

The hadronic contribution to $\Delta\alpha_{\text{QED}}$

Teseo San José^{1 a,b} Marco Cè^c Antoine Gérardin^d
Harvey B. Meyer^{a,b} Kohtaroh Miura^{a,b,e} Konstantin Ottnad^b
Andreas Risch^b Jonas Wilhelm^b Hartmut Wittig^{a,b}

^aHelmholtz Institute Mainz, 55099 Mainz, Germany and GSI Helmholtzzentrum für Schwerionenforschung, 64291 Darmstadt, Germany

^bPRISMA⁺ Cluster of Excellence and Institute for Nuclear Physics, Johannes Gutenberg University of Mainz, 55099 Mainz, Germany

^cTheoretical Physics Department, CERN, 1211 Geneva 23, Switzerland

^dAix Marseille Univ, Université de Toulon, CNRS, CPT, Marseille, France

^eKobayashi–Maskawa Institute for the Origin of Particles and the Universe (KMI), Nagoya University, Nagoya 464-8602, Japan

Mainz, 20/11/2020

We parametrize the running of the QED coupling as

$$\alpha(Q^2) = \frac{\alpha}{1 - \Delta\alpha(Q^2)}$$

The **hadronic contribution** to $\Delta\alpha$ is

$$\Delta\alpha_{\text{had}}(Q^2) = 4\pi\alpha\bar{\Pi}(Q^2),$$

In the **time-momentum representation** (TMR) [Francis et al. 2013a; Bernecker and Meyer 2011],

$$\bar{\Pi}(Q^2) = \int_0^\infty dx_0 G(x_0) \left[x_0^2 - \frac{4}{Q^2} \sin^2 \left(\frac{Qx_0}{2} \right) \right],$$

with the **electromagnetic correlator** (b-quark contribution neglected)

$$G(x) = G^{33}(x) + \frac{1}{3}G^{88}(x) + \frac{4}{9}C^{c,c}(x),$$

$$G^{33}(x) = \frac{1}{2}C^{\ell,\ell}(x),$$

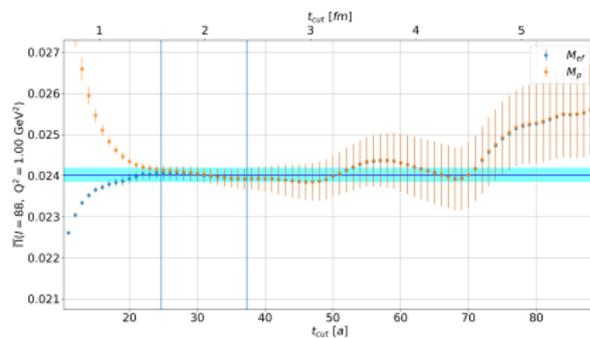
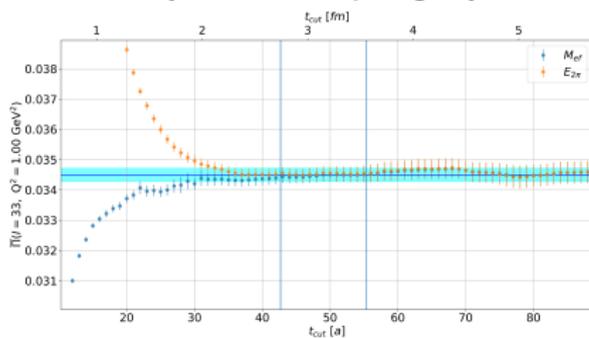
$$G^{88}(x) = \frac{1}{6} \left[C^{\ell,\ell}(x) + 2C^{s,s}(x) + 2D^{\ell-s,\ell-s}(x) \right]$$

In our study we use the same $N_f = 2 + 1$ set of CLS ensembles and similar techniques as in the $g - 2$ analysis. [A. Gerardin talk on Monday]

Error budget on ensemble E250 at 1 GeV²

Ensemble at the physical pion mass ($M_\pi \sim 130$ MeV) with $L = 6.2$ fm, $M_\pi L = 4.1$ and $96^3 \times 192$ lattice sites:

- 1 The initial **statistical error** of the renormalized, $O(a)$ -improved [Gerardin, Harris, and Meyer 2019] correlators on ensemble E250 is:
 - Isvector: 1%
 - Isoscalar: 5%
 - Charm: 0.05%
- 2 We apply the **bounding method** [Gérardin et al. 2019; Blum et al. 2018] to isovector and isoscalar only. Errors drop slightly below 1%.



