Challenges for an axion explanation of the muon g-2 measurement

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$$\Delta a_{\mu} = a_{\mu}^{\rm EXP} - a_{\mu}^{\rm SM} = (25.1 \pm 5.9) \times 10^{-10}$$

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Let us sketch the type of BSM that could explain $(g-2)_{\mu}$:

- Absence of $(g 2)_e$, $a_e^{EXP} a_e^{SM} = (4.8 \pm 3.0) \times 10^{-13}$ $\Delta a_e \sim \left(\frac{m_e}{m_{\mu}}\right)^2 \Delta a_{\mu} \approx 5.9 \times 10^{-14}$
- Constraints from EDM, $|d_e| < 1.1 \times 10^{-29} \ e \ {\rm cm}$ and $|d_{\mu}| < 1.9 \times 10^{-19} \ e \ {\rm cm}$

$$\begin{aligned} |d_e| &\sim \frac{e}{2m_{\mu}} \Delta a_{\mu} \approx 2.3 \times 10^{-22} e \text{ cm}, \\ |d_{\mu}| &\sim \frac{m_e}{m_{\mu}} d_{\mu} \approx 1.1 \times 10^{-24} e \text{ cm}. \end{aligned}$$

• Constraints from flavor violating processes $Br(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13}$

$$Br(\mu \to e\gamma) \sim \frac{6\pi^2 e^2}{G_F^2 m_\mu^4} (\Delta a_\mu)^2 \approx 2.0 \times 10^{-3}.$$

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makes the naturalness problem worse, unless embedded in a bigger theory that addresses it

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benign to the naturalness problem by themselves

Axion EFT

Axion EFT – Cont'd



$$\Delta a_{\mu}^{(1)} \propto -\frac{c_{\mu\mu}^2}{16\pi^2}, \qquad \Delta a_{\mu}^{(2)} \propto -\frac{c_{\mu\mu}c_{\gamma\gamma}\alpha}{16\pi^3}, \qquad \Delta a_{\mu}^{(3)} \propto -\frac{c_{\mu\mu}c_{ii}\alpha}{16\pi^3},$$

We consider two scenarios

•
$$c_{\mu\mu} \neq 0$$
, $c_{\gamma\gamma} \neq 0$
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$$\begin{aligned} c_{\gamma\gamma}/c_{\mu\mu} &< 0, \quad m_a \subset (40 \,\mathrm{MeV} - 200 \,\mathrm{GeV}) \\ \left| \frac{f_a}{c_{\gamma\gamma}} \right| &\lesssim (10 - 25) \,\mathrm{GeV}, \quad \left| \frac{f_a}{c_{\mu\mu}} \right| &\lesssim 100 \,\mathrm{GeV} \;. \end{aligned}$$

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Collider Bounds of Axion EFT



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• $c_{\mu\mu} \neq 0$, $c_{ee} \neq 0$



$$\begin{split} m_a \gtrsim 2 \,\text{GeV}, \quad c_{ee}/c_{\mu\mu} < 0, \\ \left| \frac{f_a}{c_{\mu\mu}} \right| \lesssim 100 \,\text{GeV}, \quad \left| \frac{f_a}{c_{ee}} \right| \lesssim 25 \,\text{GeV} \quad \text{for} \quad m_a = 5 \,\text{GeV} \end{split}$$

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Model completion and challenges

• μ is directly charged under PQ

• μ is neutral under PQ, but coupled to other fields charged under $U(1)_{PQ}$.

Outline of the ingredients:

- Scalars content: H_1 , H_2 , ...; Φ
- All break PQ
- H_1 , H_2 ... break EW
- Φ is SM singlet

 $\Rightarrow \mbox{ For } v_1 \sim v_2 \sim \ldots \sim v_{\Phi}, \quad v_{EW}^2 = v_1^2 + v_2^2 + \ldots, \quad f_a \sim v_{EW}^2 + v_{\Phi}^2. \mbox{ In } \mbox{ case of invisible axion, one only needs to make } v_{\Phi} \gg v_{EW}.$

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 \Rightarrow One solution is to have $v_1 \gg v_2, v_{\Phi}$, so the component eaten by Z is (mostly) H_1 , and the axion is (mostly) a linear combination of H_2 and Φ . This gives $v_{EW} \sim v_1$, $f_a \sim v_2, v_{\Phi}$.

