



北京航空航天大学  
BEIHANG UNIVERSITY

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# Lattice QCD calculation of the unpolarized transverse-momentum-dependent parton distributions

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Jul. 25 @ LaMET2023



Based on *hep-lat/2211.02340*

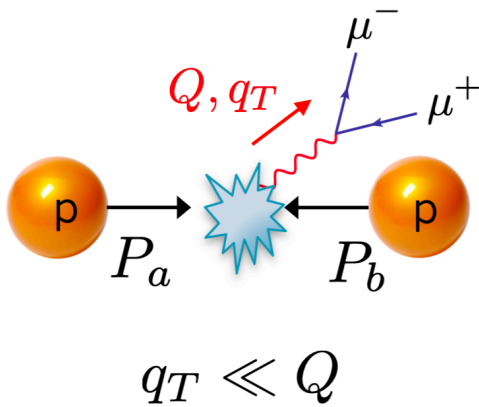
# OUTLOOK

- **Motivation and recent progresses**
- **Lattice QCD calculation of unpolarized TMDPDF**
  - **Extract TMDPDF from LaMET**
  - **Quasi TMDPDF matrix elements and their renormalization**
  - **From Quasi TMDPDF to physical TMDPDF**
  - **Numerical results and discussion**
- **Summary and outlook**

# What's and why TMDs?

## TMD processes:

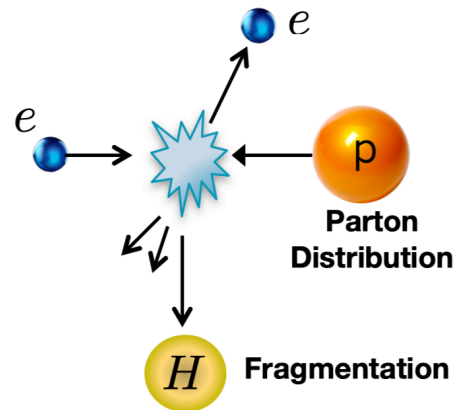
### Drell-Yan



LHC, FermiLab,  
RHIC, ...

$$\sigma \sim f_{q/P}(x, k_T) f_{q/P}(x, k_T)$$

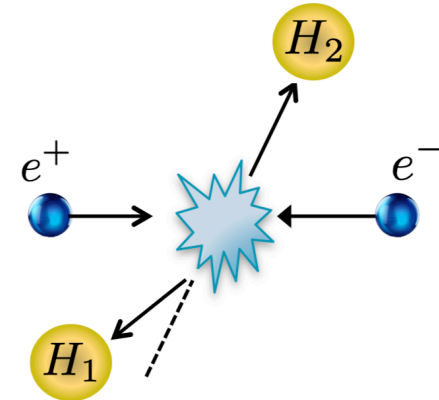
### Semi-Inclusive DIS



HERMES, COMPASS,  
JLab, EIC, ...

$$\sigma \sim f_{q/P}(x, k_T) D_{h/q}(x, k_T)$$

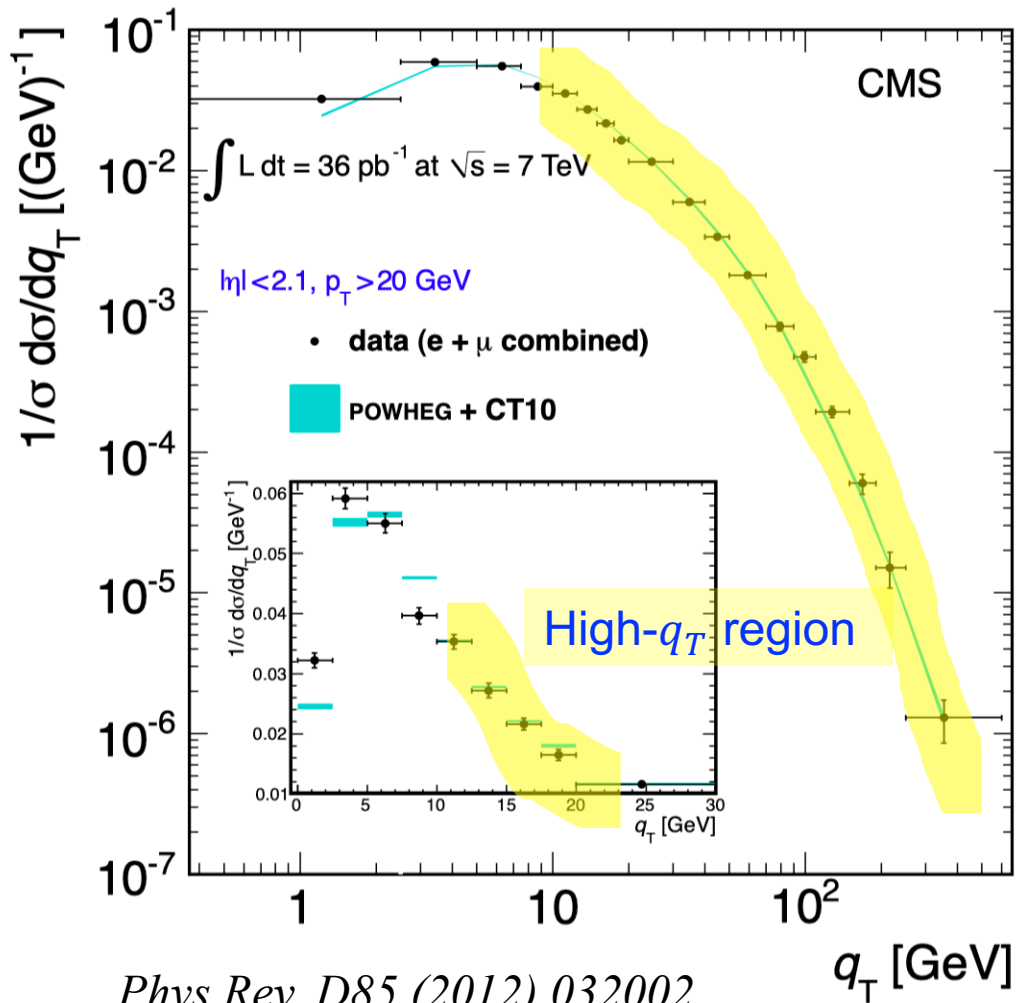
### Dihadron in e+e-



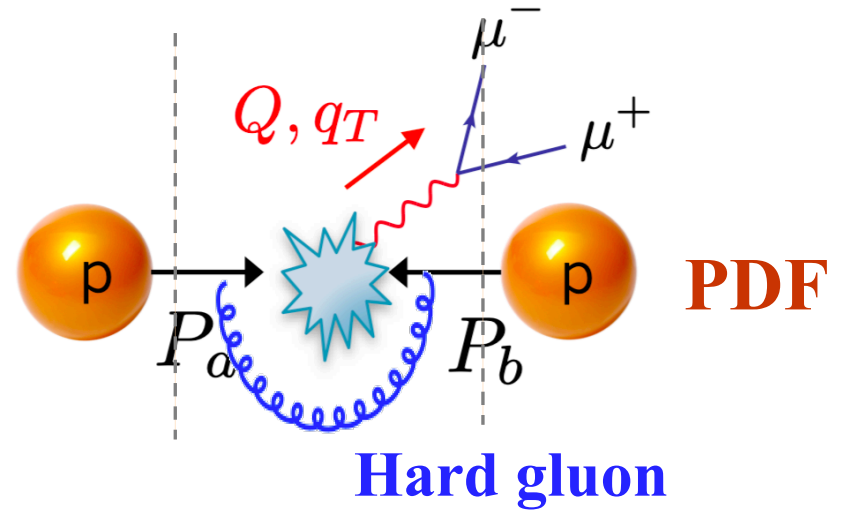
BESIII, Babar,  
Belle, ...

$$\sigma \sim D_{h_1/q}(x, k_T) D_{h_2/q}(x, k_T)$$

# Z-production $q_T$ spectrum at LHC



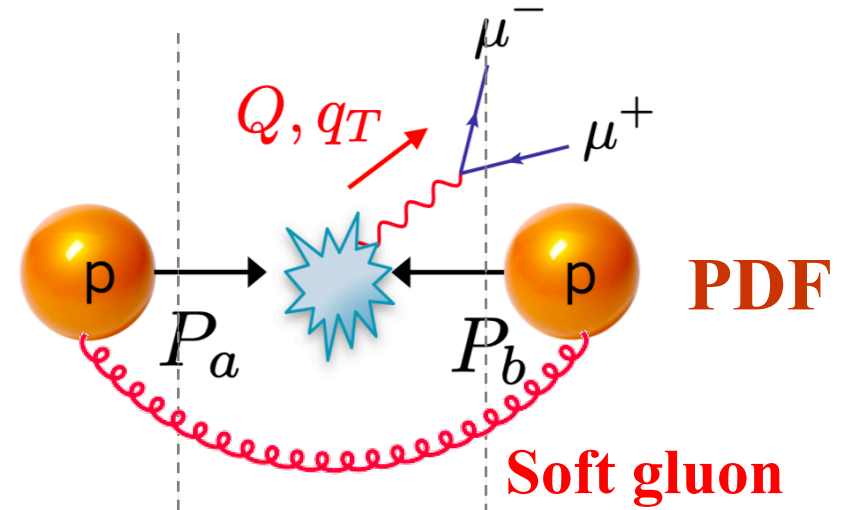
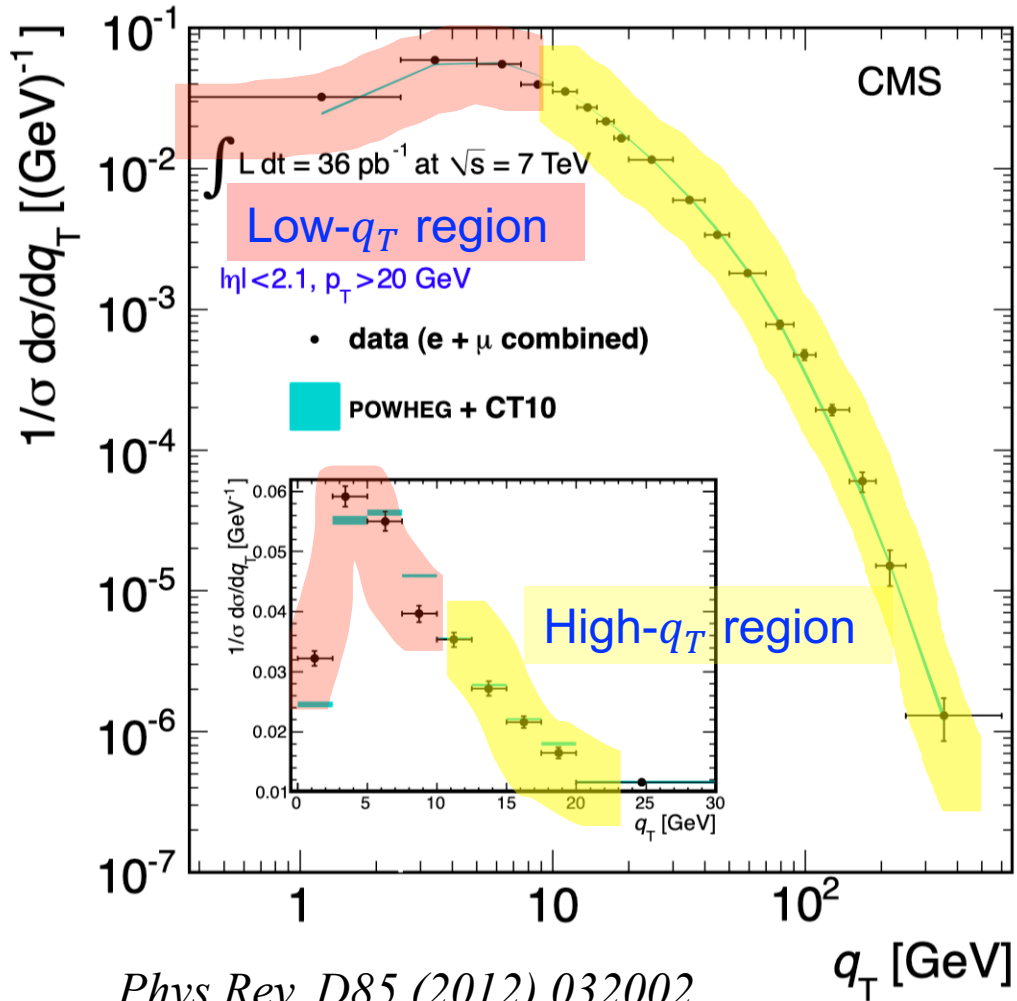
*Phys.Rev. D85 (2012) 032002.*



