

The Higgs of Baryon Number and Dark Matter

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with Pavel Fileviez Perez, Elliot Golias and Clara Murgui

[\[arXiv: 2003.09426\]](#) [JHEP 07 \(2020\) 087](#)

[\[arXiv: 2012.06599\]](#) [JHEP 02 \(2021\) 163](#)



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Aim of the talk

Discuss the Higgs and dark matter phenomenology in minimal gauge extensions of the SM that predict dark matter

1. $U(1)_B$

[Fileviez Perez, Golias, Li, Murgui, ADP 1904.01017]

[Fileviez Perez, Murgui, ADP 2003.09426]

[Fileviez Perez, ADP 2008.09116]

Gauging baryon number

- Baryon number is an accidental global symmetry in the SM
- Only broken by non-perturbative effects - $SU(2)$ instantons
- Anomalous in the Standard Model

Gauging baryon number

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$$\underbrace{U(1)_B} \quad \langle S_B \rangle \neq 0$$

Local gauge symmetry
gauge boson: Z_B

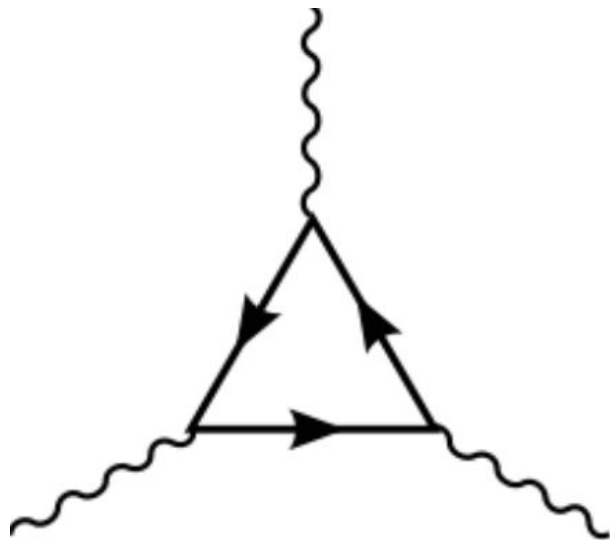
- Spontaneous breaking of baryon number
- Consistent completion of simplified models of dark matter

[Pais 1973]

[Fileviez Perez, Wise 2011]

Anomaly cancellation

- Baryon number broken in 3 units: $\Delta B = \pm 3$ interactions
 ➔ No proton decay
- Need to add new fermions to cancel anomalies



$$\mathcal{A}_1(\text{SU}(3)^2 \otimes \text{U}(1)_B), \quad \mathcal{A}_2(\text{SU}(2)^2 \otimes \text{U}(1)_B),$$
$$\mathcal{A}_3(\text{U}(1)_Y^2 \otimes \text{U}(1)_B), \quad \mathcal{A}_4(\text{U}(1)_Y \otimes \text{U}(1)_B^2),$$
$$\mathcal{A}_5(\text{U}(1)_B), \quad \mathcal{A}_6(\text{U}(1)_B^3),$$

In the SM the non-zero values are:

$$\mathcal{A}_2 = -\mathcal{A}_3 = 3/2$$

Anomaly-free model

[Fileviez Perez, Ohmer, Patel 1403.8029]

Fields	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$U(1)_B$
$\Psi_L = \begin{pmatrix} \Psi_L^+ \\ \Psi_L^0 \end{pmatrix}$	1	2	$\frac{1}{2}$	$\frac{3}{2}$
$\Psi_R = \begin{pmatrix} \Psi_R^+ \\ \Psi_R^0 \end{pmatrix}$	1	2	$\frac{1}{2}$	$-\frac{3}{2}$
$\Sigma_L = \frac{1}{\sqrt{2}} \begin{pmatrix} \Sigma_L^0 & \sqrt{2}\Sigma_L^+ \\ \sqrt{2}\Sigma_L^- & -\Sigma_L^0 \end{pmatrix}$	1	3	0	$-\frac{3}{2}$
χ_L^0	1	1	0	$-\frac{3}{2}$

- Neutral fermion required for anomaly cancellation
- Automatically stable from remnant $U(1) \rightarrow Z_2$ symmetry



DM Candidate



Simplified Dark Matter

χ : Majorana DM

Z_B : Leptophobic mediator

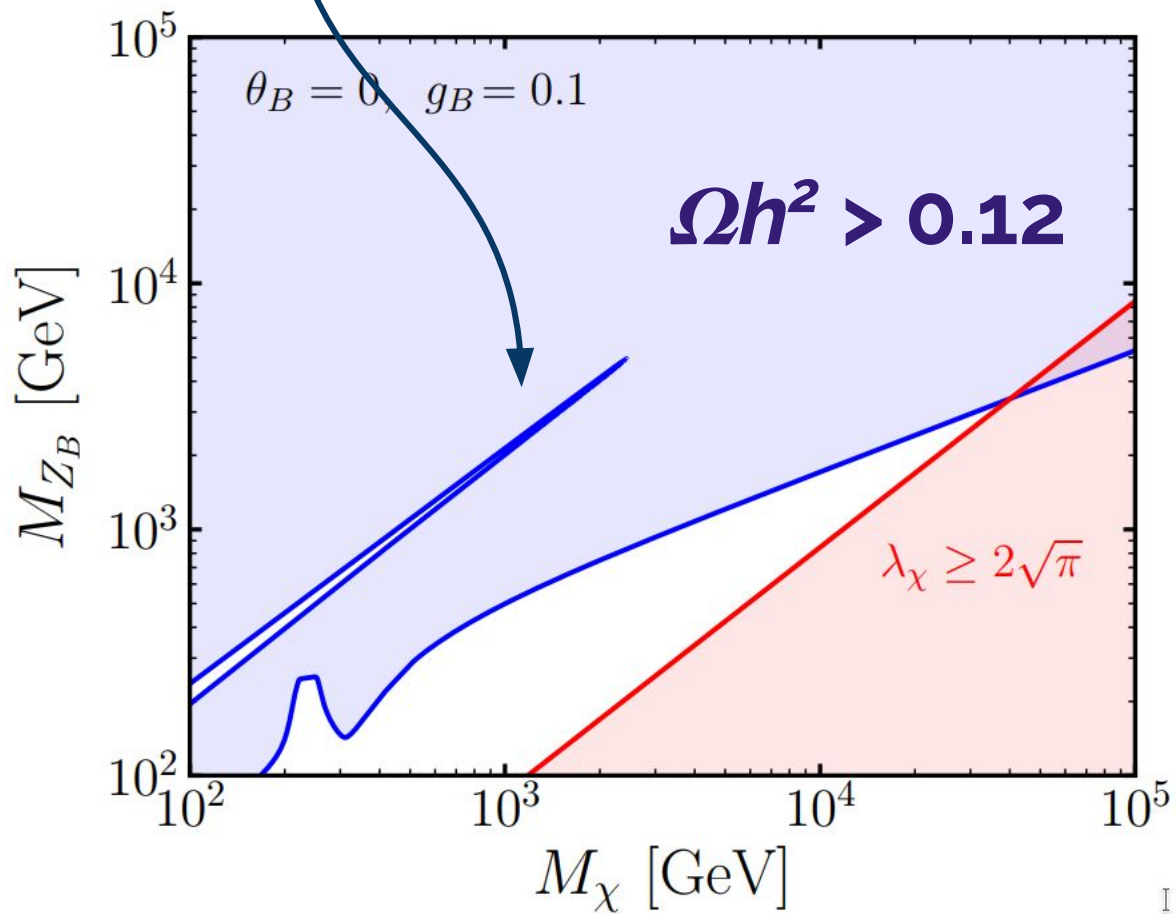
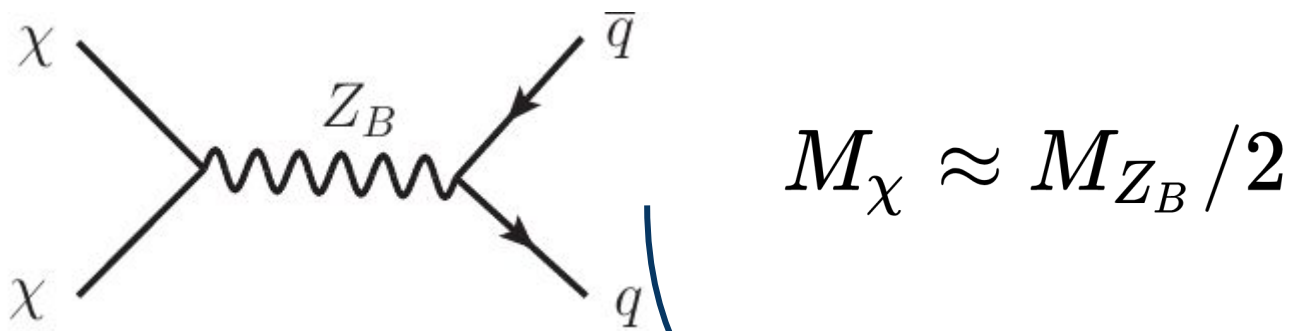
$$\mathcal{L} \supset \underbrace{\frac{3}{4} g_B \bar{\chi} \gamma^\mu \gamma^5 \chi Z_\mu^B}_{\text{Axial}} - \underbrace{\frac{1}{3} g_B \bar{q} \gamma^\mu q Z_\mu^B}_{\text{Vector}} + \frac{M_\chi}{2v_B} \sin \theta_B \bar{\chi} \chi h_1 - \frac{M_\chi}{2v_B} \cos \theta_B \bar{\chi} \chi h_2$$

Axial

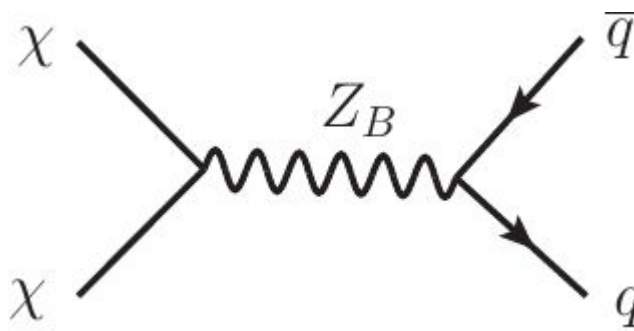
Vector

the free parameters in the model are:

