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# The strong coupling from an improved tau vector isovector spectral function

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Diogo Boito

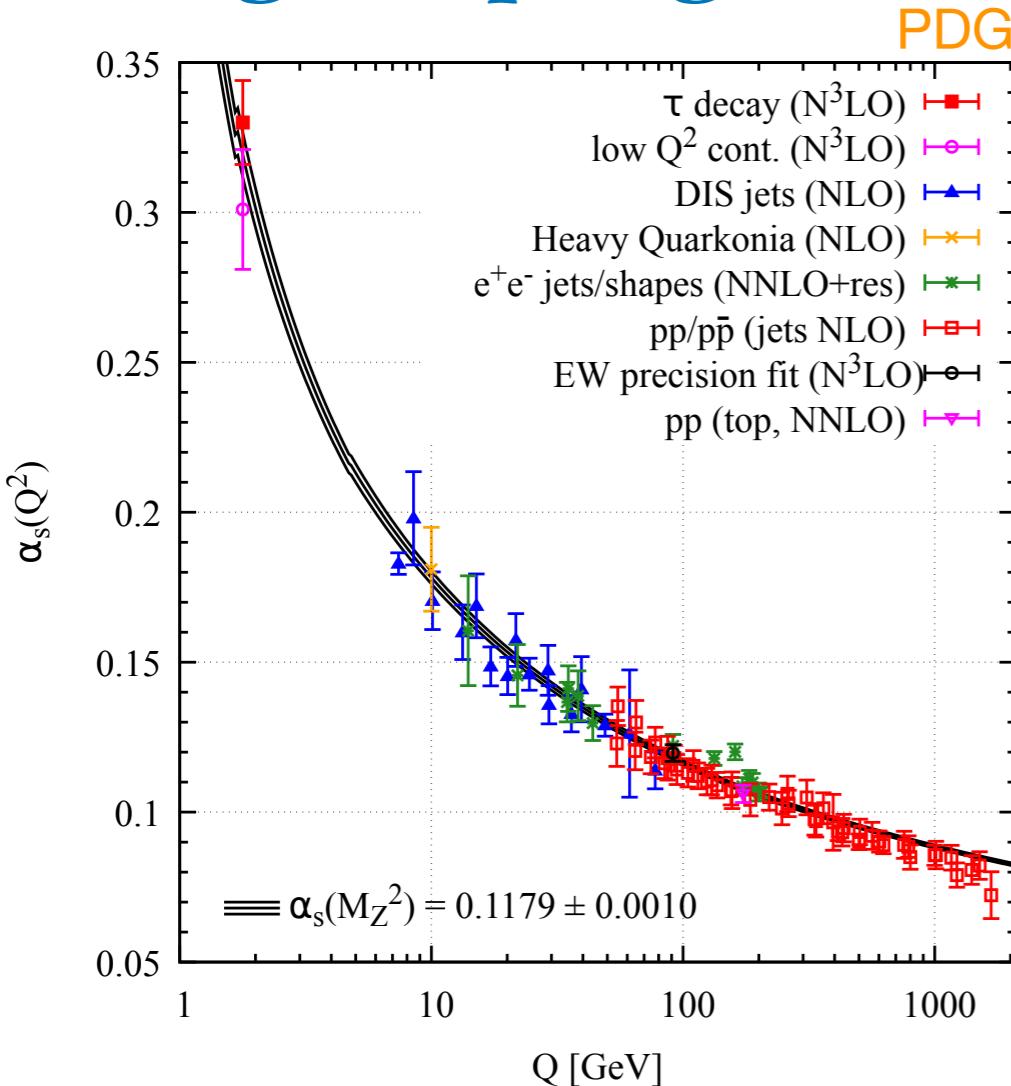
University of São Paulo  
University of Vienna

with Maarten Golterman, Kim Maltman, Santi Peris, Marcus Rodrigues and Wilder Schaaf

DB, Golterman, Maltman, Peris, Rodrigues, Schaaf, arXiv:2012.10440, PRD 103 (2021)



# strong coupling from tau decays



Lower energies

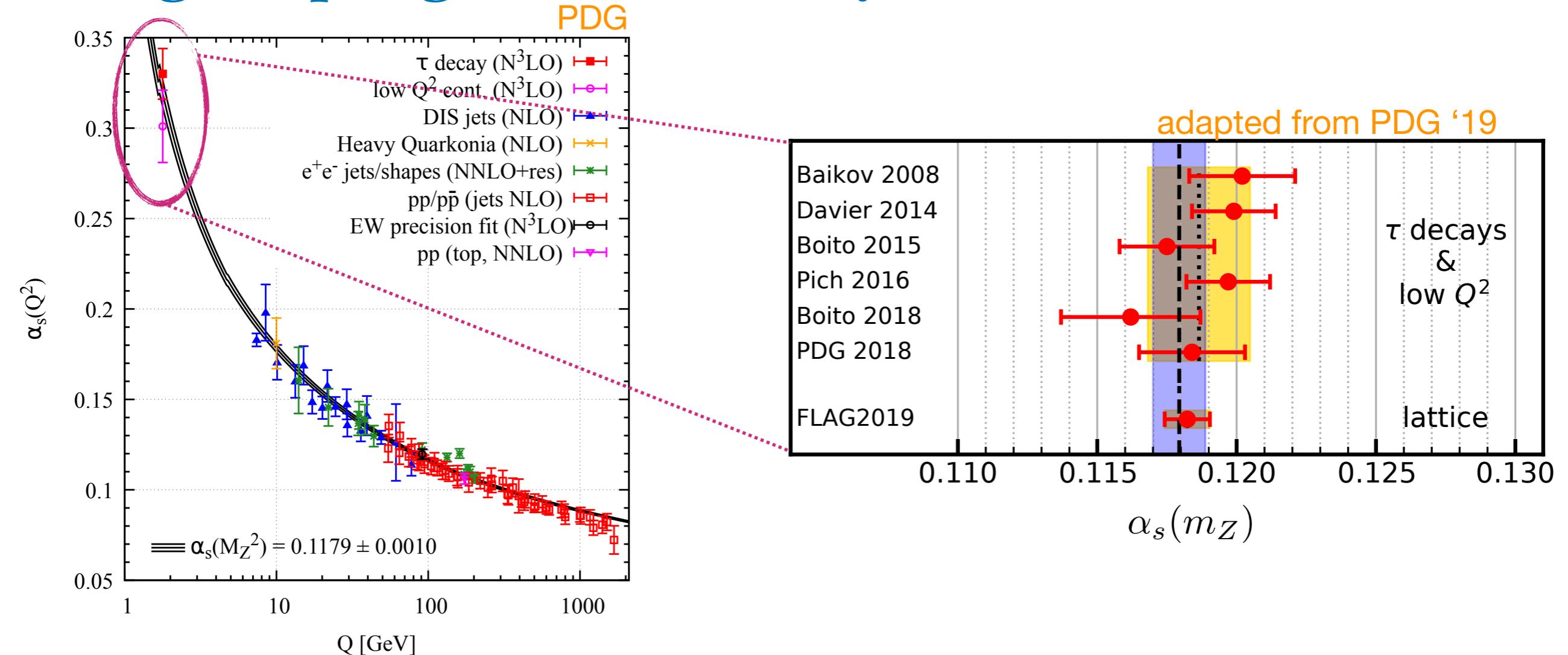
Larger coupling, more sensitivity to QCD corrections.

Larger non-perturbative physics (OPE, DVs),  
Problems with pt. theory  
(renormalons,...).

