

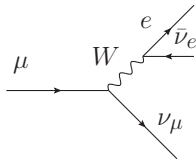
Third order correction to the muon lifetime

16th International Workshop on Tau Lepton Physics

Matteo Fael | Sept. 27, 2021

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in collaboration with [K. Schönwald](#), [M. Steinhauser](#), [Phys.Rev.D 104 \(2021\) 016003](#)



- Fundamental process to study weak interaction.
- Fermi Constant G_F extracted from muon lifetime

$$\frac{\hbar}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} F(\rho) \left[1 + H_1(\rho) \frac{\hat{\alpha}(m_\mu)}{\pi} + H_2(\rho) \left(\frac{\hat{\alpha}(m_\mu)}{\pi} \right)^2 + H_3(\rho) \left(\frac{\hat{\alpha}(m_\mu)}{\pi} \right)^3 \right]$$

with $\rho = m_e/m_\mu \simeq 1/210$.

■ $O(\alpha)$ Behrends, Finkelstein, Sirlin Phys.Rev.101 (1956) 866.

■ $O(\alpha^2)$
 van Ritbergen, Stuart Phys.Rev.Lett. 82 (1999) 488 ($m_e = 0$)
 Czarnecki, Pak, Phys.Rev.Lett. 100 (2008) 241807 ($m_e \ll m_\mu$)

■ $O(\alpha^3)$ **NEW**

MF, Schönwald, Steinhauser, Phys.Rev.D 104 (2021) 016003
 ($m_e \sim m_\mu$)
 parts of the calculation also confirmed in
 Czakon, Czarnecki, Dowling Phys.Rev.D 103 (2021) L111301.

- Fine structure constant [CODATA](#)

$$1/\alpha = 137.035\,999\,084\,(21) \text{ (0.15 ppb)}$$

from a_e : 0.35 ppb [Aoyama, Hayakawa, Kinoshita, Nio, PRL 109 \(11\) 111807](#)
from h/m in ^{87}Rb : (0.08ppb) [Morel, Yao, Cladé, Guellati-Khélifa, Nature 588 61](#)

- Fermi Constant

$$G_F = 1.166\,378\,7\,(6) \times 10^{-5} \text{ GeV}^{-2} \text{ (0.5 ppm)}$$

[Webber, et al., 2011, Phys.Rev.Lett.106, 041803; Phys.Rev.Lett.106, 079901](#)

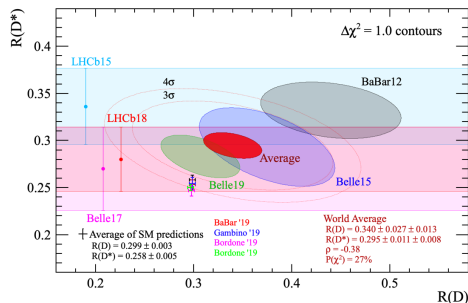
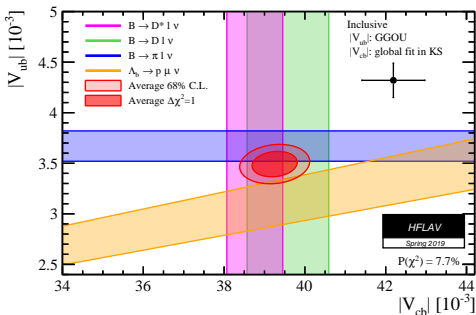
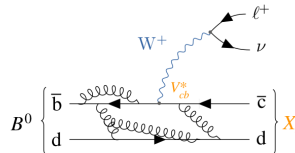
- Z-boson mass

$$M_Z = 91.1876 \pm 0.0021 \text{ GeV (23 ppm)}$$

[P.A. Zyla et al. \(Particle Data Group\), Prog. Theor. Exp. Phys. 2020, 083C01 \(2020\)](#)
FCC-ee aims at improving M_Z by factor 400 (stat).