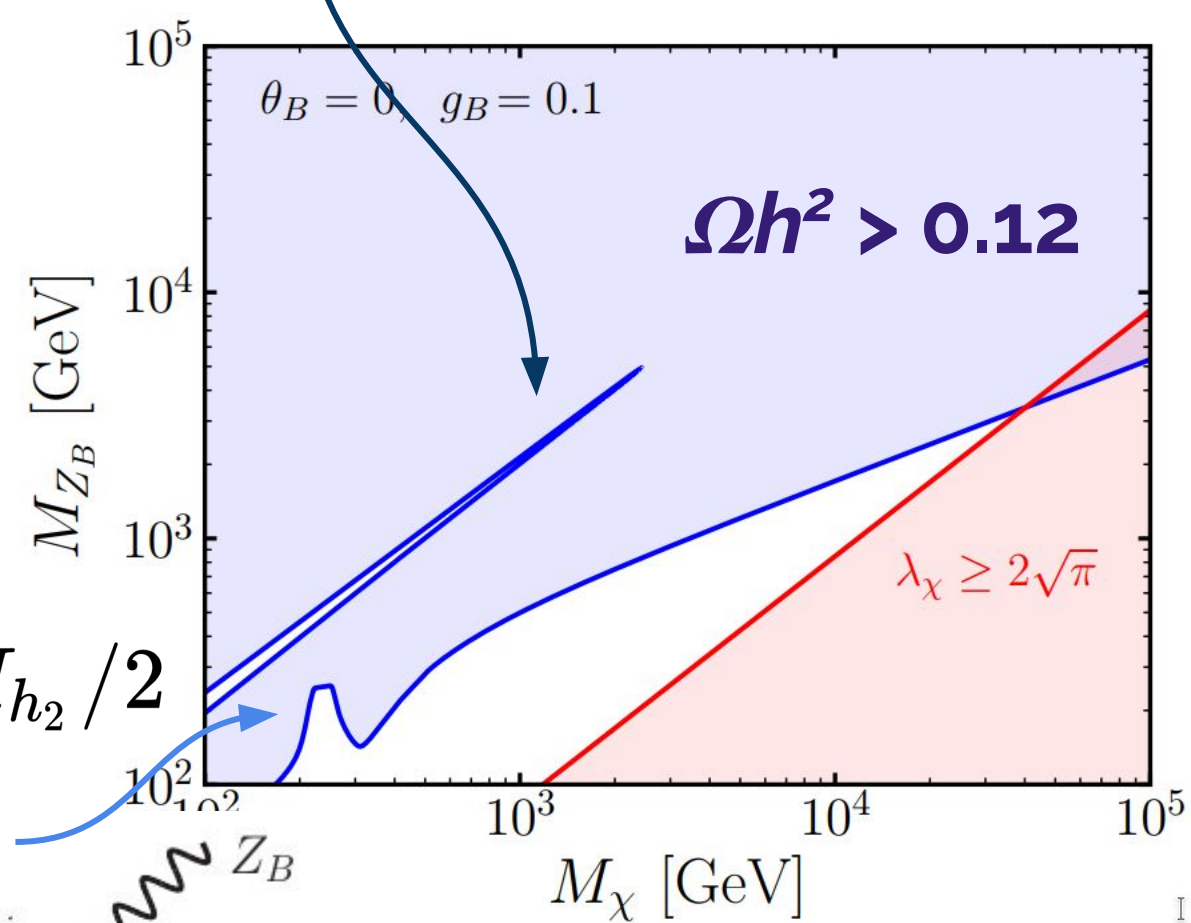
$$M_\chi, M_{Z_B}, M_{h_2}, \theta_B, g_B.$$



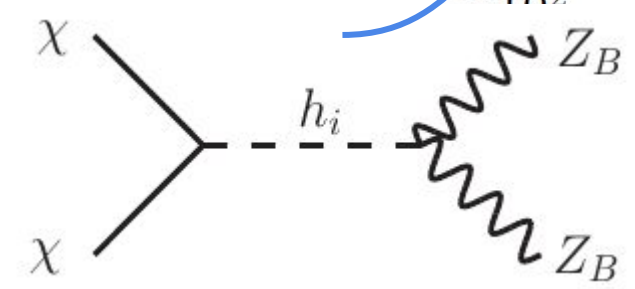
[Fileviez Perez, Golias, Li, Murgui, ADP 1904.01017]

