

An even lighter QCD axion .

PHENO 21 - May 24th 2021



Pablo Quílez Lasanta - pablo.quilez@desy.de

Based on “An even lighter QCD axion” [2102.00012](#)

“Dark matter from an even lighter QCD axion: trapped misalignment” [2102.01082](#)

In collaboration with L. Di Luzio, B. Gavela and A. Ringwald

The QCD axion

- Solves the Strong CP problem
- Excellent Dark Matter candidate

[Peccei+Quinn 77]

[Weinberg, 78]

[Wilczek, 78]

[Abbot+Sikivie, 83]

[Dine and W. Fischler, 83]

[Preskil et al, 91]

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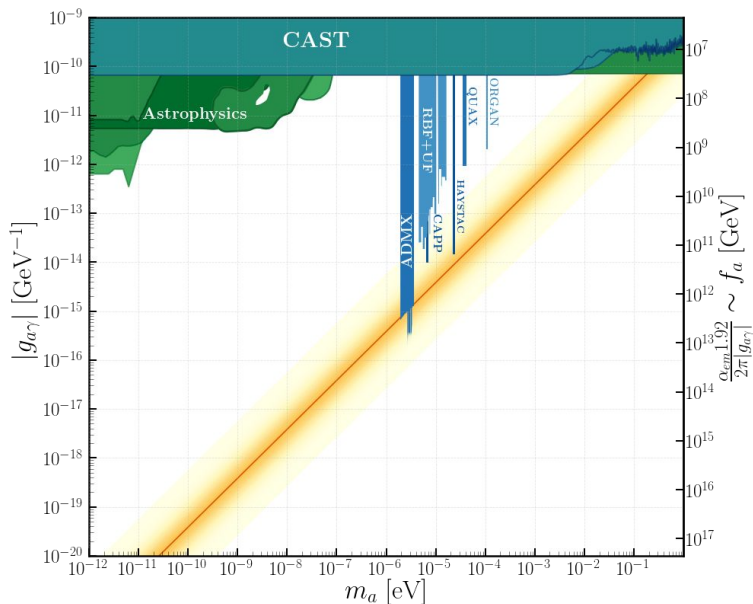
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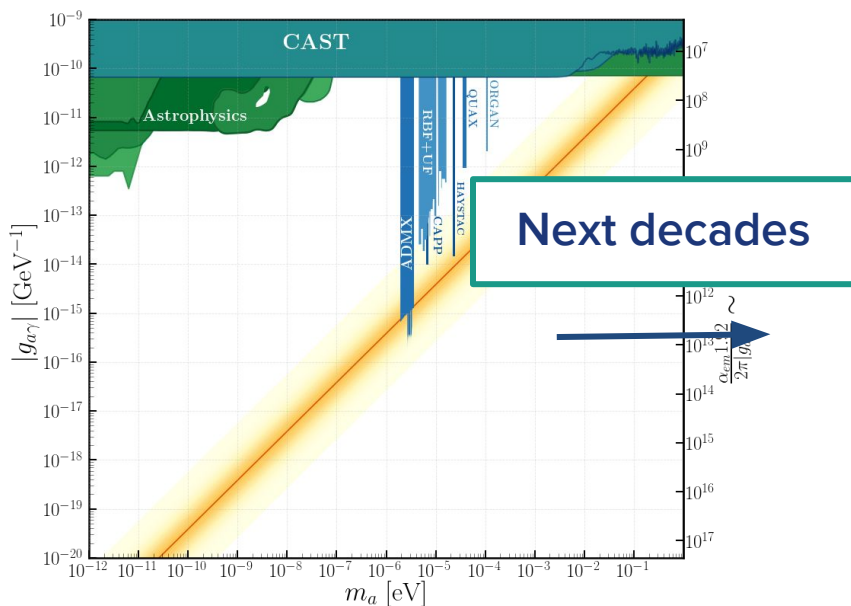
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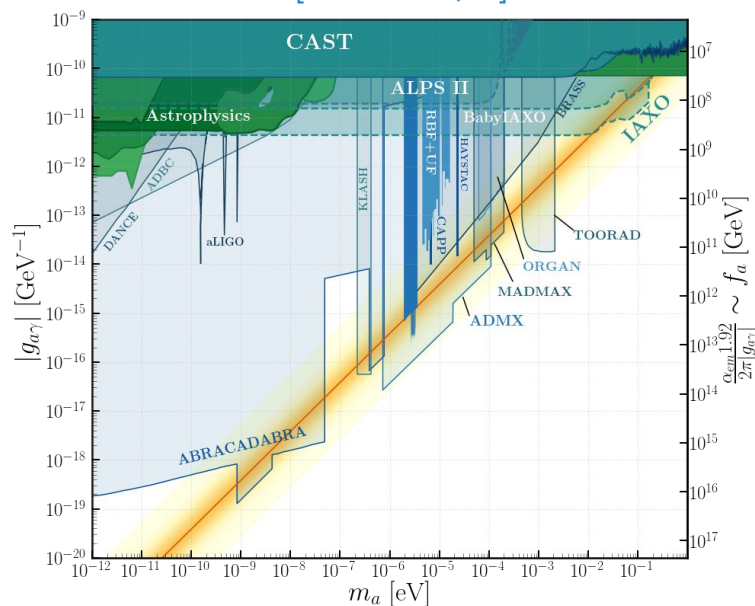
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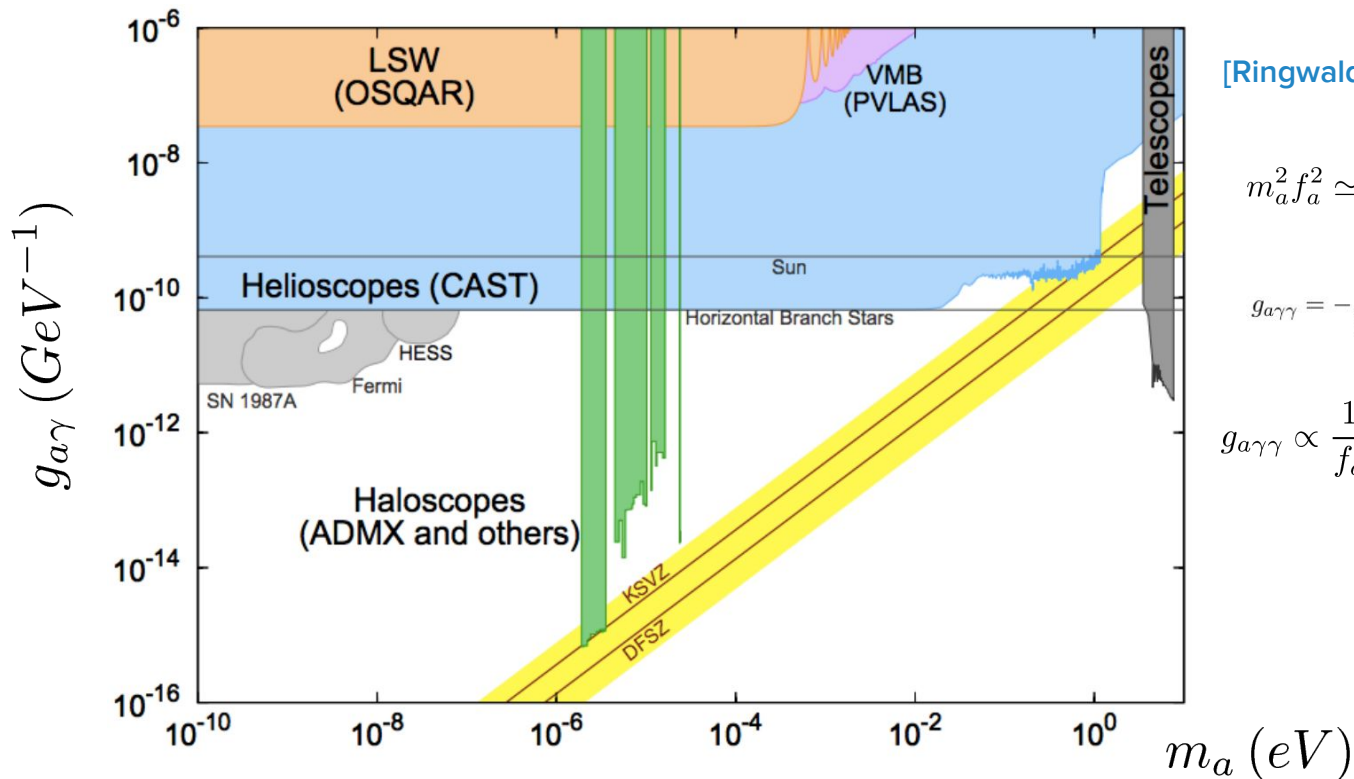
[Dine and W. Fischler, 83]

[Preskil et al, 91]



Based on **2102.00012** and **2102.01082**

Invisible axion parameter space



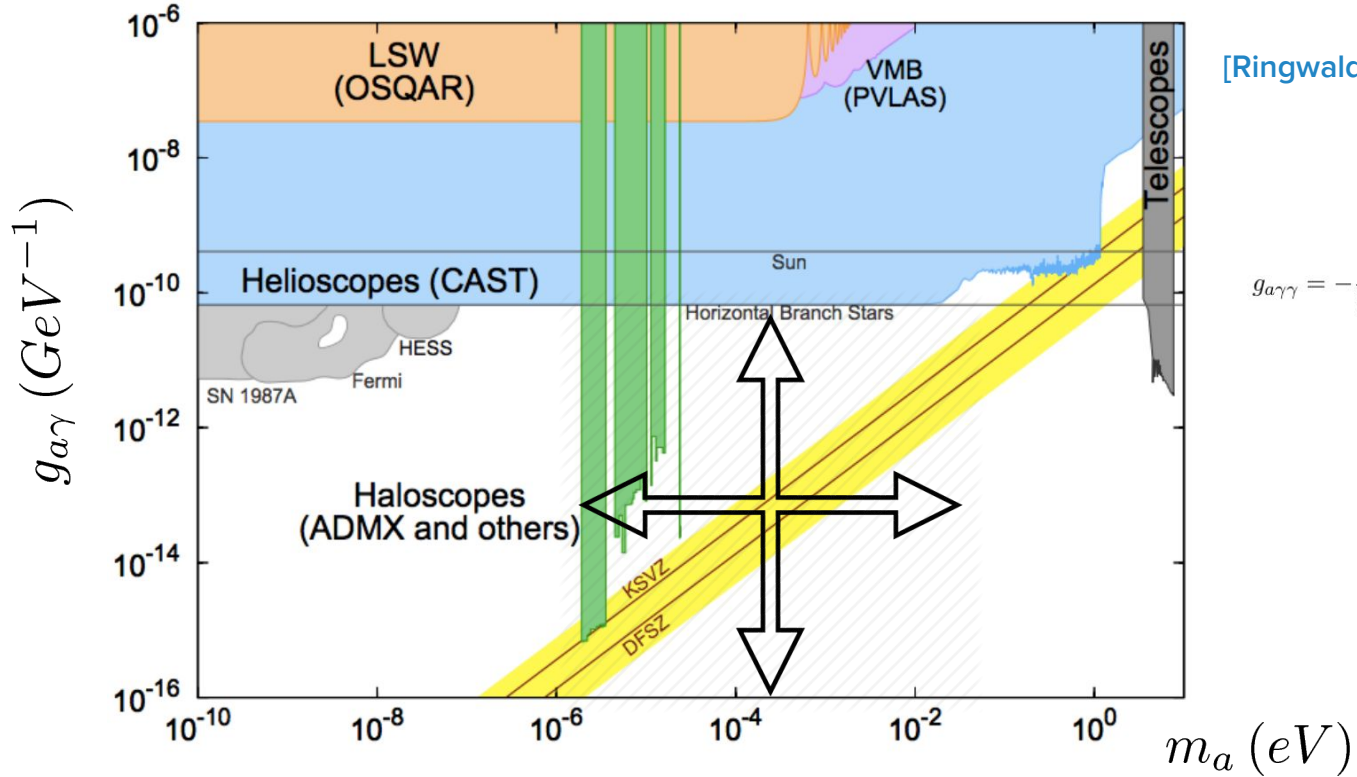
[Ringwald, PDG 17]