Higher energies

Smaller coupling, less sensitive to QCD corrections, more precision required from exp.  
Small contamination from non-perturbative physics,  
pt. series is "almost" convergent

# strong coupling from tau decays



Lower energies

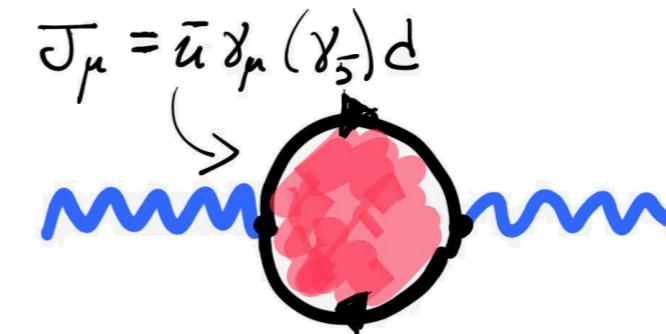
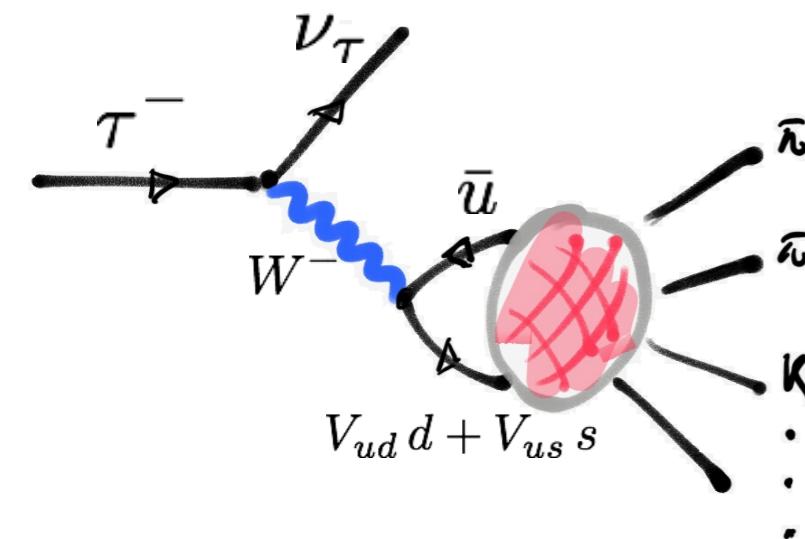
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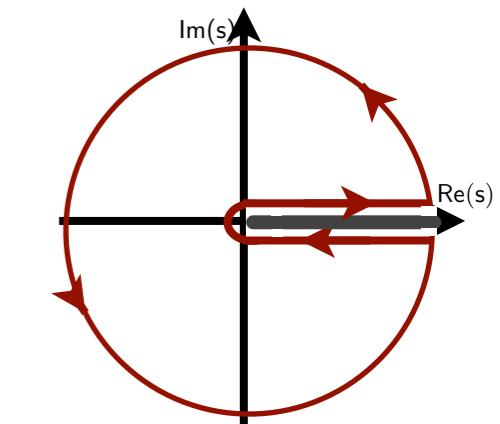


$$\Pi_{\mu\nu}(q) = i \int d^4x e^{iqx} \langle 0 | T\{ J_\mu(x) J_\nu(0)^\dagger \} | 0 \rangle$$

Massless (V&A) correlators

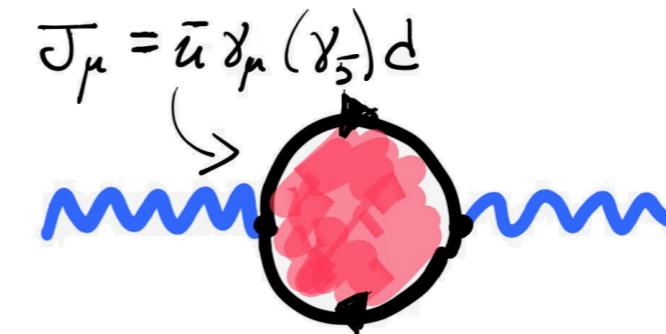
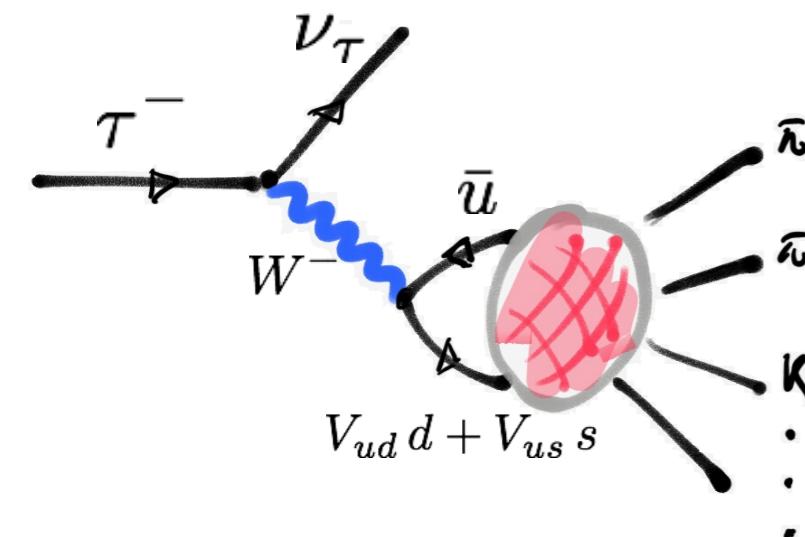
Braaten, Narison, and Pich '92

Sum rules (using Cauchy's theorem)



$$\frac{1}{s_0} \int_0^{s_0} ds w(s) \frac{1}{\pi} \text{Im} \Pi(s) = - \frac{1}{2\pi i s_0} \oint_{|z|=s_0} dz w(z) \Pi(z)$$

# strong coupling from tau decays



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Braaten, Narison, and Pich '92

Sum rules (using Cauchy's theorem)

experiment

$$\frac{1}{s_0} \int_0^{s_0} ds w(s) \frac{1}{\pi} \text{Im} \Pi(s) = - \frac{1}{2\pi i s_0} \oint_{|z|=s_0} dz w(z) \Pi(z)$$

theory

A diagram of the complex plane with the horizontal axis labeled  $\text{Re}(s)$  and the vertical axis labeled  $\text{Im}(s)$ . A contour line is shown in red, enclosing the origin in the upper half-plane. The contour has a small circular indentation on the real axis to the left of the origin, indicated by a blue arrow.

# theory overview

$$\frac{-1}{2\pi i} \oint_{|z|=s_0} dz w(z) \Pi(z) \approx S_{\text{EW}} N_c (1 + \delta^{(0)} + \delta_{\text{EW}} + \delta_{\text{OPE}} + \delta_{\text{DVs}})$$

theory

Perturbation theory (OPE)  $\longrightarrow$  
$$\sum_{n=0}^4 \left(\frac{\alpha_s}{\pi}\right)^n \sum_{k=0}^{n+1} c_{n,k} \log^k \left(\frac{-s}{\mu^2}\right) + \frac{C_4}{Q^4} + \frac{C_6}{Q^6} + \frac{C_8}{Q^8} + \dots$$

see Pich's talk

Gorishnii, Kataev, Larin '91  
Surguladze&Samuel '91

Baikov, Chetyrkin, Kühn '08

$$\alpha_s^1 \quad \alpha_s^2 \quad \downarrow \alpha_s^3 \quad \downarrow \alpha_s^4$$

**pt. correction is ~20%**

$$\delta_{\text{FO}}^{(0)} = 0.1012 + 0.0533 + 0.0273 + 0.0133 = 0.1952$$

(fixed order perturbation theory)

# theory overview

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**theory**

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Surguladze&Samuel '91

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$$\alpha_s^1 \quad \alpha_s^2 \quad \downarrow \alpha_s^3 \quad \downarrow \alpha_s^4$$

**pt. correction is  $\sim 20\%$**

$$\delta_{\text{FO}}^{(0)} = 0.1012 + 0.0533 + 0.0273 + 0.0133 = 0.1952$$

(fixed order perturbation theory)

Duality Violations

$$\rightarrow \rho_{\text{DV}}(s) = e^{-\delta - \gamma s} \sin(\alpha + \beta s)$$

Ansatz based on widely accepted assumptions about QCD: Regge behaviour and large- $N_c$ . Main expected corrections: logarithmic and powers of  $1/s$ .

