

Probing axion-like particles with $\gamma\gamma$ final states from vector boson fusion processes at the LHC

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Theoretical Origins

- The quantum chromodynamics (QCD) Lagrangian admits a CP (charge conjugation-parity) symmetry violating term, but experiments place stringent constraints on its magnitude; the cause of this suppression is unknown (the **strong CP problem**)
- In 1977, Roberto Peccei and Helen Quinn proposed a solution involving the promotion of the CP violation phase $\bar{\theta}$ to a scalar field which spontaneously broke a new global symmetry
- The quanta of this new scalar field is the **axion**

Axion Properties and Modern Status

- The axion is a *neutral spin-0 boson with negative parity* (i.e., a *pseudoscalar*)
- Strict mass-coupling relationships must hold for the axion solve the strong CP problem; axions satisfying these are denoted **QCD axions** while unconstrained neutral pseudoscalars are **axion-like particles (ALPs)**
- Light ALPs are compatible with current dark matter relic density calculations, making them dark matter candidates
- String theory (ST) has more recently predicted the **axiverse**, a collection of ALPs, incentivizing ALP study and linking ST with ALP phenomenology

The ALP Lagrangian

We adopt an effective field theory approach with cutoff scale Λ .

$$\begin{aligned}
 \mathcal{L} \supset & \frac{1}{2}(\partial_\mu a)^2 - \frac{1}{2}m_a^2 a^2 + \frac{c_1}{\Lambda} \partial_\mu a \bar{f} \gamma_\mu \gamma_5 f - \frac{c_2}{\Lambda} a G_{\mu\nu} \tilde{G}^{\mu\nu} \\
 & - \frac{c_3}{\Lambda} a F_{\mu\nu} \tilde{F}^{\mu\nu} - \frac{c_4}{\Lambda} a F_{\mu\nu} \tilde{Z}^{\mu\nu} - \frac{c_5}{\Lambda} a Z_{\mu\nu} \tilde{Z}^{\mu\nu} \\
 & + \frac{c_6}{\Lambda} (\partial_\mu a)(\partial^\mu a) \phi^\dagger \phi + \frac{c_7}{\Lambda^3} (\partial^\mu a)(\phi^\dagger i D_\mu \phi + h.c.) \phi^\dagger \phi + \dots
 \end{aligned}$$

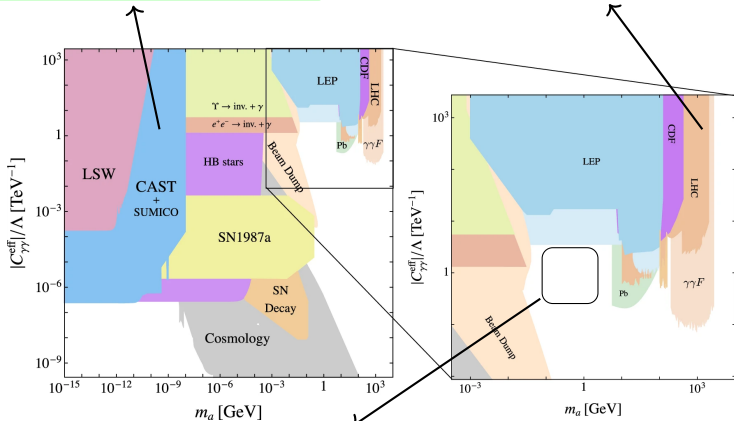
The Lagrangian is annotated with decay channels in colored boxes:

- $a \rightarrow \gamma\gamma$ (pink box, arrow from $\frac{c_3}{\Lambda} a F_{\mu\nu} \tilde{F}^{\mu\nu}$)
- $Z \rightarrow \gamma a$ (pink box, arrow from $\frac{c_4}{\Lambda} a F_{\mu\nu} \tilde{Z}^{\mu\nu}$)
- $a \rightarrow l^+ l^-$ (green box, arrow from $\frac{c_1}{\Lambda} \partial_\mu a \bar{f} \gamma_\mu \gamma_5 f$)
- $a \rightarrow gg$ (pink box, arrow from $\frac{c_2}{\Lambda} a G_{\mu\nu} \tilde{G}^{\mu\nu}$)
- $a \rightarrow ZZ$ (pink box, arrow from $\frac{c_5}{\Lambda} a Z_{\mu\nu} \tilde{Z}^{\mu\nu}$)
- $h \rightarrow aa$ (blue box, arrow from $\frac{c_6}{\Lambda} (\partial_\mu a)(\partial^\mu a) \phi^\dagger \phi$)
- $h \rightarrow Za$ (blue box, arrow from $\frac{c_7}{\Lambda^3} (\partial^\mu a)(\phi^\dagger i D_\mu \phi + h.c.) \phi^\dagger \phi$)

Introduction to ALP Research

Astrophysics (solar axions, magnetar ALP production, etc.)

