

A handle on anomalous top-Higgs couplings in top quark pair production through EW loops

based on [arXiv:2104.04277]

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Motivation

The LHC is a top quark factory:

- Top quark produced $\approx 240 \times 10^6$ times in Run-2 ($\approx 160 \text{ fb}^{-1}$)
(1000× more than at Tevatron)

Main top quark dynamics predictable at percent level:

- NNLO+NNLL QCD predictions for single and pair production
[Czakon, Fiedler, Heymes, Mitov], [Brucherseifer, Caola, Melnikov], [Berger, Gao, Yuan, Zhu],
[Czakon, Mitov, Sterman], [Beneke, Czakon, Falgari, Mitov, Schwinn], [Beneke, Falgari, Klein, Schwinn], [Kidonakis],
[Ferroglia, Pecjak, Yang], [Ferroglia, Marzani, Pecjak, Yang], [Czakon, Ferroglia, Heymes, Mitov, Pecjak, Scott]
- Decay known to NNLO QCD (\leftrightarrow prod. via NWA and beyond)
[Gao, Li, Zhu], [Brucherseifer, Caola, Melnikov], [Bevilacqua, Czakon, v. Hameren, Papadopoulos, Worek],
[Denner, Dittmaier, Kallweit, Pozzorini], [Heinrich, Maier, Nisius, Schlenk, Winter],
[Frederix, Frixione, Papanastasiou, Prestel, Torrielli], [Denner, Pellen]
- EW corrections known for production and width
[Beenakker, Denner, Hollik, Mertig, Sack, Wackerlo], [Kühn, Scharf, Uwer], [Bernreuther, Fücker, Si],
[Moretti, Nolten, Ross], [Groote, Körner, Mauser], [Basso, Dittmaier, Huss, Toggero]

Top quark sector ideal lab for New Physics searches!

Dimension-six operators in SMEFT

Deviations from the SM parametrised within EFT

| Wilson coefficient | Dimension-six operator |
|---|---|
| $\mathcal{L}^{\text{EFT}} = \mathcal{L}_{\text{SM}}^{(4)} + \sum_X \frac{C^X}{\Lambda^2} Q_X^{(6)}$ | $+ \mathcal{O}\left(\frac{Q^{(8)}}{\Lambda^4}\right)$ Scale of New Physics |

Focus on electroweak top quark interactions

E.g. (SMEFT in Warsaw basis [Dedes,Materkowska,Paraskevas,Rosiek,Suxho'17]):

$$Ztt : Q_{33}^{\varphi q1} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{q}'_3 L \gamma^\mu q'_3 L) \leftrightarrow C_{33}^{\varphi q1},$$

$$Ztt, Wtb, \chi tt, \phi tb : Q_{33}^{\varphi q3} = (\varphi^\dagger i \tau^I \overleftrightarrow{D}_\mu \varphi) (\bar{q}'_3 L \tau^I \gamma^\mu q'_3 L) \leftrightarrow C_{33}^{\varphi q3},$$

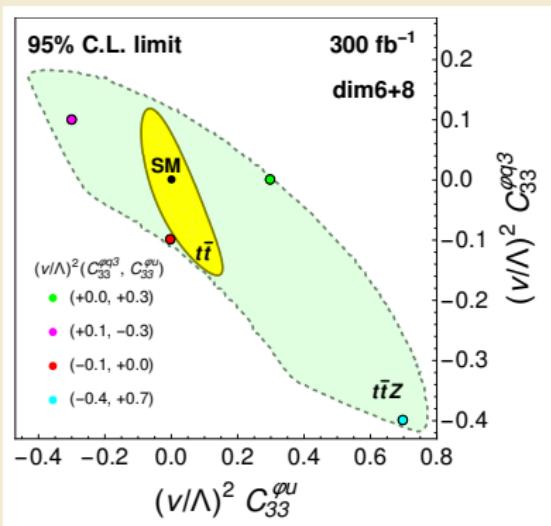
$$Ztt, \chi tt : Q_{33}^{\varphi u} = (\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{t}'_R \gamma^\mu t'_R) \leftrightarrow C_{33}^{\varphi u},$$

$$Htt : Q_{33}^{u\varphi} = (\varphi^\dagger \varphi) (\bar{q}'_3 L t'_R \tilde{\varphi}) \leftrightarrow C_{33}^{u\varphi}$$

On-shell vs. loop sensitivity [Schulze, TM'20]

Previous study: Focus on anomalous Ztt : $C_{33}^{\varphi u}$ & $C_{33}^{\varphi q3}$

- Sensitivity in $t\bar{t}Z$:
low rate (coupling, threshold, BR), known to NLO QCD ($\approx \pm 15\%$ scale unc.)
- $t\bar{t}$ sensitive via EW loops:
 $\approx 100\times$ more events, QCD background known to NNLO+NNLL ($\approx \pm 5\%$ scale unc.), Sudakov enhancement in diff. distrib., requires EW NLO calculation within SMEFT



Promising way to constrain New Physics!

