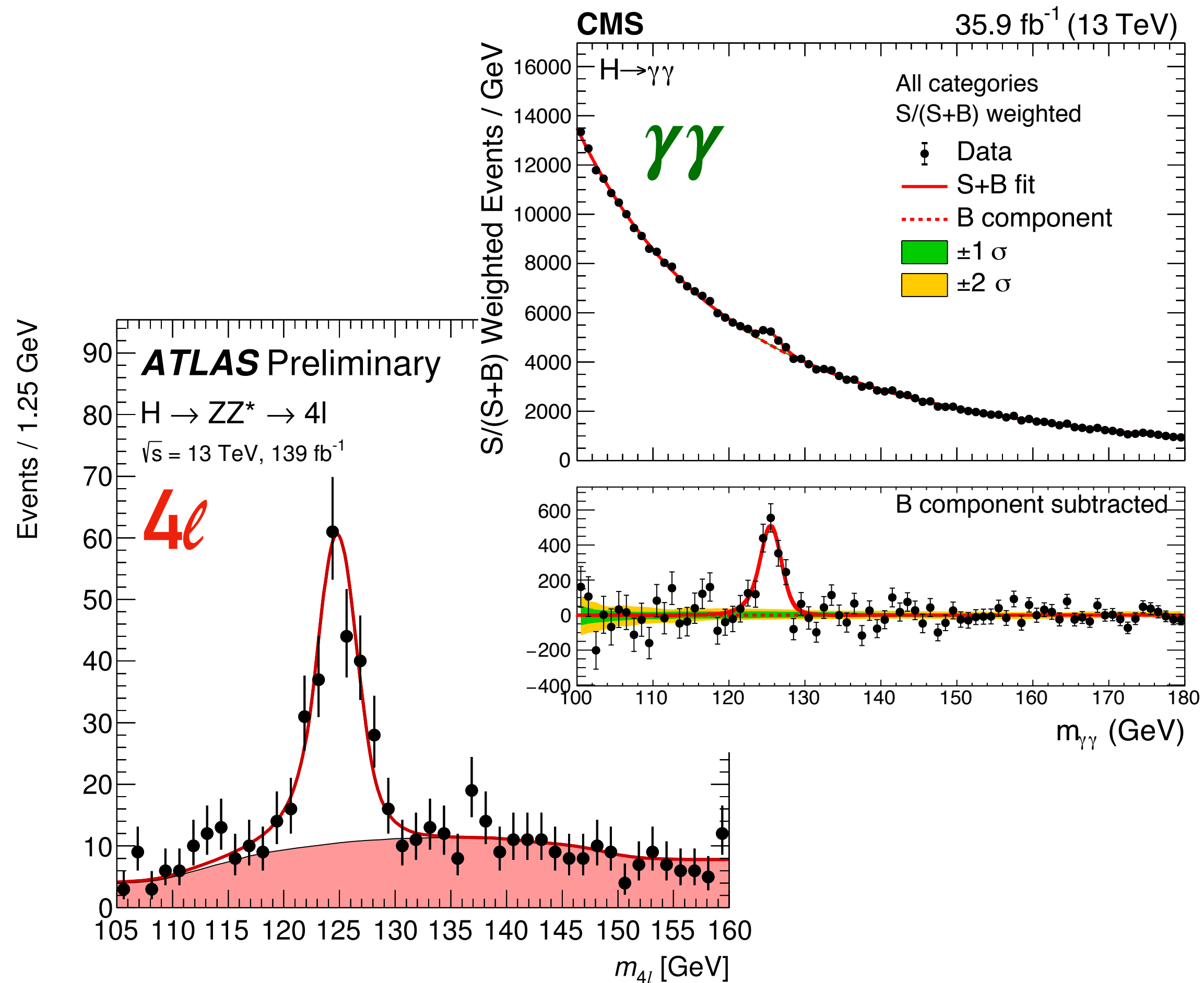
The concept oversimplified

The m_H measurement

The concept oversimplified

“Mass Peaks”

using high resolution channels ($4\ell + \gamma\gamma$)

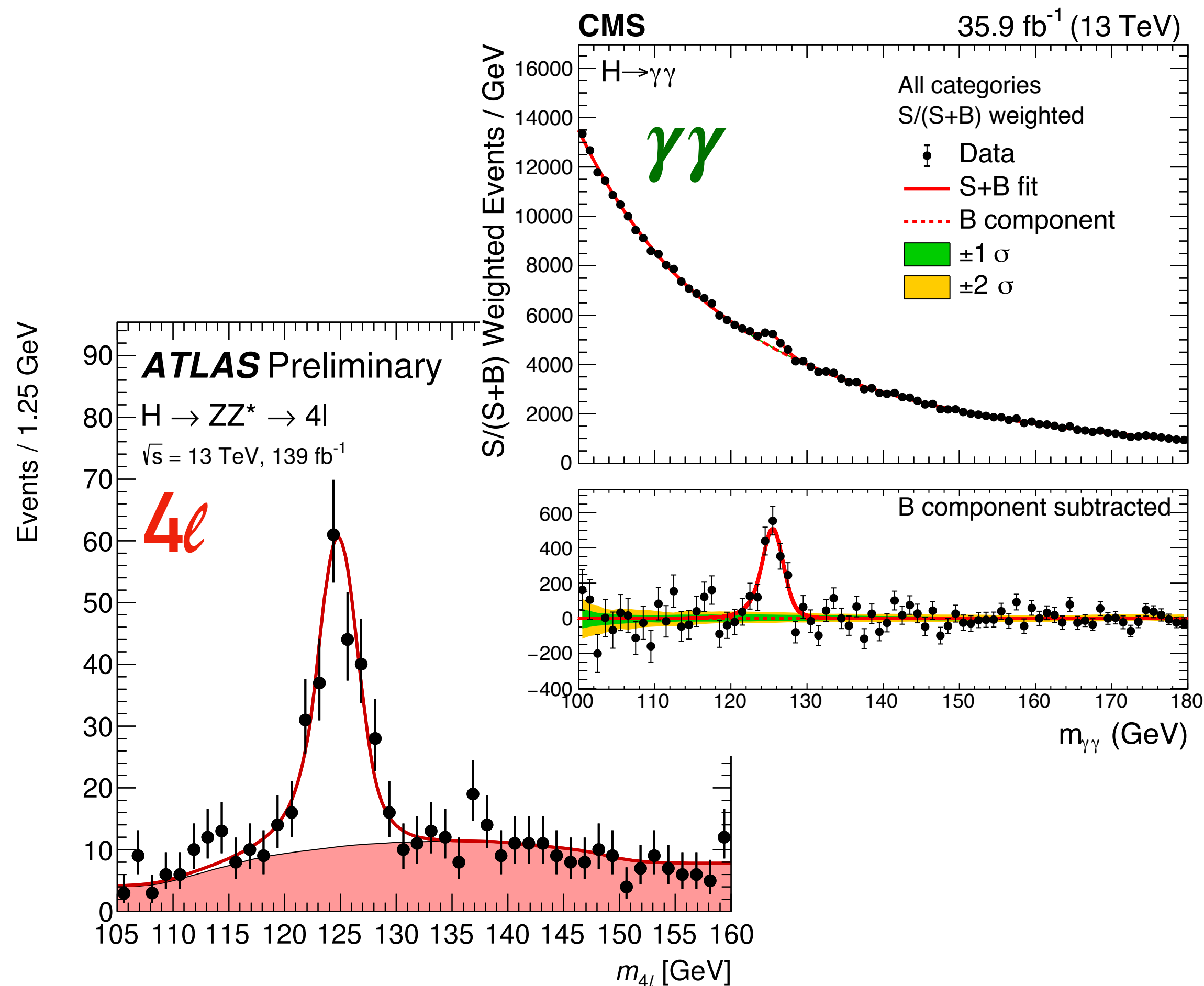


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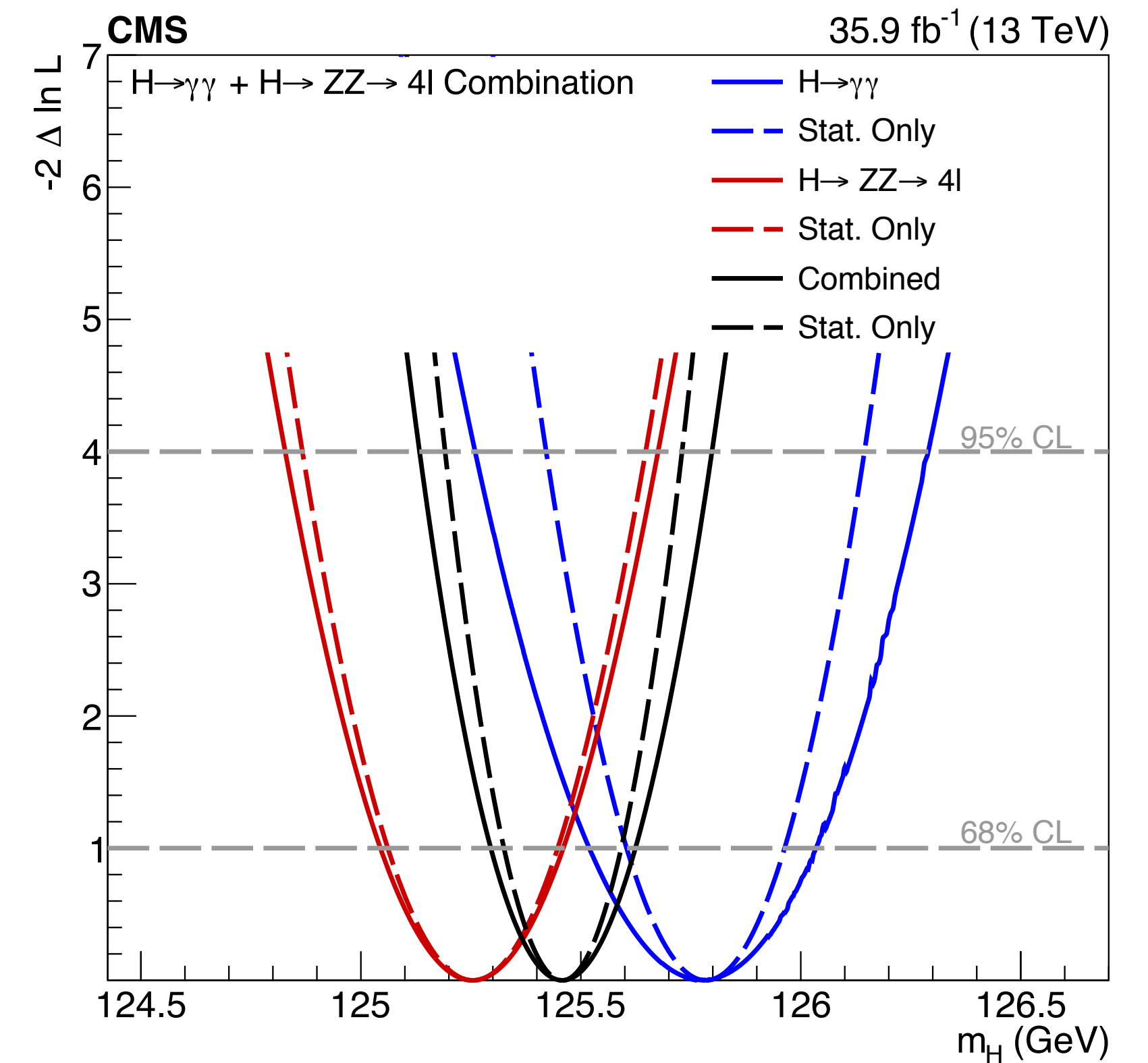
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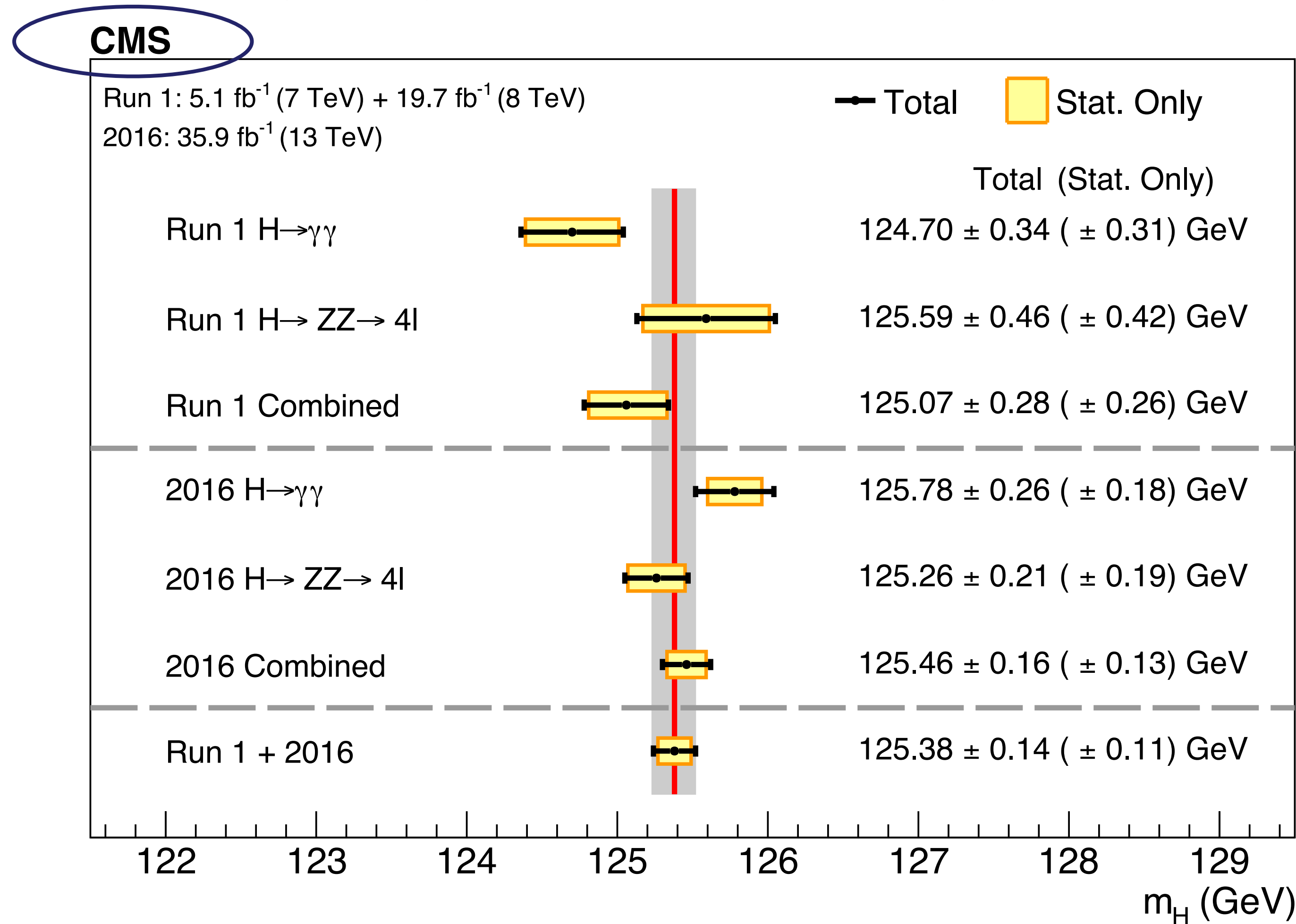


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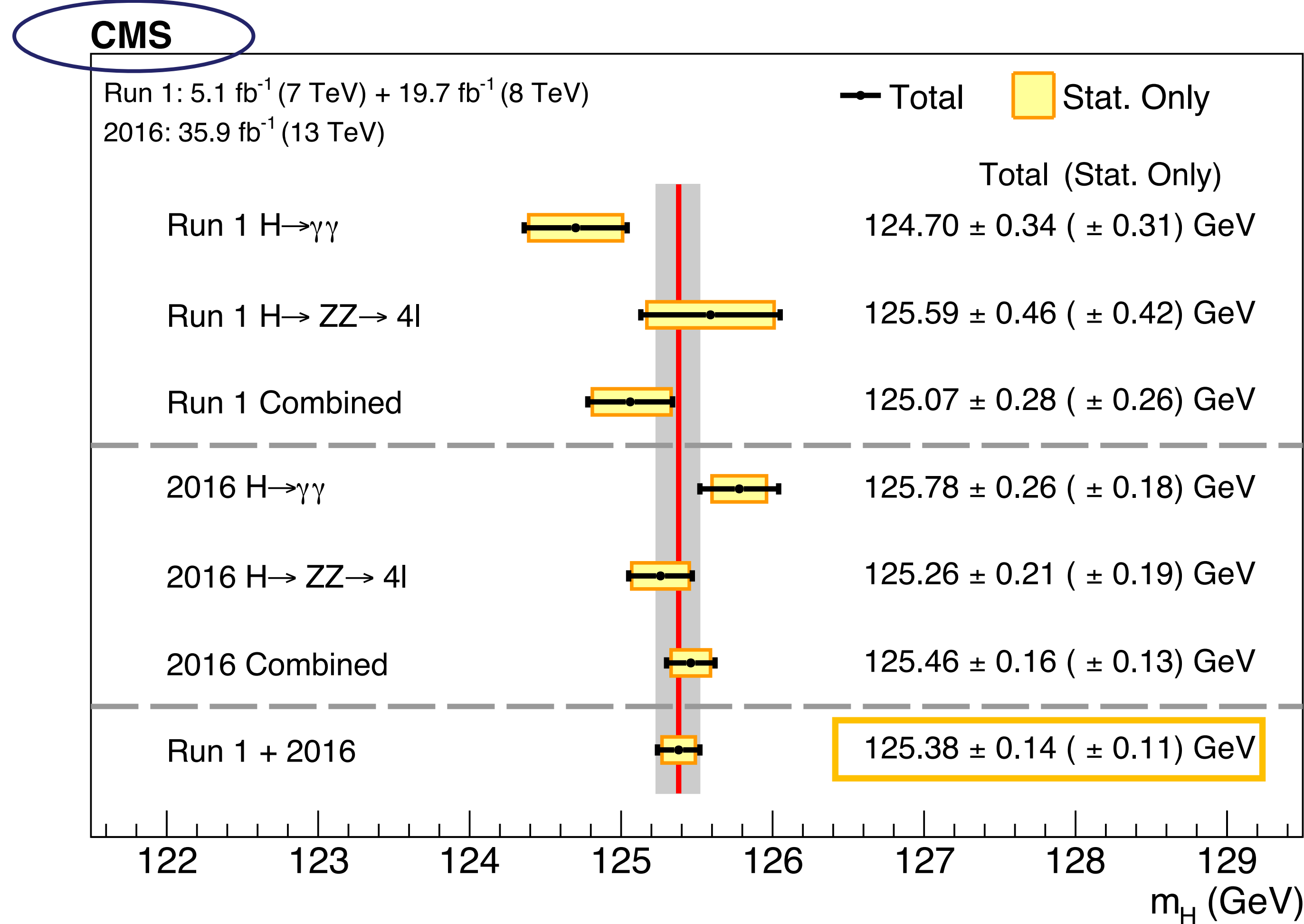
m_H : best precision to date

Combination of channels



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Uncertainty 140 MeV (0.11%)

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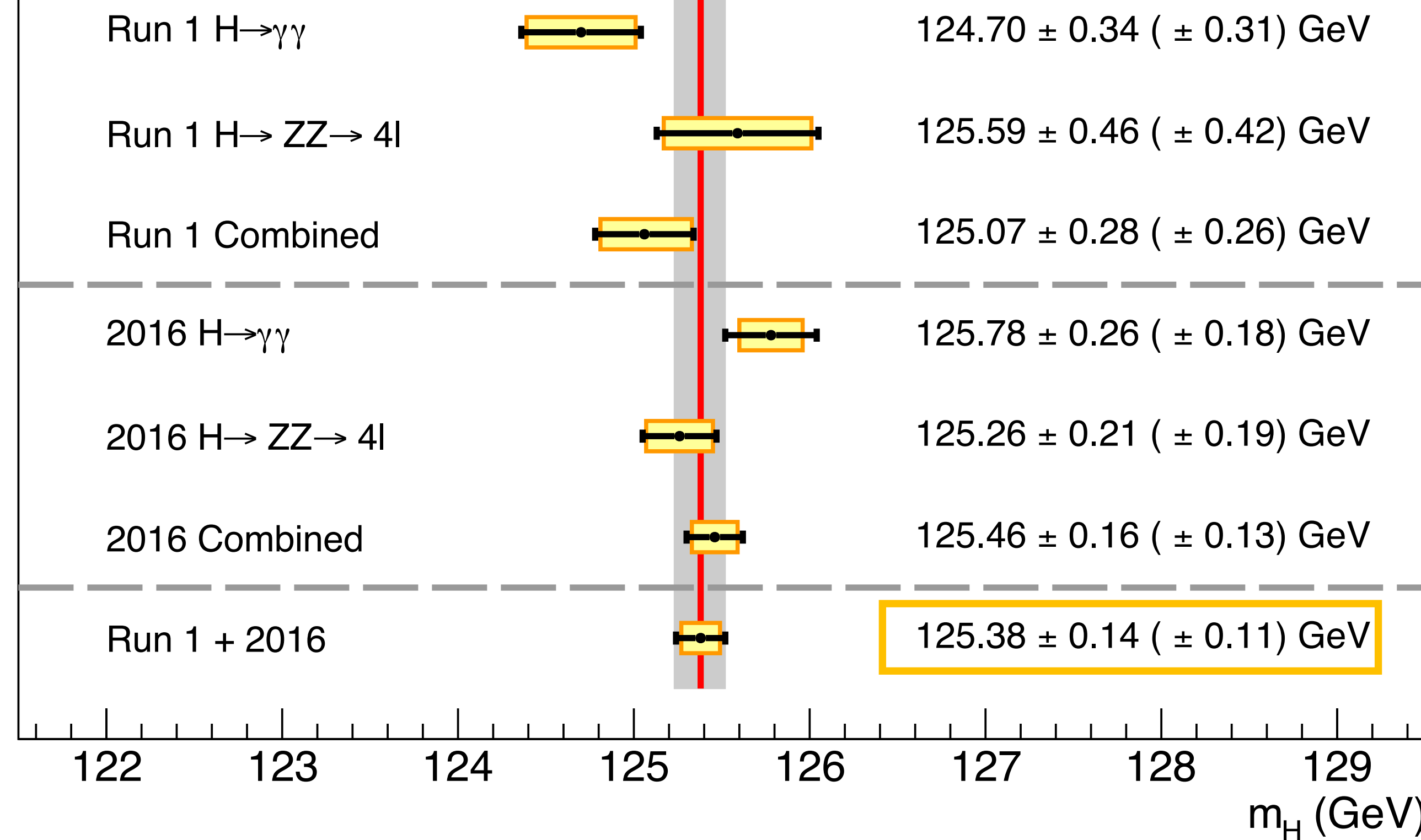
Combination of channels

CMS

Run 1: 5.1 fb⁻¹ (7 TeV) + 19.7 fb⁻¹ (8 TeV)
2016: 35.9 fb⁻¹ (13 TeV)

— Total □ Stat. Only

Total (Stat. Only)



Uncertainty 140 MeV (0.11%)

Precision driven by statistics, but photon, electron and muon scale and resolution systematics will soon become limiting!

Some measurements still based on partial Run 2 datasets: more improvements to come, and ATLAS+CMS mass combination

Projections on m_H

hopes & wishes

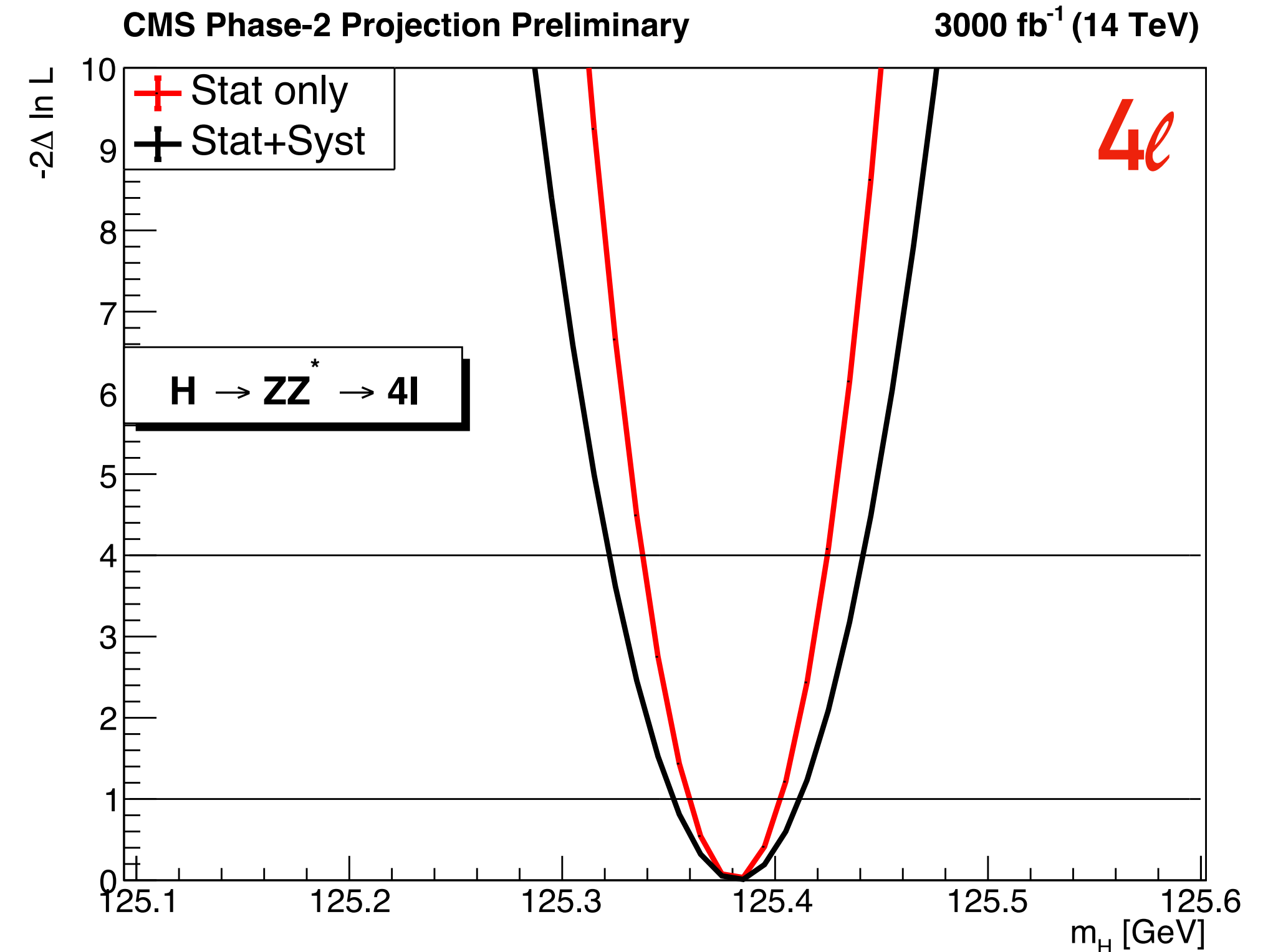


at the High-Luminosity LHC with the CMS detector

Expected m_H measurement uncertainty, given in MeV, in the two different scenarios: *Optimistic* and *Pessimistic*

$m_{4\ell}$ expected uncertainty (MeV)	inclusive	4μ	$4e$	$2e2\mu$	$2\mu2e$
<i>Optimistic</i>					
Total	26	30	105	60	67
Syst impact	16	11	64	31	32
Stat only	22	28	83	51	59
<i>Pessimistic</i>					
Total	30	32	206	107	112
Syst impact	20	15	189	94	95
Stat only	22	28	83	51	59

