

# Footprints of LQs: from $B$ to $K$ rare decays

**Luiz Vale Silva**

University of Sussex

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[arXiv:1711.xxxxx [hep-ph]], in collaboration with  
**Svjetlana Fajfer and Nejc Košnik (Institut Jožef Stefan)**

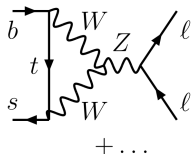
NExT Physics Meeting, RHUL

# Outline

- 1 Introduction
- 2 NP for  $B$ -anomalies
- 3 Pheno of two LQ models
- 4  $s \rightarrow d\nu\nu$  transitions

# B-physics anomalies

SM LFU (Lepton Flavor Universality) respected to a very good extent  
 Pattern of deviations w.r.t. the SM  $\rightarrow$  NP LFUV



## Neutral Currents

$$R_{K^{(*)}} = \frac{\mathcal{B}(B \rightarrow K^{(*)} \mu^+ \mu^-)}{\mathcal{B}(B \rightarrow K^{(*)} e^+ e^-)}$$

[BaBar, Belle, LHCb; LHCb'17]

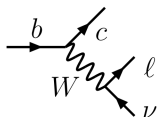
$$\sim 2.1 - 2.6\sigma \text{ (w/ SM)}$$

BRs and angular obs.  
 in  $b \rightarrow s \mu^+ \mu^-$

[LHCb, Belle, ATLAS, CMS]

$$\sim 2.2 - 2.9\sigma \text{ (w/ SM)}$$

## Charged Currents



$$R_{D^{(*)}} = \frac{\mathcal{B}(B^0 \rightarrow D^{(*)-} \tau^+ \nu)}{\mathcal{B}(B^0 \rightarrow D^{(*)-} \ell^+ \nu)}$$

[BaBar, Belle; BaBar, Belle, LHCb]

$$\sim 2 - 3.4\sigma \text{ (w/ SM)}$$

$$R_{J/\psi} = \frac{\mathcal{B}(B_c^+ \rightarrow (J/\psi) \tau^+ \nu)}{\mathcal{B}(B_c^+ \rightarrow (J/\psi) \mu^+ \nu)}$$

[LHCb'17]

$$\sim 2\sigma \text{ (w/ SM)}$$

## Possible EFT interpretations

$$\text{N.C.} : \mathcal{L}_{b \rightarrow s \ell \ell}^{NP} \supset 4 \frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \left[ \delta^\ell \bar{s} \gamma_\rho P_L b \cdot \bar{\ell} \gamma^\rho (1 - \gamma_5) \ell \right] + \text{h.c.}$$

[Altmannshofer+'17, Capdevila+'17, L.-S. Geng+'17]

$$\text{C.C.} : \mathcal{L}_{b \rightarrow c \ell \nu}^{NP} \supset -\frac{2G_F}{\sqrt{2}} V_{cb} \left[ \epsilon^\ell \bar{c} \gamma_\rho P_L b \cdot \bar{\ell} \gamma^\rho (1 - \gamma_5) \nu_\ell \right] + \text{h.c.}$$

[X.-Q. Li+'16, Alonso+'16, Celis+'17]

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$\delta^\mu(\mu_b)$  and  $\epsilon^\tau(\mu_b)$  at the level of  $\sim \mathcal{O}(10\% - 20\%)$  of the SM,  
with very different meanings: SM loop vs. SM tree

New degrees of freedom at scales  $\Lambda_{NP} \sim \mathcal{O}(1 - 100)$  TeV

# Correlation with other flavor sectors

Move to specific models to relate  $B$ - and  $K$ -decays [Crivellin+'16, Bordone+'17]

## Rare $s \rightarrow d\nu\nu$ transitions

- NA62/CERN:  $K^\pm \rightarrow \pi^\pm \nu \bar{\nu}$
- KOTO/J-PARC:  $K_L \rightarrow \pi^0 \nu \bar{\nu}$  (CP Violation)



