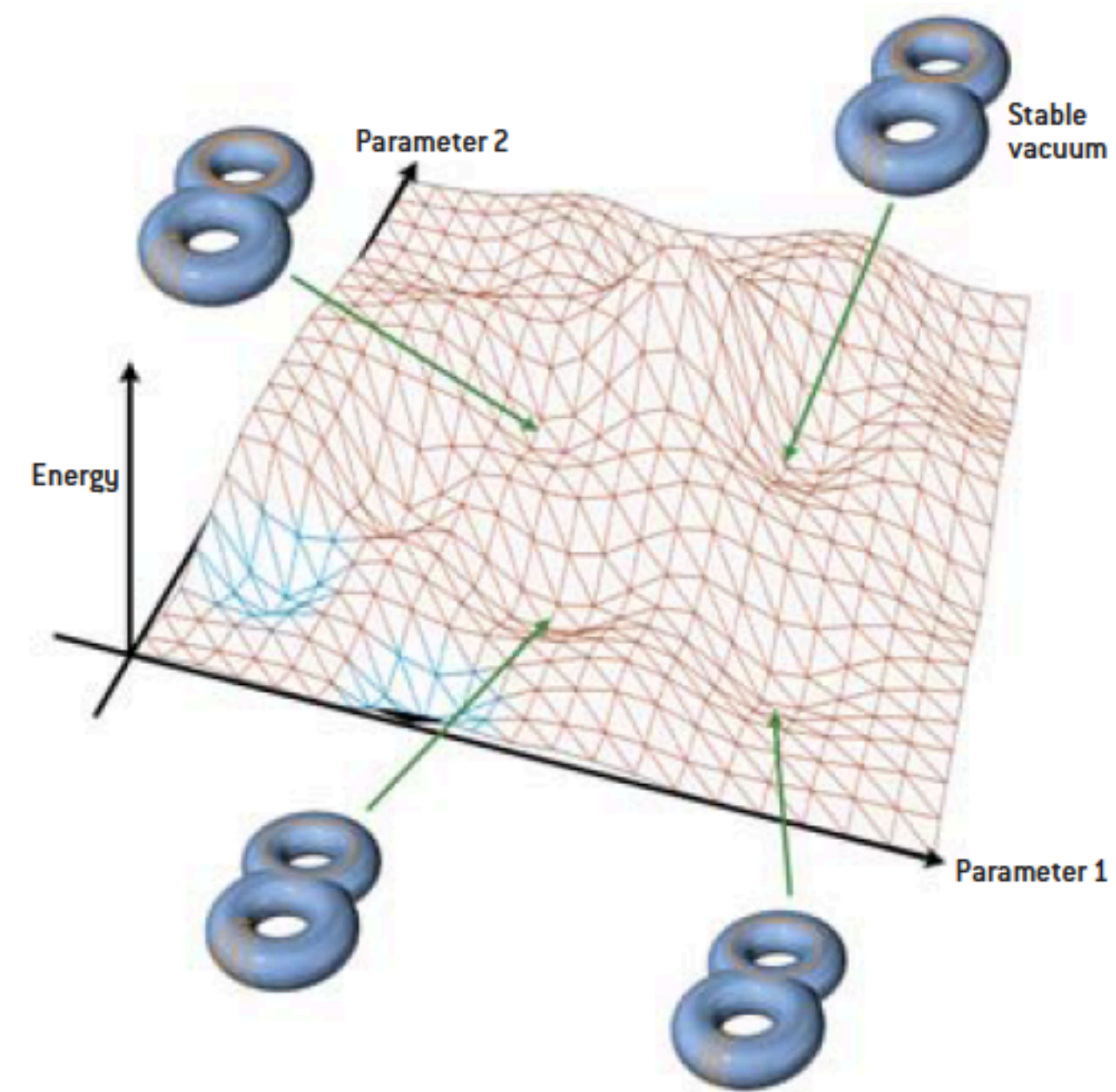


Sparticle and Higgs boson masses from the landscape: dynamical vs. spontaneous SUSY breaking

