

# Pragmatic evaluation of multi-loop multi-scale integrals



Sophia Borowka

CERN



European Research Council

Established by the European Commission

Based on:

1804.06824 SB, T. Gehrmann, D. Hulme,

Comput.Phys.Commun. 225 (2018) 1-9 C. Bogner, SB, T. Hahn,  
G. Heinrich, S. Jahn, S.P. Jones, M. Kerner, A. von Manteuffel, M. Michel, E.  
Panzer, V. Papara

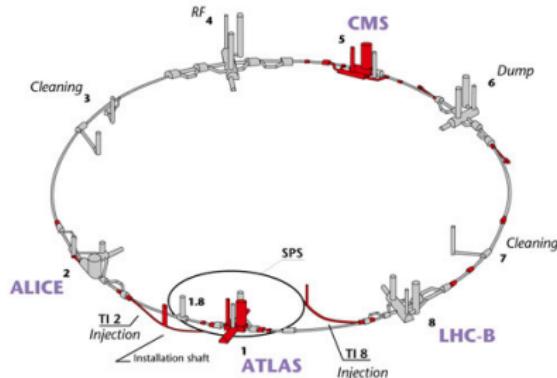
Comput.Phys.Commun. 222 (2018) 313-326 SB, G. Heinrich, S. Jahn,  
S.P. Jones, M. Kerner, J. Schlenk, T. Zirke

Joint Pheno Seminar, University Milano-Bicocca, Jun 7<sup>th</sup> 2018

# Outline

- 0) Introduction**
- 1) Database for loop integrals - Loopedia**
- 2) Analytical vs. numerical evaluation**
- 3) Evaluate integrals**
  - a) fully numerically (PYSECDEC)**
  - b) using precise analytical approximations (TAYINT)**
- 4) Summary & Outlook**

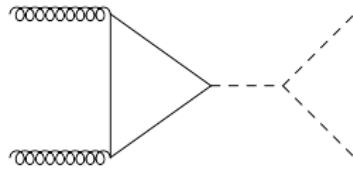
# Higher-order predictions



- ▶ we need accurate predictions at NLO and beyond
- ▶ we want differential predictions
- ▶ we want to take exact heavy quark mass dependences (and H, W, Z) into account
- ▶ phenomenological predictions involve complicated multi-loop multi-scale integrals

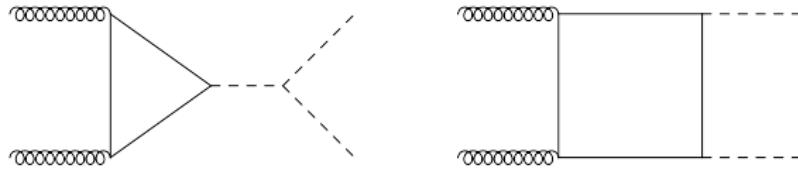
# Higgs-boson pair production in gluon fusion

The leading order is loop induced [Glover, van der Bij '88](#)



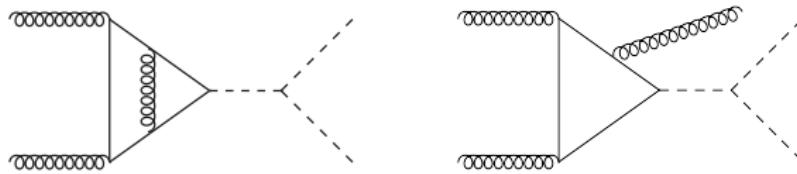
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# $gg \rightarrow hh$ @ NLO with full top mass dependence

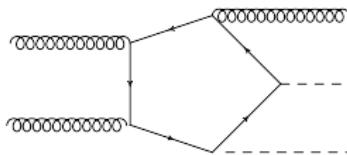
- ▶ Next-to-leading order (NLO) type diagrams



# $gg \rightarrow hh$ @ NLO with full top mass dependence

Higher order correction to box-type LO diagram also needed:

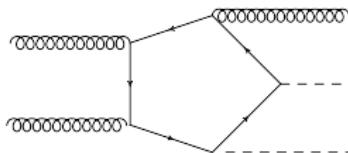
- ▶ one-loop pentagon integrals



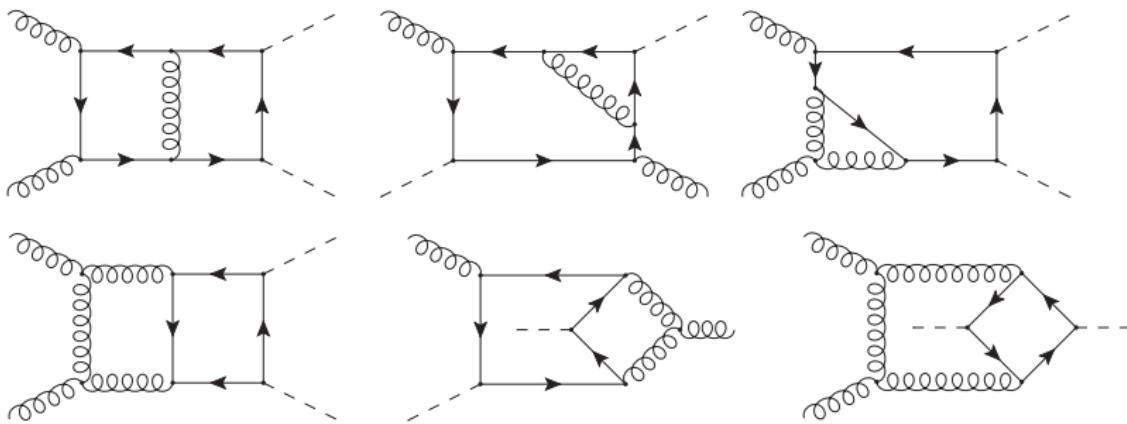
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Higher order correction to box-type LO diagram also needed:

- ▶ one-loop pentagon integrals



- ▶ two-loop integrals with 4 scales  $\hat{s}$ ,  $\hat{t}$ ,  $m_t^2$ ,  $m_h^2$



# Techniques for multi-loop multi-scale integrals

- ▶ Diverse approaches on the market to compute these
  - ▶ Feynman parametrization, Mellin-Barnes representation, differential equations, difference equations, dispersion integrals, integrals in coordinate space, gluing, experimental mathematics, combination of several methods

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- ▶ contribution lookup table: **loop encyclopedia** C. Bogner, SB, T. Hahn, G. Heinrich, S.P. Jones, M. Kerner, A. von Manteuffel, M. Michel, E. Panzer, V. Papara '17

Loopedia

# Loopedia start page

# Loopedia

Ex.: Edge list [(1,2),(2,3),(3,4)] or 1 2 2 3 2 3 3 4 — Nickel index e11|e|

Enter your graph by its edge list (adjacency list) or Nickel index

or browse:

Loops

Legs

Scales

Fulltext must contain:

must not contain:

Search

Reset

If you wish to add a new integral to the database, start by searching for its graph first.

