Precision HH predictions: QCD and EW corrections

NExT Workshop, Sussex

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Higgs self coupling

Standard Model Higgs potential:

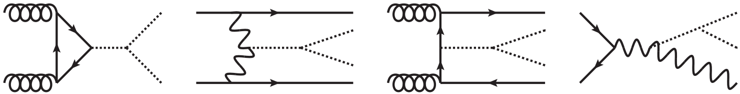
$$V(H) = \frac{1}{2} m_H^2 H^2 + \lambda v H^3 + \frac{\lambda}{4} H^4,$$

where $\lambda = m_H^2/(2v^2) \approx 0.13$.

Want to measure λ, to determine if $V(H)$ is consistent with nature.

- Challenging! Cross-section $\approx 10^{-3} \times H$ prod.
- $-3.3 < \lambda/\lambda_{SM} < 8.5$ [CMS '21]

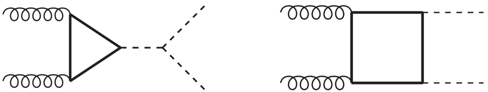
λ appears in various production channels:



- Gluon fusion – dominant, 10x
- VBF
- $t\bar{t}$ associated production
- H-strahlung

Gluon Fusion

Leading order (1 loop) partonic amplitude:



$$\mathcal{M}^{\mu\nu} \sim \mathcal{A}_1^{\mu\nu}(\mathcal{F}_{tri} + \mathcal{F}_{box1}) + \mathcal{A}_2^{\mu\nu}(\mathcal{F}_{box2})$$

- $\mathcal{F}_{tri}$ contains the dependence on λ at LO

Form factors:

- LO: known exactly [Glover, van der Bij '88]
- Beyond LO... no fully-exact (analytic) results to date
  - QCD: numerical evaluation, expansion in various kinematic limits
  - EW: first steps: HE expansion [Davies, Mishima, Schönwald, Steinhauser, Zhang '22]
    - (see also HTL considerations) [Mühlleitner,Schlenk,Spira '22]

$gg \to HH$ Beyond LO QCD

NLO QCD:

- large-m_t [Dawson,Dittmaier,Spira '98] [Grigo,Hoff,Melnikov,Steinhauser '13]
- numeric [Borowka,Greiner,Heinrich,Jones,Kerner,Schlenk,Schubert,Zirke '16] [Baglio,Campanario,Glaus,Mühlleitner,Spira,Streicher '19]
- large-m_t + threshold exp. Padé [Gröber, Maier, Rauh '17]
- high-energy expansion [Davies, Mishima, Steinhauser, Wellmann '18,'19]
- small-p_T expansion [Bonciani, Degrassi, Giardino, Gröber '18]
- small-t expansion [Davies, Mishima, Schönwald, Steinhauser '23]

NNLO QCD:

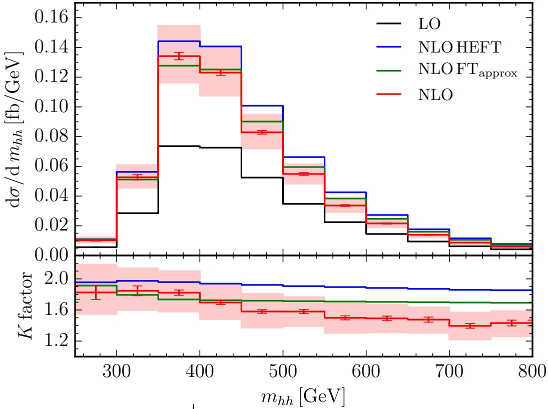
- large-m_t virtuals [de Florian, Mazzitelli '13] [Grigo, Hoff, Steinhauser '15][Davies, Steinhauser '19]
- HTL+numeric real ("FTapprox") [Grazzini, Heinrich, Jones, Kallweit, Kerner, Lindert, Mazzitelli '18]
- large-m_t reals [Davies, Herren, Mishima, Steinhauser '19 '21]
- small-t expansion ?? [work in progress: Davies, Schönwald, Steinhauser]

N3LO QCD:

- Wilson coefficient C_{HH} [Spira '16][Gerlach, Herren, Steinhauser '18]
- HTL [Chen, Li, Shao, Wang '19]

$gg \to HH$ Beyond LO QCD

[Borowka, Greiner, Heinrich, Jones, Kerner '16]

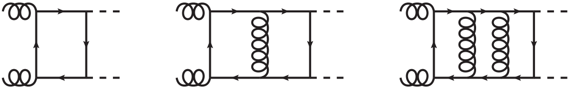


Total XS (14TeV):