➤ **High- $q_T$  region:**

Collinear factorization => PDF

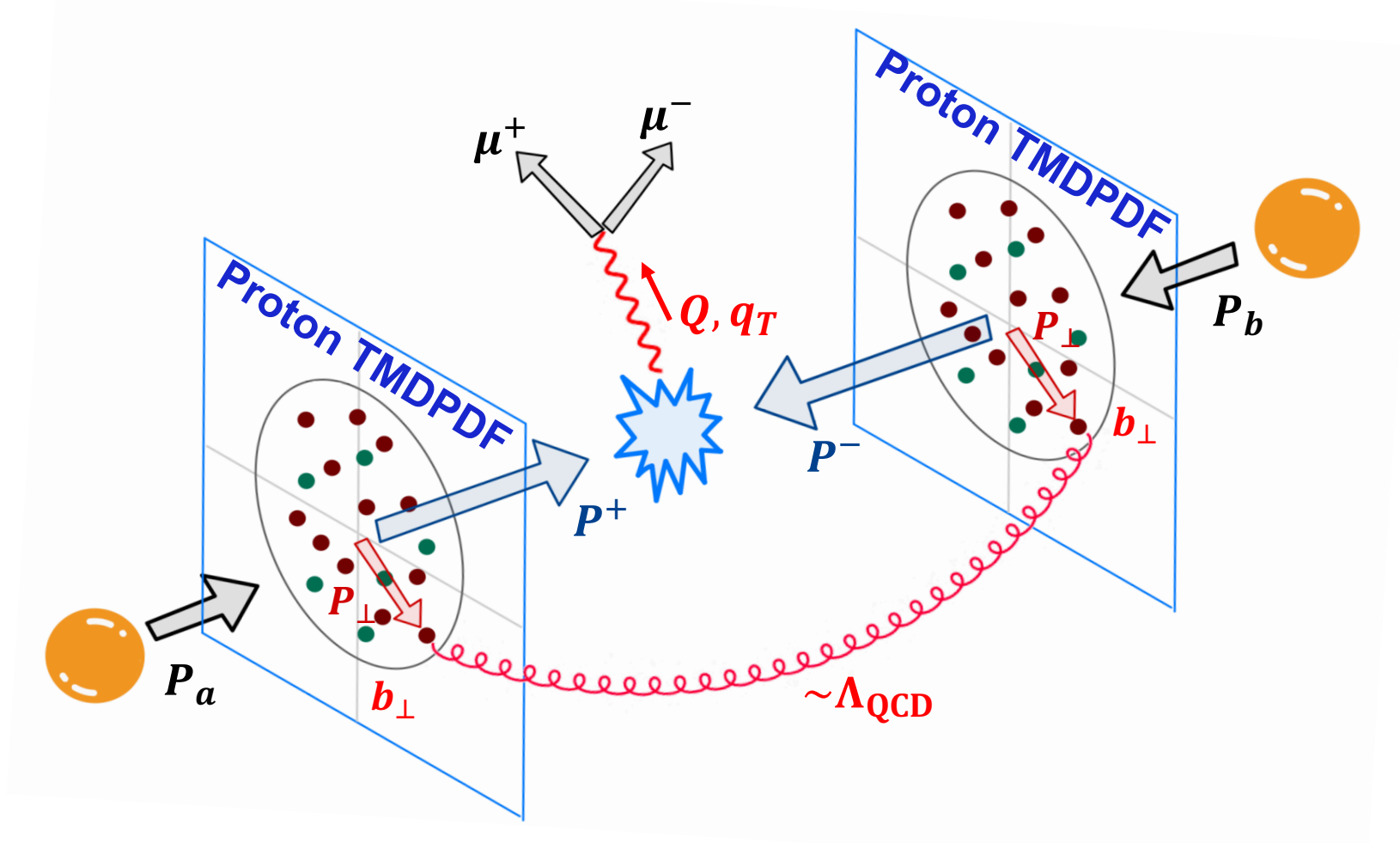
# Z-production $q_T$ spectrum at LHC



- **High- $q_T$  region:**  
Collinear factorization  $\Rightarrow$  PDF
- **Low- $q_T$  region:**  
TMD factorization  
 $\Rightarrow$  Generalize to TMDPDFs

# TMDPDFs: 3D tomography of the nucleon

- Low- $q_T$  region of Drell-Yan Process:

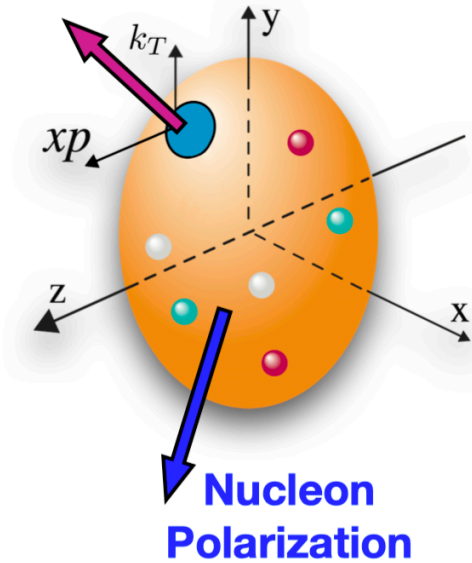


Revealing the  
confined motion of  
partons inside the nucleon



# TMDPDFs: 3D tomography of the nucleon

Quark Polarization



Leading Quark TMDPDFs



|                      |   | Quark Polarization   |   |   |
|----------------------|---|--|---|---|
|                      |   | Un-Polarized (U)   | Longitudinally Polarized (L)  | Transversely Polarized (T)  |
| Nucleon Polarization | U | $f_1 = \text{○} \cdot$<br>Unpolarized                              |   | $h_1^\perp = \text{○} \uparrow - \text{○} \downarrow$<br>Boer-Mulders   |
|                      | L |  | $g_{1L} = \text{○} \rightarrow - \text{○} \rightarrow$<br>Helicity                            | $h_{1L}^\perp = \text{○} \nearrow - \text{○} \searrow$<br>Worm-gear   |
|                      | T | $f_{1T}^\perp = \text{○} \uparrow - \text{○} \downarrow$<br>Sivers | $g_{1T}^\perp = \text{○} \uparrow \rightarrow - \text{○} \downarrow \rightarrow$<br>Worm-gear | $h_1 = \text{○} \uparrow - \text{○} \downarrow$<br>Transversity<br>$h_{1T}^\perp = \text{○} \nearrow - \text{○} \searrow$<br>Pretzelosity |

TMD Handbook, TMD Collaboration, 2304.03302

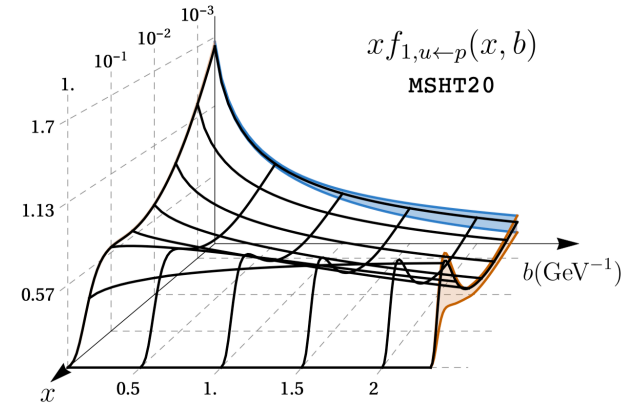
# Progress in the study of TMDs

## ➤ Theoretical analysis

- **TMD factorization, evolution and resummation:**

*Boussarie et al., TMD handbook, 2304.03302;*

*Collins, Foundations of perturbative QCD; .....*



## ➤ Phenomenological parametrizations and extractions

*u-quark unpolarized TMDPDF, 2201.07114*

- **Unpolarized:**

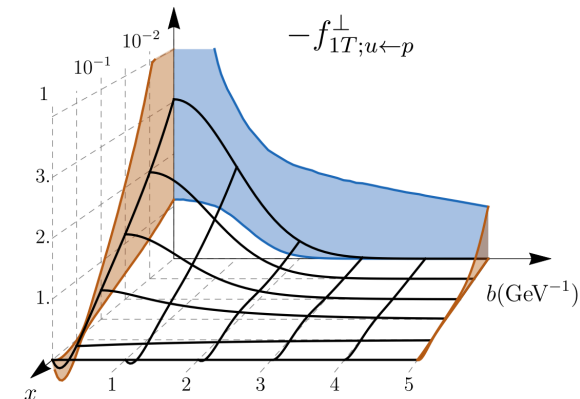
*Moos, 2305.07473; Bacchetta, 2206.07598; Bury, 2201.07114;*

*Scimem, JHEP06 (2020); Bacchetta, JHEP06 (2017); .....*

- **Sivers, Boer-Mulders:**

*Bury, PRL126 (2021), JHEP05 (2021) ; Cammarota, PRD102(2020);*

*Zhang, PRD77 (2008), Lu, PRD81 (2010) ; .....*



*u-quark Sivers function, PRL126 (2021)*



## ➤ Lattice calculations

- **Lorentz-invariant approach:** ratios of Mellin moments

*Hagler, EPL88(2009); Musch, PRD85(2012); Engelhardt, PRD93(2016); Yoon, 1601.05717, PRD96(2017); .....*

- **LaMET formalism:**

- ✓ **I: theoretical analysis of matching kernel, soft function, Collins-Soper kernel, .....**

*Rio, 2304.14440; Ji, 2305.04416, RMP93(2021), NPB955(2020), PLB811(2020);*

*Ebert, JHEP04(2022); Deng, JHEP09(2022).....*

- ✓ **II: lattice calculation of intrinsic soft function, Collins-Soper kernel, beam function, .....**

*LPC, 2306.06488, PRL125(2020); Li, PRL128(2022); LPC, PRD106(2022); Shanahan,*

*PRD104(2021); Schlemmer, JHEP08(2021); .....*

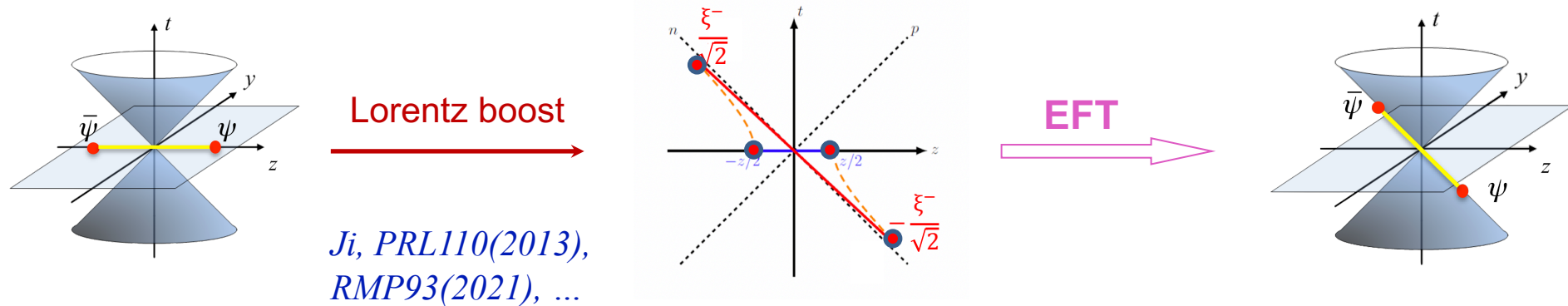
- ✓ **III: Nonperturbative renormalization, resummation, .....**

*Zhang, 2305.05212; Ji, 2305.04416; Su, NPB991(2023); LPC, PRL129(2022); 2209.01236.....*

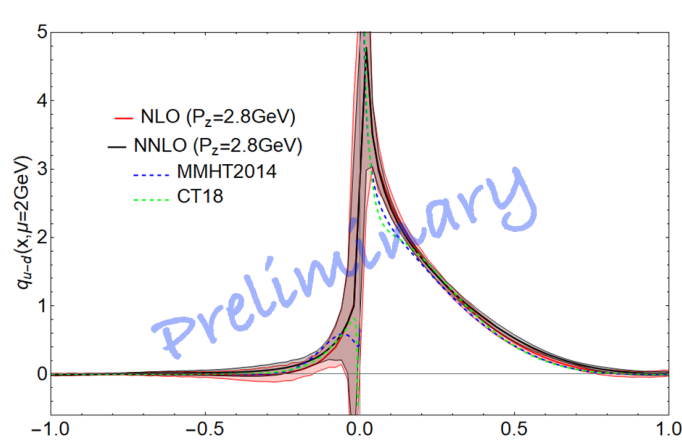
- **IV: A real lattice calculation of TMD observable?**

# Extracting TMDs in LaMET formalism

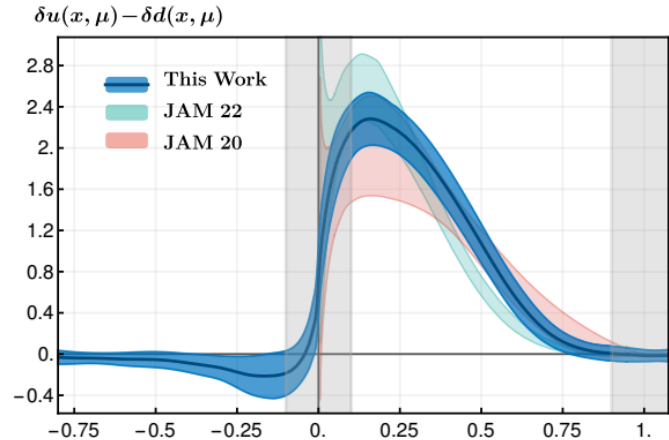
- Large-momentum effective theory: connecting Euclidean lattice and physical observables



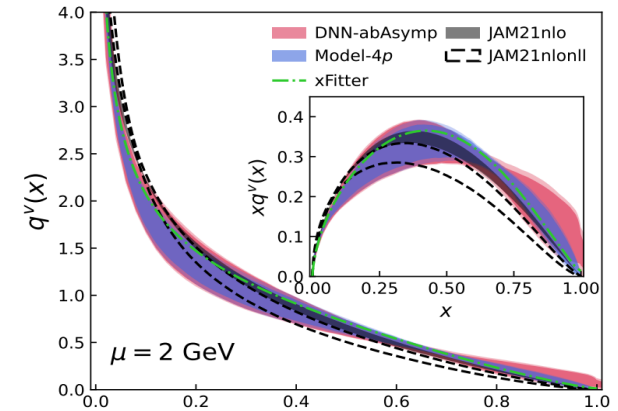
- Achieved great success in the studies of PDF:



*Proton unpolarized PDF, in preparation*

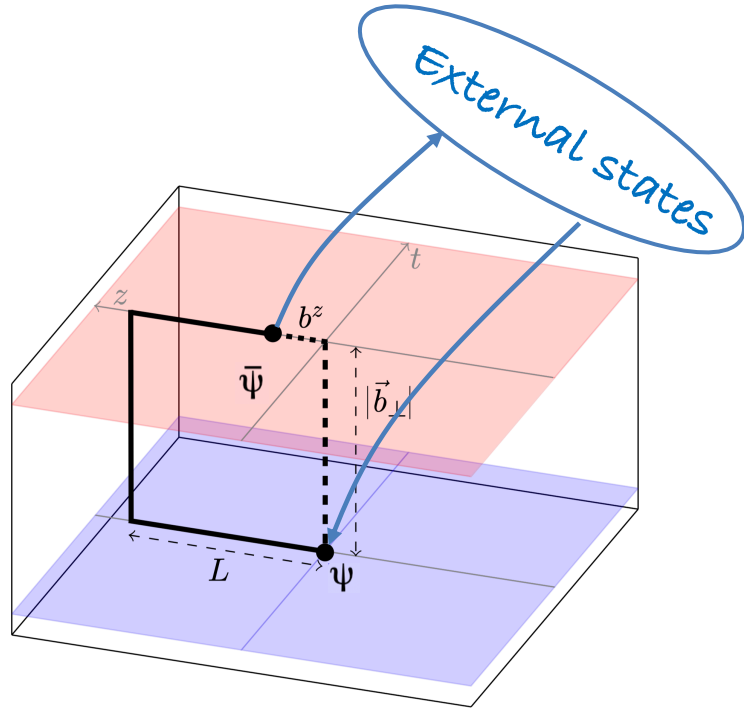


*Proton transversity PDF, 2208.08008*



*Pion valance PDF, PRD106(2022)*

- **Matching from quasi TMDs to TMDs**



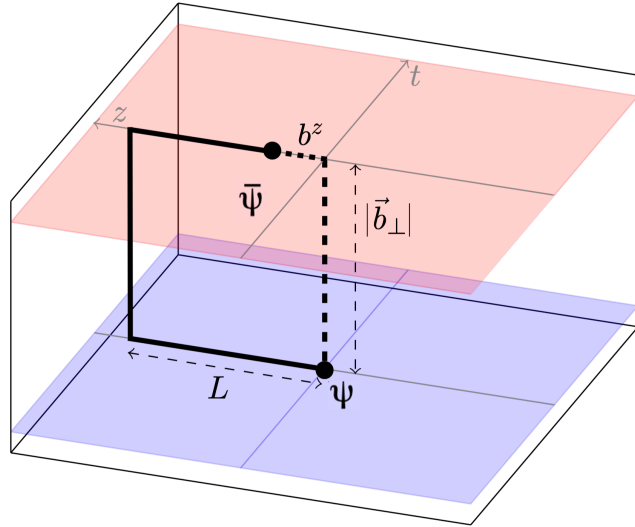
Equal-time correlators  
with staple-shaped Wilson link,  
directly calculable on lattice

- Hadronic matrix element reduced from equal-time correlators:

$$\begin{aligned} \tilde{h}_{\Gamma}^0(z, b_{\perp}, P^z) &= \lim_{L \rightarrow \infty} \langle P^z | \bar{\psi}(b_{\perp} \hat{n}_{\perp}) \Gamma \\ &\quad \times U_{\square}(b_{\perp} \hat{n}_{\perp} \leftarrow b_{\perp} \hat{n}_{\perp} + L \hat{n}_z; b_{\perp} \hat{n}_{\perp} + L \hat{n}_z \leftarrow L \hat{n}_z; L \hat{n}_z \leftarrow z \hat{n}_z) \\ &\quad \times \psi(z \hat{n}_z) | P^z \rangle \end{aligned}$$

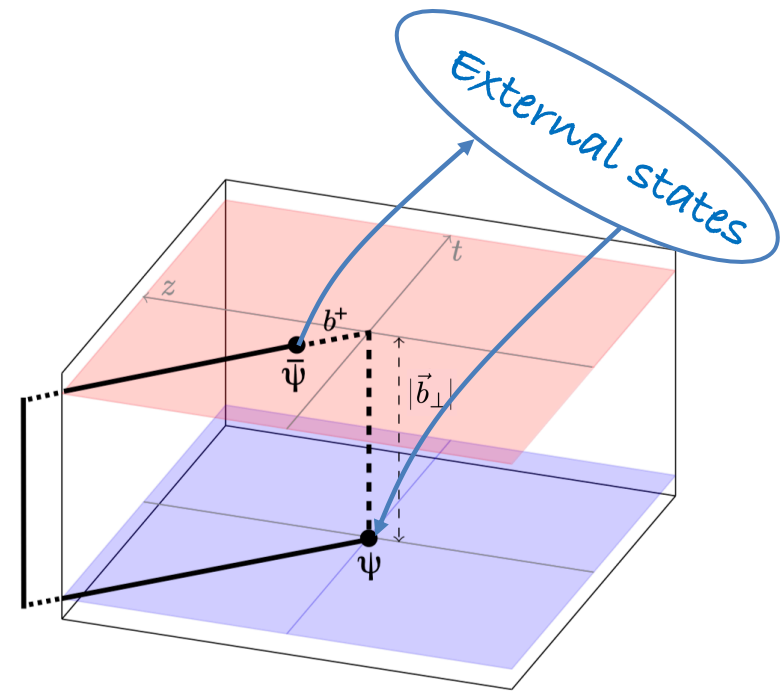
- Subtracted quasi TMDPDFs:

$$\tilde{f}_{\Gamma}(x, b_{\perp}, P^z, \mu) \equiv \lim_{\substack{a \rightarrow 0 \\ L \rightarrow \infty}} \int \frac{dz}{2\pi} e^{-iz(xP^z)} \frac{\tilde{h}_{\Gamma}^0(z, b_{\perp}, P^z, a, L)}{\sqrt{Z_E(2L + z, b_{\perp}, a)} Z_O(1/a, \mu, \Gamma)}$$



Equal-time correlators,  
directly calculable on lattice

Lorentz boost  
 $\longrightarrow$   
 $L \rightarrow \infty$



Space-like correlators,  
NO effective method for directly calculation

### Connected at large-momentum limit

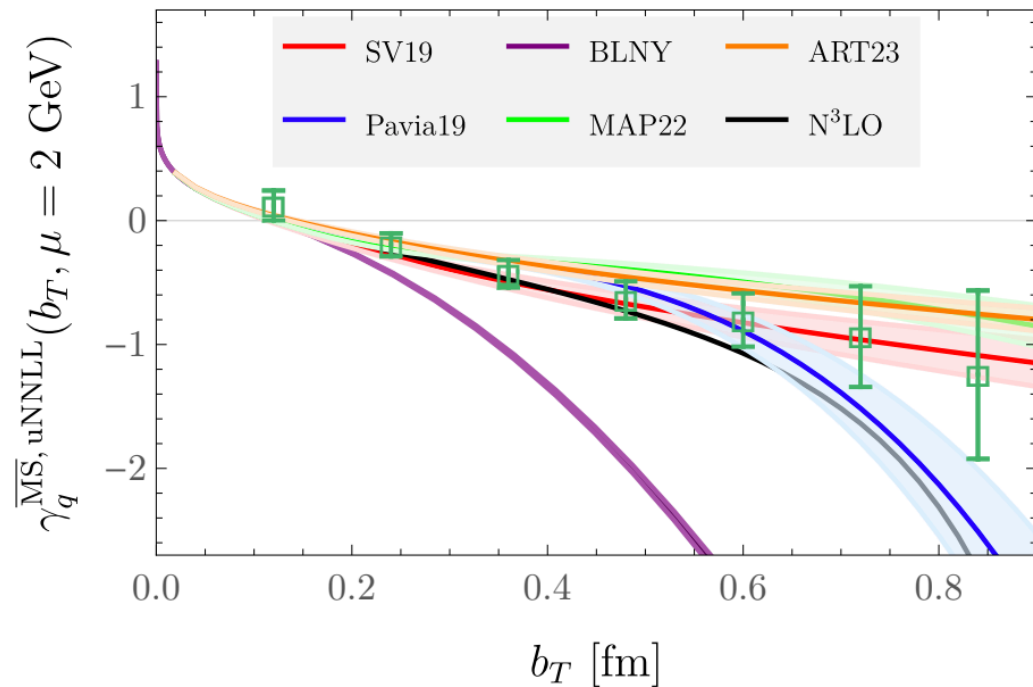
*Ji, PLB811(2020); Ebert, JHEP04(2022)*

$$\underbrace{\tilde{f}_\Gamma(x, b_\perp, \zeta_z, \mu)}_{\text{Quasi TMDPDF}} \underbrace{\sqrt{S_I(b_\perp, \mu)}}_{\text{Intrinsic soft function}} = \underbrace{H_\Gamma\left(\frac{\zeta_z}{\mu^2}\right)}_{\text{Matching kernel}} e^{\frac{1}{2} \ln\left(\frac{\zeta_z}{\mu}\right)} \underbrace{K(b_\perp, \mu)}_{\text{Collins-Soper kernel}} \underbrace{f(x, b_\perp, \mu, \zeta)}_{\text{Light-cone TMDPDF}} + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{\zeta_z}, \frac{M^2}{(P^z)^2}, \frac{1}{b_\perp^2 \zeta_z}\right)$$

# Collins-Soper kernel

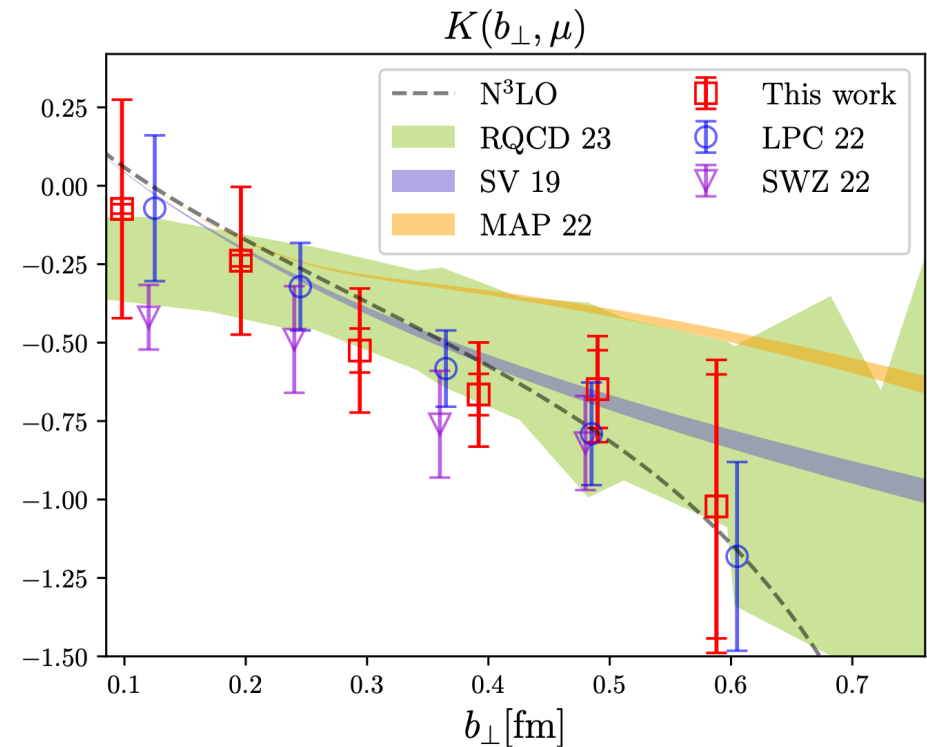
## From quasi beam function:

*Avkhadiev, 2307.12359; Shanahan, PRD104(2021), PRD102(2020); Schlemmer, JHEP08(2021); .....*



## From quasi TMDWF:

*Chu, 2306.06488, PRD106(2022); Zhang, PRL125(2020); Li, PRL128(2022); .....*



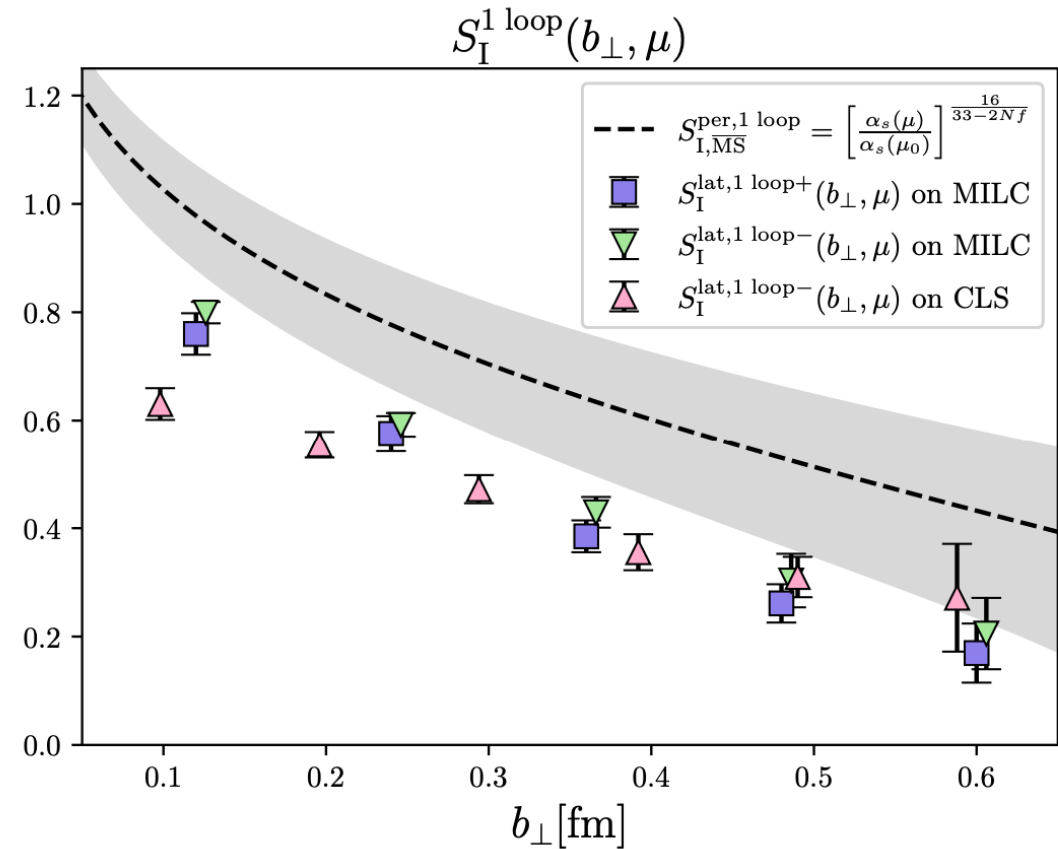
# Intrinsic soft function

From quasi TMDWF + 4-quark matrix element:

*Chu, 2302.09961; Ji, NPB955(2020);*

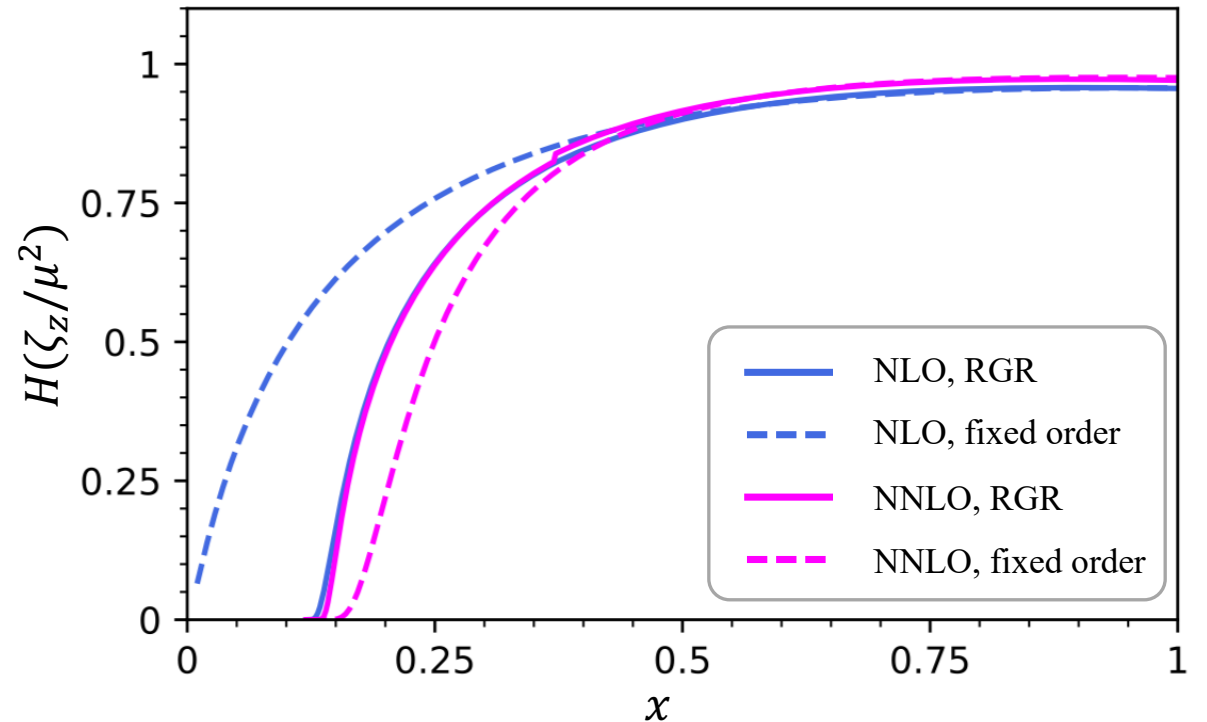
*Zhang, PRL125(2020); Li, PRL128(2022);*

.....



# Perturbative matching kernel

- **NLO:** *Ji, PLB811(2020); RMP93(2021)*
  - **NNLO:** *Rio, 2304.14440; Ji, 2305.04416*
- **Fixed order:**  $\mu = 2\text{GeV}$ ;
- **RGR:** RG evolution from lattice scale  
 $\sqrt{\zeta_z} = 2xP^z$  to  $\overline{\text{MS}}$  scale  $\mu = 2\text{GeV}$ .



# Lattice calculation of physical TMDPDF?

$$\tilde{f}_\Gamma(x, b_\perp, \zeta_z, \mu) \sqrt{S_I(b_\perp, \mu)} = H_\Gamma\left(\frac{\zeta_z}{\mu^2}\right) e^{\frac{1}{2} \ln\left(\frac{\zeta_z}{\zeta}\right)} K(b_\perp, \mu) f(x, b_\perp, \mu, \zeta) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{\zeta_z}, \frac{M^2}{(P^z)^2}, \frac{1}{b_\perp^2 \zeta_z}\right)$$

Quasi TMDPDF

Intrinsic soft function

Collins-Soper kernel



**Simulating quasi TMDPDF on a Euclidean lattice:**

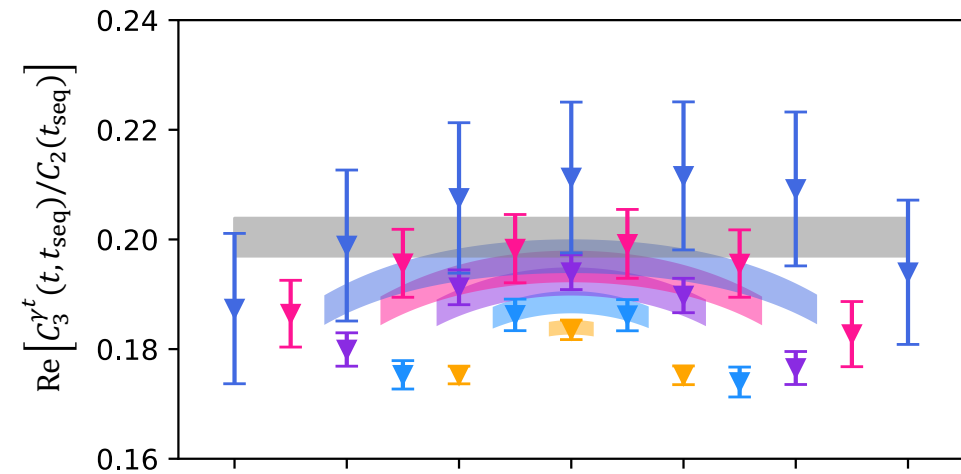
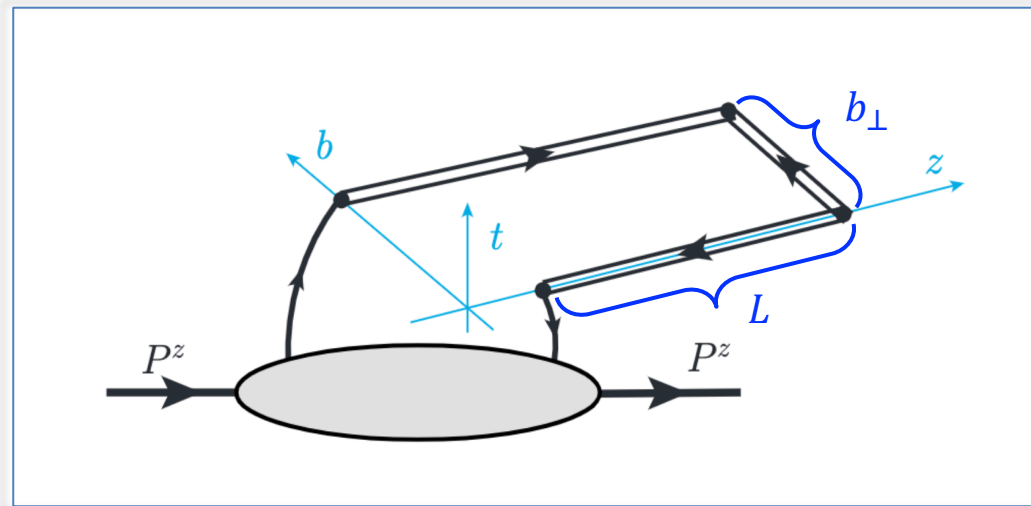
- MILC configuration:  $48^3 \times 64$ ,  $a = 0.12\text{fm}$ ;
- Pion mass:  $m_\pi^{\text{sea}} = 130\text{MeV}$ ,  $m_\pi^{\text{val}} = \{310, 220\}\text{MeV} \Rightarrow$  extrapolate to physical mass
- Large momentum:  $P^z = \{1.72, 2.15, 2.58\}\text{GeV} \Rightarrow$  extrapolate to infinity
- Saturated length of Wilson link  $L = 0.72\text{fm}$ ;
- $z_{\text{max}} = 1.44\text{fm}$ ,  $b_{\perp\text{max}} = 0.6\text{fm} \Rightarrow$  momentum distribution.



# Quasi TMDPDF matrix element

## Bare quasi TMDPDF matrix element

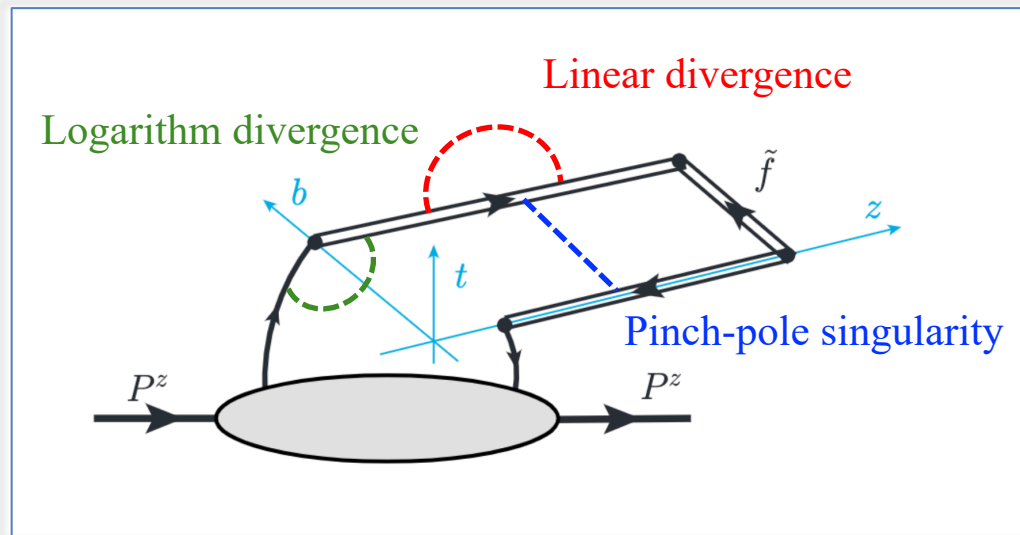
$$\tilde{h}_{\Gamma}^0(z, b_{\perp}, P^z) = \lim_{L \rightarrow \infty} \left\langle P^z \left| \bar{\psi}(b_{\perp} \hat{n}_{\perp}) \Gamma U_{\square}(b_{\perp} \hat{n}_{\perp} \leftarrow b_{\perp} \hat{n}_{\perp} + L \hat{n}_z; b_{\perp} \hat{n}_{\perp} + L \hat{n}_z \leftarrow L \hat{n}_z; L \hat{n}_z \leftarrow z \hat{n}_z) \psi(z \hat{n}_z) \right| P^z \right\rangle$$