Next step: Focus on anomalous Htt : $C_{33}^{u\varphi}$

Including $Q_{33}^{u\varphi}$ in the Lagrangian yields modified Feynman rules:

$$\Gamma_{Htt}^{\text{EFT}} = \frac{-im_t}{v} + \frac{i v^2}{\sqrt{2}\Lambda^2} \left(P_L C_{33}^{u\varphi *} + P_R C_{33}^{u\varphi} \right) \text{ /w } P_{R/L} = \frac{1}{2}(1 \pm \gamma_5)$$

Definitions: $\kappa = 1 - \frac{v}{\sqrt{2}m_t} \frac{v^2}{\Lambda^2} \text{Re}(C_{33}^{u\varphi})$, $\tilde{\kappa} = -\frac{v}{\sqrt{2}m_t} \frac{v^2}{\Lambda^2} \text{Im}(C_{33}^{u\varphi})$

$$\Rightarrow \Gamma_{Htt}^{\kappa, \tilde{\kappa}} = \frac{-im_t}{v} (\kappa + i\gamma_5 \tilde{\kappa})$$

CP even

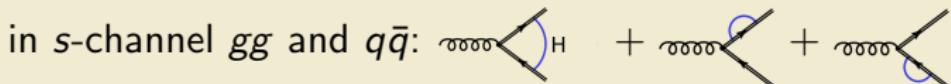
CP odd

- 0^- states inherent to, e.g., SUSY or two-Higgs-doublet models (CP violation required to explain baryon asymm. [Sakharov'91])
- Arbitrary CP-mixing possible via $\kappa, \tilde{\kappa}$
- SM recovered for $\kappa = 1$ and $\tilde{\kappa} = 0$

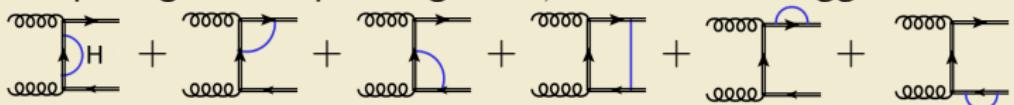
Loop sensitivity on $\kappa, \tilde{\kappa}$ in $t\bar{t}$ production

Calculational details — modified amplitudes

- Final state corrections depending on $\kappa, \tilde{\kappa}$



- Loop diagrams depending on $\kappa, \tilde{\kappa}$ in t -channel gg :



Due to coupling structure in $\Gamma_{Htt}^{\kappa, \tilde{\kappa}}$, interference of diagrams with Higgs loops and Born level either $\propto (\kappa^2 + \tilde{\kappa}^2)$ or $\propto (\kappa^2 - \tilde{\kappa}^2)$.

\Rightarrow No sensitivity on signs of $\kappa, \tilde{\kappa}$

Loop diagrams involving the Higgs boson are IR finite but contain UV poles \Rightarrow Renormalization necessary within SMEFT

One-loop amplitude is UV finite after renormalization! ✓

Loop sensitivity on $\kappa, \tilde{\kappa}$ in $t\bar{t}$ production

Calculational details — numerical evaluation

We build upon the existing SM implementation of $t\bar{t}$ -production at NLO EW in MCFM and allow for arbitrary top Yukawa couplings

- Analytic results of our calculation available as external add-on:
https://github.com/TOPAZdevelop/MCFM-8.3_EWSMEFT_ADDON
- MCFM results reproduced when setting $\kappa = 1, \tilde{\kappa} = 0$
- MCFM established framework able to produce multi-dim. kinematic distributions together with EW correction factor

$$\delta_{\text{wk}}^{\kappa, \tilde{\kappa}} = \frac{d\sigma_{\kappa, \tilde{\kappa}}^{\text{NLO EW}} - d\sigma^{\text{LO}}}{d\sigma^{\text{LO}}}$$