	σ_{LO}	σ_{NLO}	σ_{NNLO}
B-i HTL	–	$38.32^{+18.1\%}_{-14.9\%}$	$39.58^{+1.4\%}_{-4.7\%}$
FTapprox	–	$34.25^{+14.7\%}_{-13.2\%}$	$36.69^{+2.1\%}_{-4.9\%}$
Full	$19.85^{+27.6\%}_{-20.5\%}$	$32.88^{+13.5\%}_{-12.5\%}$	–

QCD Corrections

Example diagrams at LO, NLO, NNLO:



Diagrams depend on ϵ, s, t, m_t, m_H:

- analytic result very complicated
- simplify: expand in certain kinematic limits

Here I will describe two expansions:

- high-energy: description for larger p_T values $s, |t| \gg m_t^2 \gg m_H^2$
- small-t: description for smaller p_T values $s, m_t^2 \gg |t|, m_H^2$

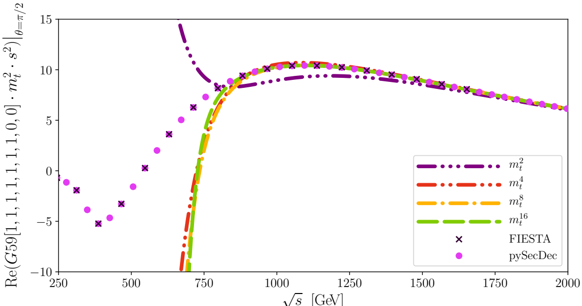
High-energy expansion

First, expand around $m_H \to 0$

- removes dependence on m_H from the integrals
- IBP-reduce to master integrals which depend on ϵ, s, t, m_t

Determine master integrals as an expansion around $m_t \to 0$:

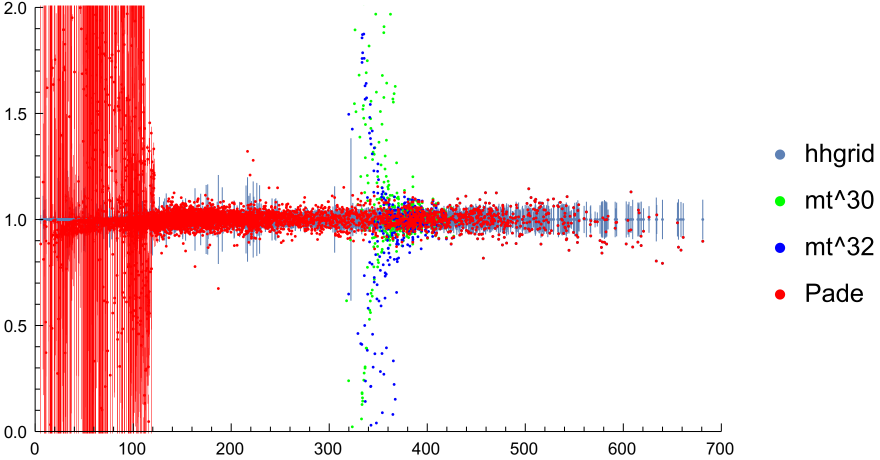
- use diff. eqns + BCs to obtain a deep expansion
- result: power series in m_t and $\log m_t$
- Padé approximants improve the region of validity of the series



High-energy expansion: "V_{fin}"

Comparison with `hhgrid`: [https://github.com/mppmu/hhgrid]

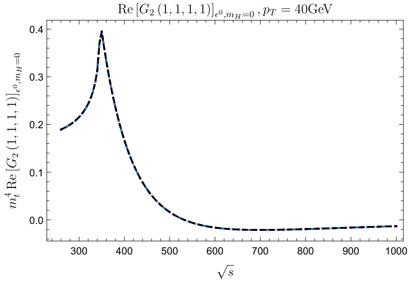
- interpolation grid of 6320 points evaluated by `pySecDec`
- grid points normalized to `hhgrid` vals. , function of p_t:

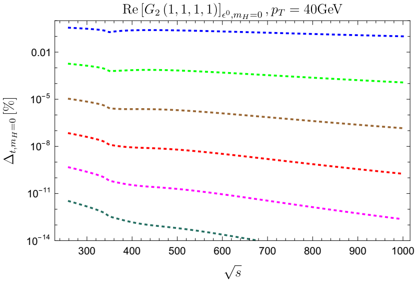


Small-t expansion

Again, first expand around $m_H \to 0$.
Then, two approaches (which give the same result at NLO):