H. Baer, V. Barger, S. Salam and H. Serce
arXiv:2103.12123

Pheno 2021 meeting, May 25, 2021

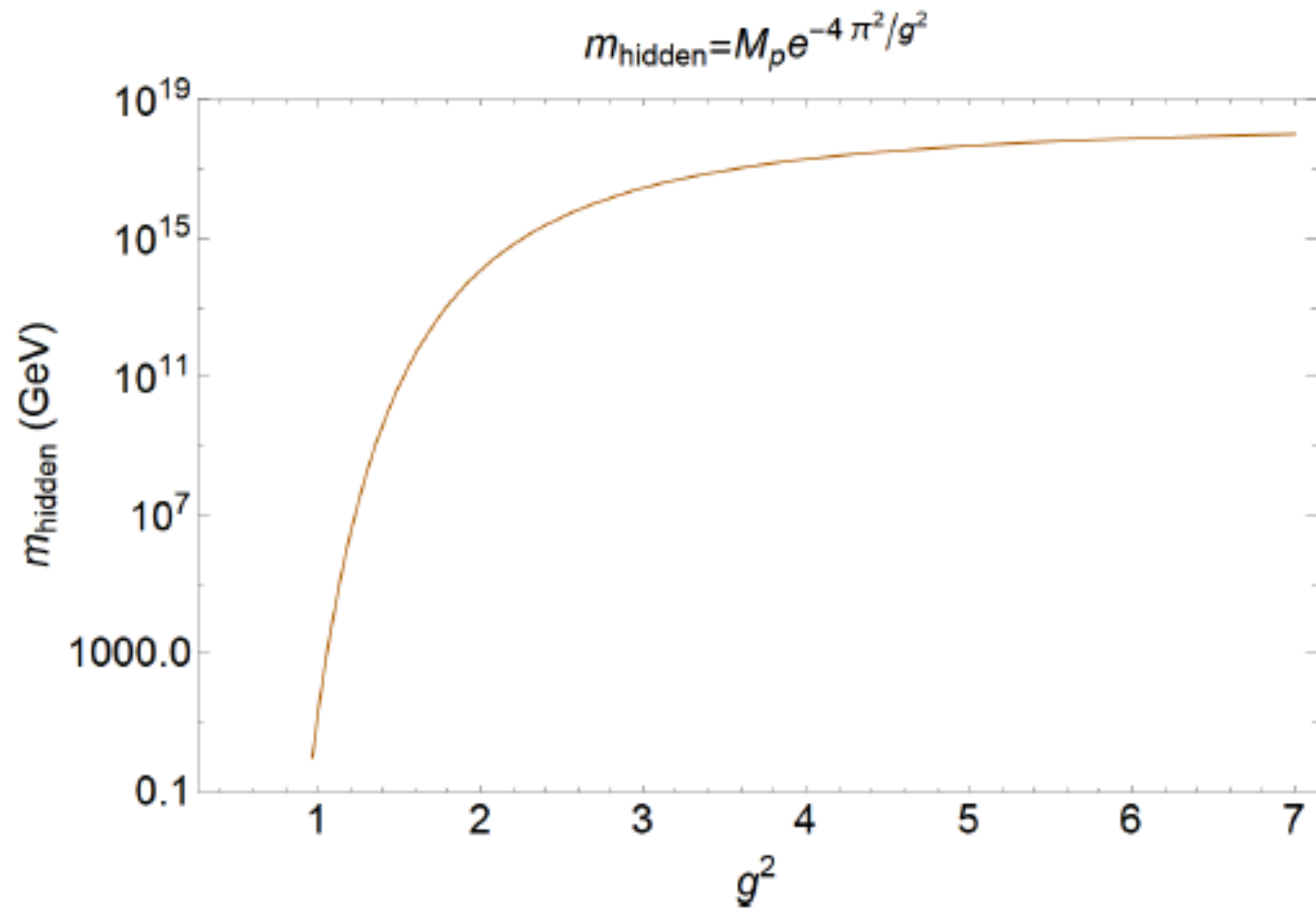
(dedicated to memory
of Joe Polchinski)



