

Black Holes and Thermoelectric Transport

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In seeking applications of holography to real materials, thermoelectric conductivities are important observables

$$\begin{pmatrix} J^i \\ Q^i \end{pmatrix} = \begin{pmatrix} \sigma^{ij} & T\alpha^{ij} \\ T\bar{\alpha}^{ij} & T\bar{\kappa}^{ij} \end{pmatrix} \begin{pmatrix} E_j \\ \zeta_j \end{pmatrix}$$

$$\zeta \leftrightarrow -(\nabla T)/T$$

The DC conductivity is particularly interesting.

Naively an IR observable, but depends on UV physics

e.g. CFT on Minkowski space:

Translational invariance \Rightarrow momentum conserved $\Rightarrow \sigma_{DC} = \infty$

More precisely, $Re[\sigma(\omega)] \sim \delta(\omega)$

General framework for dissipating momentum:

Holographic Lattices

CFT with a deformation by an operator that **explicitly** breaks translation invariance.

[Horowitz, Santos, Tong] [.....]

- For example, consider D=4 black holes with bulk scalar field

$$\phi(r, x^i) \rightarrow \frac{\phi_s(x^i)}{r} + \frac{v(x^i)}{r^2} + \dots \quad r \rightarrow \infty$$

Corresponds to a deformation of the CFT:

$$L_{CFT} \rightarrow L_{CFT} + \int dx \phi_s(x) \mathcal{O}(x)$$

Holographic lattices are interesting

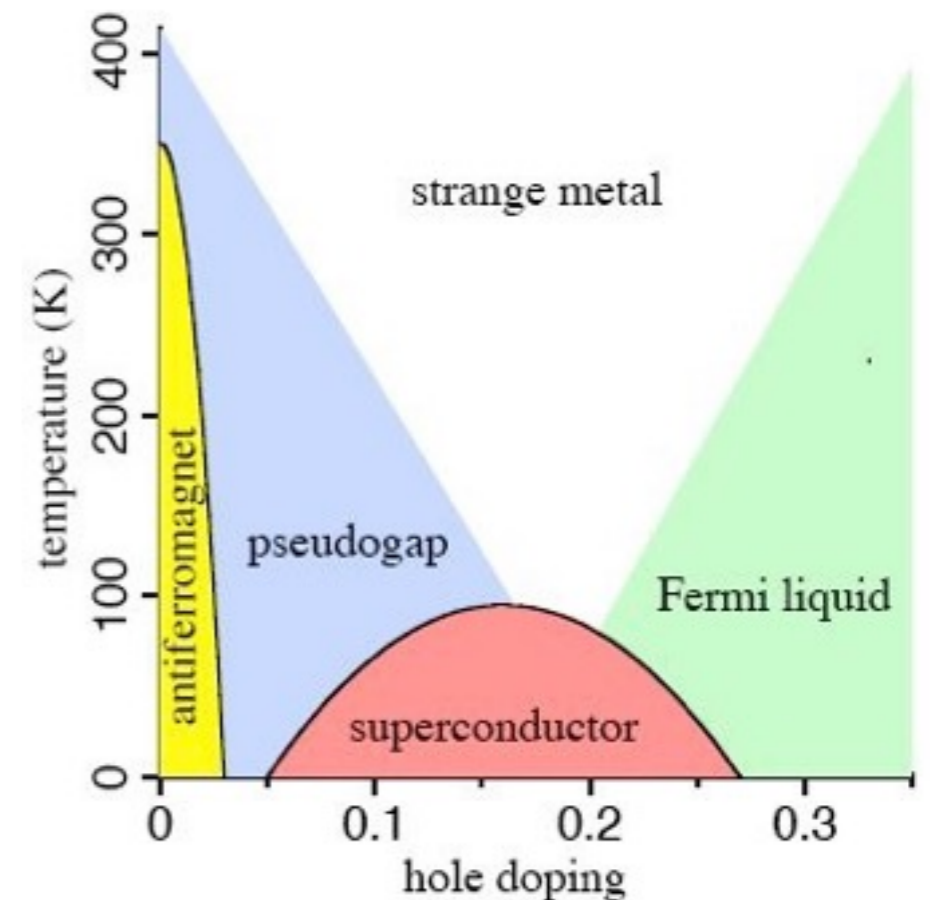
The lattice deformation can lead to:

- ‘Coherent metals’ aka Drude physics - smoothed out $\delta(\omega)$
[Hartnoll,Hoffman][...]