Stat > Syst but same magnitude



Projections on m_H

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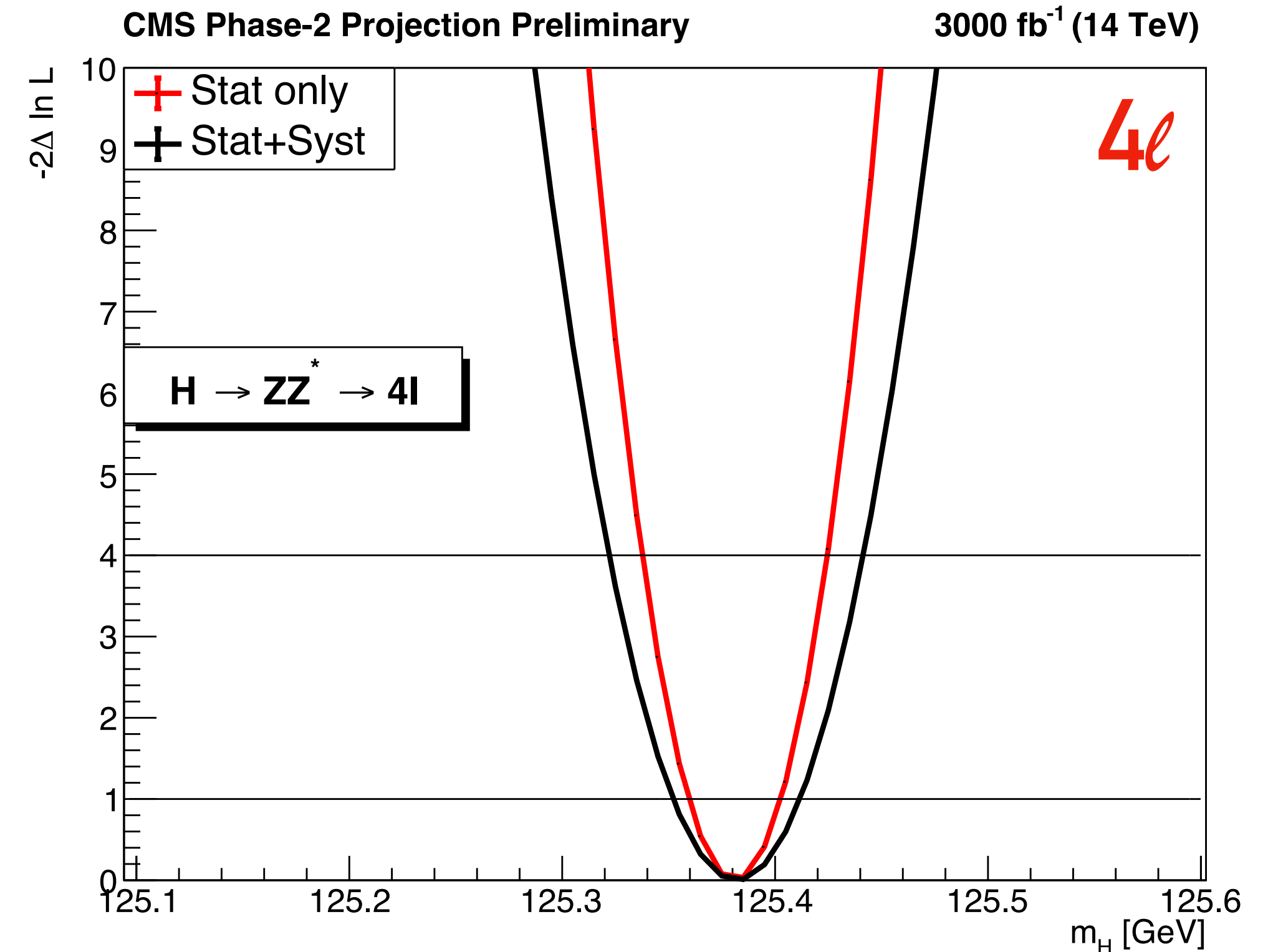


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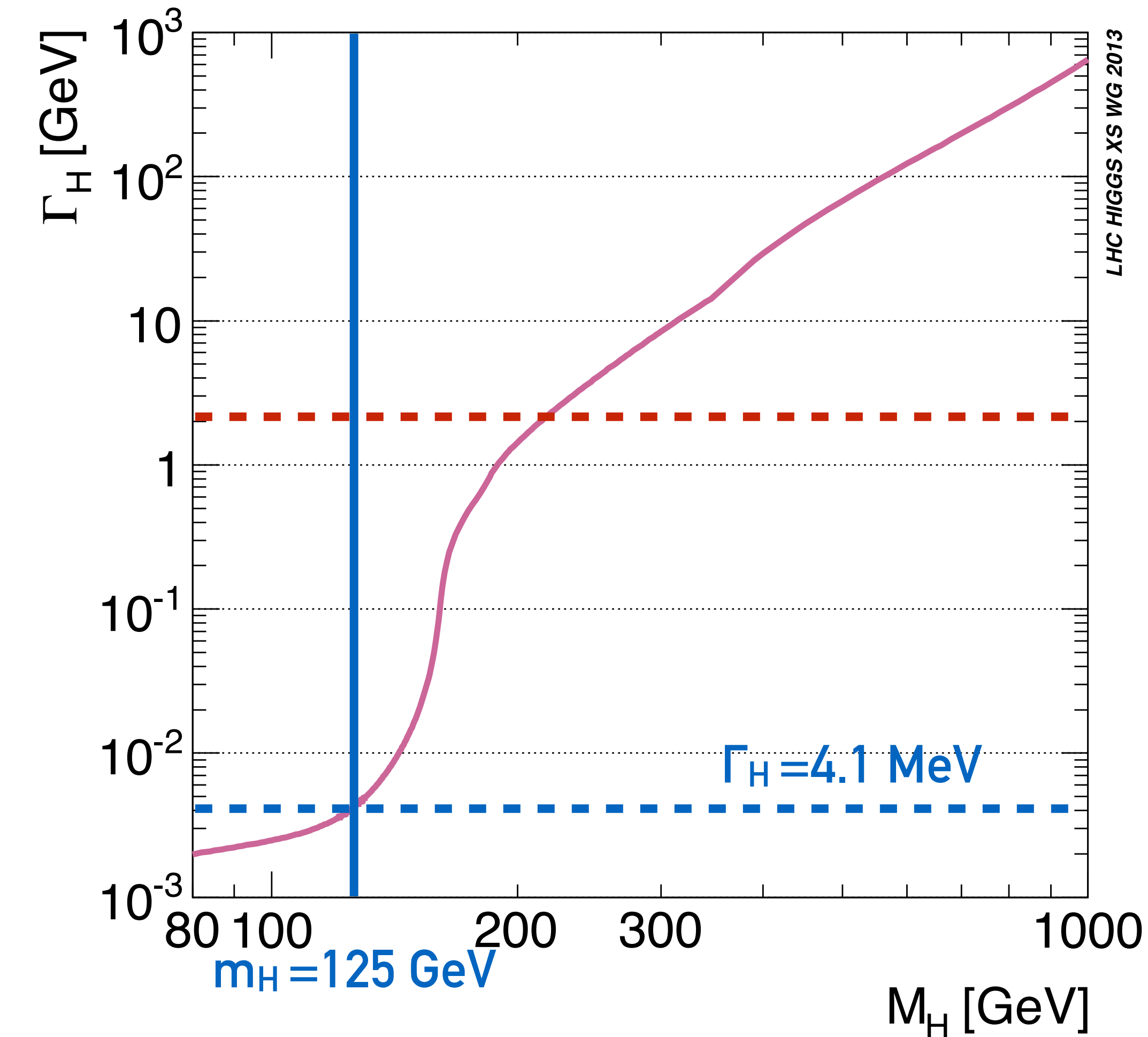
Stat > Syst but same magnitude



> than a factor 7 (210 MeV \Rightarrow 26-30 MeV) compared to the 4l current measurement

Improvements due to: a factor 10 in luminosity, the new tracker with less material, the precision and stability of the HGCal, the improvements to the barrel calorimeters, and the pileup suppression provided by the new MTD.

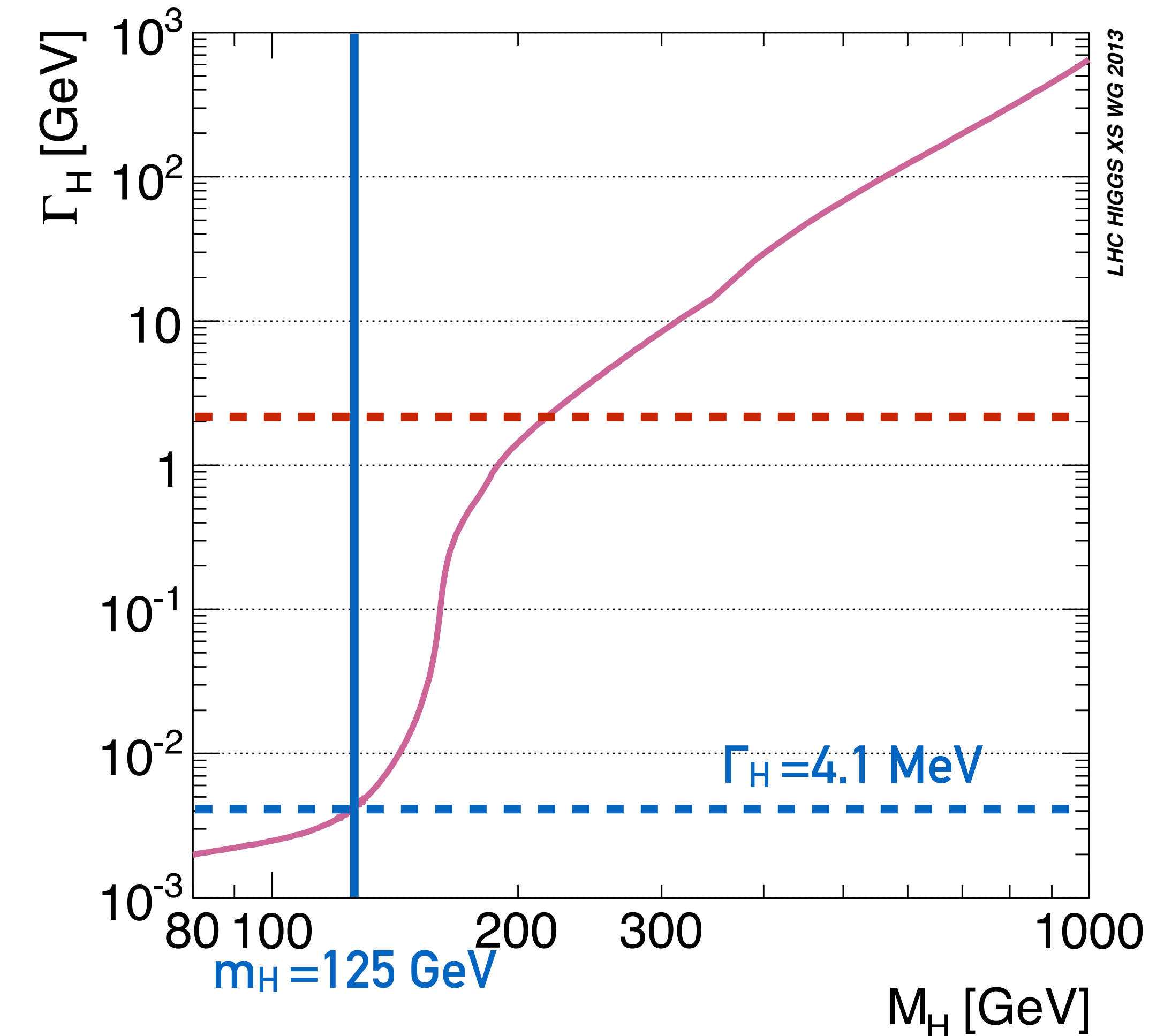
The Higgs boson width



A crucial parameter for BSM searches, in SM $c\tau_H = 48$ fm, small width $\Gamma_H = 4.1$ MeV

We have long experience with heavy EW bosons (W and Z). However, their width is $\Gamma_H \sim 2$ GeV !

The Higgs boson width



A crucial parameter for BSM searches, in SM $c\tau_H = 48$ fm, small width $\Gamma_H = 4.1$ MeV

We have long experience with heavy EW bosons (W and Z). However, their width is $\Gamma_H \sim 2$ GeV !

The direct measurements it is extremely hard! In addition, the total width is the sum of all the partial widths, on the contrary of LEP, at LHC only $\sigma \times \text{BR}$ can be measured.

The Higgs boson width : strategies

Direct measurements

on-shell line shape
lifetime

Indirect measurements

couplings
off-shell production

The Higgs boson width : strategies

Direct measurements

on-shell line shape
lifetime

Limited by detector resolutions



Indirect measurements

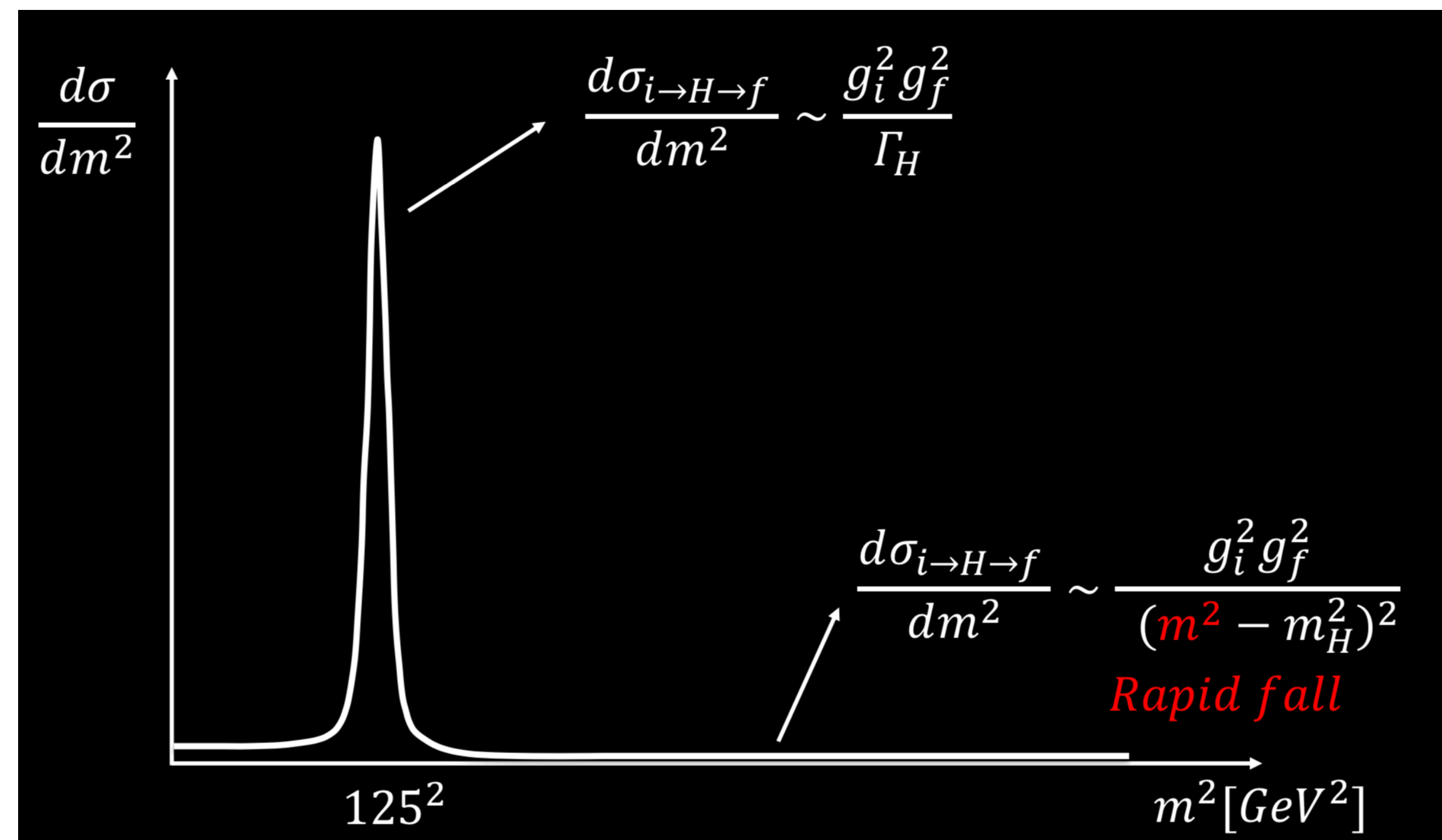
couplings
off-shell production

The way out



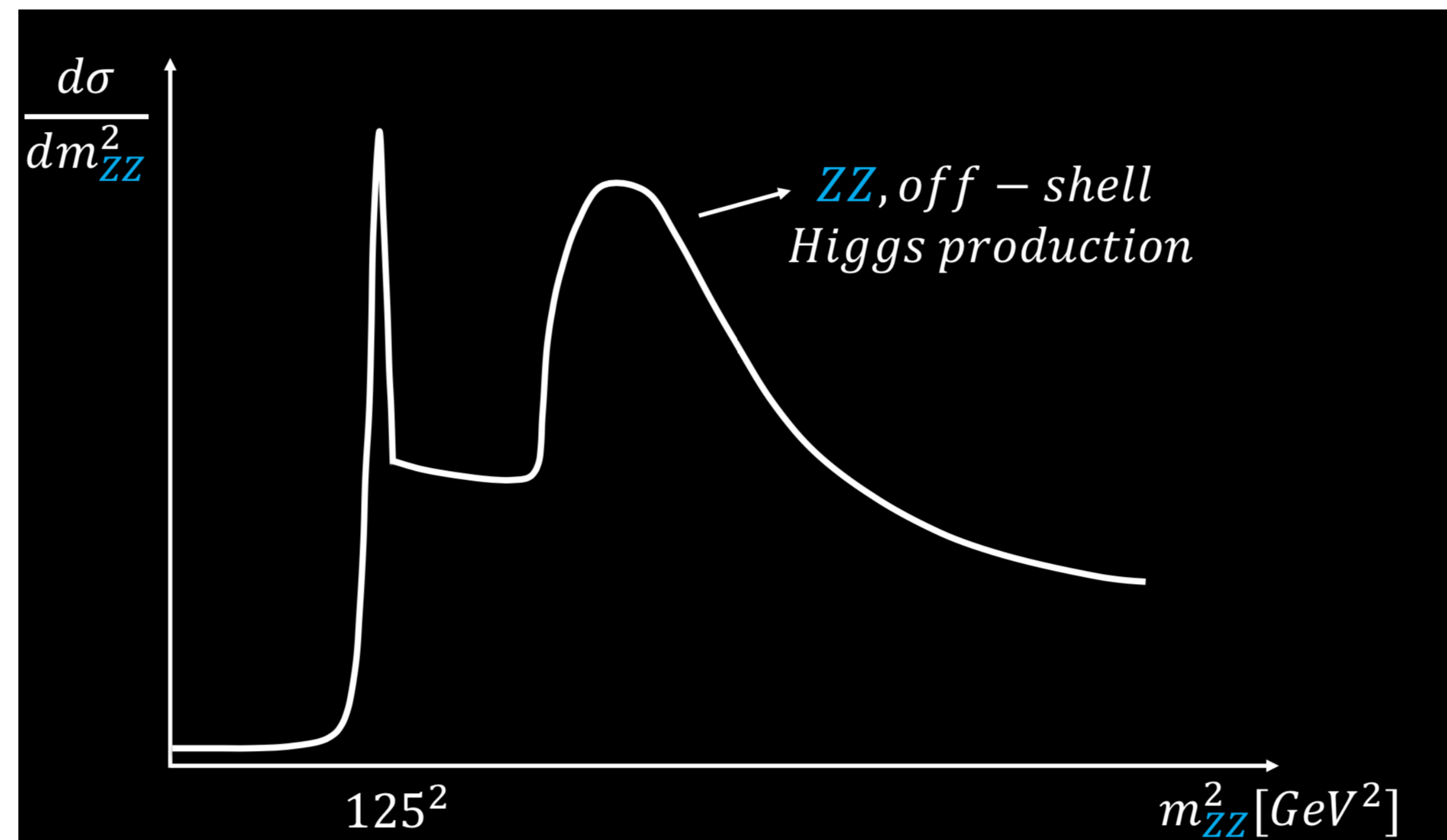
From off-shell production

Indirect measurement



The fixed-width Breit-Wigner scheme is generally good in describe the inclusive differential ($d\sigma/dm^2$) Higgs boson production

... but ...



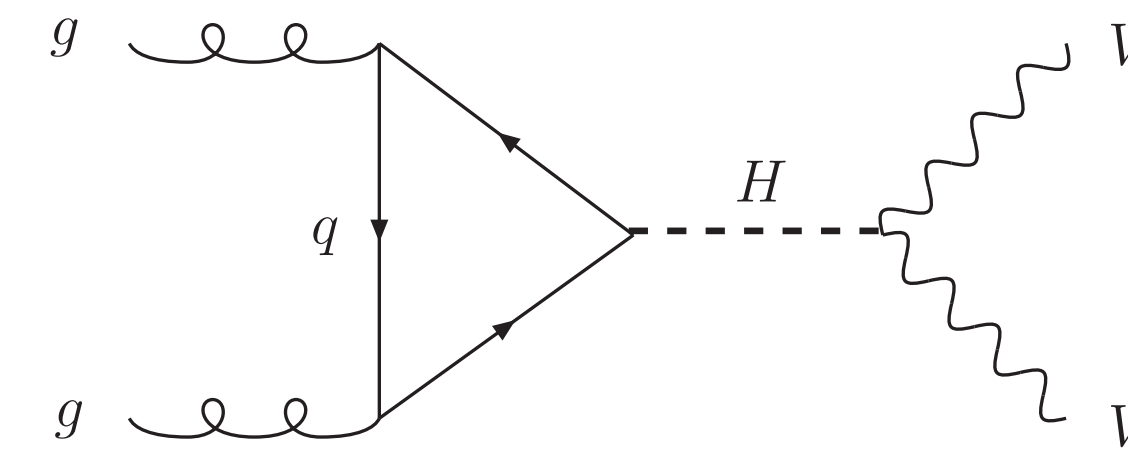
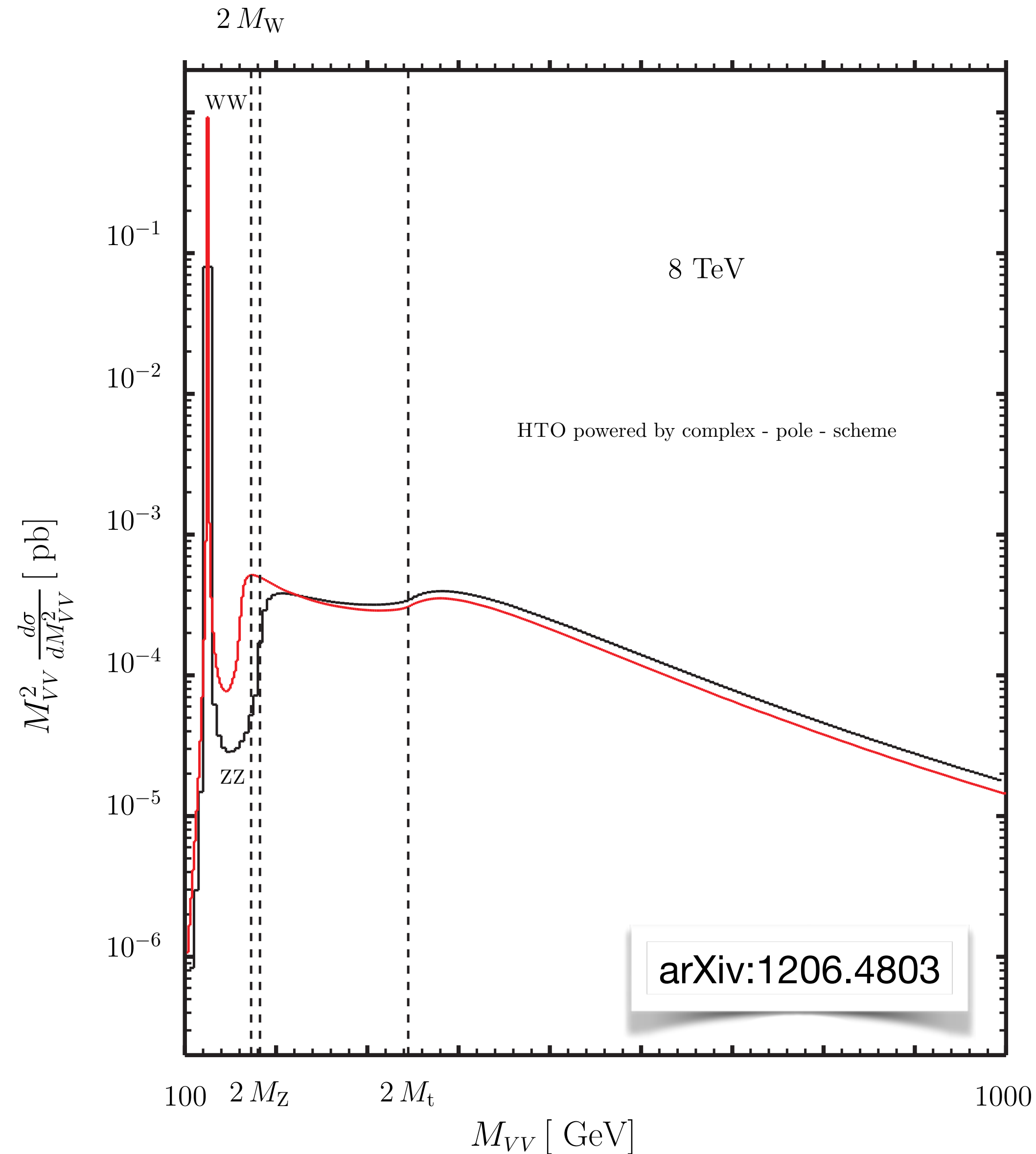
The fixed-width Breit-Wigner scheme is generally good in describe the inclusive differential ($d\sigma/dm^2$) Higgs boson production

... but ...

... as soon as we restrict to VV decay channel there is a large off-shell contribution above the VV threshold (high Higgs virtuality), it means that two q^2 propagators compensate and the cross section is enhanced.

From off-shell production

Indirect measurement



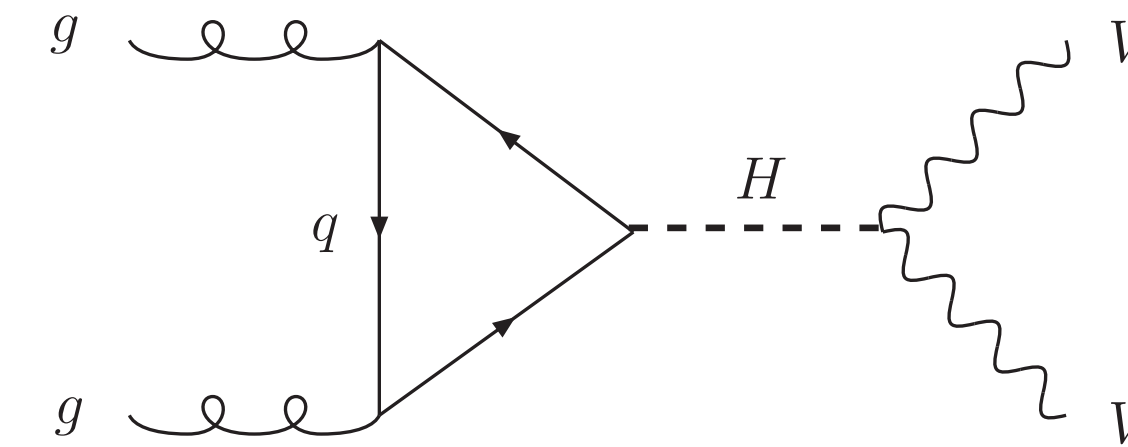
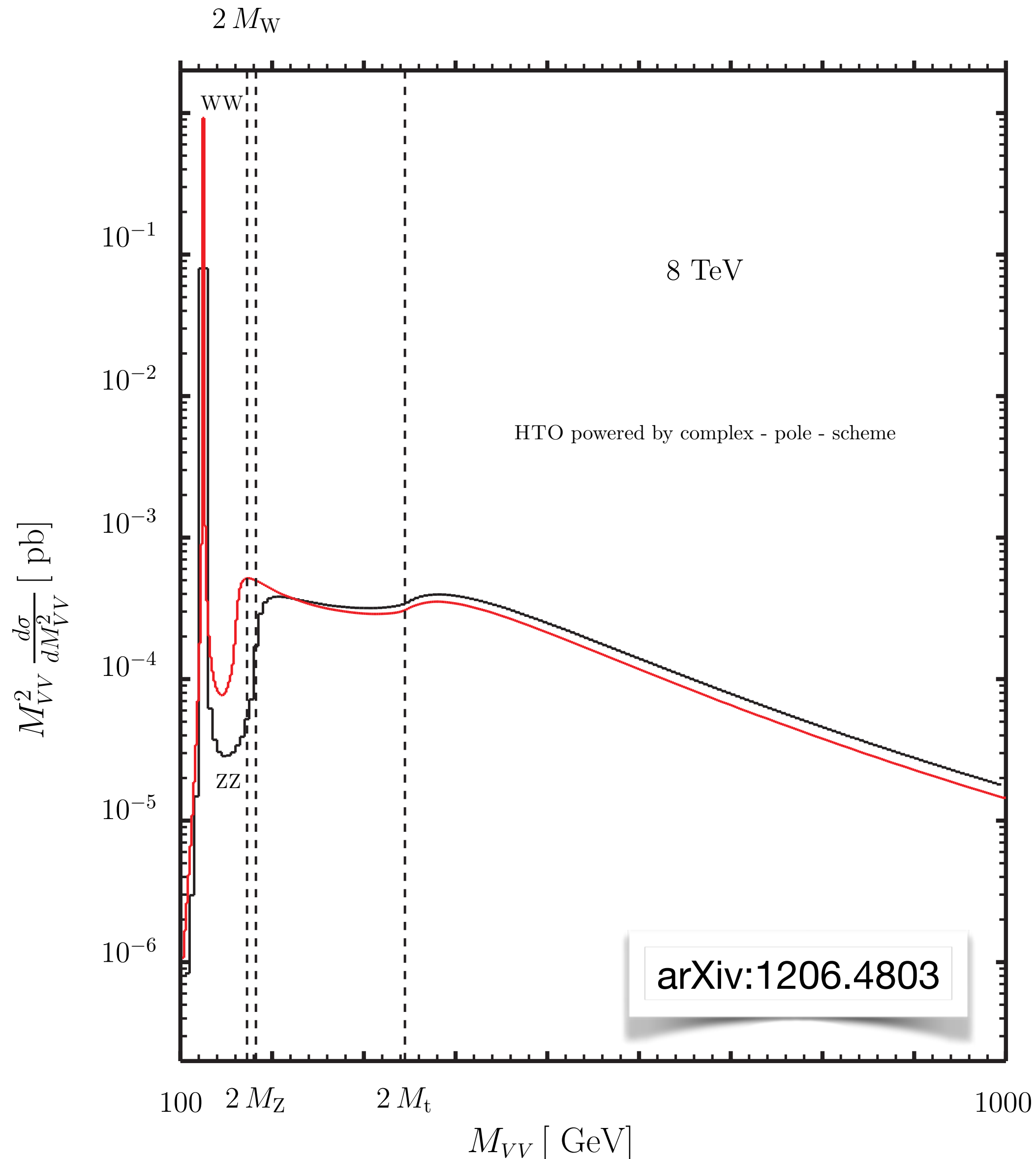
The NNLO **ZZ (black)** and **WW (red)** invariant mass distributions in $gg \rightarrow VV$ for $m_H = 125 \text{ GeV}$

$$\sigma_{gg \rightarrow H \rightarrow ZZ^*}^{\text{on-shell}} \sim \frac{g_{ggH}^2 g_{HZZ}^2}{m_H \Gamma_H}$$

$$\sigma_{gg \rightarrow H^* \rightarrow ZZ}^{\text{off-shell}} \sim \frac{g_{ggH}^2 g_{HZZ}^2}{(2m_Z)^2}$$