Semileptonic B decays

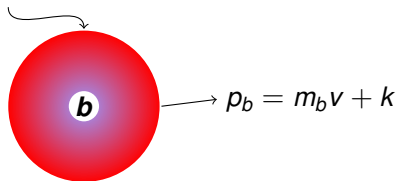
- $b \rightarrow cl\bar{\nu}_\ell$ ($\ell = e, \mu$) sensitive to $|V_{cb}|$.
- **Inclusive decays**
 - $\bar{B} \rightarrow X_c l \bar{\nu}$, with $X_c = D, D^*, D\pi, \dots$



The Heavy-Quark Expansion

$$\Gamma_{sl} = \Gamma_0 + \Gamma_{\mu\pi} \frac{\mu_\pi^2}{m_b^2} + \Gamma_{\mu G} \frac{\mu_G^2}{m_b^2} + \Gamma_{\rho D} \frac{\rho_D^3}{m_b^3} + \Gamma_{\rho LS} \frac{\rho_{LS}^3}{m_b^3} + \dots$$

quark-gluon cloud

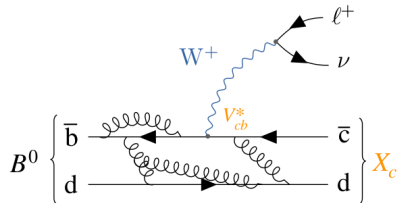


Reviews:

Benson, Bigi, Mannel, Uraltsev, Nucl.Phys. B665 (2003) 367;

Dingfelder, Mannel, Rev.Mod.Phys. 88 (2016) 035008.

- Γ_i are computed in **perturbative QCD**.
- Γ_0 : **free-quark decay** $b \rightarrow cl\nu$!
- The HQE parameters:
 $\mu_\pi, \mu_G, \rho_D, \rho_{LS} \sim \langle B | \mathcal{O}_i^{\bar{b}b} | B \rangle$
- HQE parameters are **extracted from data or lattice**.



$$\frac{\hbar}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} \left[X_0(\rho) + \sum_{n \geq 1} \left(\frac{\alpha}{\pi}\right)^n X_n(\rho) \right]$$

Possible strategies

- Exact result with m_μ and m_e only at $O(\alpha)$.

Nir, Phys.Lett.B 221 (1989) 184

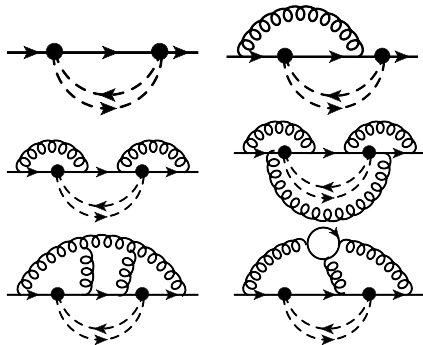
- Set $m_e = 0$ ($m_c = 0$).

van Ritbergen, Stuart

- Approximation exploiting $m_e < m_\mu$

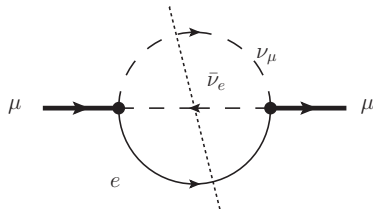
$(m_c < m_b)$.

Czarnecki, Pak



$$\frac{\hbar}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} \left[1 - 8\rho^2 - 12\rho^4 \log(\rho^2) + 8\rho^6 - \rho^8 \right]$$

where $\rho = m_e/m_\mu$.



- Optical theorem.
- Multi-loop diagrams with **two scales m_e and m_μ** .
- Use **method of regions**:

Beneke, Smirnov, NPB 522 (1998) 321; Smirnov, Springer Tracts Mod. Phys. 177

- Expansion in $\rho = m_e/m_\mu$ (or $\rho = m_c/m_b$).
Obvious choice, too difficult to extend at $O(\alpha^3)$.
- Expansion in $\delta = 1 - m_e/m_\mu$ (or $\delta = 1 - m_c/m_b$)
Crucial decoupling and simplifications of loop integrals.