Bousso & Polchinski

- In string theory, only one mass scale: $m_P \sim 2.4E18 \text{ GeV}$
- Then how do widely disparate mass scales arise, e.g. $CC \sim 10^{-120} m_P^2$ or the weak scale $m(\text{weak}) \sim 100 \text{ GeV}$ or QCD scale $\sim 1 \text{ GeV}$?
- CC: Weinberg's anthropic solution: in eternally inflating multiverse, if CC much bigger than measured value, universe would expand so quickly that structure (galaxies) wouldn't form: these 'pocket universe' would not be suitable for evolution of life (structure principle)
- QCD is different: QCD scale arises non-perturbatively (dynamically) from dimensional transmutation: QCD becomes confining at $m(\text{proton}) \sim m_P \exp(-8\pi^2/g^2)$ when $g^2 \sim 1.8$

- Weak scale: in SM, quadratic divergences $\rightarrow m(\text{weak}) \rightarrow m_P$: but can (implausibly) tune μ^2 such that $m(\text{weak}) \sim 100 \text{ GeV}$
- Weak scale SUSY stabilizes weak scale, but does not explain magnitude
- e.g. for SUGRA breaking, $W(\text{Polonyi}) = m^2(h + \beta)$ for lone superfield h : gives right answer if $m \sim 10^{11} \text{ GeV}$ and $\beta \sim m_P$ (must implausibly put in by hand)
- But maybe instead SUSY breaking dynamical (like QCD):
 $m^2 \sim m_P^2 \exp(-8\pi^2/g_{\text{hidden}}^2)$
- e.g. hidden sector gauge group $SU(N)$ becomes confining (gaugino condensation) at : $\Lambda(\text{GC}) \sim 10^{13} \text{ GeV} \Rightarrow m = \sqrt{\Lambda^3/m_P} \sim 10^{11} \text{ GeV}$



When g becomes confining $\sim 1-2$, then SUSY breaking scale uniformly distributed across the decades of possibilities:

then in landscape context, $f_{\text{SUSY}}^{\text{DSB}} \sim 1/m_{\text{soft}}$

see e.g Dine, Gorbатов, Thomas (2008)

In landscape context (used to solve CC problem),
 expect $\sim 10^{500}$ string vacua (Denef & Douglas)
 vacua distributed as:

$$dN_{vac}[m_{hidden}^2, m_{weak}, \Lambda_{cc}] = f_{SUSY} \cdot f_{EWSB} \cdot f_{cc} \cdot dm_{hidden}^2$$