From off-shell production

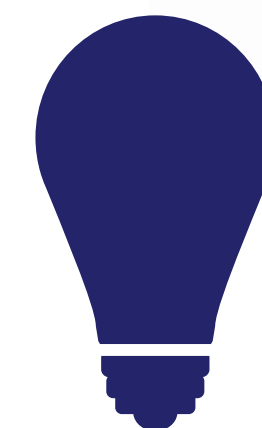
Indirect measurement



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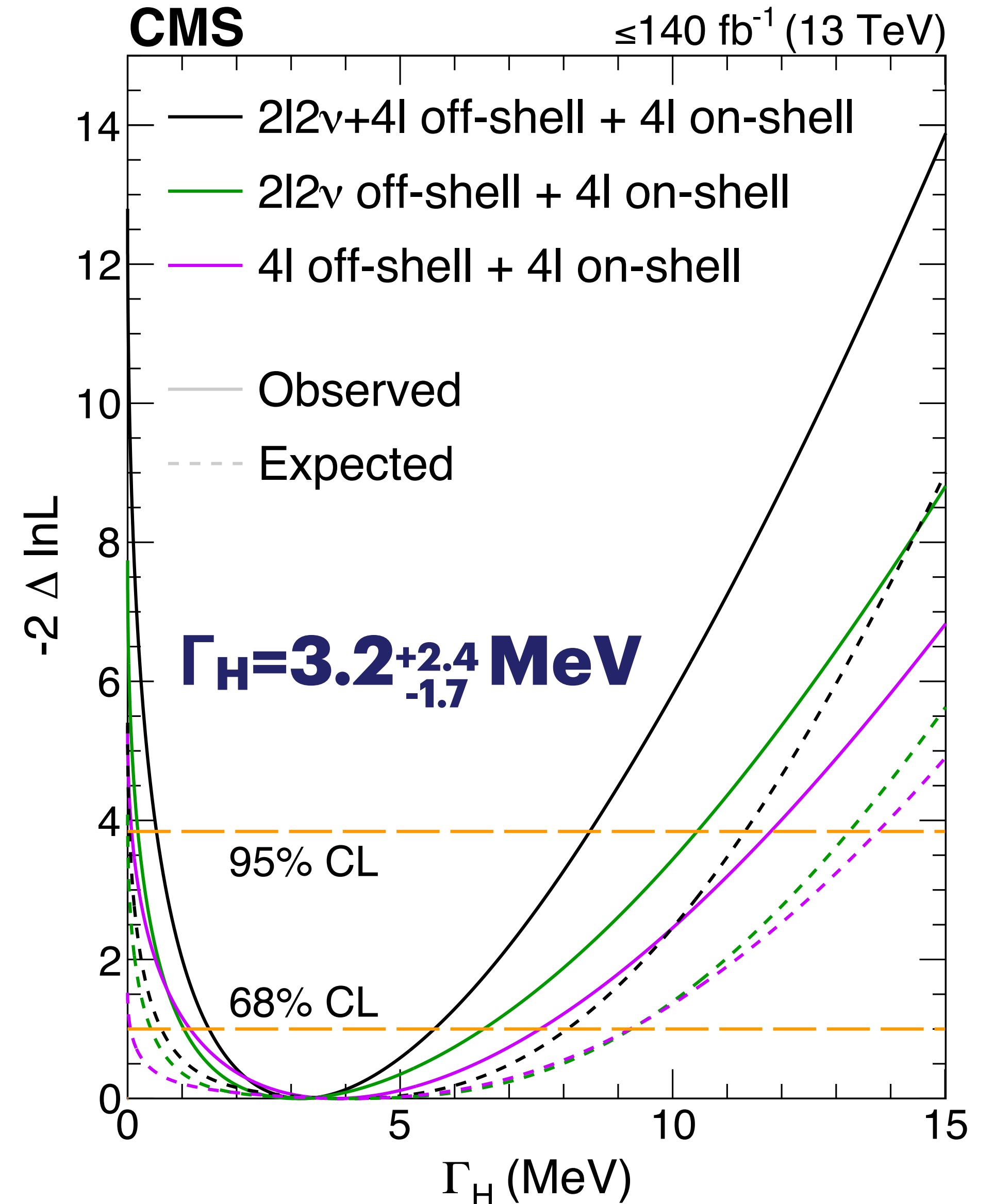
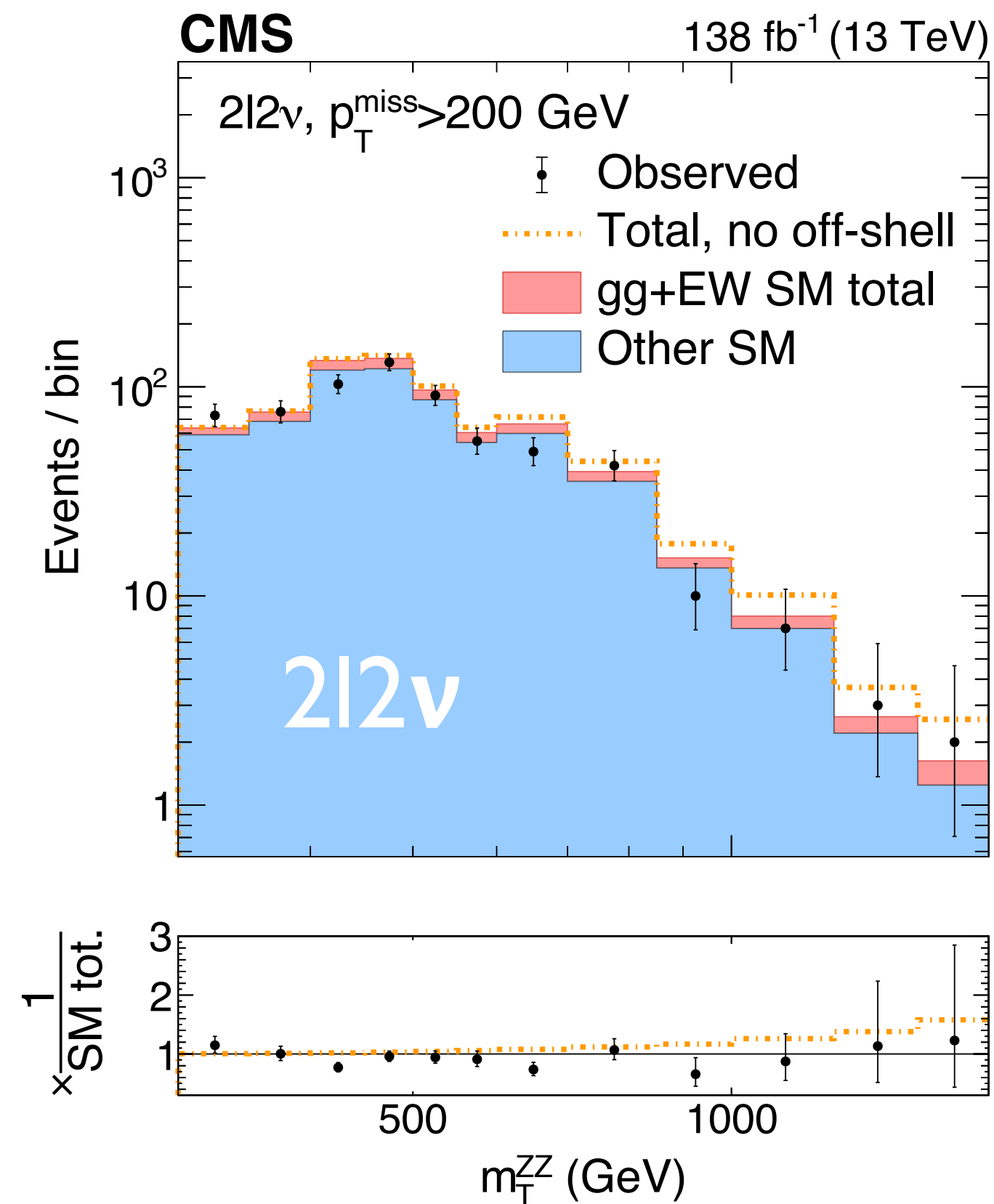
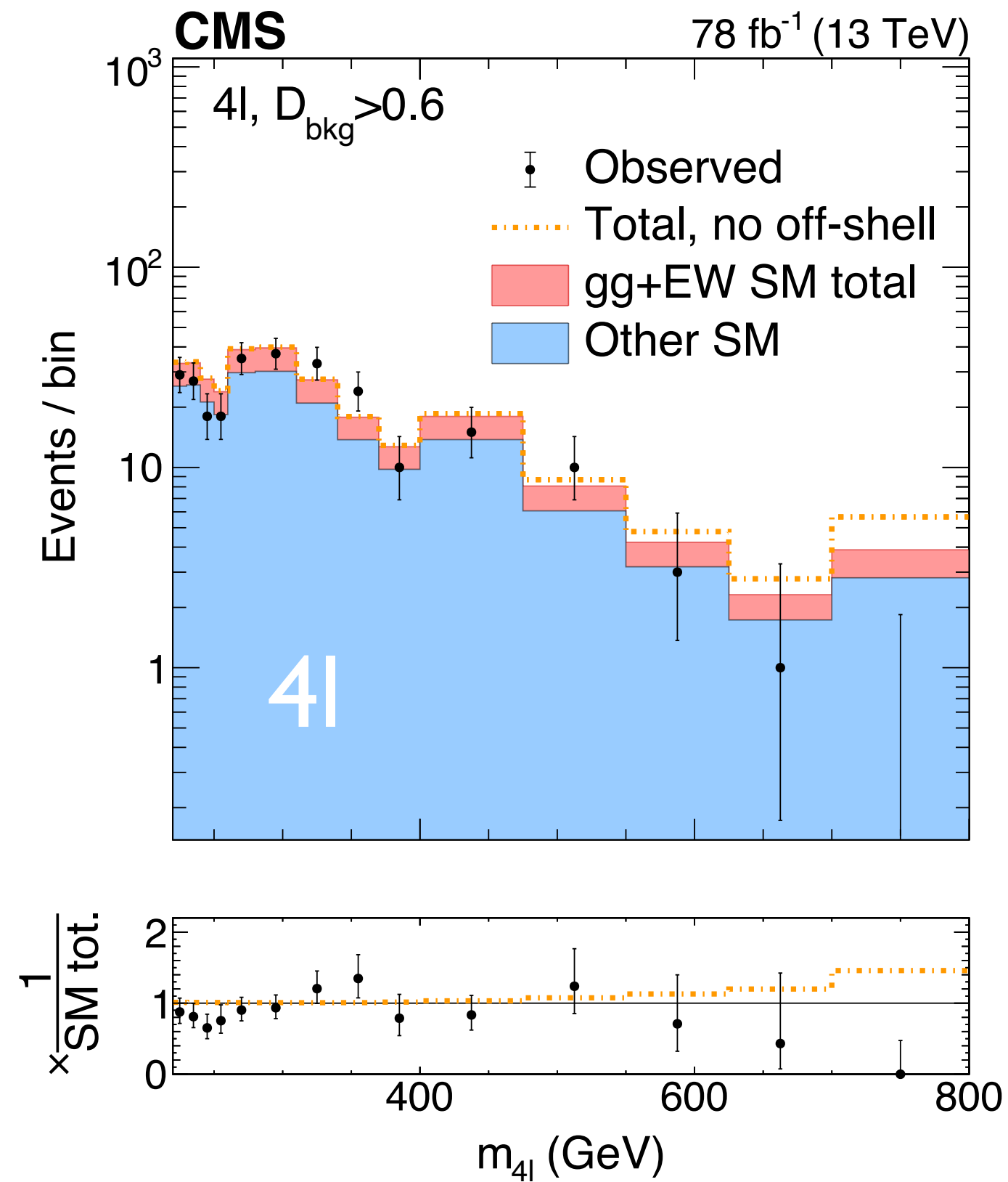
$$\sigma_{\text{on-shell}}^{gg \rightarrow H \rightarrow ZZ^*} \sim \frac{\cancel{g_{ggH}^2} \cancel{\delta_{HZZ}^2}}{m_H \Gamma_H}$$

$$\sigma_{\text{off-shell}}^{gg \rightarrow H^* \rightarrow ZZ} \sim \frac{\cancel{g_{ggH}^2} \cancel{\delta_{HZZ}^2}}{(2m_Z)^2}$$



Idea: make the ratio between the two cross sections and measure Γ_H

Measurement of the Higgs boson width

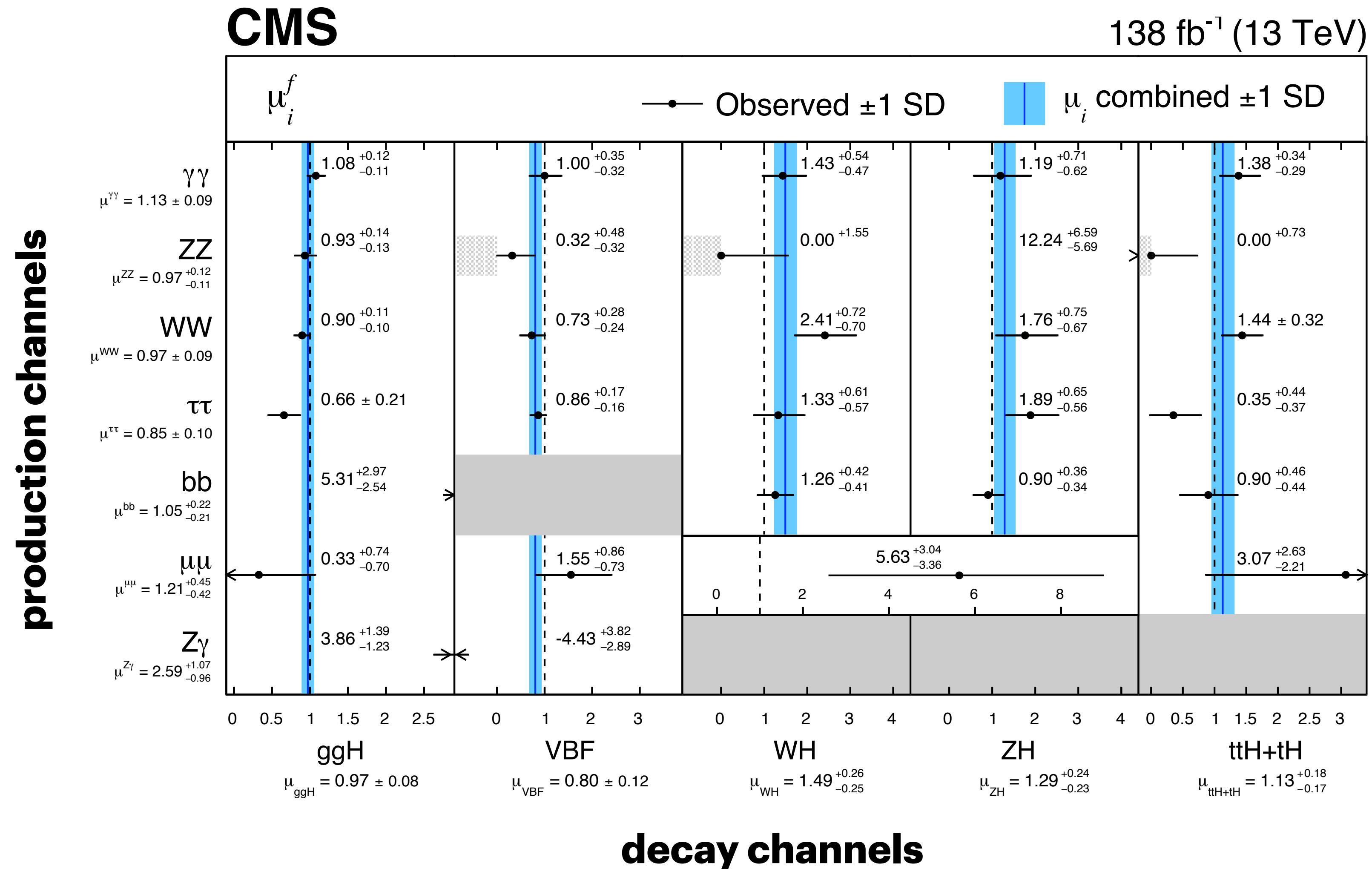


No off-shell scenario ($\mu^{\text{off-shell}}=0$) excluded $>99.9\%$ CL

\Rightarrow **3.6σ for evidence for off-shell H production**

Higgs boson couplings

After 10 years many signal strength modifiers¹⁾ measured

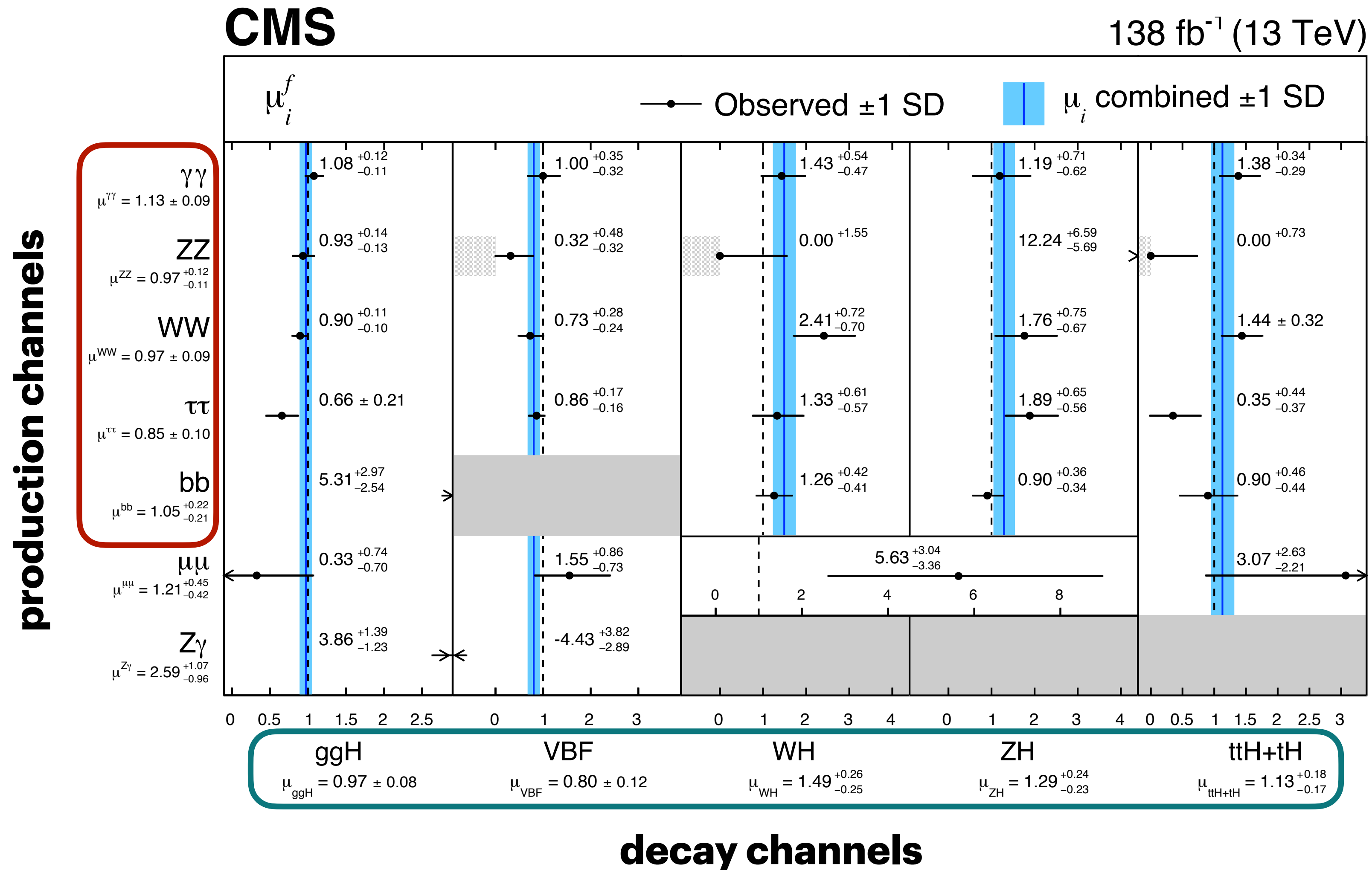


1) μ scale cross sections and branching fractions relative to the SM

Higgs boson couplings

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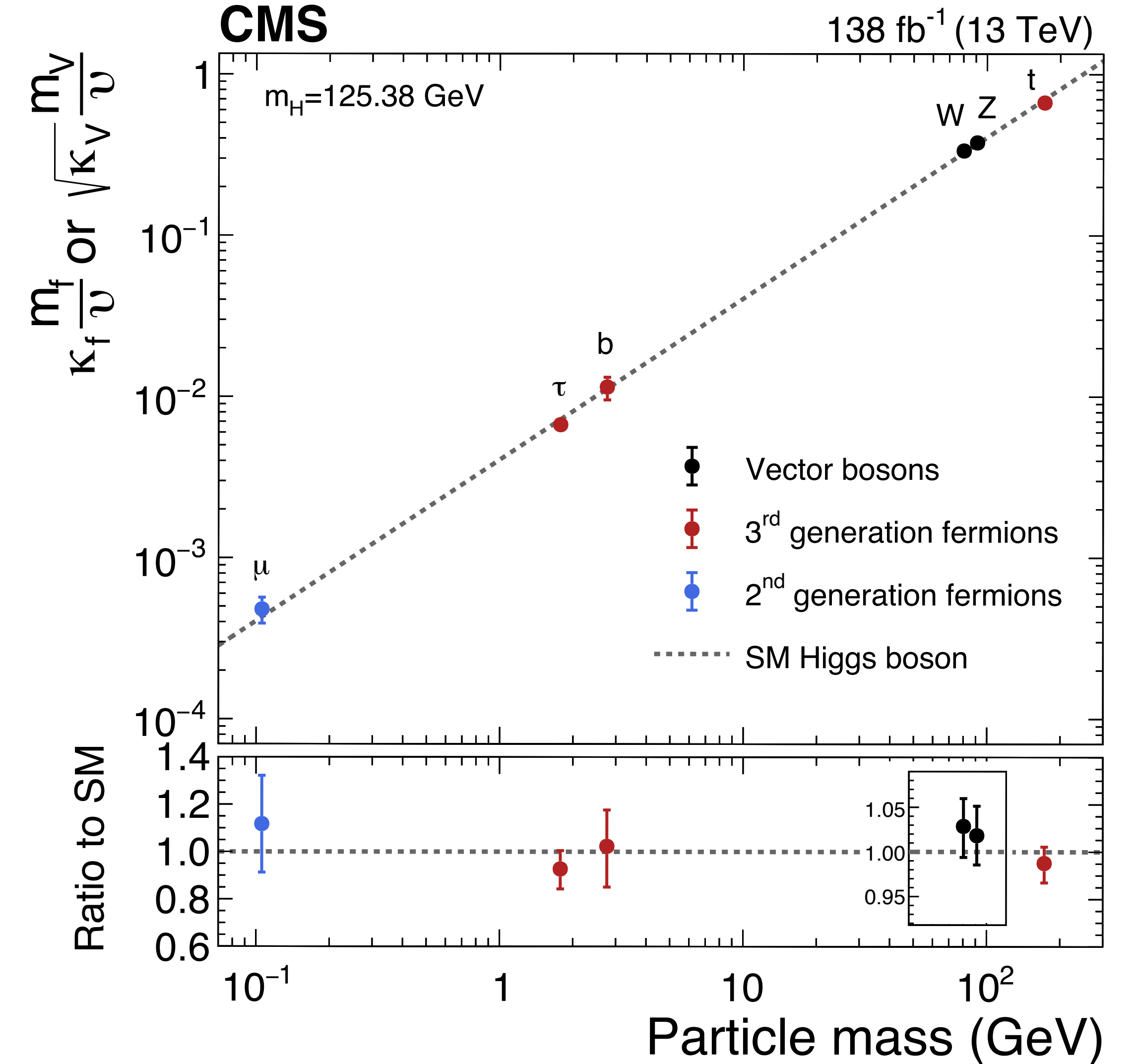
5 main production channels and **5 main** decay channels are observed



1) μ scale cross sections and branching fractions relative to the SM

Higgs boson couplings vs. mass

Remarkable agreement with the predictions of the BEH mechanism over 3 orders of magnitude of mass!



Coupling modifier k_j : parameterisation of inclusive production and decay rates
 e.g. $k_j^2 = \sigma/\sigma_{SM}$

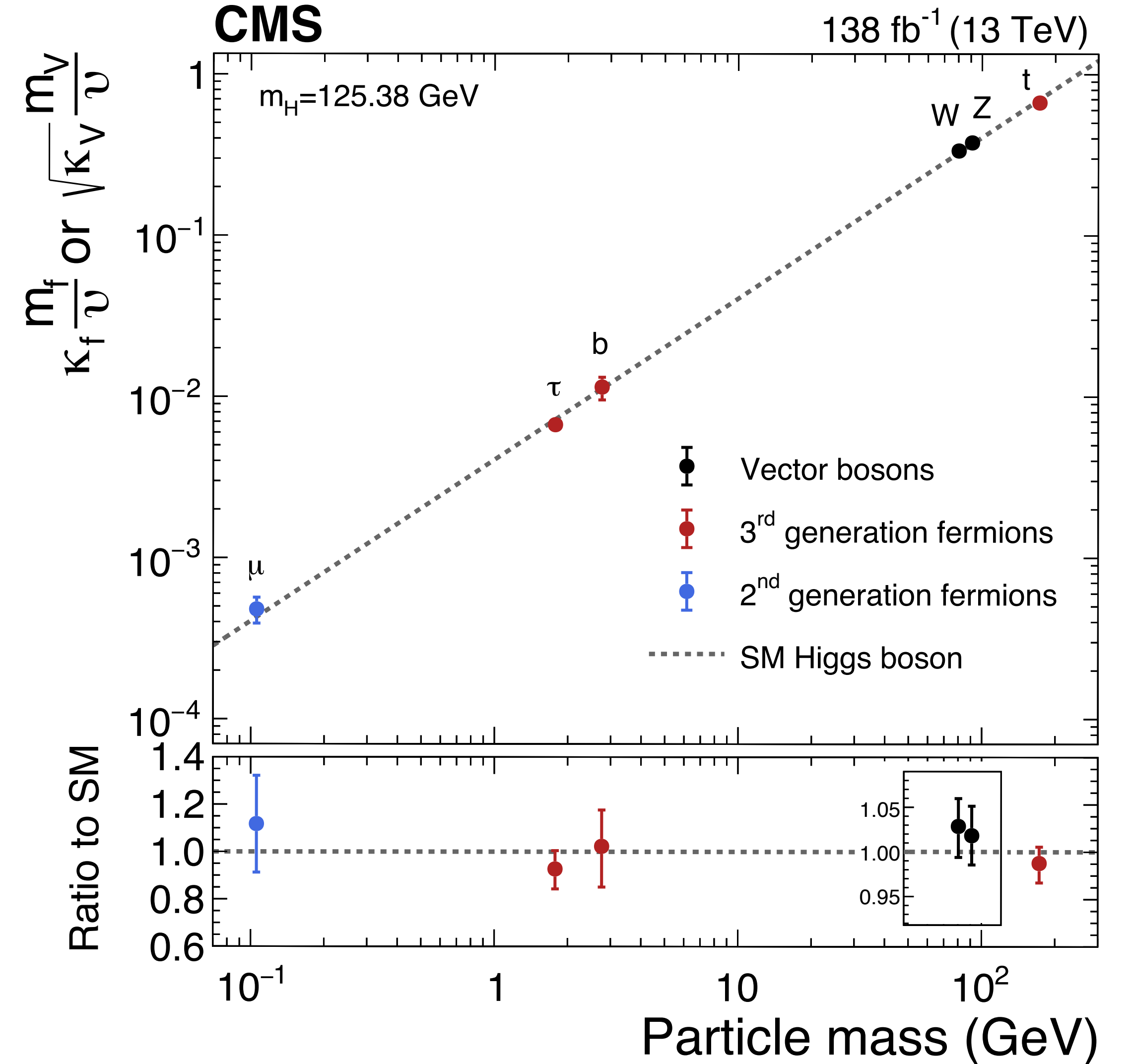
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Observation ($>5\sigma$) of coupling with 3rd gen.

Evidence ($>3\sigma$) of coupling with 2nd gen.