A toy example

$$F(m, M, \varepsilon) = \int_0^\infty dk \frac{k^{-\varepsilon}}{(k+m)(k+M)} \stackrel{\varepsilon \rightarrow 0}{=} \frac{\log(M/m)}{M-m} = \frac{\log(M/m)}{M} \sum_{n=0}^{\infty} \left(\frac{m}{M}\right)^n$$

When k is a *hard* scale $O(M)$

$$F_h(m, M, \varepsilon) = \int_0^\infty dk \frac{k^{-\varepsilon}}{(k+M)} \left[\frac{1}{k} - \frac{m}{k^2} + \dots \right] \stackrel{\varepsilon \rightarrow 0}{=} \left[-\frac{1}{\varepsilon M} + \frac{\log M}{M} \right] \sum_{n=0}^{\infty} \left(\frac{m}{M}\right)^n$$

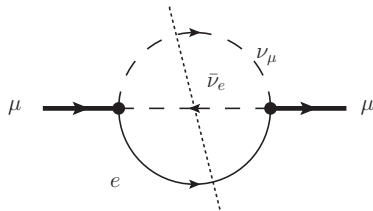
k is a *soft* scale $O(m)$

$$F_s(m, M, \varepsilon) = \int_0^\infty dk \frac{k^{-\varepsilon}}{(k+m)} \left[\frac{1}{M} - \frac{k}{M^2} + \dots \right] \stackrel{\varepsilon \rightarrow 0}{=} \left[\frac{1}{\varepsilon M} - \frac{\log m}{M} \right] \sum_{n=0}^{\infty} \left(\frac{m}{M}\right)^n$$

Muon lifetime at tree level

$$\begin{aligned}\Gamma_\mu &= \Gamma^{(hh)} + \Gamma^{(ss)} + \Gamma^{(hs)} + \Gamma^{(sh)} \\ &= \Gamma_0 \left[1 - 8\rho^2 - 12\rho^4 \log(\rho^2) + 8\rho^6 - \rho^8 \right]\end{aligned}$$

where $\rho = m_e/m_\mu$ and $\Gamma_0 = \frac{G_F^2 m_\mu^5}{192\pi^3}$

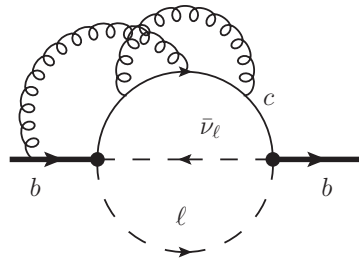


$$\Gamma^{(hh)} \sim -\frac{\rho^4}{12\varepsilon} - 24\rho^4 \log\left(\frac{\mu^2}{m_\mu^2}\right) + 1 - 8\rho^2 - 24\rho^4 + 16\rho^6 - 2\rho^8$$

$$\Gamma^{(hs)} \sim +\frac{\rho^4}{12\varepsilon} + 24\rho^4 \log\left(\frac{\mu^2}{m_\mu^2}\right) - 12\rho^4 \log(\rho^2) + 24\rho^4 - 8\rho^6 + \rho^8$$

