

Double Heavy Baryons and Corrections to Heavy Quark-Diquark Symmetry

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Outline

Introduction

Formalism

Chromo-magnetic coupling of diquarks

Conclusions & Future Directions

Introduction

- **Heavy Mesons ($\bar{Q}q$):**

- Hadrons with single **heavy** quark (anti-quark) and a light anti-quark (quark).
- Described in framework of heavy-quark effective theory (HQET) where operators scale in powers of Λ_{QCD}/m_Q .
- The light quark orbit a static source of color in $\bar{3}$ representation.

- **Double Heavy Baryons (QQq):**

- Hadrons with **two heavy** quarks (anti-quark) and a light quark (anti-quark).
- Hierarchy of energy scales characterized by the **relative velocity v** of heavy quarks:

$$m_Q > m_Q v > m_Q v^2$$

- Described in framework of nonrelativistic QCD (NRQCD) with power counting of operators determined by velocity.

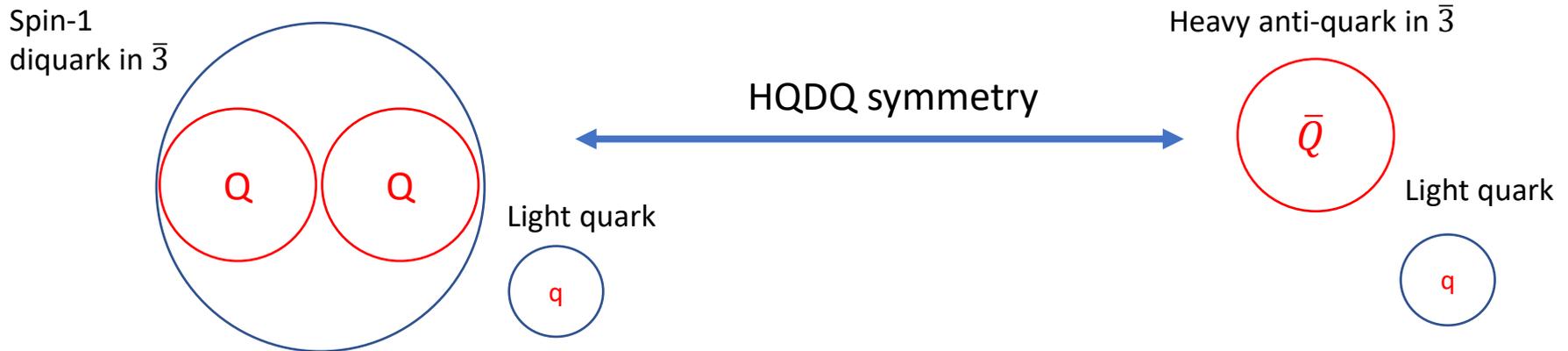
Bodwin, Braaten & Lepage (1994), Luke, Manohar & Rothstein (2000), Brambilla, Pineda & Soto (2000)

- **Attractive** Coulomb force between the heavy quarks.
- Ground state is a **spin-1** diquark in $\bar{3}$ color representation.
- Recently observed by the LHCb experiment $\Xi_{cc}^{++}(3621)$ *LHCb collaboration (2017)*

Introduction

○ Heavy quark-diquark (HQDQ) symmetry

- Relates doubly heavy baryons and heavy anti-mesons.
- In $m_Q \rightarrow \infty$ limit, the light quark (q) in both hadrons orbit around static source of color in $\bar{3}$ representation.



- Relation between hyperfine mass splitting's of lowest lying double heavy baryons and heavy anti-mesons due to $\mathcal{O}(1/m_Q)$ chromo-magnetic coupling operators.