The Loopedia Team is C. Bogner, S. Borowka, T. Hahn, G. Heinrich, S. Jones, M. Kerner, A. von Manteuffel, M. Michel, E. Panzer, V. Papara.  
Software version of 06 Nov 2017 16:01 UTC. In case of technical difficulties with this site please contact [Thomas Hahn](#) or [Viktor Papara](#).

This Web site uses the [GraphState library](#) [[arXiv:1409.8227](#)] for all graph-theoretical operations  
and the neato component of [Graphviz](#) for drawing graphs.

Loopedia is free and open to everyone. To acknowledge and support the work put into keeping Loopedia up to date, please cite [arXiv:1709.01266](#).



# Loopedia Graph Browser

Results for loops = 2, legs = 4, all scales — Row 6»



Prev

Next

Show  rows per page



# Loopedia Record Display

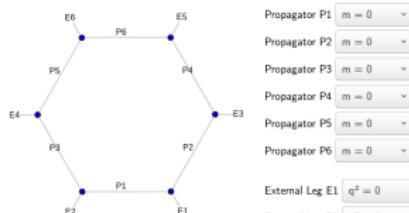
Graph e12|e3|e4|e5|e| — Masses 000|00|00|00|00|0

Edit [ ] Edit [ ] Browse [ ] Home [ ]

Edge list: (e,0|0) (0,1|0) (0,2|0) (e,1|0) (1,3|0) (e,2|0) (2,4|0) (e,3|0) (3,5|0) (e,4|0) (4,5|0) (e,5|0)

Nickel index: e12|e3|e4|e5|e5|e:000|00|00|00|00|0

Database path: 1/6/6/e12|e3|e4|e5|e5|e/8/000|00|00|00|00|0



Propagator P1  $m = 0$   
Propagator P2  $m = 0$   
Propagator P3  $m = 0$   
Propagator P4  $m = 0$   
Propagator P5  $m = 0$   
Propagator P6  $m = 0$

External Leg E1  $q^2 = 0$   
External Leg E2  $q^2 = 0$   
External Leg E3  $q^2 = 0$   
External Leg E4  $q^2 = 0$   
External Leg E5  $q^2 = 0$   
External Leg E6  $q^2 = 0$

Choose Configuration

View public records for this configuration ↓BELOW↓ or choose different configuration ↑ABOVE↑

Reference: arXiv:1104.2787

Authors: Lance J. Dixon, James M. Drummond, Johannes M. Henn

Description: The massless one-loop on-shell hexagon integral is computed by use of a differential equation. The integral is IR- and UV-finite. The result involves classical polylogarithms.

Submitter: bogner@math.hu-berlin.de

Record 1502126180.vng

added 07 Aug 2017 17:16 UTC

last modified 31 Aug 2017 22:39 UTC

Reference: arXiv:1104.2781

Authors: Vittorio Del Duca, Claude Duhr, Vladimir Smirnov

Description: The massless one-loop on-shell hexagon integral is computed in  $D = 6$  by use of Mellin-Barnes and direct integration techniques. The integral is IR- and UV-finite. The result involves

Record 1505802283.T4Pa

added 19 Jun 2017 06:24 UTC



# Analytical vs. numerical evaluation

# Analytic vs. numerical approach

## Analytical:

- + fast evaluation of result

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## Numerical:

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- + scales are not a problem
- precision vs. speed
- intuitive understanding of result harder

# Evaluate integrals fully numerically with pySecDec

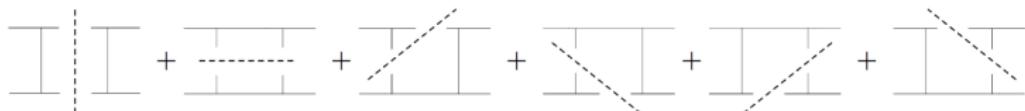
# Feynman loop integrals

- ▶ Scalar multi-loop integral in Feynman parametrization

$$G = \frac{(-1)^{N_\nu}}{\prod_{j=1}^N \Gamma(\nu_j)} \Gamma(N_\nu - LD/2) \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta(1 - \sum_{l=1}^N x_l) \frac{\mathcal{U}^{N_\nu - (L+1)D/2}(\vec{x})}{\mathcal{F}^{N_\nu - LD/2}(\vec{x}, s_{ij})}$$

with  $N_\nu = \sum_{j=1}^N \nu_j$  in  $D$  dimensions with  $L$  loops,  $N$  propagators to power  $\nu_j$

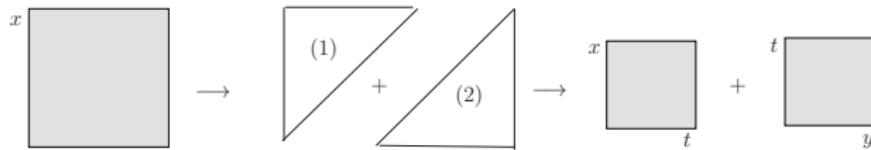
- ▶ Feynman integrals can have numerators of rank  $R$
- ▶  $\mathcal{U}$  and  $\mathcal{F}$  can be constructed via **topological cuts** or by specifying the individual propagators in momentum space



$$\mathcal{F} = -s_{12} x_1 x_3 - s_{23} x_2 x_4 - p_1^2 x_1 x_2 - p_2^2 x_2 x_3 - p_3^2 x_3 x_4 - p_4^2 x_4 x_1 .$$