Collecting data, announcements expected before 202X ( $X = 0, 1, 2$ )

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**HERE:** discuss what can be learned from these transitions  
in some specific NP contexts: leptoquarks (LQs)

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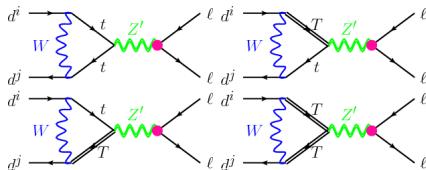


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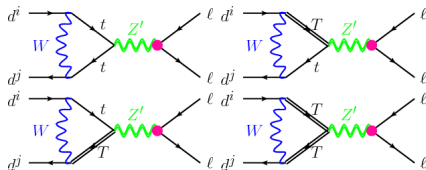
[Altmannshofer+'14, Kamenik+'17,...]



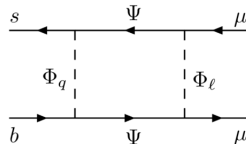
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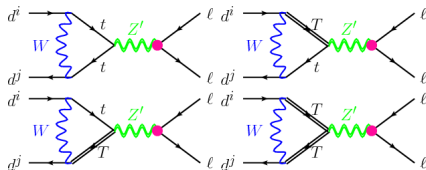
[Gripaios+'15, Arnan+'15]



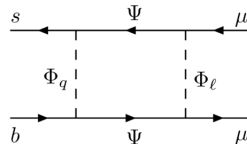
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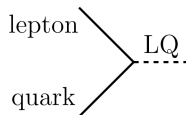
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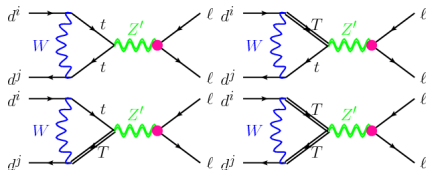
[Hiller+'14, Gripaios+'14, Becirevic+'15'16, M. Varzielas+'15, Fajfer+'15,...]



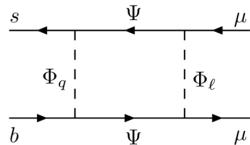
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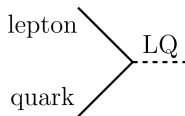
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More generally, **difficult** to accommodate both classes of LFUV anomalies,  $R_{K^{(*)}}$  and  $R_{D^{(*)}}$ , simultaneously

# Leptoquark models

- First LQs: matter unification [Pati, Salam '70's; Georgi, Glashow '74,...]

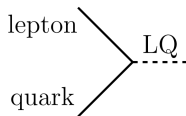
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- Spin-0 and spin-1 particle coupling to SM +  $\nu_R$  d.o.f.,  
**gauge invariant**, dim. = 4 operators:

$$\underbrace{S_3, R_2, \tilde{R}_2, \tilde{S}_1, S_1, \bar{S}_1}_{6 \text{ scalar LQs}}, \underbrace{U_3, V_2, \tilde{V}_2, \tilde{U}_1, U_1, \bar{U}_1}_{6 \text{ vector LQs}}$$

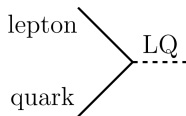




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- Vector LQs: **renormalizable** model requires larger spectrum