Arises when momentum is nearly conserved.

- Novel ‘incoherent’ metals
[Donos,JPG][Gouteraux][.....]

- Insulators and M-I transitions
[Donos,Hartnoll][Donos,JPG][.....]



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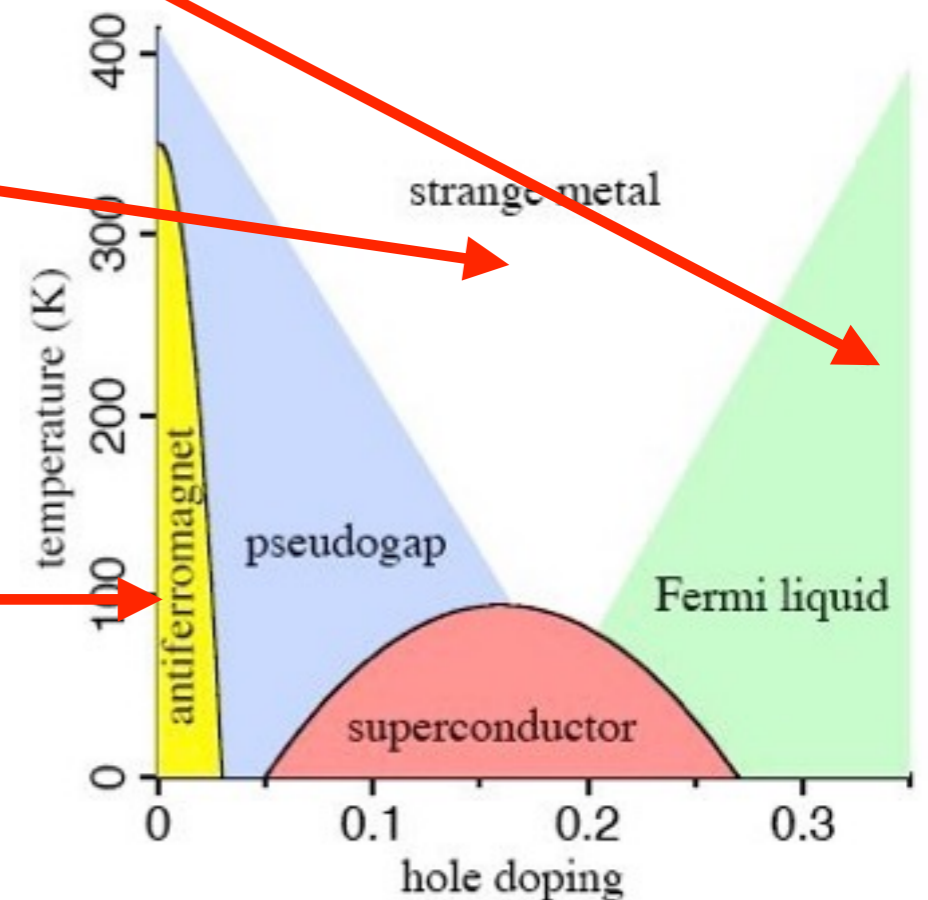
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Plan of talk

- Universal result in holography

DC conductivity matrix is universally obtained exactly by solving generalised linearised **Navier-Stokes equations** for an incompressible auxiliary fluid on the curved black hole horizon

[Donos, JPG]

Hydrostatic fluid equations give **exact** statements about correlation functions of CFT

Exact version of old membrane paradigm for black holes

- Various comments...