$$m_{\Xi^*} - m_{\Xi} = \frac{3}{4} (m_{D^*} - m_D)$$

*Savage & Wise (1990), Mehen & Fleming (2006)
Brambilla, Vairo & Rosch (2005)*

$$\Xi^*: S = 3/2$$

$$\Xi: S = 1/2$$

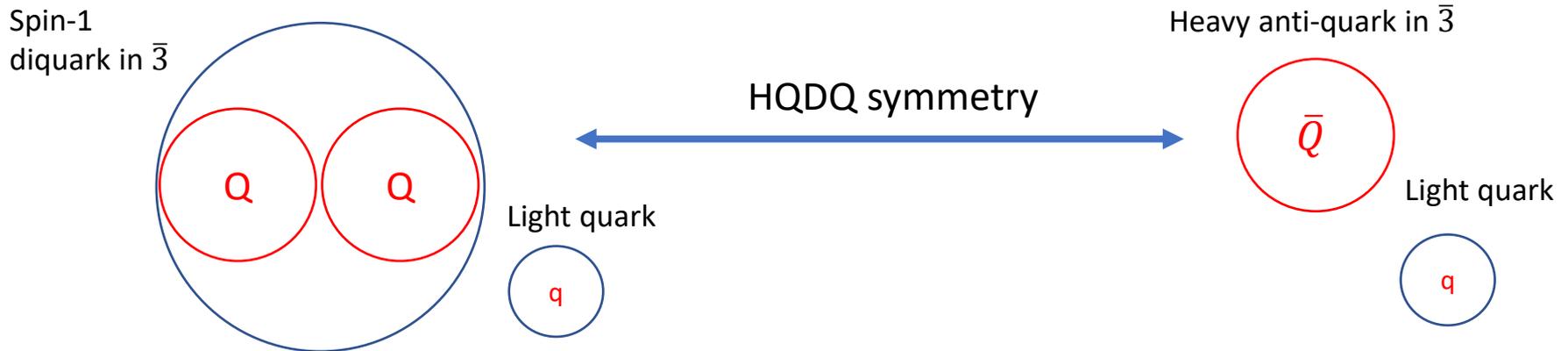
$$D^*: S = 1$$

$$D: S = 0$$

Introduction

- Heavy quark-diquark (HQDQ) symmetry

- Doubly heavy baryons and singly heavy anti-mesons are related by HQDQ symmetry.
- In $m_Q \rightarrow \infty$ limit, the light quark (q) in both hadrons orbit around static source of color in $\bar{3}$ representation.



- Relation between hyperfine mass splitting's of lowest lying double heavy baryons and heavy anti-mesons due to $\mathcal{O}(1/m_Q)$ chromo-magnetic coupling operators.

$$\Xi^*: S = 3/2$$

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$$D^*: S = 1$$

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Corrections to chromo-magnetic coupling of diquarks



Corrections to HQDQ symmetry ??

Formalism

- Interested in calculating the LO and NLO chromo-magnetic coupling of diquarks.
- Use nonrelativistic QCD (NRQCD) formalism with velocity power counting rule for operators.
- NRQCD Lagrangian :

$$\mathcal{L}_0 = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + \sum_{\mathbf{p}} \psi_{\mathbf{p}}^\dagger \left(iD^0 - \frac{(\mathbf{p} - i\mathbf{D})^2}{2m_Q} + \frac{g}{2m_Q} \boldsymbol{\sigma} \cdot \mathbf{B} \right) \psi_{\mathbf{p}} - \frac{1}{2} \sum_{\mathbf{p}, \mathbf{q}} \frac{g_s^2}{(\mathbf{p} - \mathbf{q})^2} \psi_{\mathbf{q}}^\dagger T^A \psi_{\mathbf{p}} \psi_{-\mathbf{q}}^\dagger T^A \psi_{-\mathbf{p}} + \dots ,$$

- Use color and spin Fierz identities to group the heavy quark field ψ in **spin-1, $\bar{3}$ color** and **spin-0, 6 color** representation.
- Composite diquark operators:

Spin-1, $\bar{3}$ color

$$T_{\mathbf{r}}^i = \sum_{\mathbf{p}} e^{i\mathbf{p} \cdot \mathbf{r}} \frac{1}{2} \epsilon^{ijk} (\psi_{-\mathbf{p}})_j \epsilon \sigma (\psi_{\mathbf{p}})_k$$

Spin-0, 6 color

$$\Sigma_{\mathbf{r}}^{(mn)} = \sum_{\mathbf{p}} e^{i\mathbf{p} \cdot \mathbf{r}} \frac{1}{\sqrt{2}} d_{ij}^{(mn)} (\psi_{-\mathbf{p}})_i \epsilon^T (\psi_{\mathbf{p}})_j$$