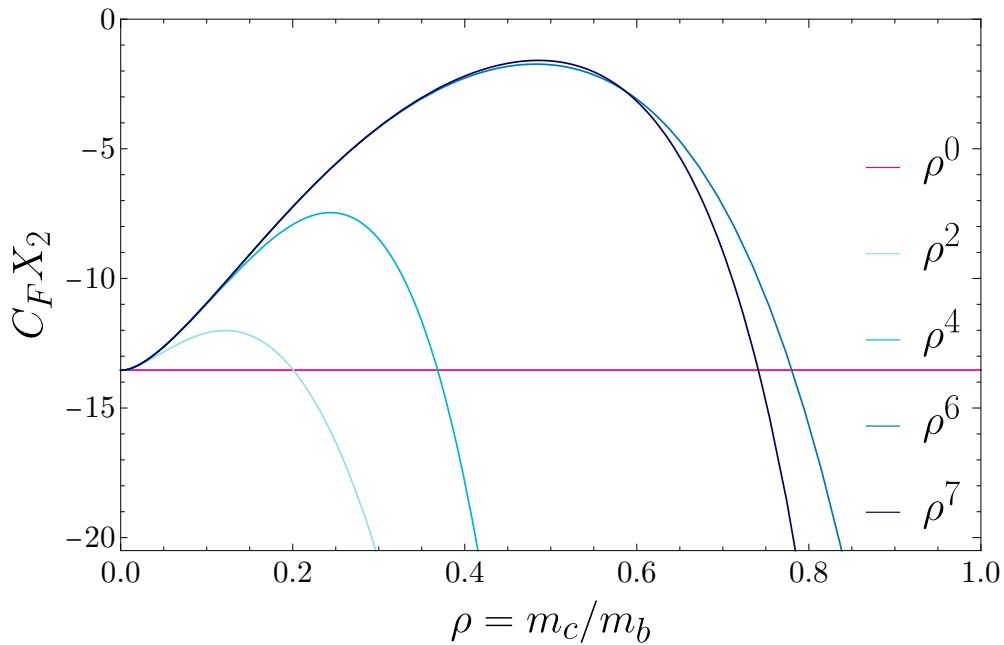
$$\Gamma^{(ss)} = \Gamma^{(sh)} = 0$$

$$\frac{\hbar}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} \left[X_0 + \frac{\alpha}{\pi} X_1 + \left(\frac{\alpha}{\pi}\right)^2 X_2 + \dots \right]$$



Czarnecki, Pak, PRD 78 (2008) 114015; PRL 100 (2008) 241807.

- Four-loop diagrams, all loop momenta can scale **hard** (m_μ) or **soft** (m_e).
- Each diagram considered up to **11 different regions**.
- The *all-hard* region reduces to **33 four-loop master integrals**.
- Expansion depth: $O(\rho^7)$ ($\rho = m_e/m_\mu$).



Towards the third order corrections

	α_s^2		α_s^3
n. diagrams	62	→	1450
n. loops	4	→	5
regions	11	→	O(20)
expansion depth	7	→	?
master integrals	33	→	?



The heavy daughter limit

Dowling, Piclum, Czarnecki, PRD 78 (2008) 074024

- Is the most natural expansion parameter sometimes also the best one?

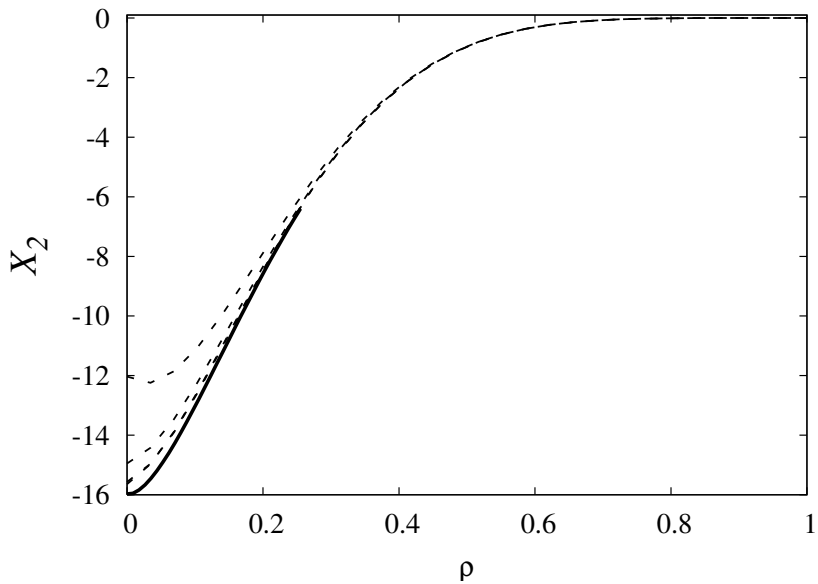
$$\frac{m_e}{m_\mu} \sim \frac{1}{210} \quad \frac{m_c}{m_b} \sim 0.3$$