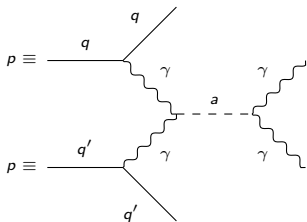
Current LHC constraints
($pp \rightarrow Z \rightarrow \gamma a$; $pp \rightarrow h \rightarrow Za, aa$)



Experimentally unconstrained;
target region

Bauer et al. (2018)

Achieving Novelty: VBF and Non-Resonant Production



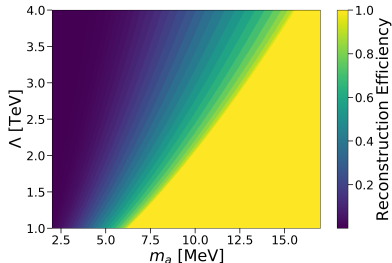
Vector Boson Fusion (VBF)

- The **vector boson fusion** topology derives merit from its distinct LHC signature
- The matrix element magnitude goes as $|\mathcal{M}|^2 \propto m^{jj} / p_T^j$ for outgoing quarks or "tagged jets" j ; maximization occurs for *energetic jets with low transverse momenta* (high pseudorapidity differences)

Non-Resonant Production of ALPs

- The ALP resonant production cross section scales as $\sigma_{\text{res}} \propto m_a^2 / \Lambda^2$ and is suppressed for $m_a \ll \Lambda$; thus *non-resonant ALP production dominates*, enabling sensitivity to MeV-scale ALPs
- With no resonant contribution, diphoton kinematics are driven only by energetic jet pair, yielding further discriminating power
- Lighter ALPs are faster and more stable; requiring ALP decay within the detector constrains the **perpendicular decay length**

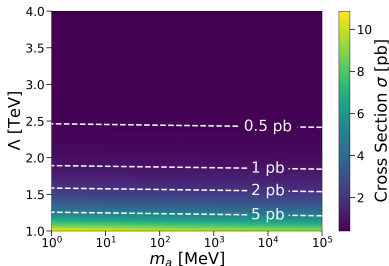
$$L_{a,\perp} = \frac{\sqrt{\gamma_a^2 - 1}}{\Gamma} \sin \theta$$



Signal Generation

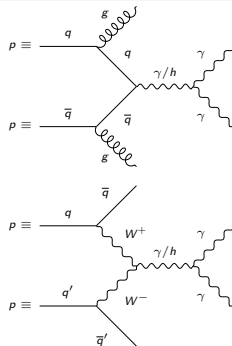
- We generate events using **MadGraph**
- Want sufficient VBF signal statistics for our event selection criteria optimization; to suppress unwanted contributions to $pp \rightarrow ajj$ ($a \rightarrow \gamma\gamma$) event generation (e.g., gg fusion, associated ALP production), we impose *MadGraph-level selections* on signal events:

$$|\Delta\eta^{jj}| > 2.4, m^{jj} > 120 \text{ GeV}$$

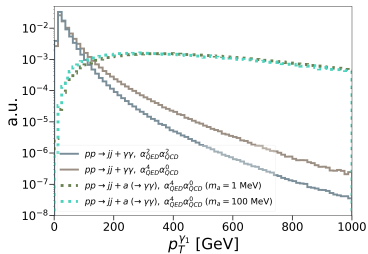
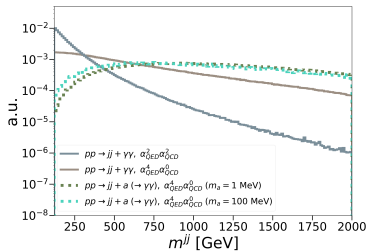
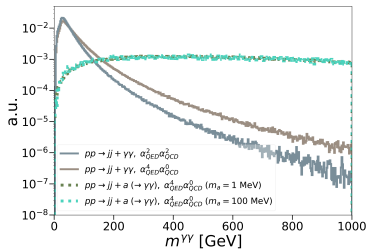
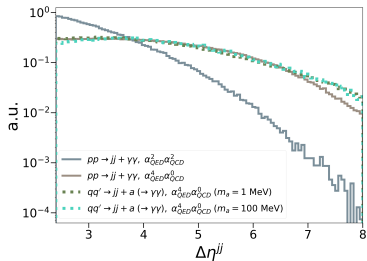


Background Generation

- The dominant Standard Model background processes are a mixed QED-QCD channel $pp \rightarrow jj\gamma\gamma$ and a pure electroweak channel $pp \rightarrow jj\gamma\gamma$ ($\alpha_{\text{QCD}} = 0$)
- Recognizing our eventual selection of high jet momentum events, we generate BG events in H_T bins to ensure sufficient high-energy statistics