Formalism

- NRQCD Lagrangian in terms of composite diquark operators :

$$\mathcal{L} = \mathcal{L}_0 + \frac{1}{2} \int d^3\mathbf{r} V^{(\bar{3})}(r) \left(\mathbf{T}_r^{i\dagger} - \sum_{\mathbf{q}} e^{-i\mathbf{q}\cdot\mathbf{r}} \epsilon_{ijk} \frac{1}{2} (\psi_{\mathbf{q}}^\dagger)_j \sigma \epsilon (\psi_{-\mathbf{q}}^\dagger)_k \right) \left(\mathbf{T}_r^i - \sum_{\mathbf{p}} e^{i\mathbf{p}\cdot\mathbf{r}} \frac{1}{2} \epsilon_{ilm} (\psi_{-\mathbf{p}})_l \sigma (\psi_{\mathbf{p}})_m \right) \\ + \frac{1}{2} \int d^3\mathbf{r} V^{(6)}(r) \left(\Sigma_r^{(mn)\dagger} - \sum_{\mathbf{q}} e^{-i\mathbf{q}\cdot\mathbf{r}} \frac{1}{\sqrt{2}} d_{ij}^{(mn)} (\psi_{\mathbf{q}})_i \epsilon (\psi_{-\mathbf{q}})_j \right) \left(\Sigma_r^{(mn)} - \sum_{\mathbf{p}} e^{i\mathbf{p}\cdot\mathbf{r}} \frac{1}{\sqrt{2}} d_{ij}^{(mn)} (\psi_{-\mathbf{p}})_i \epsilon^T (\psi_{\mathbf{p}})_j \right).$$

$$\text{Anti-triplet potential } V^{(\bar{3})} = -\frac{2}{3} \frac{\alpha_S}{r}$$

$$\text{Sextet potential } V^{(6)} = \frac{1}{3} \frac{\alpha_S}{r}$$

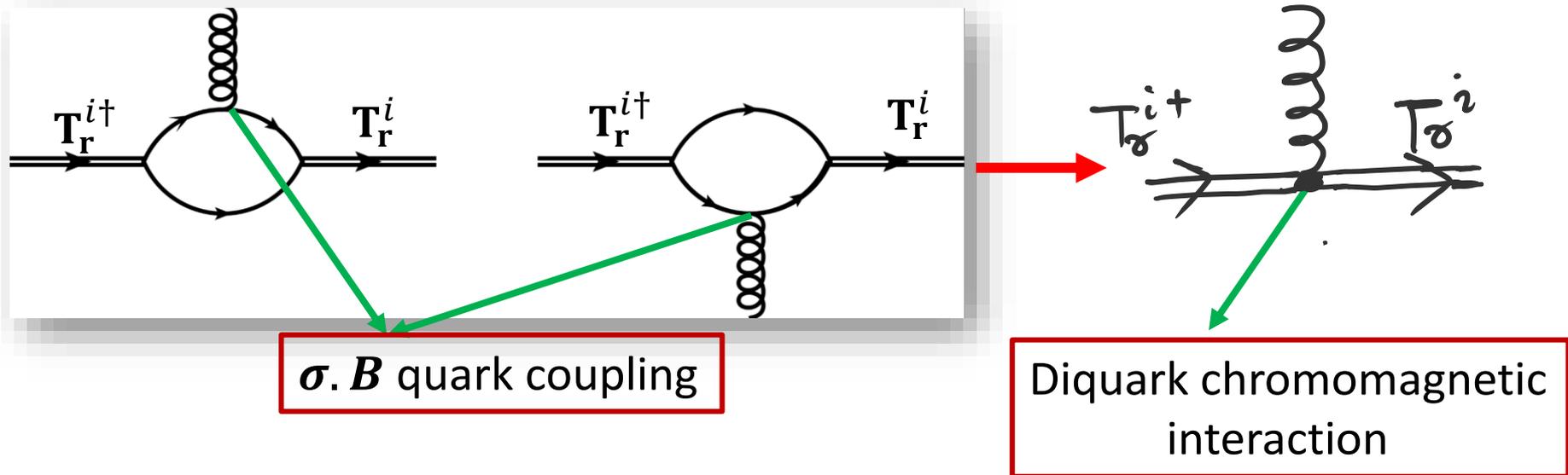
- Coupling of diquark to heavy quarks :