- Perform the expansion in the limit $m_e \sim m_\mu$ ($m_c \sim m_b$):

$$\delta = 1 - \rho = 1 - \frac{m_e}{m_\mu} \ll 1$$

- The width must behave in the the $m_e \rightarrow m_\mu$ limit as:

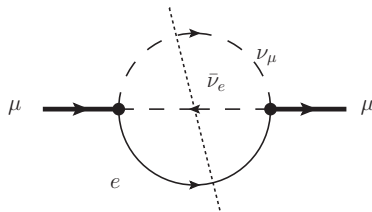
$$\Gamma_\mu \stackrel{m_e \rightarrow m_\mu}{\simeq} \frac{G_F^2}{192\pi^3} (m_\mu - m_e)^5 = \frac{G_F^2 m_\mu^5}{192\pi^3} \delta^5$$



Dowling, Piclum, Czarnecki, PRD 78 (2008) 074024

$$\begin{aligned}\Gamma_\mu &= \Gamma^{(\text{hh})} + \Gamma^{(\text{ss})} + \Gamma^{(\text{sh})} + \Gamma^{(\text{hs})} \\ &= \Gamma_0 \left[\frac{64}{5} \delta^5 - \frac{96}{5} \delta^6 + \frac{288}{35} \delta^7 + \dots \right]\end{aligned}$$

where $\rho = m_e/m_\mu$ and $\Gamma_0 = \frac{G_F^2 m_\mu^5}{192\pi^3}$



- **Crucial simplifications** in the heavy daughter limit!
- At least one electron's propagator must scale *soft* to generate an imaginary part $\log(-\delta)$.
- **Much smaller number of regions**, e.g.
 $\Gamma^{(\text{hh})} = \Gamma^{(\text{sh})} = \Gamma^{(\text{hs})} = 0$

- The heavy daughter limit converts **5 loops** \rightarrow **3 loops**!
- Loop integrals associated to ν and $\bar{\nu}$ momenta decouple and are integrated analytically at once.
- Explanation is rather technical, connected to the appearance of linear propagators:

$$\frac{1}{(p+k)^2 - m_c^2} \rightarrow \frac{1}{2p \cdot k - \delta}$$

	scaling	n. regions
$\mathcal{O}(\alpha)$	h, u	2
$\mathcal{O}(\alpha^2)$	hh , hu, uu	4
$\mathcal{O}(\alpha^3)$	hhh , uuu , huu , hhu	8

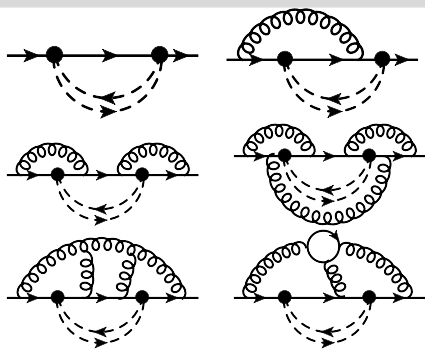
- 1450 five-loop diagrams.
- Several subtleties with FORM
 - Propagators expanded up to **10th - 12th order**.
 - Major obstacle is to keep as small as possible **the size of intermediate expressions**.
 - Efficient expansion of propagators and memory management.
- Intermediate FORM expressions up to $O(100)$ GB.
- Master integrals:
 - $O(\alpha^2)$: 3 (*ss*) and 3 (*hh*).
 - $O(\alpha^3)$: 20 (*sss*) and 19 (*hhh*).

Melnikov, van Ritbergen, Nucl.Phys.B 591 (2000) 515;
MF, Schönwald, Steinhauser, Phys.Rev.Lett. 125 (2020) 5.
- Renormalization constants at 3 loops with two massive fermions.