Consider hydrodynamic limit of holographic lattices: can extract local currents. Use this to show that **thermal backflow** can occur

DC Conductivity and Navier-Stokes

Illustrate with D=4 Einstein-Maxwell Theory

$$S = \int d^4x \sqrt{-g} \left[R + 6 - \frac{1}{4} F^2 + \dots \right]$$

Holographic lattice: static, planar black hole that preserve time reversal invariance T

- Behaviour at AdS boundary $r \rightarrow \infty$

$$ds^2 \rightarrow r^{-2} dr^2 + r^2 \left[g_{tt}^\infty(x) dt^2 + g_{ij}^\infty(x) dx^i dx^j \right]$$

$$A \rightarrow A_t^\infty(x) dt$$

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i.e. periodic sources for boundary T^{tt} , T^{ij} and J^t

- **Setup:** Perturb the black hole by DC sources E_i, ζ_i

$$\delta(ds^2) = \delta g_{\mu\nu} dx^\mu dx^\nu + 2tg_{tt}\zeta_i dt dx^i$$

$$\delta A = \delta a_\mu dx^\mu - tE_i dx^i + tA_t \zeta_i dx^i$$

- Behaviour at AdS boundary:

The **only** sources are E_i, ζ_i $\zeta \leftrightarrow -(\nabla T)/T$

- Behaviour at the horizon: perturbation is regular
- Naively: now solve full bulk perturbation. However...

- **Result 1:** Zero modes of the electric and heat currents at the horizon are equal the zero modes of those in the CFT

Extends early work of
 [Kovtun, Son, Starinets]
 [Iqbal, Liu]

Illustrate with electric currents:

Bulk equations of motion $\nabla_{\mu} F^{\mu\nu} = 0$

Define the **bulk** electric current density as $J^i = \sqrt{-g} F^{ir}$

$$\partial_i J^i = 0, \quad \partial_r J^i = \partial_j (\sqrt{-g} F^{ji})$$

At AdS boundary, $J^i(x)|_{\infty}$, is **local** current of dual CFT

Observe zero mode $\bar{J}^i = \int d^2x J^i$ is independent of radius

$$\bar{J}^i|_{\infty} = \bar{J}_0^i$$

This, in itself, is **NOT** a
membrane paradigm

- **Result 2:** Obtain local currents $J^i(x)|_0$, $Q^i(x)|_0$ on horizon as functions of E_i , ζ_i and hence DC conductivity of CFT

Use Hamiltonian decomposition of equations of motion with respect to the radial coordinate:

$$\mathcal{H} = N H + N_\mu H^\mu + \Phi C,$$

Evaluate constraints at the “stretched horizon”

Find a decoupled sector for a **subset** of the perturbation which forms a closed set of equations.

Moreover, they give $J^i|_0$ and $Q^i|_0$

[Donos, JPG]

This IS a membrane paradigm

Define

$$(v^i, p, w) \leftrightarrow (\delta g_{it}^{(0)}, \delta g_{rt}^{(0)}, \delta a_t^{(0)})$$

$$Q_0^i = 4\pi T \sqrt{h} v^i$$

$$J_0^i = \sqrt{h} [a_t^{(0)} v^i + h^{ij} (E_j + \partial_j w)]$$

Find:

$$\partial_i Q_0^i = 0$$

$$\partial_i J_0^i = 0$$

$$-2\nabla^i \nabla_{(i} v_{j)} = (4\pi T \zeta_j - \nabla_j p) + a_t^{(0)} (E_j + \nabla_j w)$$

Linear, time-independent, forced **Navier-Stokes equations** for a charged, incompressible fluid on the curved black hole horizon

Comments I

- The formalism can be generalised from Einstein-Maxwell to any theory of gravity in holography

E.g. Scalar fields ϕ give extra viscous terms appear in the Stokes equations

$$\nabla_j \phi^{(0)} \nabla_i \phi^{(0)} v^i$$

[Banks,Donos,JPG]