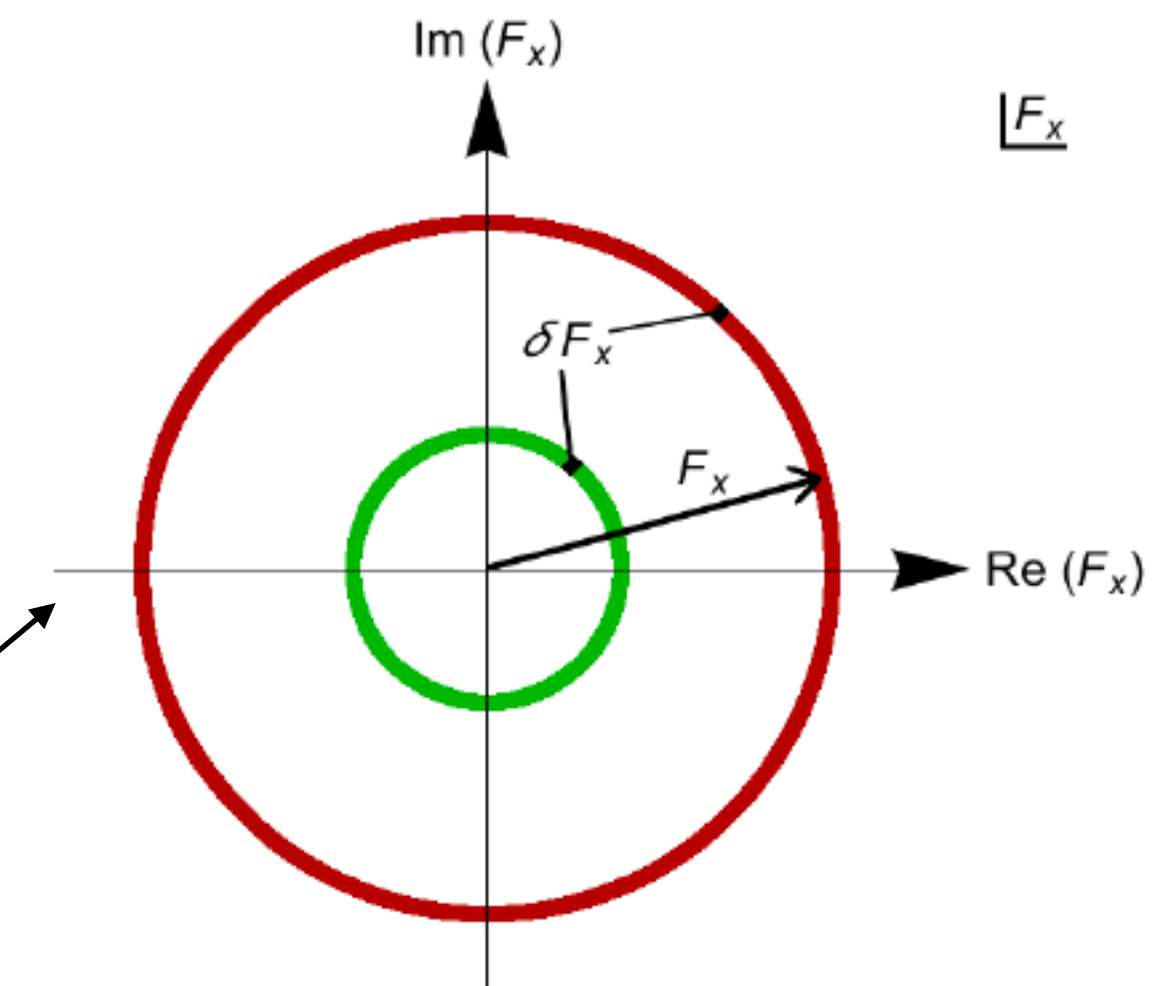
For spontaneous SUSY breaking (mass scale included, perturbative)

$$f_{SUSY}^{SSB} \sim m_{soft}^n$$

where $n = 2n_F + n_D - 1$ and n_F are the number of hidden sector SUSY breaking F -fields and n_D is the number of hidden sector D -breaking fields contributing to the overall SUSY breaking

Thus, in landscape, DSB favors low soft terms
 while SSB favors large soft terms!