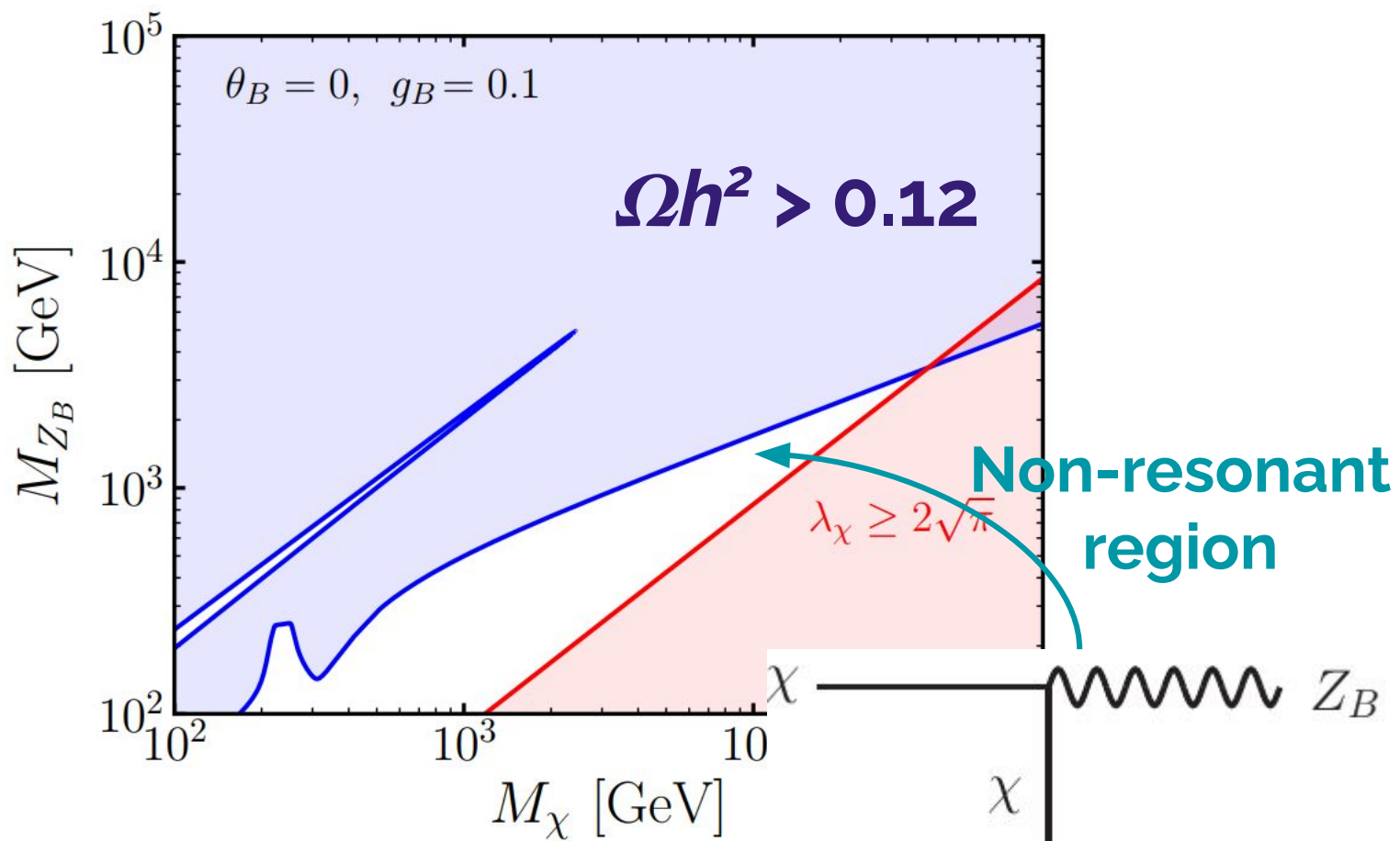


$$M_\chi \approx M_{Z_B} / 2$$

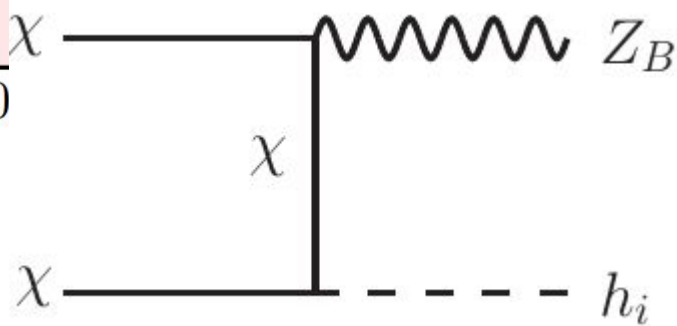


$$M_\chi \approx M_{h_2} / 2$$

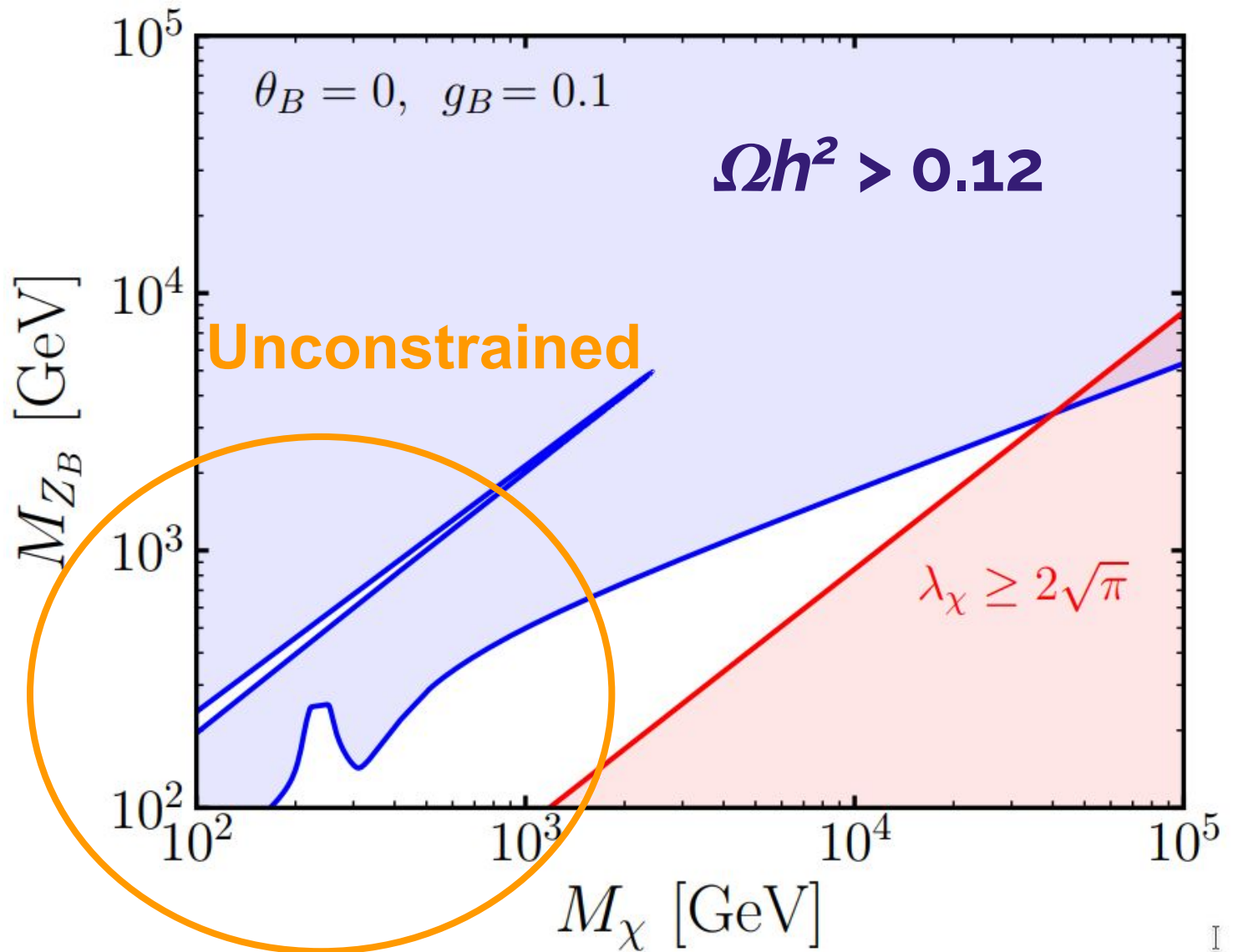




Alexis Plascencia



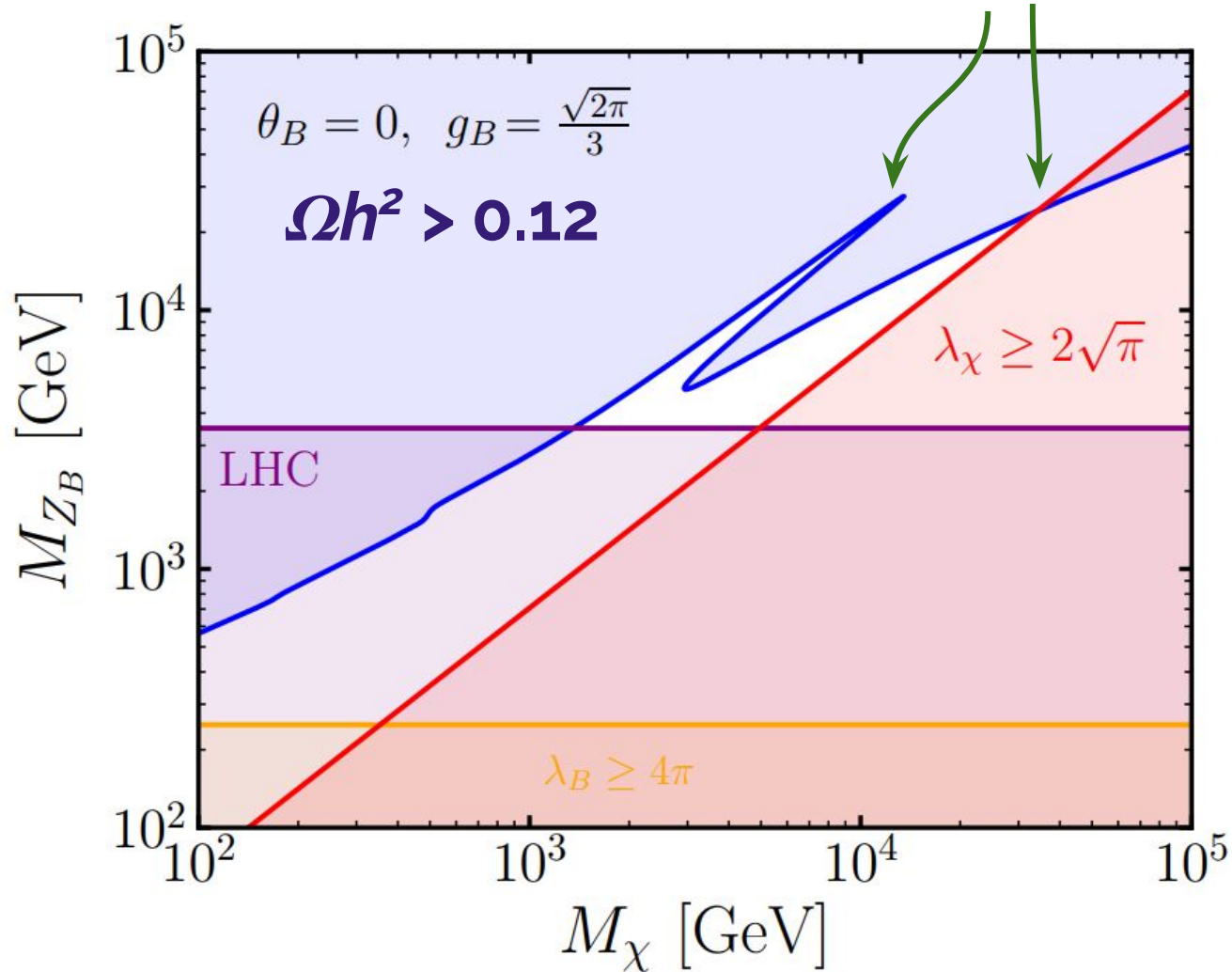
Results



[Fileviez Perez, Golias, Li, Murgui, ADP 1904.01017]

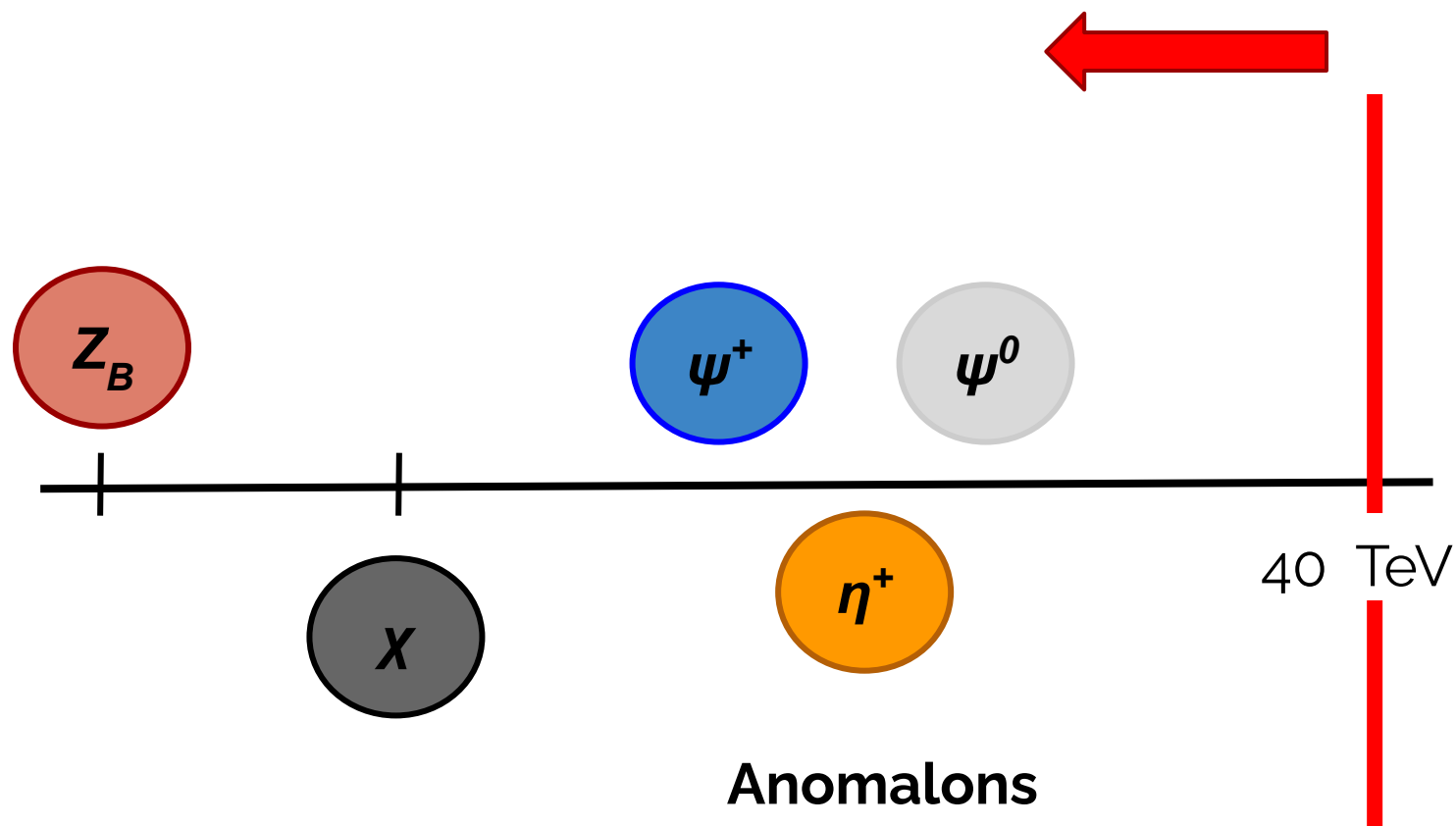
Perturbativity $g_B \leq \frac{\sqrt{2\pi}}{3} \approx 0.84$ and $\Omega h^2 \leq 0.12$