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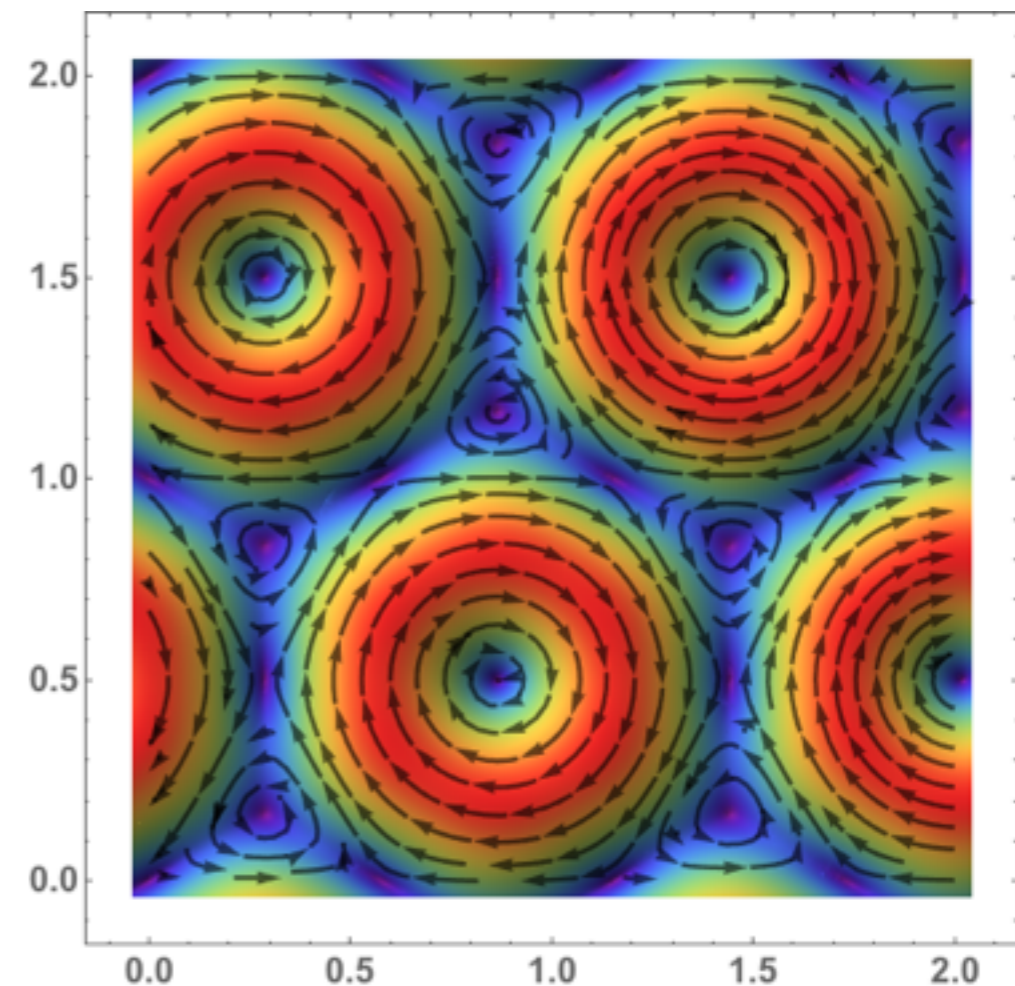
[Banks,Donos,JPG]

- Break T-reversal invariance [Donos,JPG,Griffin,Melgar]

Allow for equilibrium holographic lattices with magnetic fields but also **local magnetisation currents and heat magnetisation currents**

$$J_{\infty}^{(B)i} = \partial_j M^{(B)ij}$$

$$Q_{\infty}^{(B)i} = \partial_j M_T^{(B)ij}$$



Find: extra Lorentz and Coriolis terms in Navier-Stokes eqs

Generalise to higher derivative theories of gravity

- **Result 1:** Zero modes of the electric and heat currents at the horizon are equal the zero modes of those in the CFT

Introduce DC sources in a time independent way

Obtain result by doing a KK reduction on the time direction

- **Result 2:** Obtain local currents $J^i(x)|_0$, $Q^i(x)|_0$ on horizon as functions of E_i , ζ_i and hence DC conductivity of CFT

For the special case of Gauss-Bonnet gravity showed this is still valid and obtain generalised higher derivative Navier-Stokes

