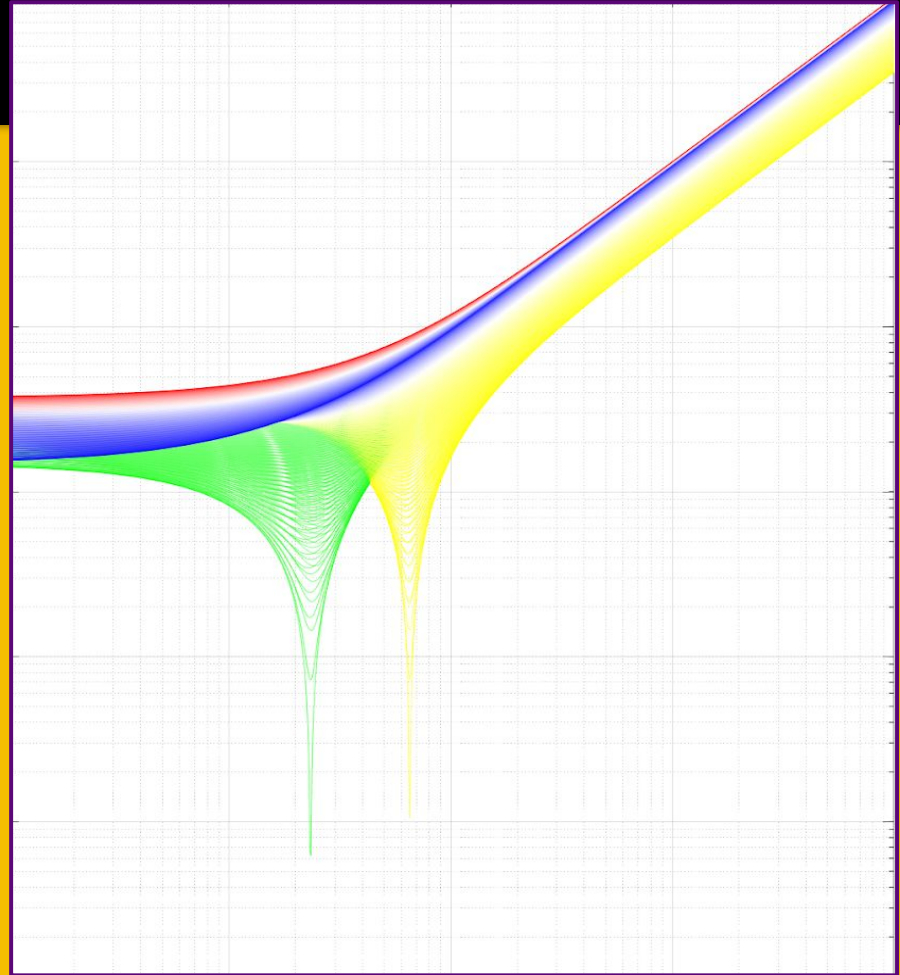


# Probing $0\nu\beta\beta$ Decay in Multiple Isotopes

**Graham Van Goffrier**

**In collaboration with M. Agostini and F. Deppisch  
University College London**

**[ucapgwg@ucl.ac.uk](mailto:ucapgwg@ucl.ac.uk)**

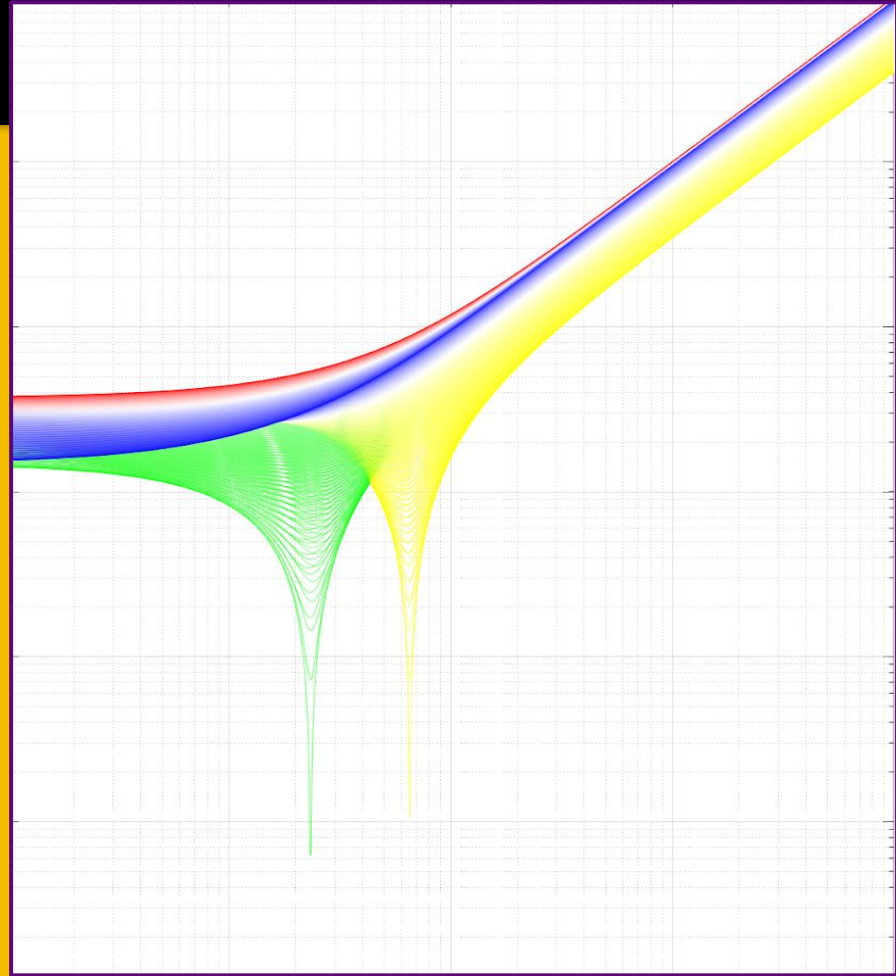


# Probing $0\nu\beta\beta$ Decay in Multiple Isotopes: Quo Vadis?

Graham Van Goffrier

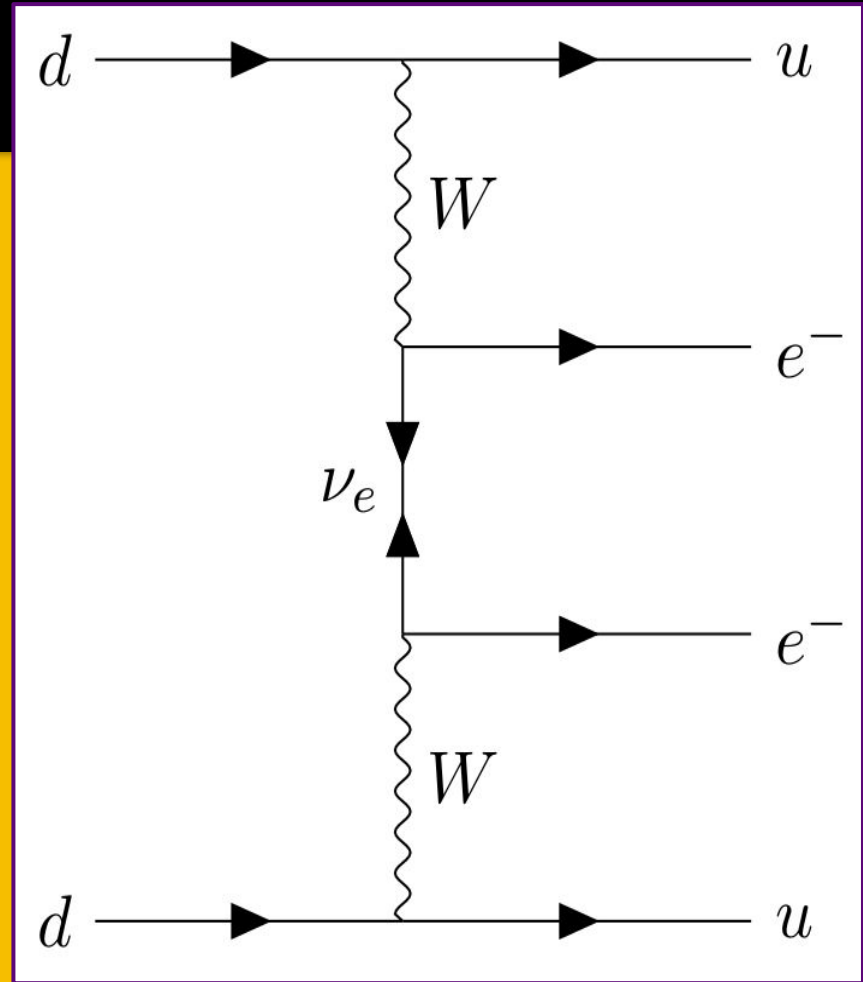
In collaboration with M. Agostini and F. Deppisch  
University College London

[ucapgwg@ucl.ac.uk](mailto:ucapgwg@ucl.ac.uk)



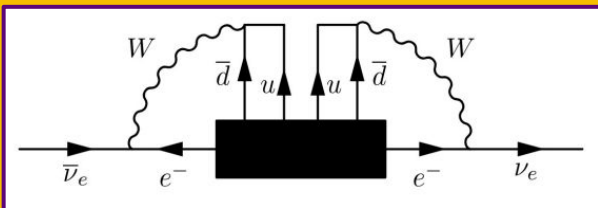
# Neutrinos and $0\nu\beta\beta$

- Flavour oscillation  $\Rightarrow$  massive neutrinos
- Cosmology, oscillations,  $\beta$ -decay provide partial answers and constraints.
- $0\nu\beta\beta$  -- just another probe? No!
  - Sensitive to LNV
  - May differentiate mass mechanisms
- Current bounds (@90%)
  - GERDA:  $T_{1/2} > 1.8 \cdot 10^{26}$  y [1]
  - KamLAND-Zen:  $T_{1/2} > 2.3 \cdot 10^{26}$  y [2]



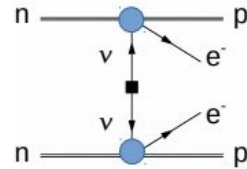
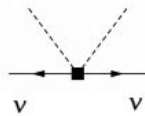
# $0\nu\beta\beta$ Mechanisms

- $\Delta L=2$  from odd-dim SM-EFT operators [Kobach, 2016]
  - 5,7 classified, 9 partially
- $\Lambda_{LNV} \gg \Lambda_{EW} \Rightarrow$  5-dim only
- Schechter-Valle Theorem:
  - $\Delta L=2 \Rightarrow$  Majorana  $\nu$  mass

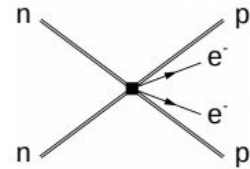


[Schechter + Valle, 1982]

Dim 5

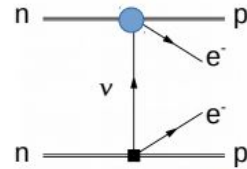
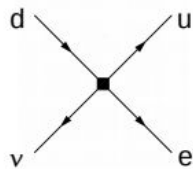


5a)

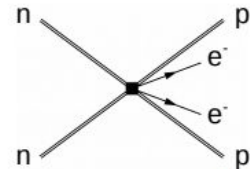


5b)

Dim 7

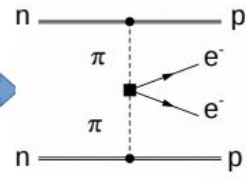
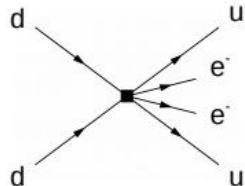


7a)

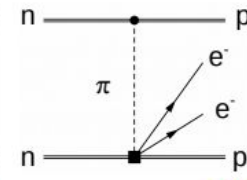


7b)

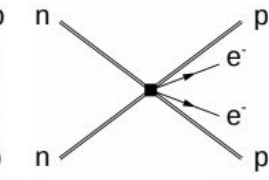
Dim 9



9a)



9b)



9c)

[Engel et al., 2022]

# Searching for Short-Range Mechanisms

- **Question:** To what extent can multi-isotope observations constrain exotic mechanisms?
  - What is the impact of (correlated) NME uncertainties, and how to fold into analysis?
- Numerous parameterised mechanisms [Deppisch et al., Phys. Rev. D 102, 095016 (2020)]
- We consider a single exotic  $0\nu\beta\beta$  mechanism with its own “heavy NME” (a la RPV SUSY):

$$T_{1/2}^{-1} = G_{iso} \left| \frac{m_{\beta\beta}}{m_e} M_{\nu,iso} + \epsilon M_{H,iso} \right|^2$$