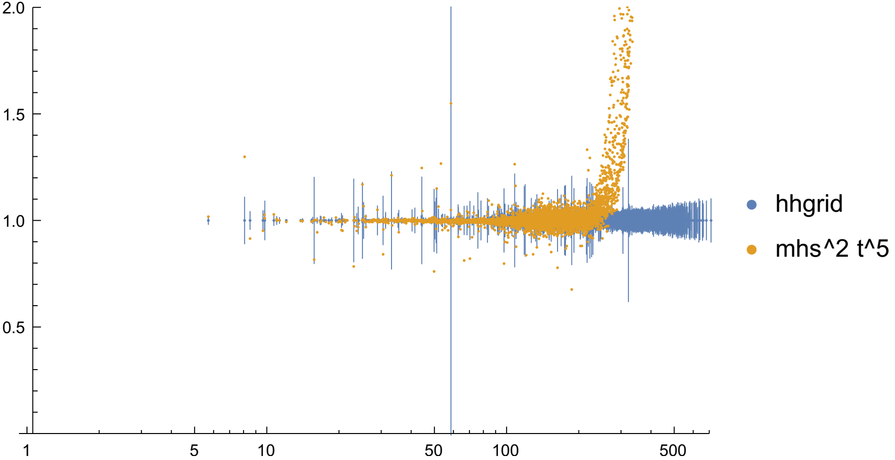
- take the IBP-reduced amplitude of the HE expansion
  - expand around $t \to 0$ instead of $m_t \to 0$
- expand un-reduced amplitude around $q_3 \to -q_1$ ($t \to 0$)
  - IBP reduce integrals which depend only on ϵ, s, m_t
  - compute resulting MIs as expansions around various s/m_t^2 values
  - can be applied at NNLO





Small-t expansion: "V_{fin}"

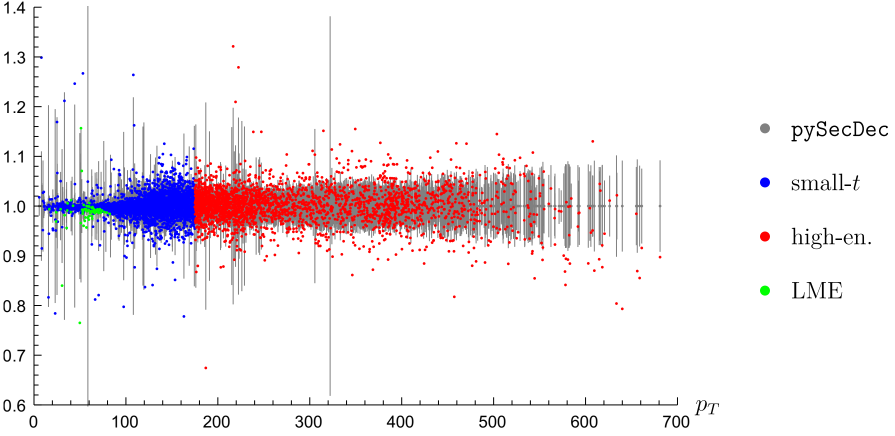
Comparison with `hhgrid`: [https://github.com/mppmu/hhgrid]



Combination: "V_{fin}"

Comparison with `hhgrid`: [https://github.com/mppmu/hhgrid]

- merge both results, switch at $p_T = 175$ GeV.
- both expansions implemented in `C++`: avg. 0.002s per point

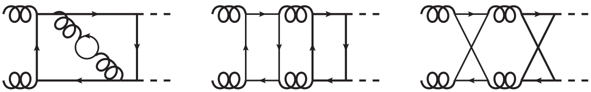


Small-t expansion: NNLO?

We would like to understand $gg \to HH$ at NNLO, due to the large NLO uncertainty from the top quark mass scheme and scale choice.

First steps: diagrams with light fermion loop, expand to $m_H^0 t^0$

[work in progress: Davies, Schönwald, Steinhauser]



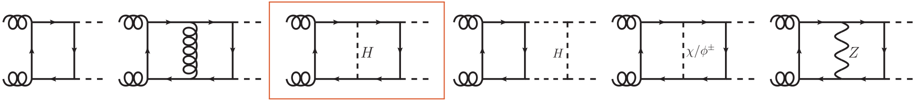
- done: IBP reduction: 177 MIs
- check: LME of MIs, agrees with literature [Davies, Steinhauser '19]
- todo: compute MIs as expansions around various s/m_t^2 values

EW Corrections

Since we look at NNLO QCD, we should also look at 2L EW corrections.