For Gauss-Bonnet gravity in D dimensions

$$\nabla_i (\delta_j^i - 4\tilde{\alpha} G_j^i) v^j = 0$$

$$-2\nabla^i (S_{ij}^{kl} \nabla_k v_l) = (\delta_j^i - 4\tilde{\alpha} G_j^i) (4\pi T \zeta_i - \nabla_i p)$$

where

$$S_{ij}^{kl} = [1 - \tilde{\alpha} 2(D-4)(D-1)] \delta_i^{(k} \delta_j^{l)} + \tilde{\alpha} \left[2h_{ij} R^{kl} + 4\delta_{(i}^{(k} R_{j)}^{l)} + 4R_i^{(k} R_j^{l)} \right]$$

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The $\tilde{\alpha}$ corrected shear viscosity

[Brigante, Liu, Myers, Shenker, Yaida]

Comments III: Two interesting perturbative limits

$$\mathcal{L} = \mathcal{L}_{CFT} + h(x)\mathcal{O}(x)$$

Take lattice to be periodic with $x \sim x + 2\pi L$

Typical amplitude h_0

Largest wave number $k_{max} = n/L$

- a) Perturbative lattice: $h_0/T^{d-\Delta} \ll 1$

Coherent metals with Drude peaks

Can prove a generalised Wiedemann-Franz law

- b) Hydrodynamic limit: $k_{max}/T \ll 1$

Fluid-gravity also valid in this limit

b) Hydrodynamic limit $k_{max}/T \ll 1$

Holographic lattices with source for $T^{\mu\nu}$ only

Equivalent to CFT on curved manifold with metric $h_{\mu\nu}$

e.g consider

$$ds^2 = -dt^2 + h_{ij}(x)dx^i dx^j$$

Two simplifications in hydro limit: [Banks,Donos,JPG,Griffin,Melgar]

- Black hole horizon metric is the UV metric h_{ij}
- Solve Navier-Stokes equations using h_{ij} gives **local** heat currents $Q^i(x)$ of CFT

Can also obtain this result from fluid gravity.

In fact this result is valid for non-holographic CFTs too

Thermal backflow for CFTs

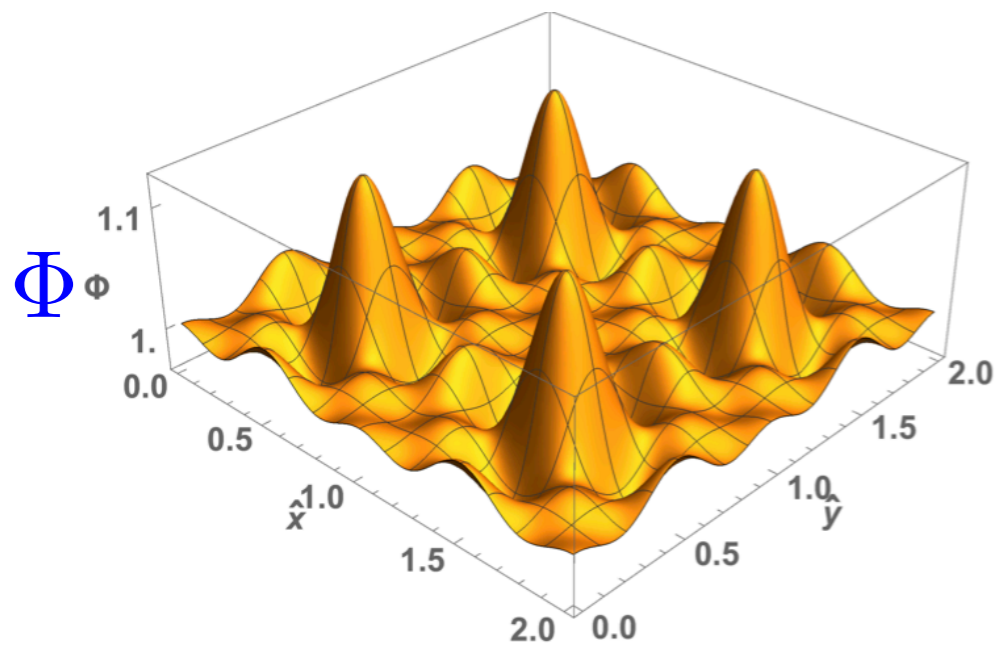
Example: CFT on metric $ds^2 = -dt^2 + \Phi(x)dx^i dx^i$