$$m_a^2 f_a^2 \simeq m_\pi^2 f_\pi^2 \frac{m_u m_d}{(m_u + m_d)^2}$$

$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{\text{em}} \left(\frac{E}{N} - 1.92(4) \right)$$

$$g_{a\gamma\gamma} \propto \frac{1}{f_a} \implies g_{a\gamma\gamma} \propto m_a$$

Are there other possibilities?

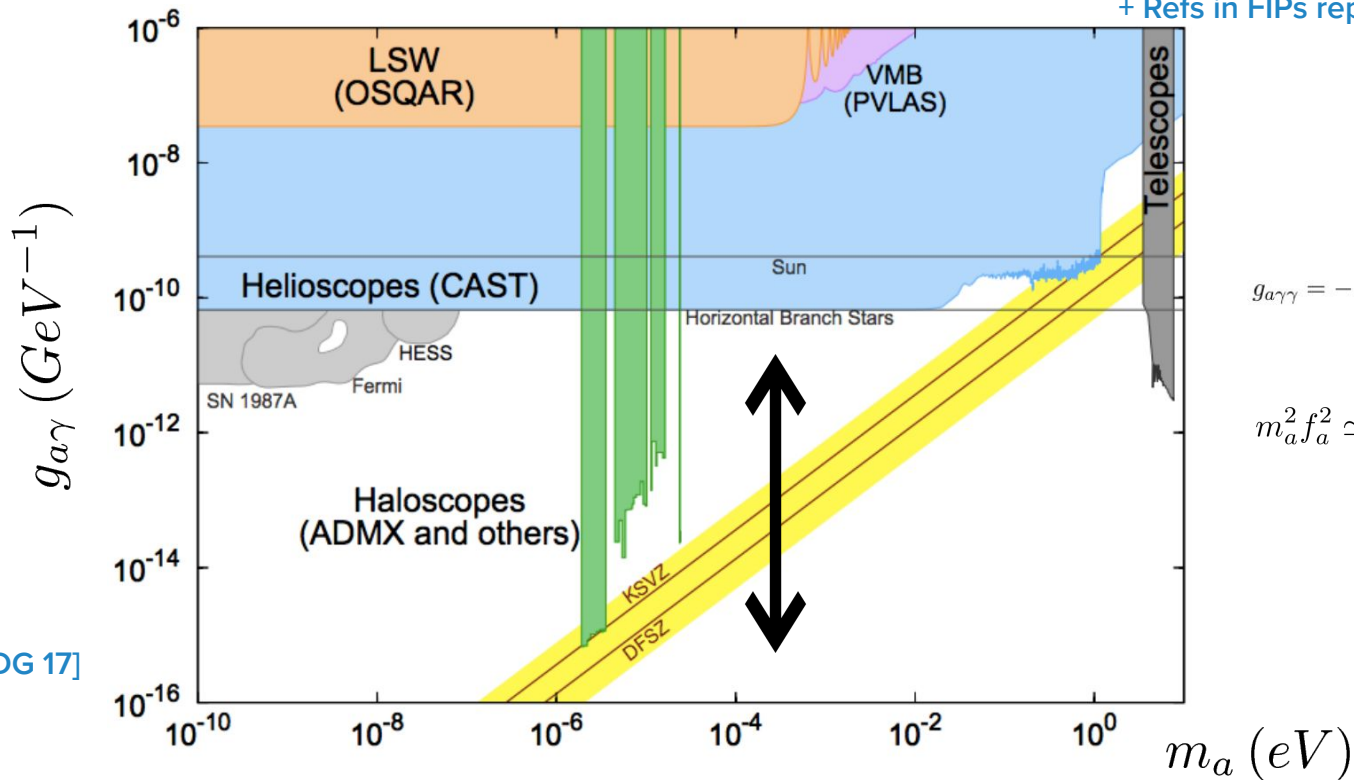


[Ringwald, PDG 17]

$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{\text{em}} \left(\frac{E}{N} - 1.92(4) \right)$$

Photophilic/photophobic

[Farina et al, 17]
 [Craig et al, 18]
 [Di Luzio+Nardi et al, 17]
 [Sokolov+Ringwald, 21] ...
 + Refs in FIPs report [2102.12143]

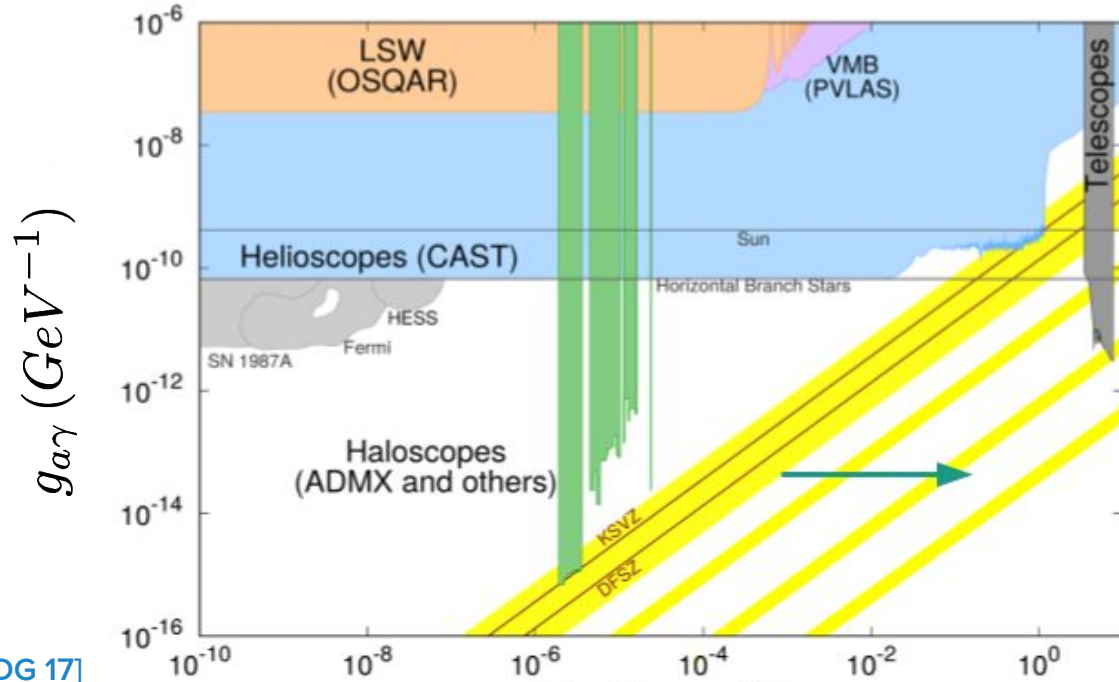


[Ringwald, PDG 17]

$$g_{a\gamma\gamma} = -\frac{1}{2\pi f_a} \alpha_{\text{em}} \left(\frac{E}{N} - 1.92(4) \right)$$

$$m_a^2 f_a^2 \simeq m_\pi^2 f_\pi^2 \frac{m_u m_d}{(m_u + m_d)^2}$$

Heavy axions



[Ringwald, PDG 17]

$$m_a^2 f_a^2 \gg m_\pi^2 f_\pi^2$$

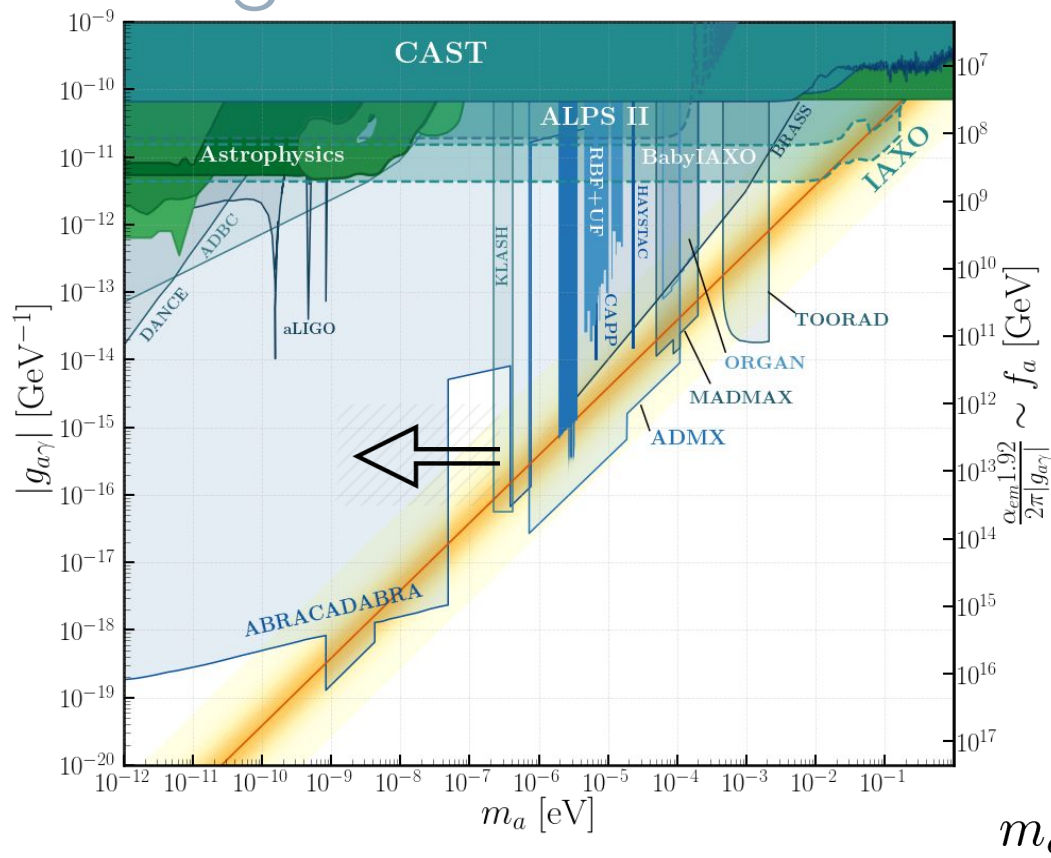


- [Rubakov, 97]
- [Bereziani et al, 01]
- [Fukuda et al, 01]
- [Hsu et al, 04]
- [Hook et al, 14]
- [Chiang et al, 16]
- [Khobadize et al,]
- [Dimopoulos et al, 16]
- [Gherghetta et al, 16]
- [Agrawal et al, 17]
- [Gaillard et al, 18]
- [Fuentes-Martin et al, 19]
- [Csaki et al, 19]
- [Gherghetta et al, 20]