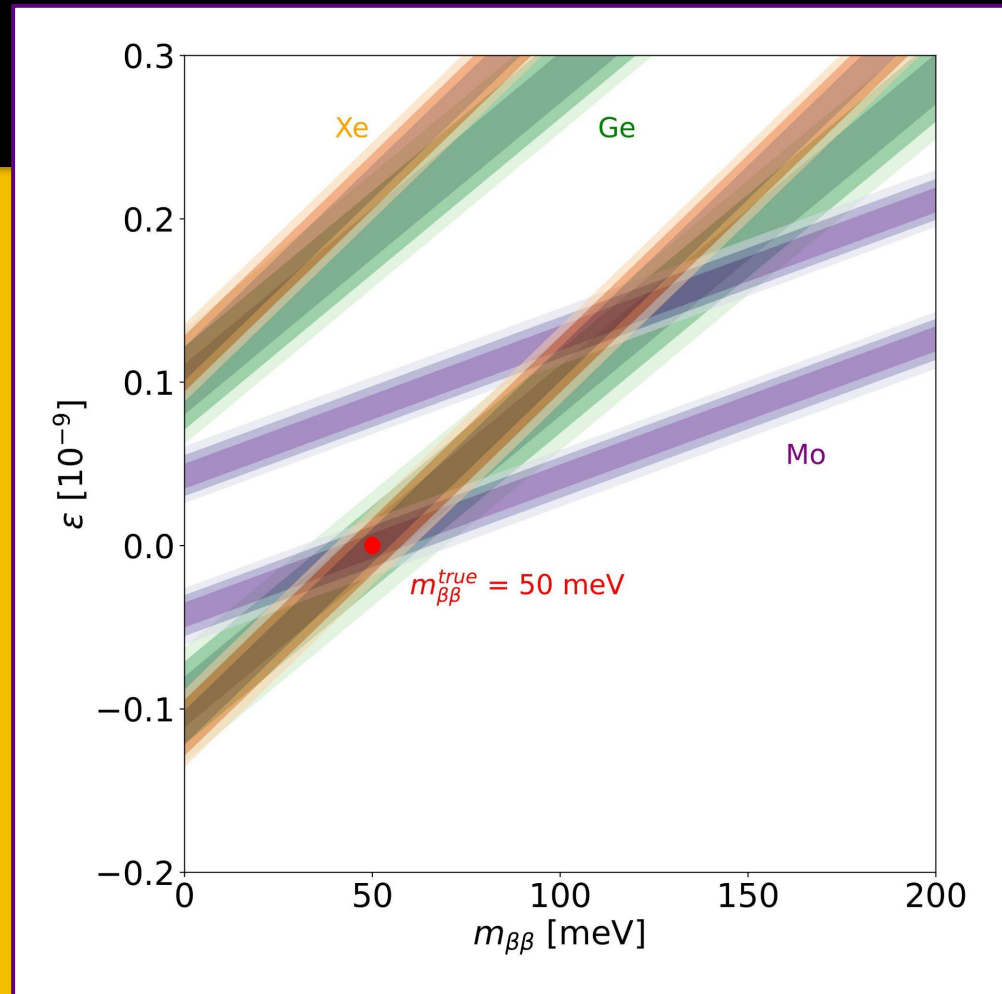
- Any fixed  $0\nu\beta\beta$  half-life corresponds to two parallel-line solutions in  $\{m_{\beta\beta}, \epsilon\}$

# Ratio Degeneracy (or not)

- Previous thinking: light/exotic NME ratio strongly degenerate across isotopes
  - QRPA in Ge/Se/Mo/Te [Faessler et al., 2011], Ge/Te/Xe [Lisi et al., 2015]
  - Big problem for multi-isotope analysis
- However, nuclear structure correlations can break this degeneracy
  - Correlations suppress NME at  $<2$  fm (in NSM/EDF) [Menéndez, 2018]
  - Therefore contribute less to exotic than to light nu exchange
  - Also verifies [F, 2011][L, 2015] by deselecting correlated nucleon states
- Degeneracy broken in IBM-2 for five distinct exotic mechanisms [Deppisch et al., 2020]
  - Strongest deviation from degeneracy in  $^{100}\text{Mo}$

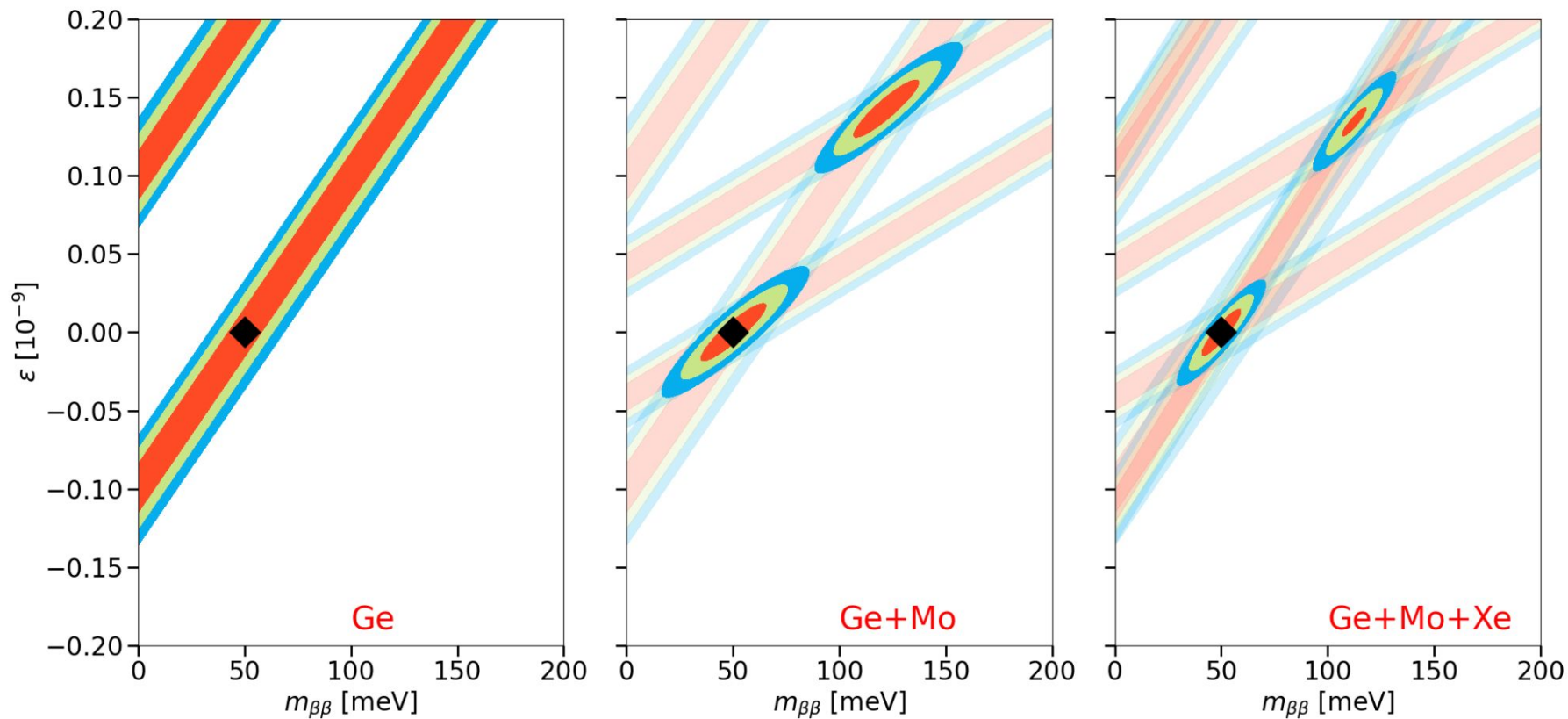
# Single Isotopes

- Proposed next-gen searches:
  - $^{76}\text{Ge}$ : LEGEND-1000
  - $^{100}\text{Mo}$ : CUPID
  - $^{136}\text{Xe}$ : nEXO
- Ratio degeneracy controls overlaps of per-isotope likelihoods
- Two-isotopes  $\Rightarrow$  secondary peak
  - Third suppresses this peak



# $\chi^2$ Analysis for Ge, Ge+Mo, Ge+Mo+Xe

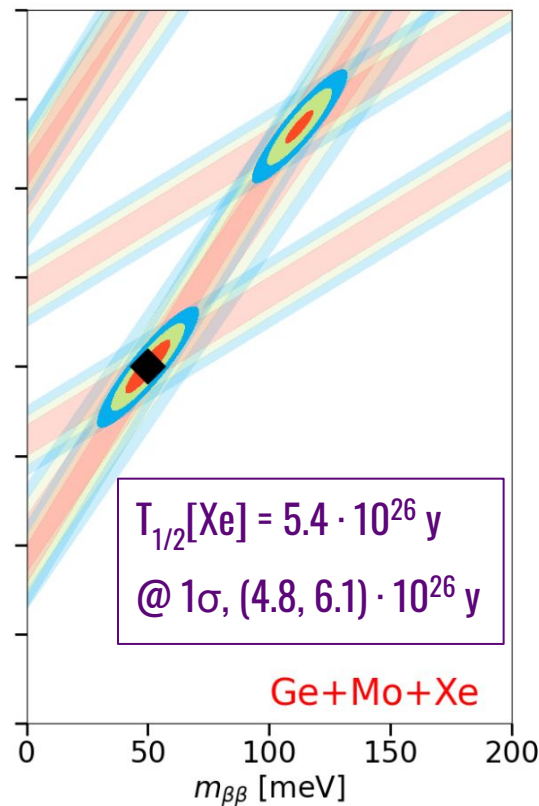
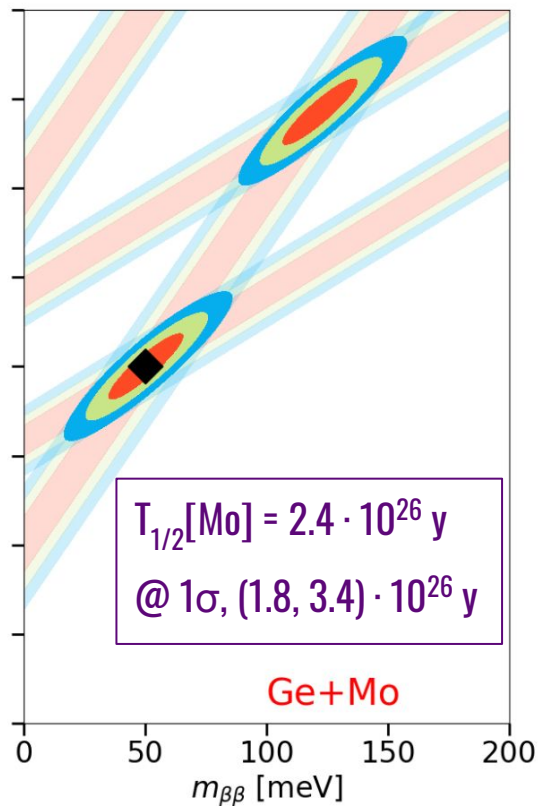
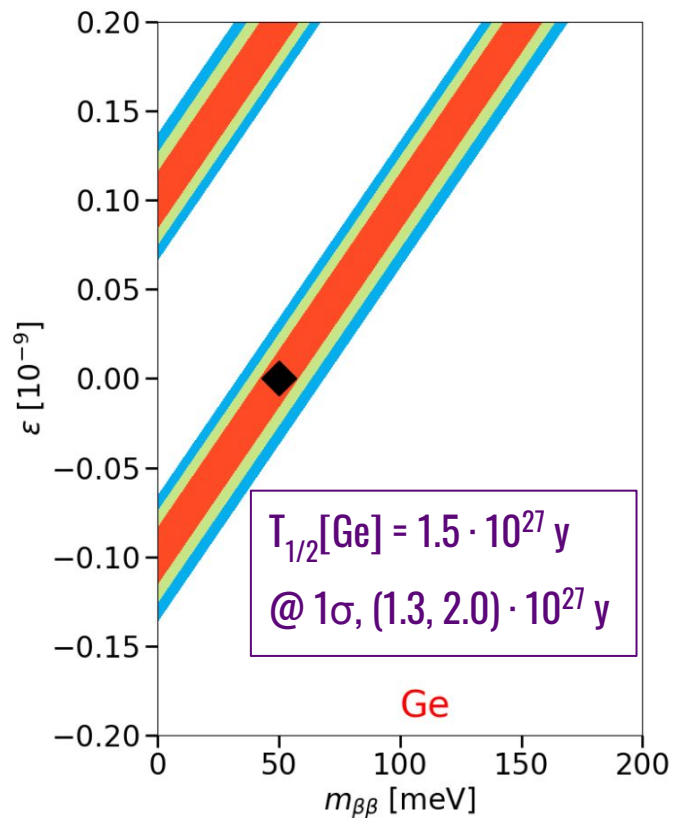
8





# $\chi^2$ Analysis for Ge, Ge+Mo, Ge+Mo+Xe

9



# Bayesian Methodology

- Update prior knowledge  $\pi(\boldsymbol{\theta})$  with likelihood  $L_{\mathbf{X}}(\boldsymbol{\theta})$  to obtain posterior knowledge  $\mathbf{p}(\boldsymbol{\theta})$ .

$$p(\boldsymbol{\theta}) = \frac{L_{\mathbf{X}}(\boldsymbol{\theta})\pi(\boldsymbol{\theta})}{\int L_{\mathbf{X}}(\boldsymbol{\theta}')\pi(\boldsymbol{\theta}')d\boldsymbol{\theta}'_H} \equiv \frac{L_{\mathbf{X}}(\boldsymbol{\theta})\pi(\boldsymbol{\theta})}{M_{\mathbf{X}}^H}$$