Equivalent to CFT with periodic isotropic spatial strain

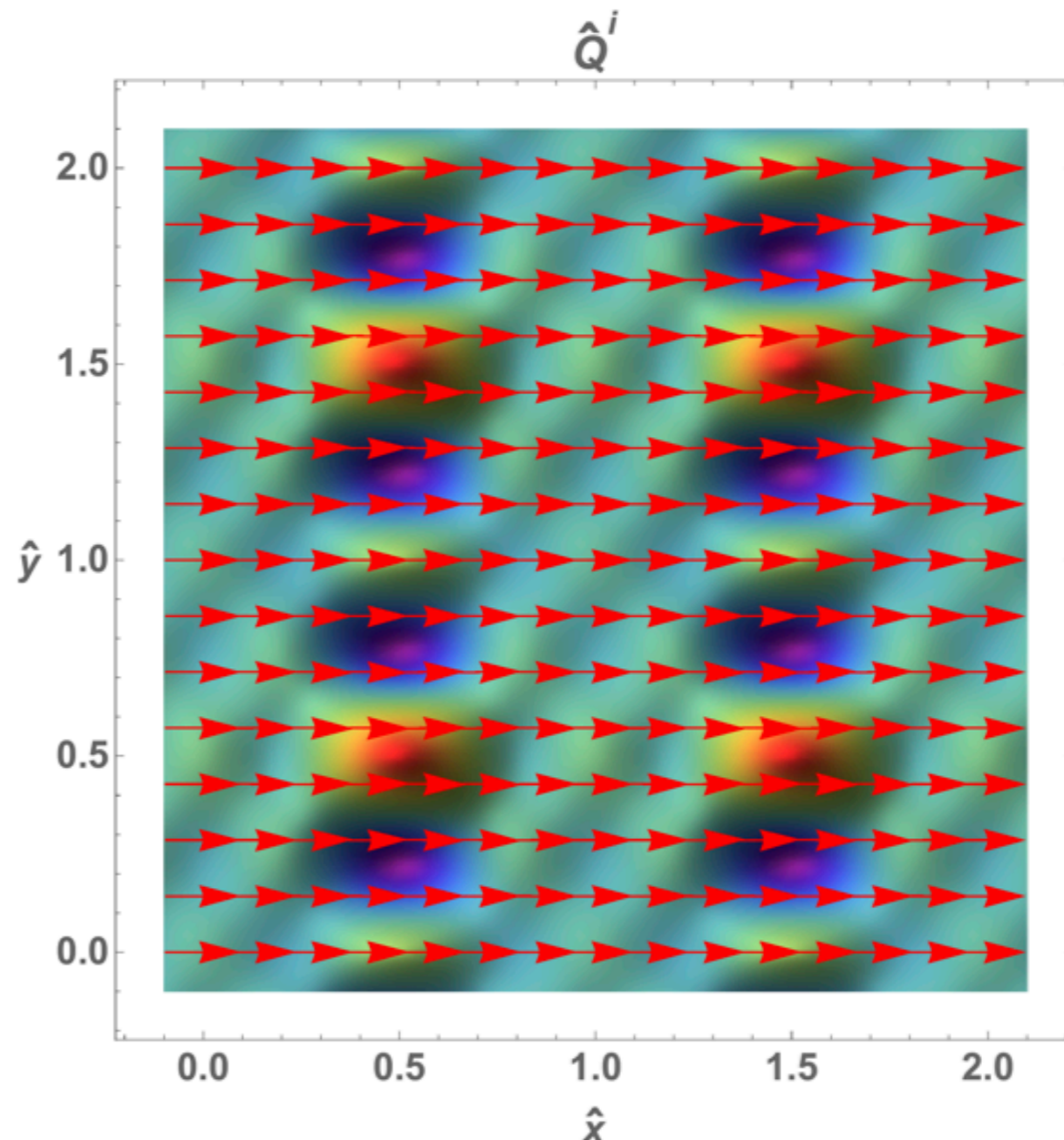
Conformally equivalent to spatially varying temperature