# The method of sector decomposition

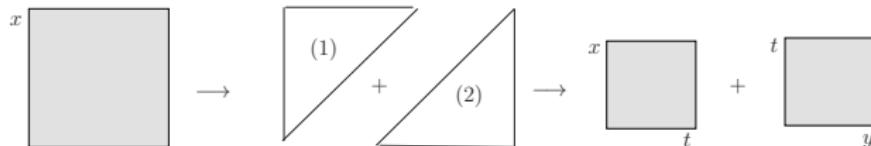
- Idea and method of sector decomposition pioneered by Hepp '66, Denner & Roth '96, Binotto & Heinrich '00



$$\begin{aligned} & \int_0^1 dx_1 \int_0^1 dx_2 \frac{1}{(x_1 + x_2)^{2+\epsilon}} \\ &= \int_0^1 dx_1 \int_0^1 dx_2 \frac{1}{(x_1 + x_2)^{2+\epsilon}} (\theta(x_1 - x_2) + \theta(x_2 - x_1)) \\ &= \int_0^1 dx_1 \int_0^{x_1} dx_2 \frac{1}{(x_1 + x_2)^{2+\epsilon}} + \int_0^1 dx_2 \int_0^{x_2} dx_1 \frac{1}{(x_1 + x_2)^{2+\epsilon}} \\ &= \int_0^1 dx_1 \int_0^1 dt \frac{x_1}{(x_1 + x_1 t)^{2+\epsilon}} + \int_0^1 dx_2 \int_0^1 d\tilde{t} \frac{1}{x_2^{1+\epsilon} (\tilde{t} + 1)^{2+\epsilon}} \end{aligned}$$

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- iterative sector decomposition is highly automatable

# Numerical evaluation using sector decomposition

Public codes:

- ▶ `sector_decomposition` (uses GiNaC) Bogner & Weinzierl '07  
supplemented with `CSectors` Gluza, Kajda, Riemann, Yundin '10  
for construction of integrand in terms of Feynman parameters
- ▶ **FIESTA\*** (uses Mathematica, C) A.V. Smirnov, V.A. Smirnov,  
Tentyukov '08 '09, A.V. Smirnov '13, A.V. Smirnov '15
- ▶ **(PY)SECDEC\*** (uses Python, C++, FORM)  
Carter & Heinrich '10; SB, Carter, Heinrich '12; SB & Heinrich '13;  
SB, Heinrich, Jones, Kerner, Schlenk, Zirke '15; SB, Heinrich, Jahn, Jones,  
Kerner, Schlenk, Zirke '17

\* Multi-scale integrals not limited to the Euclidean region  
SB, J. Carter & G. Heinrich '12; A.V. Smirnov '13

# pySecDec use cases



*Feynman  
integral*

*parametric  
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- ▶ compute single
  - ▶ **multi-loop** multi-scale integrals
  - ▶ general **parametric** functions

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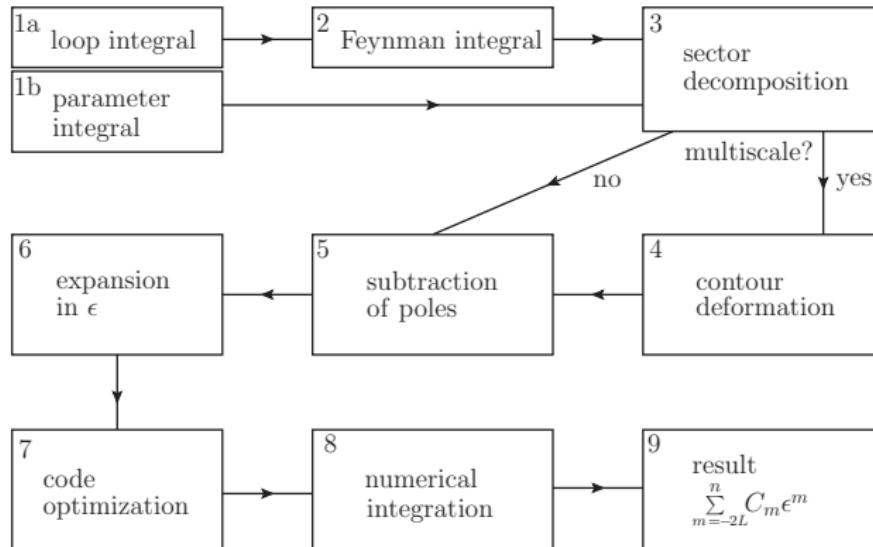


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- ▶ compute single
  - ▶ **multi-loop** multi-scale integrals
  - ▶ general **parametric** functions
- ▶ generate library of integrals
- ▶ use algebra package for symbolic manipulations on integrals

# Operational sequence of the program pySecDec



numerical integration: CUBA library Hahn '04, CQUAD GSL library

# pySecDec toolbox

## **Primary objective:**

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## Improved structure:

- ▶ dependences on open-source libraries only

NUMPY [www.numpy.org](http://www.numpy.org), SYMPY [www.sympy.org](http://www.sympy.org),

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## Download pySecDec:

- ▶ <https://github.com/mppmu/secdec/releases>

The screenshot shows the GitHub repository page for `mppmu/secdec`. The top navigation bar includes links for Features, Business, Explore, Marketplace, Pricing, and a sign-in/sign-up section. Below the header, there are buttons for Code, Issues (0), Pull requests (0), Projects (0), and Insights. The 'Releases' tab is currently selected, showing the latest release, version 1.3.1, which was released on April 24, 2013, with 12 commits since master. The release notes mention 'IPhy released this on Apr 24 · 12 commits to master since this release'. At the bottom of the page, there are links for Accept and Help.