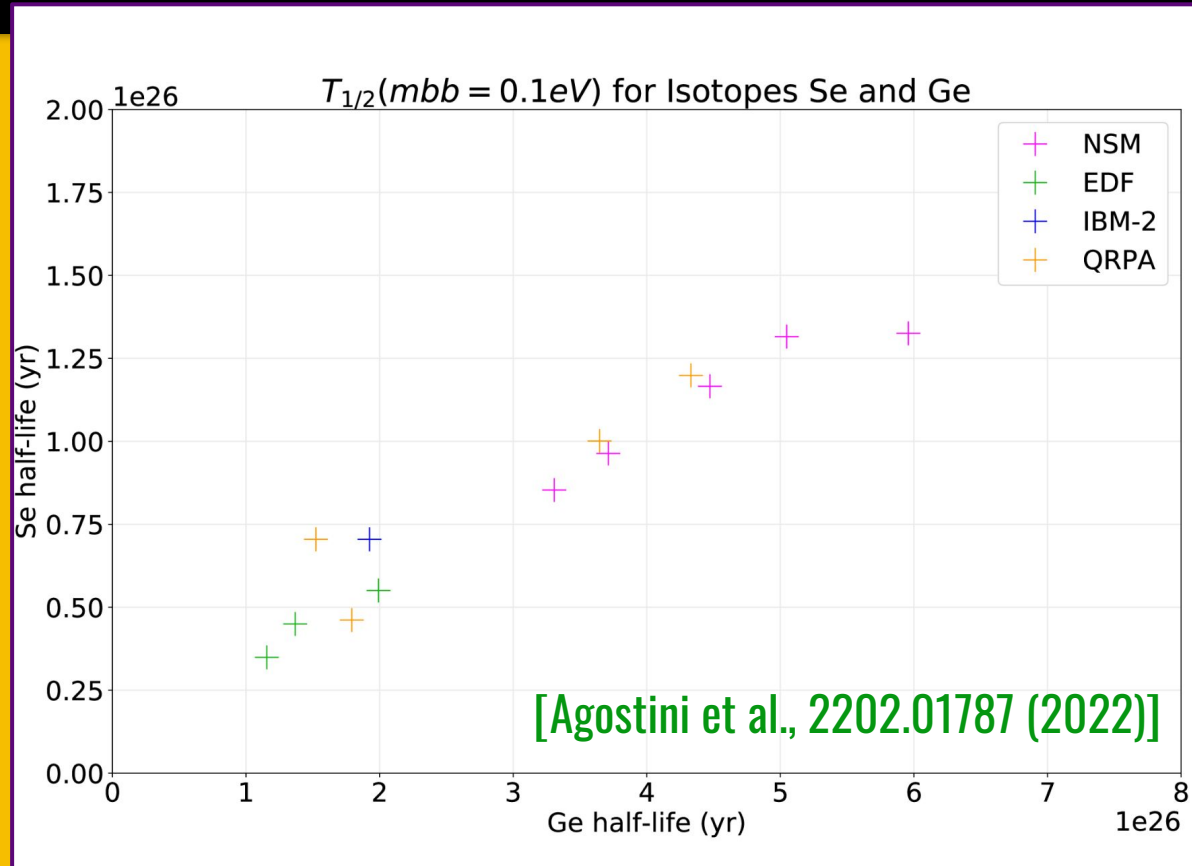
- Markov Chain Monte Carlo (MCMC) samples  $\mathbf{p}(\boldsymbol{\theta})$  using only local distributions.
  - Locality helps to combat MC rejection problems at high-dimensionalities.
- Our parameter space:  $\{m_{\beta\beta}, \varepsilon\} + \{M_{\nu}, M_H\}$  for each isotope
  - Flat priors on  $m_{\beta\beta}, \varepsilon$  [Deppisch, *GVG Phys. Rev. D* 104, 055040 (2021)]
  - What priors to use on  $\{M_{\nu}, M_H\}$ , given sample estimates?

# Nuclear Matrix Elements (NMEs)

- Recall  $m_{\beta\beta}$  and half-life:

$$T_{1/2}^{-1} = G_{0\nu} |\mathbb{M}|^2 m_{\beta\beta}^2$$

- NMEs discrepancies are correlated for (some) pairs of isotopes.
- Source: systematic uncertainties across many-body methods

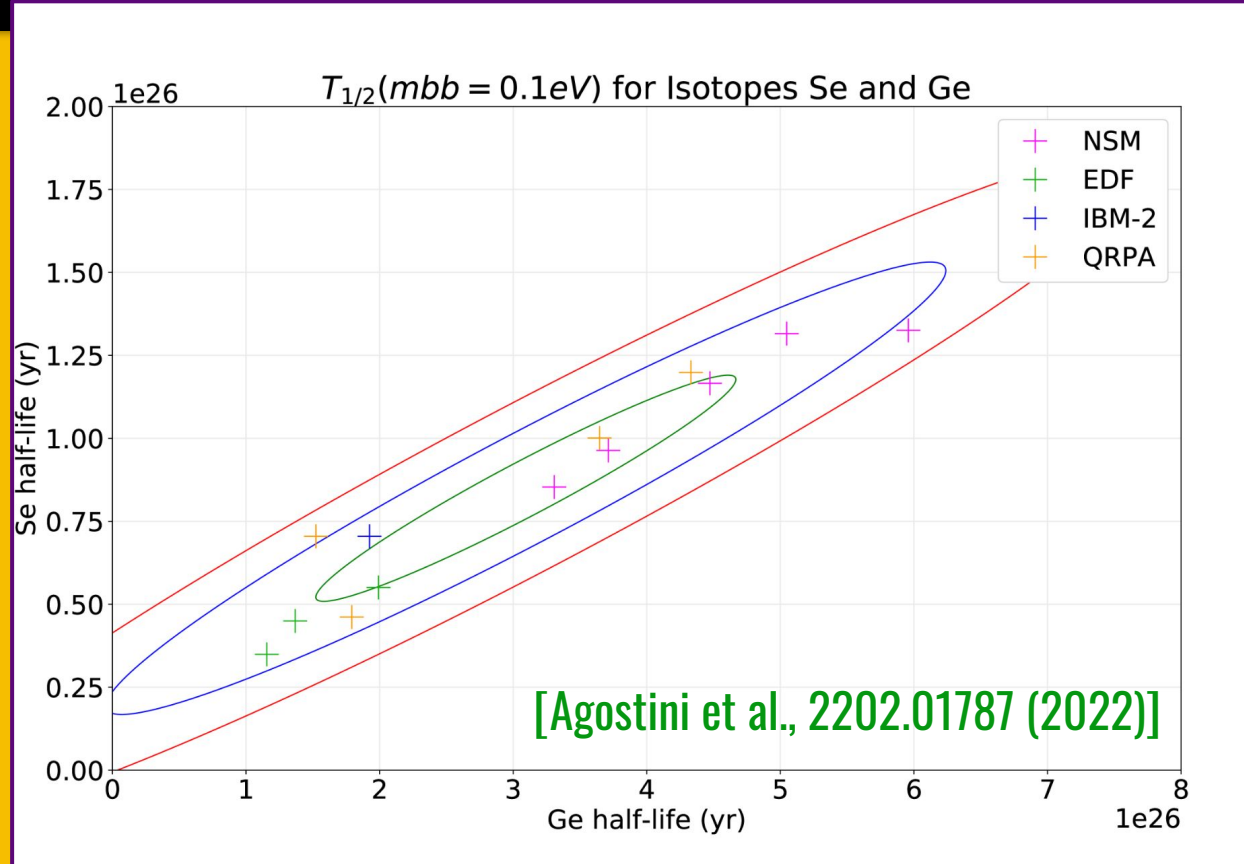


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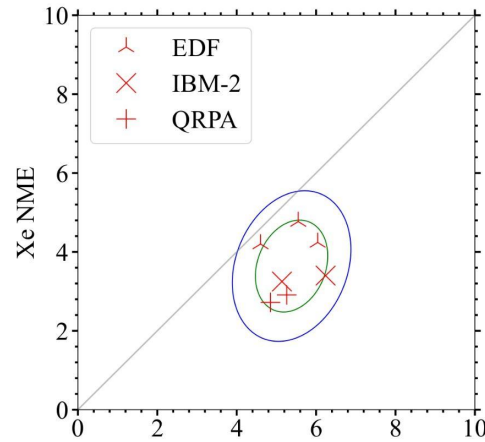
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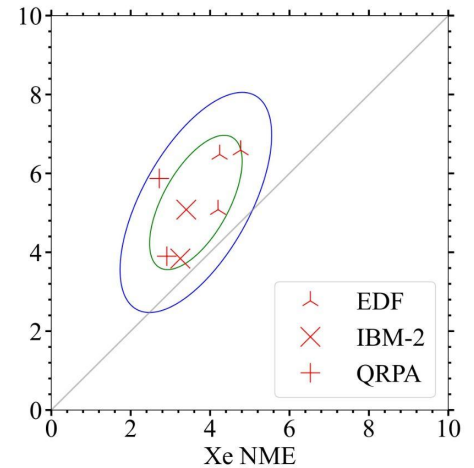
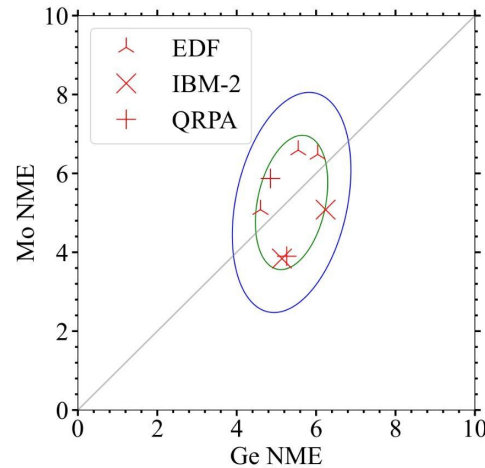
# NME Priors

- Multinormal fit on three methods
  - Seven independent estimates
- Placeholder for concrete estimates of theoretical uncertainties

$$\mu = \begin{pmatrix} 5.383 \\ 5.263 \\ 3.641 \end{pmatrix}, \quad \Sigma = \begin{pmatrix} 0.361 & 0.200 & 0.102 \\ 0.200 & 1.260 & 0.527 \\ 0.102 & 0.527 & 0.590 \end{pmatrix}$$



[Agostini et al., JHEP 172 (2023)]

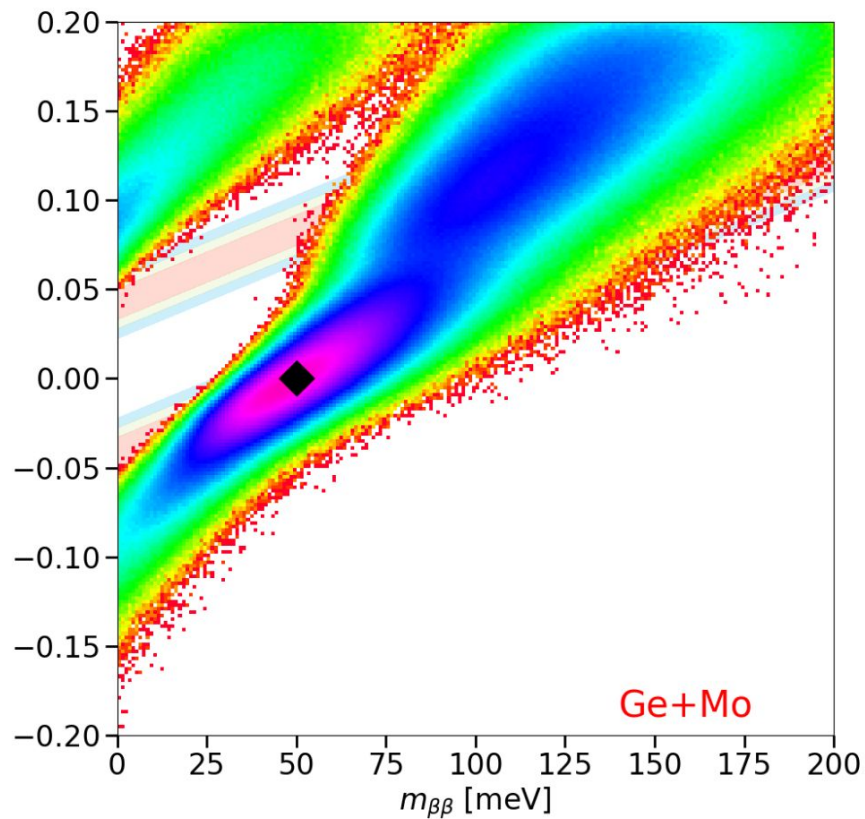
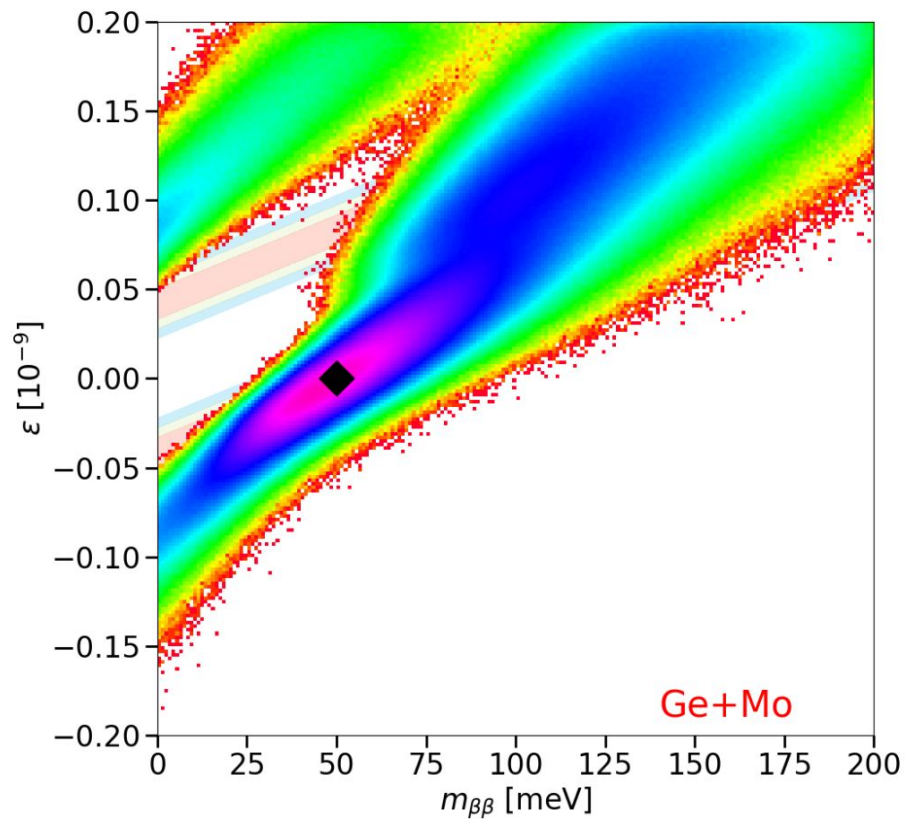