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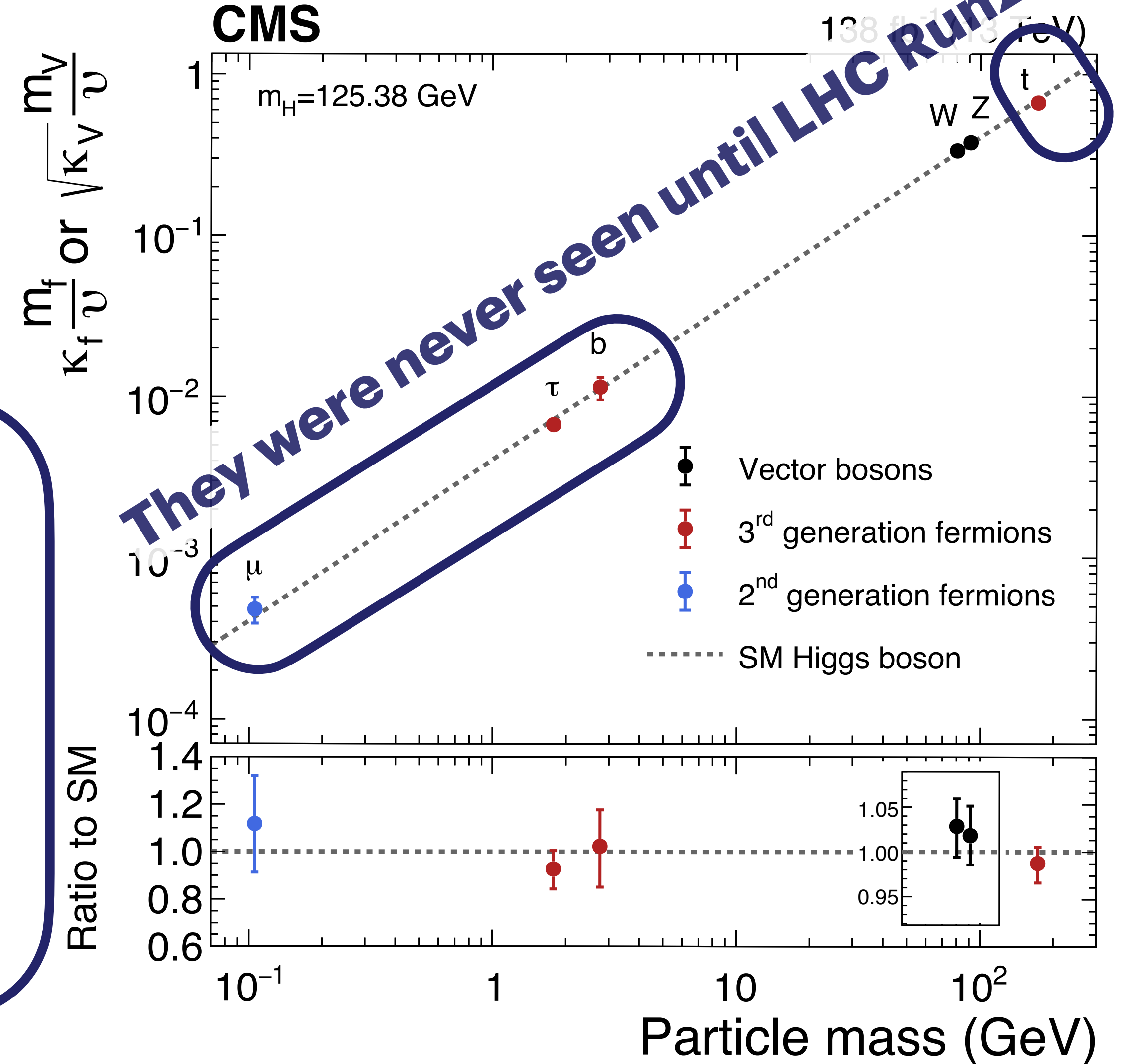
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Observation of fundamental interaction (the Yukawa) is important as observation of a fundamental particle.

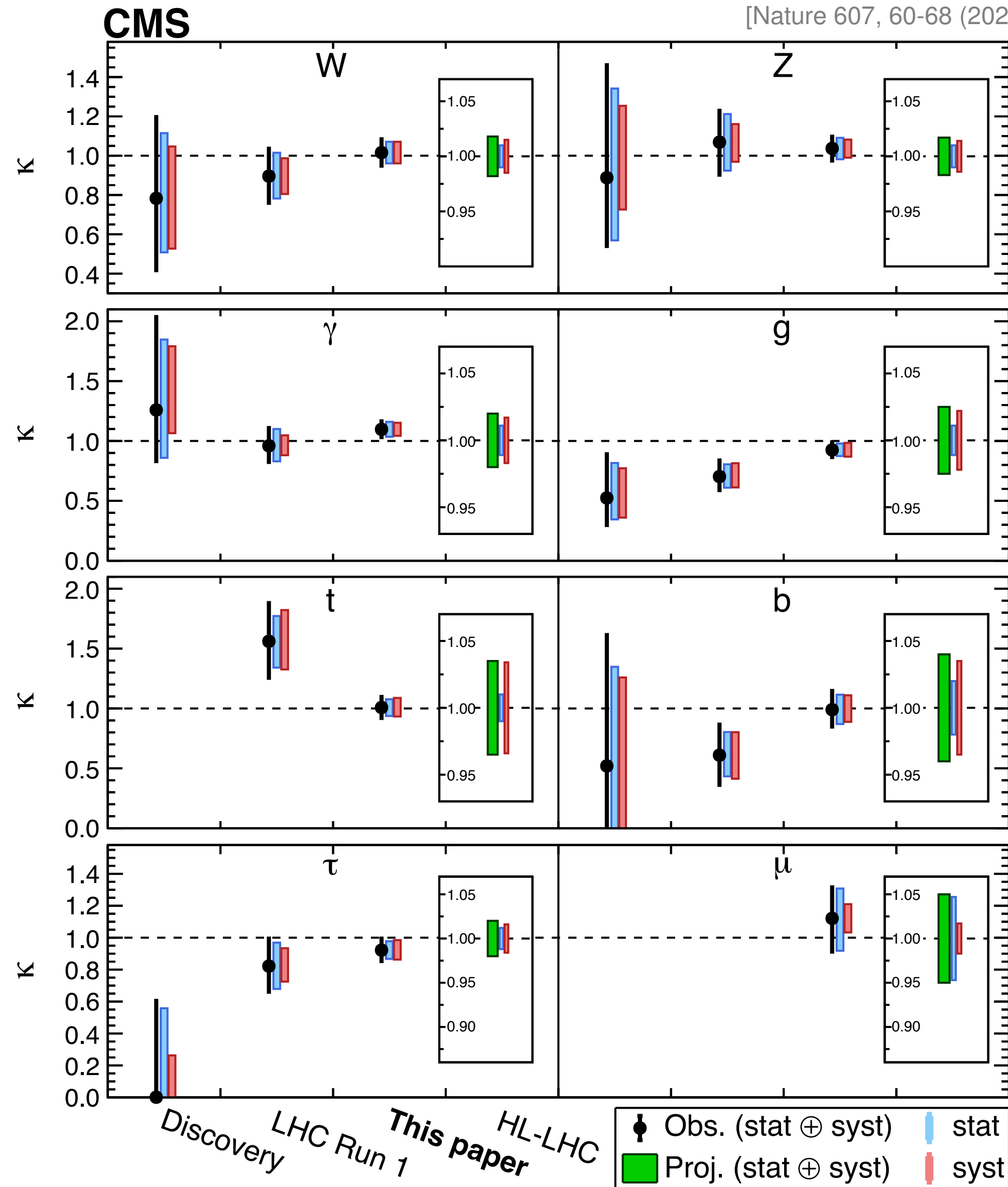


Projections of H couplings

hopes & wishes



[Nature 607, 60-68 (2022)]

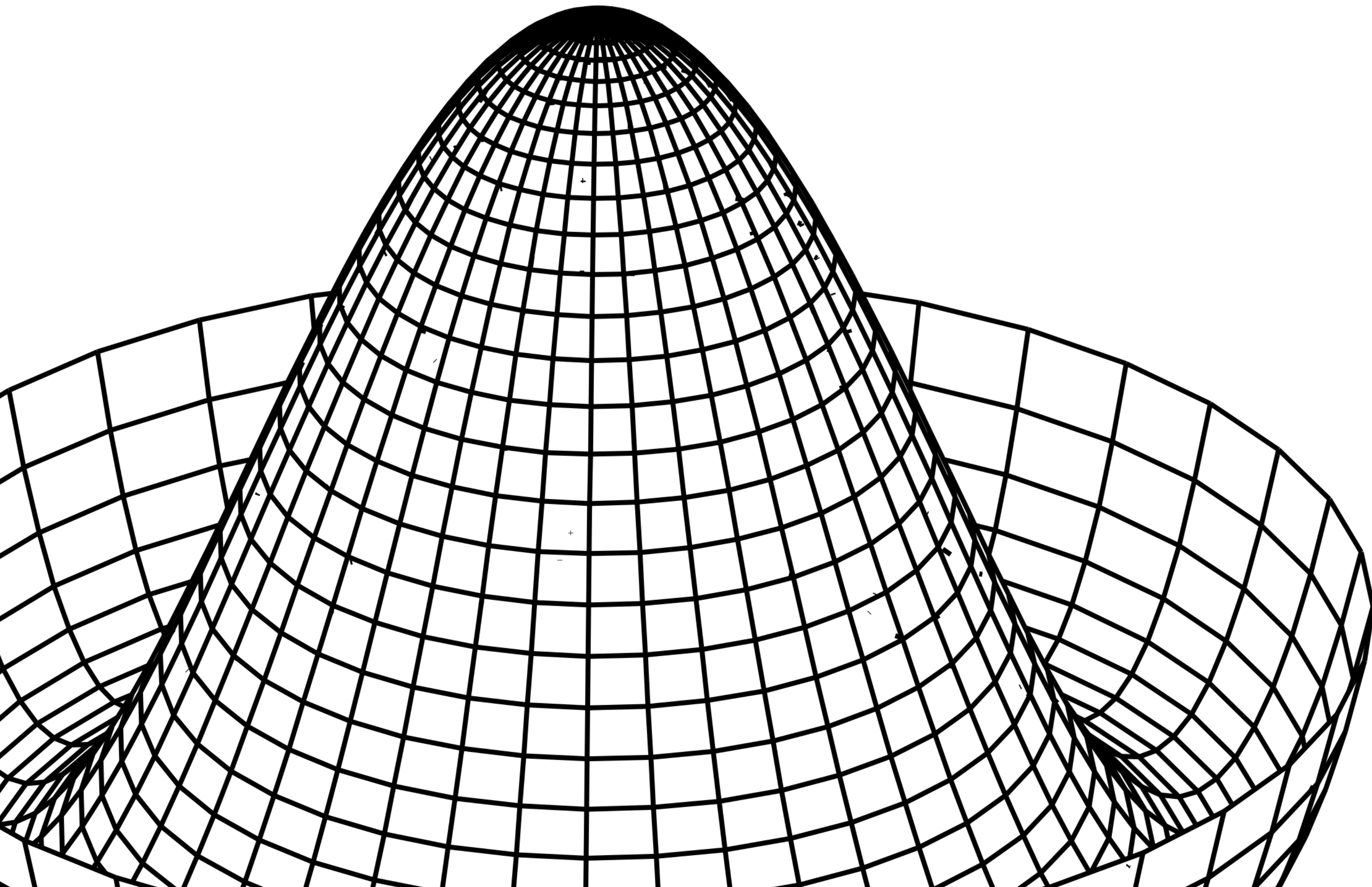


Many beyond SM scenarios predict only %-level deviations from SM !

About 20× more Higgs bosons expected from Run3 and High-Luminosity LHC

→ Harsher experimental conditions require upgrades of our detectors

H boson self-interaction



Does the Higgs boson interact with itself?

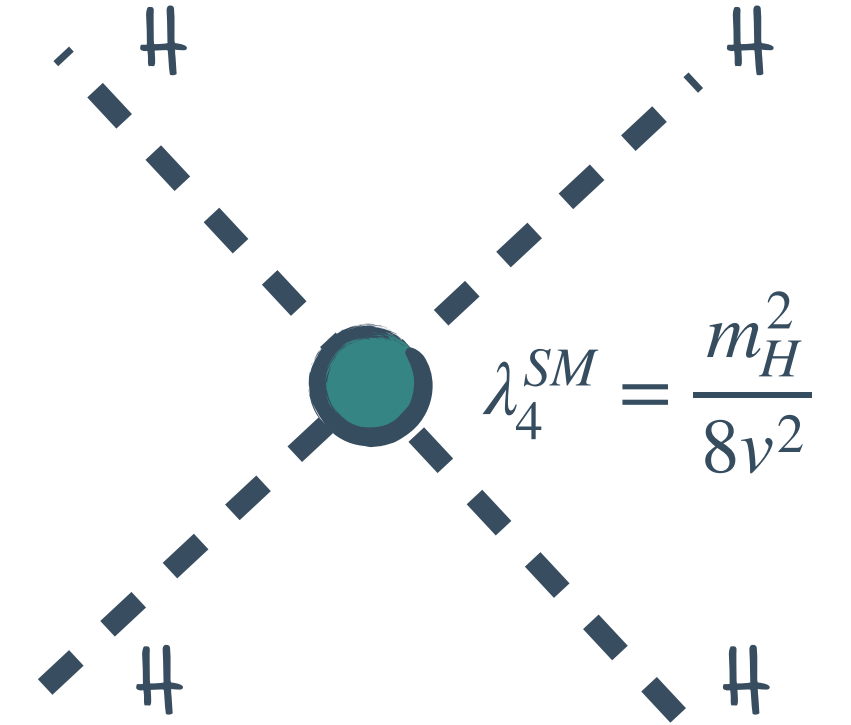
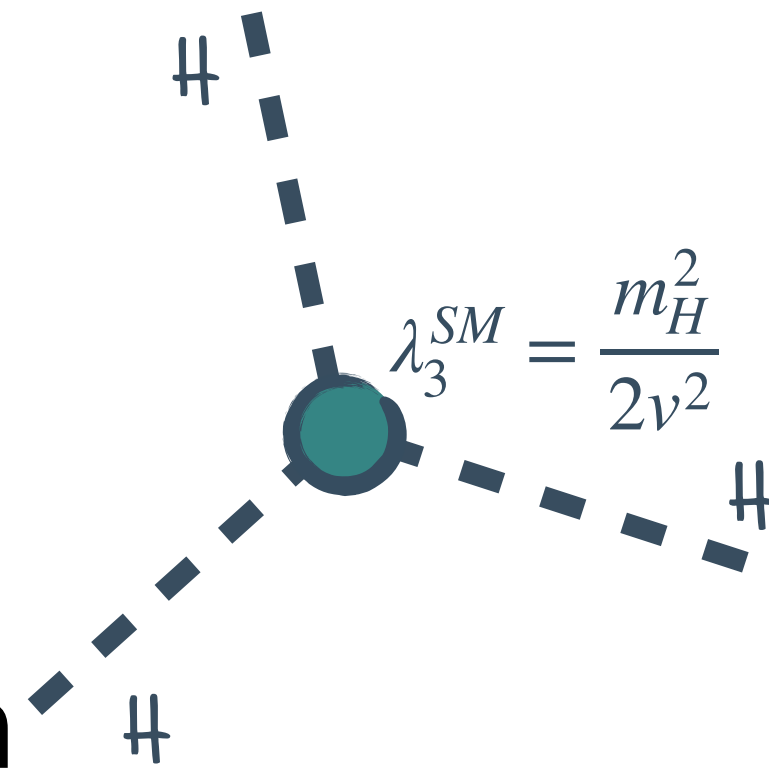
Does the Higgs boson interact with itself?

A self-interacting Higgs (as SM predicts) would be unlike anything yet seen in nature. All other interactions change particle identity.

The Higgs boson **cubic** (λ_3^{SM}) and **quartic** (λ_4^{SM}) couplings are the keys to check the EWSB. The Higgs boson potential is :

$$\mathcal{L} \subset -\frac{m_h^2}{2}h^2 - \lambda_3^{SM}vh^3 - \lambda_4^{SM}h^4$$

Direct test of cubic coupling only **with HH** production



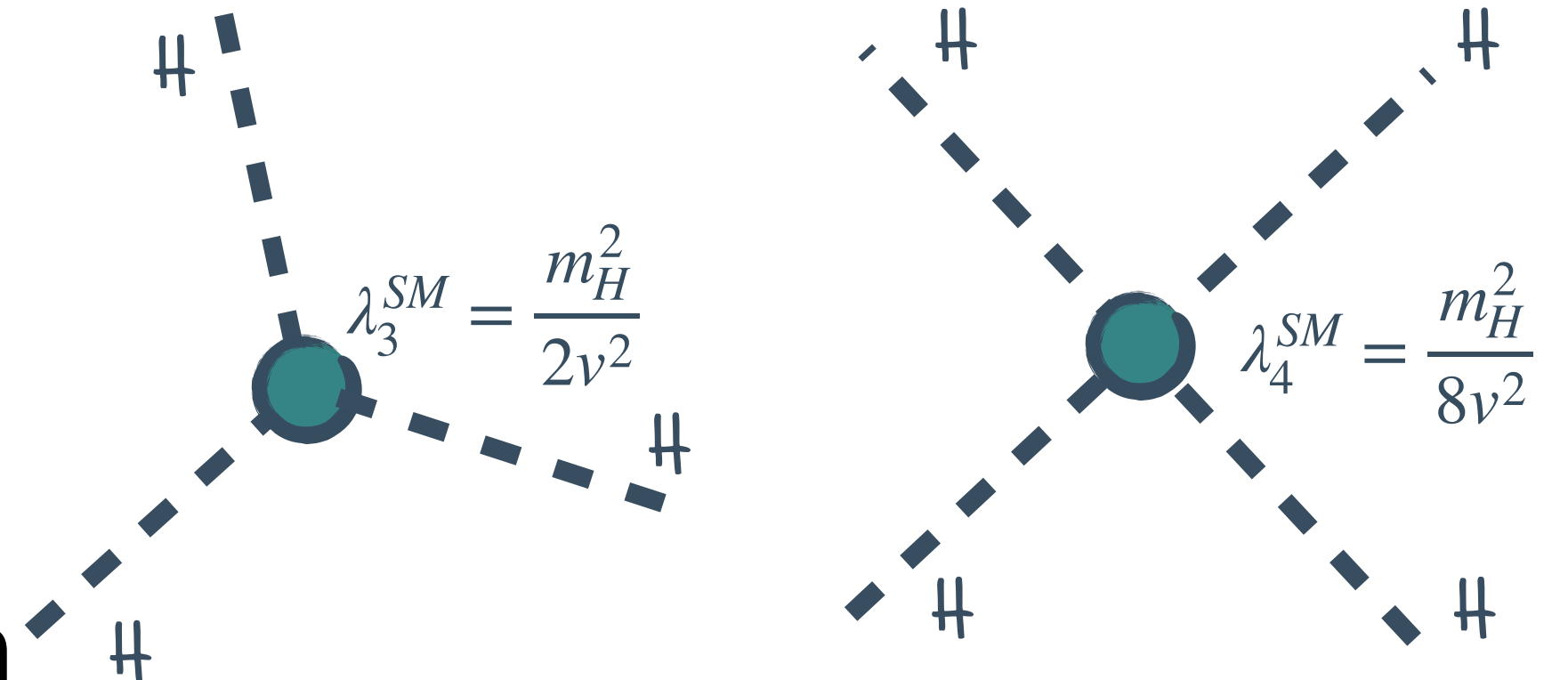
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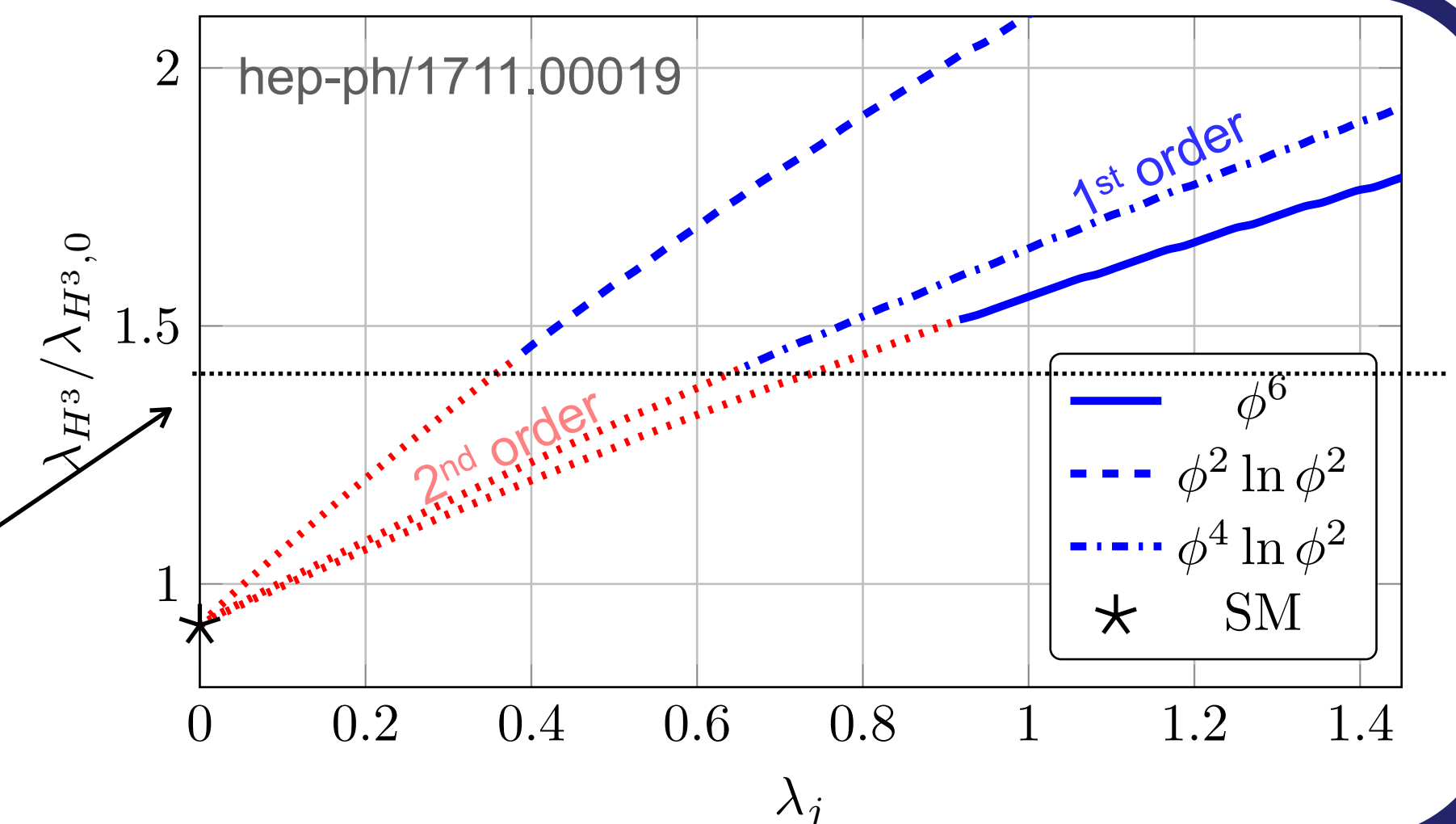
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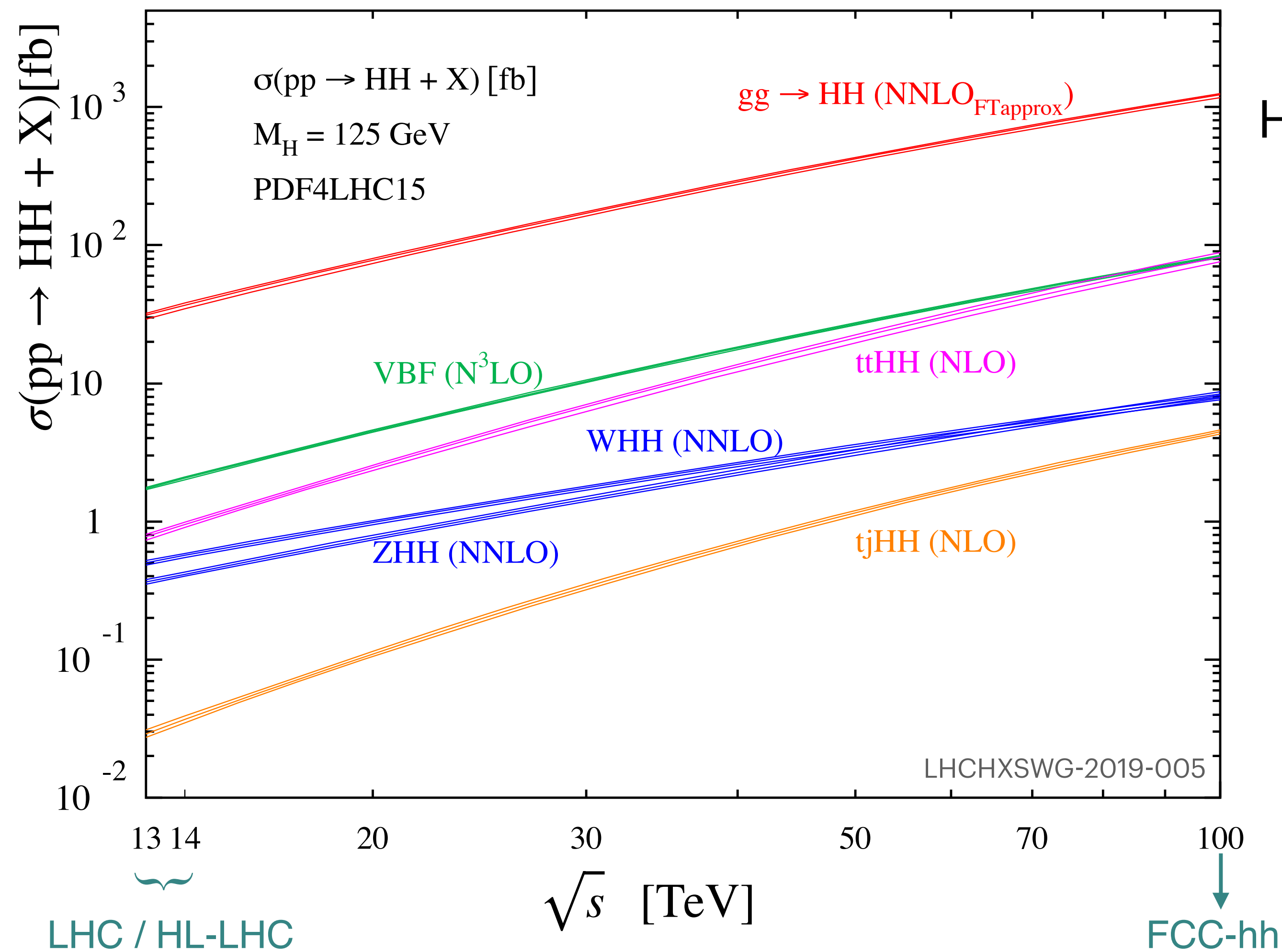
The thermal history of the Universe

Deviations from SM Higgs boson self-coupling cause a modified potential that allows **first-order electroweak phase transition** and hence an explanation of the observed matter vs anti-matter asymmetry!