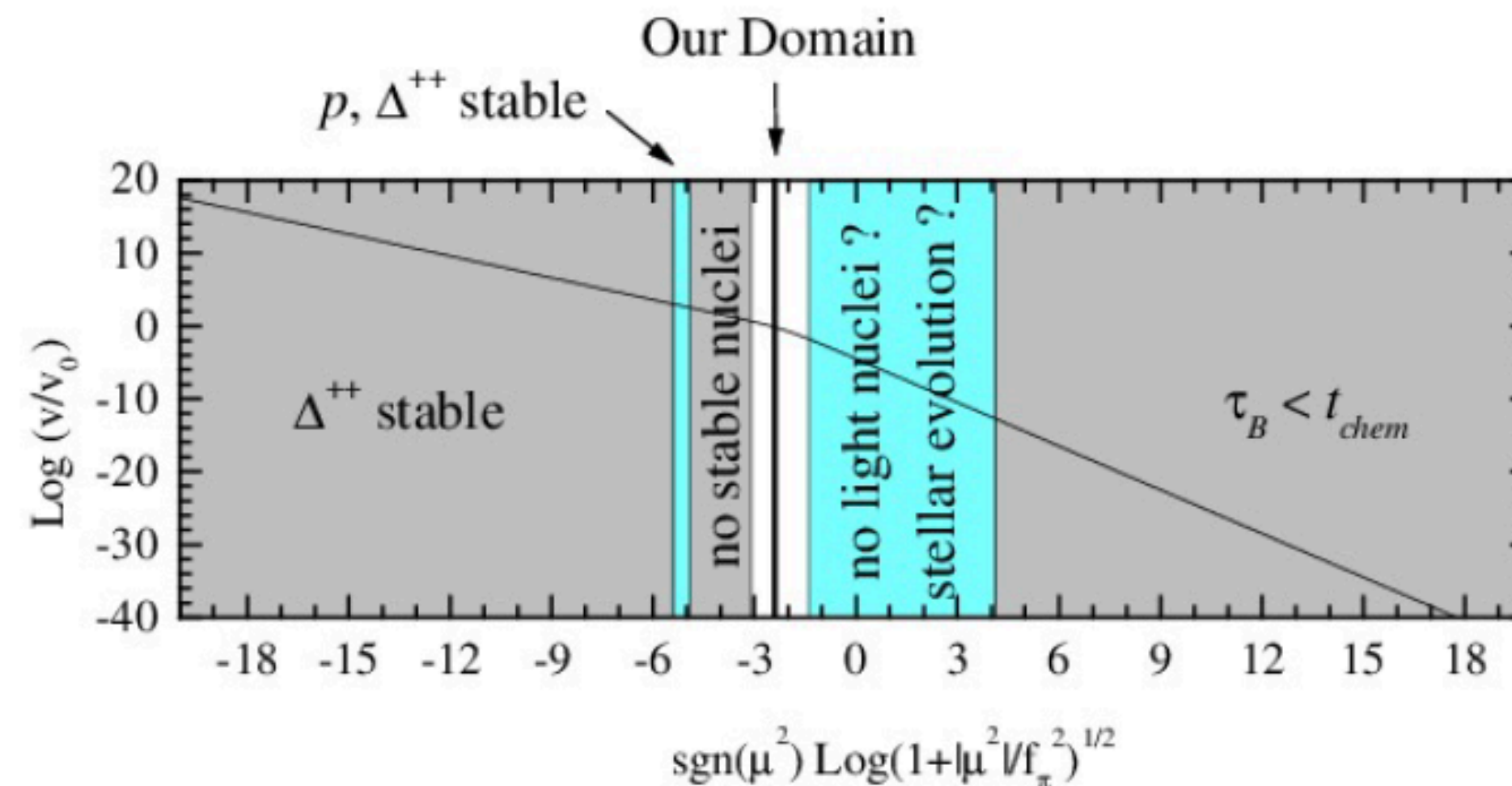
single F term distributed uniformly as complex number



[footnote: f_{cc} doesn't contribute to SUSY breaking scale determination (Douglas)]

fEWSB ?

From Agrawal, Barr, Donoghue, Seckel (ABDS, 1998):
 if pocket universe value of weak scale too displaced
 from measured value in our universe (OU) [factor 2-5],
 then complex nuclei and hence atoms will not form:
pocket universe will not sustain life as we know it!



atomic principle

also: veto CCB
 and noEWSB minima

can calculate $m(\text{weak})$
 in MSSM

$$(m_Z^{\text{PU}})^2/2 = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \simeq -m_{H_u}^2 - \Sigma_u^u - \mu^2$$

for $m(\text{weak})^{\text{PU}} < 4 * m(\text{weak})^{\text{OU}}$: then $f_{\text{EWSB}} = \Theta(30 - \Delta_{\text{EW}})$

Assume fertile patch of landscape
where MSSM is LE-EFT

Can scan over parameters in models which allow $DEW < 30$: e.g.