→ Give an upper bound on the scale



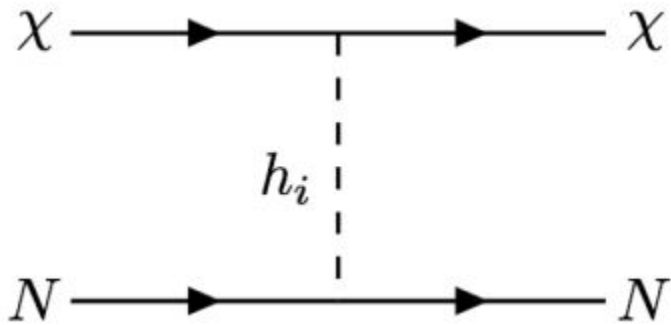
Upper bound on baryon number breaking scale

All masses connected to v_B , and hence, there is an upper bound for the full model



Direct Detection

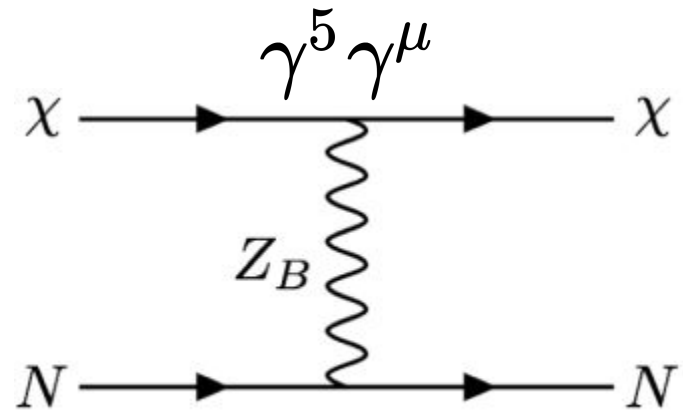
$$\sigma_{\chi N}^{\text{TOT}} = \sigma_{\chi N}(h_i) + \sigma_{\chi N}^0(Z_B)v^2$$



suppressed by Higgs mixing

$$\theta < 0.3 \quad \text{for } M_{H_2} > 200 \text{ GeV}$$

For lighter M_{H_2} stronger bound



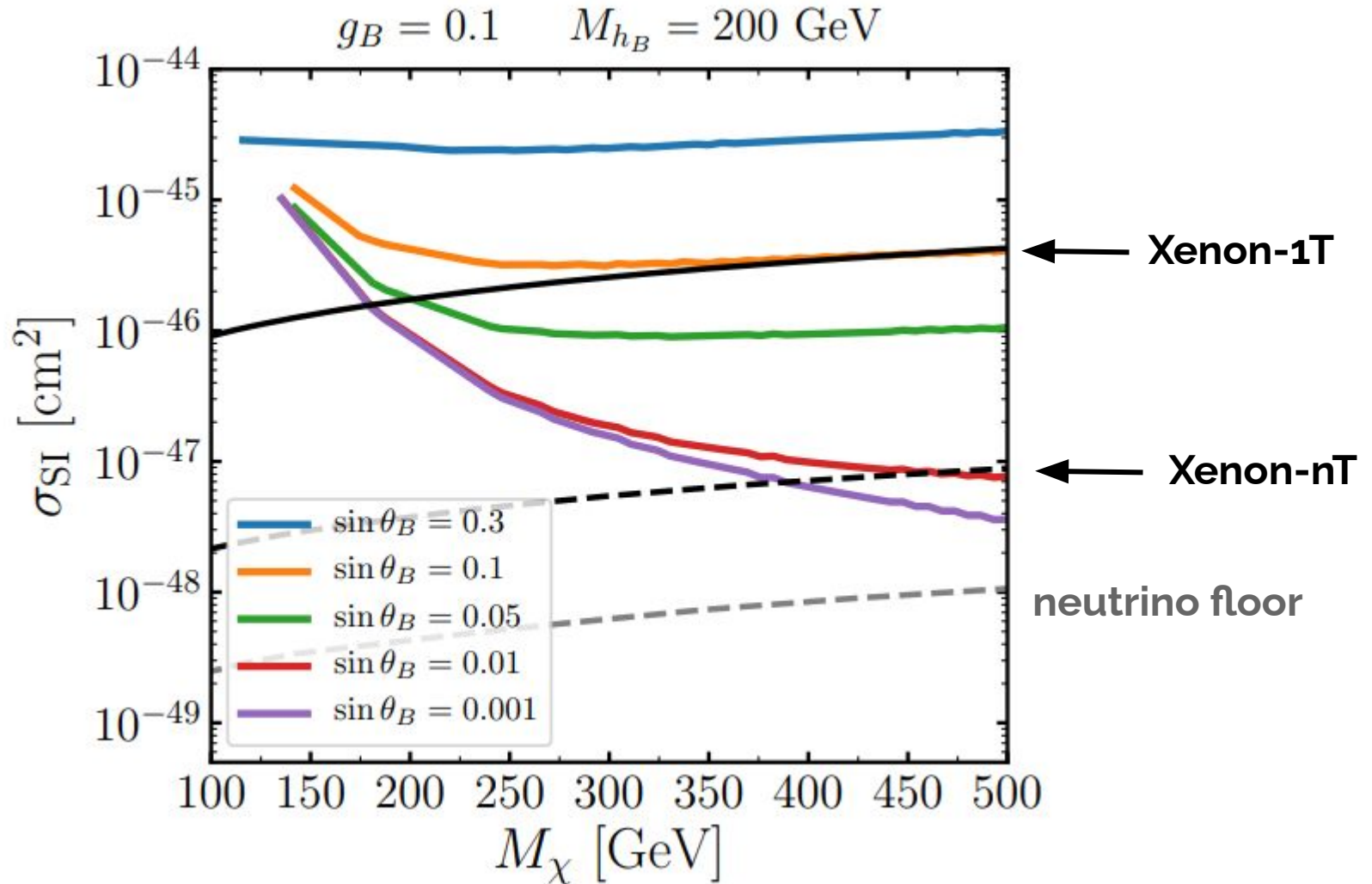
Due to axial coupling,

velocity suppressed $v \sim 10^{-3}$

[Ilnicka, Robens, Stefaniak 2018]

Direct Detection

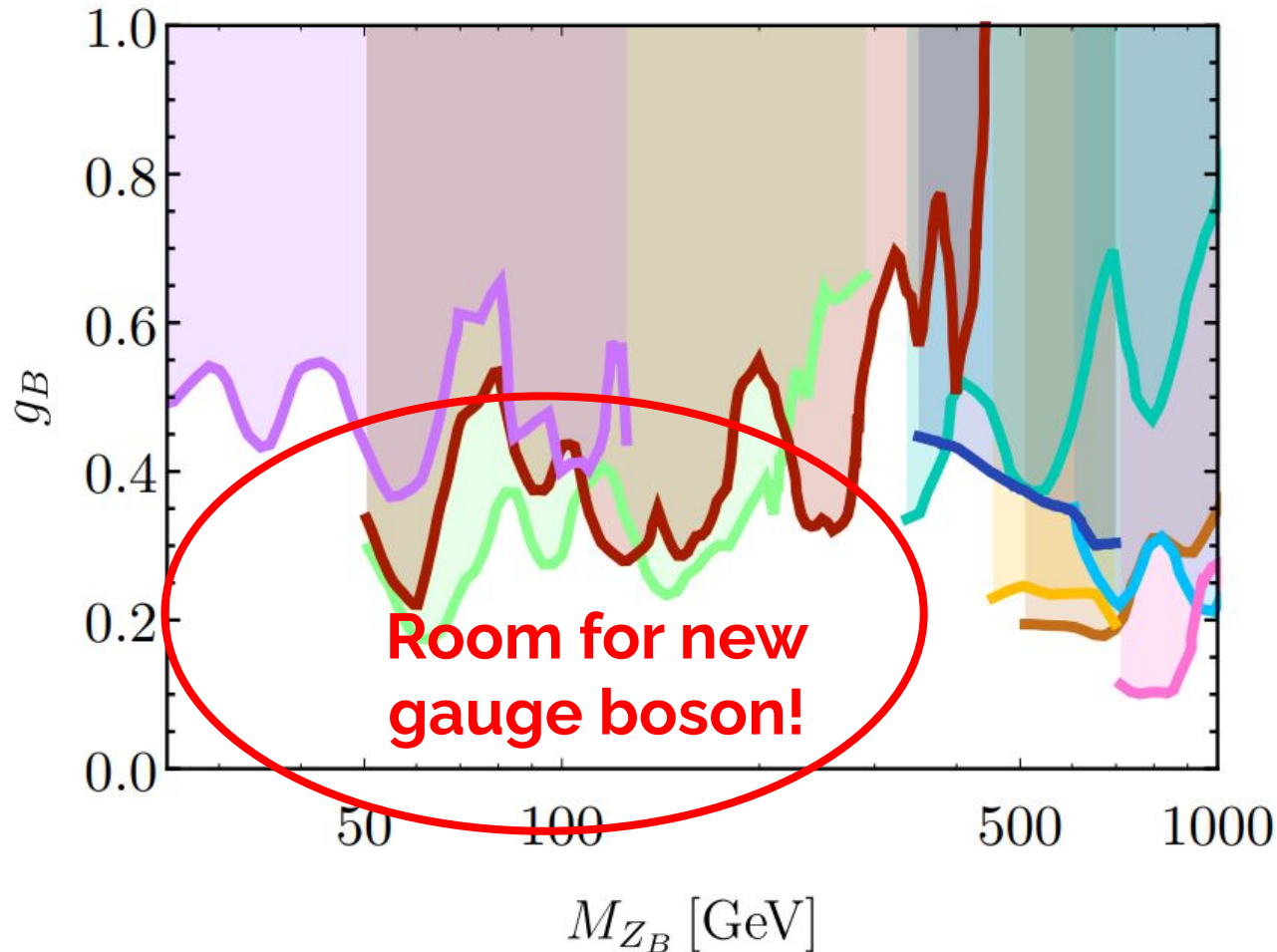
All points satisfy: $\Omega h^2 = 0.12 \pm 0.0022$