# Structure of LQ contributions to neutral currents

Measurements:  $R_K/R_K^{SM}$  and  $R_{K^*}/R_{K^*}^{SM} < 1$

Couplings of scalar LQs to SM fermions

LQs	down-type interac.	chiral structure	@ tree-level $R_K/R_K^{SM}, R_{K^*}/R_{K^*}^{SM}$
$S_3 = (\bar{\mathbf{3}}, \mathbf{3}, 1/3)$	$\bar{d}_L^c \nu_L, \bar{d}_L^c \ell_L$	$\bar{s} \gamma_\rho P_L b \cdot \bar{\ell} \gamma^\rho P_L \ell$	$< 1, < 1$
$R_2 = (\mathbf{3}, \mathbf{2}, 7/6)$	$\bar{d}_L \ell_R$	$\bar{s} \gamma_\rho P_L b \cdot \bar{\ell} \gamma^\rho P_R \ell$	$\approx 1, \approx 1$
$\tilde{R}_2 = (\mathbf{3}, \mathbf{2}, 1/6)$	$\bar{d}_R \ell_L, \bar{d}_R \nu_L$	$\bar{s} \gamma_\rho P^R b \cdot \bar{\ell} \gamma_\rho P_L \ell$	$< 1, > 1$
$\tilde{S}_1 = (\bar{\mathbf{3}}, \mathbf{1}, 4/3)$	$\bar{d}_R^c \ell_R$	$\bar{s} \gamma_\rho P_R b \cdot \bar{\ell} \gamma^\rho P_R \ell$	$\approx 1, \approx 1$
$S_1 = (\bar{\mathbf{3}}, \mathbf{1}, 1/3)$	$\bar{d}_L^c \nu_L$		$= 1, = 1$

(w/ SM +  $\nu_R$ , also  $\bar{S}_1 = (\bar{\mathbf{3}}, \mathbf{1}, -2/3)$  and new couplings)

→ Tree-level:  $S_3$

→ Loop-level: other LQs can also imply  $R_K/R_K^{SM}, R_{K^*}/R_{K^*}^{SM} < 1$

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# Specific LQ models

→ In the following, we consider and compare:

- a  $R_2$  model (“Doublet model”)
- a  $S_3$  model (“Triplet model”)

[Becirevic+'17]

[Doršner+'17]

→ Detailed phenomenological studies available

→ **Substantial differences:** e.g.,  $b \rightarrow sll$  at different orders

## $R_2$ model: features

Interactions with SM fermions:

$$\mathcal{L}_{R_2}^Y = (Vg_R)_{ij} \bar{u}^i P_R e^j R_2^{5/3} + (g_R)_{ij} \bar{d}^i P_R e^j R_2^{2/3} \\ + (g_L)_{ij} \bar{u}^i P_L \nu^j R_2^{2/3} - (g_L)_{ij} \bar{u}^i P_L e^j R_2^{5/3} + \text{h.c.}$$

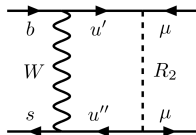
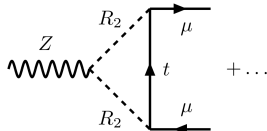
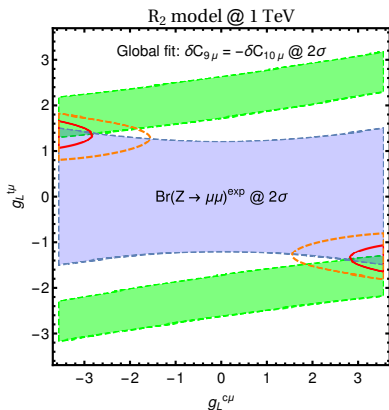
$$g_R = 0_{3 \times 3}, \quad g_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & g_L^{c\mu} & g_L^{c\tau} \\ 0 & g_L^{t\mu} & g_L^{t\tau} \end{pmatrix}, \quad m_{R_2}$$

[Becirevic+'17]

- Avoid **tree-level contributions** to  $B$ -decays w/ the *wrong* chirality
- $R_{D^{(*)}}$  not addressed (e.g.,  $g_R^{b\tau} \neq 0$  strongly constrained)
- Consistently avoid first generation couplings
- **No tree-level** contribution to  $s \rightarrow d\nu\nu$

# $R_2$ model: phenomenology

- $g_L^{c\tau} \approx 0$ ,  $g_L^{t\tau} \approx 0$  for large  $g_L^{c\mu}$ ,  $g_L^{t\mu}$  due to  $\tau \rightarrow \mu\gamma$



$$g_L^{c\mu}, g_L^{t\mu} < \sqrt{4\pi}$$

- $(g - 2)_\mu$  worsen by  $\sim 1\sigma$
- Collider bounds:  $m_{R_2} \gtrsim 650$  GeV (assuming  $t\nu, t\tau$  dominate)