DB, Caprini, Golterman, Maltman, Peris, PRD '18

# theory: FOPT vs CIPT

Fixed Order (FO) or Contour Improved (CI) lead to different  $\alpha_s$  values

see Hoang's talk

theoretical uncertainty?

$$\alpha_s^1 \quad \alpha_s^2 \quad \alpha_s^3 \quad \alpha_s^4$$

$$\delta_{\text{FO}}^{(0)} = 0.1012 + 0.0533 + 0.0273 + 0.0133 = 0.1952$$

$$\delta_{\text{CI}}^{(0)} = 0.1375 + 0.0262 + 0.0104 + 0.0072 = 0.1814$$

# theory: FOPT vs CIPT

Fixed Order (FO) or Contour Improved (CI) lead to different  $\alpha_s$  values

see Hoang's talk

theoretical uncertainty?

$$\alpha_s^1 \quad \alpha_s^2 \quad \alpha_s^3 \quad \alpha_s^4$$

$$\delta_{\text{FO}}^{(0)} = 0.1012 + 0.0533 + 0.0273 + 0.0133 = 0.1952$$

~~$$\delta_{\text{CI}}^{(0)} = 0.1375 + 0.0262 + 0.0104 + 0.0072 = 0.1814$$~~

Discrepancy between FOPT and CIPT  
 linked to an incoherent treatment of the OPE  
 (previously assumed to be the same for both prescriptions)

Hoang and Regner '20.'21

We will not quote results from CIPT in this talk

# analysis strategy

$$\frac{-1}{2\pi i} \oint_{|z|=s_0} dz w(z) \Pi(z) \approx S_{\text{EW}} N_c (1 + \delta^{(0)} + \delta_{\text{EW}} + \delta_{\text{OPE}} + \delta_{\text{DVs}})$$

theory

Desired properties from the choice of weights

1. Good perturbative behaviour.
2. Small condensate contributions.
3. Suppression of DVs.

# analysis strategy

$$\frac{-1}{2\pi i} \oint_{|z|=s_0} dz w(z) \Pi(z) \approx S_{\text{EW}} N_c (1 + \delta^{(0)} + \delta_{\text{EW}} + \delta_{\text{OPE}} + \delta_{\text{DVs}})$$

theory

Desired properties from the choice of weights

- 1. Good perturbative behaviour.
- 2. Small condensate contributions.
- 3. Suppression of DVs.

## Choice of weights

$w_0(y) = 1$

Tiny condensate contributions, sensitive to DVs

$w_2(y) = 1 - y^2$

Only D=6

$w_3(y) = (1 - y)^2(1 + 2y)$

Only D=6 and 8 Tau kinematical Moment ( $R_\tau$ )

$w_4(y) = (1 - y^2)^2$

Only D=6 and 10

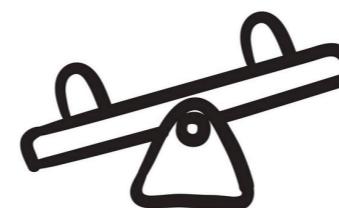
# analysis strategy

Suppression of DVs comes with the price of additional (unknown) higher dim. contributions from the OPE.

## DV strategy

DB, M. Golterman, K. Maltman, S. Peris,  
M. V. Rodrigues and W. Schaaf, 2012.10440

- Accept some DVs, strongly suppress contamination on the OPE side.



## Truncated OPE strategy

A Pich, A. Rodriguez-Sanchez 1605.06830  
Davier, Höcker, Malaescu, Yuan, Zhang 1312.1501

[see Golterman's talk](#)

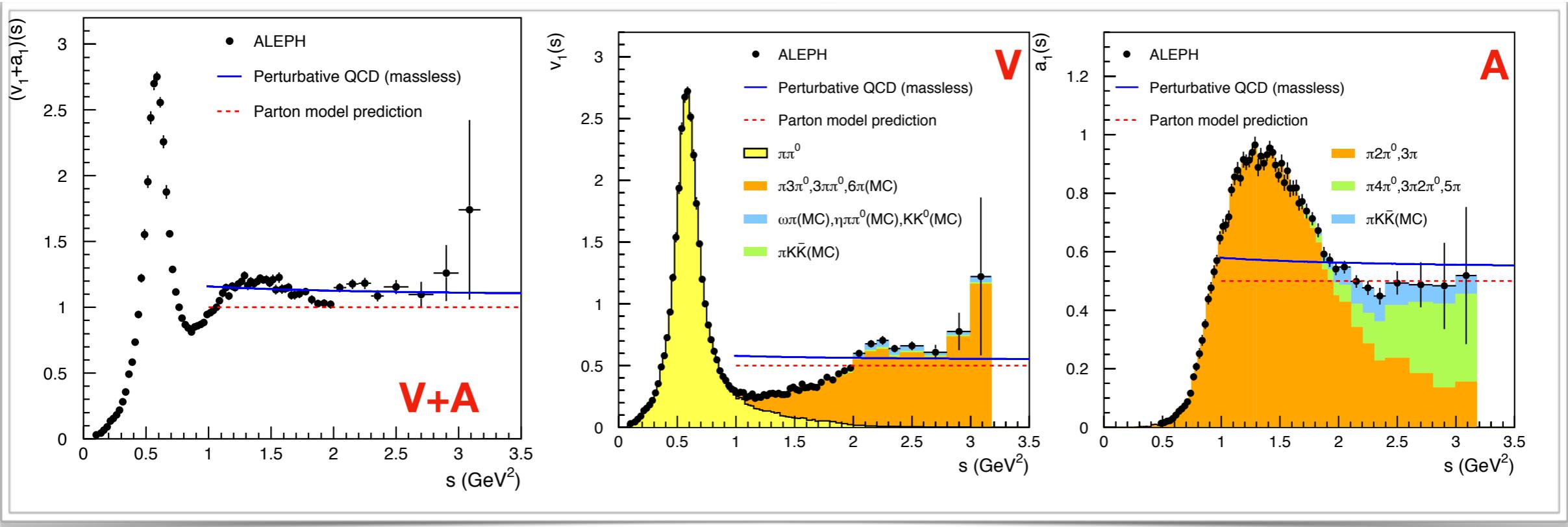
- Suppress DVs but need to ignore the higher order contributions on the OPE side (too many parameters).

(Serious issues with the truncation of the OPE)  
DB, M. Golterman, K. Maltman, S. Peris '16 '19