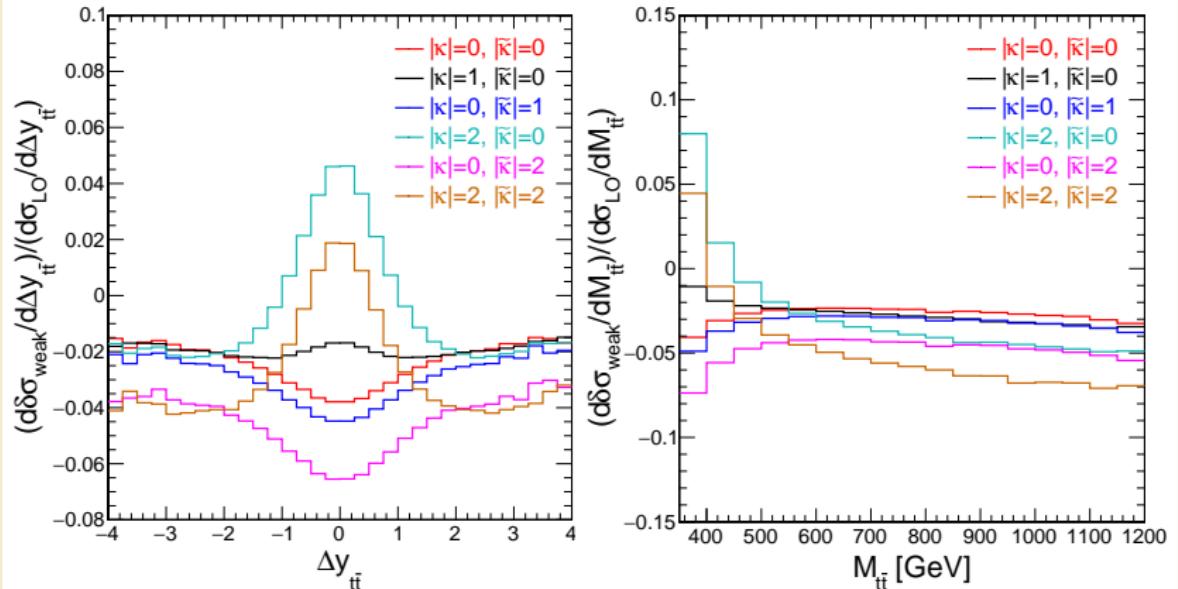
- $\delta_{\text{wk}}^{\kappa, \tilde{\kappa}}$ can be used to reweight distributions generated by LO+PS+DetectorSim. generators

Note: Higgs loops with scalable κ (CP even only) available in HATHORv2.1 and resp. study done by CMS [Phys.Rev.D100 7,(2019)072007]

Loop sensitivity on $\kappa, \tilde{\kappa}$ in $t\bar{t}$ production

Dependence of the shapes on $\kappa, \tilde{\kappa}$

$\Delta y_{t\bar{t}}$ and $M_{t\bar{t}}$ sensitive to CP structure of top Yukawa coupling



Event simulation and analyses

Loop sensitivity: $t\bar{t}$ (semileptonic channel)

- Top quark pair events simulated by MadGraph5_v2.6.4+Pythia8.1+Delphes3 (CMS setting)
- Simulated events are normalized to $\sigma_{t\bar{t}}^{\text{NNLO QCD}} = 832^{+40}_{-46} \text{ pb}$
- 2-dim distribution of $M_{t\bar{t}}$ and $\Delta y_{t\bar{t}}$ filled with events
- 2-dim distrib. can be reweighted with $\delta_{\text{wk}}^{\kappa, \tilde{\kappa}}$ from modified MCFM
- Main background: single top, $V+\text{jets}$ and QCD multijets

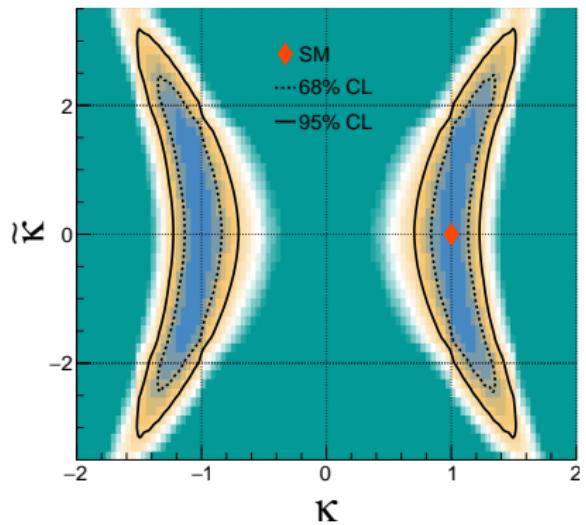
On-shell benchmark: e.g. tHW (full hadr. channel and $H \rightarrow \gamma\gamma$)

- Events with top and antitop quarks simulated by JHU Generator+Pythia8.1+Delphes3 (CMS setting)
- Differential cross sections for different hypotheses implemented in JHU Generator
- MELA (matrix element likelihood approach): 3 discriminants
- Main background: $t\bar{t}H$