2. LHC Phenomenology

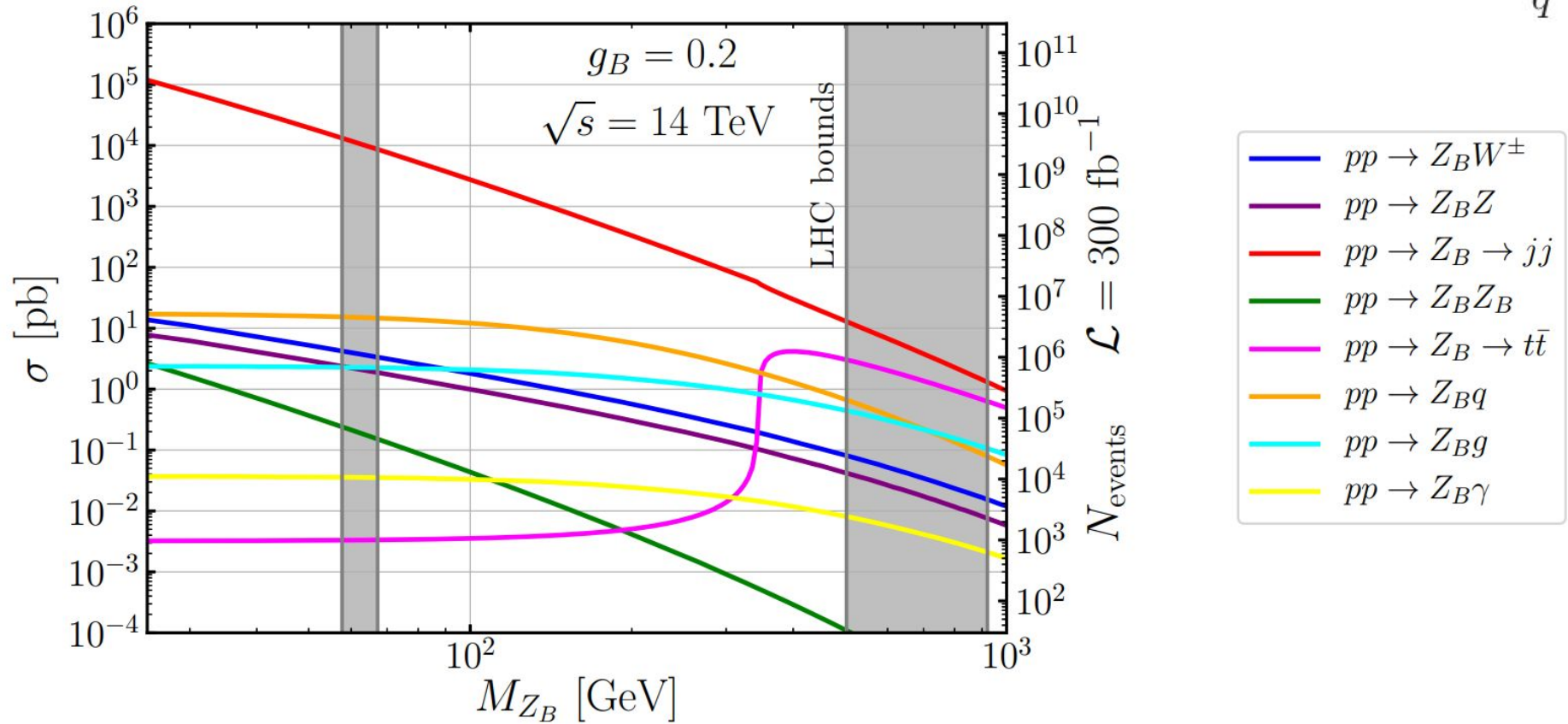
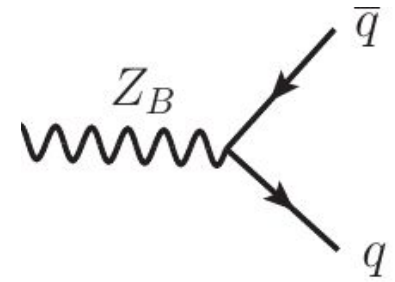
LHC bounds on leptophobic gauge boson

- Di-jet searches at CMS and ATLAS - Run I & II



[Fileviez Perez, Murgui, ADP 2003.09426]

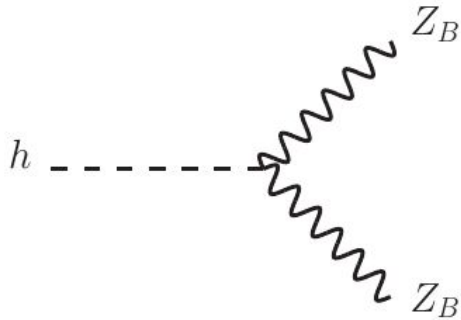
LHC production cross-sections



[Fileviez Perez, Golias, Murgui, ADP 2003.09426]

Exotic Higgs decays

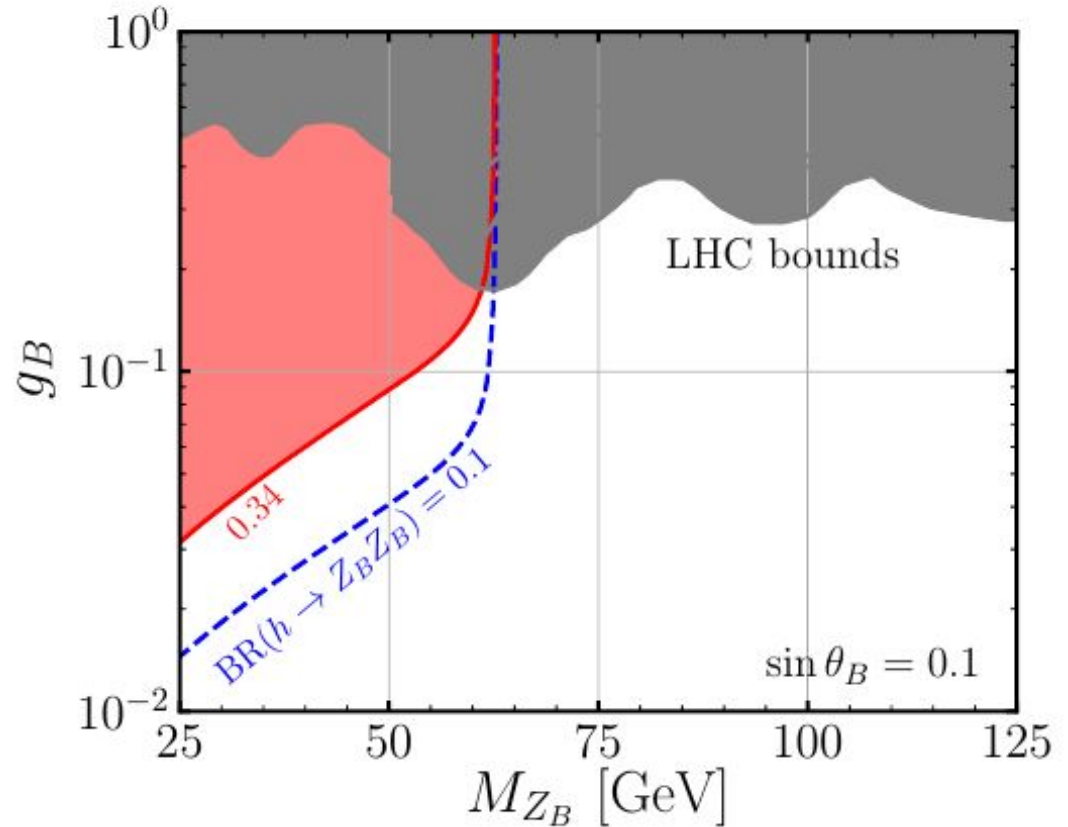
When $M_{Z_B} \leq M_h/2$:



$$hZ_B^\mu Z_B^\nu : 2i \frac{M_{Z_B}^2}{v_B} g^{\mu\nu} \sin \theta_B,$$

$$\text{BR}(h \rightarrow \text{BSM}) \leq 0.34$$

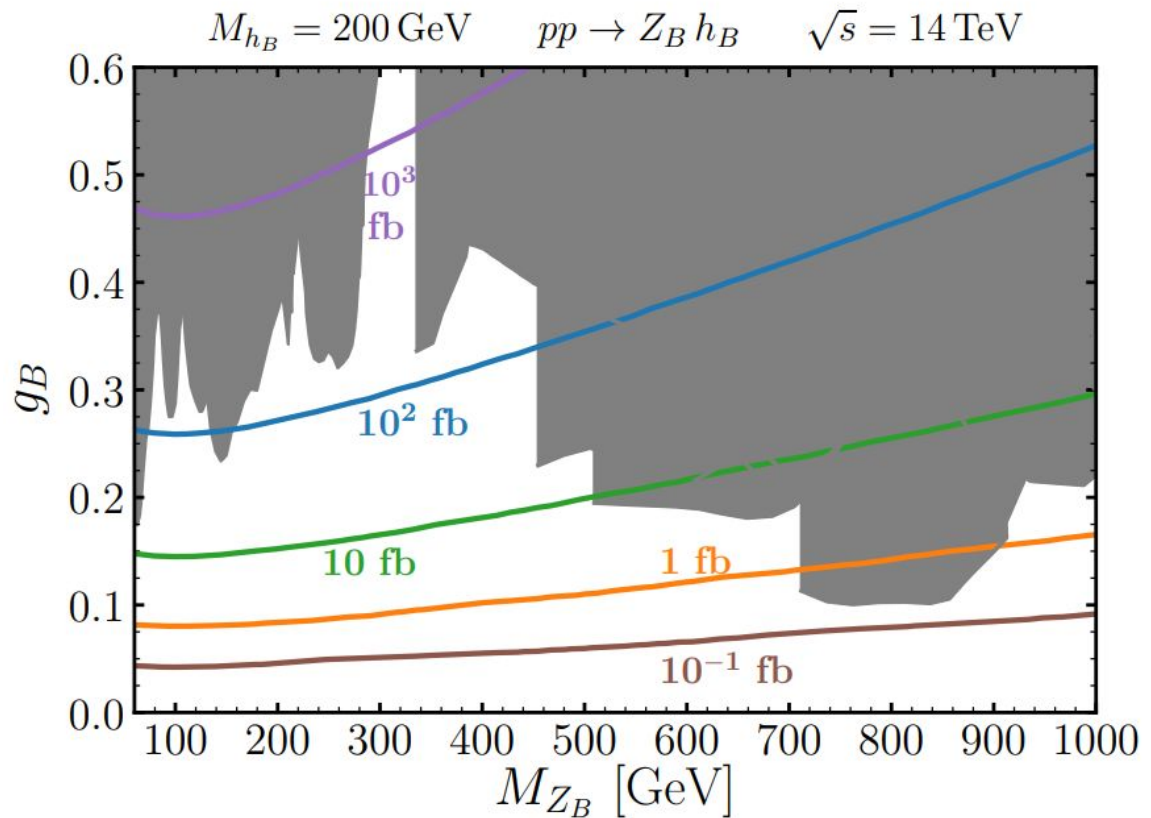
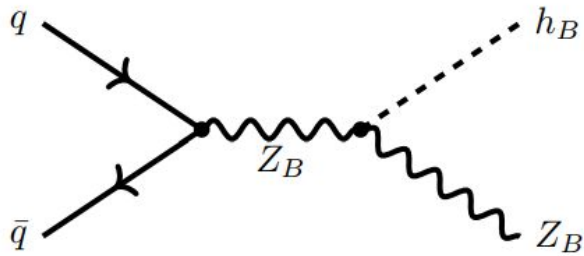
[ATLAS & CMS 1606.02266]