$$\mathbf{T}_r^i \longrightarrow \begin{array}{l} \mathbf{p}, j, \alpha \\ \mathbf{-p}, k, \beta \end{array} \quad -ie^{-i\mathbf{p}\cdot\mathbf{r}} V^{(\bar{3})}(r) \frac{\epsilon_{ijk}}{2} (\sigma \epsilon)_{\alpha\beta}$$

$$\Sigma_r^{(mn)} \longrightarrow \begin{array}{l} \mathbf{p}, i, \alpha \\ \mathbf{-p}, j, \beta \end{array} \quad -ie^{-i\mathbf{p}\cdot\mathbf{r}} V^{(6)}(r) \frac{d_{ijk}^{(mn)}}{\sqrt{2}} \epsilon_{\alpha\beta}$$

Chromo-magnetic coupling

- $\mathcal{O}(v^2)$ diagrams contributing to the chromo-magnetic coupling of diquarks:



- Leading order chromo-magnetic Lagrangian of diquarks :

$$\mathcal{L}_{\sigma \cdot \mathbf{B}} = \frac{g}{2m_Q} \int d^3\mathbf{r} \, i \mathbf{T}_r^{i\dagger} \cdot \mathbf{B}^c \bar{T}_{ij}^c \times \mathbf{T}_r^j.$$

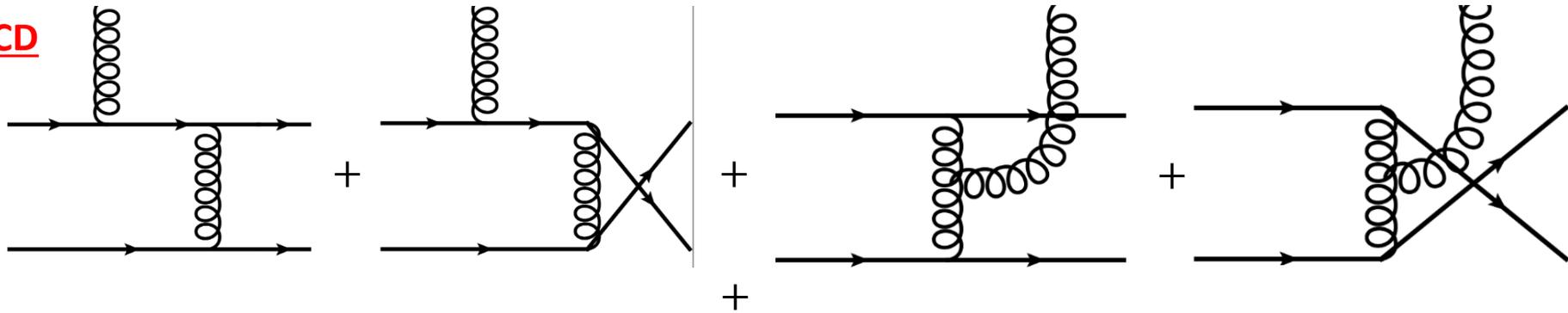
Mehen & Fleming (2006)
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- $\mathcal{O}(1/m_Q)$ correction to heavy spin symmetry.
- Responsible for hyperfine mass splittings in the ground state of doubly heavy baryons.

$1/m_Q \times \alpha_s/r$ correction

- Corrections to the chromo-magnetic coupling of diquarks from **effective five-point** contact operator that couples **4 heavy quark** with a $\sigma \cdot B$ vertex.
- Effective 5-point operator obtained from matching diagrams in QCD and NRQCD.