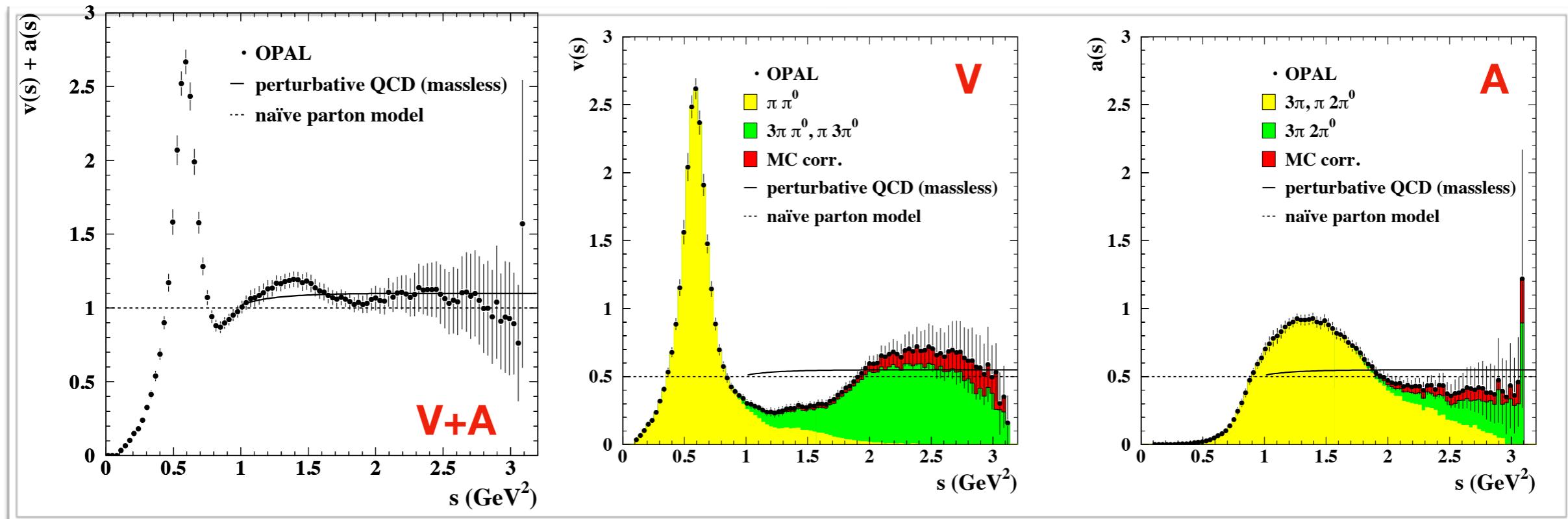
# Data

# inclusive hadronic tau decay data

Davier et al [ALEPH] '14

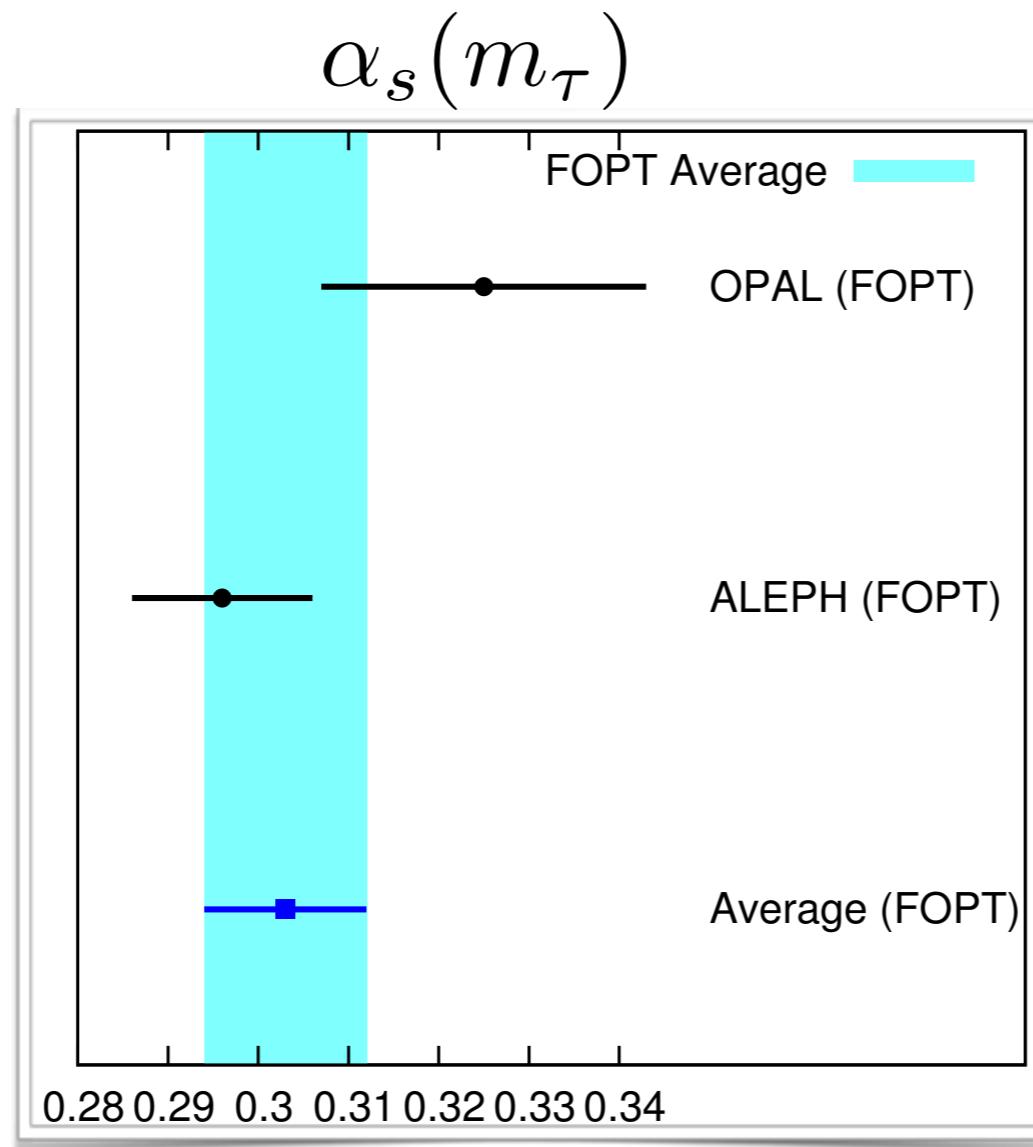


Ackerstaff et al [OPAL] '98 updated for recent values of branching fractions Boito et al '11, '21



# tau decay data and the strong coupling

The two data sets lead to compatible values for  $\alpha_s(m_\tau)$  : average



DB, Golterman, Jamin, Mahdavi,  
Maltman, Osborne, Peris, '12

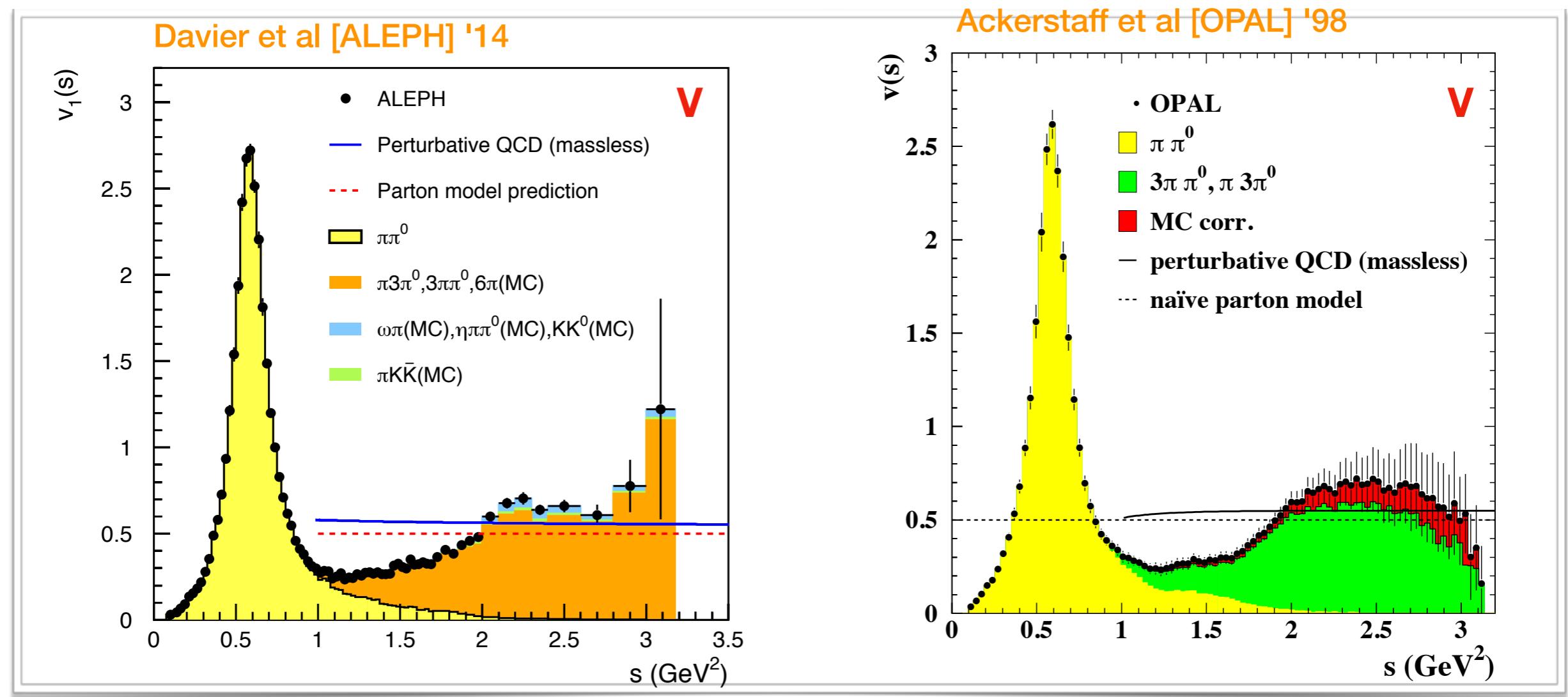
DB, Golterman, Maltman, Osborne, Peris, '15

Are the data sets compatible? Can we combine them?