# Install pySecDec

- ▶ **Install:**

```
tar xzvf pySecDec-1.2.2.tar.gz  
cd pySecDec-1.2.2  
make
```

- ▶ **Prerequisites:**

- ▶ python 2.7 or 3
- ▶ python libraries numpy & sympy
- ▶ C++ compiler

- ▶ **Optional prerequisites:**

- ▶ geometric decomposition strategies:  
NORMALIZ Bruns, Ichim, Roemer, Soeger '12
- ▶ NEATO <http://www.graphviz.org>

- ▶ **All other dependences are shipped for easier installation**

# Comparison of pySecDec to SecDec 3 & Fiesta

	PYSECDEC time (s) (algebraic, num.)	SECDEC 3 time (s) (algebraic, num.)	FIESTA time (s) (algeb., num.)
triangle 2L	(40.5, 9.6)	(56.9, 28.5)	(211.4, 10.8)
triangle 3L	(110.1, 0.5)	(131.6, 1.5)	(48.9, 2.5)
box2L6 <sub>(Eucl.)</sub>	(8.2, 0.2)	(4.2, 0.1)	(4.9, 0.04)
box2L6 <sub>(Phys.)</sub>	(21.5, 1.8)	(26.9, 4.5)	(115.3, 4.4)
box2L7 <sub>(invprop)</sub>	(345.7, 2.8)	(150.4, 6.3)	(21.5, 8.8)

4-core Intel(R) Core(TM) i7-4770 CPU @ 3.40GHz

- ▶ relative accuracy goal set to  $10^{-2}$
- ▶ algebraic part usually takes longer, but is only done once in (PY)SECDEC

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→ pySecDec is ready for many more applications! ←

# TayInt - evaluate integrals using precise analytical approximations

# Quasi-finite Basis

Panzer '14; von Manteuffel, Panzer, Schabinger '14

(Scalar) Multi-loop Feynman integral in Feynman parametrization:

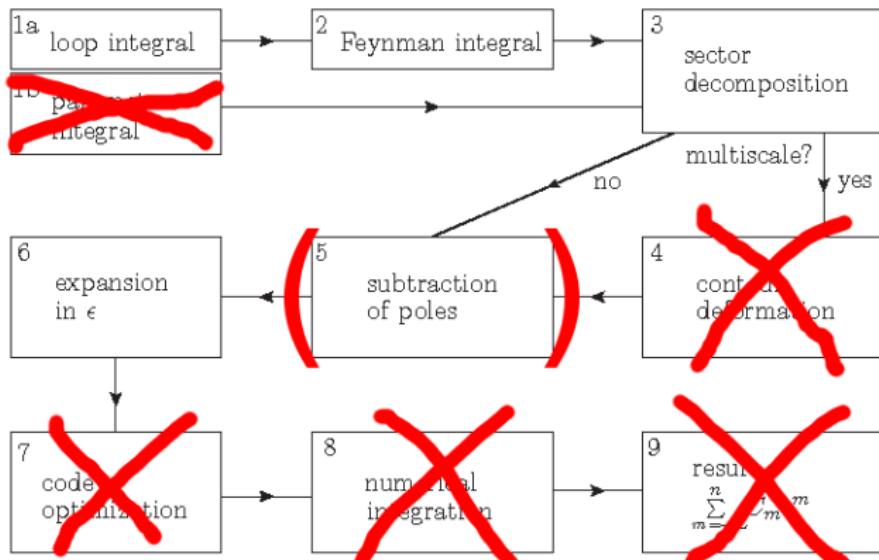
$$G = \frac{(-1)^{N_\nu}}{\prod_{j=1}^N \Gamma(\nu_j)} \Gamma(N_\nu - LD/2) \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta(1 - \sum_{l=1}^N x_l) \frac{\mathcal{U}(\vec{x})^{N_\nu - (L+1)D/2}}{\mathcal{F}(\vec{x}, s_{ij})^{N_\nu - LD/2}}$$

with  $N_\nu = \sum_{j=1}^N \nu_j$  in  $D$  dimensions with  $L$  loops,  $N$  propagators to power  $\nu_j$

- ▶ Appearance of UV and IR divergences depend on  $N_\nu$ ,  $L$  and  $D$
- ▶ Add dots on propagators and/or shift the dimension until the integral is finite
- ▶ facilitates numerical integration due to lack of subtraction terms
- ▶ implemented in Reduze 2 von Manteuffel, Panzer, Schabinger '14

# TayInt uses part of SecDec 3

- ▶ use part of SecDec 3 to generate iterated sub-sectors for finite integrals

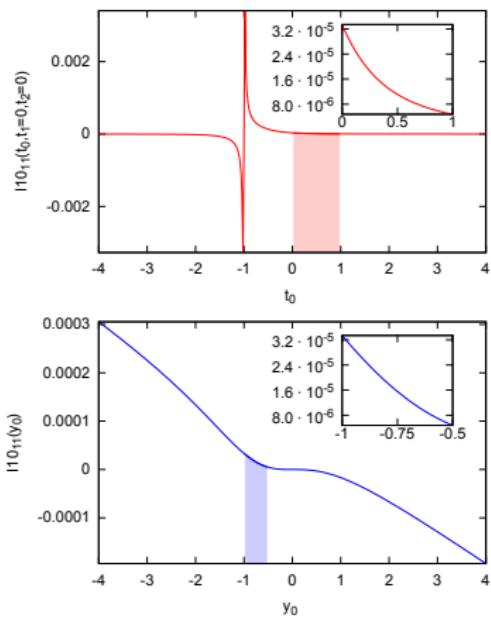


# Conformal mapping

- ▶ conformal mapping

$$t_0 = \frac{-1 - y_0}{y_0} ,$$

- ▶ below threshold: only preparatory step before Taylor expansion
- ▶ above threshold: continuation to complex plane needed