MF, Schönwald, Steinhauser, JHEP 10 (2020) 087.

$$\frac{\hbar}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} \left[X_0 + \sum_{n \geq 1} \left(\frac{\hat{\alpha}(m_\mu)}{\pi} \right)^n X_n \right]$$

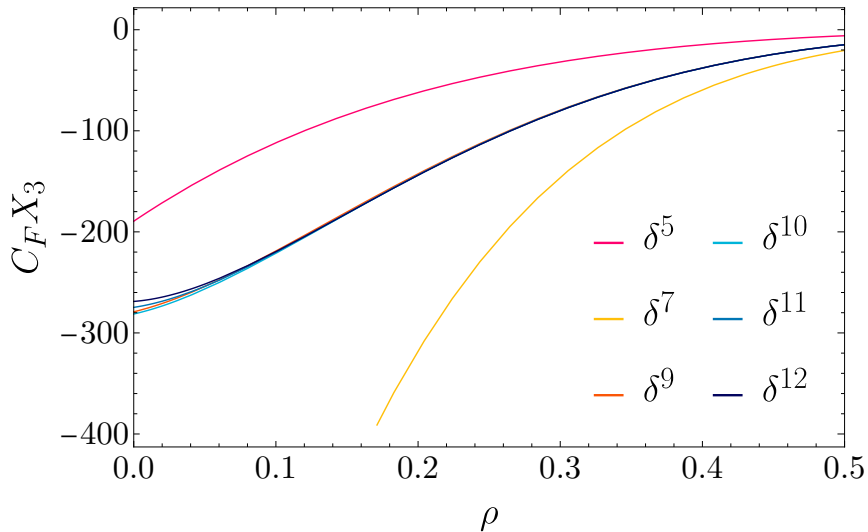
- NEW:** X_3 up to δ^{12} (first 8 terms).



$$X_3 = \sum_{m \geq 5} x_{3,m} \delta^m = \delta^5 \left[\frac{256}{3} a_4 - 128 \zeta(5) + \frac{56}{5} \pi^2 \zeta_3 + \frac{1984}{45} \zeta_3 + \frac{452}{675} \pi^4 - \frac{3652}{405} \pi^2 \right. \\ \left. - \frac{18451}{270} + \frac{32}{9} \log^4(2) - \frac{352}{45} \pi^2 \log^2(2) - \frac{32}{45} \pi^2 \log(2) \right] + O(\delta^6)$$

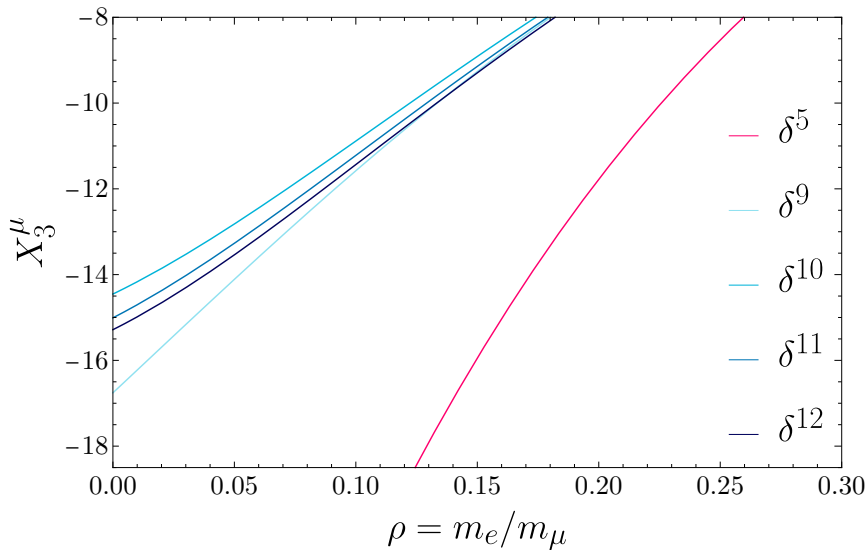
MF, Schönwald, Steihauser, Phys.Rev.D 104 (2021) 016003