$m_0(1, 2), m_0(3), m_{1/2}, A_0, \tan \beta, \mu, m_A$ (NUHM3)

- $m_0(1, 2) : 0.1 - 60$ TeV,
- $m_0(3) : 0.1 - 20$ TeV,
- $m_{1/2} : 0.5 - 10$ TeV,
- $A_0 : -50 - 0$ TeV,
- $m_A : 0.3 - 10$ TeV,

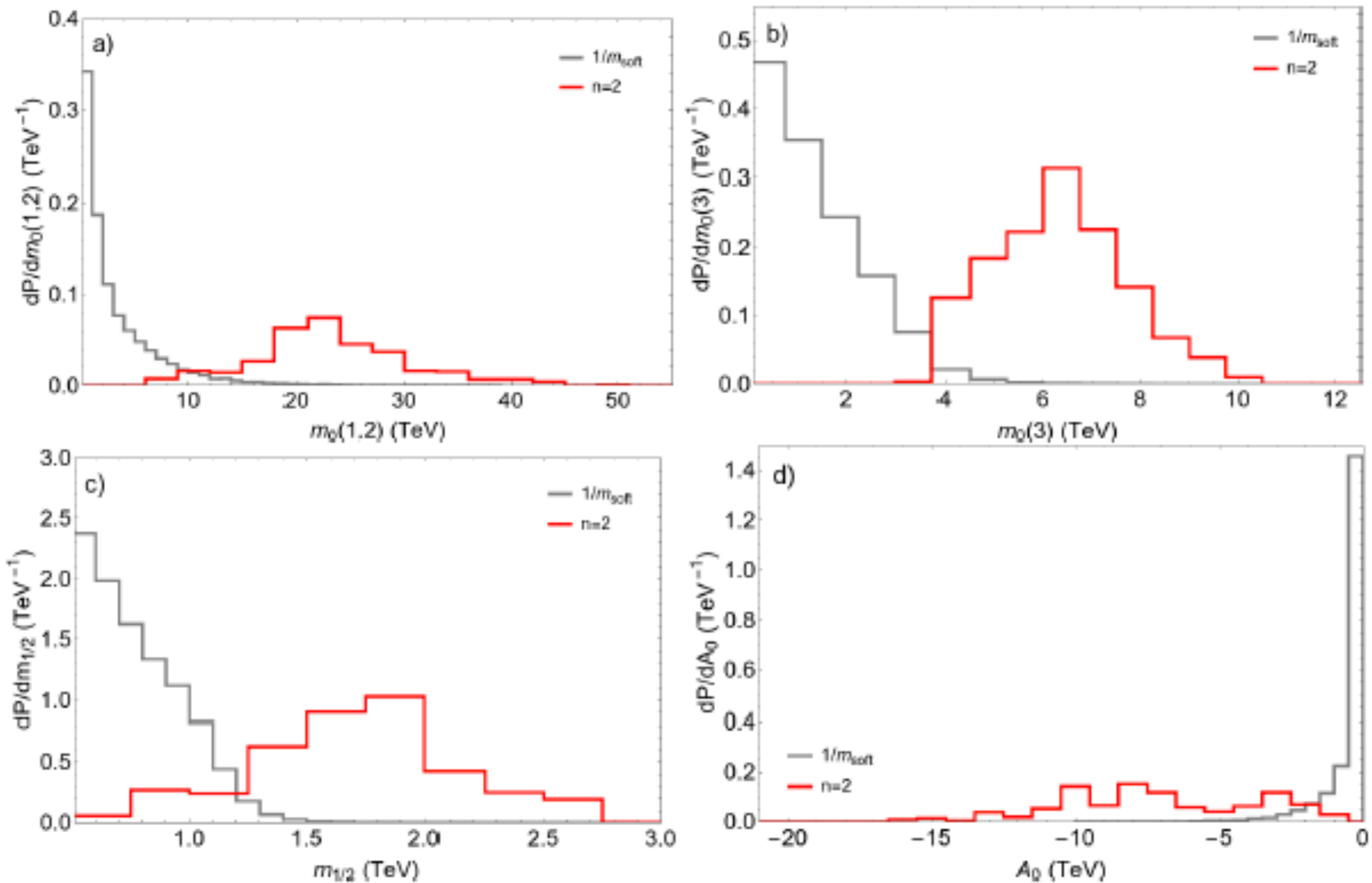
with $\mu = 150$ GeV while $\tan \beta : 3 - 60$