$$m_a \text{ (eV)}$$

Based on **2102.00012** and **2102.01082**

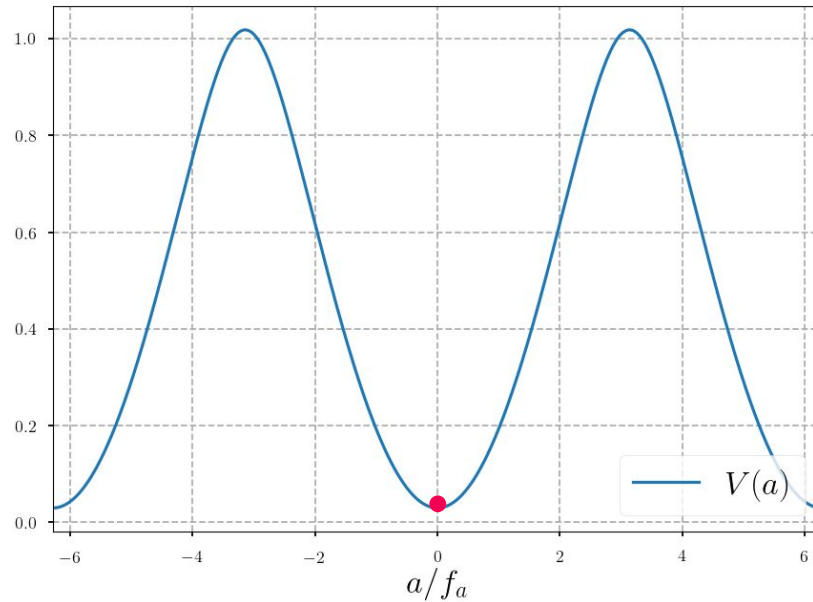
What about lighter axions?



Based on **2102.00012** and **2102.01082**

Axion potential

$$\mathcal{L} \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} \quad \longrightarrow \quad V(a) = -m_\pi^2 f_\pi^2 \sqrt{1 - \frac{4m_u m_d}{(m_u + m_d)^2} \sin^2\left(\frac{a}{2f_a}\right)}$$



[Di Vecchia +Veneziano,80]
[Leutwyler+Smilga, 92]
[di Cortona et al, 15]

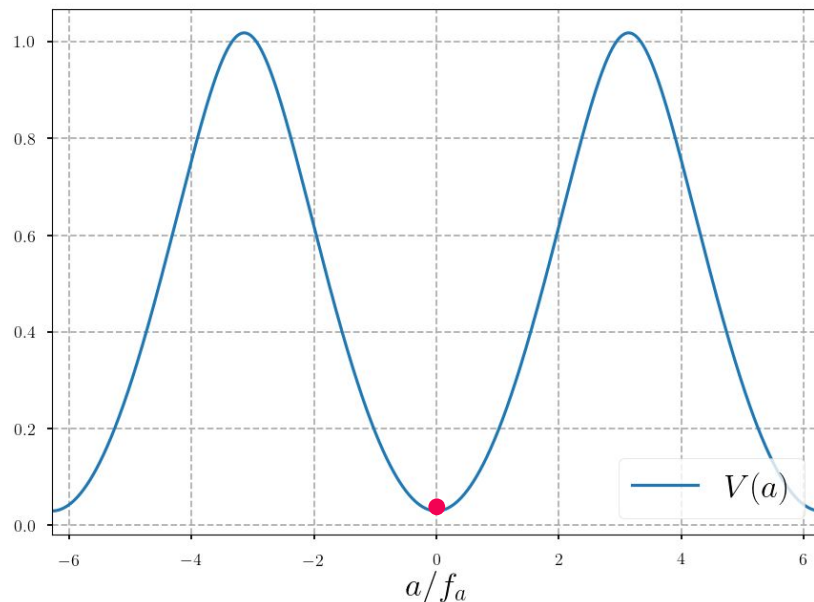
$$\bar{\theta}_{\text{eff}} = \langle \bar{\theta} - \frac{a}{f_a} \rangle = 0$$

Axion potential

$$\mathcal{L} \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G}$$



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[di Cortona et al, 15]

- Alignment
- Cancellation

$$\bar{\theta}_{\text{eff}} = \langle \bar{\theta} - \frac{a}{f_a} \rangle = 0$$

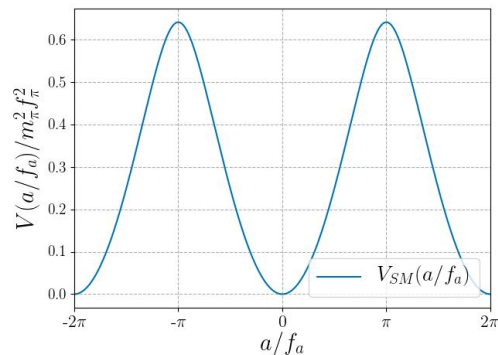
The Z_2 case: Mirror world

$$Z_2 : \quad \text{SM} \longrightarrow \text{SM}'$$
$$a \longrightarrow a + \pi f_a$$

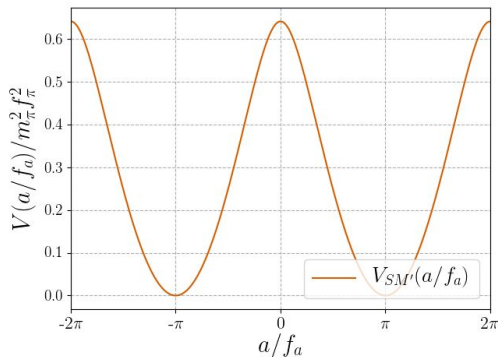
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{SM}'} + \frac{\alpha_s}{8\pi} \left(\frac{a}{f_a} - \theta \right) G\tilde{G} + \frac{\alpha_s}{8\pi} \left(\frac{a}{f_a} - \theta + \pi \right) G'\tilde{G}'$$

What about lighter axions?

SM

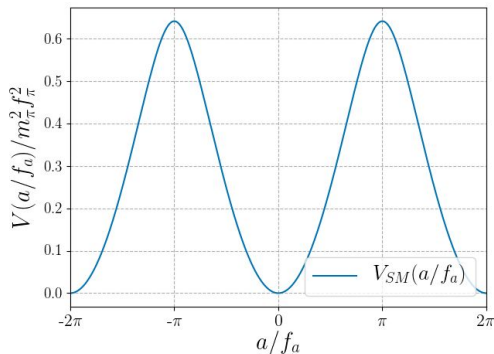


SM'

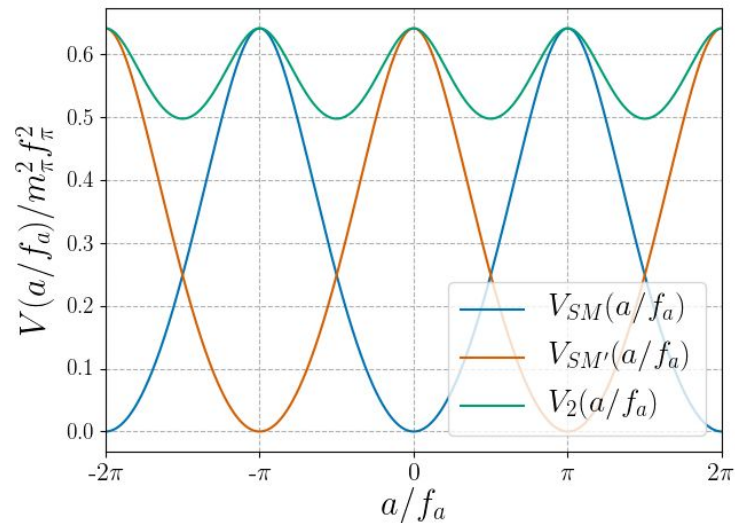
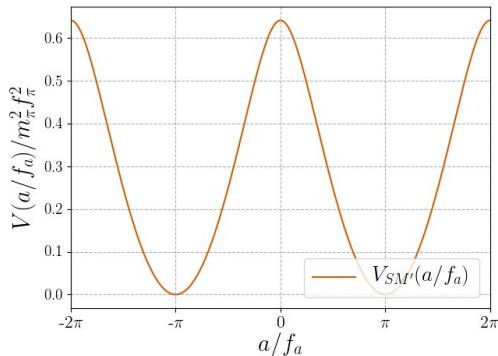


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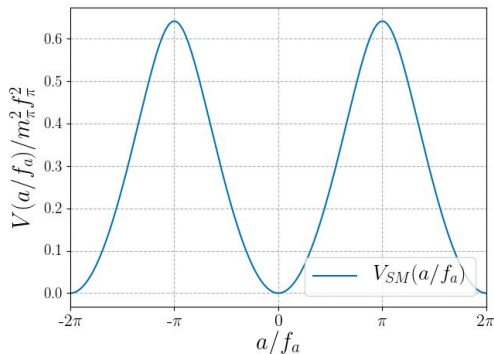


SM'