[Fileviez Perez, Golias, Murgui, ADP 2020]

The Higgs of Baryon Number h_B

Associated production does not depend on scalar mixing



[Fileviez Perez, Murgui, ADP 2012.06599]

The Higgs of Baryon Number h_B

We calculate full one-loop decays

Second Higgs can have a large branching ratio into pair of photons

Anomaly-canceling fermions run in the loop!

• $h_B \rightarrow \gamma\gamma$:

$$\Gamma(h_B \rightarrow \gamma\gamma) = \frac{\alpha^2}{64\pi^3 M_{h_B}^5} \left| \cos \theta_B \sum_{F^+} \frac{3g_B M_{F^+}^2}{M_{Z_B}} F_{F^+} + \frac{\sin \theta_B}{v_0} \left(\sum_{f^+} N_c^f Q_f^2 m_{f^+}^2 F_{f^+} - F_W \right) \right|^2.$$

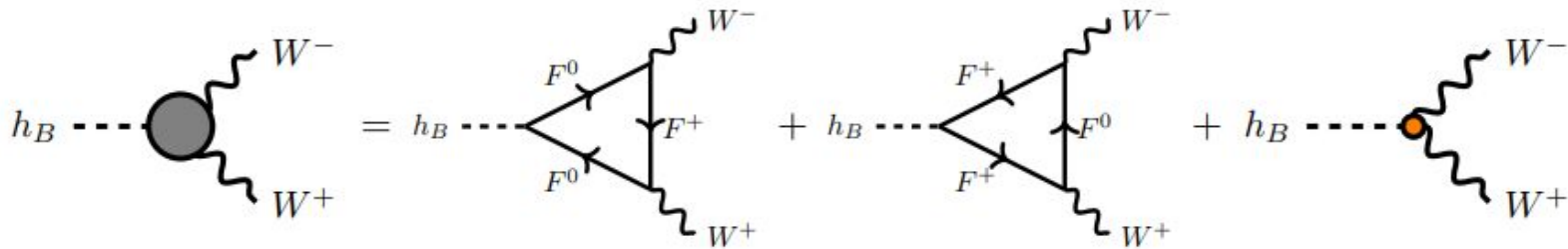
[Fileviez Perez, Murgui, ADP 2012.06599]

The Higgs of Baryon Number h_B

- $h_B \rightarrow WW$:

small mixing with SM h

large mixing



$$\Gamma(h_B \rightarrow WW) = \frac{\sqrt{M_{h_B}^2 - 4M_W^2}}{16\pi M_{h_B}^2 M_W^4} \left(\cos^2 \theta_B \frac{9g_B^2}{M_{Z_B}^2} \left| \sum_F g_{WF}^2 B_F[W] \right|^2 + 2(M_{h_B}^2 - 2M_W^2) \times \right.$$

$$\left. \text{Re} \left\{ \cos \theta_B \frac{3g_B}{M_{Z_B}} \sum_F g_{WF}^2 B_F[W] \left(\cos \theta_B \frac{3g_B}{M_{Z_B}} \sum_F g_{WF}^2 C_F^*[W] + \frac{\sin \theta_B}{v_0} M_W^2 \right) \right\} \right.$$

$$\left. + \left| \cos \theta_B \frac{3g_B}{M_{Z_B}} \sum_F g_{WF}^2 C_F[W] + \frac{\sin \theta_B}{v_0} M_W^2 \right|^2 (M_{h_B}^4 - 4M_{h_B}^2 M_W^2 + 12M_W^4) \right).$$

[Fileviez Perez, Murgui, ADP 2012.06599]

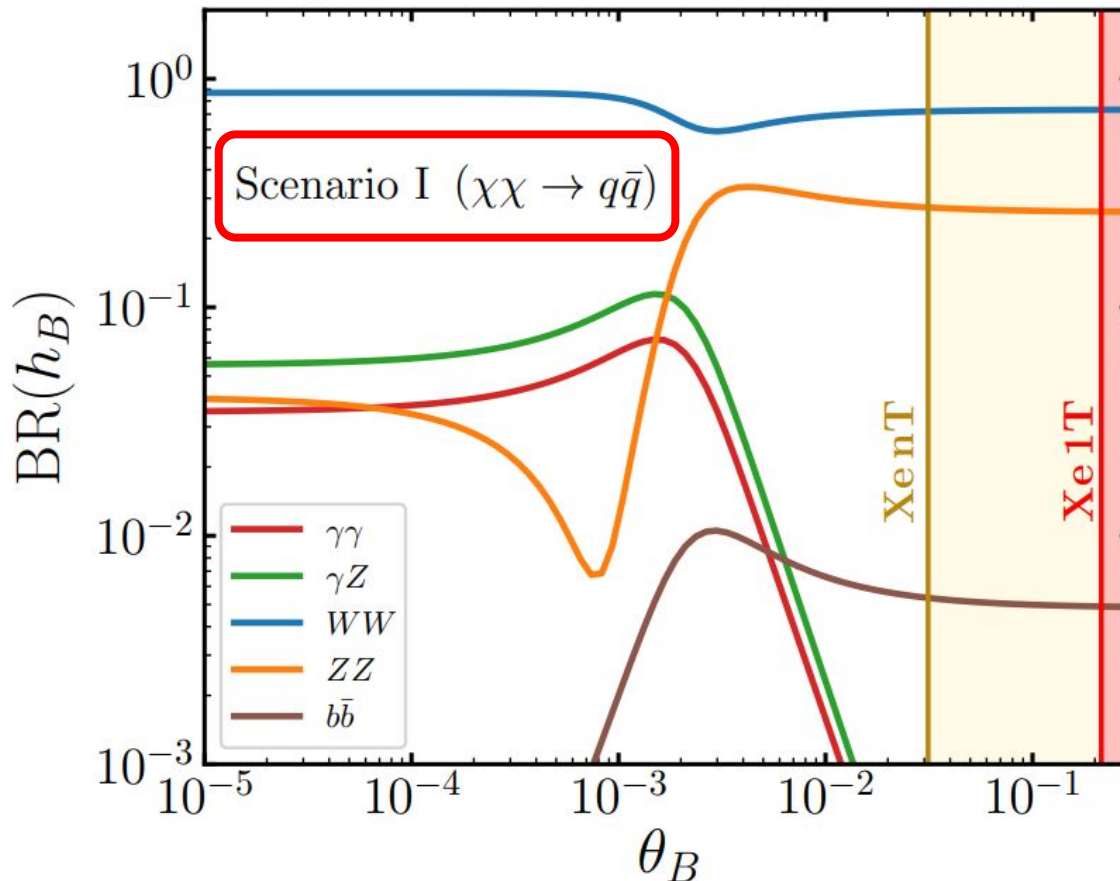
The Higgs of Baryon Number h_B

Correlation between dark matter and the properties of the second Higgs

$$g_B = 0.1 \quad M_{h_B} = 200 \text{ GeV} \quad M_{Z_B} = 436 \text{ GeV}$$

$$M_\chi = 190 \text{ GeV} \quad M_\Psi = 250 \text{ GeV} \quad M_\Sigma = 400 \text{ GeV}$$

$$\Omega_{\text{DM}} h^2 = 0.12$$



$$\text{Br}(h_B \rightarrow \gamma\gamma) \simeq 5\%$$

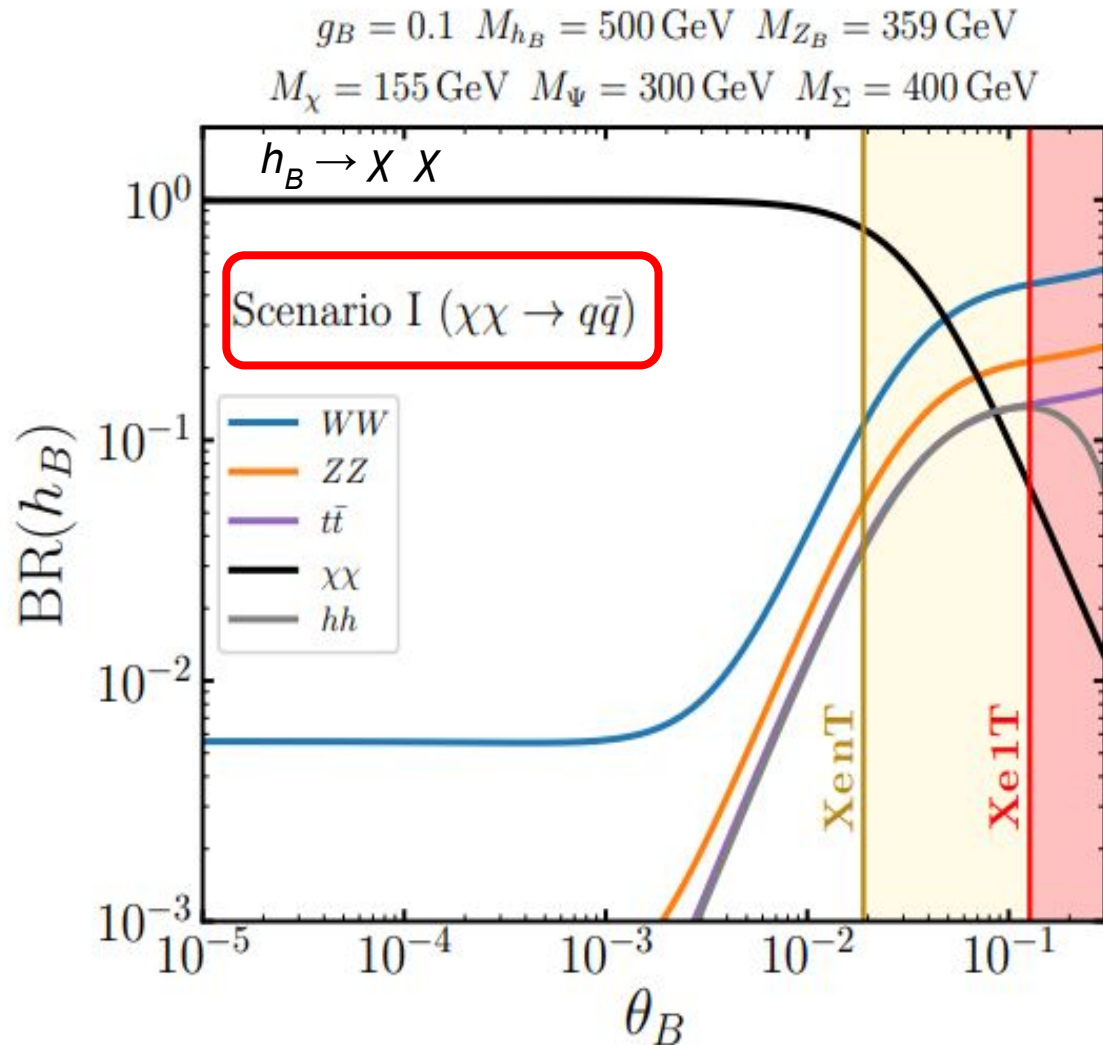
Large di-photon decay!