confirmed C_F^3 , $C_F^2 N_H$ and $C_F N_H^2$ color factors up to δ^9 Czakon, Czarnecki, Dowling, Phys.Rev.D 103 (2021) L111301



$$C_F X_3(\rho = 0.28) = -91.2 \pm 0.4 \quad (0.4\%)$$

MF, Schönwald, Steinhauser, Phys.Rev.D 104 (2021) 1, 016003



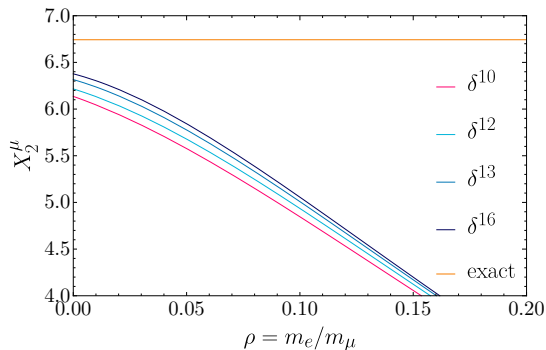
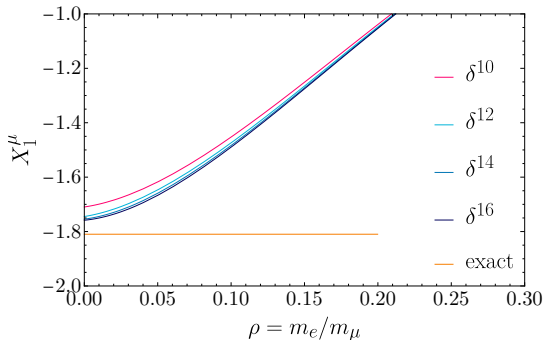
$$X_3^\mu = -15.3 \pm 2.3 \quad (15\%)$$

MF, Schönwald, Steinhauser, Phys.Rev.D 104 (2021) 1, 016003

Previous estimate: $X_3^\mu \simeq -20$

Ferrogli, Ossola, Sirlin, Nucl.Phys.B 560 (1999) 23

Theoretical uncertainties



order	X_n/X_{exact}	$x_{n,12}/X_n$
α	0.96	0.6 %
α^2	0.92	1.0%
α^3	~ 0.85	1.8 %

$$\frac{\hbar}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^2} F(\rho) (1 + \Delta q)$$

$$\Delta q = -(4\,234\,530|\hat{\alpha} + 36\,332|\hat{\alpha}^2 + 200|\hat{\alpha}^3 \pm 29|\delta\hat{\alpha}^3 \pm 11|\delta_{\text{had}}) \times 10^{-9}$$

- $\delta(1 + \Delta q) = 0.031 \text{ ppm}$

- $\frac{1}{2} \frac{\delta\tau_\mu}{\tau_\mu} = 0.5 \text{ ppm}$

- $\frac{5}{2} \frac{\delta m_\mu}{m_\mu} = 0.05 \text{ ppm}$

- at $O(\alpha^2)$: $\delta(1 + \Delta q) = 0.17 \text{ ppm}$

van Ritbergen, Stuart, Nucl.Phys.B 564 (2000) 343

Sirlin, Ferroglia, Rev.Mod.Phys. 85 (2013) 1

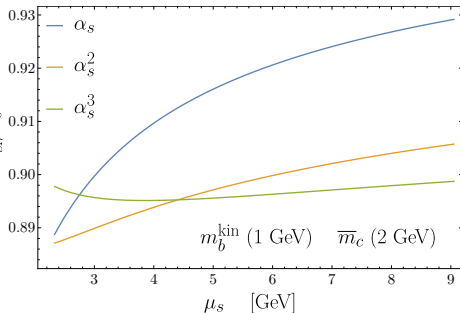
- Precise $\overline{\text{MS}}$ -on shell conversion of α
up to 4 loops