# Uncorrelated

# vs.

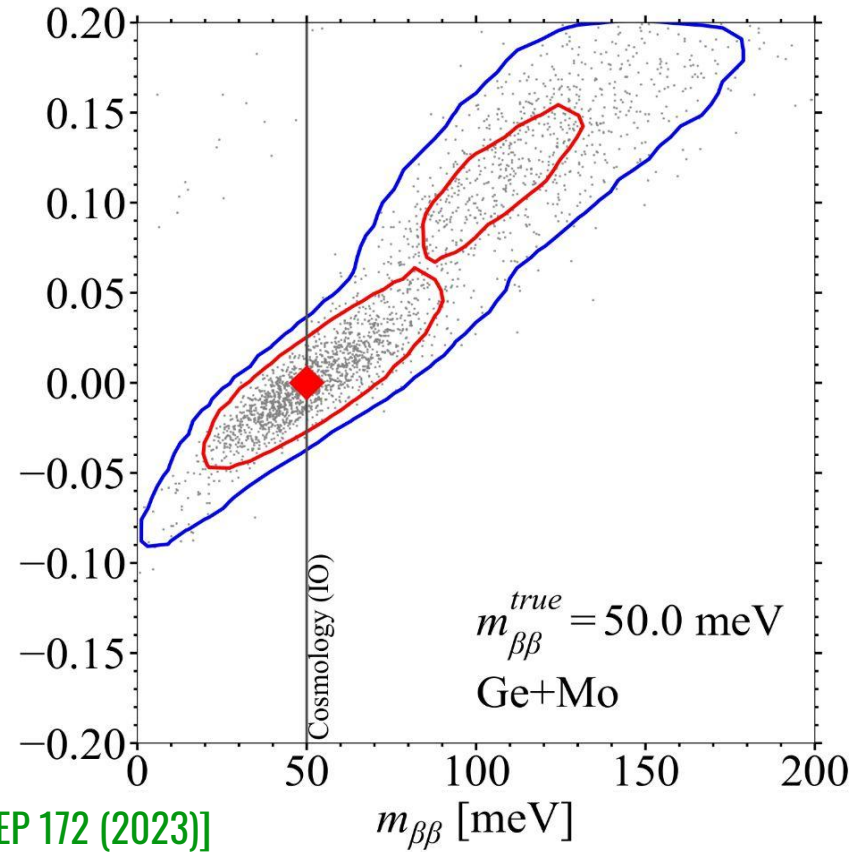
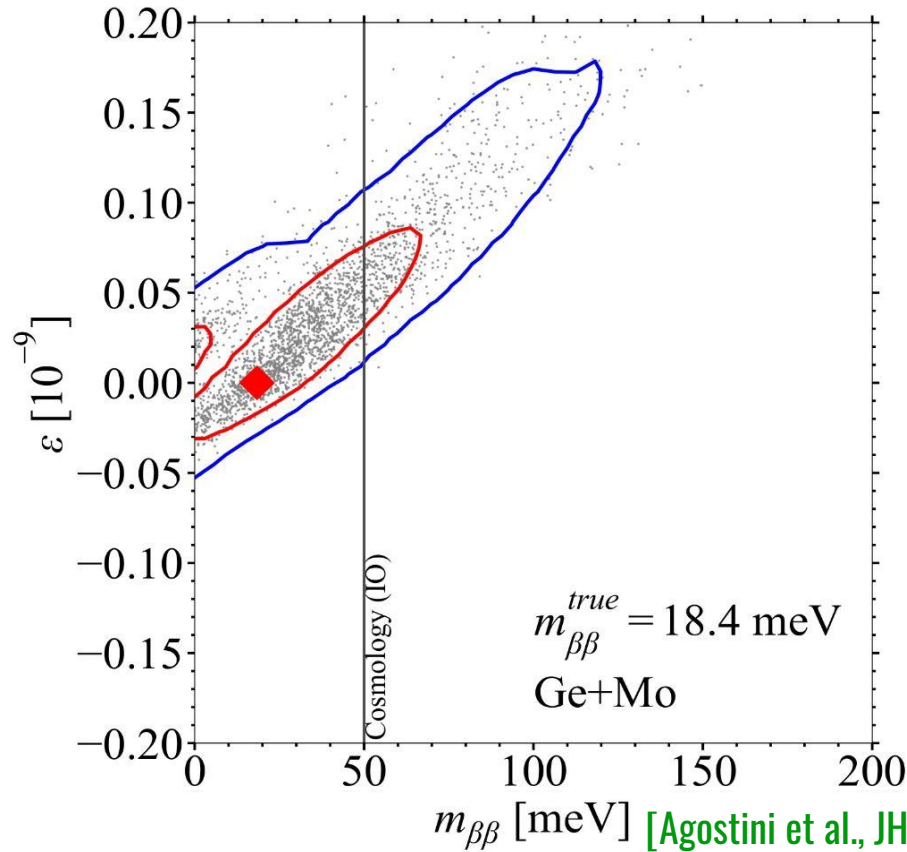
# Correlated

14



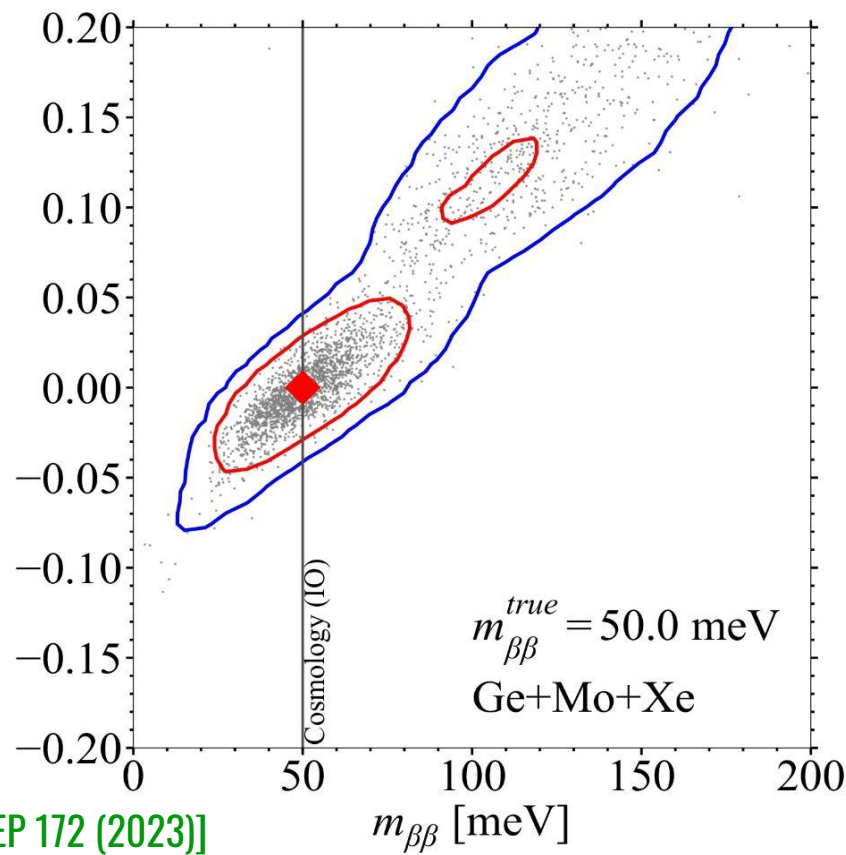
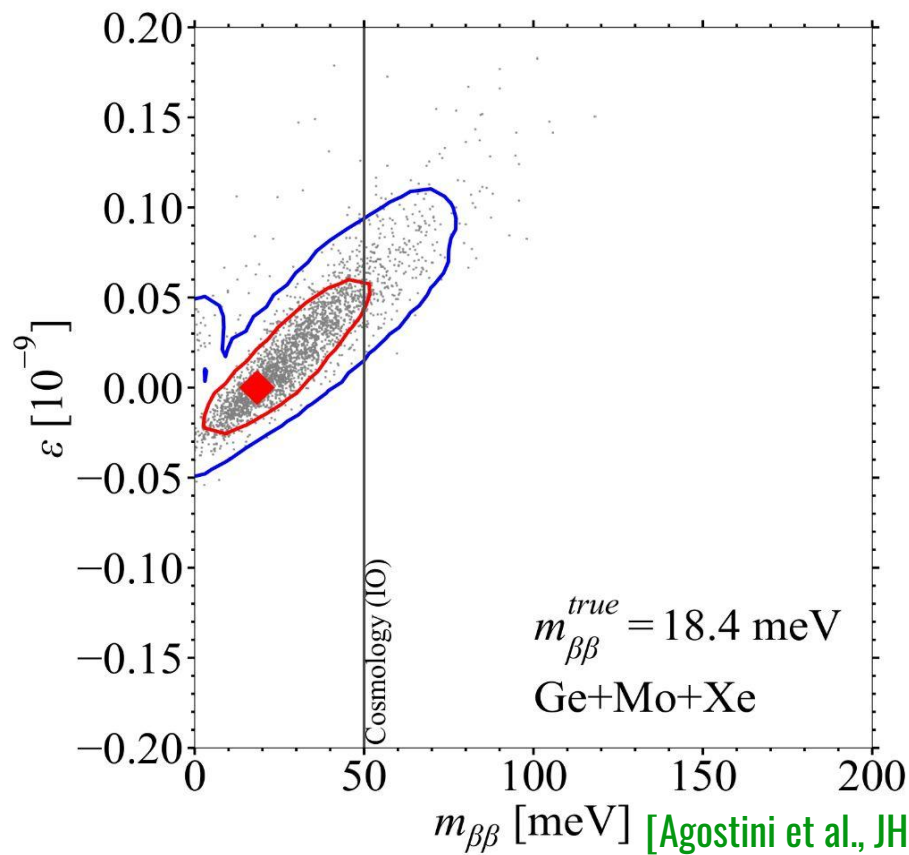
# Ge+Mo,

# Current NMEs



# Ge+Mo+Xe,

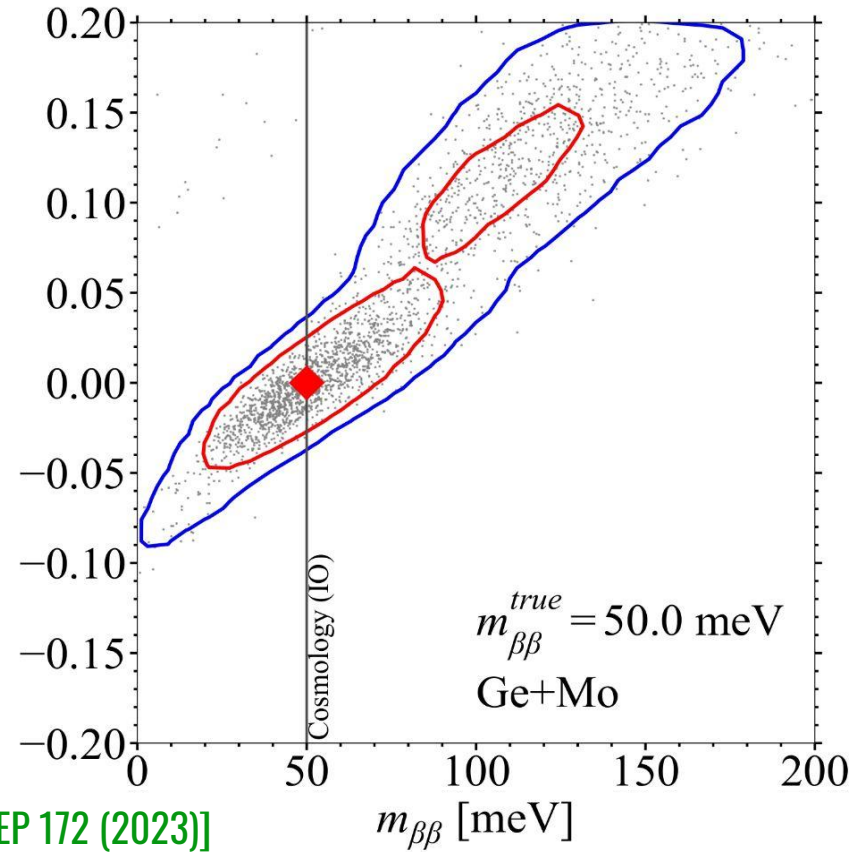
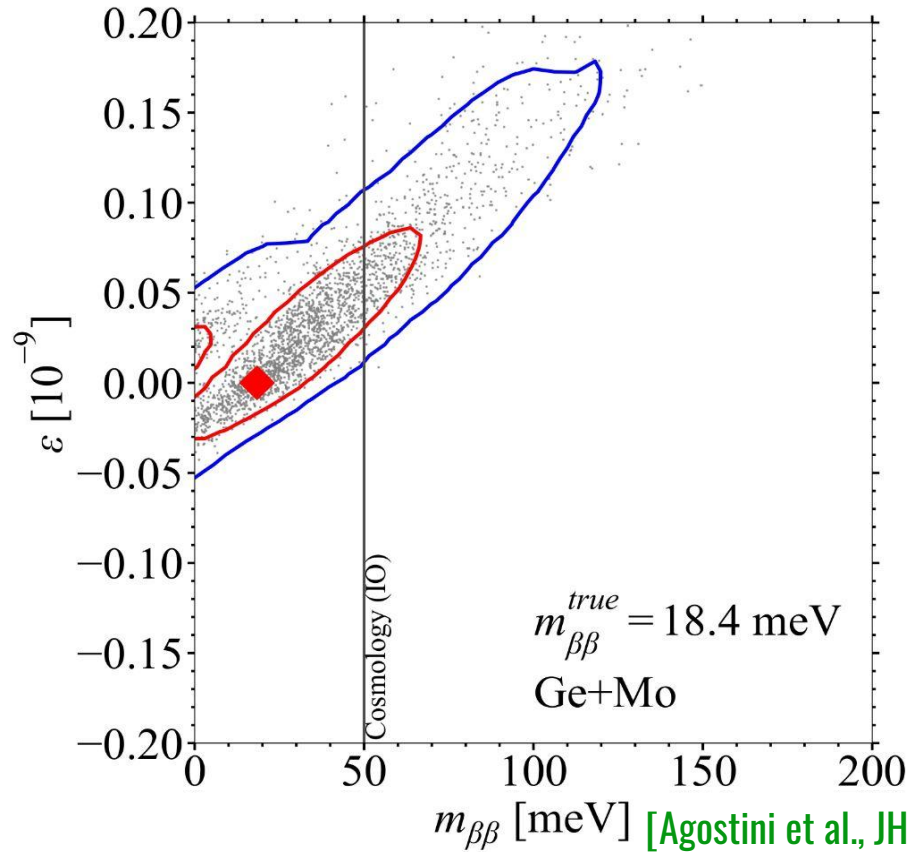
# Current NMEs