Results

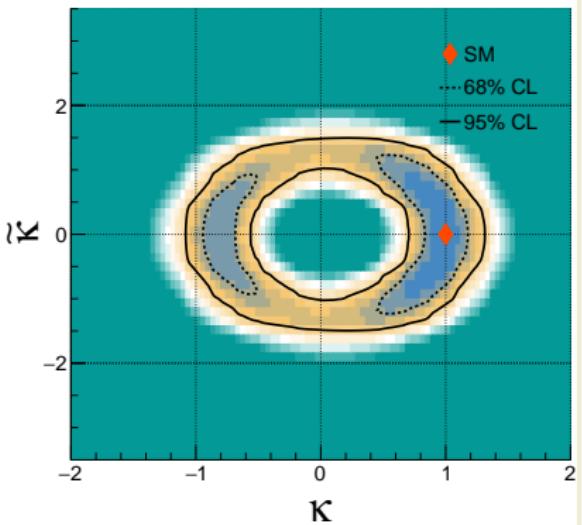
Two dimensional likelihood scans of κ and $\tilde{\kappa}$

Loop sensitivity



$t\bar{t}$ at 300 fb^{-1}

On-shell benchmark



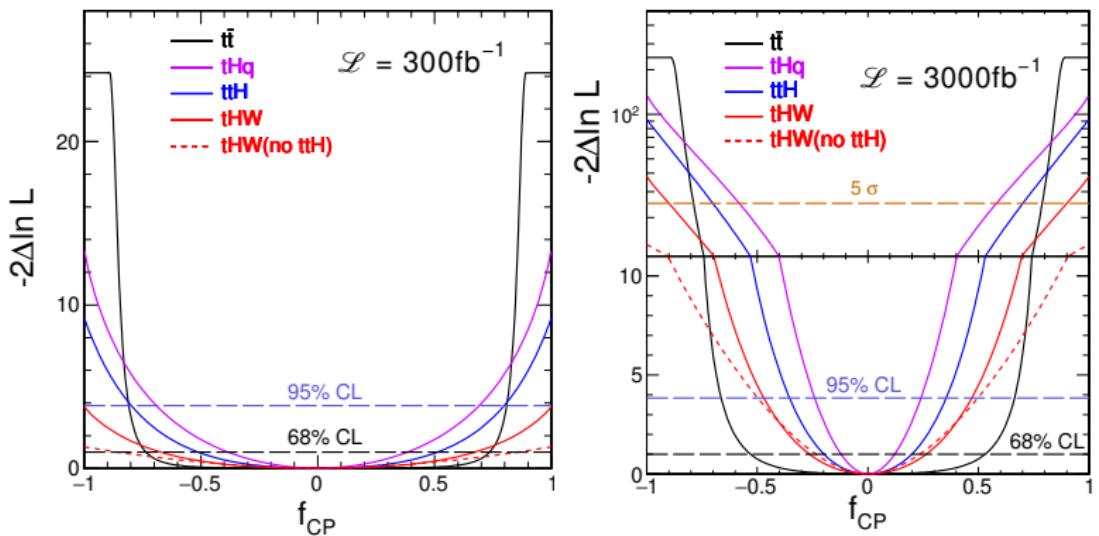
tHW at 300 fb^{-1}

Results

Fractional contribution of the CP-odd component

$$f_{CP} = \frac{|\tilde{\kappa}|^2}{|\kappa|^2 + |\tilde{\kappa}|^2} \text{sign} \left(\frac{\tilde{\kappa}}{\kappa} \right)$$

Likelihood scans of f_{CP}



$t\bar{t}H$ and tHq results taken from [Gritsan, Röntsch, Schulze, Xiao '16]

Conclusions

- **Probing CP structure of top Yukawa coupling through EW loops**
 - Calculation of EW correction to $pp \rightarrow t\bar{t}$
 - NP included via SMEFT parametrisation of Htt coupling
- **Comparison to direct on-shell probes, e.g., $pp \rightarrow tHW$** with full-fledged simulation chain
 - $t\bar{t}$ expected to exclude $|f_{CP}| > 0.81$ for 300 fb^{-1} and $|f_{CP}| > 0.67$ for 3000 fb^{-1} at 95% CL
 - tHW can exclude pure pseudo-scalar model at 2σ for 300 fb^{-1} and $|f_{CP}| > 0.48$ for 3000 fb^{-1} at 95% CL
 - tHq gives most stringent 95% CL exclusion: $|f_{CP}| > 0.68$ for 300 fb^{-1} and $|f_{CP}| > 0.22$ for 3000 fb^{-1}
 - $t\bar{t}$ best probe to exclude purely CP-odd top Yukawa coupling
- $t\bar{t}$ and tHW with anomalous top Yukawa couplings available as add-on to MCFM or via JHU Generator respectively