## $S_3$ model: features

Interactions with SM fermions:

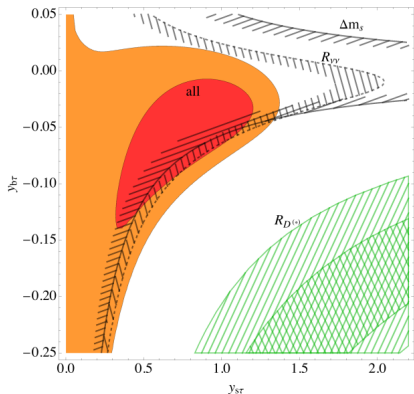
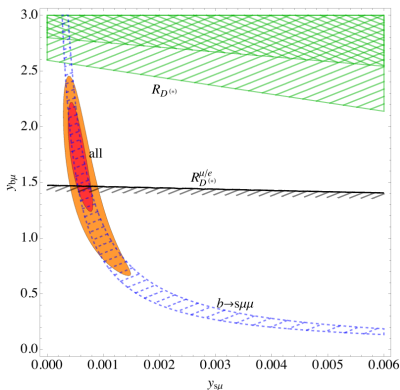
$$\mathcal{L}_{S_3}^Y \equiv -y_{ij} \bar{d}_L^{Ci} \nu_L^j S_3^{1/3} - (V_{CKMY}^*)_{ij} \bar{u}_L^{Ci} e_L^j S_3^{1/3} \\ - \sqrt{2} y_{ij} \bar{d}_L^{Ci} e_L^j S_3^{4/3} + \sqrt{2} (V_{CKMY}^*)_{ij} \bar{u}_L^{Ci} \nu_L^j S_3^{-2/3} + \text{h.c.}$$

$$y = \begin{pmatrix} 0 & 0 & 0 \\ 0 & y_{s\mu} & y_{s\tau} \\ 0 & y_{b\mu} & y_{b\tau} \end{pmatrix}, \quad V_{CKMY}^* \approx \begin{pmatrix} 0 & \lambda_{CKM} y_{s\mu} & \lambda_{CKM} y_{s\tau} \\ 0 & y_{s\mu} & y_{s\tau} \\ 0 & y_{b\mu} & y_{b\tau} \end{pmatrix}, \quad m_{S_3}$$

[Doršner+'17]

- Single coupling matrix for up- and down-type processes
- With the **choice**  $y_{d\mu} = \mathbf{0} \Rightarrow$  no  $s \rightarrow d\nu\nu$  @ tree-level

# $S_3$ model: phenomenology



[Doršner+'17]

- $R_{D^{(*)}}^{\mu/e} = \frac{\mathcal{B}(B \rightarrow D^{(*)} \mu \nu)}{\mathcal{B}(B \rightarrow D^{(*)} e \nu)}$ ,  $R_{\nu\nu}^{(*)} = \frac{\mathcal{B}(B \rightarrow K^{(*)} \nu \nu)}{\mathcal{B}(B \rightarrow K^{(*)} \nu \nu)_{\text{SM}}}$ ,  $\Delta m_s(B_s^0 \bar{B}_s^0)$
- Collider bounds on LQ pair and  $\tau\tau$  production satisfied



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# Overview

## Experimental values

$$\mathcal{B}_{\text{exp}}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) < 3.35 \times 10^{-10} \quad @ \quad 90 \% \text{ CL} \quad \text{[BNL-E787, E949]}$$

$$\mathcal{B}_{\text{exp}}(K_L \rightarrow \pi^0 \nu \bar{\nu}) < 2.6 \times 10^{-8} \quad @ \quad 90 \% \text{ CL} \quad \text{[KEK-E391a]}$$