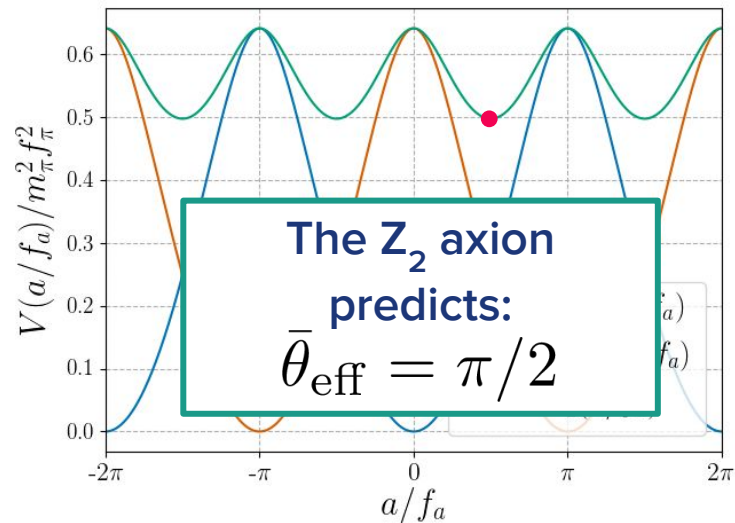
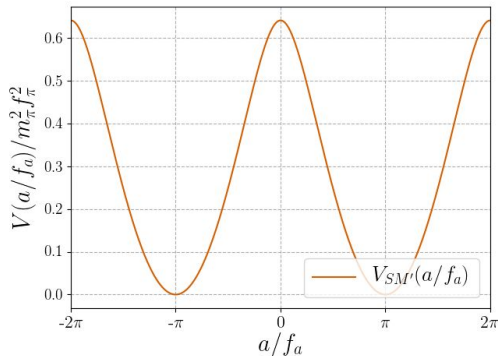


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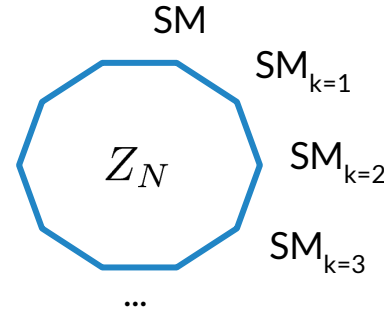
SM'



Z_N axion: N-mirror worlds

[Hook, 18]

$$Z_N : \text{SM} \longrightarrow \text{SM}^k$$
$$a \longrightarrow a + \frac{2\pi k}{N} f_a$$



- The axion realizes the Z_N non-linearly.
- N degenerate worlds with the same couplings as in the SM except for the theta parameter

$$\mathcal{L} = \sum_{k=0}^{N-1} \left[\mathcal{L}_{\text{SM}_k} + \frac{\alpha_s}{8\pi} \left(\theta_a + \frac{2\pi k}{N} \right) G_k \tilde{G}_k \right] + \dots$$

Z_N axion: N-mirror worlds

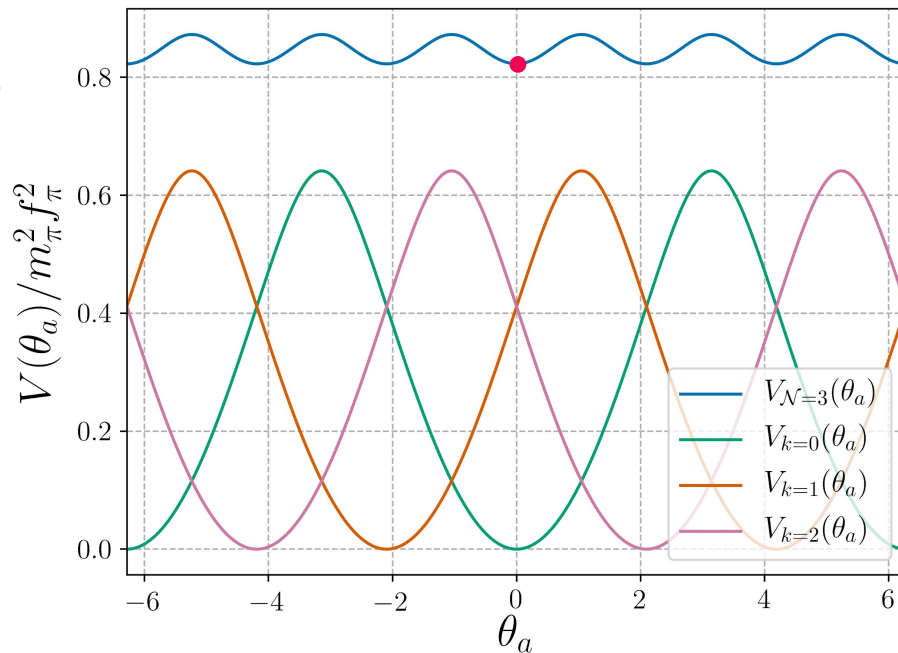
[Hook, 18]

→ N needs to be odd.

Example: Z_3

$$\mathcal{L} = \sum_{k=0}^{N-1} \left[\mathcal{L}_{\text{SM}_k} + \frac{\alpha_s}{8\pi} \left(\theta_a + \frac{2\pi k}{N} \right) G_k \tilde{G}_k \right] + \dots$$

$$\frac{m_a(N)}{m_a(N=1)} \sim \frac{4}{2^{N/2}}$$



Z_N axion: N-mirror worlds

[Hook, 18]

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Example: Z_3

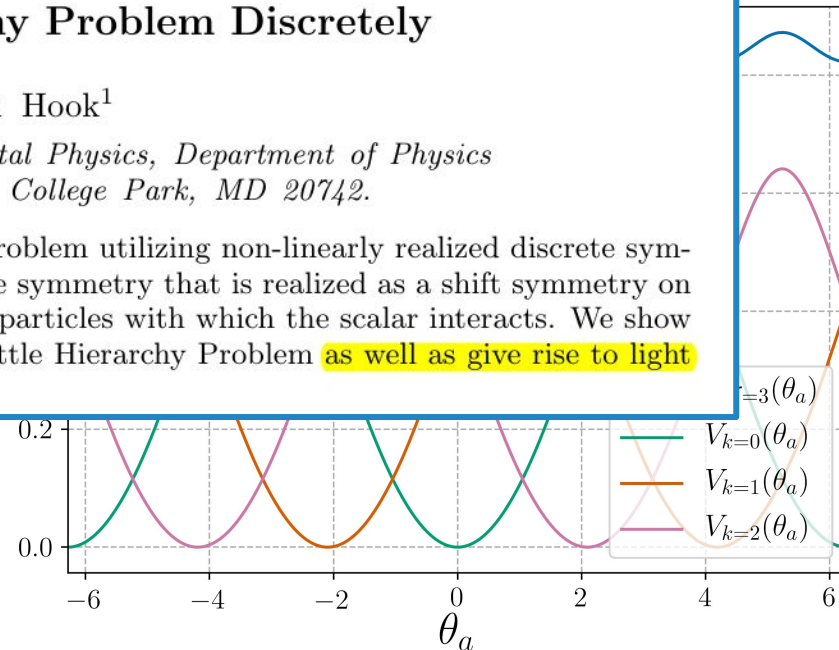
Solving the Hierarchy Problem Discretely

Anson Hook¹

¹Maryland Center for Fundamental Physics, Department of Physics
University of Maryland, College Park, MD 20742.

We present a new solution to the Hierarchy Problem utilizing non-linearly realized discrete symmetries. The cancelations occur due to a discrete symmetry that is realized as a shift symmetry on the scalar and as an exchange symmetry on the particles with which the scalar interacts. We show how this mechanism can be used to solve the Little Hierarchy Problem as well as give rise to light axions.

$$\frac{m_a(N)}{m_a(N=1)} \sim \frac{4}{2^{N/2}}$$



Why exp. suppressed?

$$V_{\mathcal{N}}(a) = - \sum_{k=0}^{N-1} m_{\pi}^2 f_{\pi}^2 \sqrt{1 - \frac{4m_u m_d}{(m_u + m_d)^2} \sin^2 \left(\frac{a}{2f_a} + \frac{\pi k}{N} \right)}$$

→ One would expect:

$$m_a^2 f_a^2 \sim \mathcal{N} m_{\pi}^2 f_{\pi}^2$$

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→ One would expect:

$$\cancel{m_a^2 f_a^2} \approx \cancel{N m_{\pi}^2 f_{\pi}^2}$$

→ Let's understand the cancellation:

$$V_{\mathcal{N}}(\theta_a) = \sum_{k=0}^{N-1} V \left(\theta_a + \frac{2\pi k}{N} \right)$$

\uparrow
 $\theta_a \equiv \frac{a}{f_a}$

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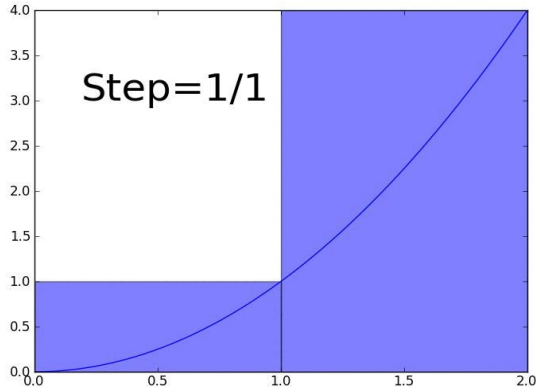
~~$$m_a^2 f_a^2 \sim N m_{\pi}^2 f_{\pi}^2$$~~

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$$\theta_a \equiv \frac{a}{f_a}$$

Source: Wikipedia



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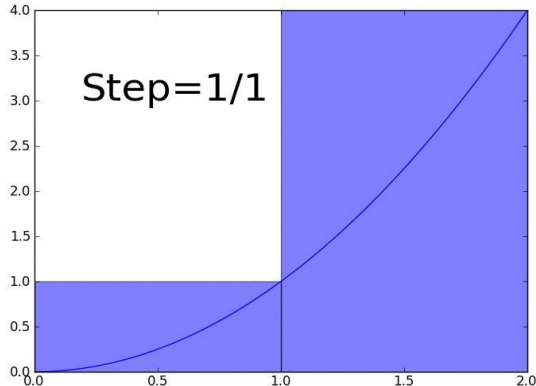
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Does not depend
on the axion!
 $= \text{cte}$

Source: Wikipedia



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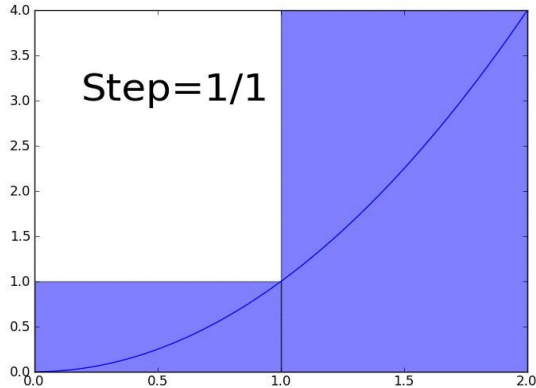
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$$\theta_a \equiv \frac{a}{f_a}$$

Source: Wikipedia



Does not depend
on the axion!
 $= \text{cte}$

The axion potential is
contained in the
subleading terms

Why exp. suppressed?