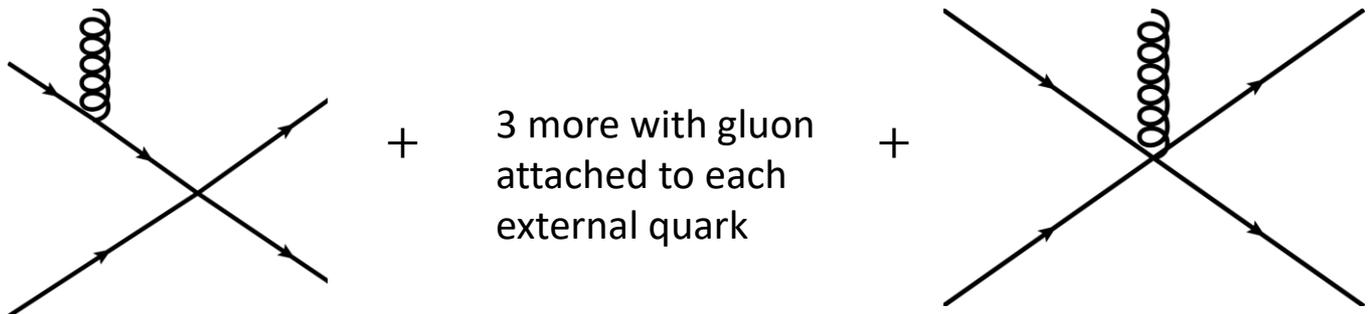
QCD



6 more with gluon attached to each external quark

NRQCD

=



3 more with gluon attached to each external quark

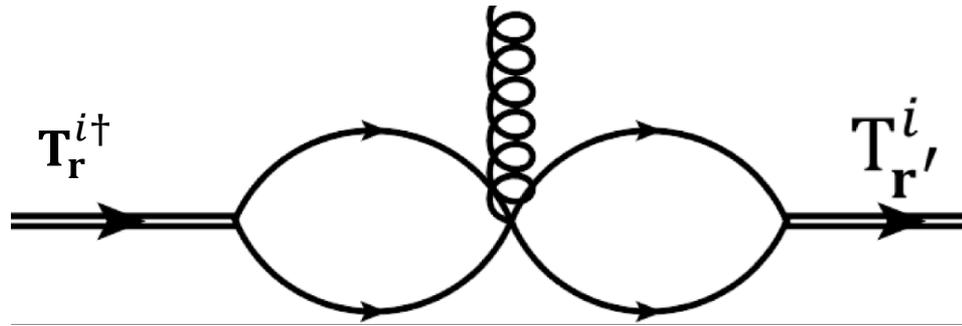
Effective 5-point operator

$1/m_Q \times \alpha_s/r$ correction

- Effective Lagrangian for 5-point contact operator:

$$\mathcal{L}_{eff} = -\frac{g^3}{2} \frac{1}{2m_Q} \sum_{\mathbf{P}_1, \mathbf{P}_2, \mathbf{P}_3, \mathbf{P}_4} \psi_{\mathbf{P}_4}^\dagger T^a \psi_{\mathbf{P}_2} \psi_{\mathbf{P}_3}^\dagger (T^a T^c + T^c T^a) \frac{\boldsymbol{\sigma} \cdot \mathbf{B}^c}{2m_Q} \psi_{\mathbf{P}_1} \frac{1}{(\mathbf{P}_4 - \mathbf{P}_2)^2}$$

- 5-point contact operator contributing to the chromo-magnetic coupling of diquarks:



$\mathcal{O}(v^4)$ diagram
in velocity power
counting of NRQCD

- Corrections to chromo-magnetic coupling of diquarks from 2-loop diagram:

$$\mathcal{L}'_{\sigma \cdot \mathbf{B}} = \frac{ig\alpha_s}{6m_Q^2} \int d^3\mathbf{r} \mathbf{T}_r^{i\dagger} \cdot \frac{1}{r} \mathbf{B}^c \bar{T}_{ij}^c \times \mathbf{T}_r^j$$

- Next-to-leading order (NLO) Lagrangian: $\mathcal{O}(1/m_Q^2 \times \alpha_s/r)$ corrections to heavy spin-symmetry.

$1/m_Q \times \alpha_s/r$ correction

- In $m_Q \rightarrow \infty$, ground state spatial wavefunction of diquarks in doubly heavy baryons:

$$\phi(\mathbf{r}) = \left(\frac{1}{\pi a_0^3} \right)^{1/2} e^{-r/a_0}$$

$$\text{Bohr Radius: } a_0 = 3/\alpha_s m_Q$$

- Chromo-magnetic coupling of diquarks from 2-loop diagram:

$$\mathcal{L}'_{\sigma.B} = \frac{ig\alpha_s}{6m_Q^2} \int d^3\mathbf{r} \mathbf{T}_r^{i\dagger} \cdot \frac{1}{r} \mathbf{B}^c \bar{T}_{ij}^c \times \mathbf{T}_r^j$$