Pre-Selection Kinematics



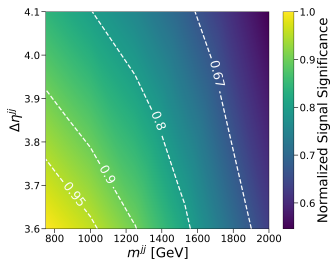
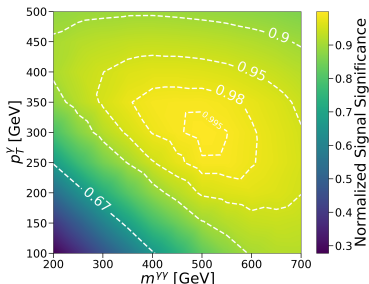
Optimizing Event Selection Criteria

Process

- We adopt the following signal significance (SS) metric; note our conservative estimation of systematic error

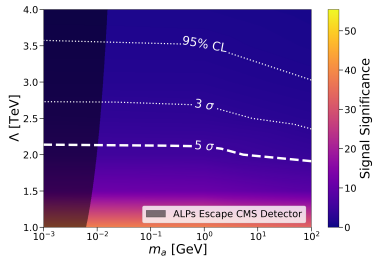
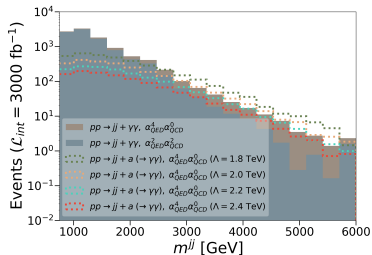
$$\frac{S}{\sqrt{S + B + (0.25 \cdot (S + B))^2}}$$

- Using this metric, we optimize event selection criteria on two kinematic variables simultaneously by sampling SS on a grid



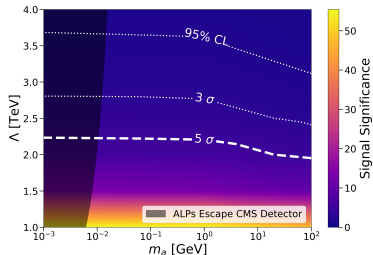
Criterion	$\gamma_1\gamma_2j_1j_2$
Central Selections	
$ \eta^\gamma $	< 2.5
p_T^γ	> 30 GeV
$p_T^{\gamma_1}$	> 300 GeV
$m^{\gamma\gamma}$	> 500 GeV
$N(\ell), N(b)$	$= 0$
VBF Selections	
p_T^j	> 30 GeV
$ \eta^j $	< 5.0
$\Delta R^{\gamma j}$	> 0.4
$N(j)$	≥ 2
$\eta^{j_1} \cdot \eta^{j_2}$	< 0
$ \Delta\eta^{jj} $	> 3.6
m^{jj}	> 750.0 GeV

Results: Signal Significance in the Parameter Space

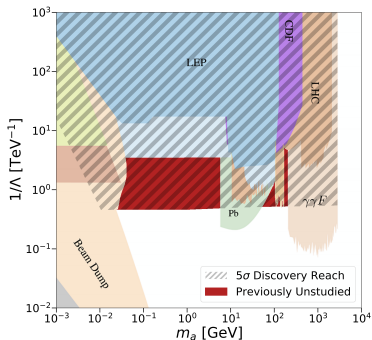


Comments

- On the right we depict the signal significance achieved by our selections as a function of m_a and Λ for two integrated luminosities: 150 fb^{-1} (LHC run II, top) and 3000 fb^{-1} (high luminosity LHC, bottom)
- We have discovery potential for a significant range of ALP masses ($\sim \text{MeV}$ scale to TeV scale) in the region $\Lambda \lesssim 2.25 \text{ TeV}$



Discussion and Summary



Discussion

- We overlay our discovery region on the plot of existing ALP constraints shown at the beginning of this talk
- In particular, we see that our methodology constrains a significant portion of the parameter space and broadens the LHC constraint region, including unprecedented lower mass/weak coupling scenarios

Summary

- We pursue a phenomenological study of ALPs, a class of particles well motivated by modern problems in the Standard Model as well as by string theory
- While ALPs are probed in a variety of settings, we take interest in the high mass, strong coupling scenario and employ a collider approach
- The unique detector signature of the VBF topology and the domination of non-resonant ALP production together provides several kinematic variables with distinct discrimination power
- Consequently, an optimization of event selection criteria yields discovery potential in a substantial region of the ALP parameter space
- In particular, our approach makes novel contributions to the extent of LHC constraints on the ALP parameter space, including the incorporation of previously unstudied regions

Thank you!