$$ds^2 = -\Phi(x)^{-1}dt^2 + dx^i dx^i$$

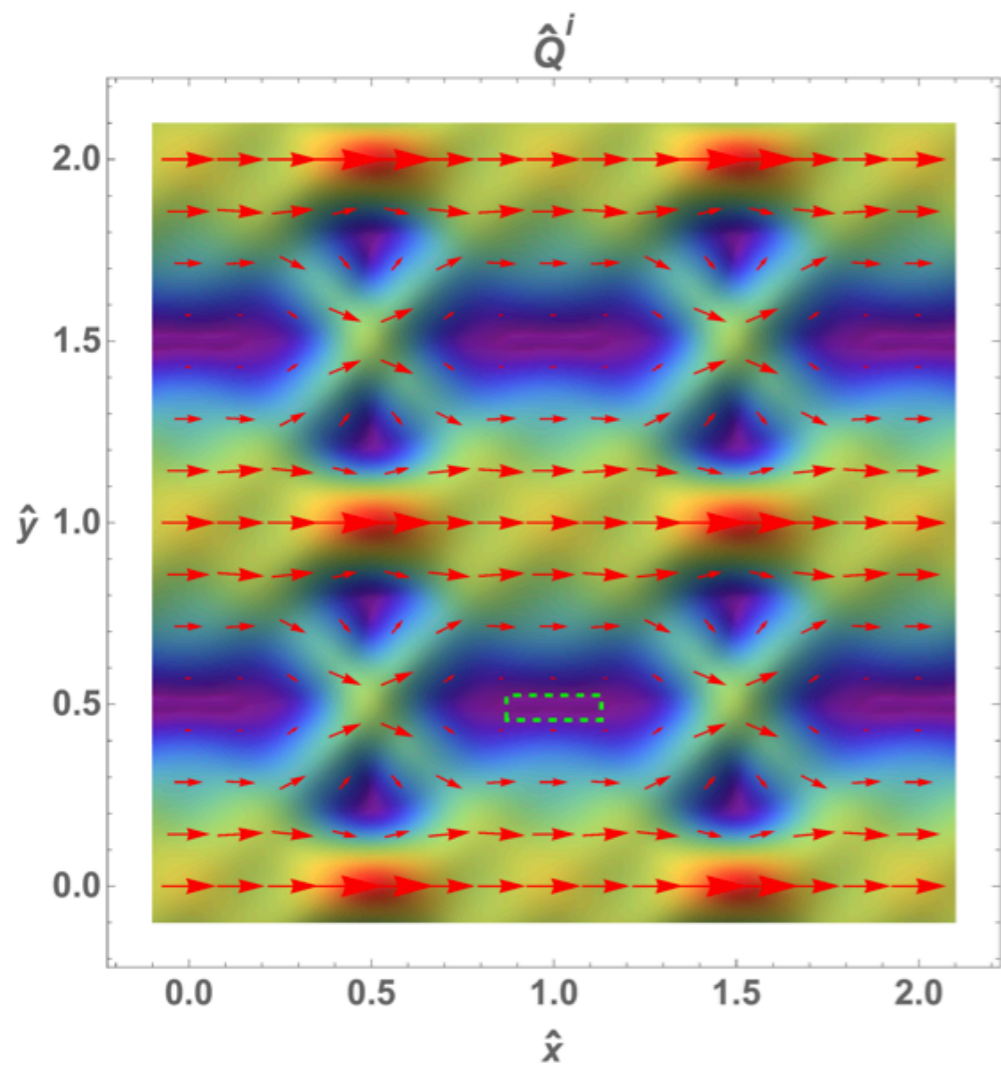
Small amplitude deformation Φ



“Ohmic” response

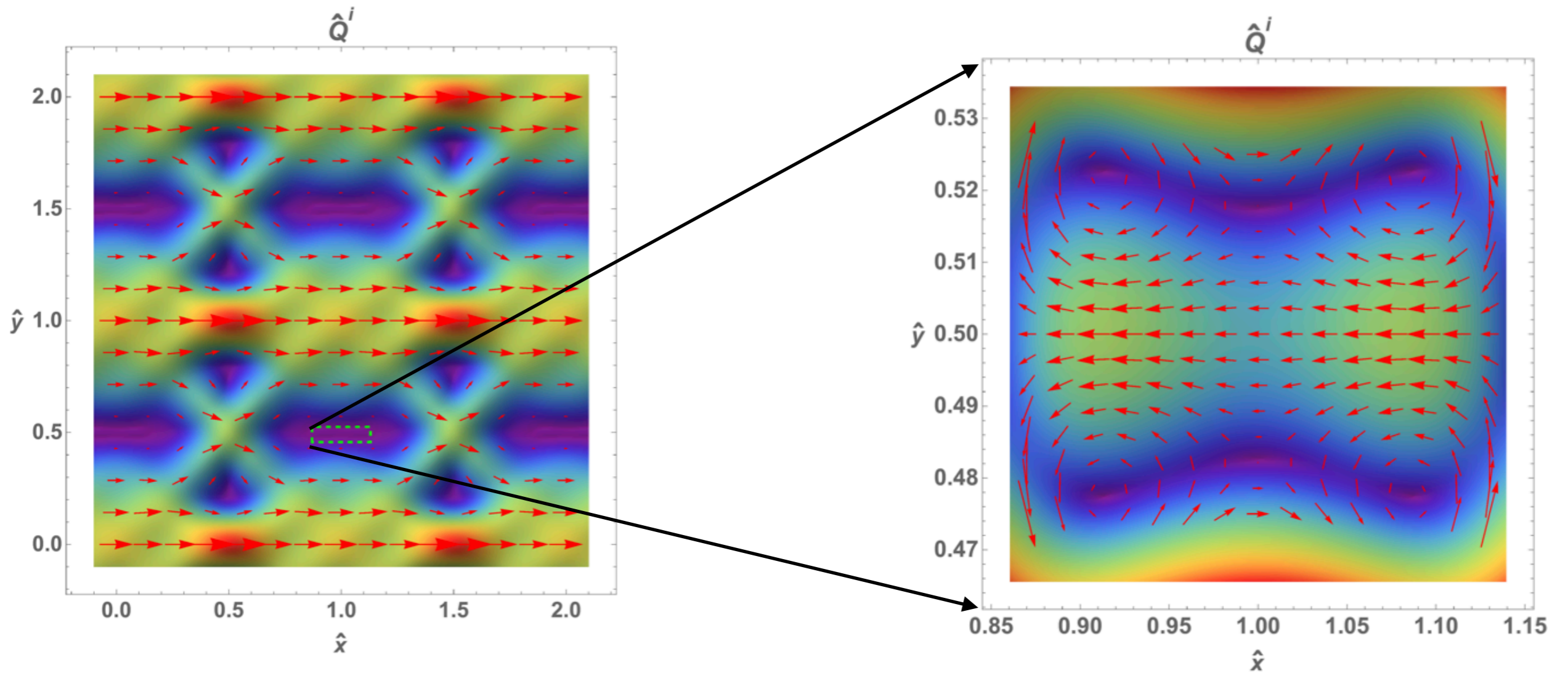


Increase amplitude of Φ



Increase amplitude of Φ

Thermal back flow!



Perhaps observable in graphene?

c.f. electric backflow discussed by

[Levitov, Falkovich] [Bandurin et al] [Moll, Kushwaha, Nandi, Schmidt, MacKenzie]

Final Comments

- Solve Navier-Stokes on black hole horizons
 - Gives, universally, DC thermoelectric conductivities for all T
 - Exact version of membrane paradigm
- Valid for higher derivative gravity
- Solve the Navier-Stokes equations perturbatively:
 - a) Perturbative lattices - small amplitude
 - b) Hydrodynamic limit - long wavelength.
 - Can connect to fluid gravity and hence relate our membrane paradigm to fluid gravity
 - General DC result very obscure in fluid-gravity
 - Predict thermal backflow: observable?