This is a much more difficult computation:

- 2L QCD: 118 Feynman diagrams
- 2L EW: 3810 Feynman diagrams



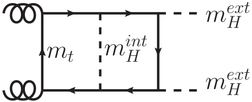
There are also more scales to deal with, compared to the QCD.

- start with $\alpha_s\alpha_t^2$ diagrams with internally propagating Higgs
  - expansion parameter $\alpha_t = \alpha m_t^2/(2s_W^2 m_W^2) \sim \alpha_s/2$
  - only planar integrals in this subset [Davies, Mishima, Schönwald, Steinhauser, Zhang '22]

EW Corrections: High-energy expansion

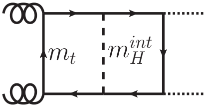
Again, full diagrams depend on many variables:

- ϵ, s, t, m_t, m_H



As before, expand around $m_H^{ext} = 0$:

- integral still depends on m_H^{int}



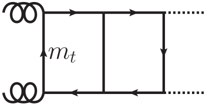
Expand in m_H^{int} also, two ways to do it:

- B: $s, |t| \gg m_t^2 \sim {m_H^{int}}^2 \gg {m_H^{ext}}^2$,
- A: $s, |t| \gg m_t^2 \gg {m_H^{int}}^2 \sim {m_H^{ext}}^2$.

High-energy Expansion "B"

Option B: expand around $m_H^{int} \approx m_t$,

- simple Taylor expansion, easy to implement



Write Higgs propagator as: $\frac{1}{p^2 - m_H^2} = \frac{1}{p^2 - m_t^2(1 - [2-\delta]\delta)}$

- expand around $\delta \to 0$ where $\delta = 1 - m_H/m_t \approx 0.28$.

This yields new integral families compared to the QCD computation:

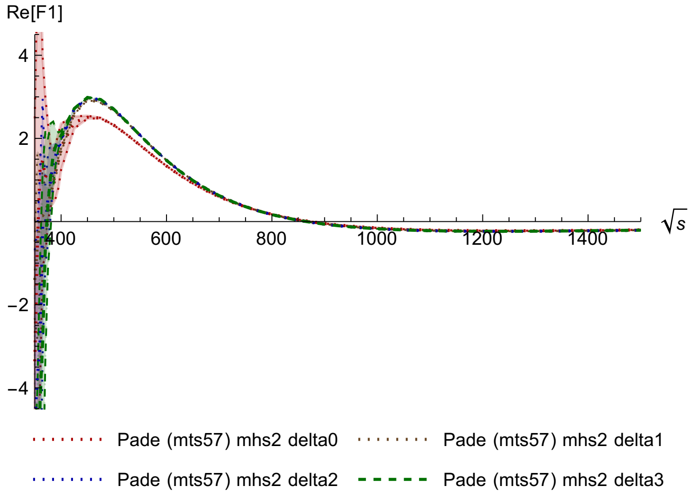
- all lines have the mass m_t,
- IBP reduce and compute the MIs in the high-energy limit

Expand to $(m_H^{ext})^4$ and δ^3.

Convergence of delta expansion (“B”)

Re(F_{box1}), fixed $\cos\theta = 0$, expansion “B” Padé (to $(m_H^2)^2\delta^{\{0,1,2,3\}}$):

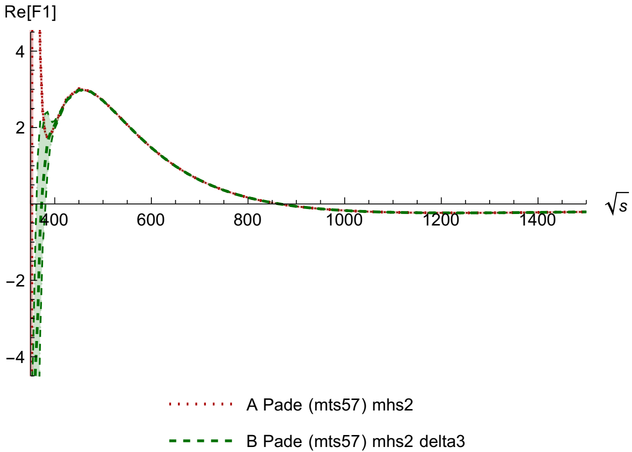
- δ^2 and δ^3 terms differ by at most 0.5% for $\sqrt{s} \geq 400$GeV



Comparison of “A”, “B” expansions

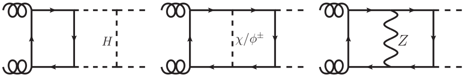
Re(F_{box1}), fixed $\cos\theta = 0$, best “A” and “B” Padé

- “A”, “B” differ by at most 2% for $\sqrt{s} \geq 400$GeV,
- 0.1% for $\sqrt{s} \geq 500$GeV



EW Corrections: next steps

High-energy: same idea, for remaining diagram classes:



Diagrams with more internal Higgs:

- delta expansion and IBP reduction done
- todo: compute master integrals

Diagrams with charged goldstone or W exchange:

- new integral topologies where top-quark loop doesn't close
- not yet studied

Small-t: not yet studied

Conclusion

Multi-scale multi-loop integrals are hard.