NA62 and KOTO: anticipated accuracies of 10%

## Theoretical predictions in the SM

$$\mathcal{B}_{SM}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = 0.882_{-0.098}^{+0.092} \times 10^{-10} \quad (\sim 10\%) \quad \text{[CKMfitter, preliminary]}$$

$$\mathcal{B}_{SM}(K_L \rightarrow \pi^0 \nu \bar{\nu}) = 0.314_{-0.018}^{+0.017} \times 10^{-10} \quad (\sim 5\%) \quad \text{[CKMfitter, preliminary]}$$

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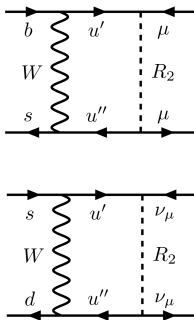
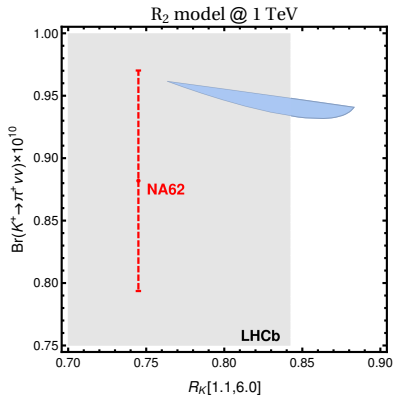
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( $sd$ )( $\ell\ell$ ) transitions in the SM: large  $\mathcal{B}(K^\pm \rightarrow \pi^\pm \ell^+ \ell^-)$  points to large long-distance effects from  $K \rightarrow \pi \gamma^*$

# Results for the $R_2$ model

Max. enhancement of 9% for  $K^\pm \rightarrow \pi^\pm \nu\nu$  and 5% for  $K_L \rightarrow \pi^0 \nu\nu$



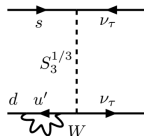
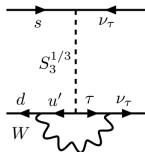
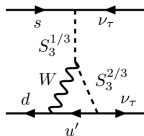
Effects induced by muon couplings,  $g_L^{c\mu}$ ,  $g_L^{t\mu}$

# Preliminary results for the $S_3$ model

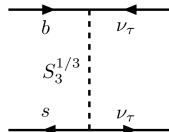
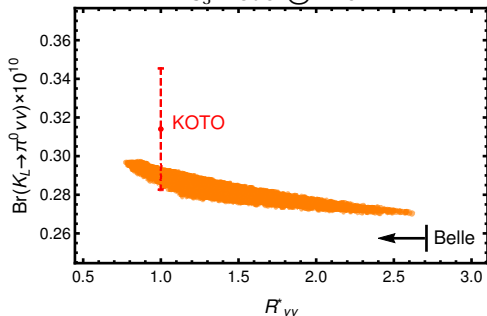
$\sim$  max. suppression of 10% for  $K^\pm \rightarrow \pi^\pm\nu\nu$  and 14% for  $K_L \rightarrow \pi^0\nu\nu$

tau couplings:

$$y_{d\tau} = [(1.9 + 0.7i)y_{b\tau} - (0.08 + 0.03i)y_{s\tau}] \times 10^{-4}$$



$S_3$  model @ 1 TeV



Preliminary: sub-region of  $\{g_{s\mu}, g_{b\mu}, g_{s\tau}, g_{b\tau}\} 1\sigma$

# What do we learn?

To correlate  $B$ - and  $K$ -physics, two LQ models w/ real couplings:

$$R_2 : \{g_L^{c\mu}, g_L^{t\mu}, g_L^{c\tau}, g_L^{t\tau}\}, S_3 : \{y_{C\mu}, y_{t\mu}, y_{C\tau}, y_{t\tau}\} \quad (m_{R_2} = m_{S_3} = 1 \text{ TeV})$$