# Analytic continuation and partitioning

- analytic continuation of integrand

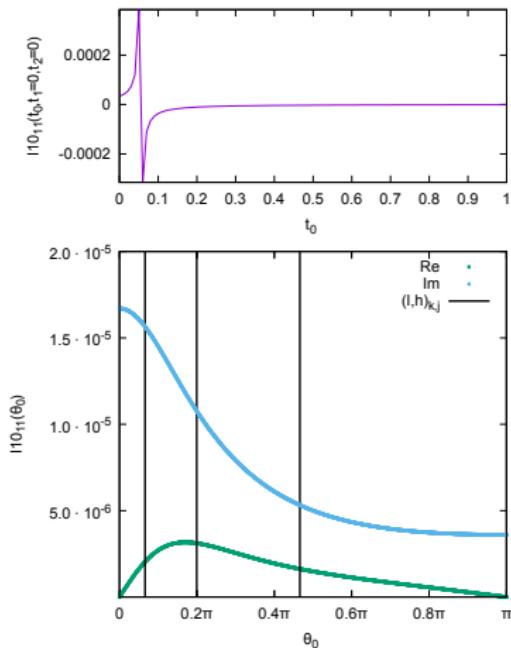
$$\tilde{G}_I^F(t_j) \rightarrow \tilde{G}_I^F(t'_j) = \tilde{G}_I^F\left(\frac{1}{2} + \frac{1}{2} e^{i\theta_j}\right)$$

- partition the integral

$$\prod_j \int_0^{\pm\pi} d\theta_j G_I^F(\theta_j)$$

$$= \prod_j \sum_{k=1}^n \int_{l_{k,j}}^{h_{k,j}} d\theta_j G_I^F(\theta_j) \quad (1)$$

with  $h_{n,j} = \pm\pi$  and  $l_{1,j} = 0$ ,



# Choice of integration contour

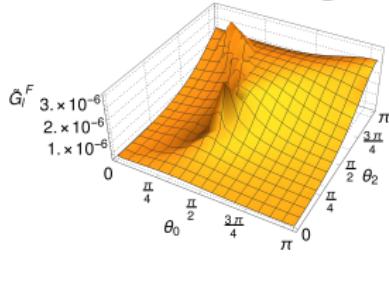
- ▶ sample over  $A$  values for  $\theta_j$ ,  $\gamma$  sets of kinematic values for all contour configurations  $o(j) = \pm\pi$

$$\bar{m}_I(\Theta_{o(1), \dots, o(J)}^A) = \frac{1}{A} \sum_{a=1}^A \frac{1}{\gamma} \sum_{i=1}^{\gamma} \text{Abs} \left[ \frac{1}{J} \sum_{j=1}^J \left( \frac{\partial}{\partial \theta_j} \tilde{G}_I^F(\theta_j, \mathcal{K}_i) \right) \right] \Big|_{\{\theta_j\} \rightarrow \Theta_{o(1), \dots, o(J)}^a}$$

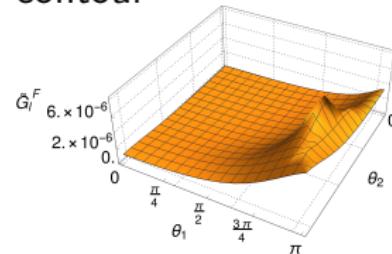
- ▶ decide which integration variable to integrate out exactly
- ▶ decide which contour to choose

# Choice of integration contour

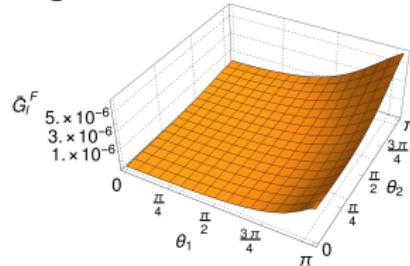
- ▶ choose contour with TAYINT but not exact integration



- ▶ choose exact integration variable with TAYINT but not contour



- ▶ choose exact integration variable and contour with TAYINT

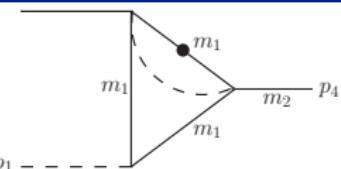
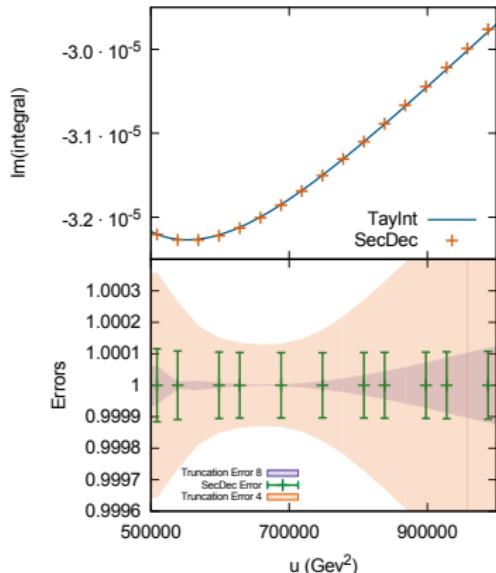
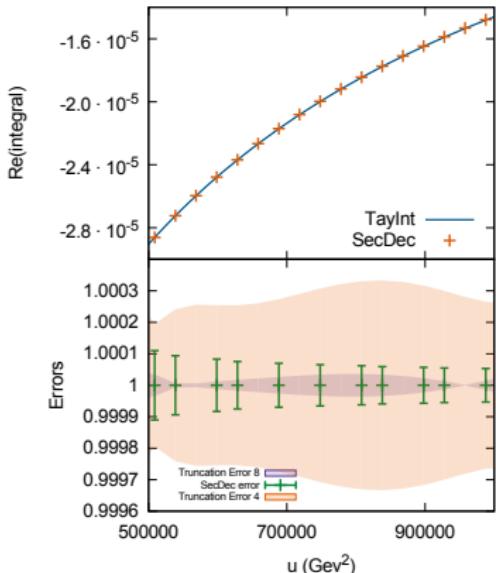


# Summary of TayInt steps

U1: reduce the Feynman Integral to a quasi finite basis	
U2: perform a sector decomposition on the finite integrals in the basis	
below threshold	above threshold
BT1: $t_j \rightarrow y_j$	OT1: $t_j \rightarrow \theta_j$ , generate $\mathcal{K}$ samples
BT2: Taylor expand the integrand and integrate	OT2: find contour configuration $\Theta_{o(1), \dots, o(J)}$
	OT3: perform one fold integrations
	OT4: find optimum resultant surface $\Theta_{o(1), \dots, o(J-1)}$
	OT5: determine partition $\mathcal{P}_j$
	OT6: Taylor expand and integrate