We need to probe size of modification down to 1.4, the expected uncertainty of the measurement should be $\mathcal{O}(10\%)$

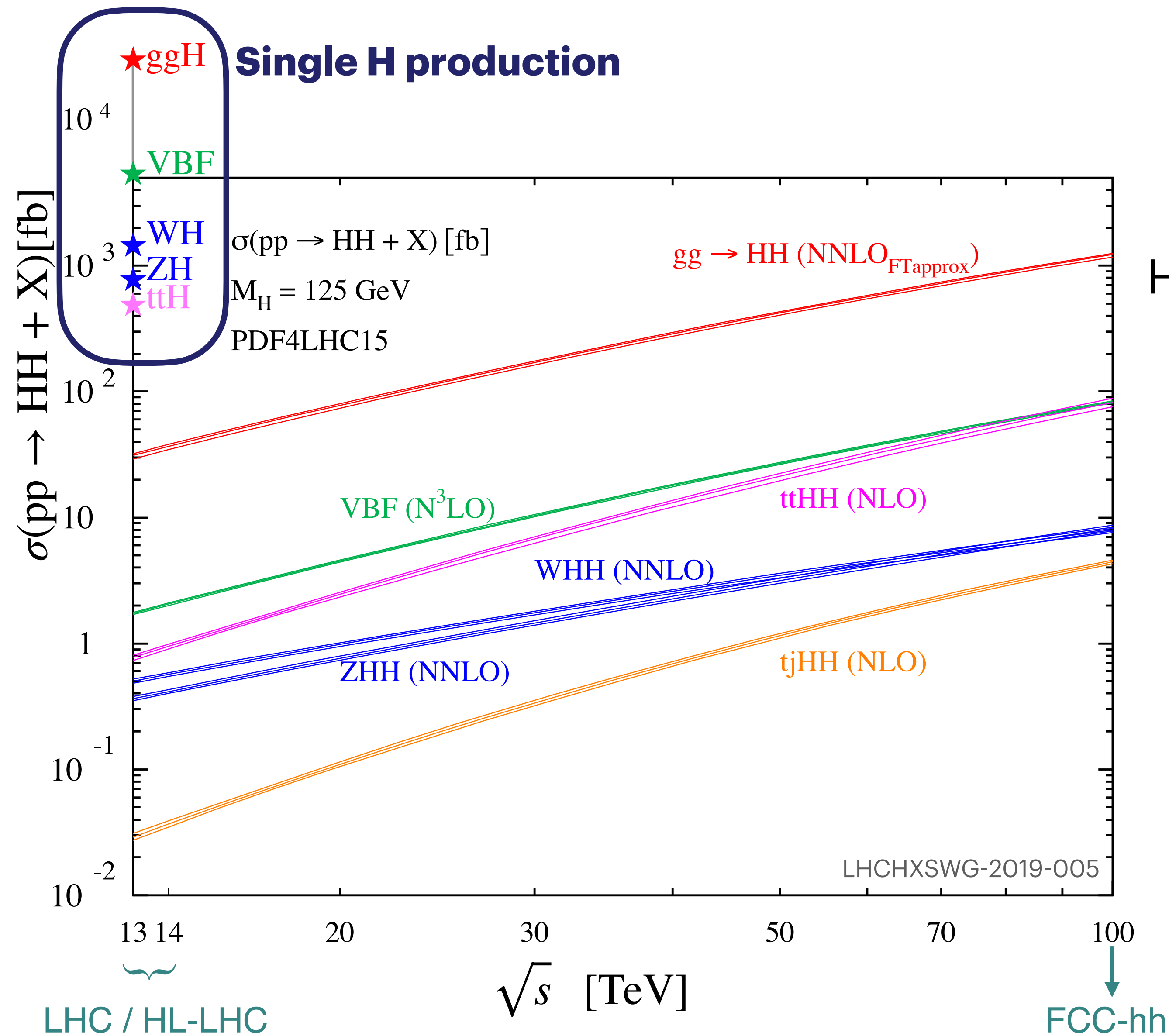


HH production cross sections within SM



HH production channels similar to H production,
but there is a very important difference

HH production cross sections within SM



HH production channels similar to H production, **but** there is a very important difference

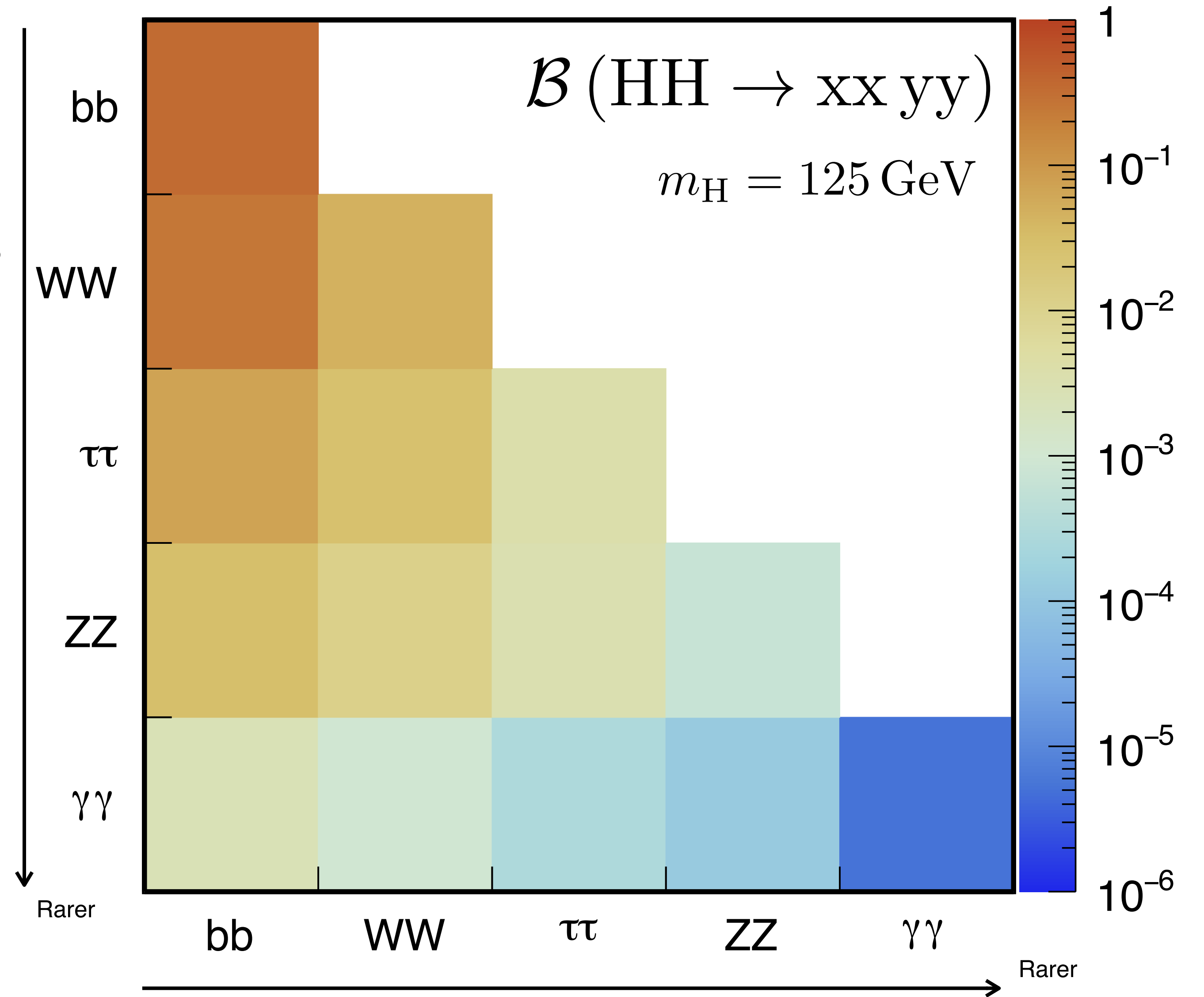
$$\sigma(pp \rightarrow HH) \sim \frac{\sigma(pp \rightarrow H)}{1000}$$

Higgs boson pairs are predicted to be **1000×** rarer than single Higgs

HH final states within SM

A **very** rich set of final states

ATLAS/CMS results with LHC Run 2 covers

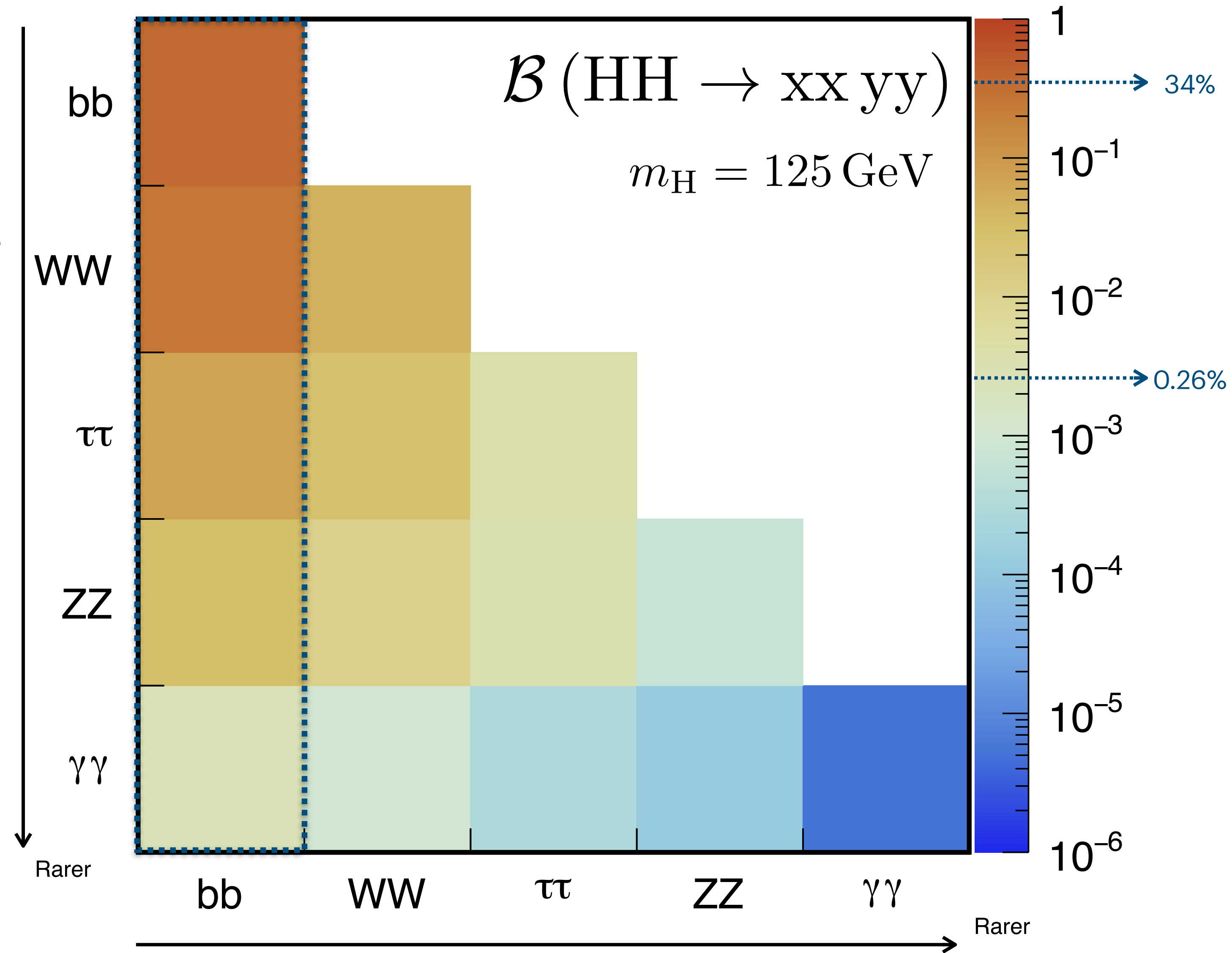


HH final states within SM

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ATLAS/CMS results with LHC Run 2 covers

≥ 1 H decaying to b quarks channels



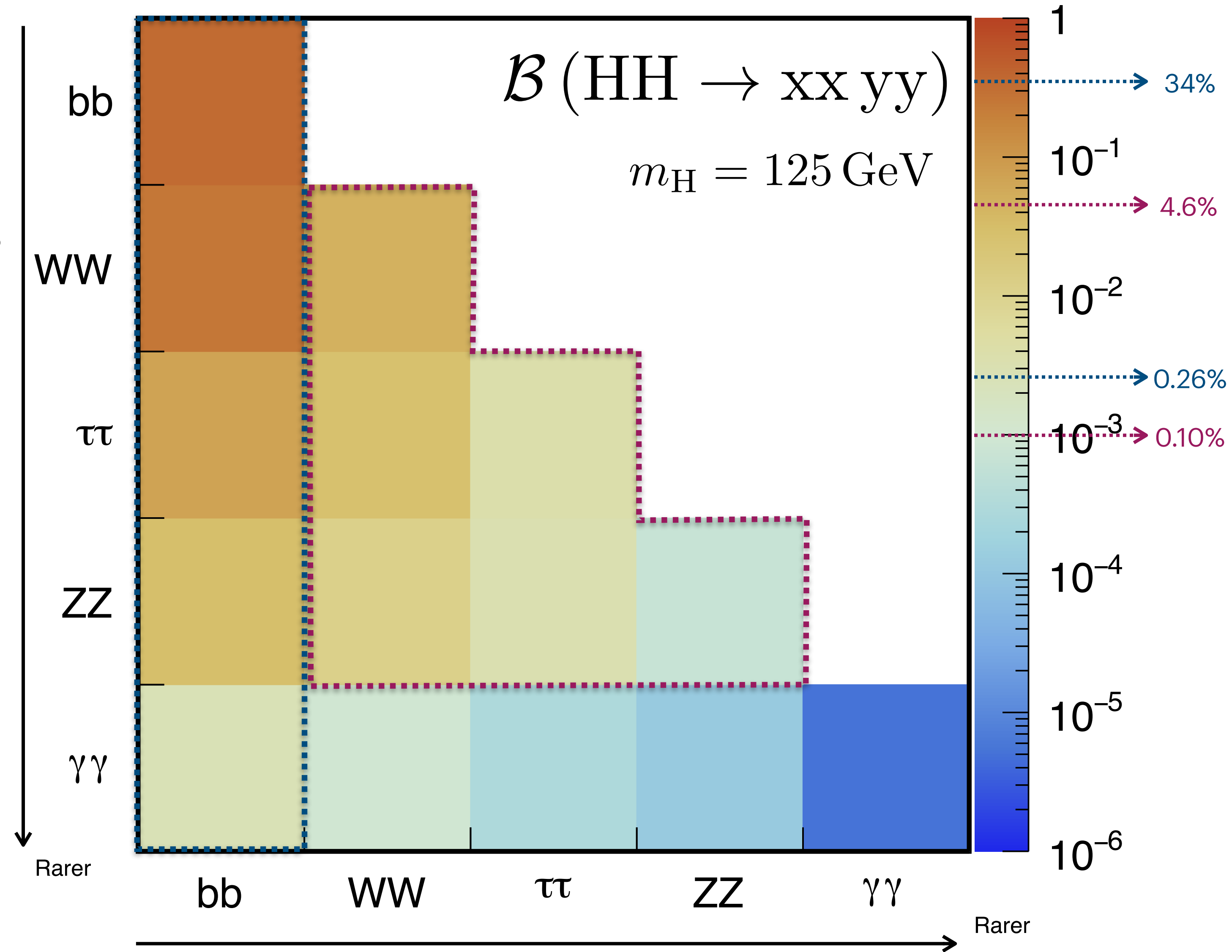
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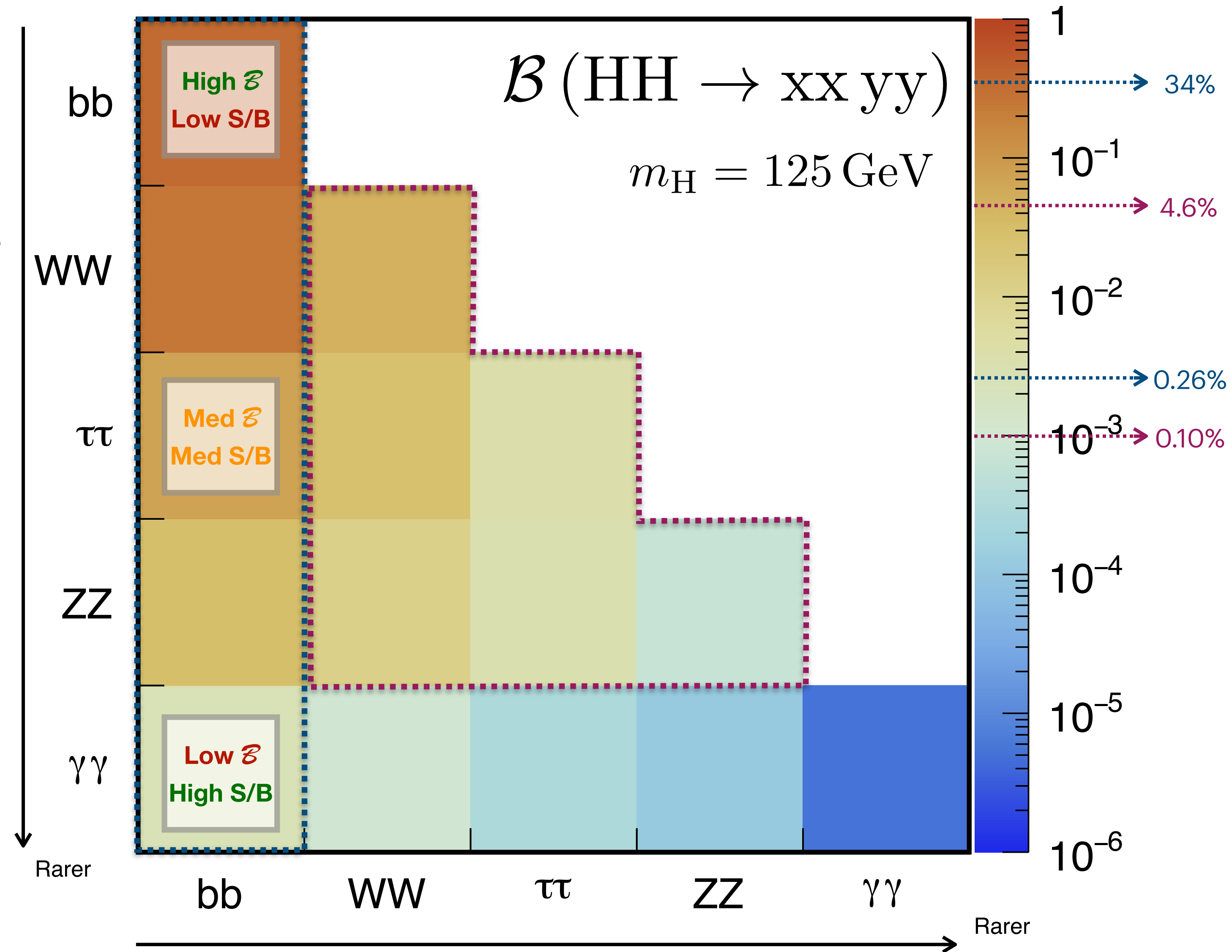
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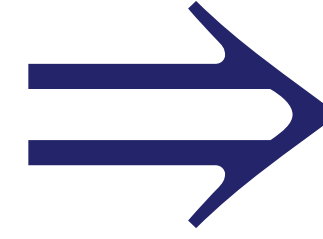
multilepton channels

Not a single **“golden”** channel
but (at least) three **“silver”** bullets



Putting all together

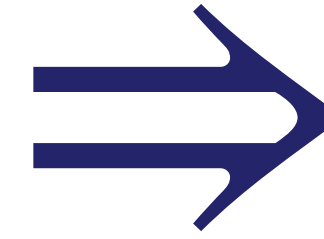
Not a single “golden” channel but various contributions to the overall sensitivity



Combinations are key

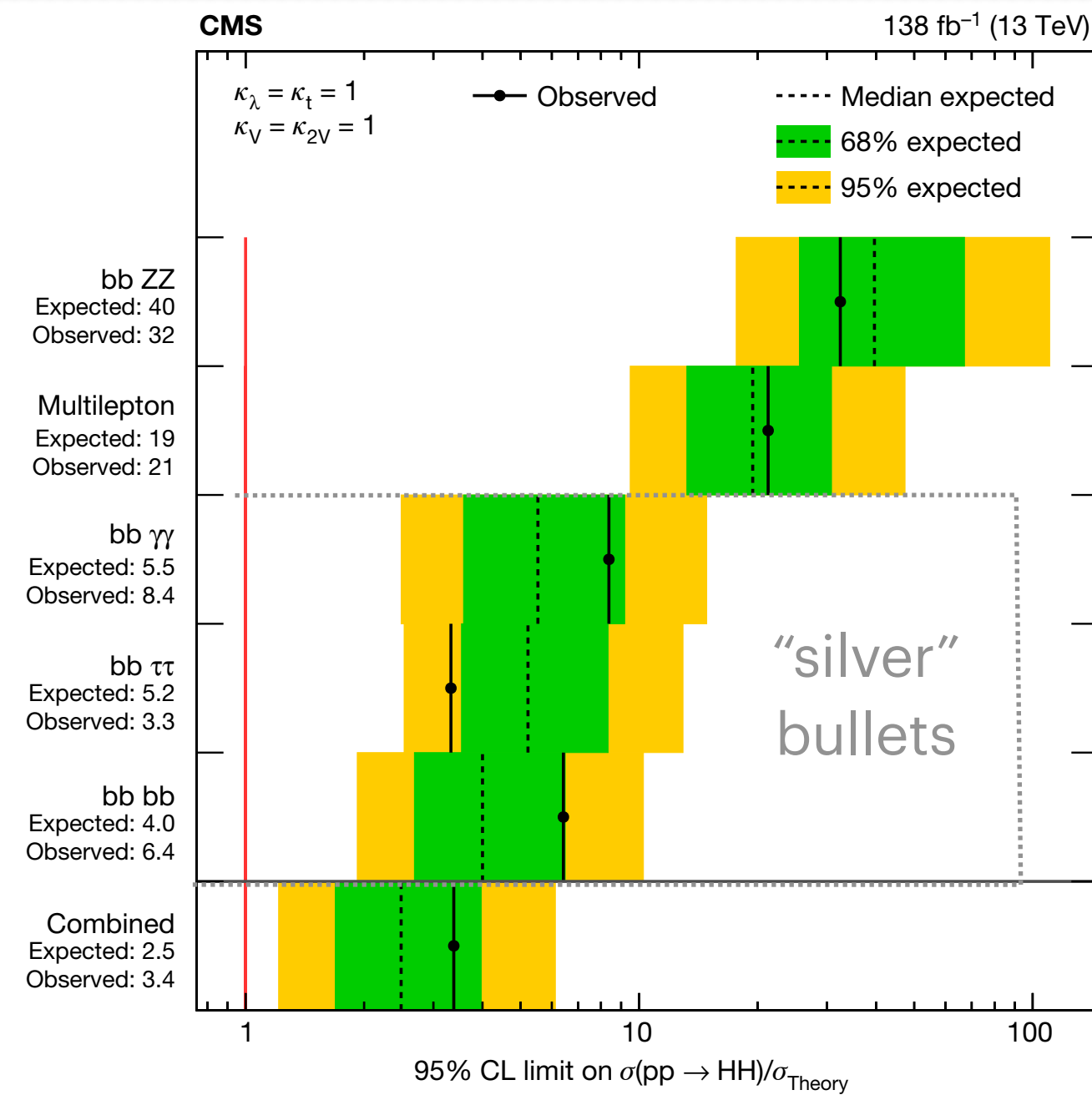
Putting all together : HH production

Not a single “golden” channel but various contributions to the overall sensitivity



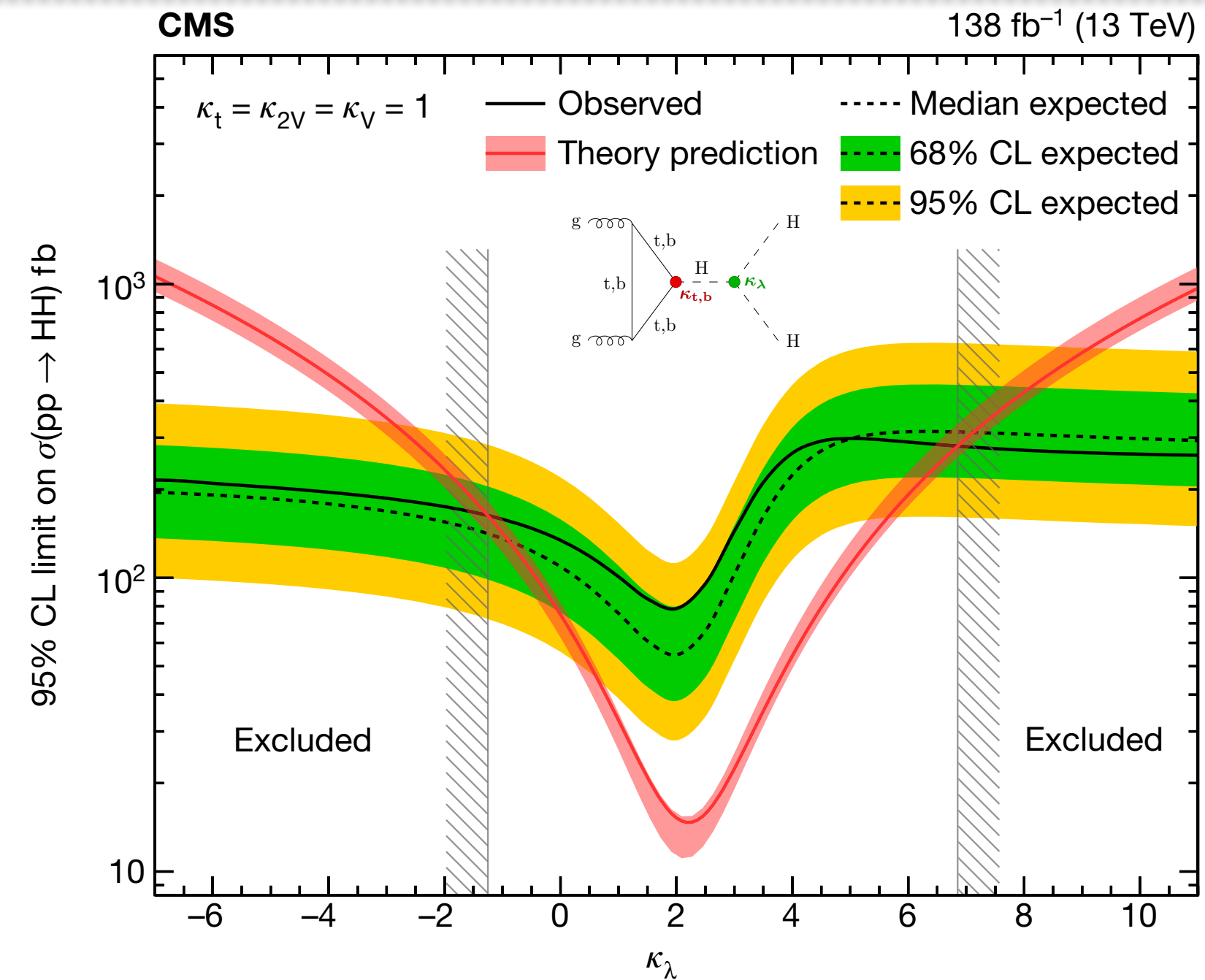
Combinations are key

CMS



Combination →

CMS

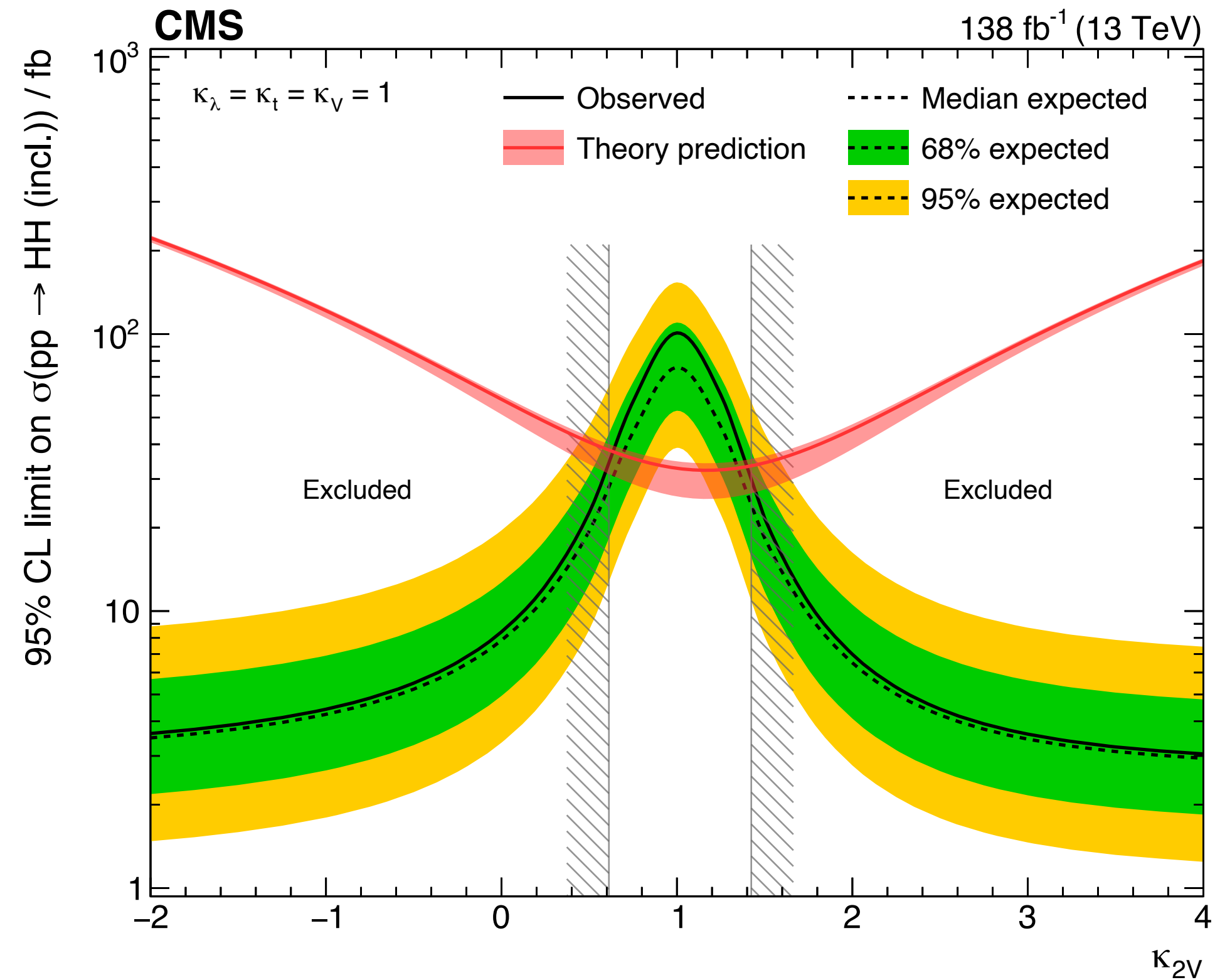
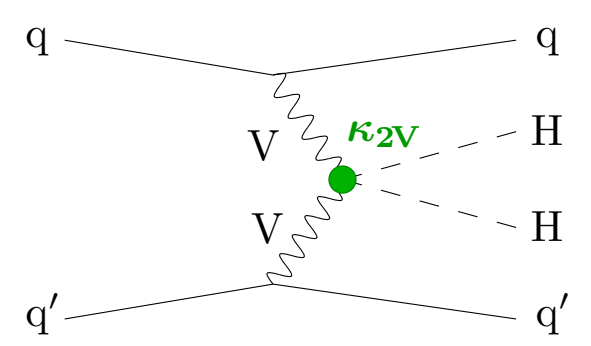


ATLAS

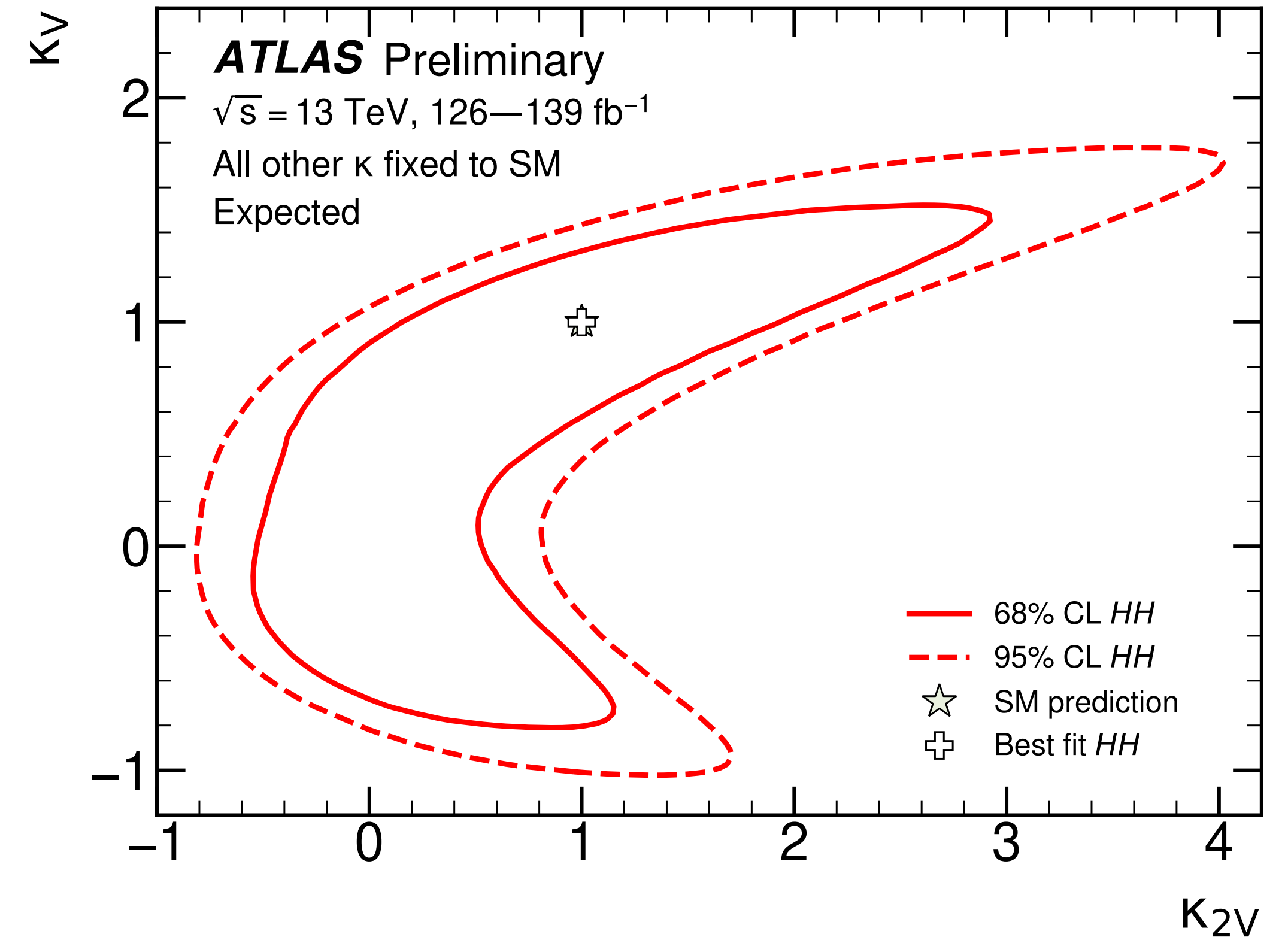
95% CL limit on $\sigma_{HH}/\sigma^{SM}_{HH}$ 2.4(2.9) obs.(exp.)