Can the inclusive spectral functions be improved with recent data?

# anatomy of the data sets

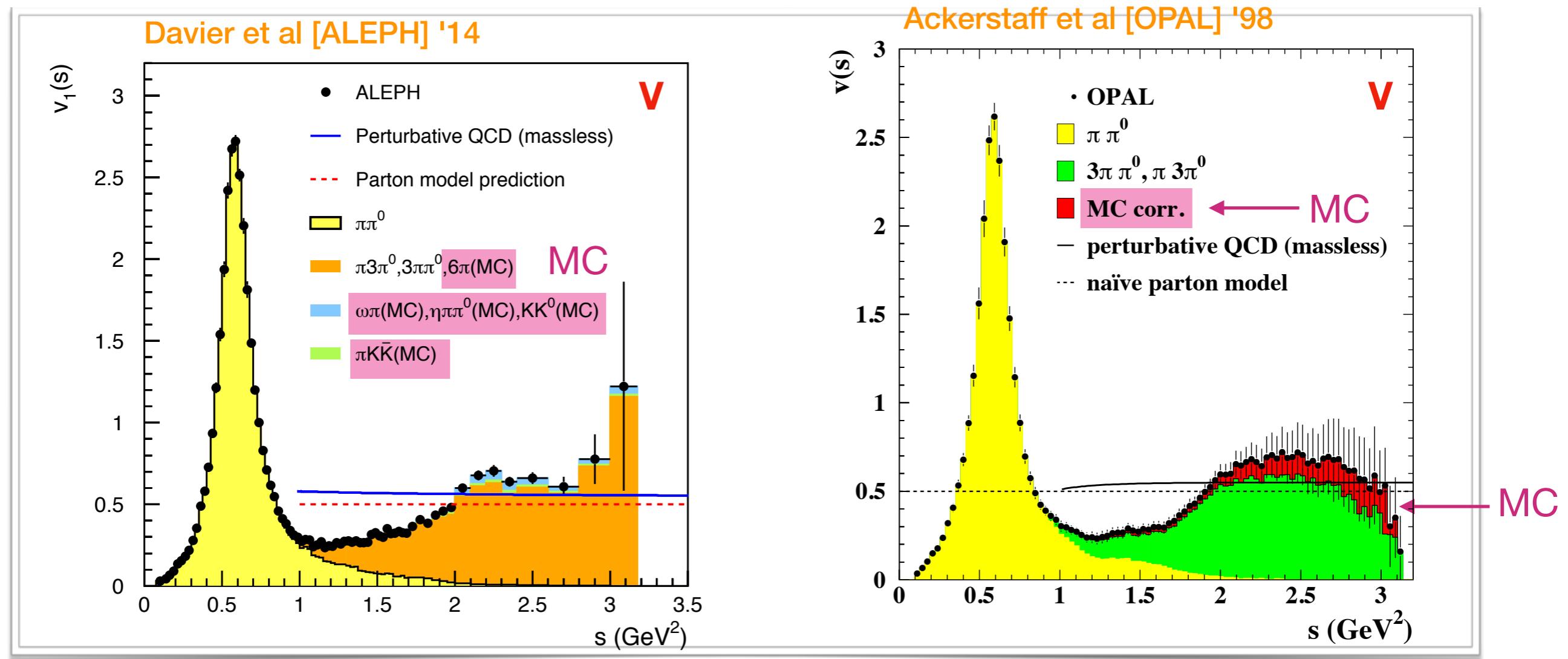
- V channel dominated by  $\tau \rightarrow 2\pi + \nu_\tau$  and  $\tau \rightarrow 4\pi + \nu_\tau$
- “Residual” channels subdominant (but important for  $\alpha_s$ !)
- Monte Carlo (MC) inputs for several channels



Recently measured channels in  $e^+e^-$  can be used to improve the vector channel

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# new vector isovector spectral function

- Combined data for  $2\pi$  and  $4\pi$  channels from ALEPH & OPAL

Data combination: same algorithm used in R-data combination for muon g-2.

Keshavarzi, Nomura, Teubner '18

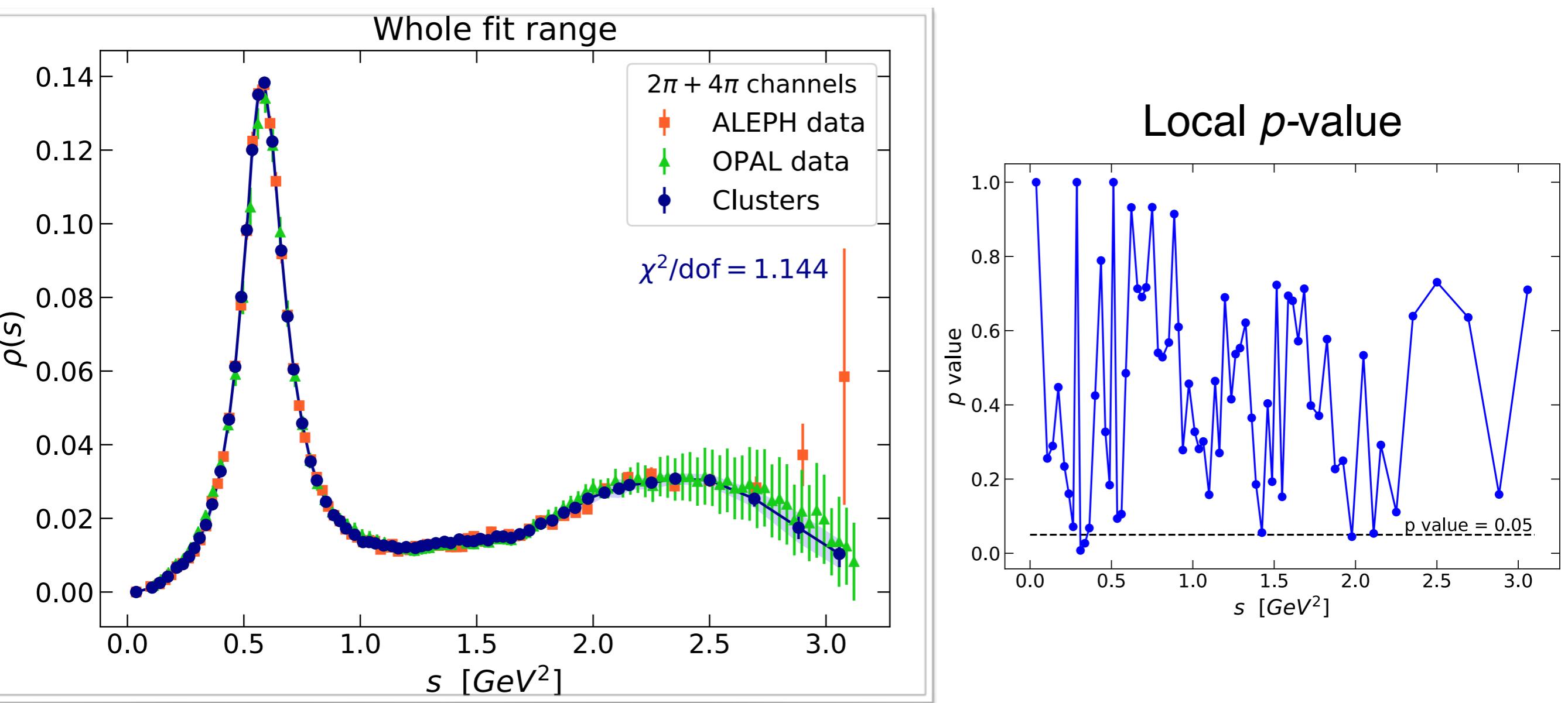
- Exp. data only: 7 residual channels from  $e^+e^-$  using CVC (conserved vector current) and BaBar data for  $\tau \rightarrow KK_S\nu_\tau$

No Monte Carlo inputs; IB corrections to CVC negligible

- Results updated for recent branching ratio measurements

# new vector isovector spectral function

Combination of  $2\pi + 4\pi$  channels  
 Good  $\chi^2$  both locally and globally, no  $\chi^2$  inflation needed