Baikov, Chetyrkin, Kuhn, Sturm, Nucl.Phys.B 867 (2013) 182

Lessons for $B \rightarrow X_{cl}\nu$

$$\Gamma_{sl} = \frac{G_F^2 m_b^5 |V_{cb}|^2}{192\pi^3} F(\rho) \left[1 + \sum_n Y_n \left(\frac{\alpha_s(m_b)}{\pi} \right)^n \right] \Gamma_{sl}/\Gamma_0$$

- n=1 Jezabek, Kühn, Jezabek, Kuhn, NPB 314 (1989) 1
- n=2 Melnikov, PLB 666 (2008) 336; Pak, Czarnecki, PRD 78 (2008) 114015.
- n=3 Fael, Schönwald, Steinhauser, hep-ph/2011.13654



$$m_b^{\text{OS}} : m_c^{\text{OS}} \quad 1 - 1.78 \left(\frac{\alpha_s}{\pi} \right) - 13.1 \left(\frac{\alpha_s}{\pi} \right)^2 - 163.3 \left(\frac{\alpha_s}{\pi} \right)^3$$

$$\bar{m}_b(\bar{m}_b) : \bar{m}_c(3 \text{ GeV}) \quad 1 + 3.07 \left(\frac{\alpha_s}{\pi} \right) + 13.3 \left(\frac{\alpha_s}{\pi} \right)^2 + 62.7 \left(\frac{\alpha_s}{\pi} \right)^3$$

$$m_b^{\text{kin}}(1 \text{ GeV}) : \bar{m}_c(2 \text{ GeV}) \quad 1 - 1.24 \left(\frac{\alpha_s}{\pi} \right) - 3.65 \left(\frac{\alpha_s}{\pi} \right)^2 - 1.0 \left(\frac{\alpha_s}{\pi} \right)^3$$

- Fit BR and moments from B factories

- Global fit strategy from 2014

Gambino, Schwanda, *Phys.Rev.D* 89 (2014) 014022

Alberti, Gambino, Healey, Nandi, *Phys.Rev.Lett.* 114 (2015) 6, 061802

- $O(\alpha_s^3)$ semileptonic width

- $O(\alpha_s^3)$ relation between $\bar{m}_b - m_b^{\text{kin}}$

MF, Schönwald, Seinhauser, *Phys.Rev.Lett.* 125 (2020) 052003;

Phys.Rev.D 103 (2021) 1, 014005

- Precise input from lattice

$$\bar{m}_c(3 \text{ GeV}) = 0.988 (7) \text{ GeV}$$

$$\bar{m}_b(\bar{m}_b) = 4.198 (12) \text{ GeV}$$

$$\longrightarrow m_b^{\text{kin}} = 4.56 (19) \text{ GeV}$$

FLAG2019

$$\begin{aligned} |V_{cb}| &= 42.16 (30)_{\text{th}} (32)_{\text{exp}} (25)_{\Gamma} \times 10^{-3} \\ &= 42.16 (50) \times 10^{-3} \end{aligned}$$

Bordone, Capdevila, Gambino, [hep-ph/2107.00604](https://arxiv.org/abs/hep-ph/2107.00604)

error improvement of 34% compared to 2014.

- QED $O(\alpha^3)$ to the muon lifetime
- QCD α_s^3 corrections to $\Gamma(b \rightarrow X_c \ell \nu)$.
- Heavy daughter limit: small parameter $\delta = 1 - m_e/m_\mu$ and $\delta = 1 - m_c/m_b$.
- $\Delta q^{(3)}$ with relative 15% uncertainty. Theory error on Δq reduced to 0.03 ppm.
- 1% theory uncertainty in semileptonic decays width.
- Improved extraction of $|V_{cb}|$ inclusive by about 34%.