95% CL limit on $k_\lambda \in [-0.6; +6.6]$

Putting all together : VVHH interaction

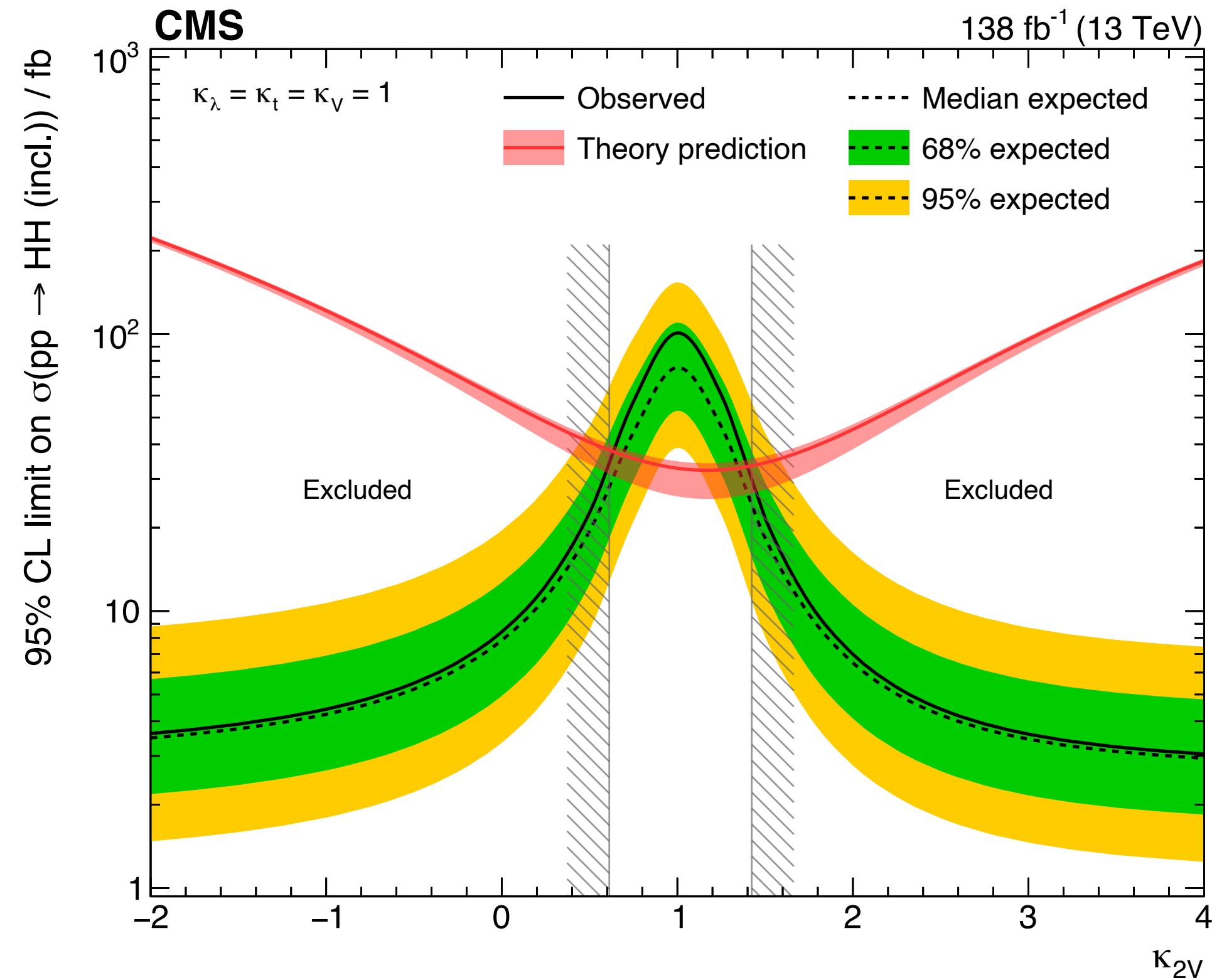
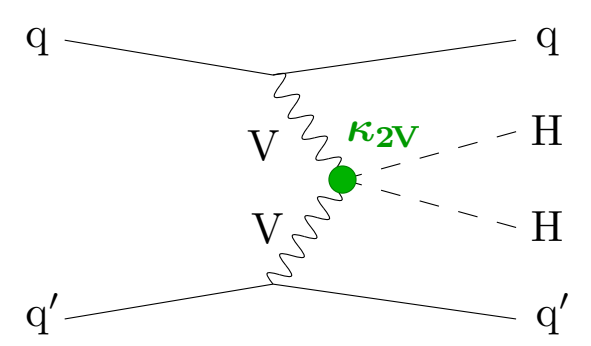


$k_{2V} = 0$ is excluded, with a significance of 6.6 s.d. assuming all other couplings to be SM

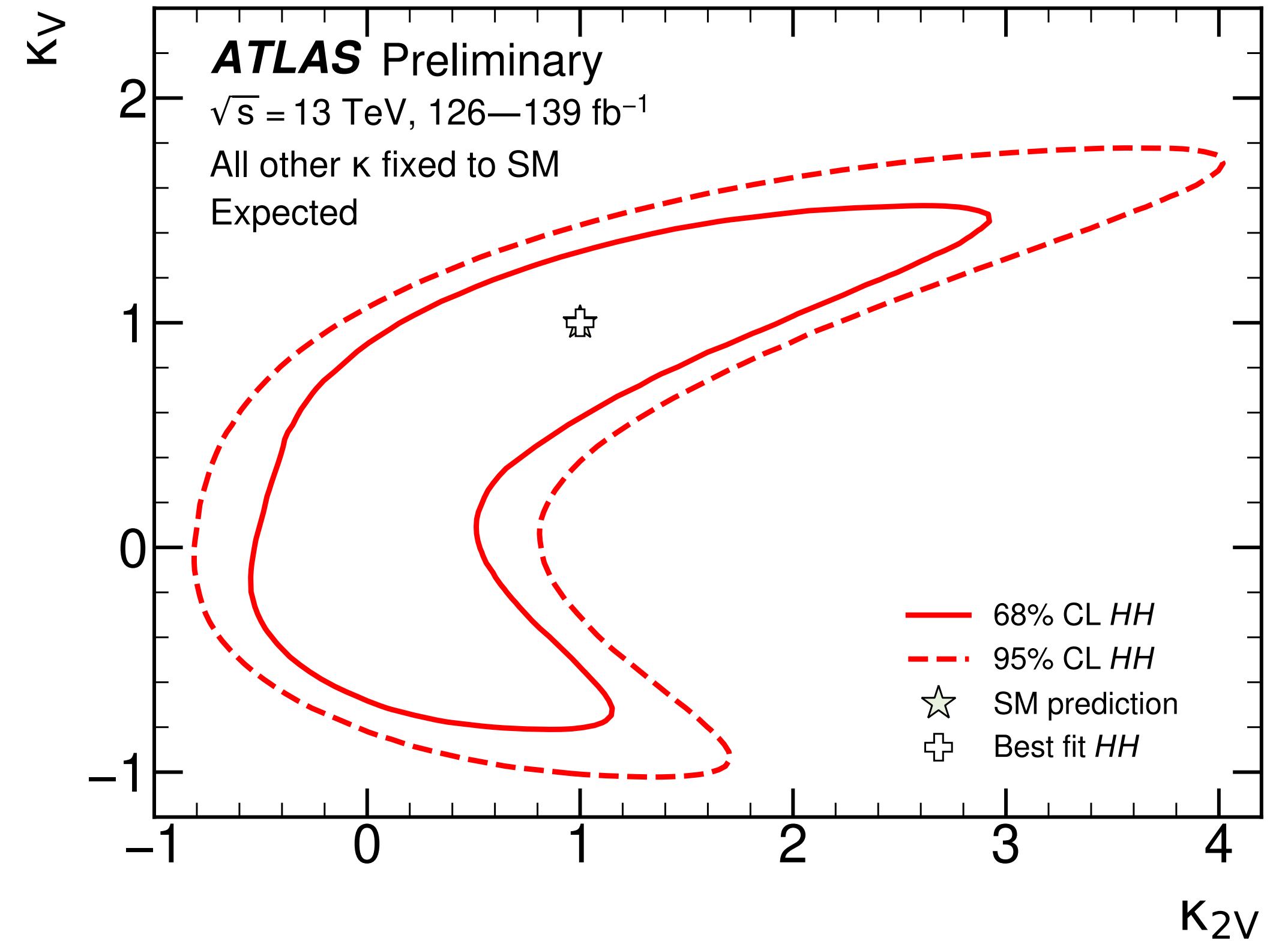


Constrained in the k_V and k_{2V} plane

Putting all together : VVHH interaction



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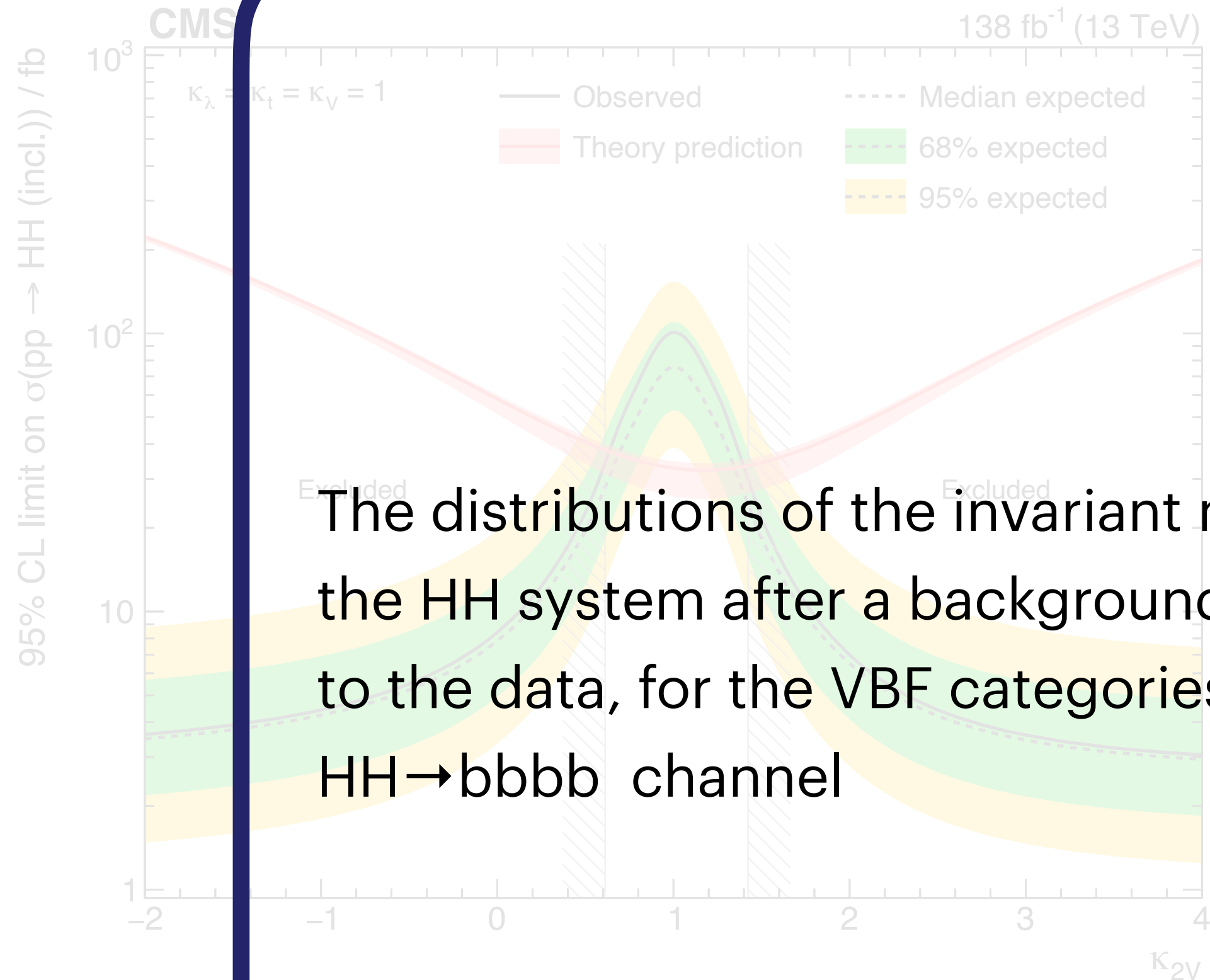
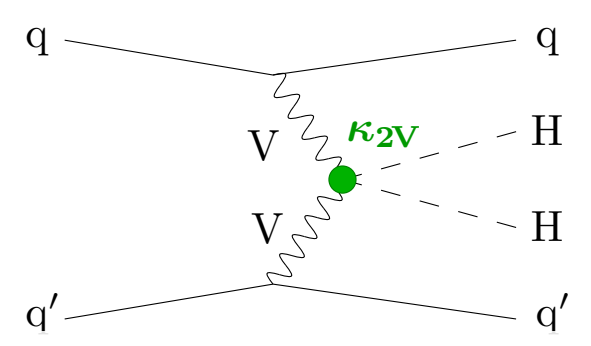


Constrained in the k_V and k_{2V} plane

Establishing the existence of the quartic coupling **VVHH**

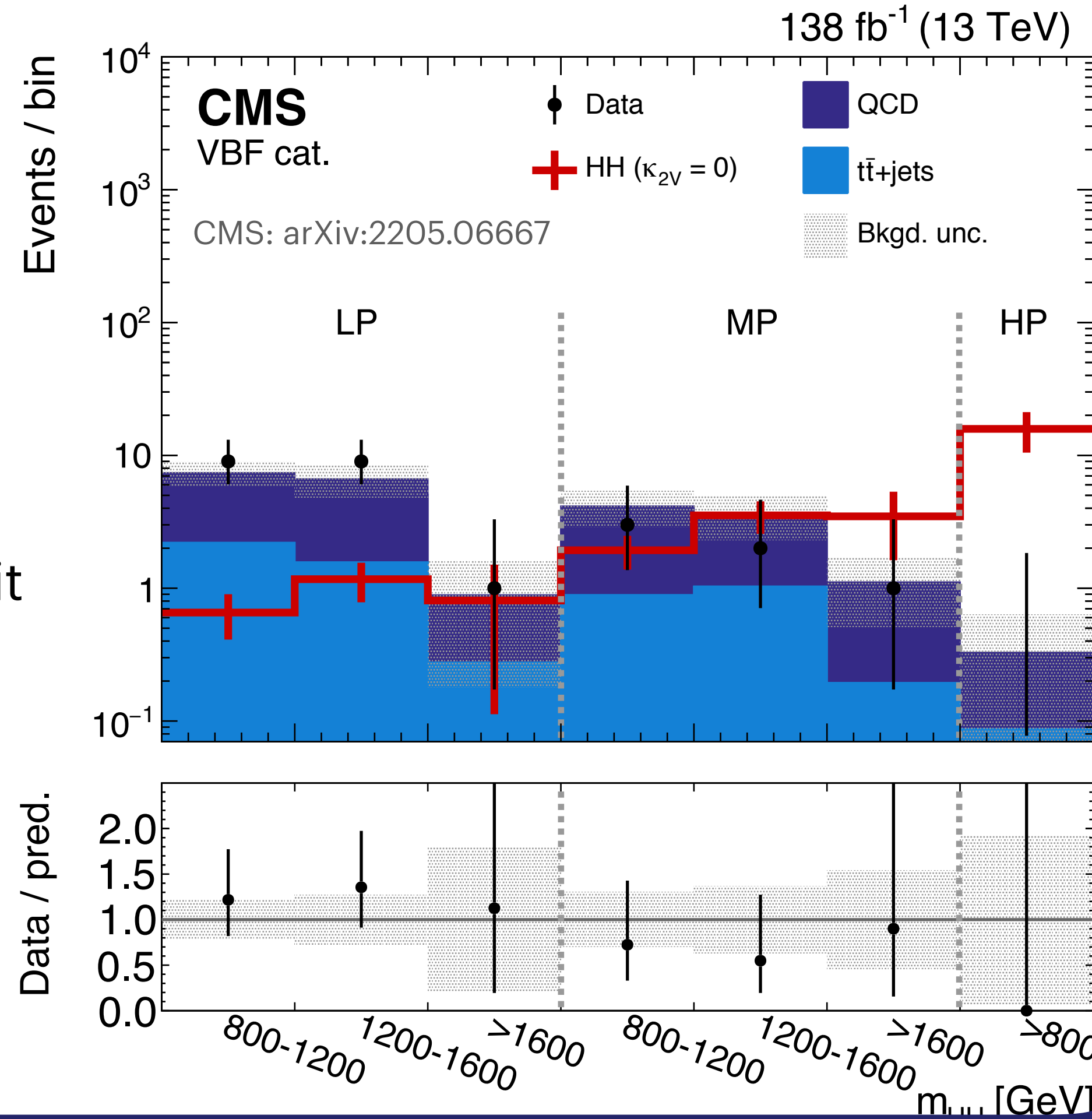
Key role of (HH→bbbb) boosted searches and ML for H→bb decay ID

Putting all together : VVHH interaction



The distributions of the invariant mass of the HH system after a background-only fit to the data, for the VBF categories of the HH→bbbb channel

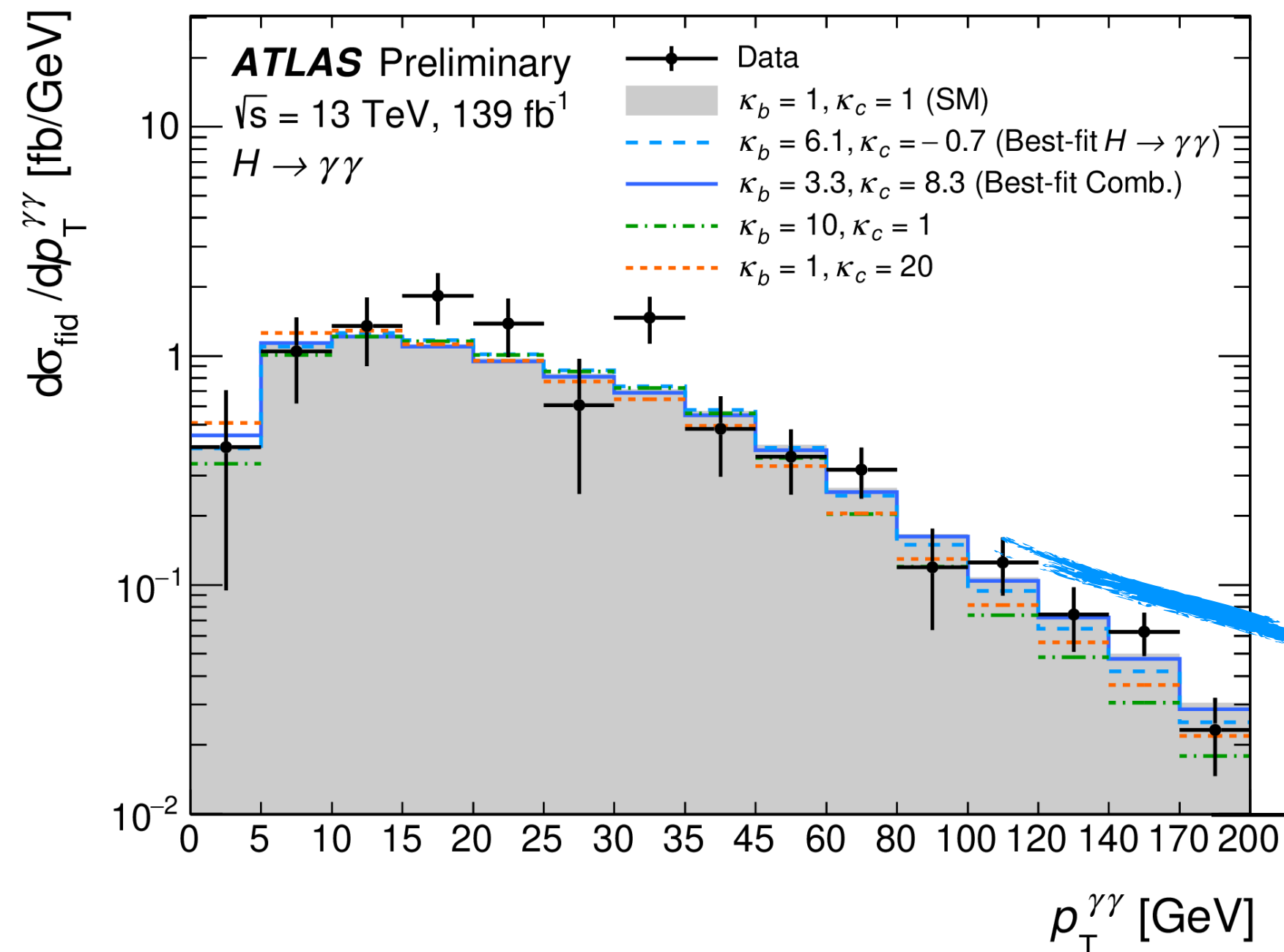
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Establishing the existence of the quartic coupling **VVHH**