DB, Golterman, Maltman, Peris, Rodrigues and Schaaf, arXiv:2012.10440

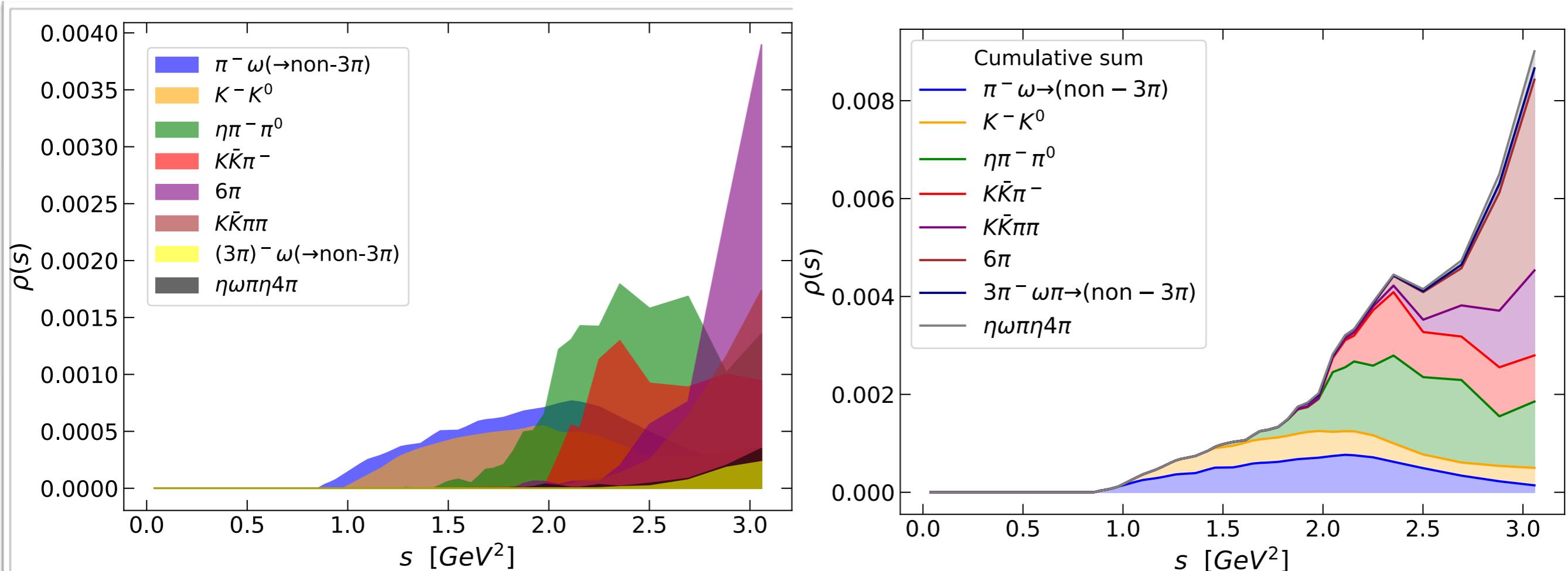
# new vector isovector spectral function

7 residual channels extracted from  $e^+e^-$  data + BaBar data for  $\tau \rightarrow KK_S\nu_\tau$

Dramatic improvement in errors for higher multiplicity modes (near end point)

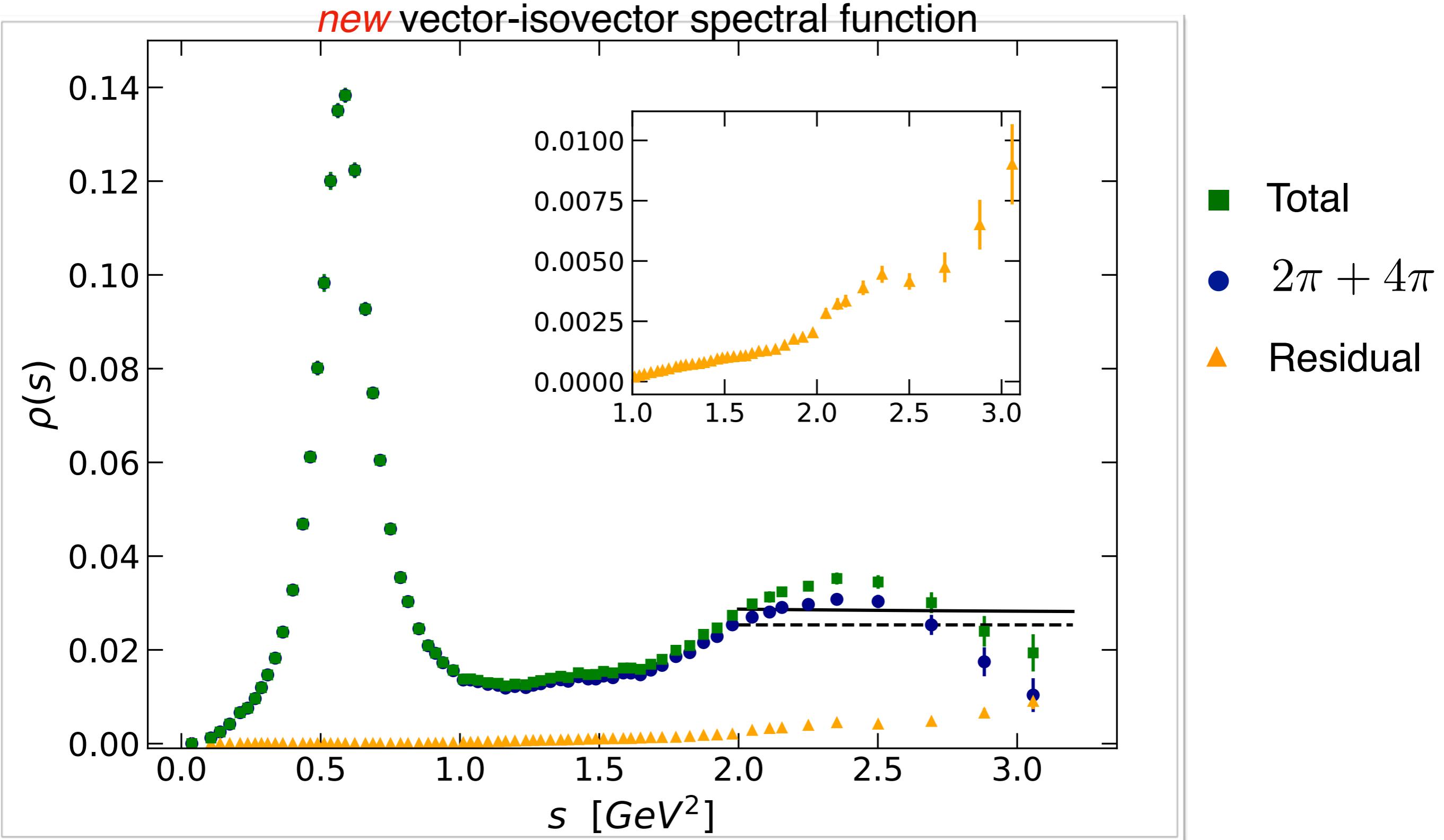
## No Monte Carlo input

Original data sets from: BABAR, CMD-3 and SND (results from 16 papers)



# new vector isovector spectral function

Combined  $2\pi + 4\pi$  (ALEPH and OPAL) + residual channels from data  
 99.95% of the Branching Fraction covered