- Extracted from 3- and 2-point functions

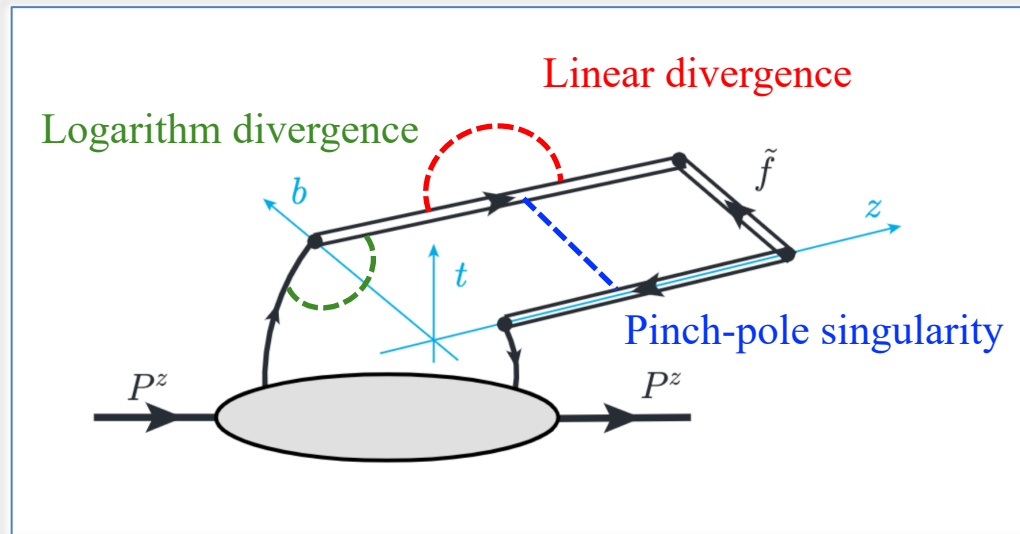
# Quasi TMDPDF matrix element

## 1. Divergences in bare quasi TMDPDF

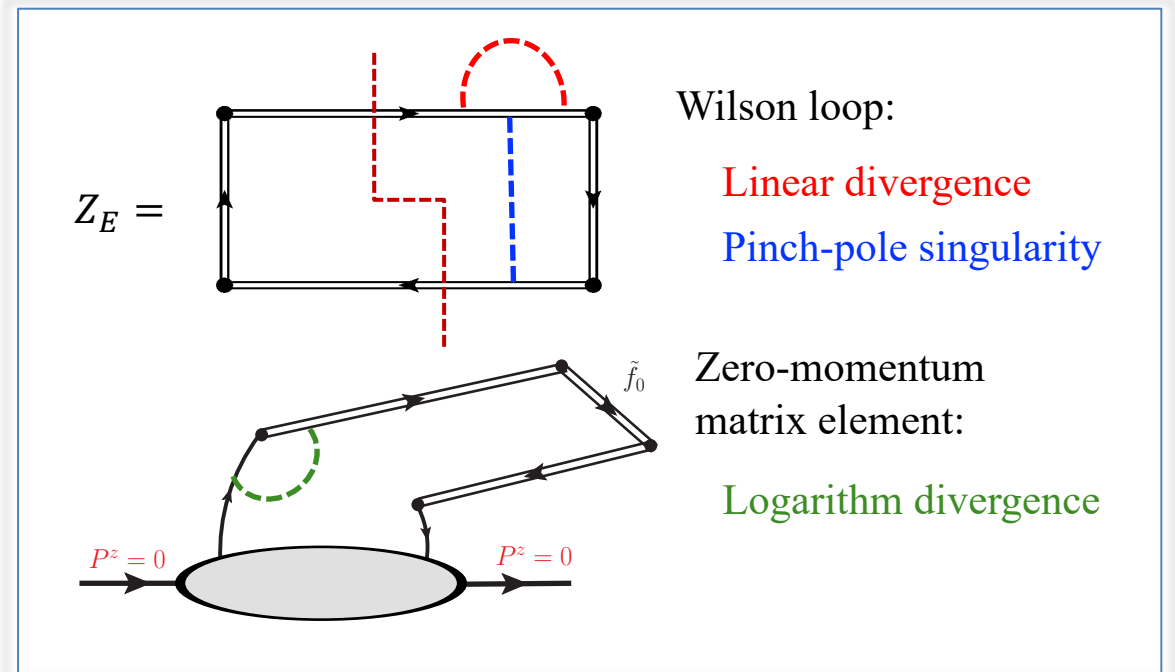


# Quasi TMDPDF matrix element

## 1. Divergences in bare quasi TMDPDF



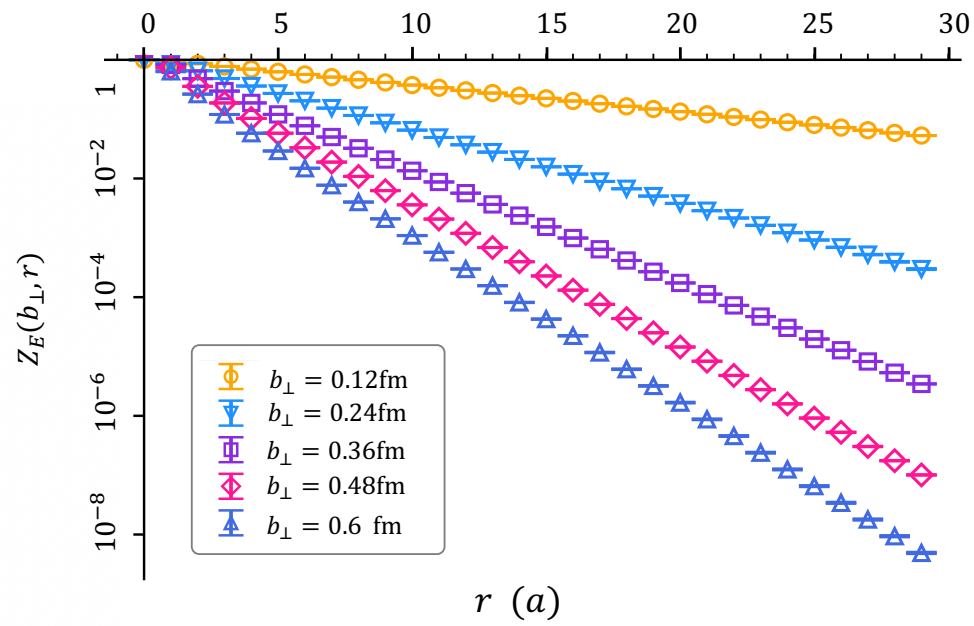
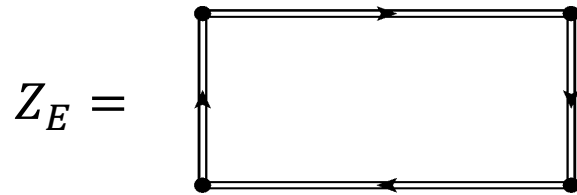
## 2. Renormalization



*Ji, PRL120(2018), NPB964(2021), PLB257(1991); Zhang, PRD95(2017), NPB939(2019); Ishikawa, PRD96(2017); Green, PRL121(2018); Huo, NPB969(2021); Chen, NPB915(2017); Musch, PRD83(2011); .....*

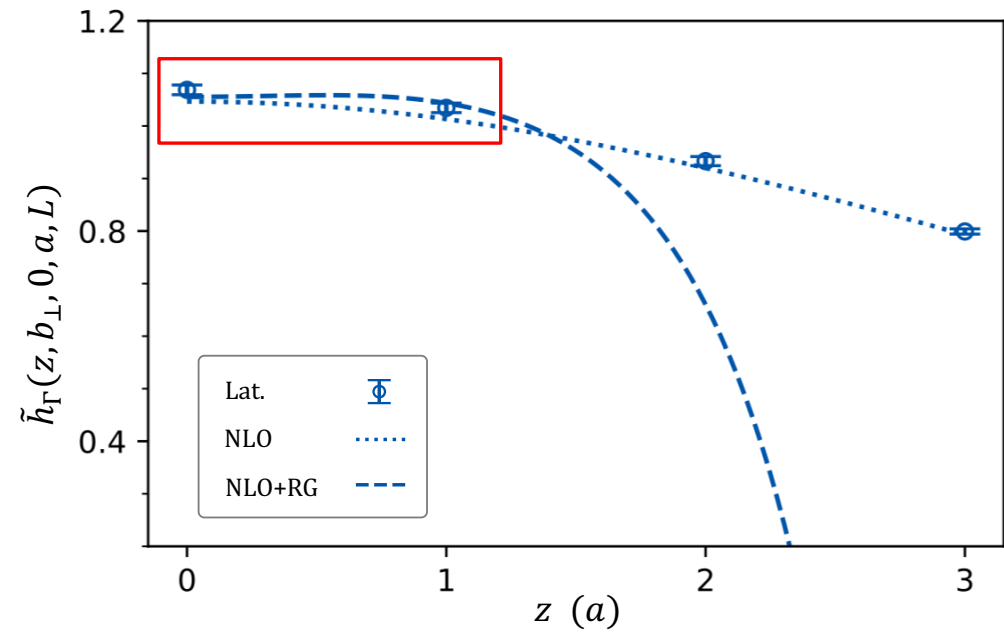
# Quasi TMDPDF matrix element

- Wilson loop



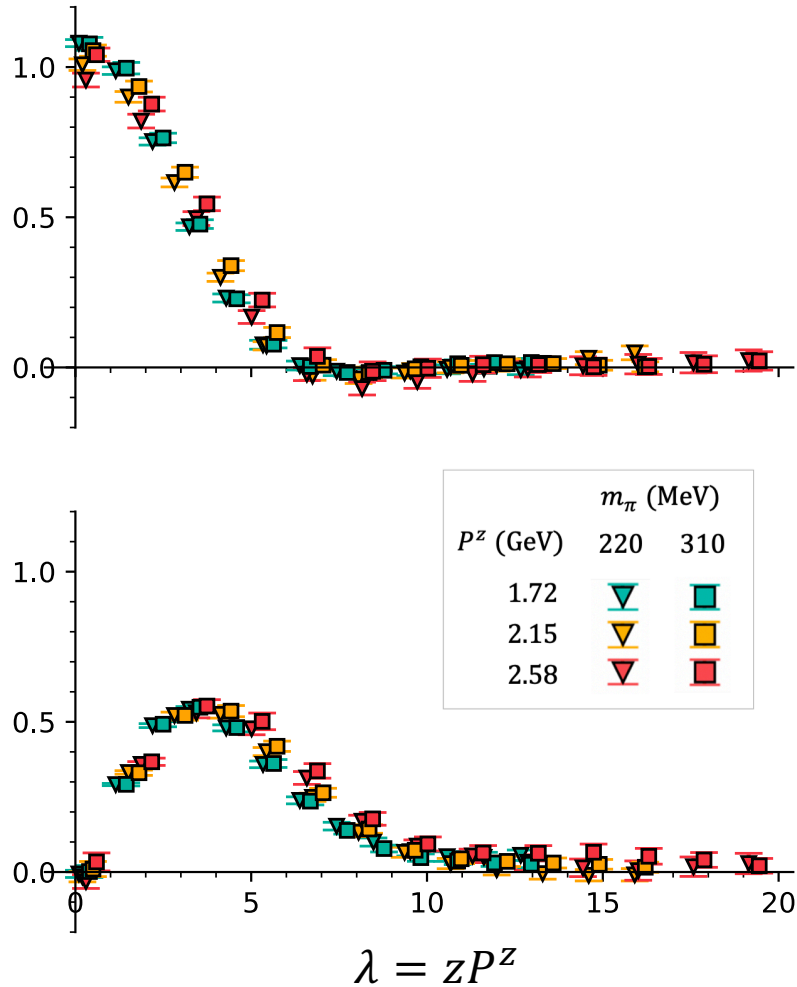
- Logarithmic divergences factor

$$Z_O(1/a, \mu, \Gamma) = \lim_{L \rightarrow \infty} \frac{\tilde{h}_{\Gamma}^0(z, b_{\perp}, 0, a, L)}{\sqrt{Z_E(2L + z, b_{\perp}, a)} \tilde{h}_{\Gamma}^{\overline{\text{MS}}}(z, b_{\perp}, \mu)}$$