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Higgs boson in BSM searches



▲ define the scale of new physics



The role of H in searches for BSM

Few examples

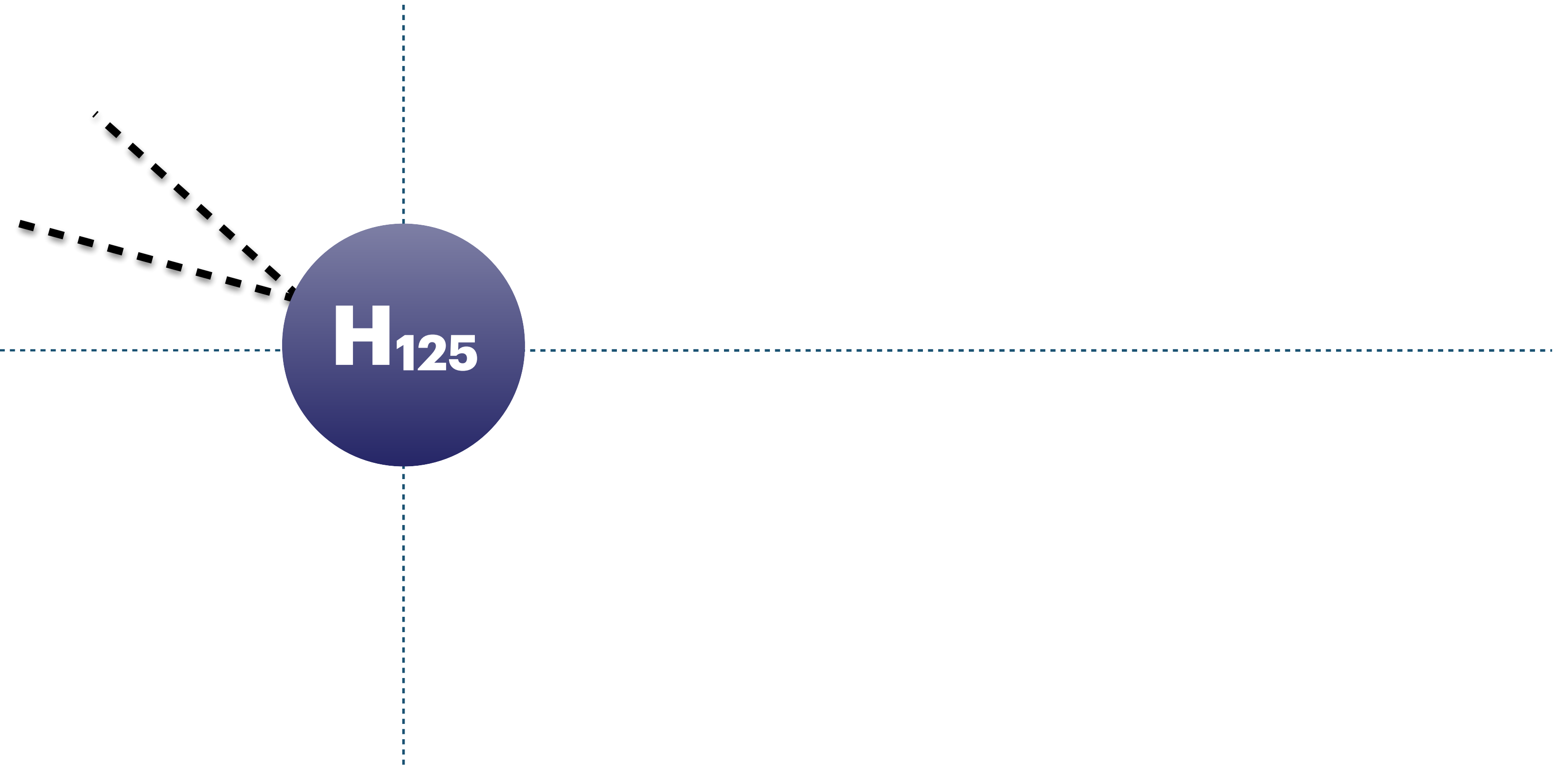


The role of H in searches for BSM

Few examples

invisible decay

H_{125} invisible
and undetected decays



The role of H in searches for BSM

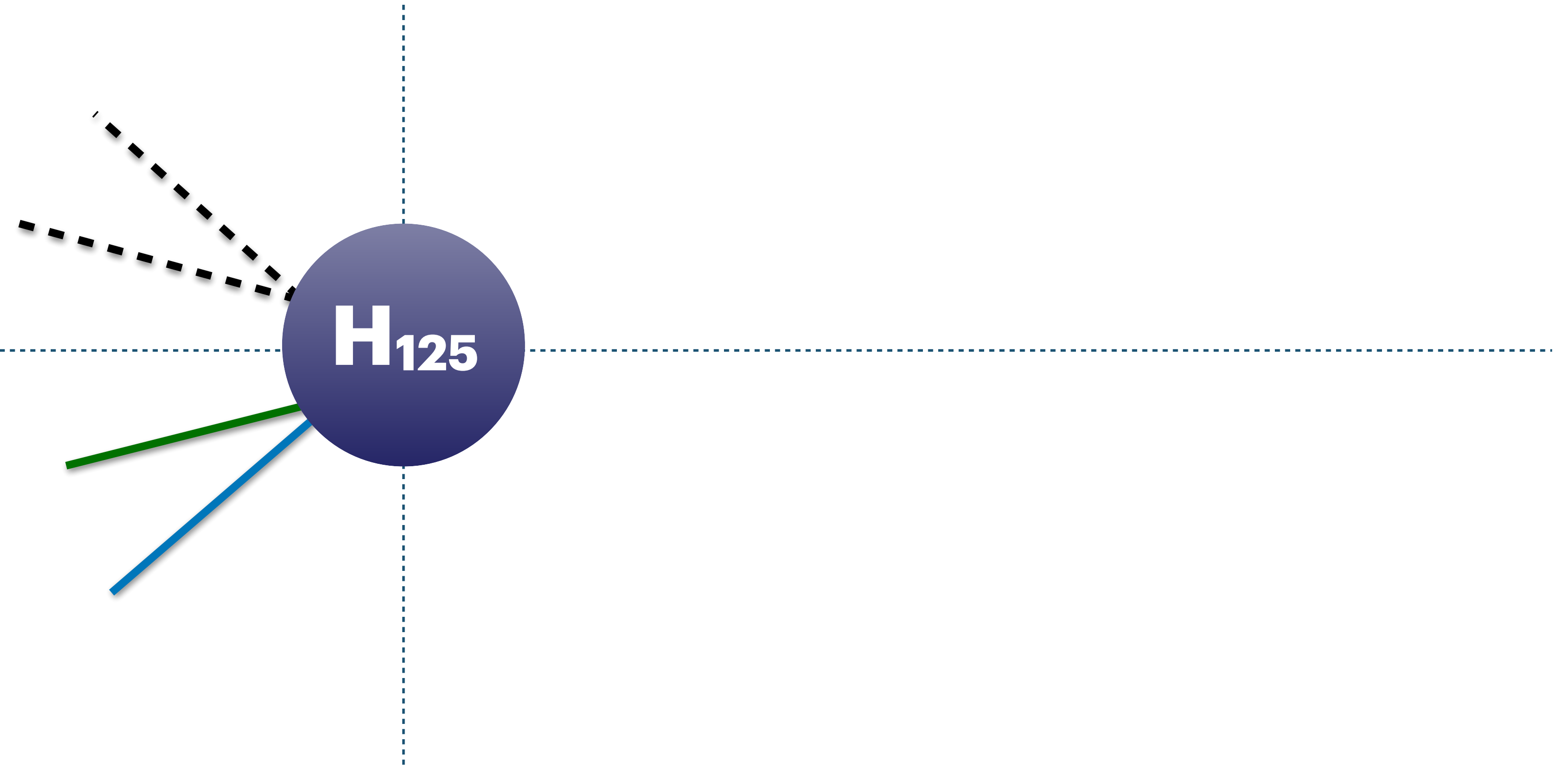
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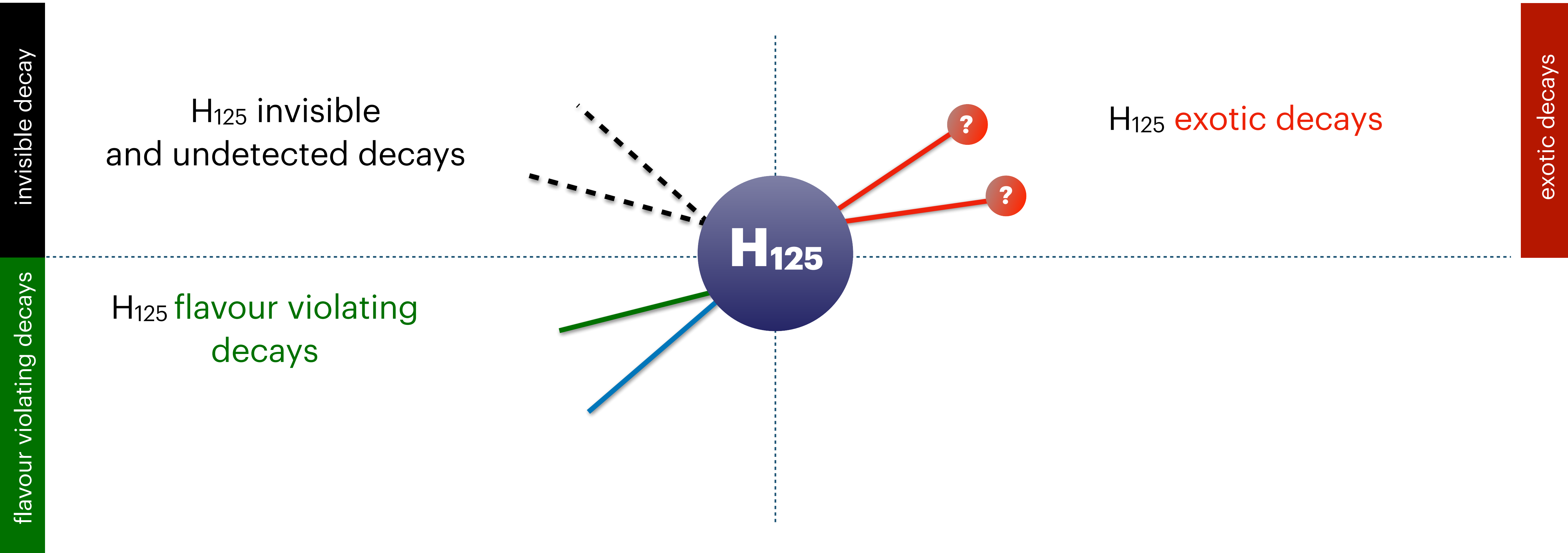
flavour violating decays

H_{125} flavour violating
decays



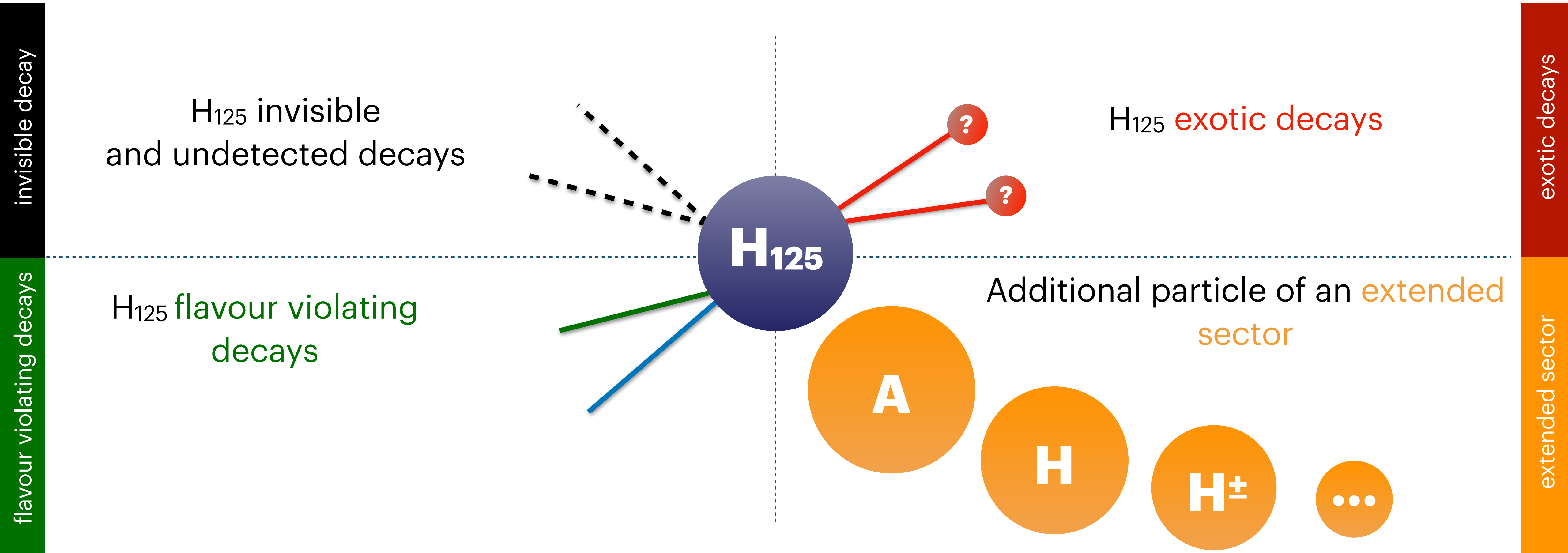
The role of H in searches for BSM

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The role of H in searches for BSM

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Higgs boson to invisible decay

The expected SM H branching fraction to invisible decay (\mathcal{B}_{inv}) is 0.12% due to $H \rightarrow ZZ^* \rightarrow \nu\bar{\nu}\nu\bar{\nu}$

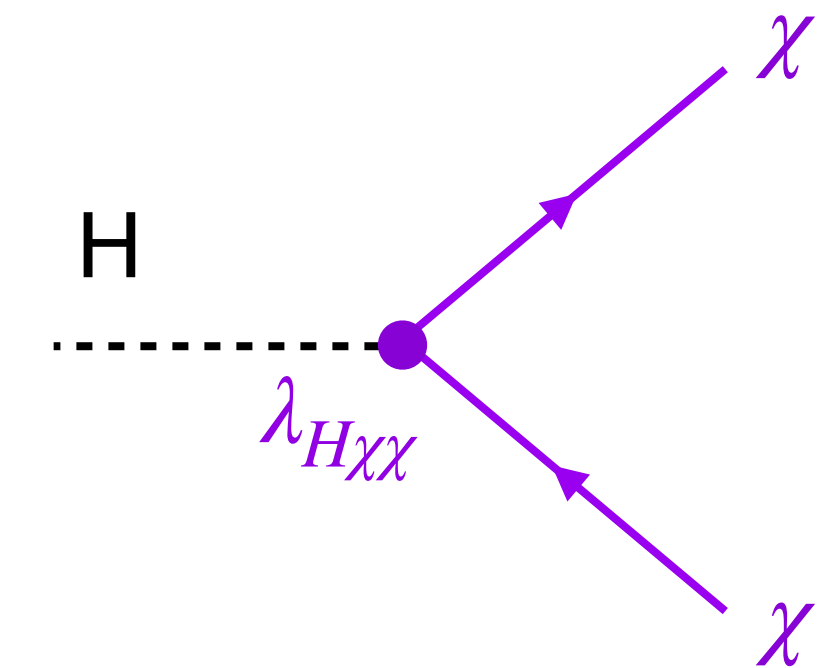
Several BSM scenario \Rightarrow anomalous and sizeable values, \mathcal{B}_{inv} is significantly enhanced.

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In one class of models H decay in a pair of stable WIMPs.
They represent a simple extension of the SM to provide a Dark Matter (DM) candidate and are able to predict the observed relic DM density via s -channel $\chi\chi \rightarrow f\bar{f}$ annihilation.



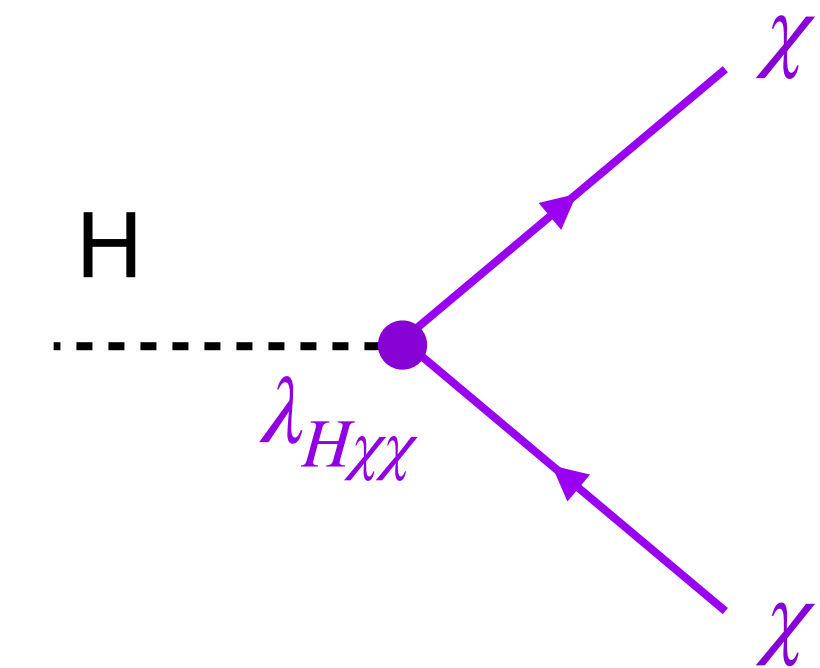
The solution of the DM problem could be found within the Higgs sector.

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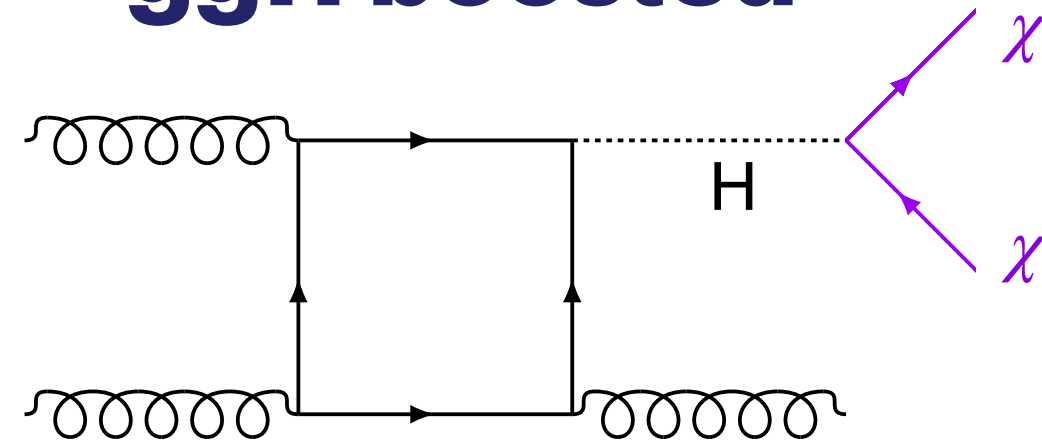
The solution of the DM problem could be found within the Higgs sector.

Common signature : significant missing transverse momentum from the Higgs boson decay.

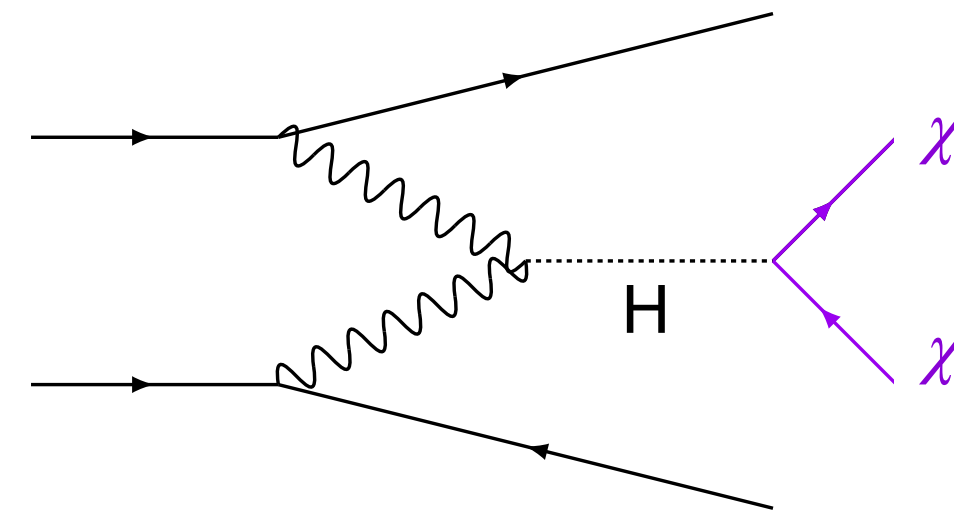
Identify the event : profit of visible particles recoiling against the Higgs boson.

Search for Higgs boson to invisible decay

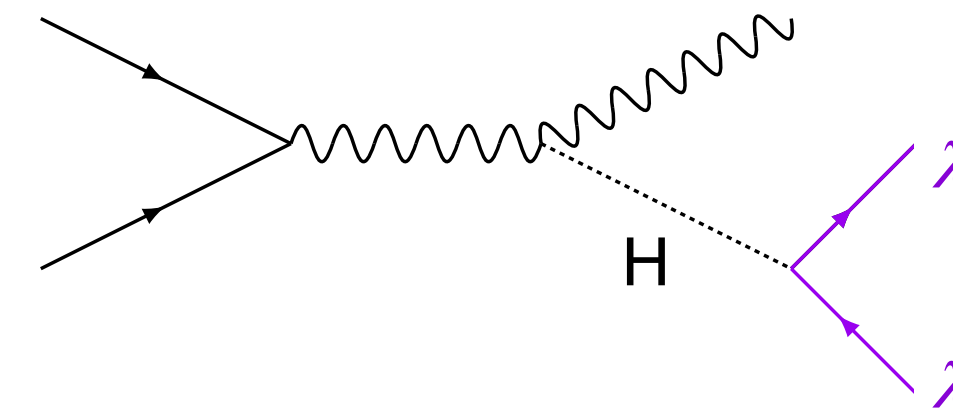
ggH boosted



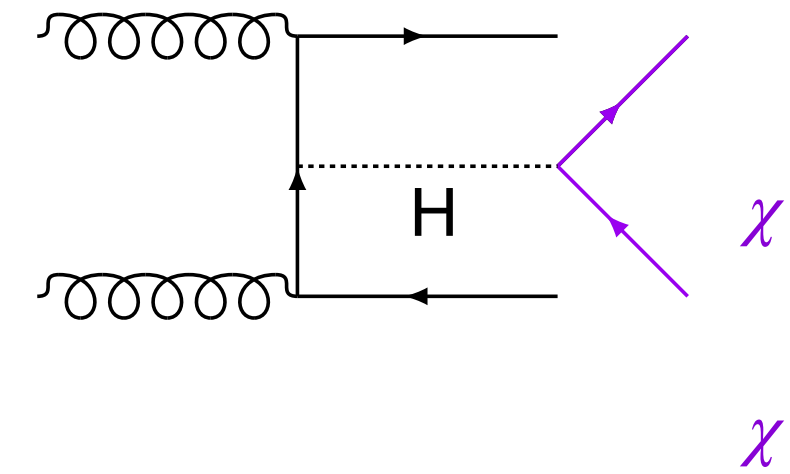
VBF



Associated production



ttH



ATLAS and CMS probe all production mechanisms

Run2 analyses



Phys.Rev.D 103 (2021) 11, 112006



JHEP 11 (2021) 153

arXiv:2202.07953

Eur.Phys.J.C 82 (2022) 2, 105

Phys.Rev.D 105 (2022) 092007

Phys.Lett.B 829 (2022) 137066

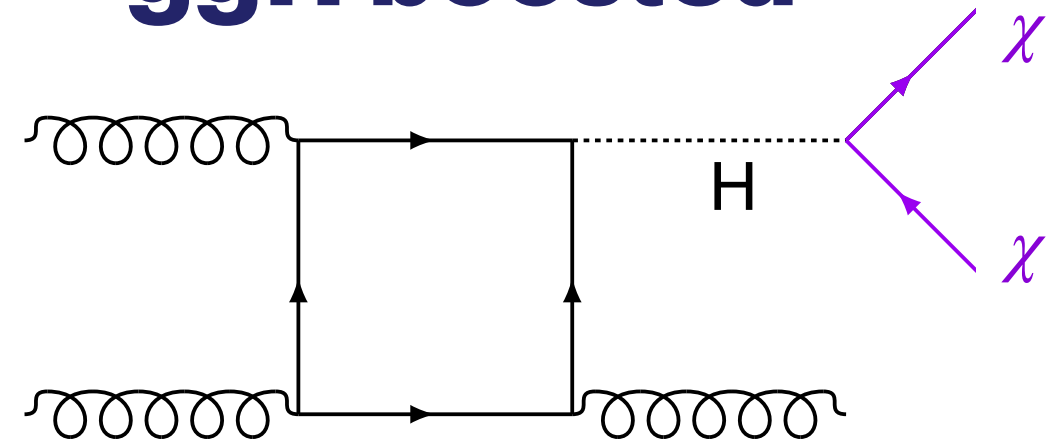
Eur.Phys.J.C 81 (2021)

JHEP 11 (2021) 153

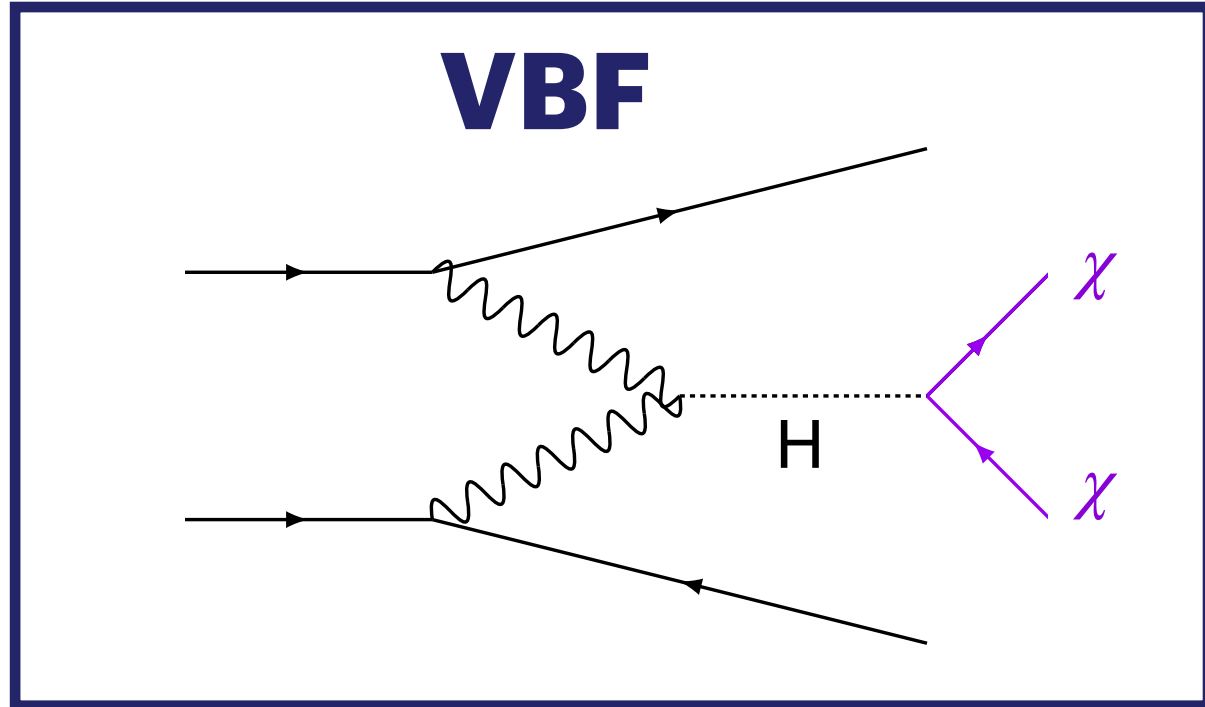
ATLAS-CONF-2022-007

Search for Higgs boson to invisible decay

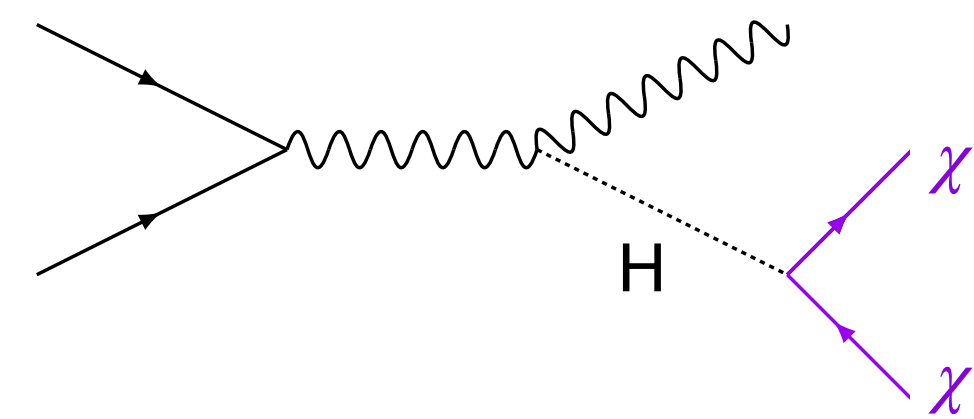
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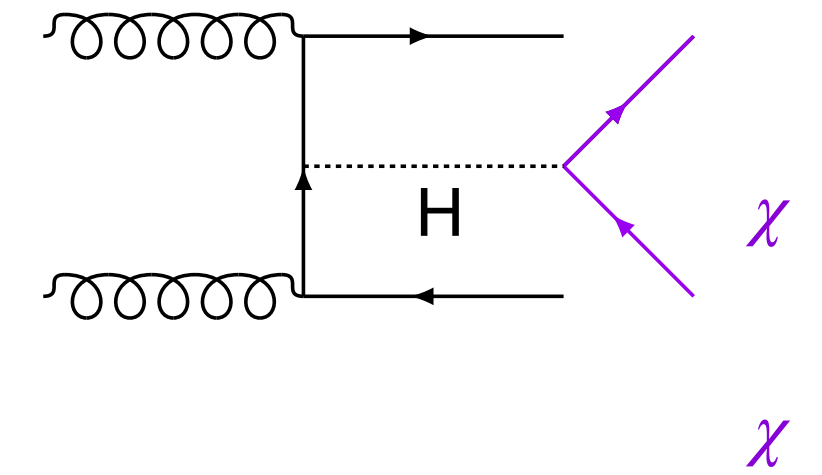
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ATLAS-CONF-2022-007

The VBF production mechanism drives the overall sensitivity in the direct search for invisible decays of the Higgs boson, thanks to its large production cross section and **distinctive event topology**

VBF Higgs boson to invisible decay



arXiv:2202.07953

Phys.Rev.D 105 (2022) 092007

Strategy

2 jets with large angular separation $\Delta\eta_{jj}$ and large invariant mass m_{jj}

Veto on other objects (leptons/photons)

High missing transverse momentum (trigger constraint) \rightarrow reject QCD

Low $|\Delta\phi_{jj}| \rightarrow$ reject QCD

\Rightarrow Main remaining backgrounds: $Z(\nu\nu) +$ jets and $W(l\nu) +$ jets (strong and VBF productions)

VBF Higgs boson to invisible decay

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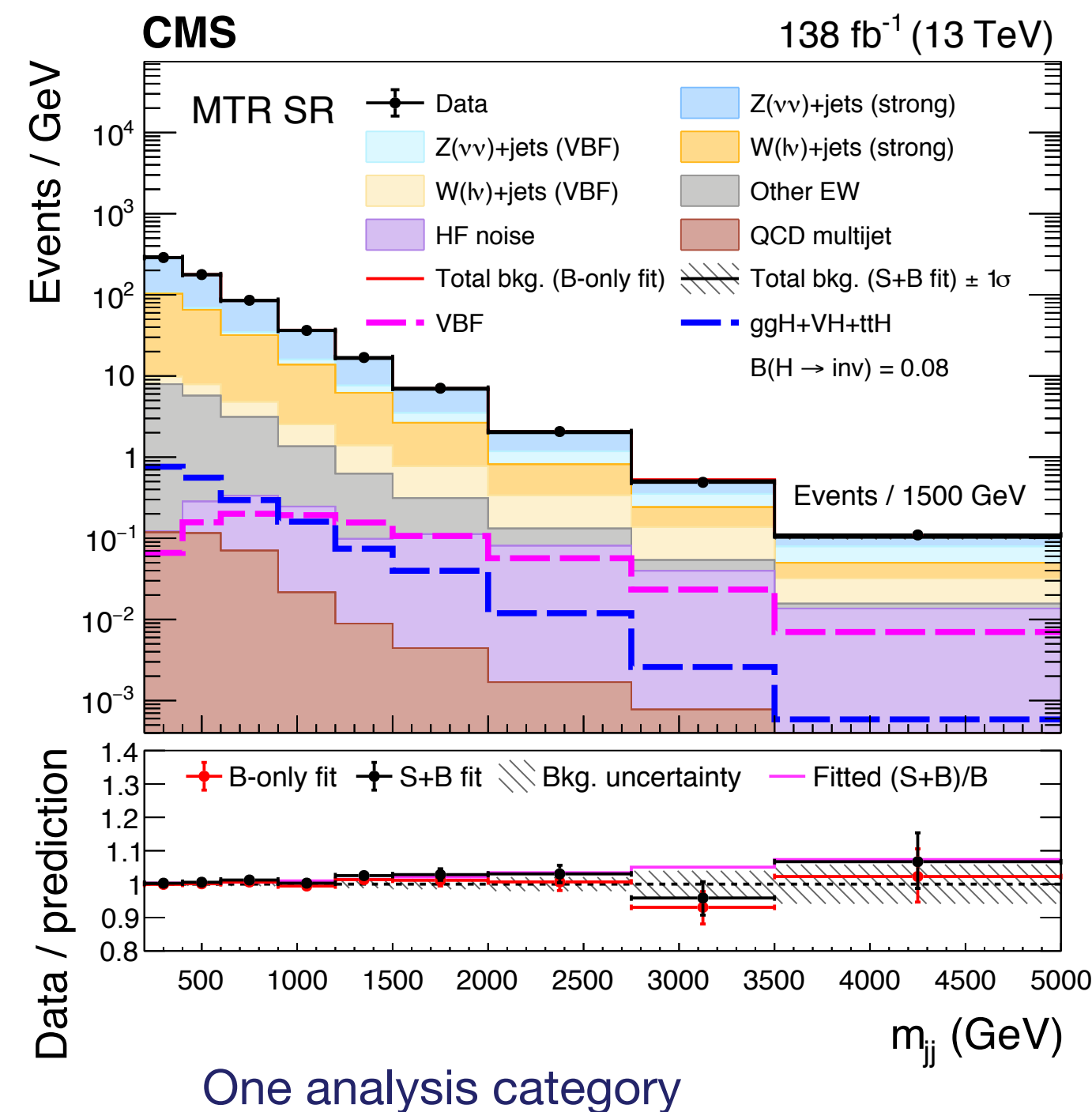
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VBF Higgs boson to invisible decay