Suppression from Heavy quark mass

Enhancement from the Coloumb exchange

$$\left\langle \frac{1}{r} \right\rangle = \frac{\alpha_s m_Q}{3}$$

$1/m_Q \times \alpha_s/r$ correction

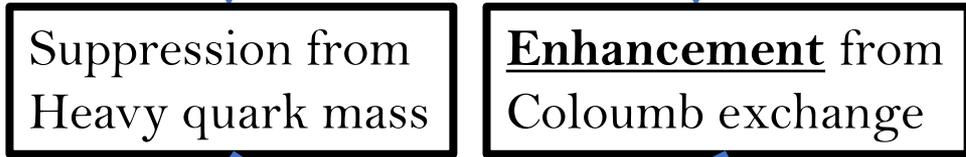
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$\left\langle \frac{1}{r} \right\rangle = \frac{\alpha_s m_Q}{3}$

Overall $\mathcal{O}(\alpha_s^2/m_Q)$ order chromo-magnetic coupling of diquarks !!!

Correction to HQDQ symmetry

- HQDQ symmetry relates mass hyperfine splittings of ground state double heavy baryons and heavy anti-mesons.
- Correction to HQDQ symmetry due to the $\mathcal{O}(\alpha_s^2/m_Q)$ order chromo-magnetic coupling of diquark from effective 5-point operator.
- Relation between mass hyperfine splittings:

Hyperfine splitting
in heavy meson

$$\underbrace{m_{P^*} - m_P}_{\text{Hyperfine splitting in heavy meson}} = \frac{4}{3} (m_{\Sigma^*} - m_{\Sigma}) \left(1 + \frac{\alpha_s}{12m_Q} \left\langle \frac{1}{r} \right\rangle \right),$$
$$= \frac{4}{3} (m_{\Sigma^*} - m_{\Sigma}) \left(1 + \frac{\alpha_s^2}{9} \right)$$

- Correction to hyperfine splitting of diquarks from $\mathcal{O}(\alpha_s^2/m_Q)$ coupling :

Bottom quark $\approx 1.4 \times 10^{-2}$

Charm quark $\approx 3.1 \times 10^{-2}$

Correction to HQDQ symmetry

- HQDQ symmetry relates mass hyperfine splittings double heavy baryons and heavy anti-mesons.
- Correction to HQDQ symmetry due to the $\mathcal{O}(\alpha_s^2/m_Q)$ order chromo-magnetic coupling of diquark from effective 5-point operator.
- Relation between mass hyperfine splittings:

$$\begin{aligned} m_{P^*} - m_P &= \frac{4}{3} (m_{\Sigma^*} - m_{\Sigma}) \left(1 + \frac{\alpha_s}{12m_Q} \left\langle \frac{1}{r} \right\rangle \right), \\ &= \frac{4}{3} (m_{\Sigma^*} - m_{\Sigma}) \left(1 + \frac{\alpha_s^2}{9} \right) \end{aligned}$$

- Correction to hyperfine splitting of diquarks from $\mathcal{O}(\alpha_s^2/m_Q)$ coupling :

Bottom quark $\approx 1.4 \times 10^{-2}$

Charm quark $\approx 3.1 \times 10^{-2}$

Small correction to hyperfine splittings of doubly charm / bottom quark systems !!

Conclusions

- Corrections to HQDQ symmetry from effective **5-point contact operator**.
- Chromo-magnetic coupling of diquark from 5-point contact operator:

$$\mathcal{L}'_{\sigma.B} = \frac{ig\alpha_s}{6m_Q^2} \int d^3\mathbf{r} \mathbf{T}_r^{i\dagger} \cdot \frac{1}{r} \mathbf{B}^c \bar{T}_{ij}^c \times \mathbf{T}_r^j.$$

Naïve expectation: $\mathcal{O}(1/m_Q^2)$ correction

Actually: $\mathcal{O}(\alpha_s^2/m_Q)$ correction due to $\frac{1}{r}$ factor.

- **Small** corrections to HQDQ symmetry from $\mathcal{O}(\alpha_s^2/m_Q)$ coupling :

Bottom quark $\approx 1.4 \times 10^{-2}$ Charm quark $\approx 3.1 \times 10^{-2}$

Future Directions

- RG evolution coupling constant of effective 5-point operator due to one loop ?? .

Thank You!!