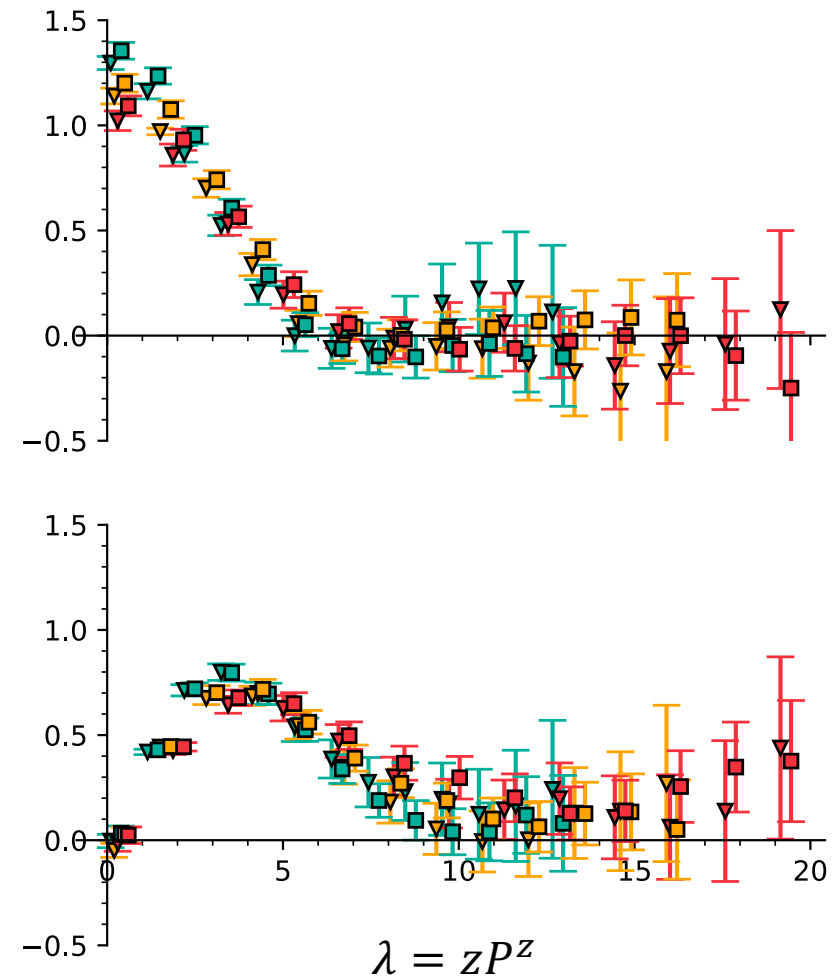


# Quasi TMDPDF matrix element

Quasi TMDPDF matrix element at  $b_{\perp} = 2a$



Quasi TMDPDF matrix element at  $b_{\perp} = 5a$



# Quasi TMDPDF matrix element

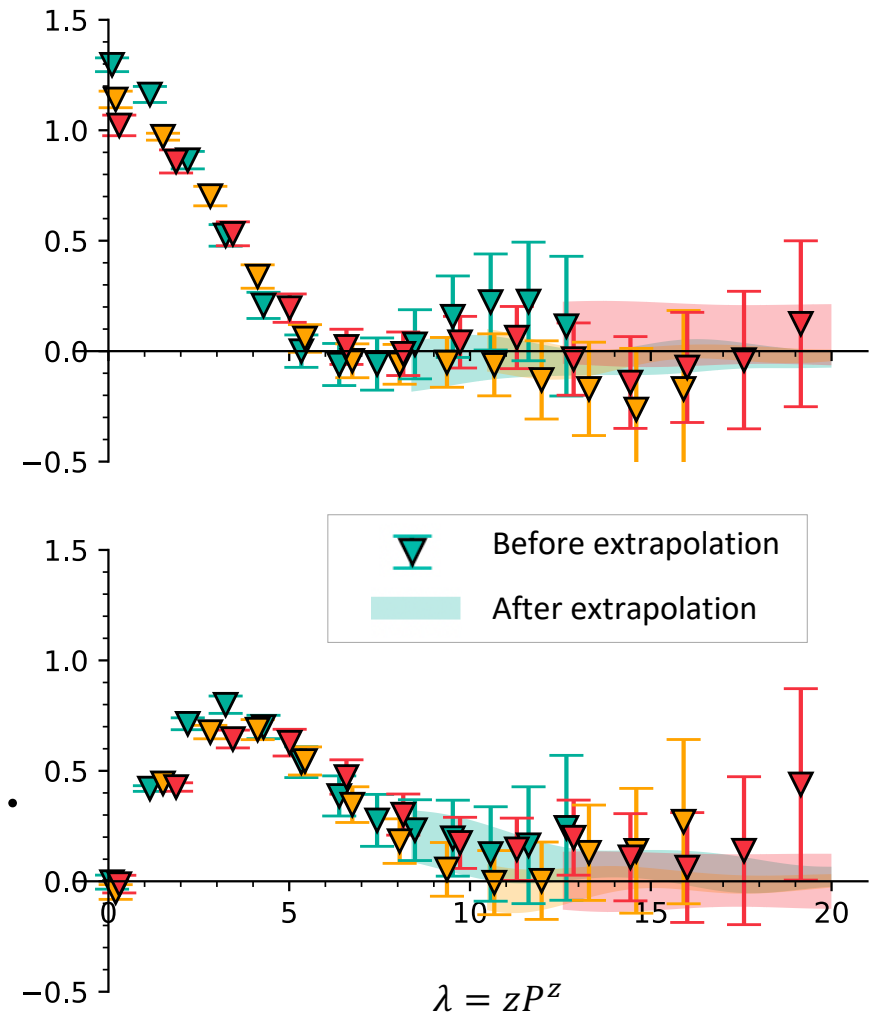
## Physical based extrapolation:

$$\tilde{h}_{\text{extra}}(\lambda) = \left[ \frac{c_1}{(-i\lambda)^{n_1}} + e^{i\lambda} \frac{c_2}{(i\lambda)^{n_2}} \right] e^{-\lambda/\lambda_0}$$

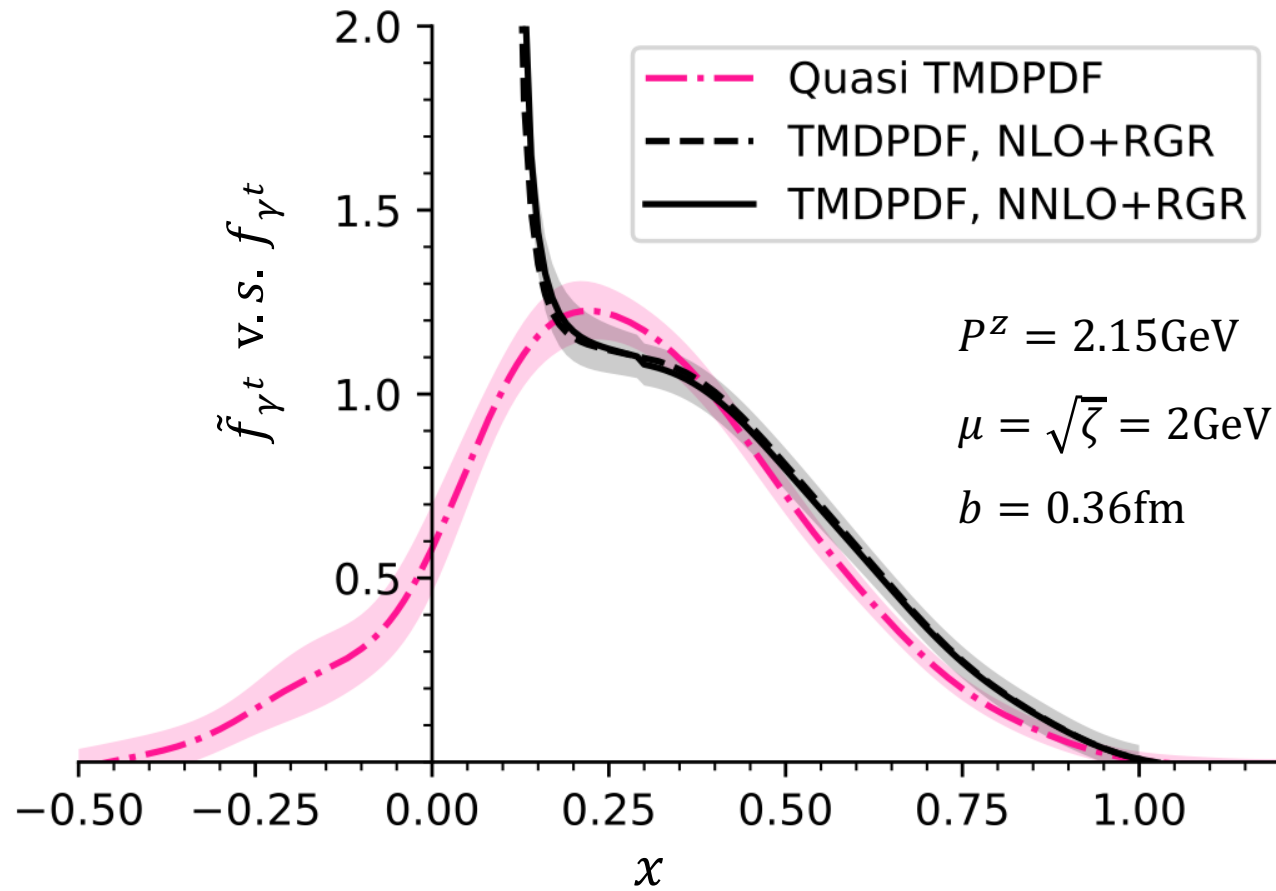
- end point power-law behavior  $x^a(1-x)^b$ ;
- correlation function has a finite correlation length  $\lambda_0$ .

😓 End point region in momentum space will be affected .....

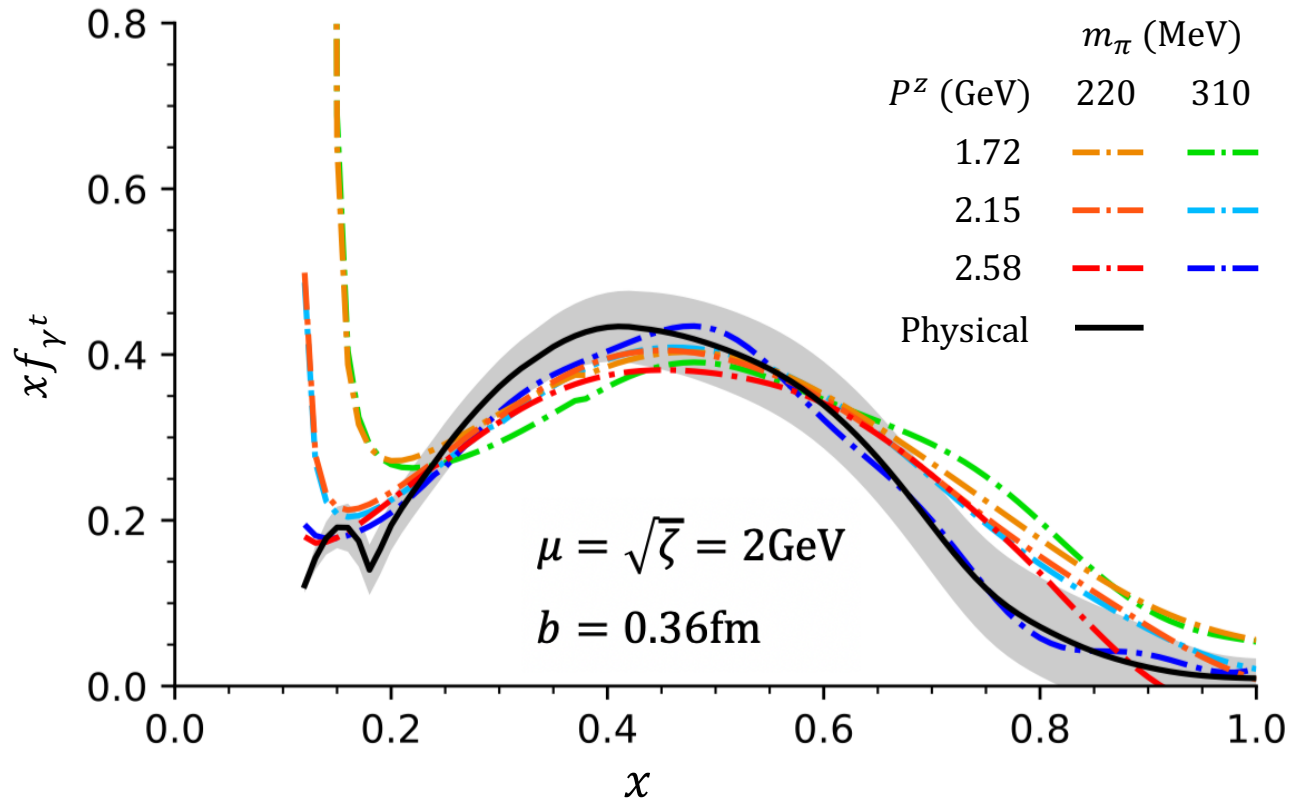
😊 LaMET is not reliable in that region (power correction)



# From Quasi TMDPDF to TMDPDF



# Physical TMDPDF



## Physical extrapolation:

$$f^{\text{lat}}(x, b_\perp, \mu, \zeta; m_\pi, P^z) =$$

$$f^{\text{phy}}(x, b_\perp, \mu, \zeta) * [1 + \text{Corrections}]$$

$$d_0(m_\pi^2 - m_{\pi, \text{phy}}^2) + \frac{d_1}{(P^z)^2} \quad \text{Main form}$$

$$d_0(m_\pi^2 - m_{\pi, \text{phy}}^2)^2 + \frac{d_1}{(P^z)^2}$$

$$d_0(m_\pi^2 - m_{\pi, \text{phy}}^2) + \frac{d_1}{(P^z)^2} + \frac{d_2}{P^z}$$

Used for estimating systematical errors



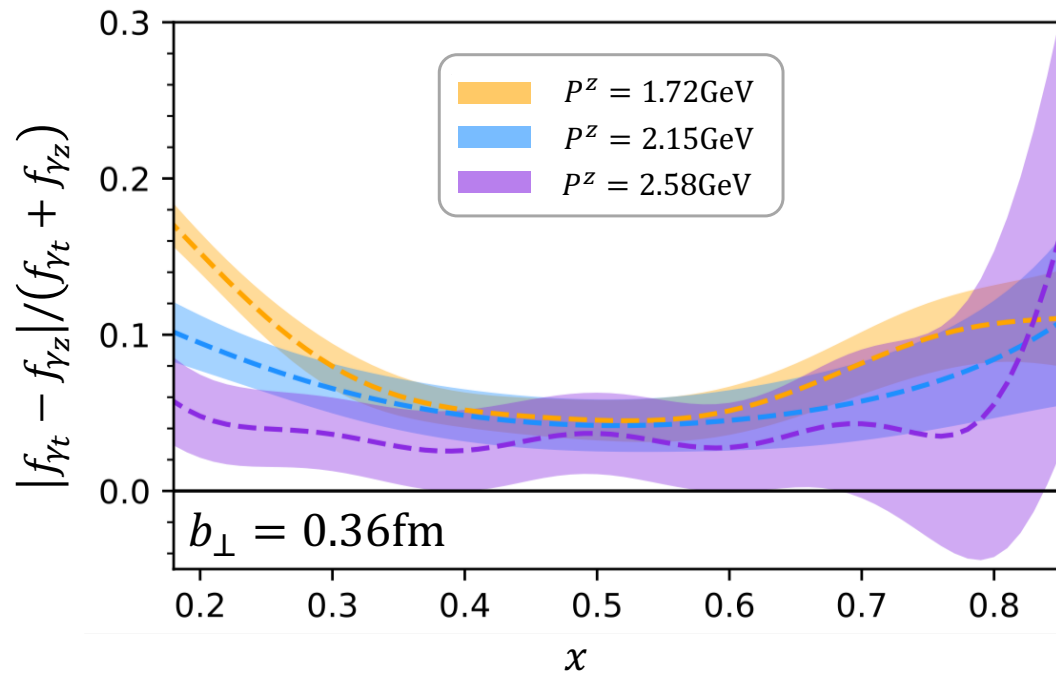
# Error estimation

After Lorentz boost:

Leading power

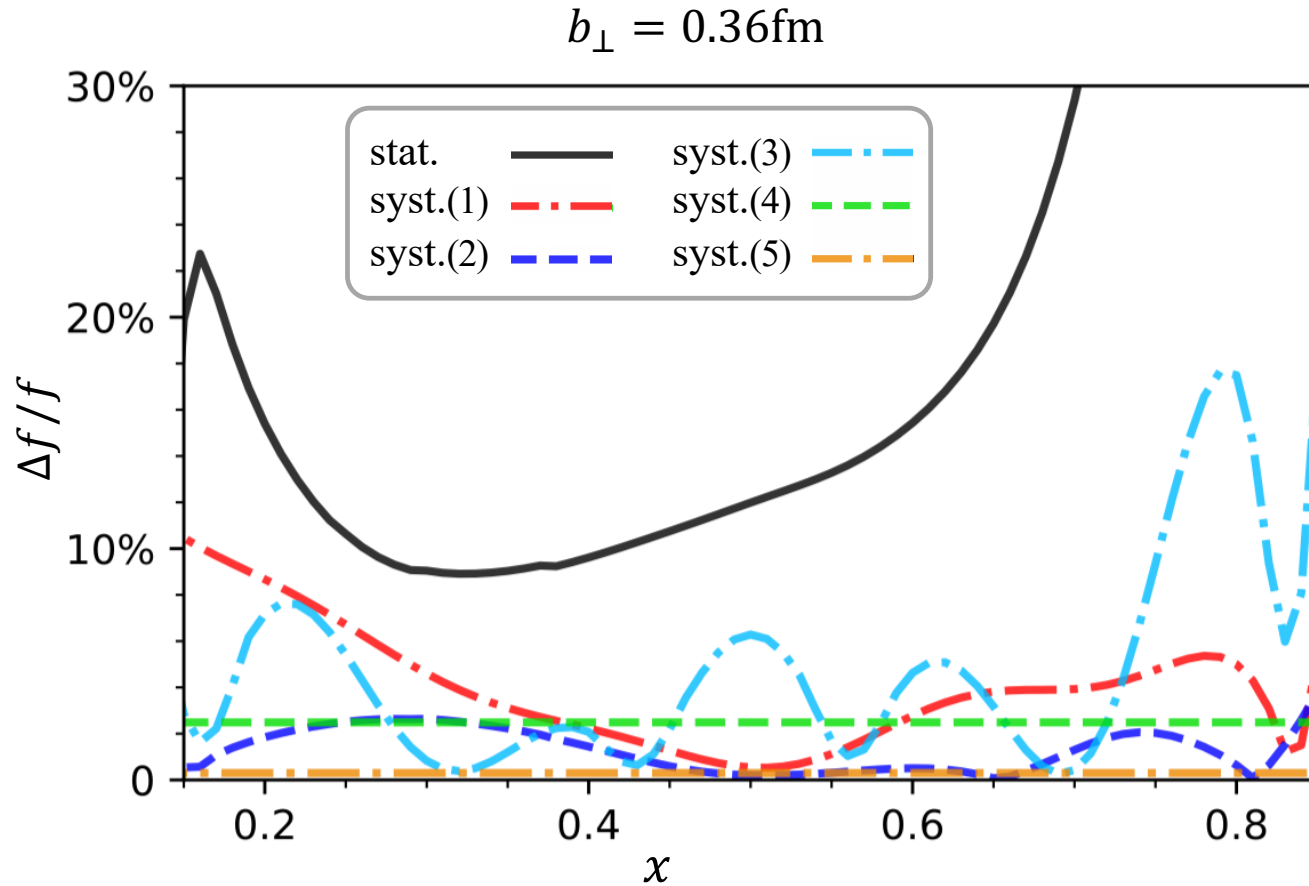
Higher power

$$\begin{aligned}\bar{\psi}(z)\gamma^t\psi(0) &= \frac{1}{2}\bar{\psi}(z)\gamma^+\psi(0) + \frac{1}{2}\bar{\psi}(z)\gamma^-\psi(0) \\ \bar{\psi}(z)\gamma^z\psi(0) &= \frac{1}{2}\bar{\psi}(z)\gamma^+\psi(0) - \frac{1}{2}\bar{\psi}(z)\gamma^-\psi(0)\end{aligned}$$



- Ratios denote the deviations from light-like correlator with specific  $P^z$ ;
- Ratio becomes smaller with  $P^z$  increasing.

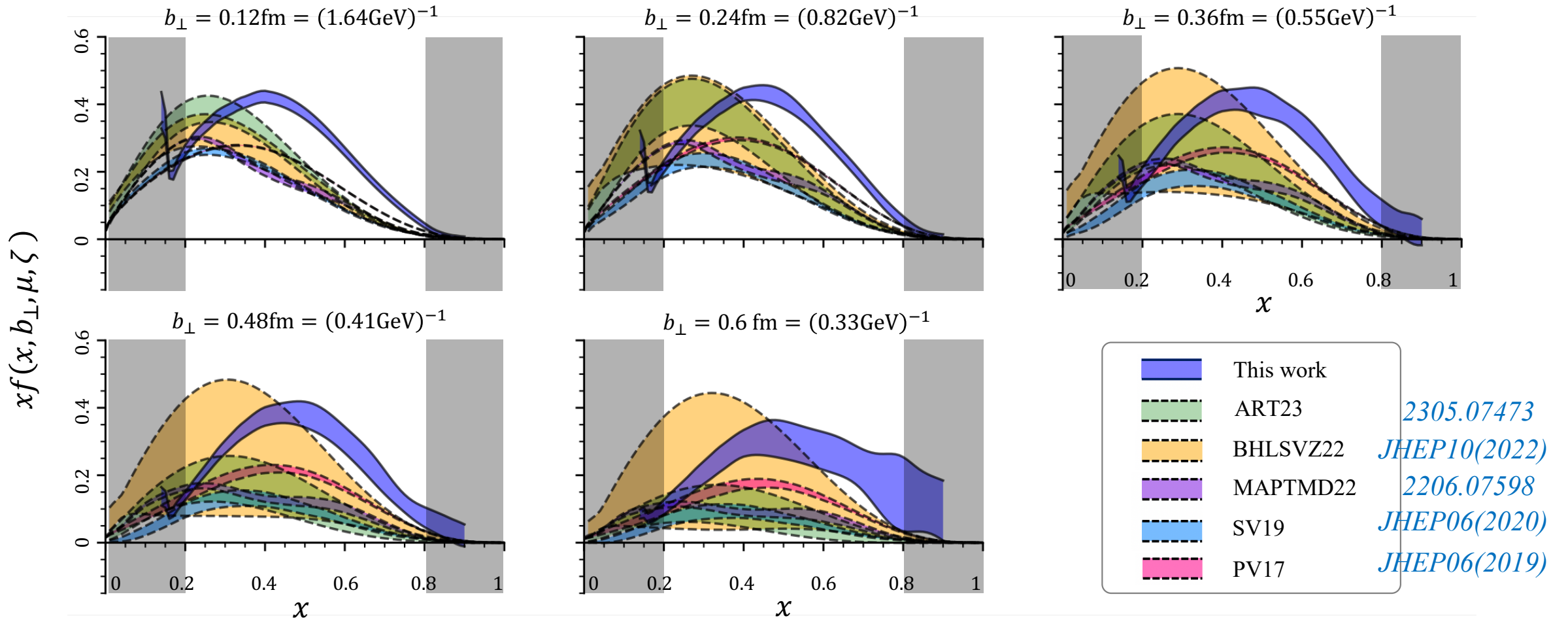
# Error estimation



## All errors:

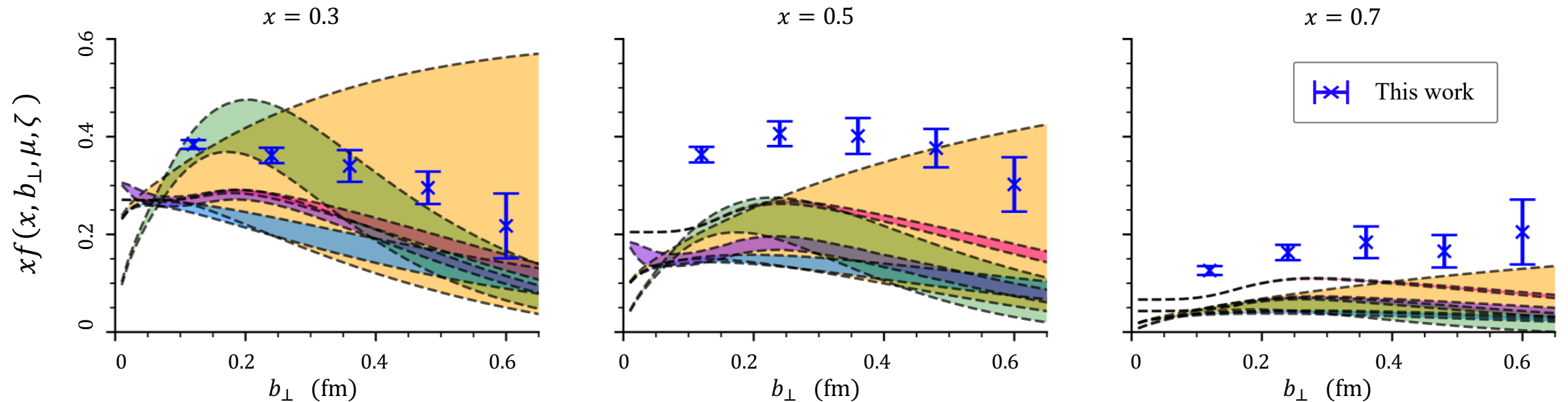
- **Statistical error;**
- (1) From power correction**
- (2) From physical extrapolation**
- (3) From  $\lambda$ -extrapolation**
- (4) From soft function**
- (5) From Collins-Soper kernel**

# Final results and discussion



# Final results and discussion

Compare the  $b_{\perp}$ -dependence of lattice and phenomenological results:



## Final results and discussion



**Lattice discretization and finite-volume systematics are still absent in this preliminary work...**

- It is a challenging work for calculating the TMDPDF at small lattice spacing
- From the previous experience of PDF ([Lin, 2011.14971](#)), we can roughly estimate that:

**Finite-volume effect is less than 1%;**

**Discretization effects overall within 2 standard deviations.**

- Furthermore, need more efforts to achieve the precise calculation on finer lattice.

## Summary and outlook

**We present the lattice QCD calculation of TMDPDF at first attempt:**

- ✓ **The state-of-the-art techniques in renormalization and extrapolation on the lattice;**
- ✓ **The latest perturbative kernel up to 2-loop with RG evolution;**
- ✓ **Physical extrapolation include chiral-continuum and infinity momentum;**
- ✓ **Comparable results with phenomenological global fits.**

## Summary and outlook

While there is still much room for further improvement:

- 🤔 **Better control of uncertainties: *more statistic, first and most important*;**
- 🤔 **Larger  $b_{\perp}$  (up to nucleon radius?):**  
**More statistic, improved algorithm, better lattice extrapolation, .....**
- 🤔 **Continuum extrapolation: more lattice spacings;**
- 🤔 **Theoretical improvement: power correction (small- $x$  region), operator mixing, .....**