Strategy

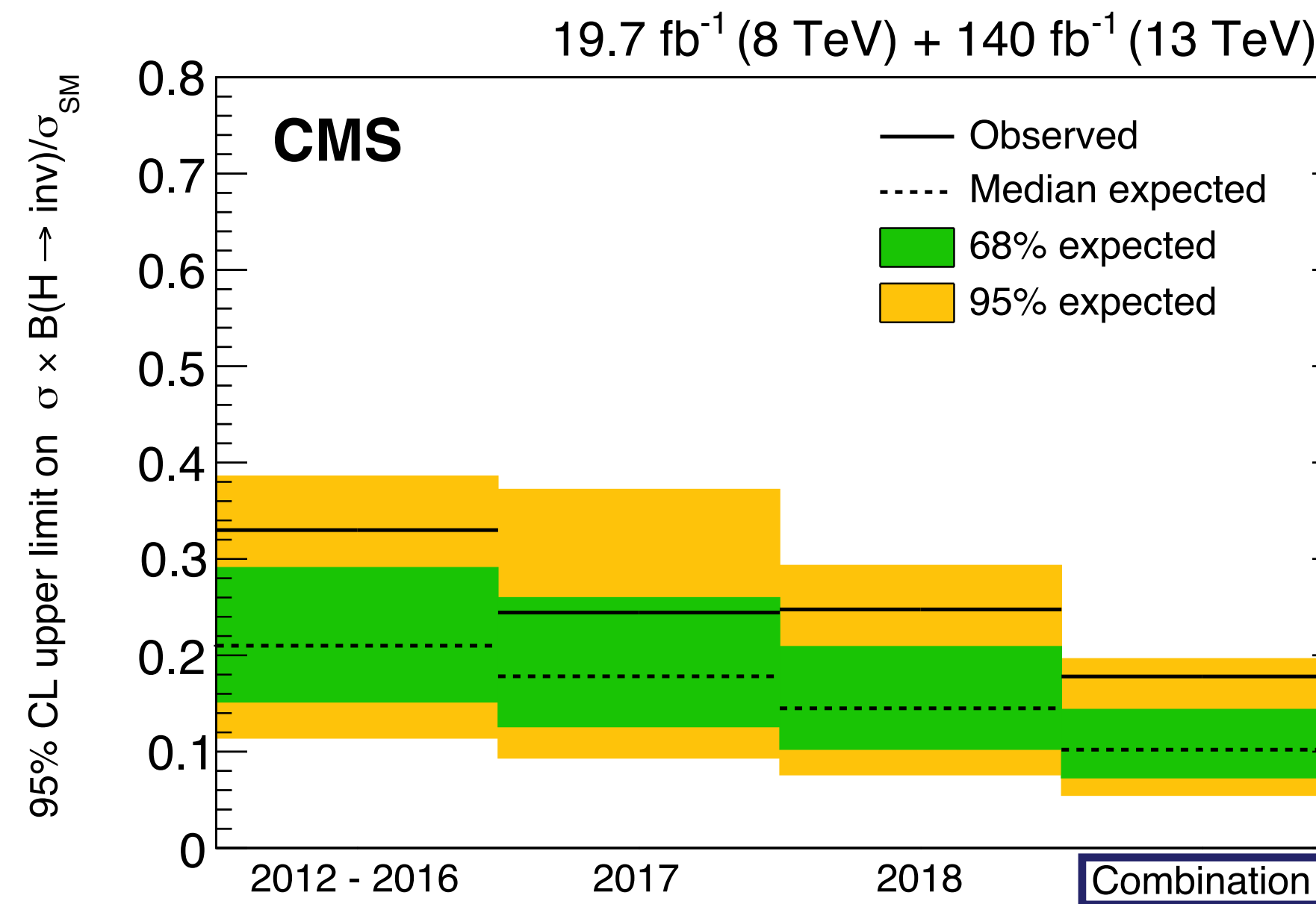
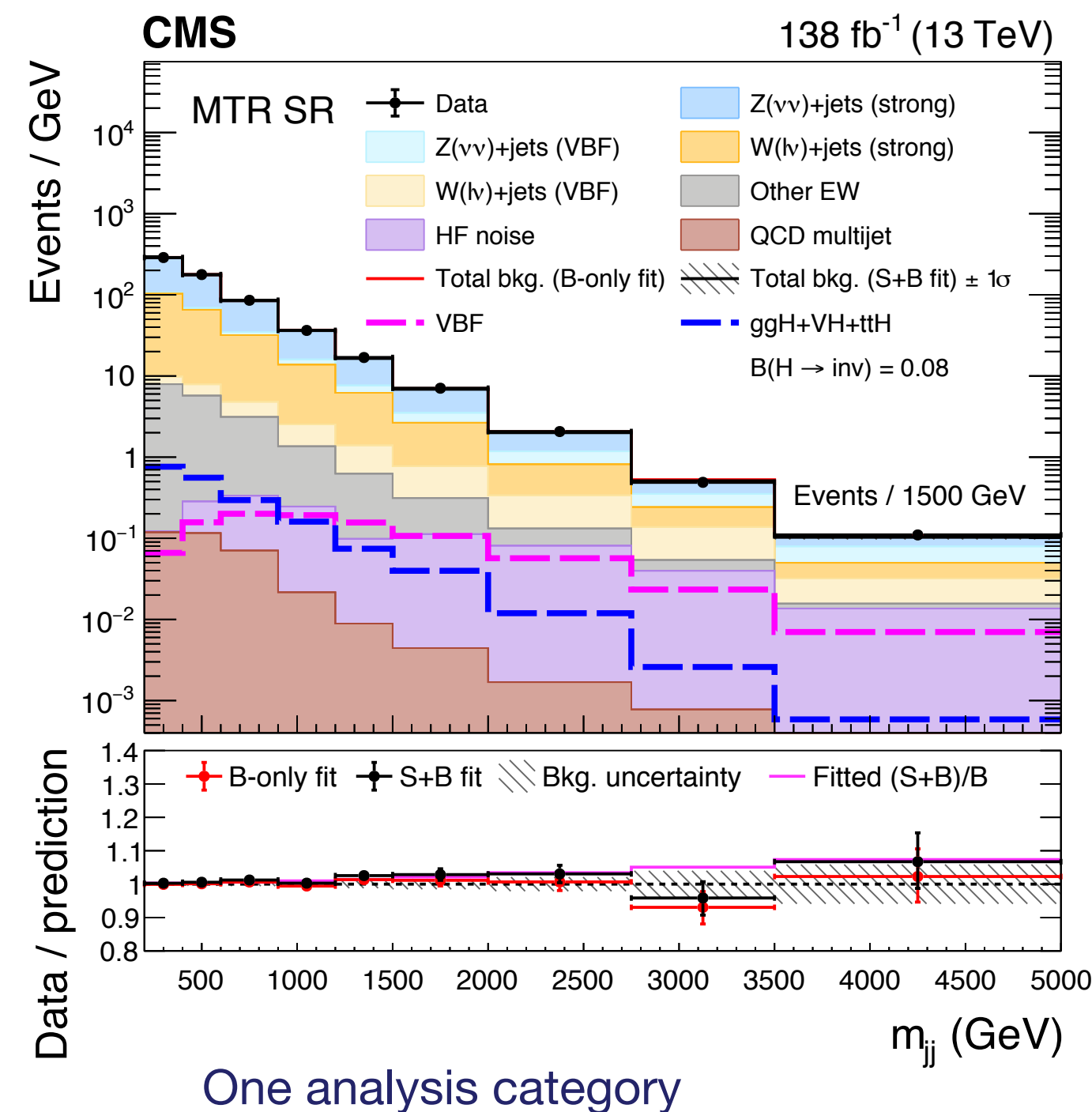
2 jets with large angular separation $\Delta\eta_{jj}$ and large invariant mass m_{jj}

Veto on other objects (leptons/photons)

High missing transverse momentum (trigger constraint) \rightarrow reject QCD

Low $|\Delta\phi_{jj}| \rightarrow$ reject QCD

\Rightarrow Main remaining backgrounds: $Z(\nu\nu) +$ jets and $W(l\nu) +$ jets (strong and VBF productions)



Results

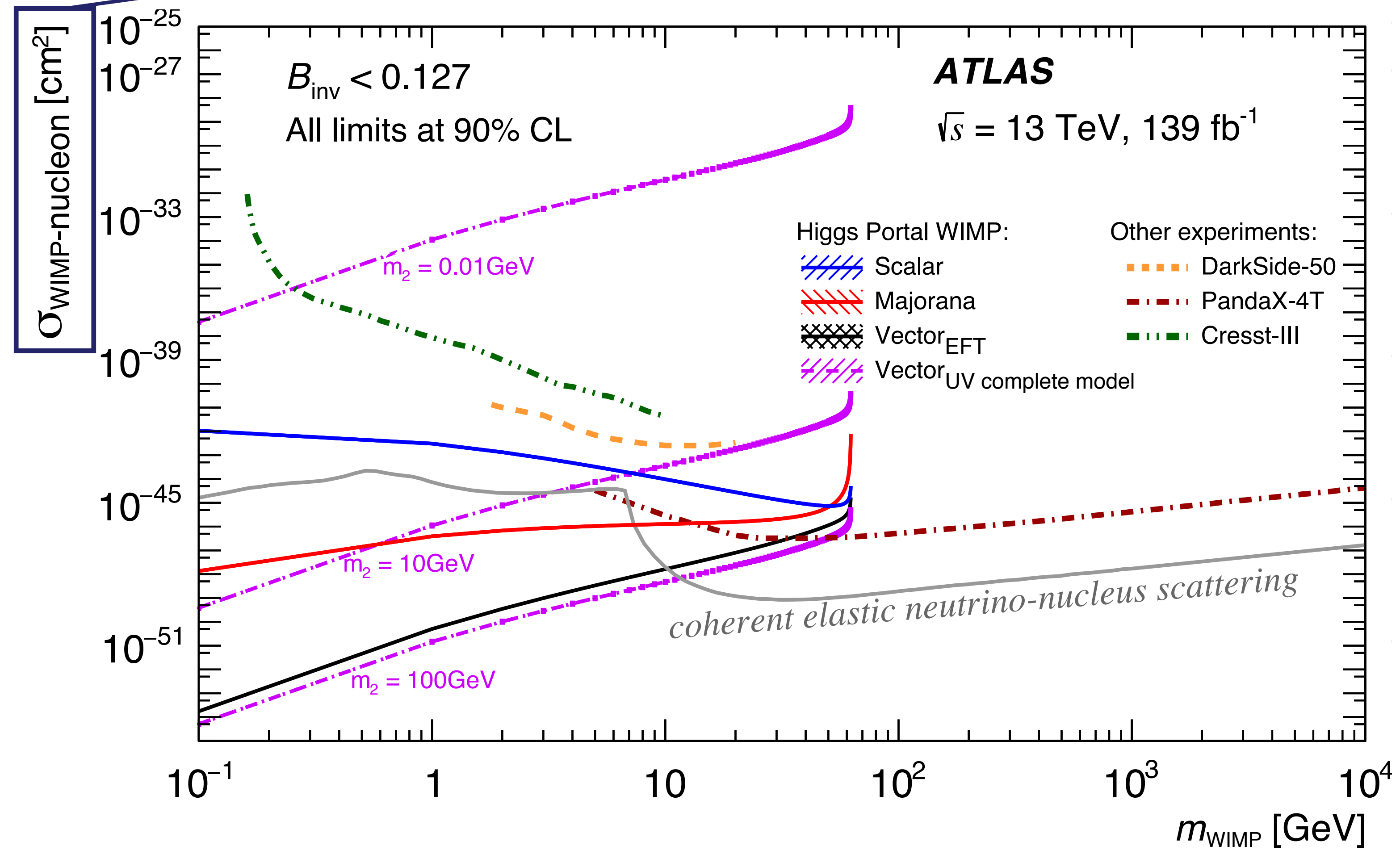
obs(ext) UL on \mathcal{B}_{inv}

0.145(0.103)

0.18(0.10)

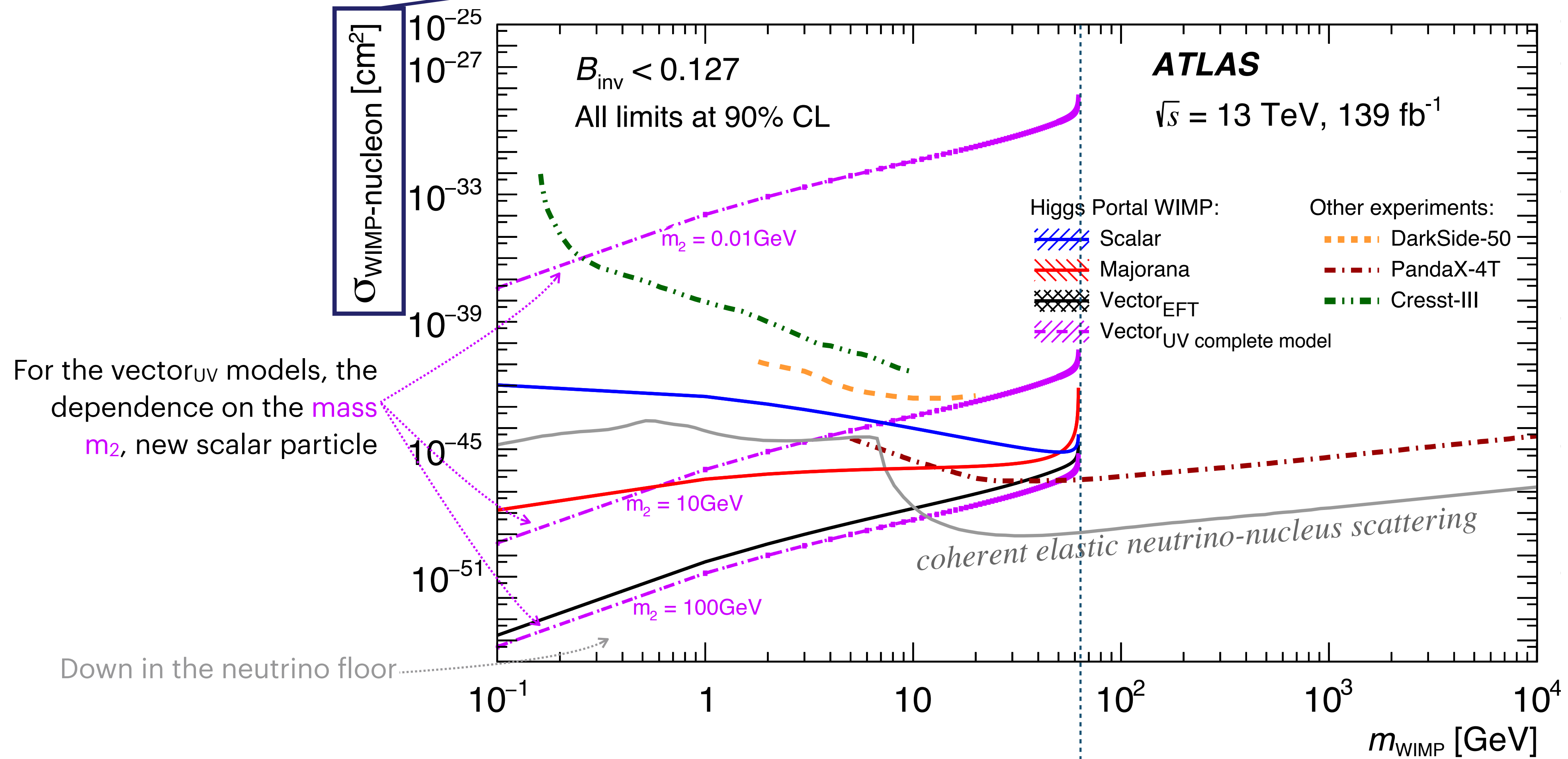
Interpretations

Upper limits on the spin-independent **WIMP-nucleon cross section**



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Upper limits on the spin-independent **WIMP-nucleon cross section**



Kinematical threshold
 $m_{\text{WIMP}} = m_{\text{DM}} = m_{H/2} = 62.5 \text{ GeV}$

Outperforms direct searches experiments for low m_{DM}

Lepton-Flavour Violating decays

$H \rightarrow e\mu/\mu\tau/e\tau$ are forbidden in the SM but takes place through the LFV Yukawa couplings $Y_{ij} \neq (m_i/v)\delta_{ij}$ arising in two Higgs doublet models, extra dimensions, models with flavor symmetries, models of compositeness, ...

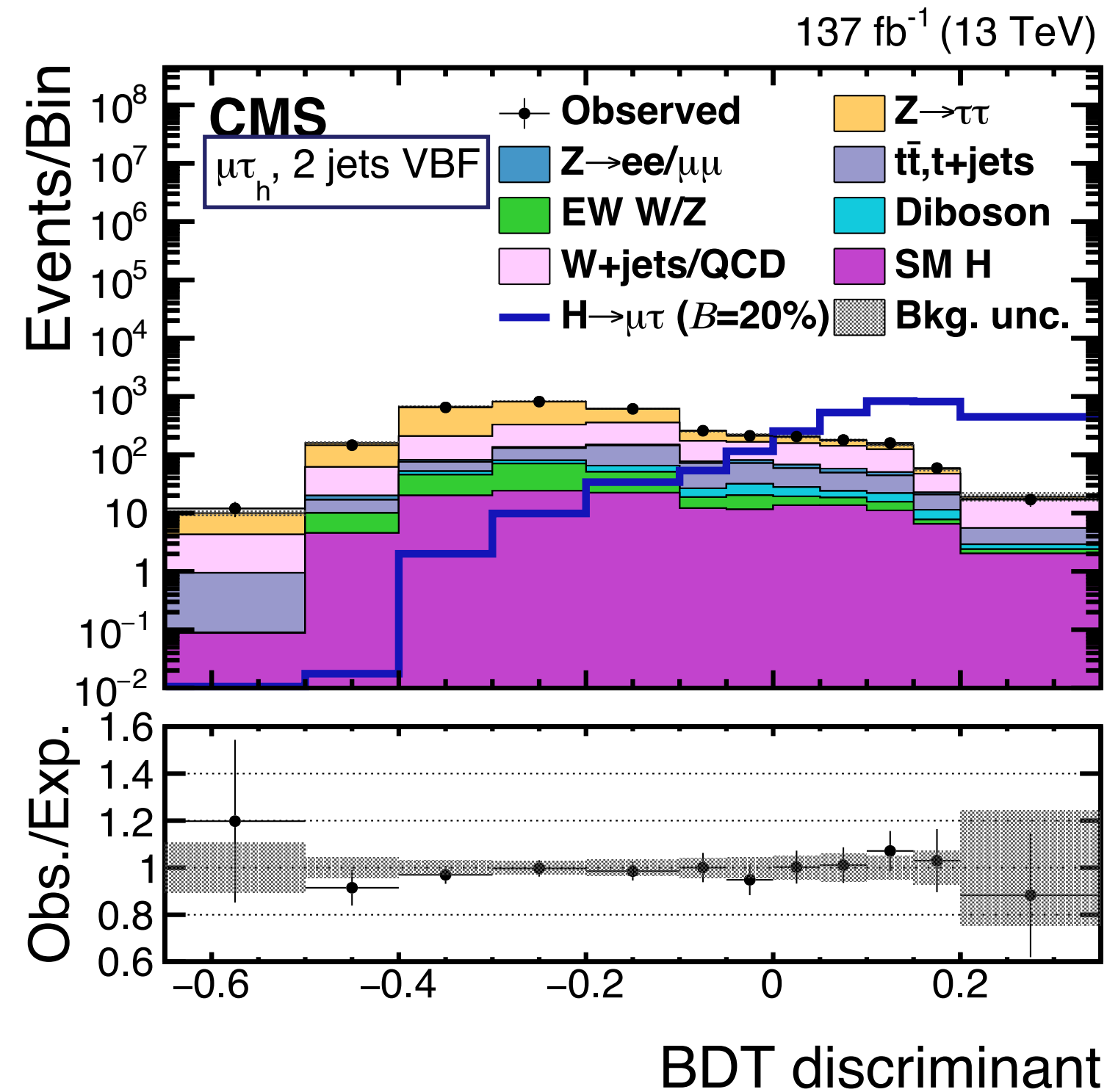
$$\begin{pmatrix} Y_{ee} & Y_{e\mu} & Y_{e\tau} \\ & Y_{\mu\mu} & Y_{\mu\tau} \\ & & Y_{\tau\tau} \end{pmatrix}$$

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Focus on $Y_{e\tau}$ and $Y_{\mu\tau}$ ($Y_{e\mu}$ strongly constrained by $\mu \rightarrow e\gamma$)



Channels: $e\tau_h, e\tau_\mu, \mu\tau_h, \mu\tau_e$

Jet categories: 0j, 1j, 2j (ggH), VBF

BDTs to discriminate signal

Joint fit to BDT outputs

Most sensitive category

Lepton-Flavour Violating decays



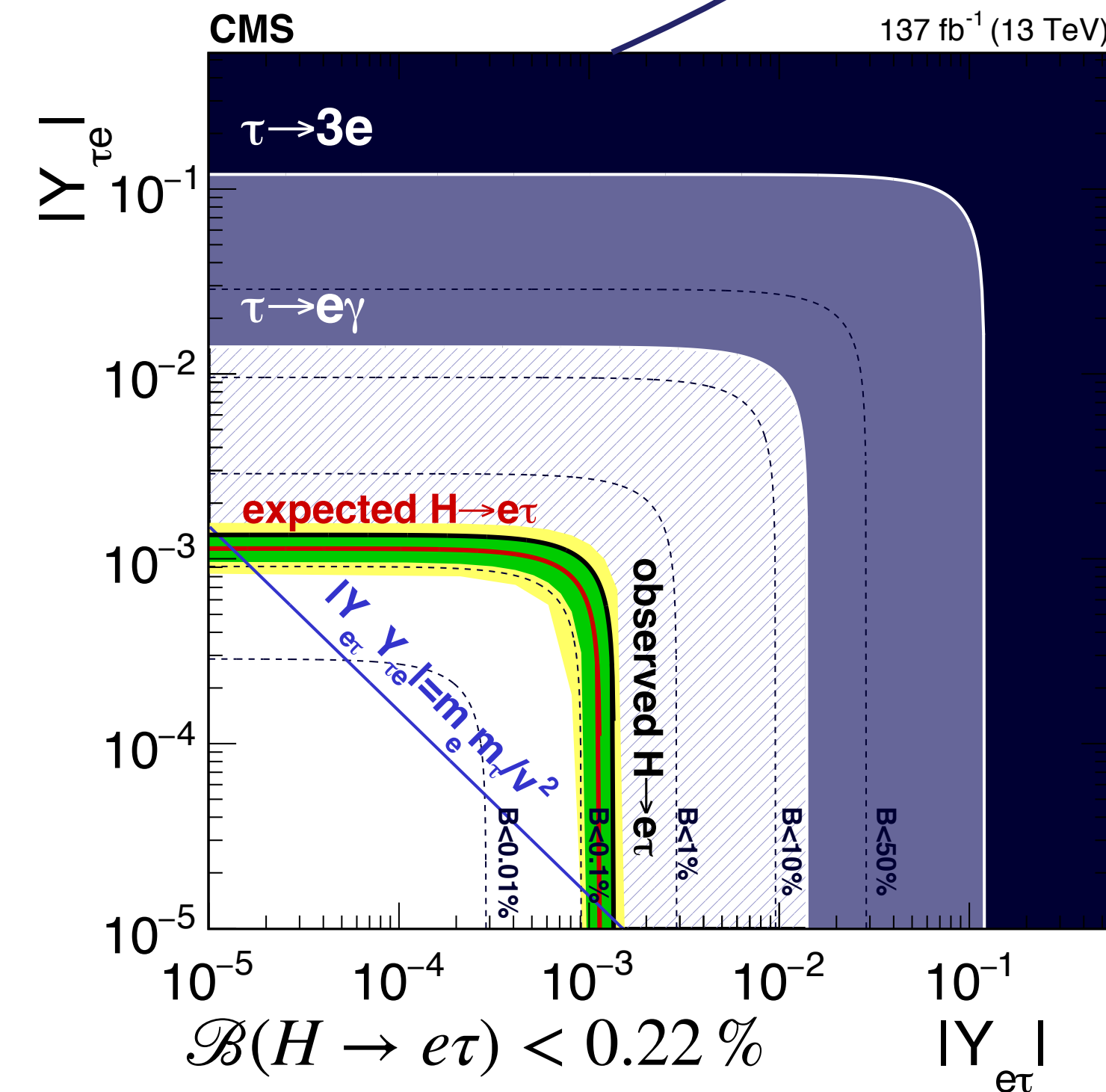
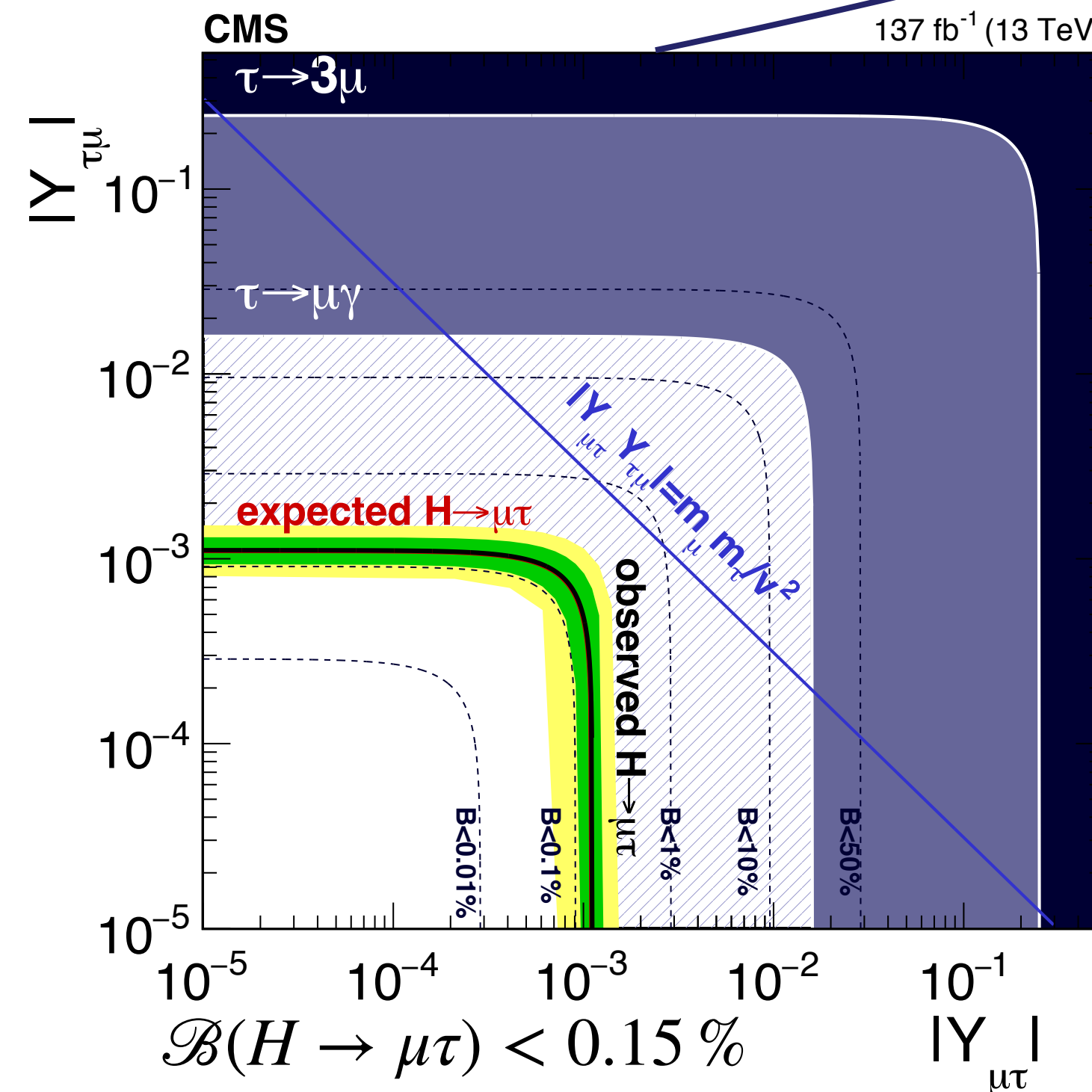
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The upper limits on $\mathcal{B}(H \rightarrow e\tau)$ and $\mathcal{B}(H \rightarrow \mu\tau)$ are used to put constraints on $Y_{e\tau}$ and $Y_{\mu\tau}$

Better than constraints from other experiments and for $Y_{\mu\tau}$ within the naturalness limit $|Y_{\mu\tau}Y_{\tau\mu}| < \frac{m_\mu m_\tau}{v^2}$



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The Higgs boson **remains** compatible with SM predictions.

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The coming decades are crucial to understand it and make use of it in exploring nature.