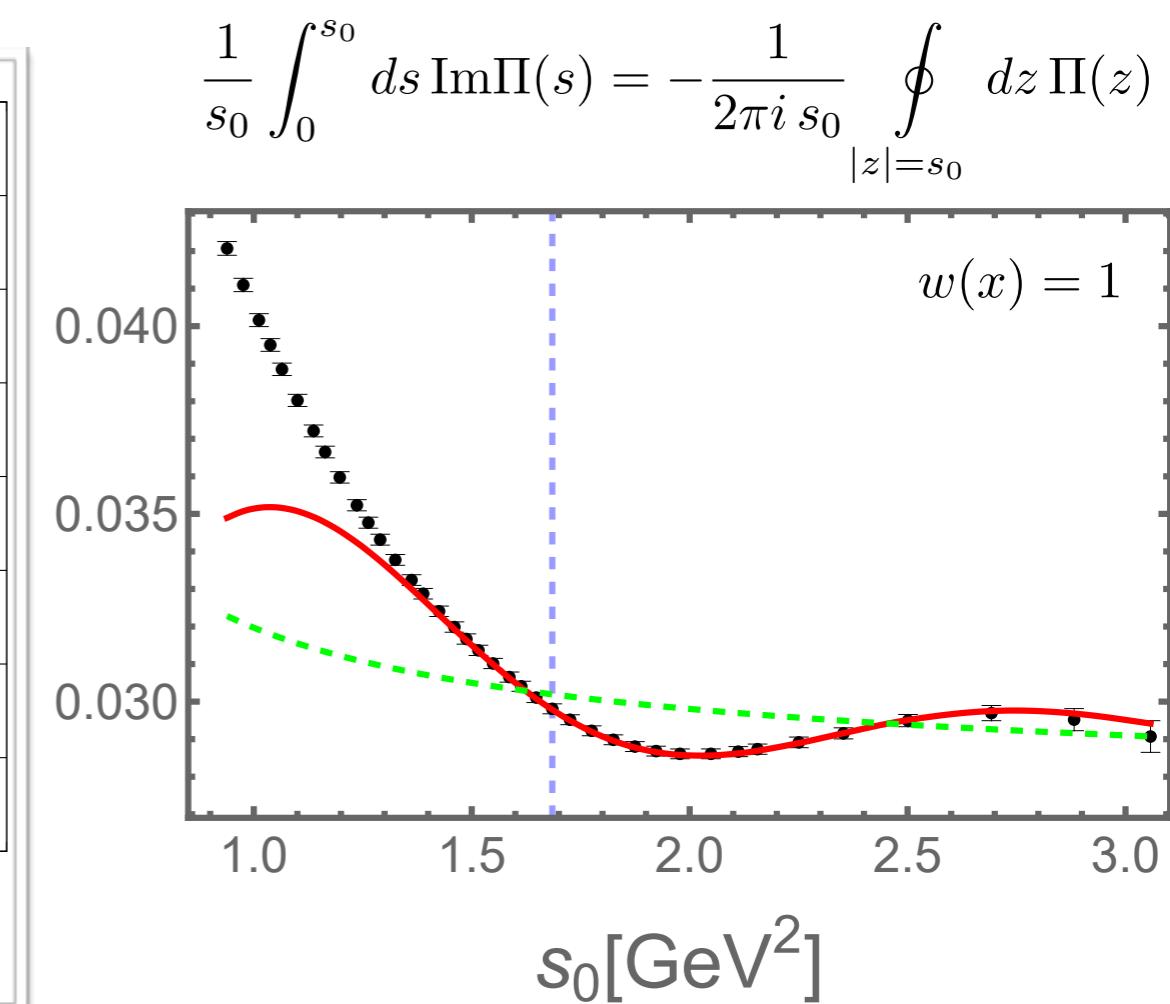
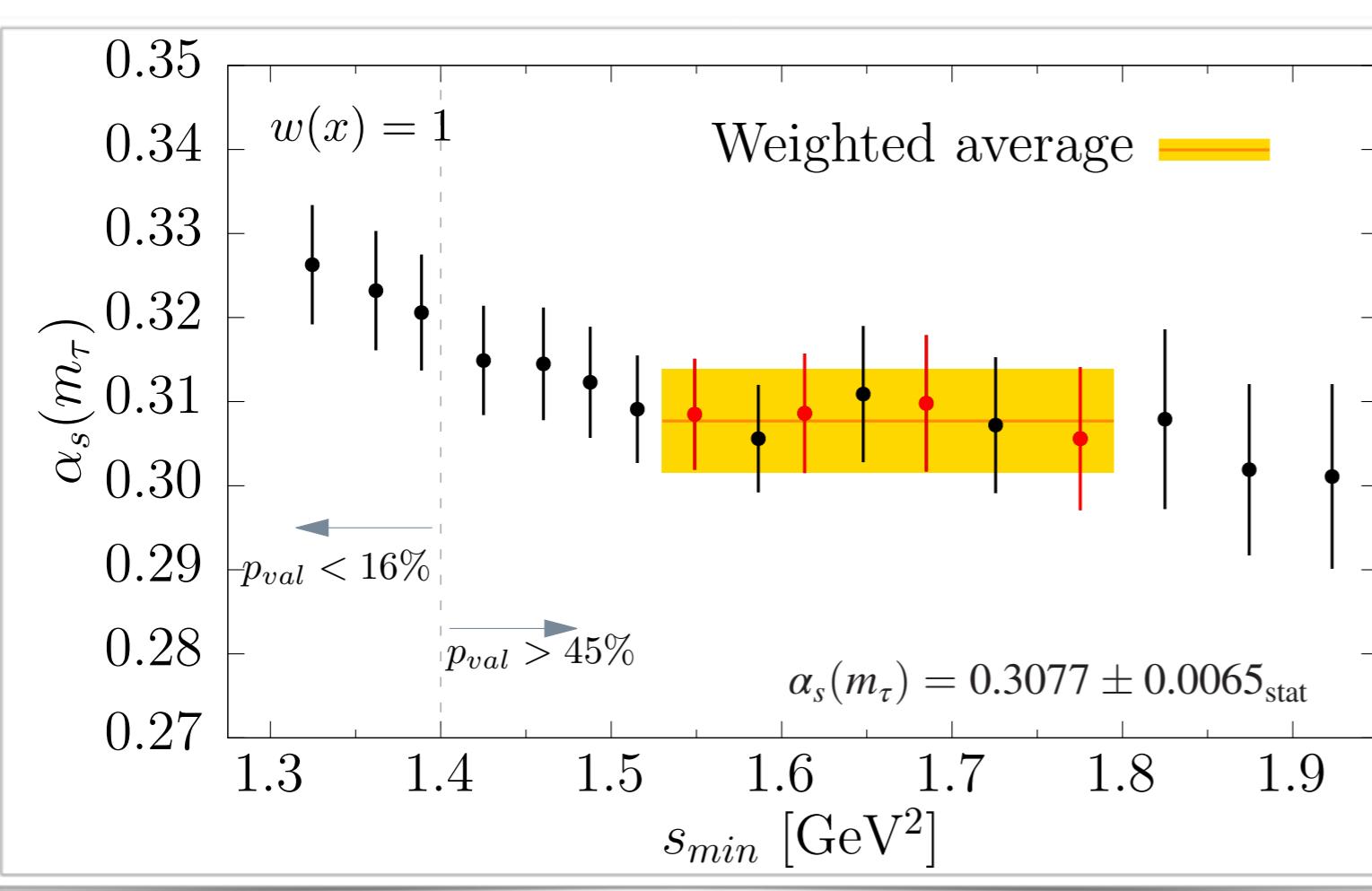
# Results

# strong coupling from the new spectral function

Several fits, single moments or in combination

Many fit windows:  $[s_{\min}, m_\tau^2]$

Consistency between different fits ( $\alpha_s$ , condensates, DV params.)