# Results for 3-scale integral I10

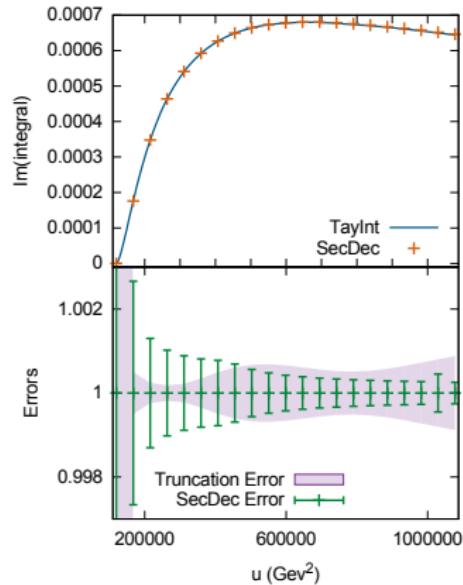
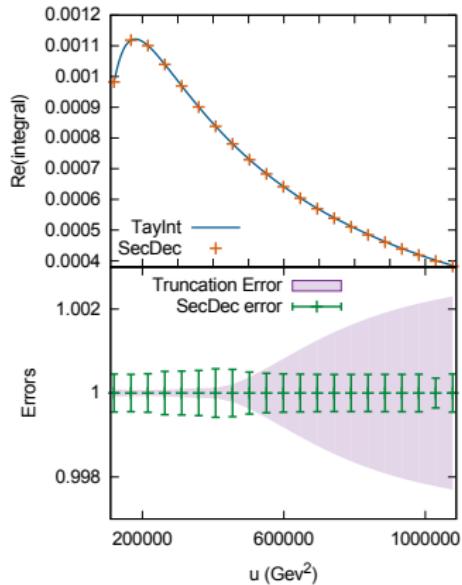
$$m_1 = 173, m_2 = \frac{1}{\sqrt{2}}m_1$$



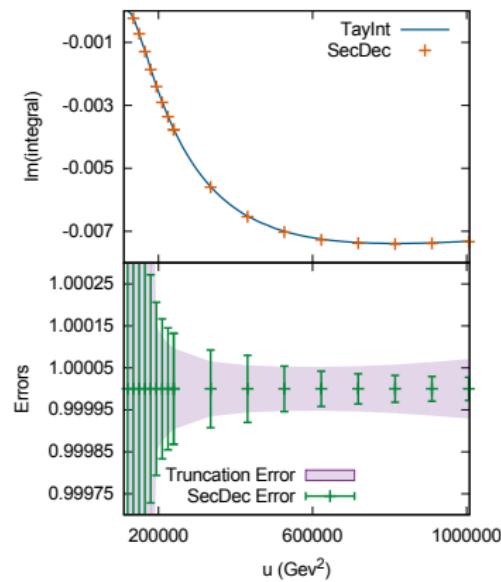
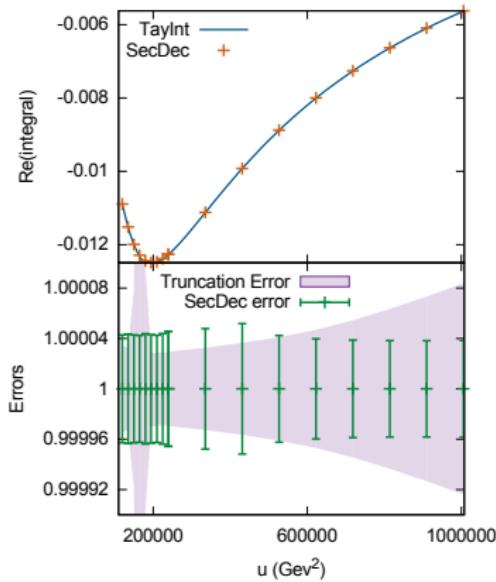
- ▶ 6<sup>th</sup> order Taylor expansion, using a 4-fold and 8-fold partitioning

# Results for 3-scale integral I10 $\mathcal{O}(\epsilon^1)$

- ▶ higher orders in  $\epsilon$  do not pose a problem
- ▶ 6<sup>th</sup> order Taylor expansion, 4-fold partitioning



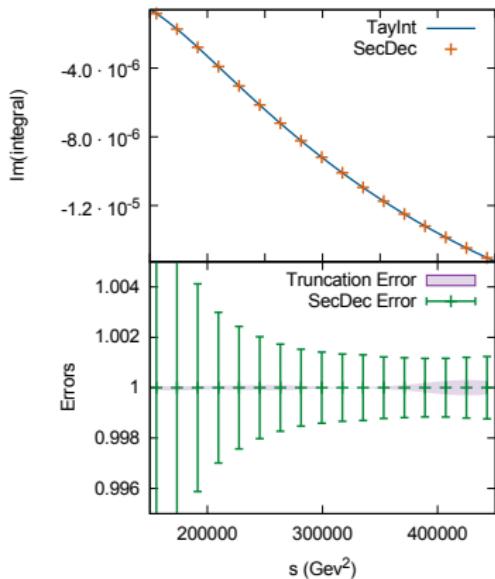
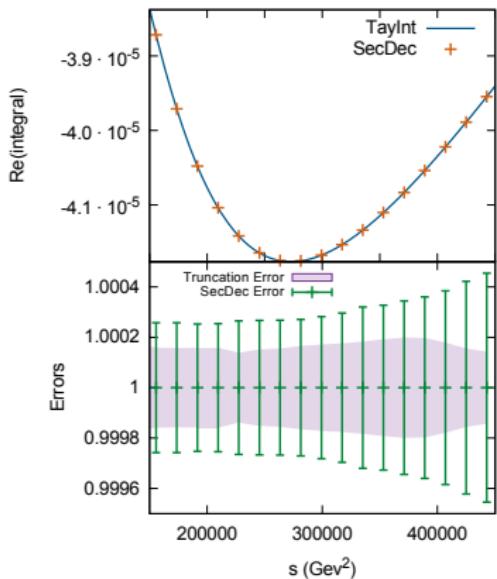
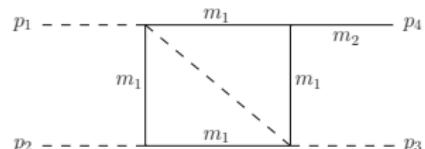
# Results for 3-scale integral I10 $\mathcal{O}(\epsilon^2)$