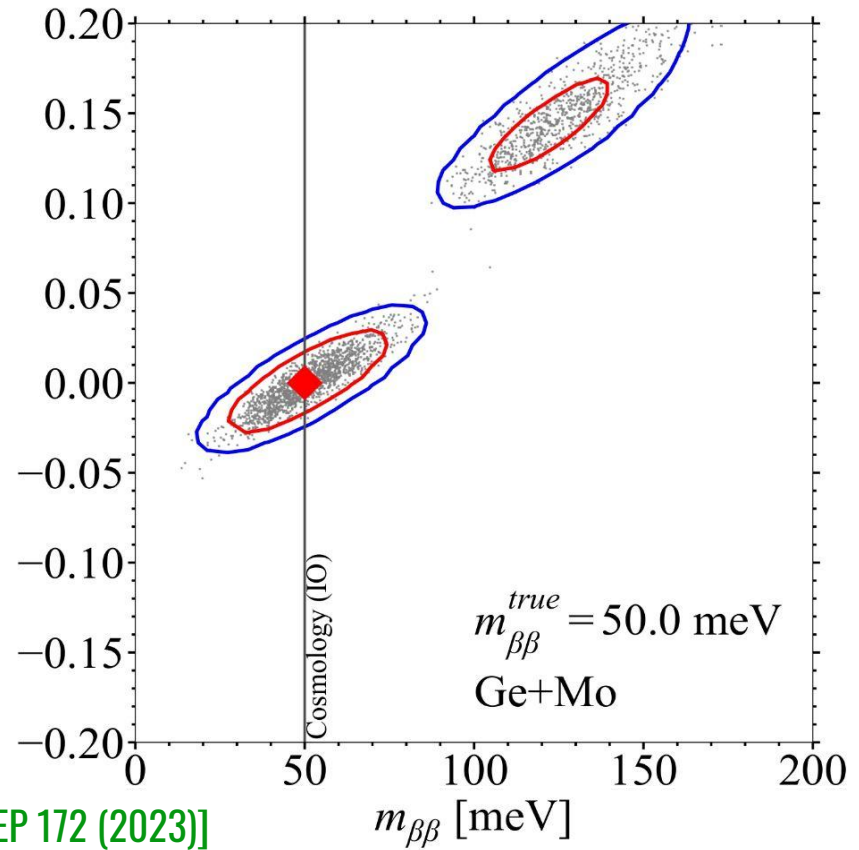
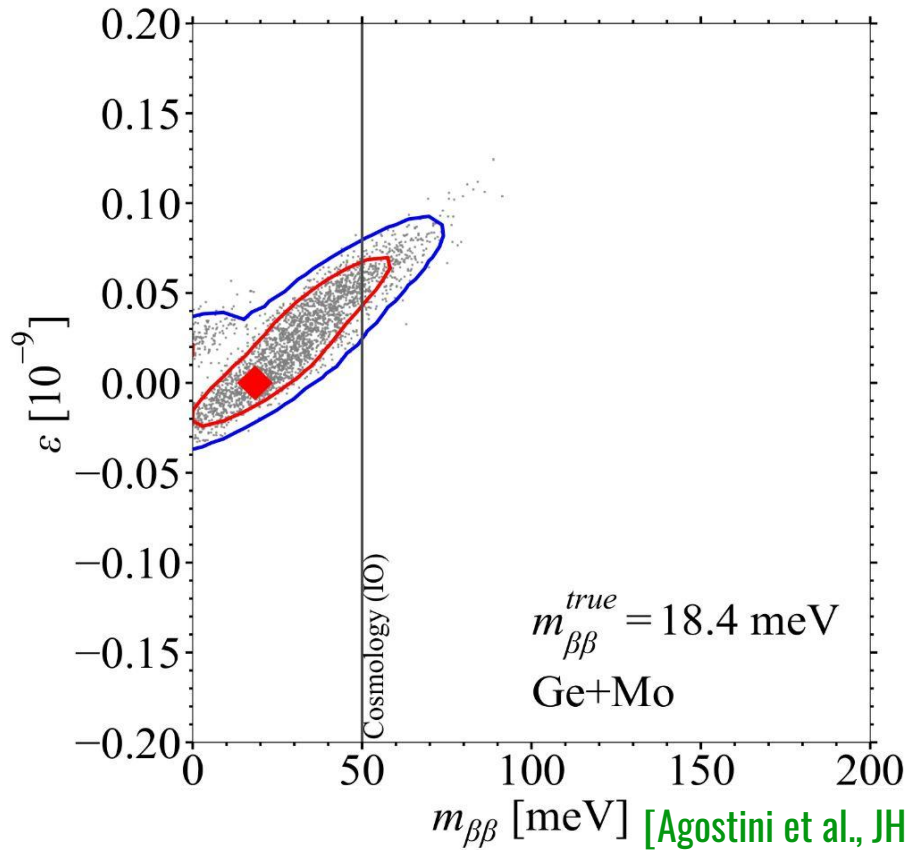
Ge+Mo,

Current NMEs



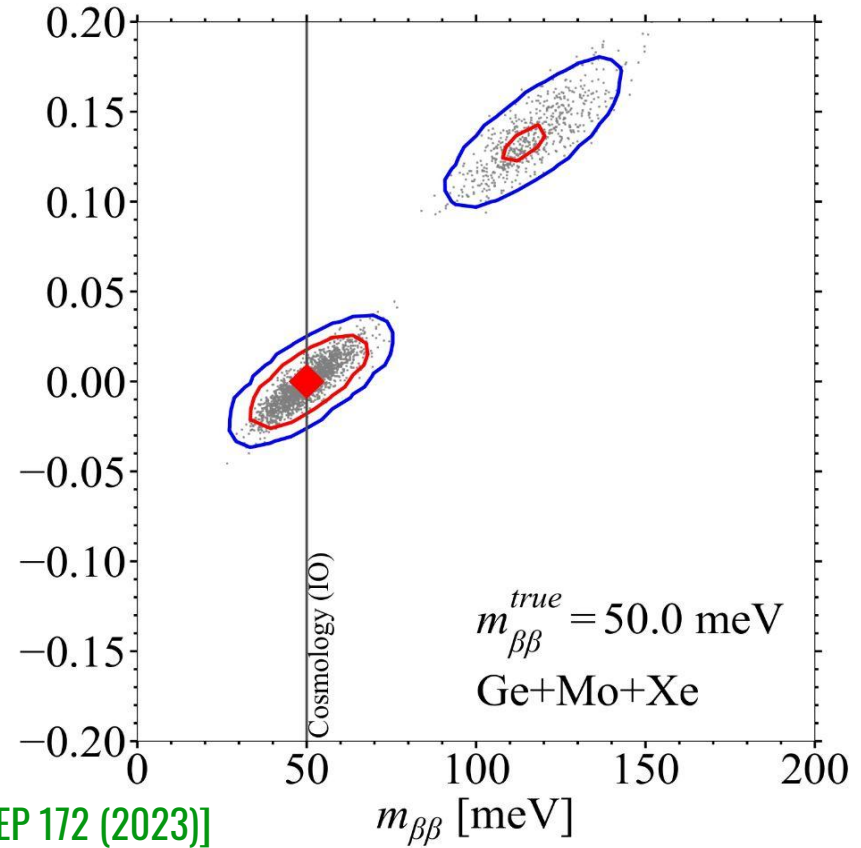
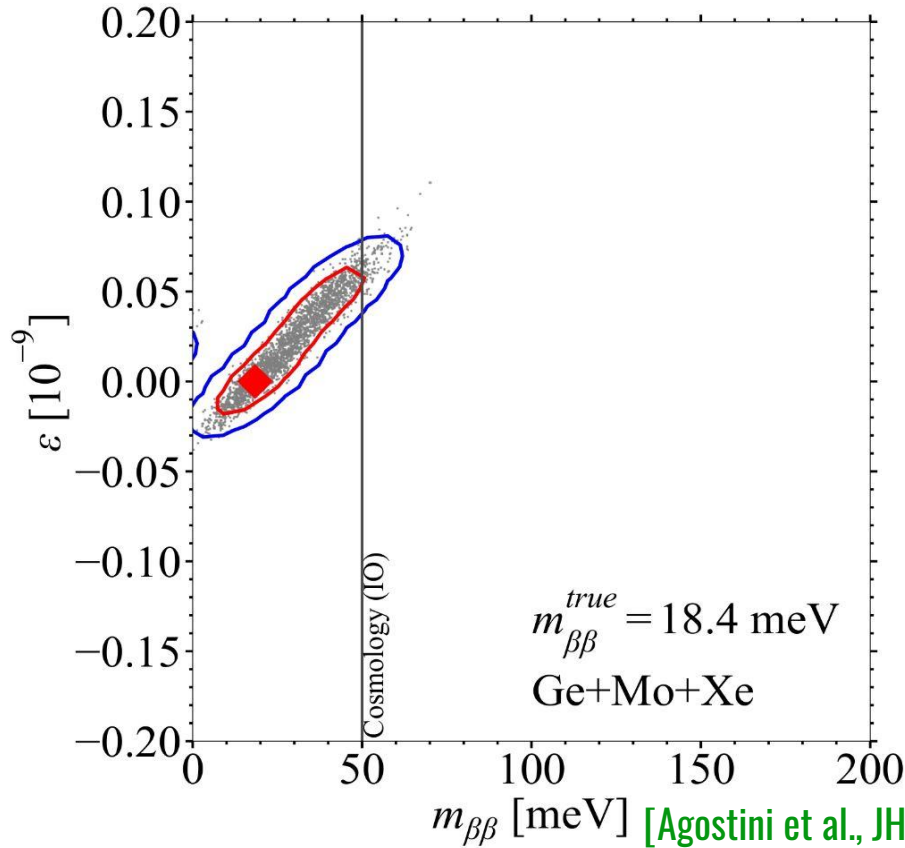
# Ge+Mo,