Field	$SU(2)_L$	$U(1)_Y$	$U(1)_{PQ}$
H_l	2	$-\frac{1}{2}$	+1
H_q	2	$+\frac{1}{2}$	+1
Φ	1	0	+1
U^c	1	$-\frac{2}{3}$	-1
D^c	1	$+\frac{1}{3}$	+1
E^c	1	+1	-1

$$\begin{aligned} V_{scalar} &= V_0(|H_l|, |H_q|, |\Phi|, |H_lH_q|) + \lambda_{ql\Phi} H_l H_q \Phi^{\dagger 2} \\ V_{int} &= y_u H_q Q U^c + y_d H_q^{\dagger} Q D^c + y_e H_l L E^c + \text{h.c.} \end{aligned}$$

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Identify the pNGB's

- $(v_q^2 \theta_q v_l^2 \theta_l)$ is eaten by Z
- $(\theta_l+\theta_q-2\theta_\Phi)$ gets massive due to $\lambda_{ql\Phi}H_lH_q\Phi^{\dagger 2}$
- light axion mode a

$$a = \frac{1}{f_a} \left(v_{\Phi}^2 \theta_{\Phi} + 2 \frac{v_q^2 v_l^2}{v_{EW}^2} (\theta_l + \theta_q) \right)$$
$$\approx \frac{1}{\sqrt{v_{\Phi}^2 + 4v_l^2}} (v_{\Phi}^2 \theta_{\Phi} + 2v_l^2 \theta_l) + \mathcal{O}\left(\frac{v_{l,\Phi}}{v_{EW}}\right)$$

$$\begin{aligned} \mathcal{L}_{af} \supset (\partial_{\mu}\theta_{l})E^{c\dagger}\bar{\sigma}^{\mu}E^{c} + (\partial_{\mu}\theta_{q})U^{c\dagger}\bar{\sigma}^{\mu}U^{c} - (\partial_{\mu}\theta_{q})D^{c\dagger}\bar{\sigma}^{\mu}D^{c} \\ \mapsto \underbrace{\frac{2}{f_{a}}\frac{v_{q}^{2}}{v_{\rm EW}^{2}}}_{-c_{\mu\mu}}(\partial_{\mu}a)E^{c\dagger}\bar{\sigma}^{\mu}E^{c} + \frac{2}{f_{a}}\frac{v_{l}^{2}}{v_{\rm EW}^{2}}(\partial_{\mu}a)\left(U^{c\dagger}\bar{\sigma}^{\mu}U^{c} - D^{c\dagger}\bar{\sigma}^{\mu}D^{c}\right), \end{aligned}$$

$$\mathcal{L}_{a\gamma} \supset 3(\theta_l + \theta_q) \frac{\alpha}{4\pi} F_{\mu\nu} \tilde{F}^{\mu\nu}$$
$$\mapsto \underbrace{\frac{6}{f_a} \frac{v_q^2 + v_l^2}{v_{\rm EW}^2}}_{c_{\gamma\gamma}} \frac{\alpha}{4\pi} a F_{\mu\nu} \tilde{F}^{\mu\nu} = \frac{6}{f_a} \frac{\alpha}{4\pi} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

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Observations:

- $v_q \sim v_{\rm EW} > v_l \sim v_{\Phi}$, axions dominantly couples to ℓ and γ
- no a g coupling b/c U^c and D^c have opposite PQ
- $c_{\mu\mu}$ has the opposite sign as $c_{\gamma\gamma}$

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- Close to alignment limit, s.t. couplings are SM-like
- Radial mode of Φ could mix with SM higgs, $h \to aa$
- Extra ingredients

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Lepton-specific 2HDM: 1106.0034, PRD41 3421, 1207.4835, 1305.2424, Global EW fit on 2HDM: 1803.01853 • μ is directly charged under PQ

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- ℓ doesn't couple to axion directly
- new vector-fermion, χ , is PQ-charged
- ℓ mass-mixes with χ
- mass-mixing $\ll \chi$ mass
- ℓ inherits "fractional" PQ charge

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$$\mathcal{L}_{\text{eff}} \supset -\frac{\partial_\mu a}{2f_a} \left(\left| \frac{y_{Ei} f_a}{M_i} \right|^2 E_i^{c\dagger} \bar{\sigma}^\mu E_i^c \right)$$

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1402.7065

- Axions are benign candidates for light modes
- From EFT point of view, axion could address $(g-2)_{\mu}$
- Yet, completing the model leads to serious challenges:
 - it requires large axion coupling, $f/c_{ae}, f/c_{a\mu}, f/c_{a\gamma} \gtrsim (100 \text{ GeV})^{-1}$
 - light charged modes $\sim \mathcal{O}(10) \mathcal{O}(100) \text{ GeV}$
 - light neutral modes that mixes with Higgs
 - direct contribute to $(g-2)_{\mu}$ from extra ingredient