- oscillatory behaviour around zero of two subsector integrands lead to larger TAYINT error

# Results for 4-scale integral I39

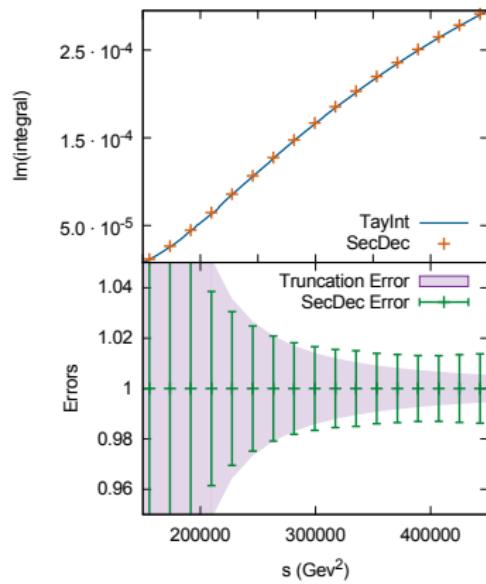
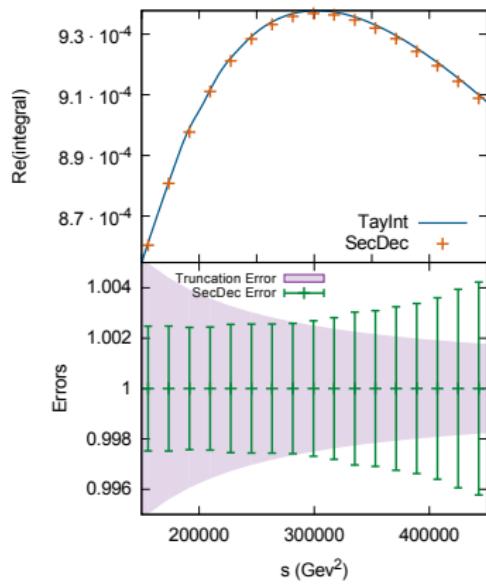
$$u = 2 \cdot 173^2, m_1 = 173, m_2 = \frac{1}{\sqrt{2}} m_1$$



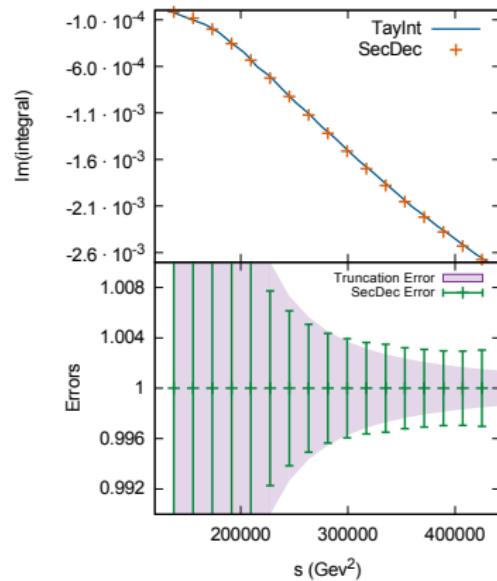
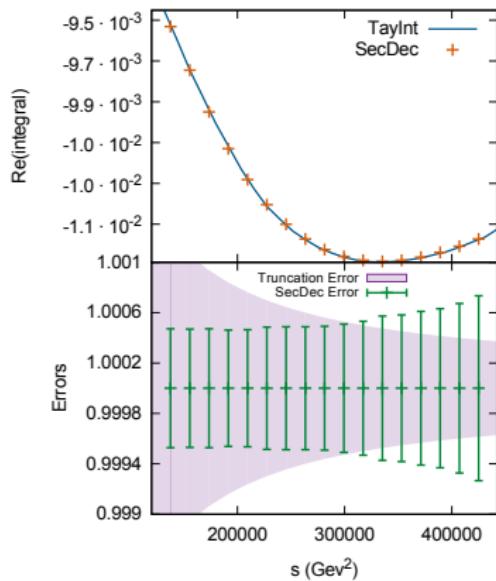
- ▶ 4<sup>th</sup> order Taylor expansion, 8-fold partitioning

# Results for 4-scale integral I39 $\mathcal{O}(\epsilon^1)$

- ▶ higher orders in  $\epsilon$  do not pose a problem, but no exact integration possible
- ▶ 4<sup>th</sup> order Taylor expansion, 6-fold partitioning



# Results for 4-scale integral I39 $\mathcal{O}(\epsilon^2)$



- ▶ 4<sup>th</sup> order Taylor expansion, 8-fold partitioning

# Summary & Outlook

## Summary

- ▶ Loopedia is a community-driven database for loop integral results
- ▶ PYSECDEC can be used for direct numerical integration of multi-loop multi-scale Feynman integrals
- ▶ TAYINT is a new method of accurate analytical approximations of multi-loop multi-scale Feynman integrals
- ▶ TAYINT shown to work for 4-scale integrals with 5 propagators

## Outlook

- ▶ Increase automation of algorithm, optimize generated file size
- ▶ application to more integrals entering phenomenological predictions