0.1% for SM Higgs

[Fileviez Perez, Murgui, ADP 2012.06599]

$$M_\chi < M_{h_B} / 2$$

Dominant invisible decay

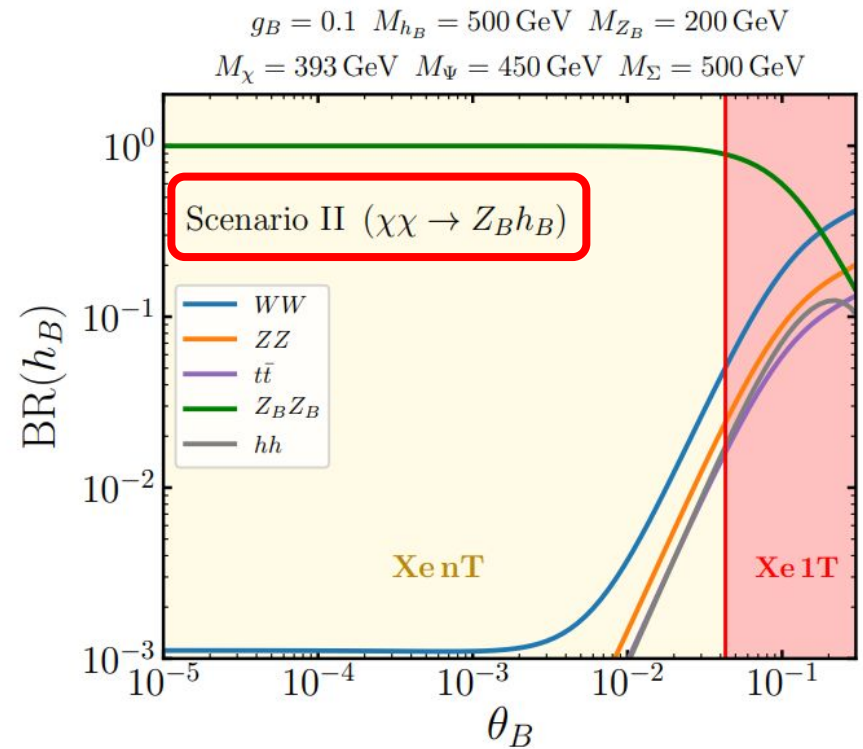
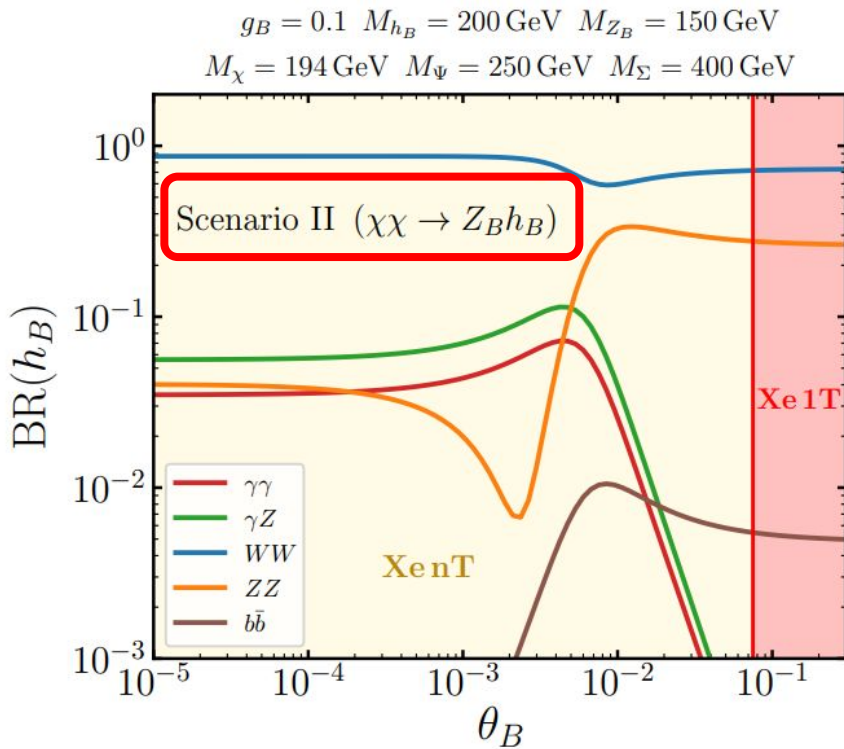


[Fileviez Perez, Murgui, ADP 2012.06599]

The Higgs of Baryon Number h_B

Correlation between dark matter and the properties of the second Higgs

$$h_B \rightarrow Z_B Z_B$$

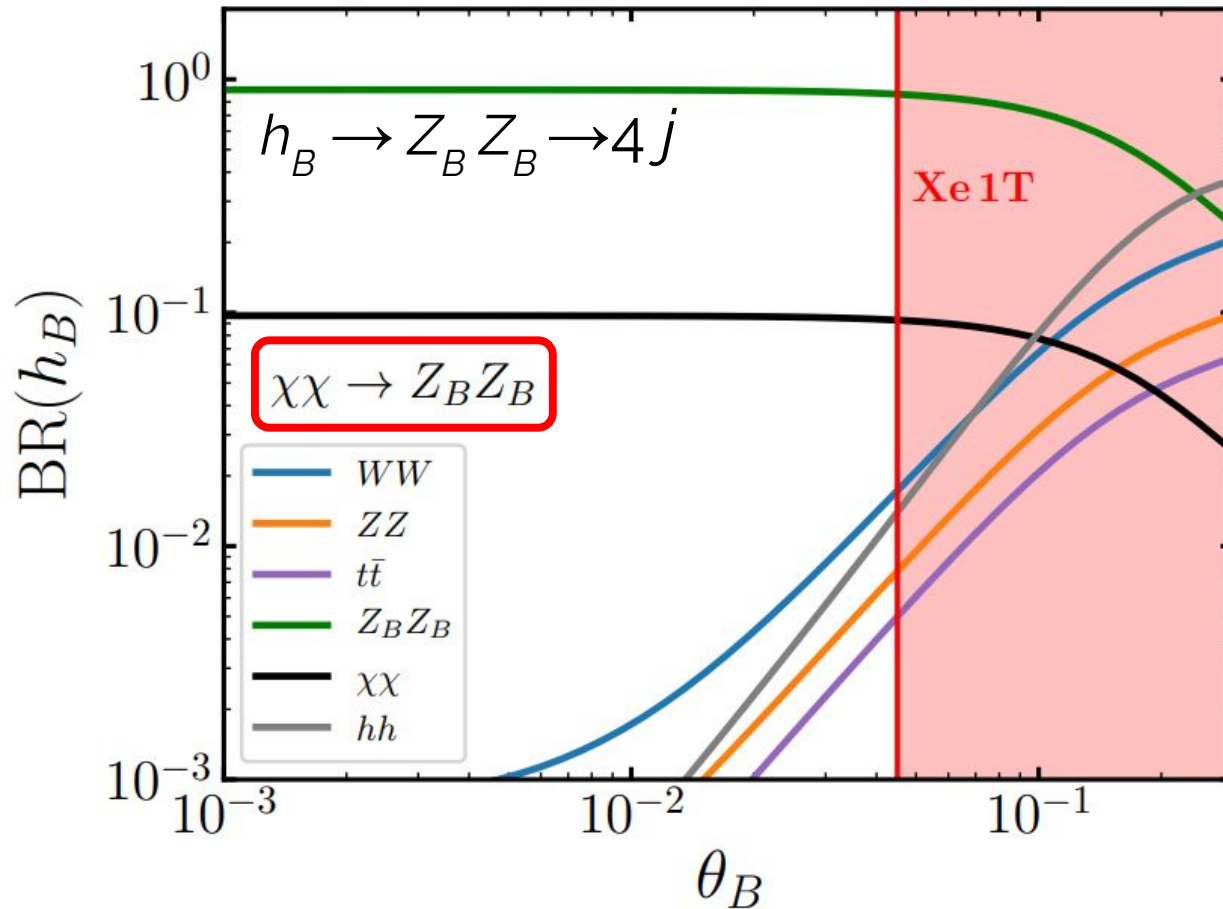


[Fileviez Perez, Murgui, ADP 2012.06599]