- 3 **Finite volume** correction, Hansen-Patella method [Hansen and Patella 2020, 2019] and MLLGS-method [Francis et al. 2013b; Meyer 2011; Lellouch and Luscher 2001]
 - Isvector: $+36 \times 10^{-5}$ (1%).
- 4 The size of the **discretization effects** is $\sim 5\%$

After chiral and continuum extrapolation, at the isospin symmetric physical point,

$Q^2[\text{GeV}^2]$	$\bar{\pi}^{33} \times 10^5$	$\bar{\pi}^{88} \times 10^5$	$\bar{\pi}^c \times 10^5$	$\Delta\alpha^{\gamma\gamma} \times 10^6$	$\Delta\alpha_{\text{error}}^{\gamma\gamma} \times 10^6$
1.00	3264 (18)(30)(6)	739 (5)(9)(0)	176.6 (0.8)(3.6)(0.2)	3832 (19)(39)(6)(28)	52 (1.4)
6.00	5815 (42)(30)(29)	1549 (14)(10)(8)	820.6 (2.7)(13.1)(0.9)	7506 (50)(49)(34)(34)	85 (1.1)

1 Statistical error

- a^3 lattice artifacts included for $Q^2 > 1 \text{ GeV}^2$.
- It grows with the energy.

2 Scale setting error

- Estimated using standard propagation of errors.

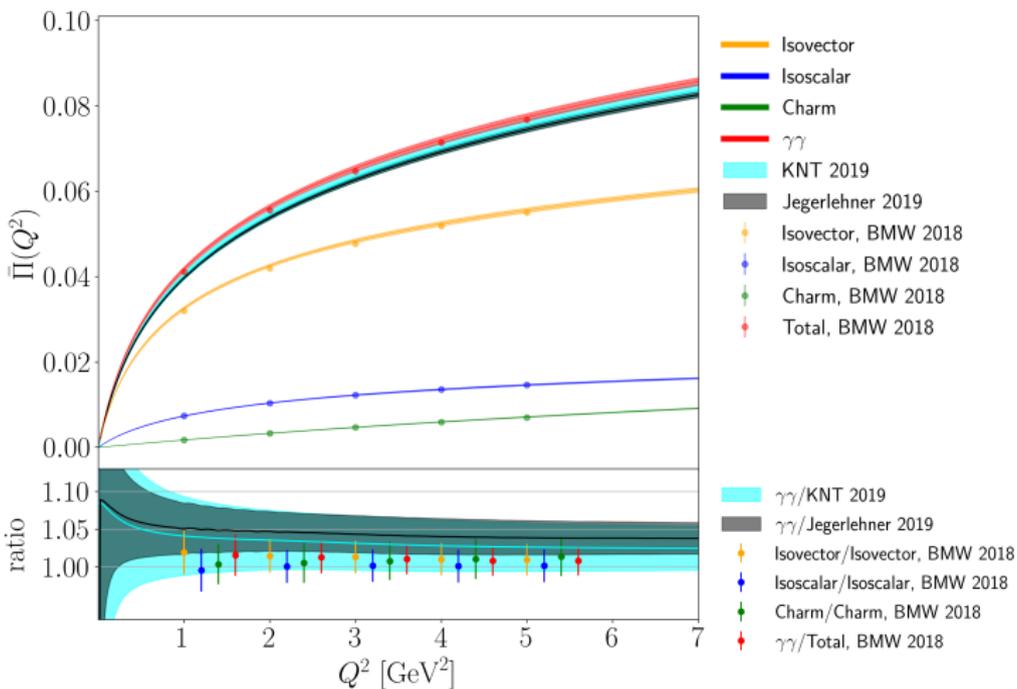
3 Extrapolation error

- Estimated excluding ensembles with $M_\pi > 400 \text{ MeV}$.
- Subleading.

4 Missing IB corrections error [A. Risch talk on Wednesday]

- Based on computation done on ensembles N200 and H102.

Results and comparison



[Keshavarzi, Nomura, and Teubner 2020; Jegerlehner 2019; Borsanyi et al. 2018]