$$V_{\mathcal{N}}(a) = - \sum_{k=0}^{N-1} m_{\pi}^2 f_{\pi}^2 \sqrt{1 - \frac{4m_u m_d}{(m_u + m_d)^2} \sin^2 \left(\frac{a}{2f_a} + \frac{\pi k}{N} \right)}$$

→ One v

→ Let's

$V_{\mathcal{N}}(\theta)$

The total $Z_{\mathcal{N}}$ axion potential is contained in the error committed in approximating the Riemann sum by an integral:

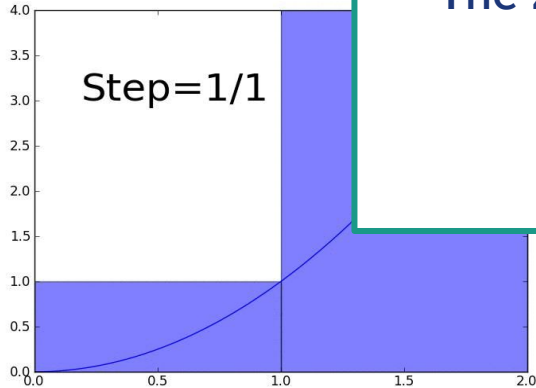
$$E_{\mathcal{N}}(V) = \int_0^{2\pi} V(x) dx - \frac{2\pi}{\mathcal{N}} \sum_{k=0}^{\mathcal{N}-1} V \left(\theta_a + \frac{2\pi k}{\mathcal{N}} \right)$$

The $Z_{\mathcal{N}}$ axion mass is exponentially suppressed:

$$\frac{m_a^2 f_a^2}{m_{\pi}^2 f_{\pi}^2} \propto z^{\mathcal{N}} \equiv \left(\frac{m_u}{m_d} \right)^{\mathcal{N}} \sim 2^{-\mathcal{N}}$$

potential is
in the
terms

Source: Wikipedia



Compact analytical formula

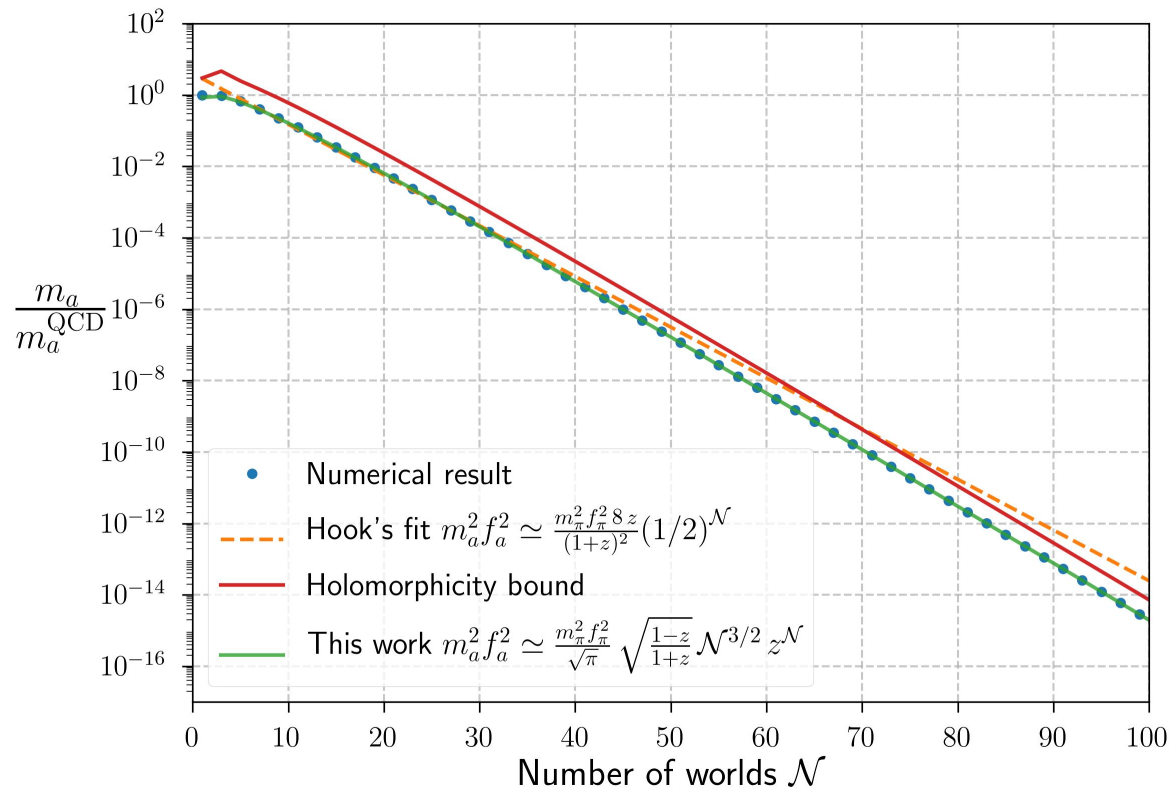
- Using Fourier decomposition and Gauss hypergeometric functions we managed to show that:
- ◆ The total Z_N axion potential approaches a cosine:

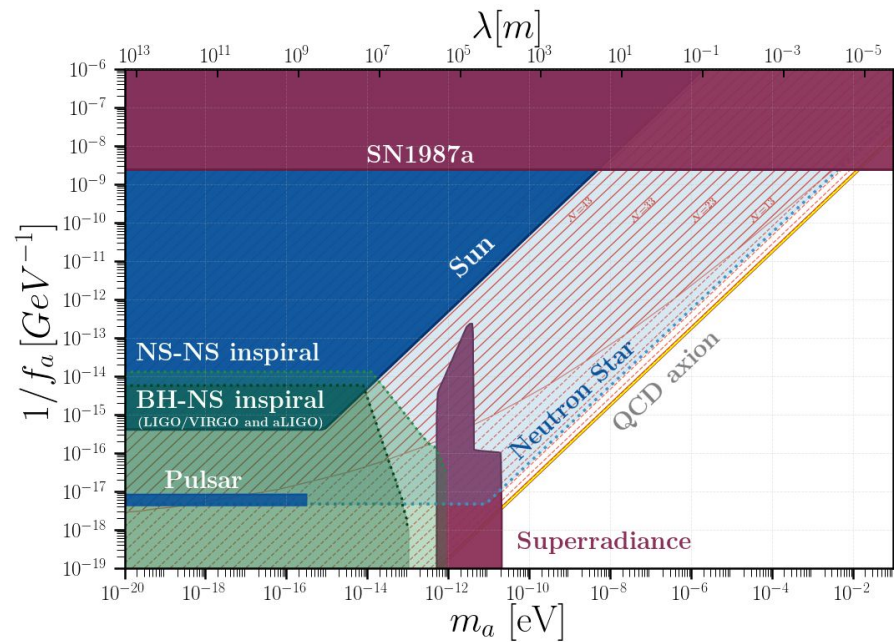
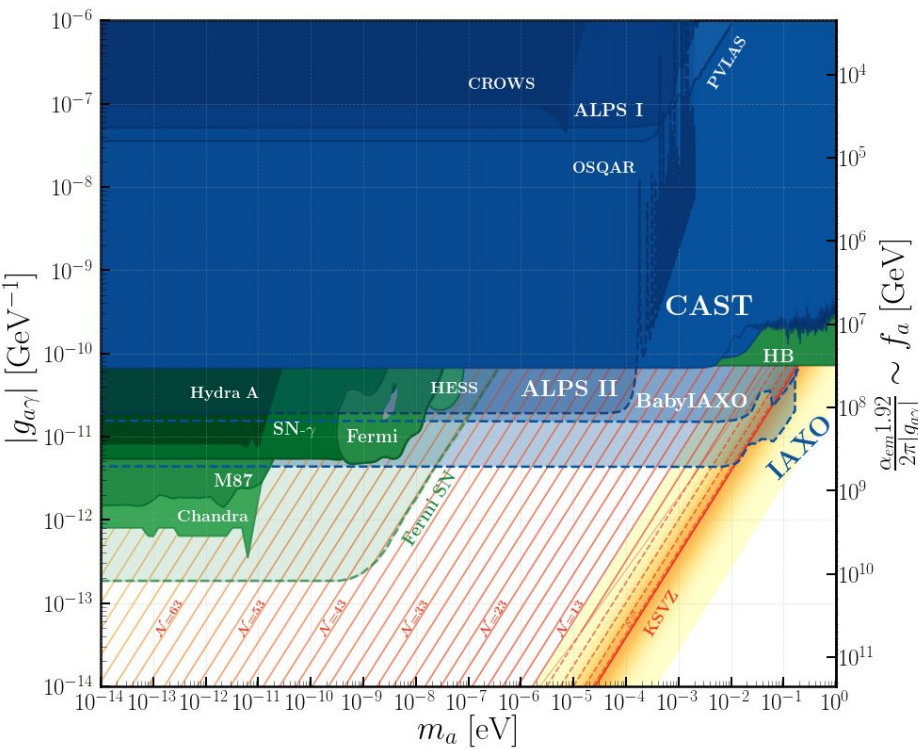
$$V_{\mathcal{N}}(\theta_a) \simeq -\frac{m_a^2 f_a^2}{\mathcal{N}^2} \cos(\mathcal{N}\theta_a)$$

- ◆ Compact analytical formula for the axion mass

$$m_a^2 f_a^2 \simeq \frac{m_\pi^2 f_\pi^2}{\sqrt{\pi}} \sqrt{\frac{1-z}{1+z}} \mathcal{N}^{3/2} z^{\mathcal{N}}$$

Z_N axion mass

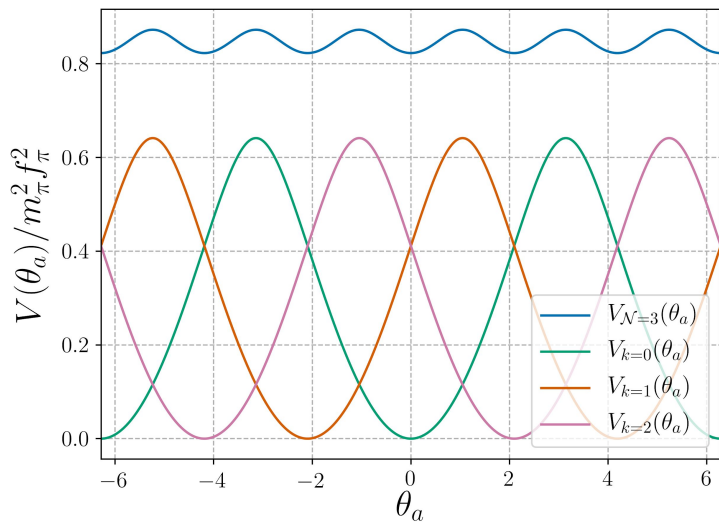




Caveat I

→ There are N minima: we only solve the strong CP with $1/N$ prob

$$\theta_a = \{\pm 2\pi\ell/\mathcal{N}\} \quad \text{for } \ell = 0, 1, \dots, \frac{\mathcal{N}-1}{2},$$



$$\bar{\theta} \lesssim 10^{-10}$$



$1/\mathcal{N}$ probability

Dark matter from the Z_N axion

Trapped misalignment

Mirror world cosmology

→ Mirror worlds need to be colder than SM due to N_{eff} bounds:

-
BBN: $N_{\text{eff}} = 2.89 \pm 0.57$,

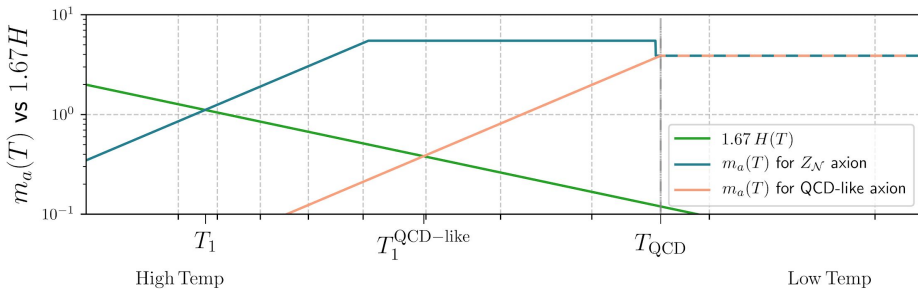
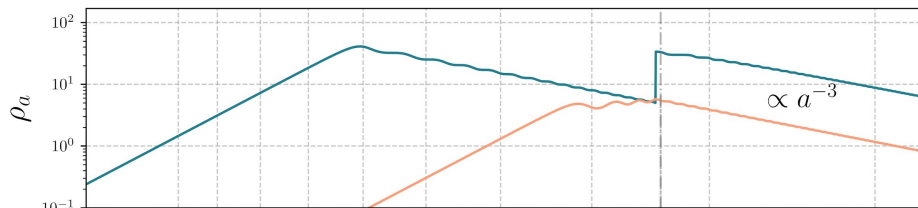
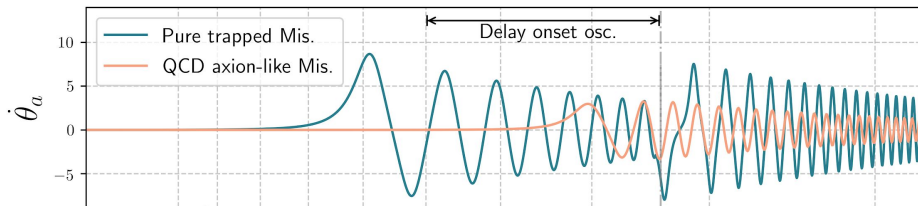
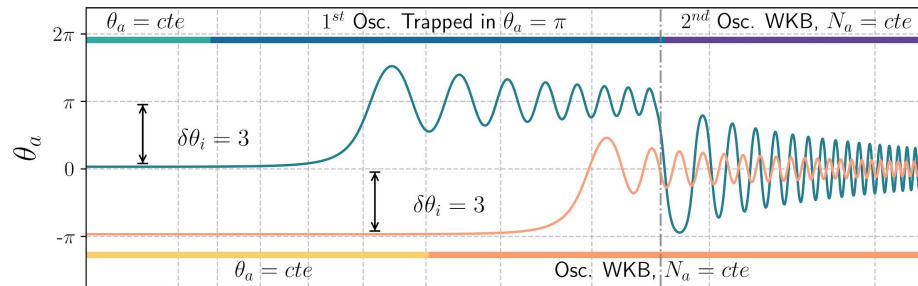
CMB: $N_{\text{eff}} = 2.99^{+0.34}_{-0.33}$.

$$\frac{T'}{T} < \frac{0.51}{(\mathcal{N} - 1)^{1/4}},$$

→ Above $T \geq \Lambda_{\text{QCD}}$ the SM contribution is suppressed which:

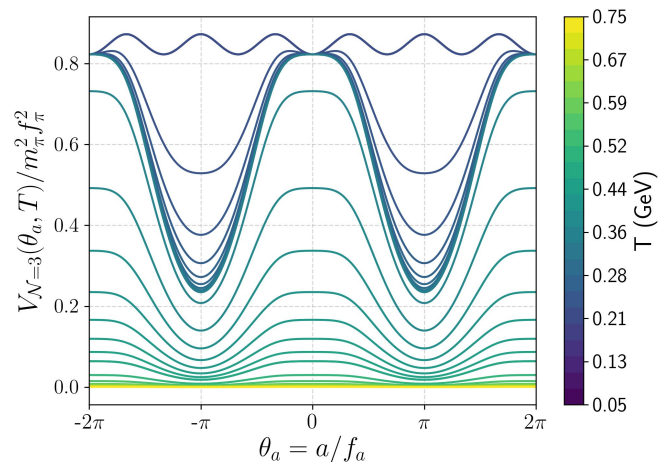
- ◆ Breaks the cancellation
- ◆ The minimum is in π

$$V_{\mathcal{N}}(\theta_a) \simeq -V_{SM}(\theta_a)$$



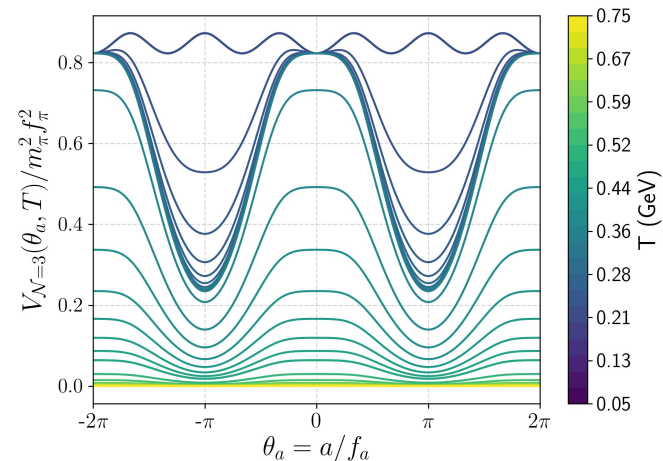
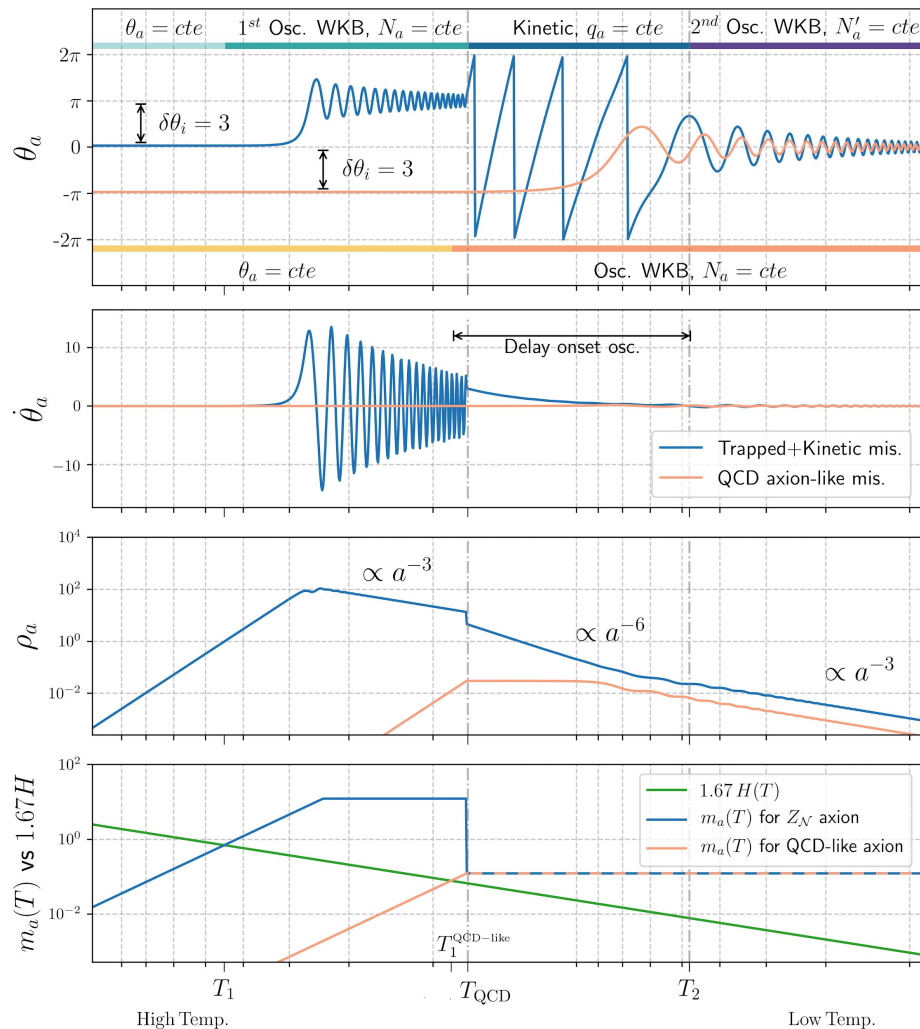
Trapped misalignment mechanism