*Thank you for your attention!*

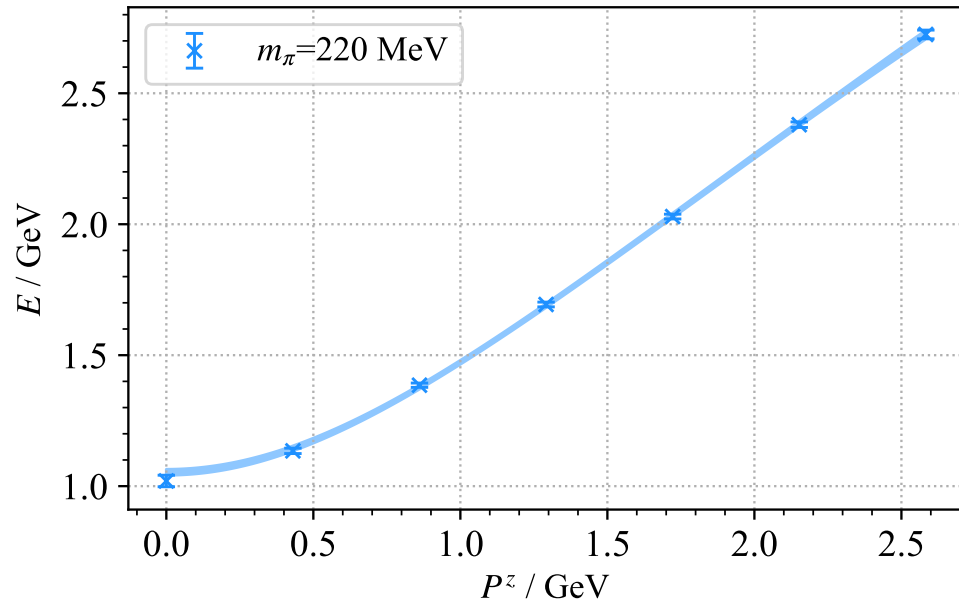
# Backup slides



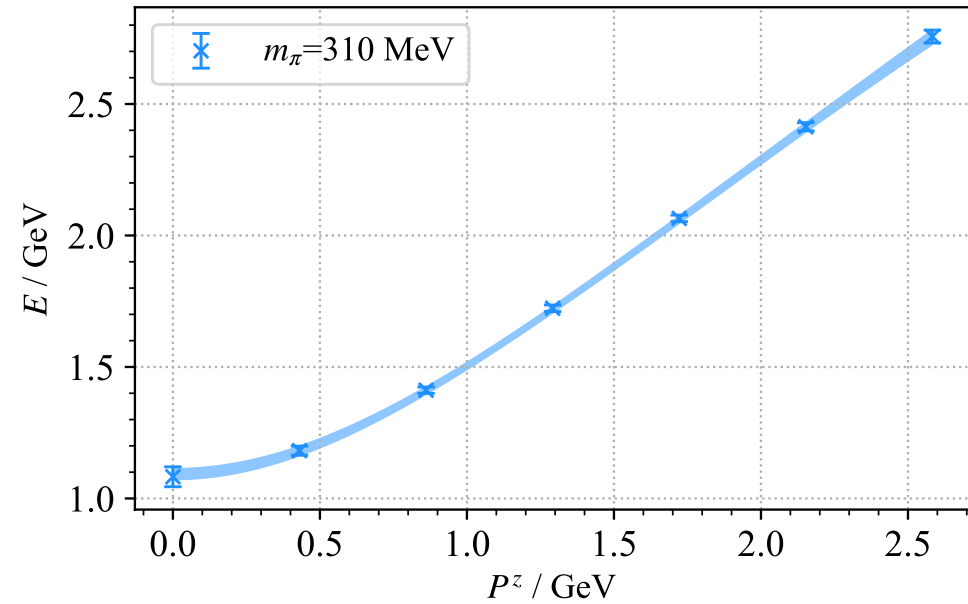
# Dispersion relation

$$E = \sqrt{m^2 + c_1(P^z)^2 + c_2(P^z)^4 a^2}$$

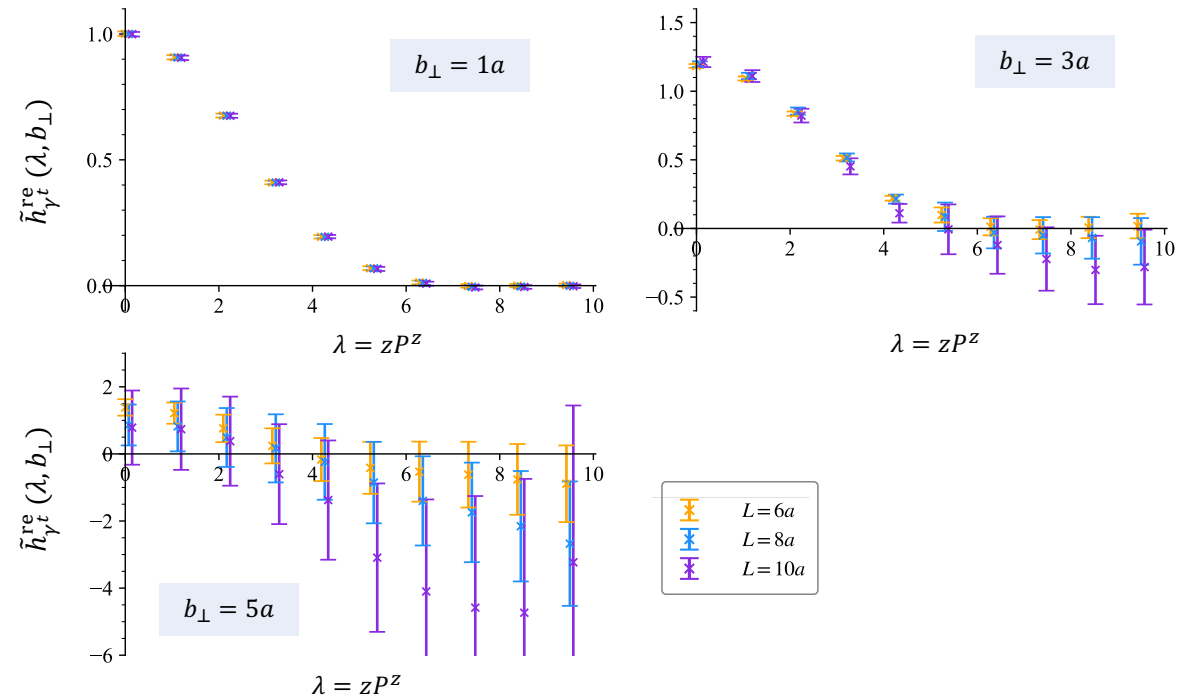
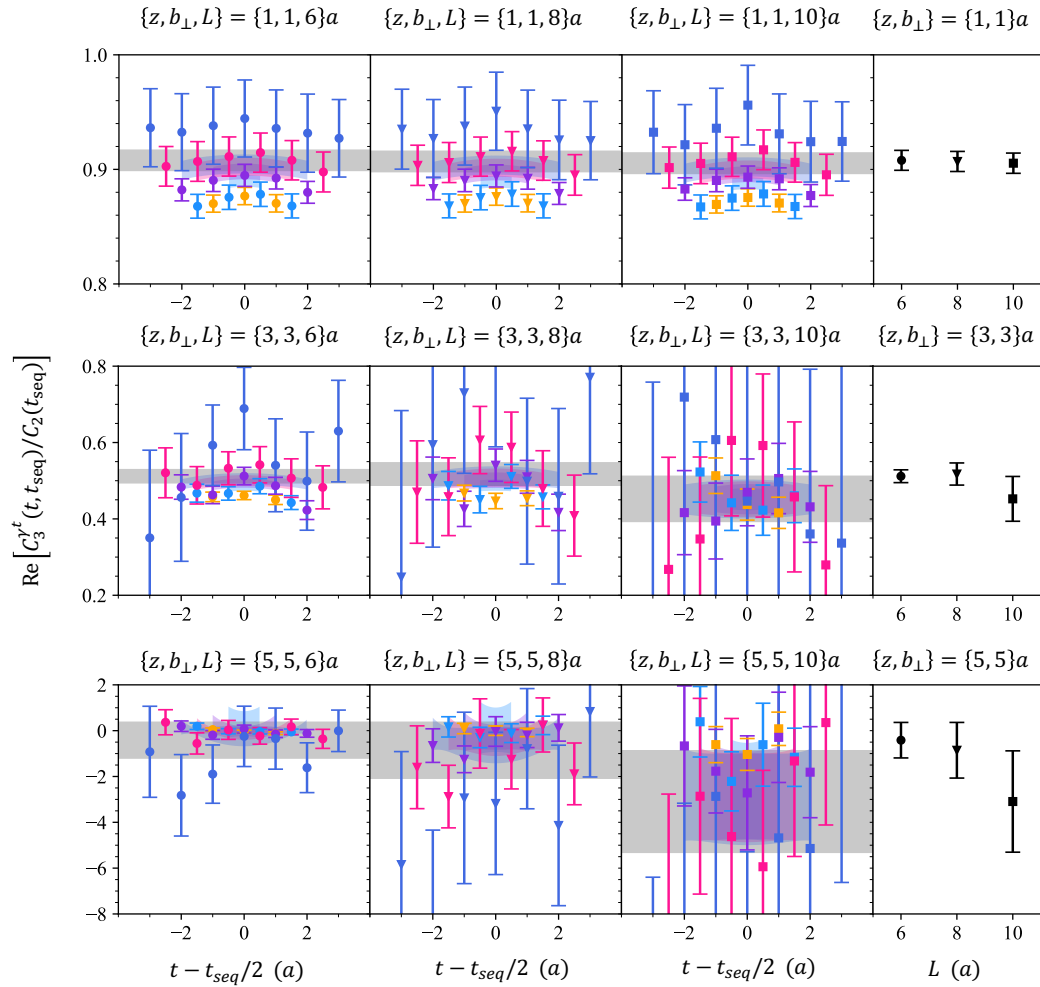
$$c_1 = 1.081(20), \quad c_2 = -0.0548(96)$$



$$c_1 = 1.087(28), \quad c_2 = -0.053(13)$$

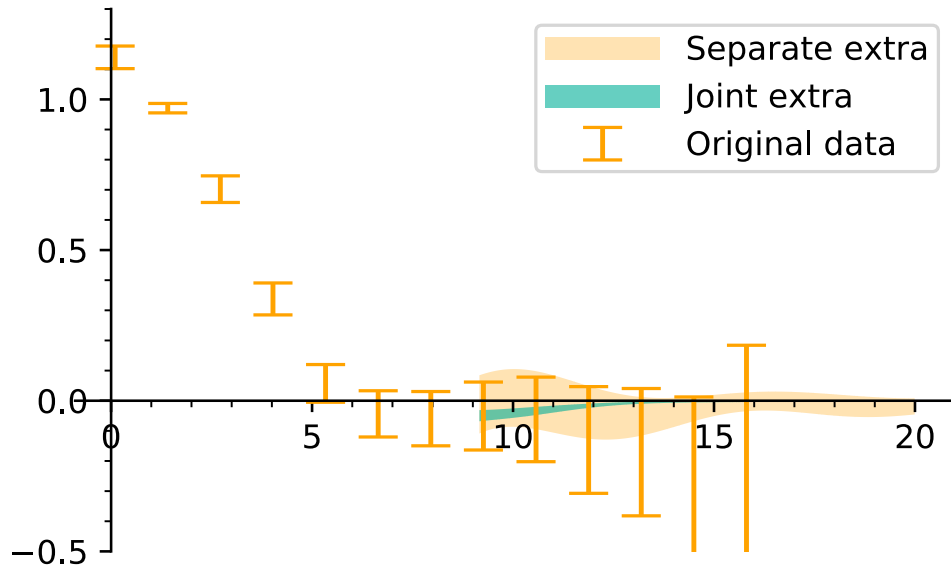


# L-dependence



# $\lambda$ -extrapolation

$$\tilde{h}_{\text{extra}}(\lambda) = \left[ \frac{c_1}{(-i\lambda)^{n_1}} + e^{i\lambda} \frac{c_2}{(i\lambda)^{n_2}} \right] e^{-\lambda/\lambda_0}$$



The **power-law behavior** and **correlation length** for each  $b_{\perp}$  should be similar,

but the joint fit will give a strict limit for large- $b_{\perp}$  cases:

|           | $m_1$     | $m_2$     | $n_1$     | $n_2$    | $\lambda_0$ | $\chi^2/\text{d.o.f}$ |
|-----------|-----------|-----------|-----------|----------|-------------|-----------------------|
| $b = 1a$  | -8.8(3.7) | 0.25(50)  | 0.909(39) | 1.13(74) | 2.63(38)    | 1.0                   |
| $b = 2a$  | -6.3(2.9) | -3.9(5.5) | 0.943(61) | 2.37(68) | 4.1(1.1)    | 1.1                   |
| $b = 3a$  | -66(60)   | -78(76)   | 0.89(12)  | 1.71(31) | 2.42(85)    | 1.4                   |
| $b = 4a$  | -8.0(4.4) | -3.3(2.9) | 0.801(78) | 1.55(38) | 4.3(1.6)    | 0.75                  |
| $b = 5a$  | -8.5(10)  | -3.8(5.3) | 0.84(16)  | 1.22(44) | 4.4(2.8)    | 0.57                  |
| Joint fit | -         | -         | 0.887(28) | 1.65(12) | 2.53(28)    | 1.2                   |

# Perturbative matching kernel and RG resummation

- **Fixed-order perturbative results up to the 2-loop level:**

$$h^{(1)}\left(\frac{\zeta_z}{\mu^2}\right) = \frac{\alpha_s C_F}{2\pi} \left( -2 + \frac{\pi^2}{12} + \ln \frac{\zeta_z}{\mu^2} - \frac{1}{2} \ln^2 \frac{\zeta_z}{\mu^2} \right),$$

$$h^{(2)}\left(\frac{\zeta_z}{\mu^2}\right) = \alpha_s^2 \left[ c_2 - \frac{1}{2} \left( \gamma_C^{(2)} - \beta_0 c_1 \right) \ln \frac{\zeta_z}{\mu^2} - \frac{1}{4} \left( \Gamma_{\text{cusp}}^{(2)} - \frac{\beta_0 C_F}{2\pi} \right) \ln^2 \frac{\zeta_z}{\mu^2} - \frac{\beta_0 C_F}{24\pi} \ln^3 \frac{\zeta_z}{\mu^2} \right]$$

- **RG equation of the matching kernel:**

$$\mu^2 \frac{d}{d\mu^2} \ln H\left(\frac{\zeta_z}{\mu^2}\right) = \frac{1}{2} \Gamma_{\text{cusp}}(\alpha_s) \ln \frac{\zeta_z}{\mu^2} + \frac{\gamma_C(\alpha_s)}{2},$$

**and its solution:**

$$H\left(\frac{\zeta_z}{\mu^2}\right) = H\left(\frac{\zeta_z}{\mu_0^2}\right) \exp \left[ \int_{\mu_0}^{\mu} \frac{d\mu}{\mu} \left( \Gamma_{\text{cusp}}^{(1)} \ln \frac{\zeta_z}{\mu^2} \alpha_s(\mu) + \gamma_C^{(1)} \alpha_s(\mu) + \Gamma_{\text{cusp}}^{(2)} \ln \frac{\zeta_z}{\mu^2} \alpha_s^2(\mu) \right. \right. \\ \left. \left. + \gamma_C^{(2)} \alpha_s^2(\mu) + \Gamma_{\text{cusp}}^{(3)} \ln \frac{\zeta_z}{\mu^2} \alpha_s^3(\mu) + \gamma_C^{(3)} \alpha_s^3(\mu) + \Gamma_{\text{cusp}}^{(4)} \ln \frac{\zeta_z}{\mu^2} \alpha_s^4(\mu) \right) \right].$$