- Given a model: must rely on the **complementarity** of the two channels
- The comparison of  $R_2$  and  $S_3$  illustrates possible ways to **discriminate** models addressing the  $B$ -anomalies:  
e.g., suppression/enhancements of  $s \rightarrow d\nu\nu$
- **Important modulations** ( $\gtrsim$  theo. unc.) also for  $S_1 + S_3$ :  
suppression of  $\sim 24\%$  for  $K^\pm \rightarrow \pi^\pm\nu\nu$ ,  $\sim 34\%$  for  $K_L \rightarrow \pi^0\nu\nu$

[Crivellin+'17]

# Conclusion

## Important exp. progress in the coming years in $K$ -physics

- The specific correlation among  $B$ - and  $K$ -physics depends on a specific NP model
- LQ models: tentative class of extensions of the SM for interpreting the  $b \rightarrow s\ell\ell$  anomalies
- The complementary study of  $s \rightarrow d\nu\nu$  decays may favor a specific LQ models

Thanks!



## Couplings of scalar LQs to SM fermions

LQs	Structures of the couplings	
	down-type	up-type
$S_3 = (\bar{\mathbf{3}}, \mathbf{3}, 1/3)$	$\bar{d}_L^c \nu_L, \bar{d}_L^c l_L,$	$\bar{u}_L^c \nu_L, \bar{u}_L^c l_L$
$R_2 = (\mathbf{3}, \mathbf{2}, 7/6)$	$\bar{d}_L l_R,$	$\bar{u}_L l_R, \bar{u}_R l_L, \bar{u}_R \nu_L$
$\tilde{R}_2 = (\mathbf{3}, \mathbf{2}, 1/6)$	$\bar{d}_R l_L, \bar{d}_R \nu_L$	
$\tilde{S}_1 = (\bar{\mathbf{3}}, \mathbf{1}, 4/3)$	$\bar{d}_R^c l_R$	
$S_1 = (\bar{\mathbf{3}}, \mathbf{1}, 1/3)$	$\bar{d}_L^c \nu_L,$	$\bar{u}_L^c l_L, \bar{u}_R^c l_R$

## Couplings of vector LQs to SM fermions

LQs	Structures of the couplings	
	down-type	up-type
$U_3 = (\mathbf{3}, \mathbf{3}, 2/3)$	$\bar{d}_L \nu_L, \bar{d}_L \ell_L,$	$\bar{u}_L \nu_L, \bar{u}_L \ell_L$
$V_2 = (\bar{\mathbf{3}}, \mathbf{2}, 5/6)$	$\bar{d}_R^c \nu_L, \bar{d}_R^c \ell_L, \bar{d}_L^c \ell_R,$	$\bar{u}_L^c \ell_R$
$\tilde{V}_2 = (\bar{\mathbf{3}}, \mathbf{2}, -1/6)$		$\bar{u}_R^c \ell_L, \bar{u}_R^c \nu_L$
$\tilde{U}_1 = (\mathbf{3}, \mathbf{1}, 5/3)$		$\bar{u}_R \ell_R$
$U_1 = (\mathbf{3}, \mathbf{1}, 2/3)$	$\bar{d}_L \ell_L, \bar{d}_R \ell_R,$	$\bar{u}_L \nu_L$

# Results for the $S_3$ model + $y_{d\mu}$

- Relax the initial requirement of  $y_{d\mu} = 0$
- A real  $y_{d\mu}^{\text{tree}}$  saturates the experimental bound of  $K^\pm \rightarrow \pi^\pm \nu \nu$ :  
 $|y_{d\mu}|$  below  $\sim 3 \times 10^{-4}$
- Much stronger than  $\tau \rightarrow \mu + K_S^0$

[Davidson+'10]

$$y = \begin{pmatrix} 0 & y_{d\mu} & 0 \\ 0 & y_{s\mu} & y_{s\tau} \\ 0 & y_{b\mu} & y_{b\tau} \end{pmatrix}$$