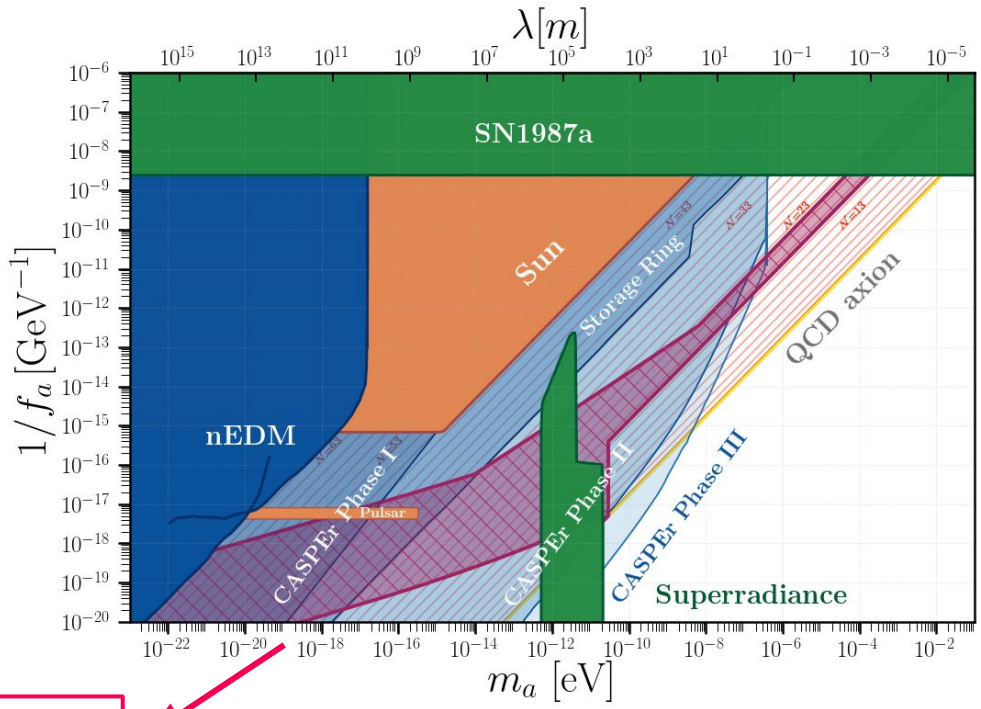
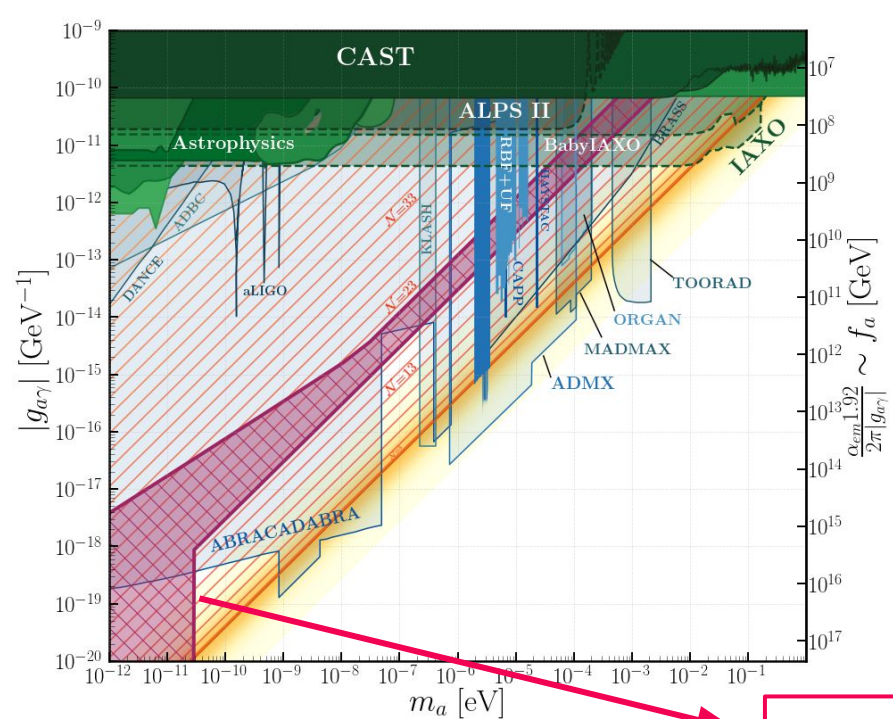
- Compare trapped (blue) with usual misalignment (orange)
- At high temperatures the axion is trapped in the wrong minimum
- The onset of oscillations is delayed
- Less dilution = more dark matter



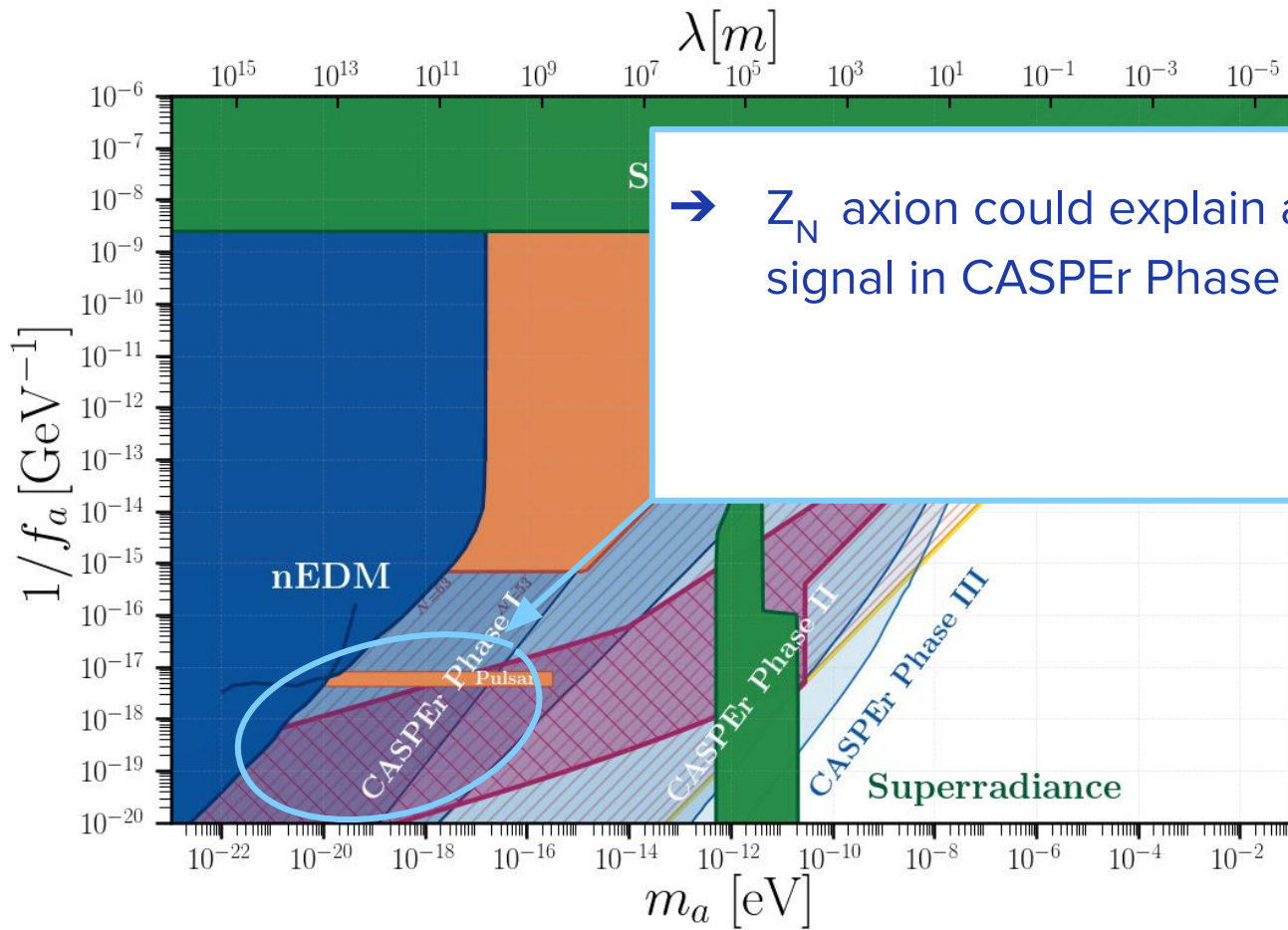
Trapped+kinetic mechanism

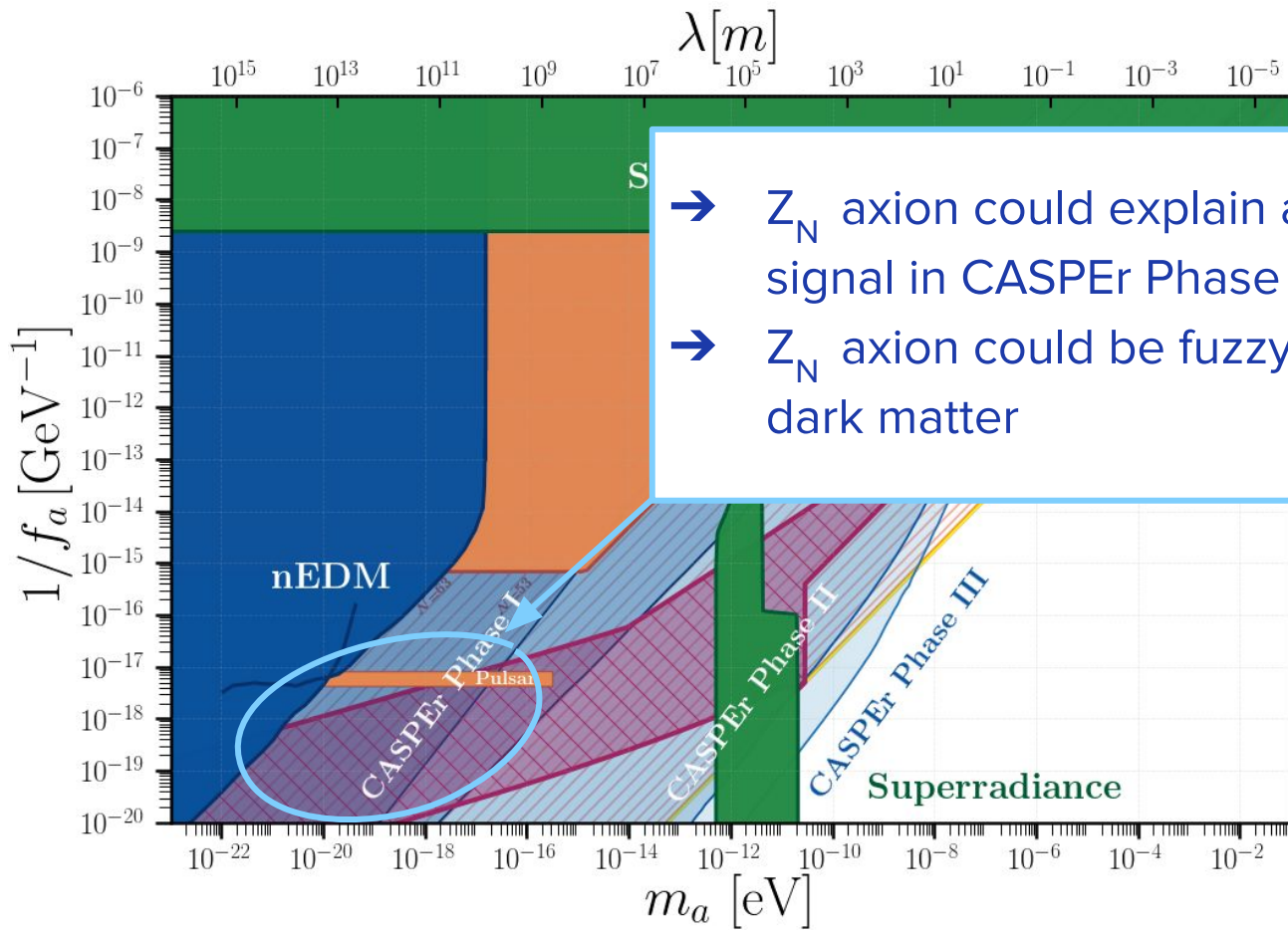
- Compare trapped+kinetic (blue) with usual misalignment (orange)
- After trapping the axion has enough kinetic energy to overcome the barriers
- The onset of oscillations is delayed even further
- Less dilution = more dark matter





Z_N Axion DM





Conclusions I

- ▷ Z_N axion as strong CP solution:
 - The QCD axion can be even lighter
 - UV completions: KSVZ is PQ protected
 - Motivates interesting regions accessible by ALPS II, BabyIAXO, IAXO...

$$m_a^2 f_a^2 \simeq \frac{m_\pi^2 f_\pi^2}{\sqrt{\pi}} \sqrt{\frac{1-z}{1+z}} \mathcal{N}^{3/2} z^{\mathcal{N}}$$

Conclusions II

- ▷ Z_N axion as Dark matter
 - It can explain DM in large regions of $\{m_a, f_a\}$ $3 \leq \mathcal{N} \leq 65$.
 - Novel production mechanism: trapped mis.
 - It can source kinetic mis.
 - First fuzzy dark matter that solves the strong CP
 - First axion model that could explain potential signal in CASPEr Electric Phase I

Caveats and outlook

- ▷ Caveats and outlook
 - N worlds is non-minimal: extra dimensions? strings?
 - Solve the strong CP with $1/N$ prob.
 - Trapped misalignment: only zero mode
 - Non-linearities?
 - Axion fragmentation?
 - Trapped in other scenarios...