The Higgs of Baryon Number h_B

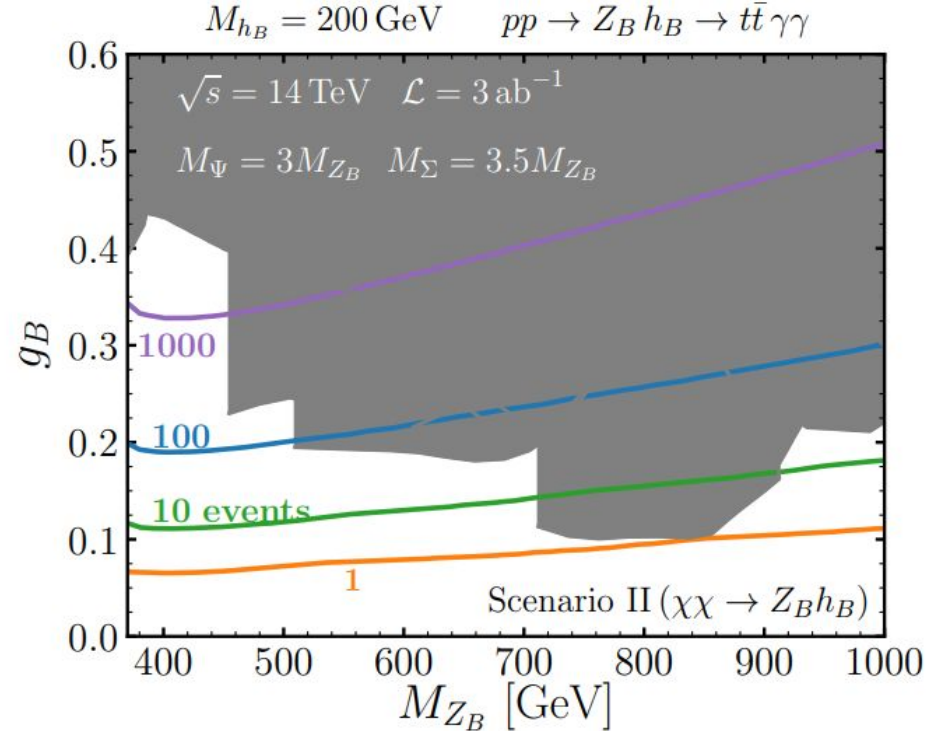
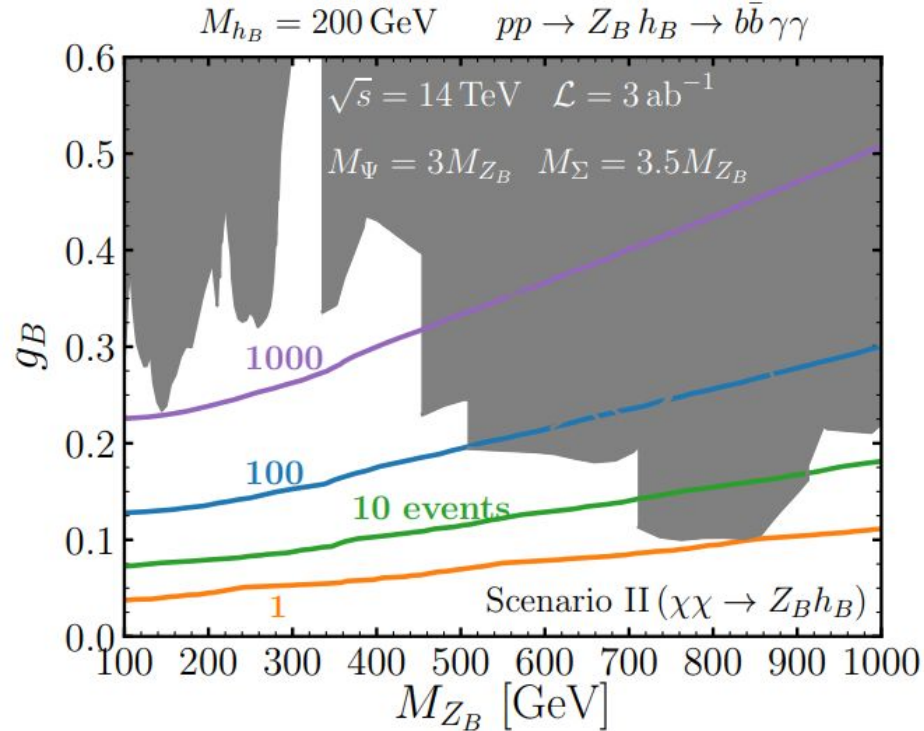
$$g_B = 0.1 \quad M_{h_B} = 500 \text{ GeV} \quad M_{Z_B} = 137 \text{ GeV}$$

$$M_\chi = 200 \text{ GeV} \quad M_\Psi = 300 \text{ GeV} \quad M_\Sigma = 400 \text{ GeV}$$



[Fileviez Perez, Murgui, ADP 2012.06599]

The Higgs of Baryon Number h_B



[Fileviez Perez, Murgui, ADP 2012.06599]

Conclusions

- Promoting baryon and/or lepton number to local symmetries predicts a new sector from anomaly cancellation
- Leptophobic mediator and second Higgs leads to interesting phenomenology at the LHC
- These new fermions lead have implications for the properties of the second Higgs $\text{Br}(h_B \rightarrow \gamma\gamma) \simeq 5\%$
- Not overproducing dark matter gives an upper bound on the full theory at the multi-TeV scale

Thank you!

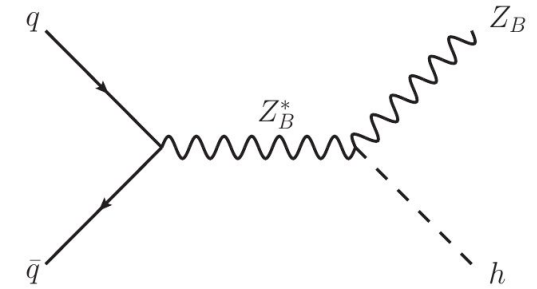
Back-up

Model II

Fields	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$U(1)_L$
$\Psi_L = \begin{pmatrix} \Psi_L^0 \\ \Psi_L^- \end{pmatrix}$	1	2	$-\frac{1}{2}$	$-\frac{3}{2}$
$\Psi_R = \begin{pmatrix} \Psi_R^0 \\ \Psi_R^- \end{pmatrix}$	1	2	$-\frac{1}{2}$	$\frac{3}{2}$
η_R^-	1	1	-1	$-\frac{3}{2}$
η_L^-	1	1	-1	$\frac{3}{2}$
χ_R^0	1	1	0	$-\frac{3}{2}$
χ_L^0	1	1	0	$\frac{3}{2}$

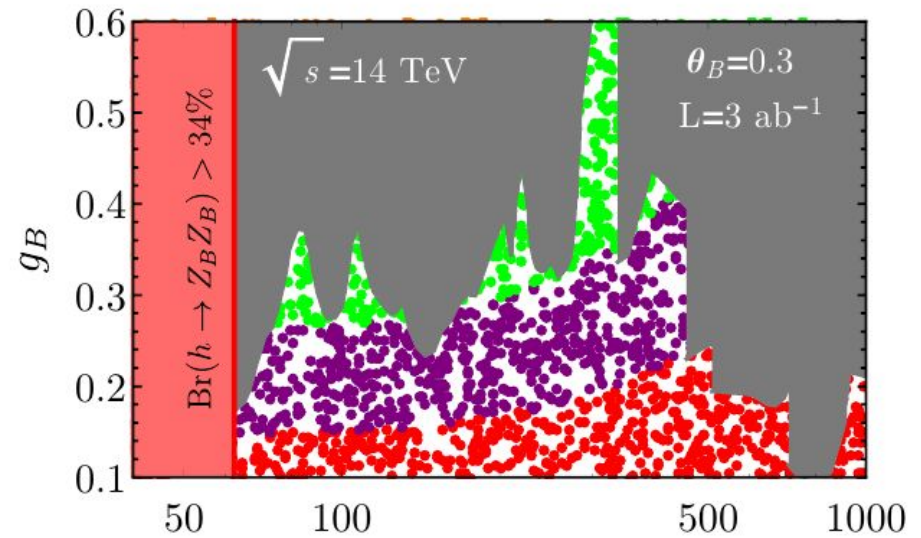
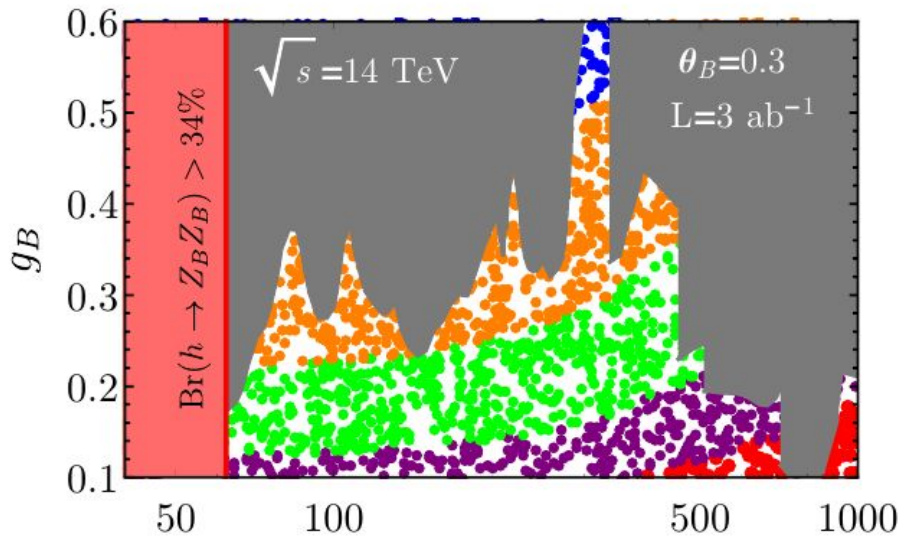
[Duerr, Fileviez Perez & Wise 2013]

Associated Higgs HB HB!! Production



$$pp \rightarrow Z_B h \rightarrow b\bar{b} b\bar{b}$$

$$pp \rightarrow Z_B h \rightarrow \gamma\gamma b\bar{b}$$



- $N_{\text{events}} > 10^5$ ● $10^4 < N_{\text{events}} < 10^5$ ● $10^3 < N_{\text{events}} < 10^4$
- $10^2 < N_{\text{events}} < 10^3$ ● $10 < N_{\text{events}} < 10^2$ ● $N_{\text{events}} < 10$

[Fileviez Perez, Murgui, ADP 2003.09426]