# Future NMEs



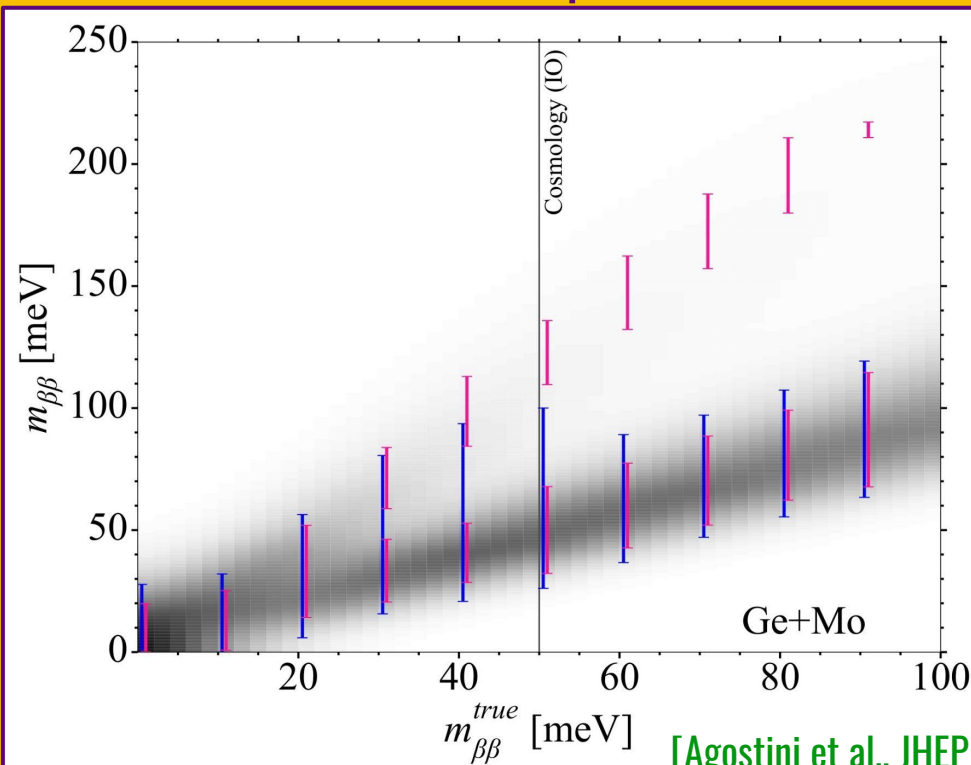
# Ge+Mo+Xe,

# Future NMEs

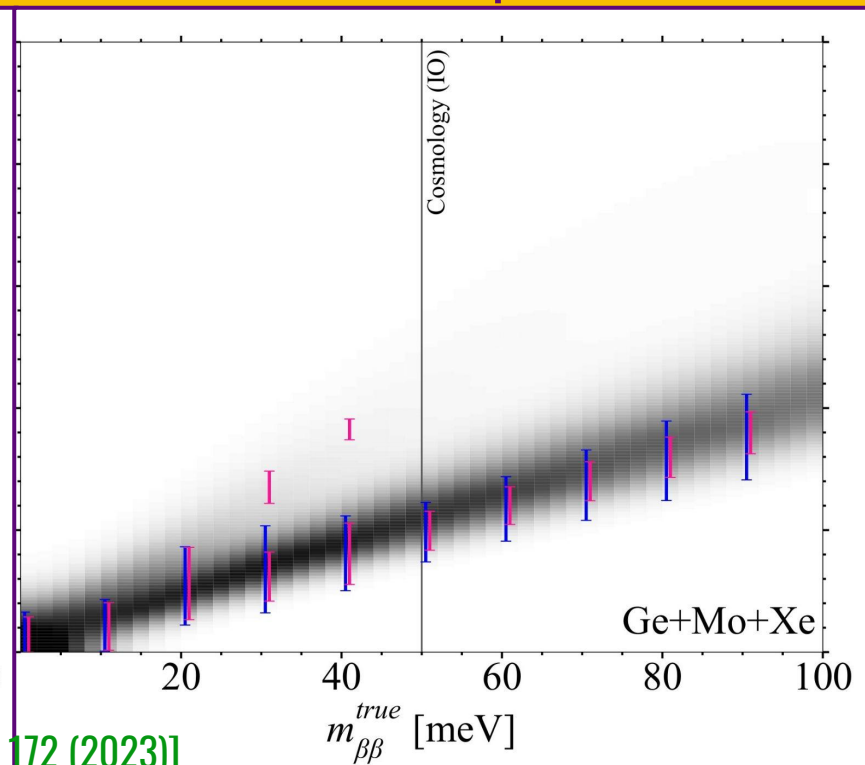


# $m_{\beta\beta}$ - Marginalized Statistics

## Two isotopes



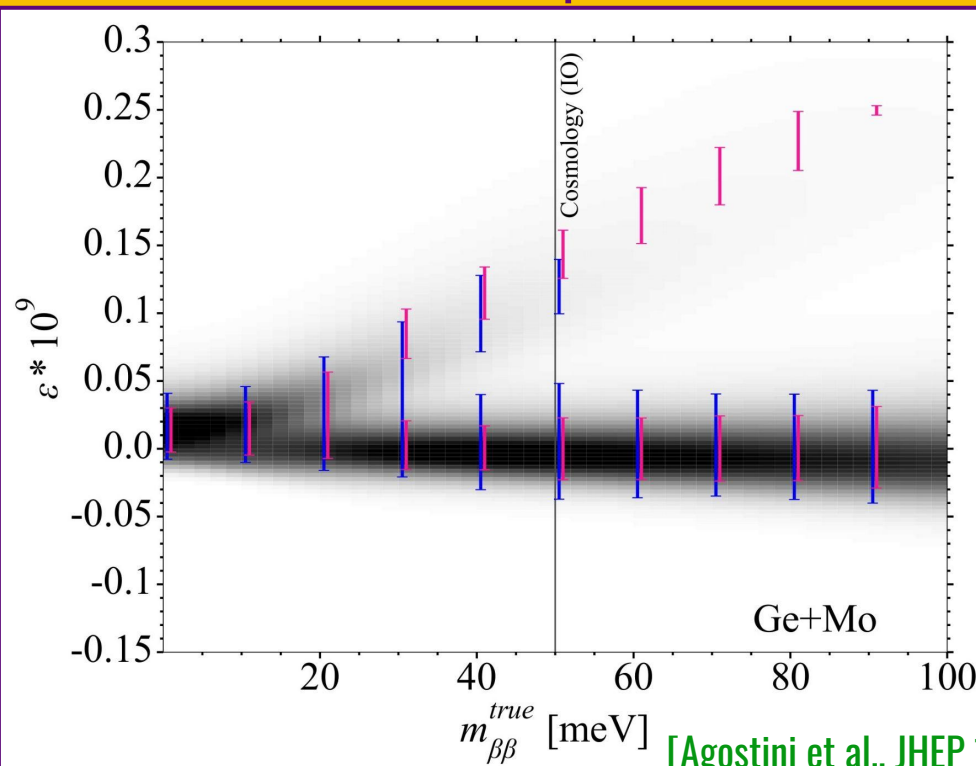
## Three isotopes



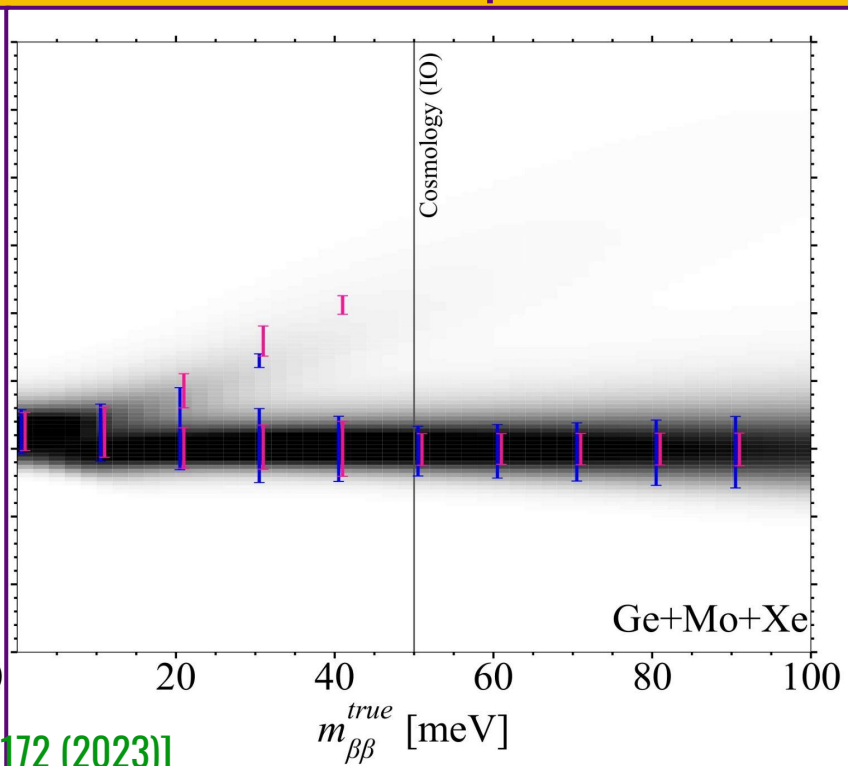
[Agostini et al., JHEP 172 (2023)]

# $\epsilon$ - Marginalized Statistics

## Two isotopes



## Three isotopes



[Agostini et al., JHEP 172 (2023)]

# Impact

- Multi-isotope observations put multi-mechanism constraints within reach!
- Two-isotope combinations can inform joint estimates on  $\{m_{\beta\beta}, \varepsilon\}$ , excluding null values.
  - A third isotope can suppress anomalous secondary peak, contingent on its slope.
  - Extra peaks also excluded by cosmology in many scenarios. [3]
- These results hold even when taking account of NME uncertainties and correlations.
  - NME correlations indeed improve inference, but more weakly than a third isotope.
  - Improving NME uncertainties of crucial value.

THANK YOU FOR LISTENING!

# Works Cited

[1] G. collaboration, M. Agostini, G. R. Araujo, A. M. Bakalyarov, M. Balata, I. Barabanov et al., Final results of gerda on the search for neutrinoless double- $\beta$  decay, 2020.

[2] KamLAND-Zen Collaboration, A. Gando, Y. Gando, T. Hachiya, A. Hayashi, S. Hayashida, H. Ikeda et al., Search for majorana neutrinos near the inverted mass hierarchy region with kamland-zen , Phys. Rev. Lett. 117 (Aug, 2016) 082503.

[3] M. Agostini, G. Benato and J. Detwiler, Discovery probability of next-generation neutrinoless double-beta decay experiments , Phys. Rev. D96 (2017) 053001, [1705.02996].

Abgrall, N., et al. 'LEGEND-1000 Preconceptual Design Report'. Phys. Rev. C 97 (7 2021): 045503. Web.

Adhikari, G., et al. 'nEXO: Neutrinoless double beta decay search beyond  $10^{28}$  year half-life sensitivity'. Phys. Rev. C 97 (6 2021): 045503. Web.

Armstrong, W. R., et al. 'CUPID pre-CDR'. Phys. Rev. C 97 (7 2019): 045503. Web.

Barea, J., J. Kotila, and F. Iachello. '0 $\nu\beta\beta$  and 2 $\nu\beta\beta$  nuclear matrix elements in the interacting boson model with isospin restoration'. Phys. Rev. C 91 (2015): 034304. Web.

Coraggio, L. et al. 'The calculation of the neutrinoless double-beta decay matrix element within the realistic shell model'. Phys. Rev. C 101.4 (2020): 044315. Web.

Fang, Dong-Liang, Amand Faessler, and Fedor Šimkovic. '0 $\nu\beta\beta$  nuclear matrix element for light and heavy neutrino mass mechanisms from deformed quasiparticle random-phase approximation calculations for 76Ge, 82Se, 130Te, 136Xe, and 150Nd with isospin restoration'. Phys. Rev. C 97 (2018): 045503. Web.

Horoi, Mihai, and Andrei Neacsu. 'Shell model predictions for 124Sn double- $\beta$  decay'. Phys. Rev. C 93 (2016): 024308. Web.

Hyvarinen, Juhani, and Jouni Suhonen. 'Nuclear matrix elements for 0 $\nu\beta\beta$  decays with light or heavy Majorana-neutrino exchange'. Phys. Rev. C 91 (2015): 024613. Web.

López Vaquero, Nuria, Tomás R. Rodríguez, and J. Luis Egido. 'Shape and pairing fluctuations effects on neutrinoless double beta decay nuclear matrix elements'. Phys. Rev. Lett. 111 (2013): 142501. Web.

Menéndez, J. 'Neutrinoless  $\beta\beta$  decay mediated by the exchange of light and heavy neutrinos: The role of nuclear structure correlations'. J. Phys. G 45 (2018): 014003. Web.

Mustonen, M. T., and J. Engel. 'Large-scale calculations of the double- $\beta$  decay of 76Ge,130Te, 136Xe, and 150Nd in the deformed self-consistent Skyrme quasiparticle random-phase approximation'. Phys. Rev. C 87 (2013): 064302. Web.

Rodríguez, Tomas R., and G. Martínez-Pinedo. 'Energy density functional study of nuclear matrix elements for neutrinoless  $\beta\beta$  decay'. Phys. Rev. Lett. 105 (2010): 252503. Web.

Šimkovic, Fedor, Adam Smetana, and Petr Vogel. '0 $\nu\beta\beta$  nuclear matrix elements, neutrino potentials and SU(4) symmetry'. Phys. Rev. C 98 (2018): 064325. Web.

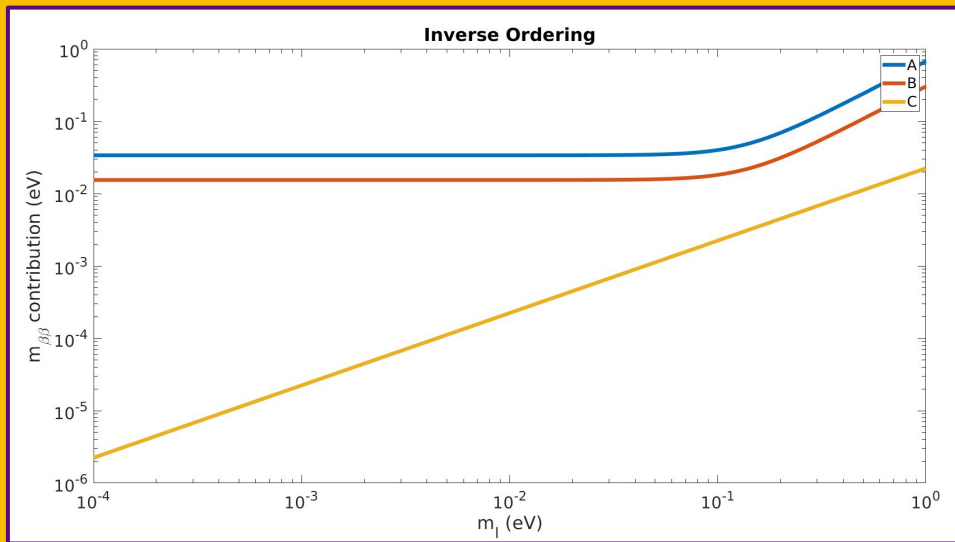
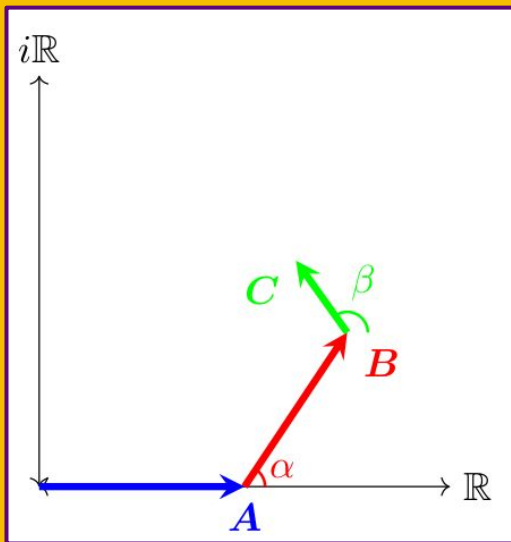
Song, L. S. et al. 'Nuclear matrix element of neutrinoless double- $\beta$  decay: Relativity and short-range correlations'. Phys. Rev. C 95 (2017): 024305. Web.