Thank you

Pablo Quílez Lasanta

Why exponentially suppressed?

$$z \equiv m_u/m_d$$

$$E_{\mathcal{N}}(V) = \int_0^{2\pi} V(x) dx - \frac{2\pi}{\mathcal{N}} \sum_{k=0}^{\mathcal{N}-1} V\left(\theta_a + \frac{2\pi k}{\mathcal{N}}\right)$$

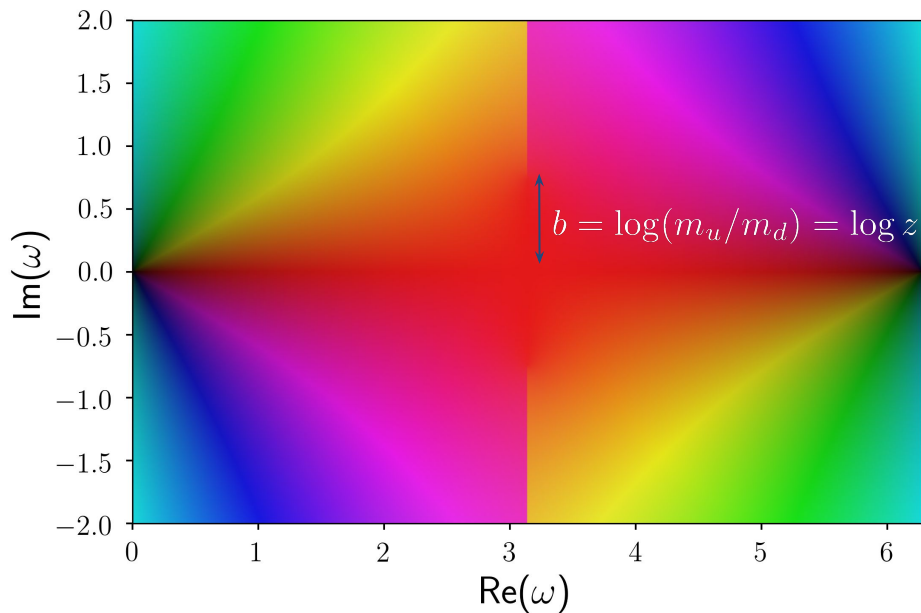
Theorem 9.28 Let $f : \mathbb{R} \rightarrow \mathbb{R}$ be analytic and 2π -periodic. Then there exists a strip $D = \mathbb{R} \times (-b, b) \subset \mathbb{C}$ with $a > 0$ such that f can be extended to a holomorphic and 2π -periodic bounded function $f : D \rightarrow \mathbb{C}$. The error for the rectangular rule can be estimated by

$$|E_{\mathcal{N}}(V)| \leq \frac{4\pi M}{e^{\mathcal{N}b} - 1},$$

where M denotes a bound for the holomorphic function f on D .

$$\frac{m_a^2 f_a^2}{m_{\pi}^2 f_{\pi}^2} \propto z^{\mathcal{N}} \sim 2^{-\mathcal{N}}$$

$$V(\omega) = -m_{\pi}^2 f_{\pi}^2 \sqrt{1 - \frac{4m_u m_d}{(m_u + m_d)^2} \sin^2\left(\frac{\omega}{2}\right)}$$



Can the QCD axion be fuzzy Dark Matter?

→ Fuzzy dark matter: light boson with $m_a \sim 10^{-22}$ eV, $\lambda_c \sim kpc$

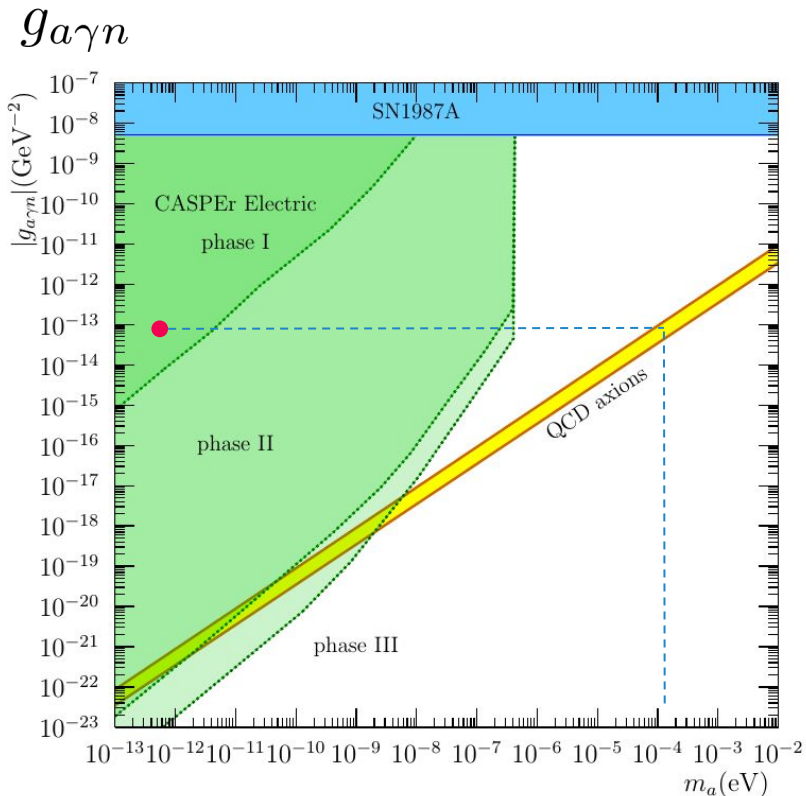
White paper [1904.09003]

$$m_a^2 f_a^2 \simeq m_\pi^2 f_\pi^2 \frac{m_u m_d}{(m_u + m_d)^2}$$

$$m_a \sim 10^{-22} \text{ eV} \implies f_a \sim 10^{28} \text{ GeV} \gg M_{\text{Pl}}$$

**NO, a canonical axion would have
transplanckian decay constant**

Could CASPER Phase I detect an axion?



[Irastorza+Redondo, 18]

Based on **2102.00012** and **2102.01082**

$$\mathcal{L} \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G}$$

$$\delta\mathcal{L} \equiv -\frac{i}{2} \frac{0.011 e}{m_n} \frac{a}{f_a} \bar{n} \sigma_{\mu\nu} \gamma_5 n F^{\mu\nu}$$

$$\equiv g_{a\gamma n}$$

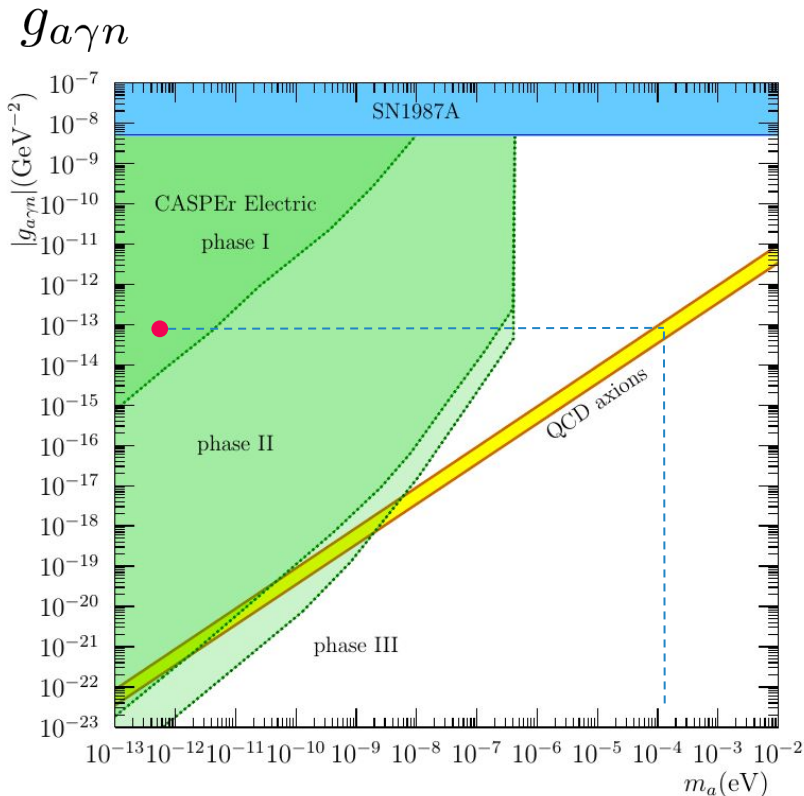
$$m_a^2 f_a^2 \simeq m_\pi^2 f_\pi^2 \frac{m_u m_d}{(m_u + m_d)^2}$$

Coupling to the
nEDM

Axion mass

m_a (eV)

Could CASPER Phase I detect an axion?



[Irastorza+Redondo, 18]

Based on **2102.00012** and **2102.01082**

$$\mathcal{L} \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \dots$$

$$\delta\mathcal{L} \equiv -\frac{i}{2} \frac{0.011 e a}{m_n f_a} \bar{n} \sigma_{\mu\nu} \gamma_5 n F^{\mu\nu} \equiv g_{a\gamma n}$$

$$m_a^2 f_a^2 \simeq m_\pi^2 f_\pi^2 \frac{m_u m_d}{(m_u + m_d)^2}$$

Coupling to the
nEDM

$$m_a^2 f_a^2 \simeq \frac{m_\pi^2 f_\pi^2}{\sqrt{\pi}} \sqrt{\frac{1-z}{1+z}} \mathcal{N}^{3/2} z^{\mathcal{N}}$$

Axion mass

m_a (eV)

Temperature dependence