# strong coupling from the new spectral function

Consistency between different fits

mom.	$\alpha_s$	$c_6[\text{GeV}^6]$
$w_0$	0.3077(65)	—
$w_0 \& w_2$	0.3091(69)	-0.0059(13)
$w_0 \& w_3$	0.3080(70)	-0.0070(12)
$w_0 \& w_4$	0.3079(70)	-0.0068(12)

$$w_0(y) = 1$$

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Final value

pt. series truncation, scale variation



$$\alpha_s(m_\tau) = 0.3077 \pm 0.0065_{\text{stat}} \pm 0.0038_{\text{pert}}$$

$$= 0.3077 \pm 0.0075 \quad (n_f = 3, \text{FOPT})$$

# Final result

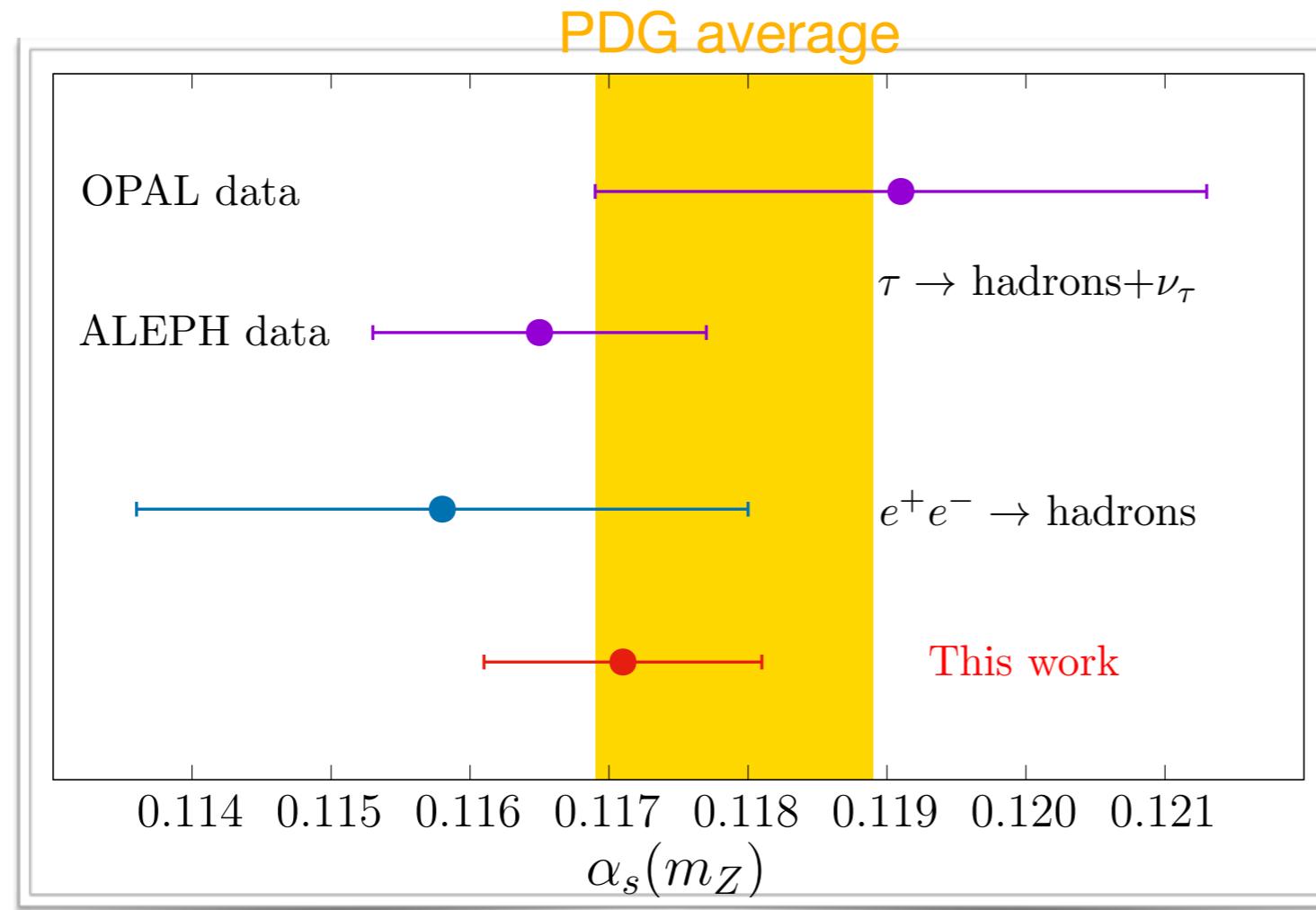
## Results at $m_\tau$

$$\begin{aligned}\alpha_s(m_\tau) &= 0.3077 \pm 0.0065_{\text{stat}} \pm 0.0038_{\text{pert}} \\ &= 0.3077 \pm 0.0075 \quad (n_f = 3, \text{FOPT})\end{aligned}$$

## Results evolved to $m_Z$

$$\alpha_s(m_Z) = 0.1171 \pm 0.0010$$

$(\overline{\text{MS}}, N_f = 5)$



DB, Golterman, Jamin, Mahdavi, Maltman, Osborne, Peris, '12

DB, Golterman, Maltman, Osborne, Peris, '15

DB, Golterman, Keshavarzi, Maltman, Nomura, Peris, Teubner '18

DB, Golterman, Maltman, Peris, Rodrigues, Schaaf, '21

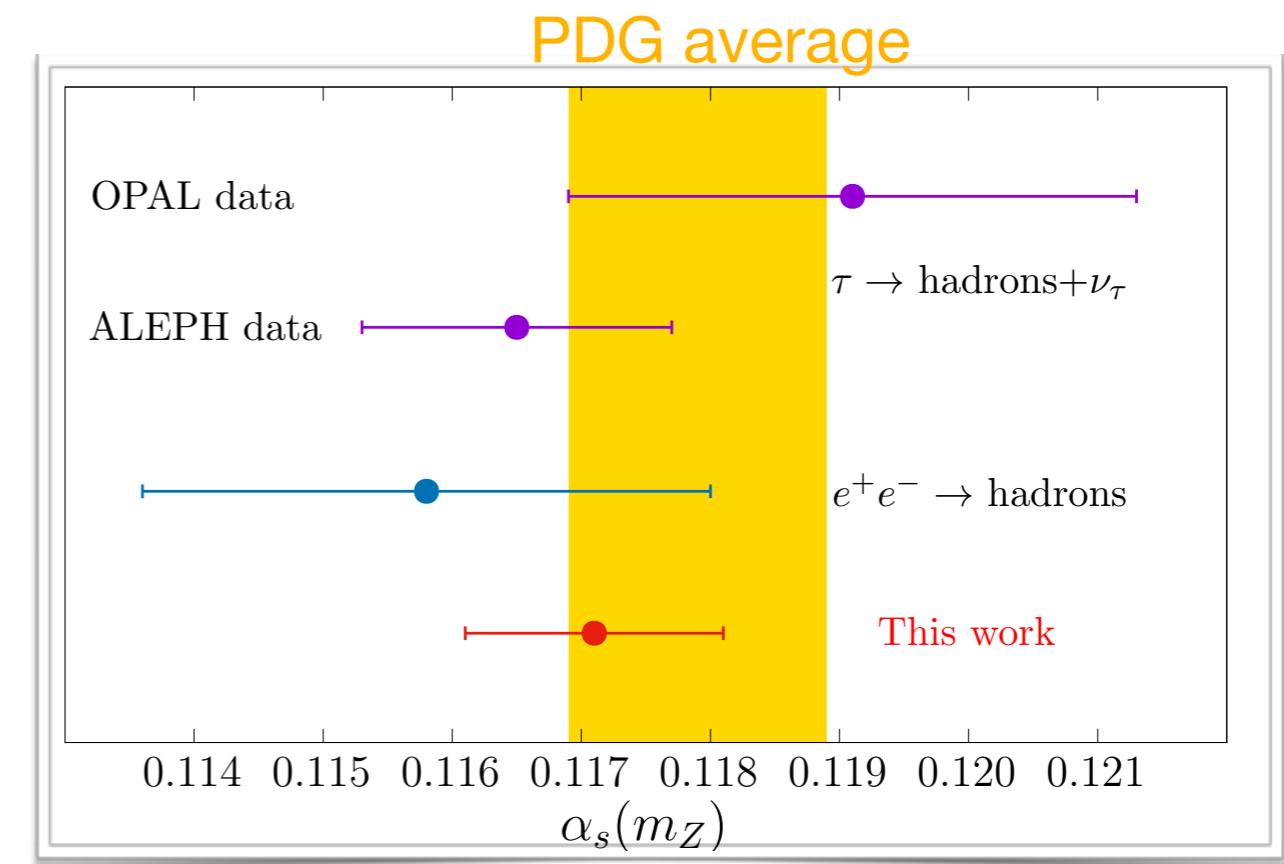
# conclusions

- Vector channel is special: CVC allows improvement near tau kin. end point.
- New vector isovector spectral function purely based on data, **no MC input**.
- Analysis can be improved with new data for the  $2\pi + 4\pi$  channels only!



- Improvements of this type not possible for the axial channel (no axial photon).
- Final result from the new vector spectral function is competitive.

$$\alpha_s(m_Z) = 0.1171 \pm 0.0010$$





# Extra