# Backup I: Neutrinos and $0\nu\beta\beta$

- $m_{\beta\beta}$  dependent on neutrino masses and two Majorana phases  $\alpha$  and  $\beta$ :

$$m_{\beta\beta} = \left| c_{12}^2 c_{13}^2 m_1 + s_{12}^2 c_{13}^2 m_2 e^{i\alpha} + s_{13}^2 m_3 e^{i\beta} \right| = \left| A + B e^{i\alpha} + C e^{i\beta} \right|$$

- With IO hierarchy, small  $m_3$  causes  $C \ll A, B$ .



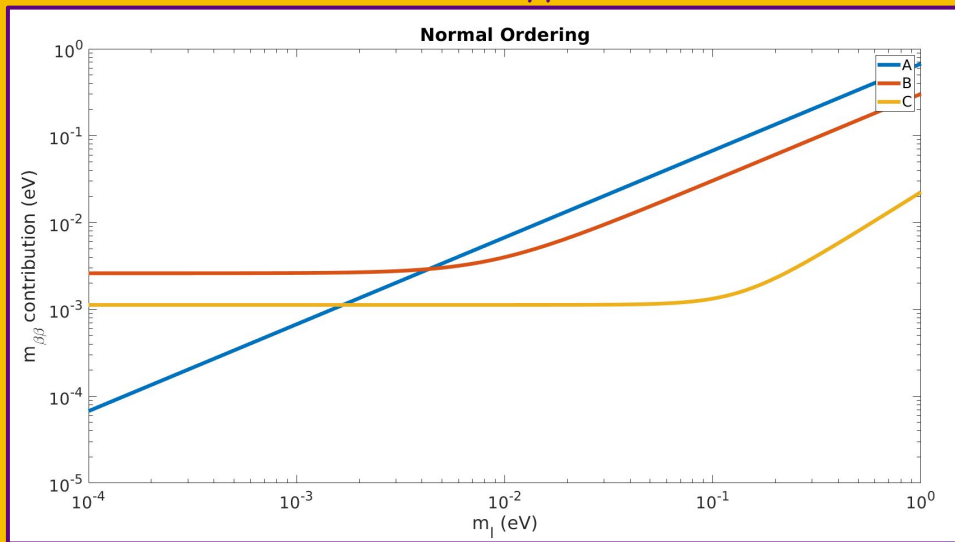
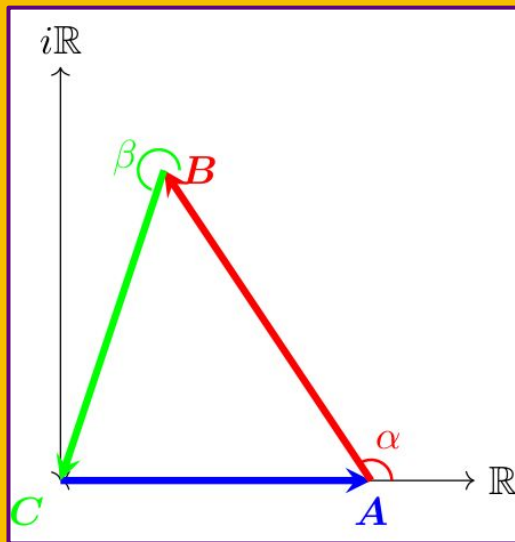


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- With NO hierarchy, possible to tune parameters and drive  $m_{\beta\beta} \rightarrow 0$ .



# Backup II: Least-Informative Priors

- **Information theory:** we want posterior to arise from experiment, not our prior biases.
- Information gain from prior to posterior captured in the Kullback-Leibler divergence.
- Define Least-Informative Prior (LIP) as prior which maximises expected information gain.

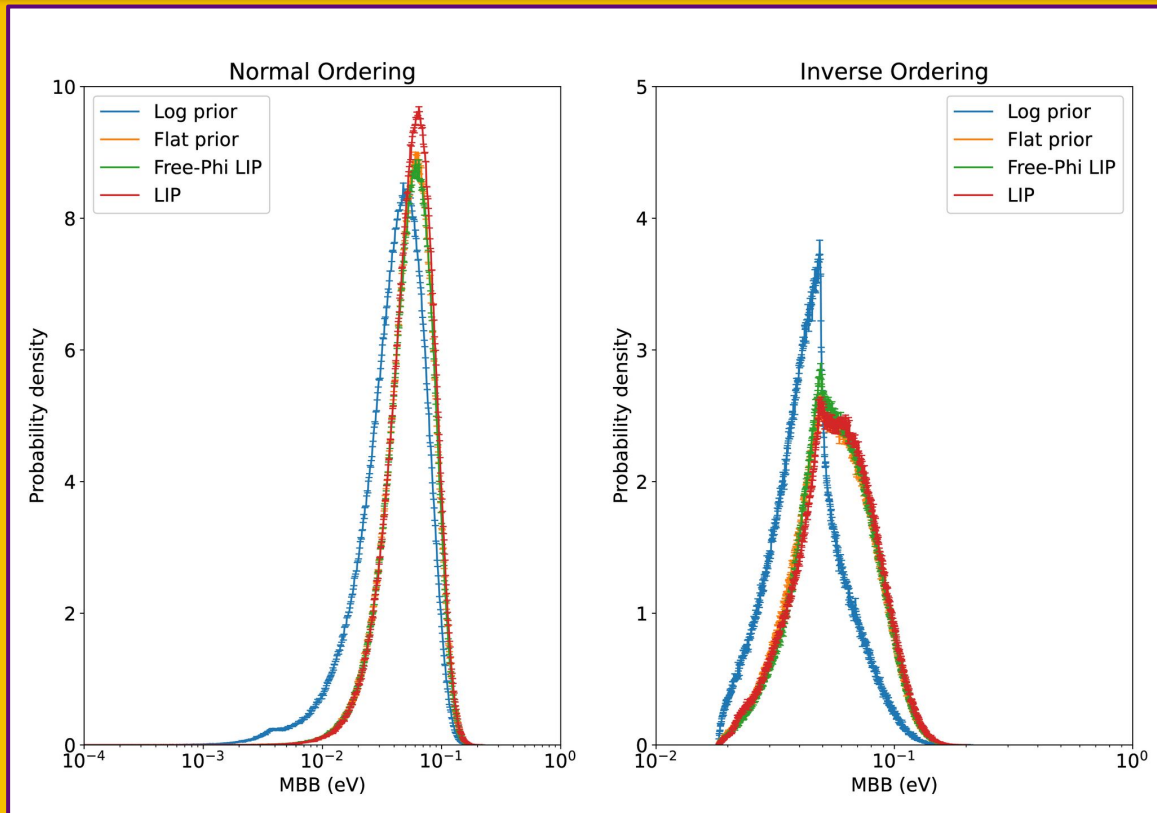
$$\langle D_{KL}(p|\pi) \rangle_{data} = \int dX \int d\theta p(\theta|X) \log \left( \frac{p(\theta|X)}{\pi(\theta)} \right)$$

- No prior is bias-free, but the LIP can serve as a reference with which to compare others.

# Aside: Least-Informative Priors

- Assess choice of priors on  $m_{\beta\beta}$  by comparing to LIPs.
- LEGEND-200 (n=1) modelled by Poisson likelihood.
- LIPs nearly reproduce posteriors for flat priors.