- expand!

Expansions give a good description for HH at NLO QCD:

- high-energy + small-t covers whole phase space
- implemented in `C++`: to be made public

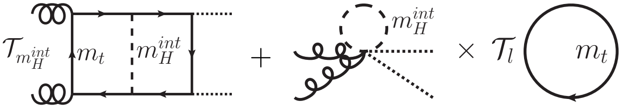
First steps:

- NNLO QCD: small-t expansion of light-fermion diagrams
  - to come: remaining diagrams
  - deeper expansion?
- EW: high-energy expansion of y_t^4 diagrams
  - to come: remaining diagrams
  - small-t expansion

High-Energy Expansion "A"

Option A: asymptotic expansion around $m_H^{int} = 0$:

- two sub-graphs:



The two-loop subgraph is a Taylor expansion of the Higgs propagator:

- results in integrals with a massless internal line, scales s, t, m_t.
- IBP reduce with `FIRE` and `Kira` [Smirnov '15] [Klappert,Lange,Maierhöfer,Usovitsch '21]
- these coincide with the QCD Master Integrals – reuse the old results [Davies,Mishima,Steinhauser,Wellmann '18,'19]

The massive tadpoles are easily computed by `MATAD`. [Steinhauser '00]

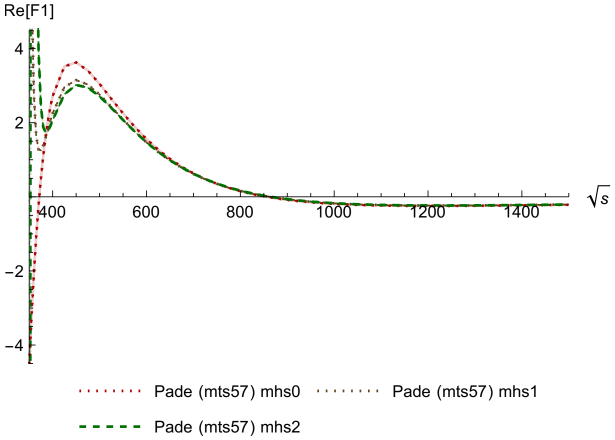
The asymp. expansion procedure is done by `exp` and `FORM` [Harlander,Seidelsticker,Steinhauser '97] [Ruijl,Ueda,Vermaseren '17]

We expand to quartic order: $(m_H^{int})^a \, (m_H^{ext})^b, \;\; 0 \leq (a + b) \leq 4$.

Convergence of asymptotic expansion ("A")

Re(F_{box1}), fixed $\cos\theta = 0$, expansion "A" Padé (to $(m_H^2)^{\{0,1,2\}}$):

- $(m_H^2)^1$ and $(m_H^2)^2$ terms differ by at most 5% for $\sqrt{s} \geq 400$GeV



Padé-Improved High-Energy Expansion

The MIs for both EW methods are computed as an exp. in $m_t \ll s, |t|$.

The expansions diverge for $\sqrt{s} \sim 750$GeV ("A"), $\sqrt{s} \sim 1000$GeV ("B").

The situation can be improved using Padé Approximants:

- approximate a function using a rational polynomial:

$$f(x) \approx [n/m](x) = \frac{a_0 + a_1 x + a_2 x^2 + \cdots + a_n x^n}{1 + b_1 x + b_2 x^2 + \cdots + b_m x^m},$$

  where a_i, b_j coefficients are fixed by the series coefficients of $f(x)$.

We compute a set of various Padé Approximants:

- combine to give a central value and error estimates
- a deeper input expansion $\rightarrow$ larger $n + m$ $\rightarrow$ smaller error
- here, m_t^{120} exp. allows for very high-order Padé Approximants

High-Energy Expansion and Padé Approximation

Re(F_{box1}), fixed $\cos\theta = 0$, expansion "B" (to $(m_H^2)^2\delta^3(m_t^2)^{\{15,16,56,57\}}$):

- m_t expansion diverges (strongly) around $\sqrt{s} \sim 1000$GeV