upper limits set beyond
anthropic upper limits:

lower limits set

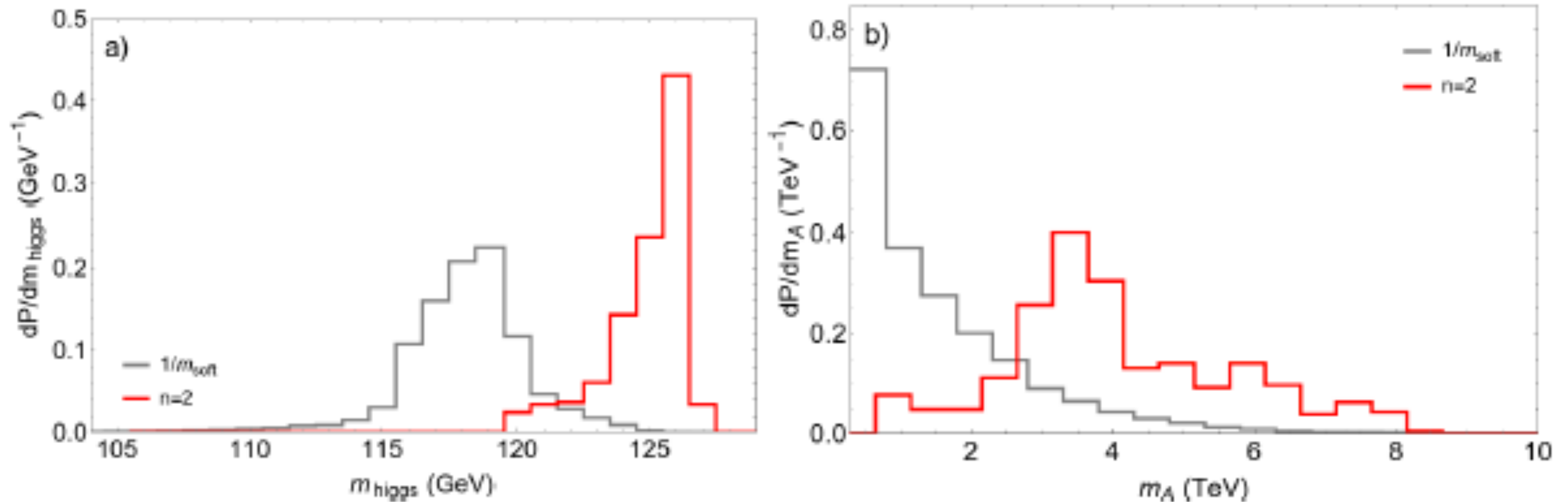
to compare against

previous scans, but can be lower yet



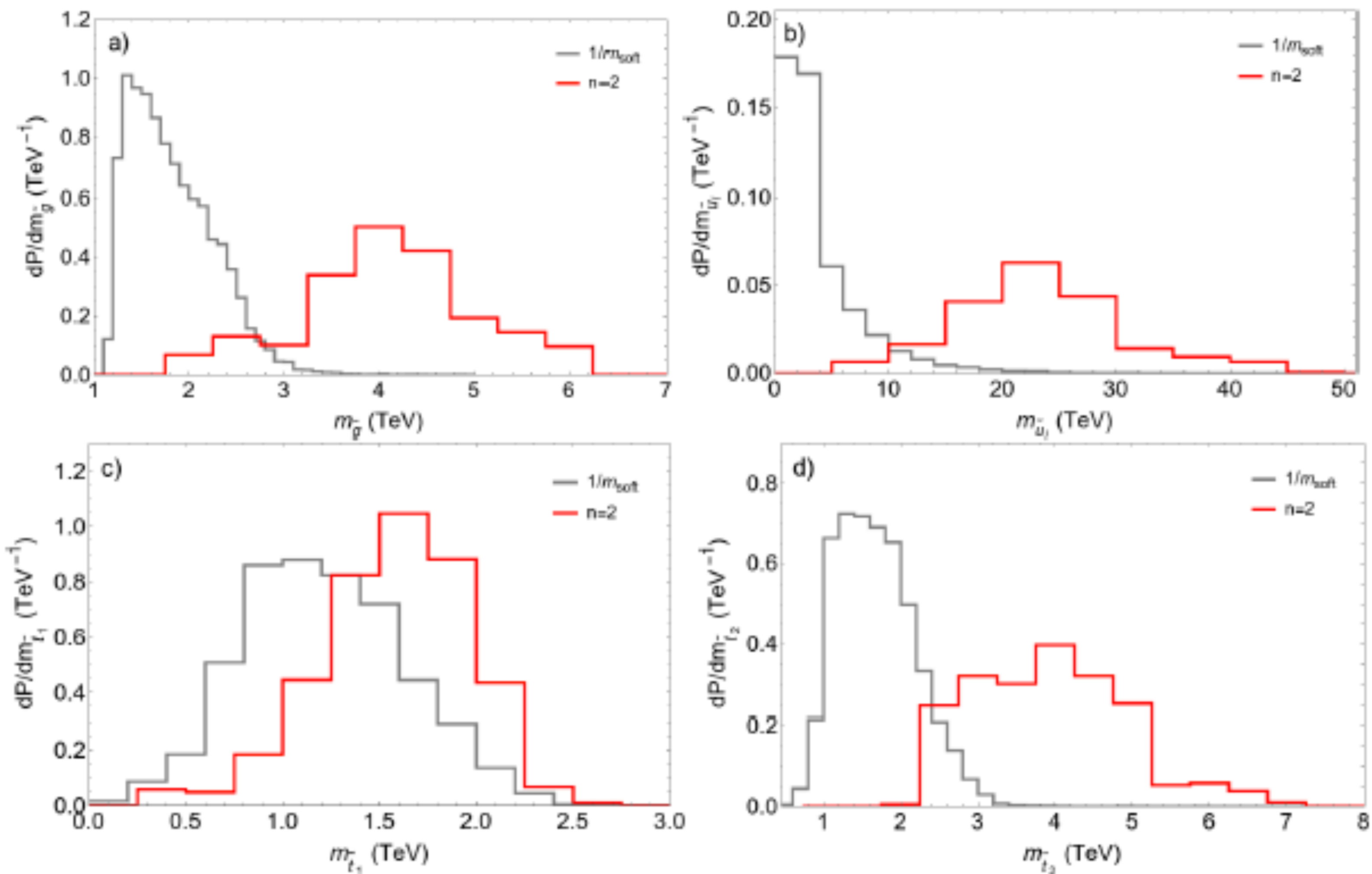
$n=2$ from
KKLT stabilization;
see Broeckel et al.

as expected, DSB (gray) prefers small soft terms
while SSB (red) prefers large



Higgs masses: DSB \Rightarrow $m(h)\ll 125$ GeV while SSB prefers $m(h)\sim 125$ GeV

DSB \Rightarrow highly mixed Higgs while SSB prefers decoupled Higgs



DSB => sparticles masses below LHC limits;
 SSB prefers => sparticles masses above LHC limits!

[smaller lower scan limits make matters worse for DSB]

Conclusions:

- DSB beautiful theory:
- DSB might explain exponential suppression of weak scale
- DSB predicts $m(h) \ll 125$ GeV and excluded sparticles
- SSB in landscape: $m(h) \sim 125$ GeV and sparticles $>$ LHC limits
- then, exponential suppression of weak scale arises as does the CC: weak scale as big as possible such that atomic principle (existence of atoms) still holds