NO:  $m_{\beta\beta} \sim 58$  meV    IO:  $m_{\beta\beta} \sim 45$  meV

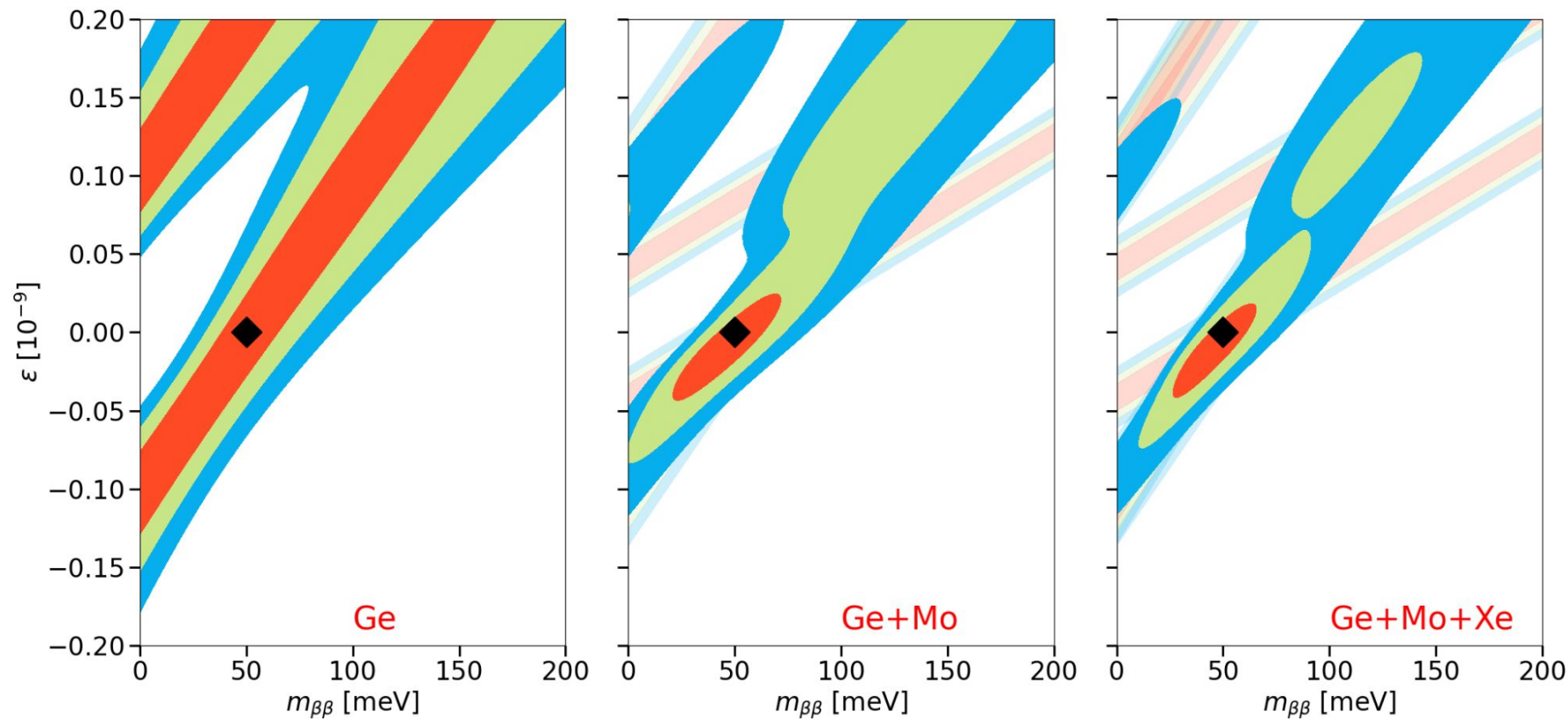


# Aside: Least-Informative Priors

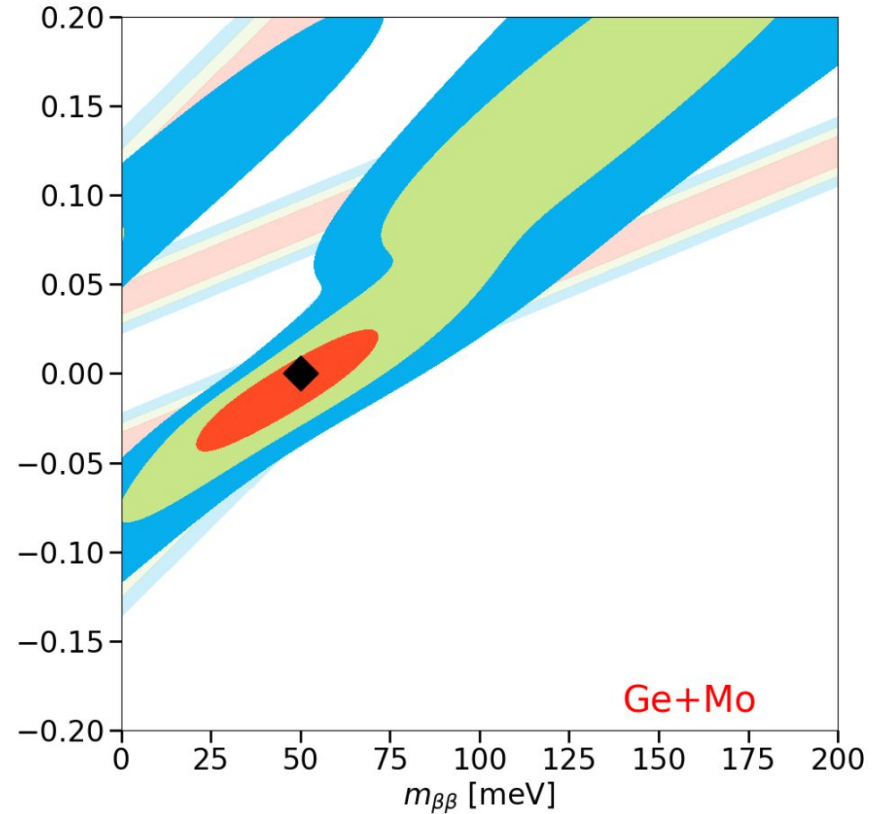
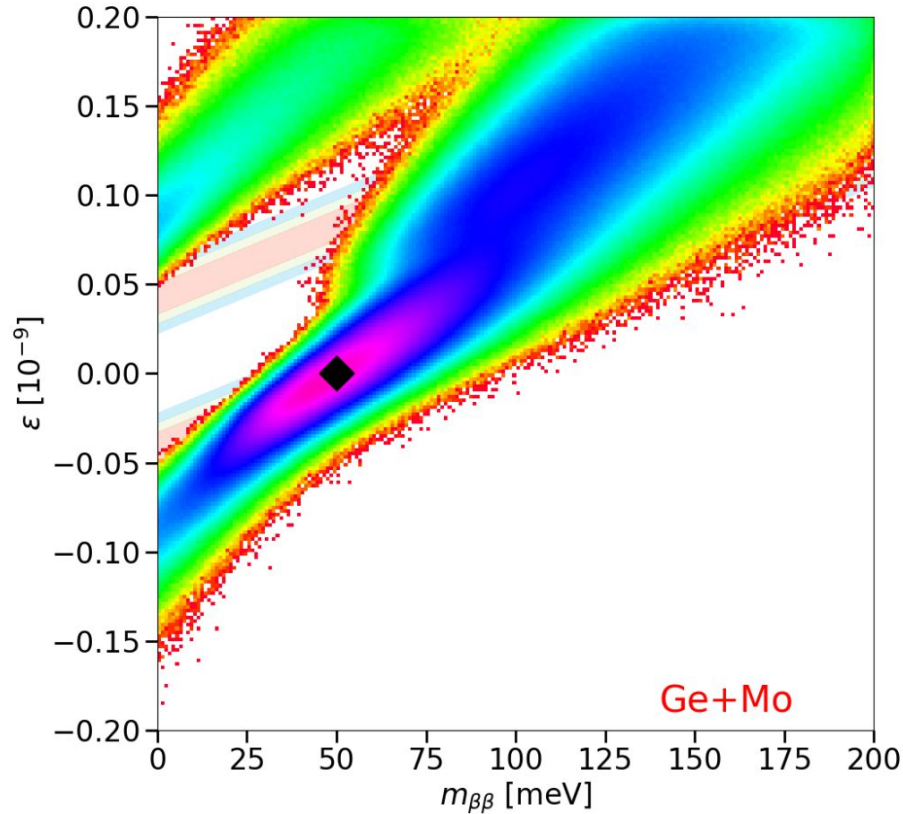
- Maximal information gain of LIP confirmed by computing Kullback-Leibler divergence.
- Further confirms sufficiency of flat prior for our purposes.

Prior	NO, $n = 0$	IO, $n = 0$	NO, $n = 1$	IO, $n = 1$
Flat $m_l$ , flat $\Phi$	20.63	20.61	20.63	20.73
Log-flat $m_l$ , flat $\Phi$	14.50	14.53	20.50	15.54
Free-phi LIP	21.90	21.94	21.48	21.75
Full LIP	<b>22.09</b>	<b>22.08</b>	<b>21.70</b>	<b>21.97</b>

# $\chi^2$ Analysis with Uncorrelated NME Uncertainties<sup>29</sup>



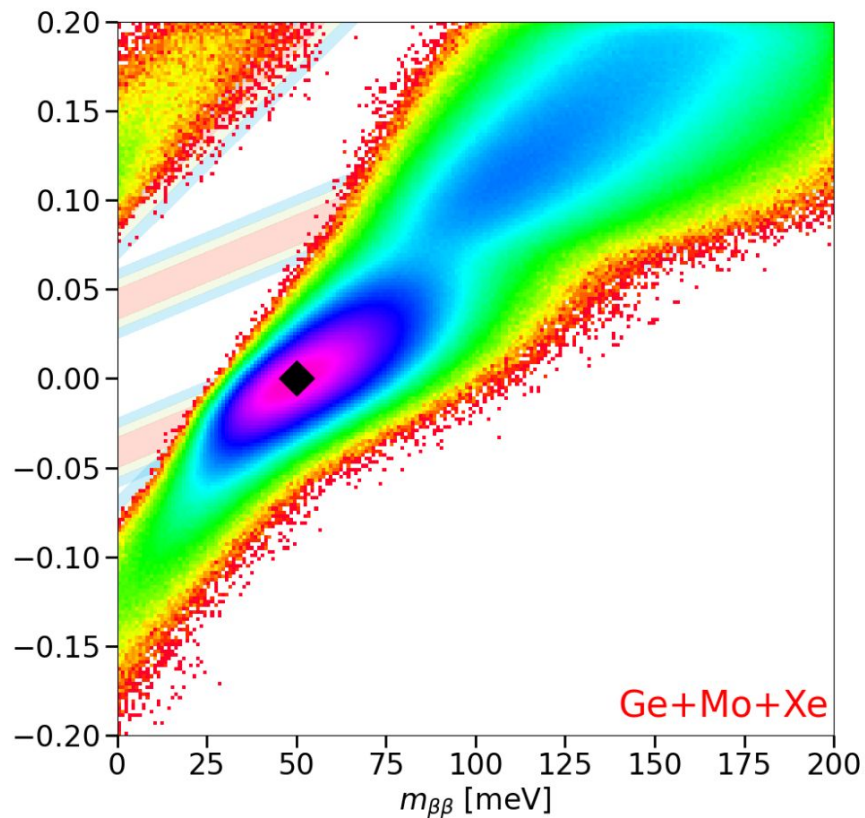
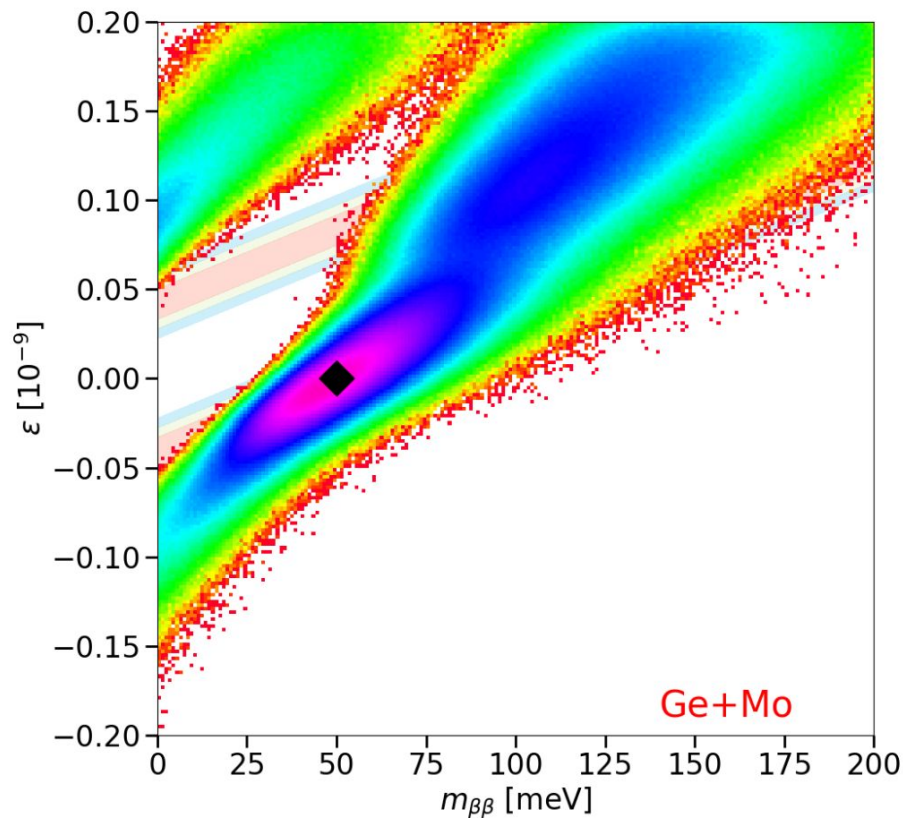
# MCMC with Uncorrelated NME Errors



# Ge+Mo

# vs.

# Ge+Mo+Xe



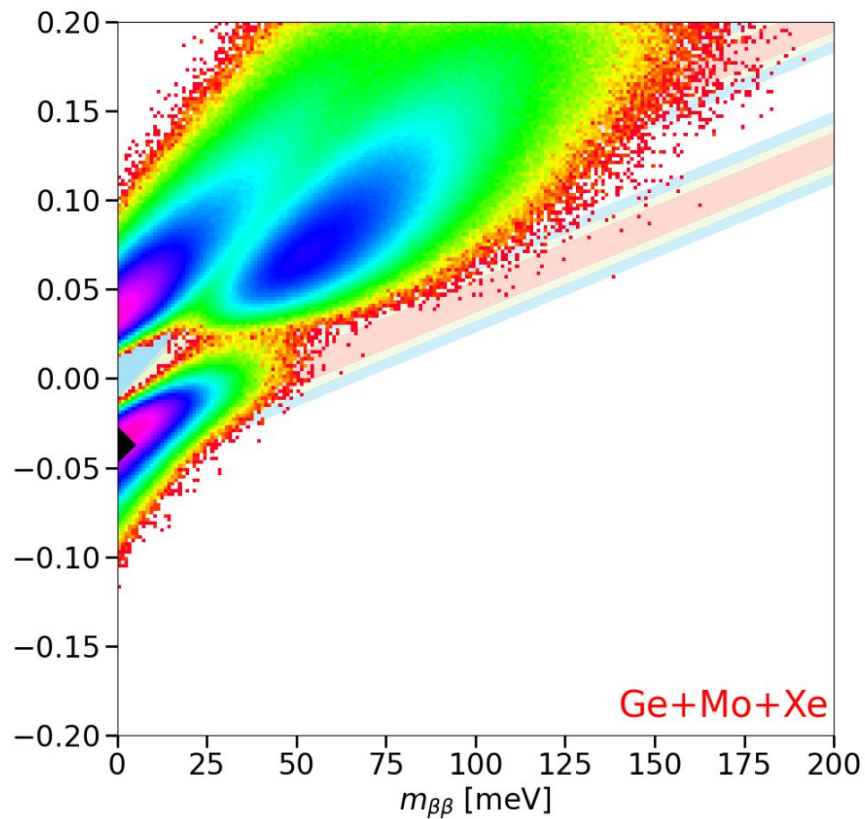
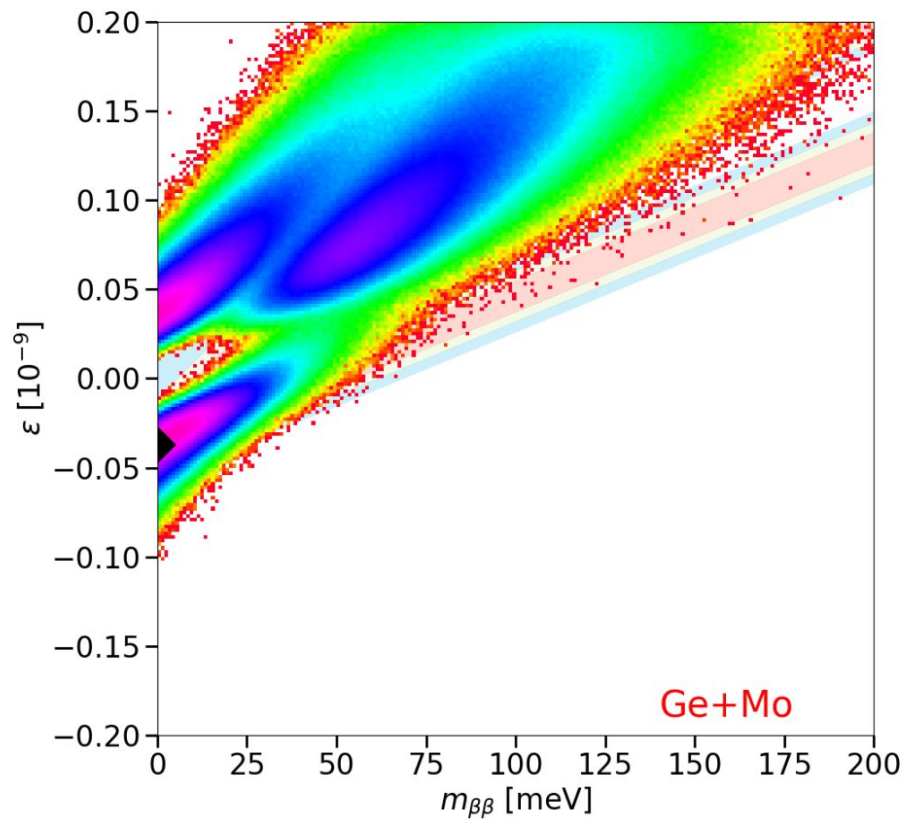
$\varepsilon \neq 0$ 

Ge+Mo

vs.

Ge+Mo+Xe

32





$\epsilon = -0.037$ 

vs.

 $\epsilon = -0.099$ 