# Backup

# Input for pySecDec

- ▶ no input files needed, direct python interface
- ▶ many usage examples provided, e.g.

```
#! /usr/bin/env python
from pySecDec.loop_integral import loop_package
import pySecDec as psd

li = psd.loop_integral.LoopIntegralFromGraph(
internal_lines = [['m',[3,4]],[['m',[4,5]],[['m',[3,5]],[[0,[1,2]],[[0,[4,1]],[[0,[2,5]]],external_lines = [['p1',1],['p2',2],['p3',3]],

replacement_rules = [
    ('p1*p1', 0),
    ('p2*p2', 0),
    ('p3*p3', 's'),
    ('p4*p4', 0),
    ('p1*p2', 's/2'),
    ('p2*p3', '-s/2'),
    ('p1*p3', '-s/2'),
    ('m**2', 'msq')
]
)

Mandelstam_symbols = ['s']
mass_symbols = ['msq']
```

# pySecDec as a library: example snippet

```
#include "yyyy_bubble/yyyy_bubble.hpp"
#include "yyyy_box6Dim/yyyy_box6Dim.hpp"

/*
 * pySecDec Master Integrals
 */

// one loop bubble
yyyy_bubble::nested_series_t<secdecutil::UncorrelatedDeviation<yyyy_bubble::integrand_return_t>> bubble(yyyy_bubble::real_t uORT)
{
    using namespace yyyy_bubble;

    const std::vector<real_t> real_parameters{uORT};
    const std::vector<complex_t> complex_parameters{};

    // optimize contour
    const std::vector<nested_series_t<yyyy_bubble::integrand_t>> integrands = yyyy_bubble::make_integrands(real_parameters, complex_parameters
        // The number of samples for the contour optimization, the minimal and maximal deformation parameters, and the decrease factor can be
        // optionally set here as additional arguments.
    );

    // add integrands of sectors (together flag)
    const yyyy_bubble::nested_series_t<yyyy_bubble::integrand_t> summed_integrands = std::accumulate(++integrands.begin(), integrands.end(), *int

    // define the integrator
    auto integrator = secdecutil::cuba::Vegas<std::complex<double>>();
    integrator.flags = 2; // verbose output
    integrator.epsrel = 1e-5;
    integrator.epsabs = 1e-7;
    integrator.maxeval = 1e7;

    // integrate
    return secdecutil::deep_apply(summed_integrands, integrator.integrate) * yyyy_bubble::prefactor(real_parameters, complex_parameters);
}

...
/*
 * numerical amplitude using pySecDec Master Integrals
 */
secdecutil::Series<secdecutil::UncorrelatedDeviation<std::complex<double>>> yyyy_numerical(double s, double
{
    return -8.*(
        1. + (t*t + u*u)/s * box6Dim(t,u) +
        (t-u)/s*( bubble(u)-bubble(t) )
    );
}
```

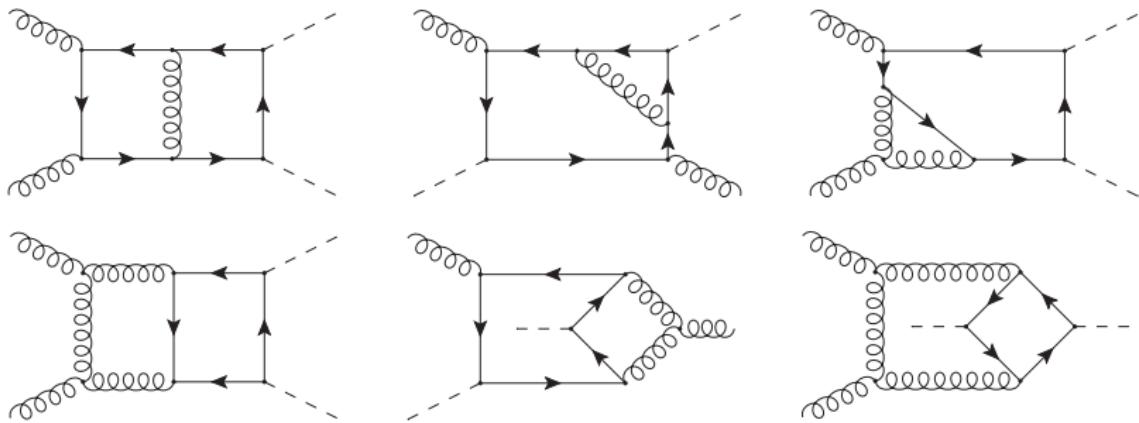
A bit of extra motivation to  
use pySecDec

# $gg \rightarrow hh$ @ NLO with full top mass dependence

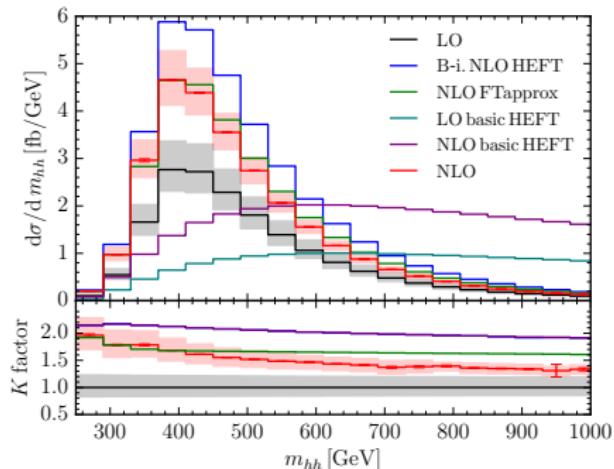
SB, N. Greiner, G. Heinrich, S.P. Jones, M. Kerner, J. Schlenk, U. Schubert, T. Zirke '16

Some of most complicated diagrams entering calculation:

- ▶ two-loop integrals, with numerators (4 independent mass scales:  $\hat{s}$ ,  $\hat{t}$ ,  $m_t^2$ ,  $m_h^2$  or 3 ratios), e.g.



# Exact result vs approximations at 100 TeV



- ▶ NLO HEFT good approximation for  $m_{hh} < 2m_t$
- ▶ scale uncertainties of HEFT and FT<sub>approx</sub> do not enclose central value of full result in  $m_{hh}$  tail → HEFT breaks down
- ▶ fast evaluation using interpolated grid Heinrich, Jones, Kerner, Luisoni, Vryonidou '17; Jones, Kuttimalai '17