

A seeded electroweak phase transition

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Asymptotic Safety meets Particle Physics & Friends - 18.12.2024



SB, Mariotti [2203.16450] PRL
SB, Jinno, Konstandin, Rubira, Stomberg [2302.06952] JCAP
Agrawal, SB, Mariotti, Nee [2312.06749] JHEP
SB, Mariotti [2405.08060] SciPost



Introduction



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Symmetries are **restored** at high temperatures/early times

Spontaneous breaking while the Universe expands and cools down

Introduction

⇒ Cosmological phase transitions

Key to address open questions: **baryogenesis**

Aftermath directly observable in **GWs**

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Evidence for **new** fundamental **physics**



Fig. from Schmitz [2002.04615] JHEP



Nucleation theory

 Assume thermal fluctuations in homogeneous spacetime:

$$\phi(\mathbf{x},\tau) = \phi(r), \quad r = |\mathbf{x}|$$

• Tunneling rate per unit volume given by O(3) action S_3/T

$$\gamma_V \sim T^4 \exp(-S_3/T)$$

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Coleman 1977 (PRD) Callan, Coleman 1977 (PRD) Linde 1983 (NPB)







Coleman 1977 (PRD) Callan, Coleman 1977 (PRD) Linde 1983 (NPB)

- False vacuum is homogeneous field configuration
- Critical bubble is a sphere
- Nucleation probability is the same everywhere





What about impurities?



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A supercool experiment

82K views 3 yr ago ...more





MONOPOLE AND VORTEX DISSOCIATION AND DECAY OF THE FALSE VACUUM

Paul Joseph STEINHARDT

Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138, USA

Received 17 February 1981

"If monopole (or vortex) solutions exist for a metastable or false vacuum, a finite density of monopoles (or vortices) can act as impurity sites that trigger inhomogeneous nucleation and decay of the false vacuum."

Cosmic separation of phases

Edward Witten* Institute for Advanced Study, Princeton, New Jersey 08540 (Received 9 April 1984)

"In particle physics it is often assumed that phase transitions are nucleated by thermal fluctuations. In practice, [...] except in very pure, homogeneous samples, **phase transitions are often nucleated by various forms of impurities and inhomogeneities of nonthermal origin**."

	Impurities in the early universe
	Yutaka Hosotani
Department of Physic	cs, University of Pennsylvania, Philadelphia, Pennsylvania 19 (Received 1 November 1982)
"Now one ha universe real to take place most cases th we consider o	is to ask the following question: Is the early Ily sufficiently pure in order for supercooling ? The aim of this paper is to show that in he early universe is very pure. [] In this pape ordinary particles as impurities."

"What if the transition was nucleated by impurities? In this case **the mean spacing between bubbles has nothing to do with free energies** of nucleation and is simply the spacing between the relevant impurities."



The nature of impurities

Compact objects and gravitational effects



Fig. from Oshita, Yamada, Yamaguchi [1808.01382], PLB Simone Blasi - Asymptotic & Friends

Primordial density fluctuations



Fig. from Jinno, Konstandin, Rubira, van de Vis, [2108.11947], JCAP

The nature of impurities

Topological defects

Domain walls



Fig. from Agrawal, SB, Mariotti, Nee [2312.06749] JHEP

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Fig. from SB, Mariotti, [2405.08060] SciPost

Topological classification

Defect	Dimension	Homotopy	
Domain walls	2	$\pi_0(\mathcal{M})$	
Strings	1	$\pi_1(\mathcal{M})$	





• SM + scalar singlet with $\mathbb{Z}_2: S \to -S$

$$V = -\frac{1}{2}(\mu^2 - c_h T^2)h^2 + \frac{1}{4}\lambda h^4$$
$$-\frac{1}{2}(m^2 - c_s T^2)S^2 + \frac{1}{4}\eta S^4$$
$$+\frac{1}{2}\kappa h^2 S^2$$

 $^*\mathbb{Z}_2$ breaking terms destabilize the wall network and are set to zero in the following

See e.g. Espinosa, Gripaios, Konstandin, Riva [1110.2876] JCAP

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Competition between homogenous and lacksquareseeded nucleation for 2nd step





At T_n^{seed} bubbles may nucleate on the walls



Competition between homogenous and seeded nucleation for 2nd step





At T_n^{seed} bubbles may nucleate on the walls

- Define nucleation rate per unit surface
- Stricter nucleation condition (only on sub manifold)

Lazarides, Shafi, Kibble 1982, PRD Perkins, Vilenkin 1992, PRD



Only O(2) symmetry



- "Exact" lacksquare
- Physical picture?
- Which initial conditions for the algorithm?

- Limited validity lacksquare
- Intuitive picture ${\color{black}\bullet}$
- Simple calculation lacksquare



- 2. Thin wall approximation
- 3. Kaluza-Klein decomposition
 - Quantitative results
 - Still intuitive
 - Initial conditions for num. algorithms and cross-checks

System of PDEs

False vacuum is non-trivial as it contains a domain wall (which depends on *z*)

$$\frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{\partial^2 \phi}{\partial z^2} = \frac{\partial V}{\partial \phi}, \quad \phi =$$
$$S(\infty, z) = S_{DW}(z), \quad S(r, \pm \infty) =$$

 $h(\infty, z) = h(\rho, \pm \infty) = 0$



Thin wall approximation



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R

 R^2

• Expand the fields around the domain wall background:

$$S = S_{DW}(z) + \sum_{k} s_{k} \left(x\right)$$
$$h = \sum_{k} h_{k} \left(x_{\mu}\right) \phi_{k}(z)$$

$$x_{\mu} = t, x, y$$

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 $x_{\mu} \bigg) \sigma_k(z)$



• Eigenspectrum of excitations on the wall:



Seeded tunneling as homogenous nucleation in lower dimension



Seeded tunneling as homogenous nucleation in lower dimension





Comparison



- Seeded transition is faster in all the two-step parameter space
- New parameter space becomes viable thanks to seeded tunneling





portal coupling κ

Domain wall network mimicked by Ising model Spectrum shifted to IR with enhanced amplitude \bullet





• Numerical simulation (Langevin equation)



What about other defects?

SB, Mariotti [2405.08060] SciPost

QCD axion strings T f_a Strings form at PQ phase transition Strings connected by axion domain walls QCD String—wall network collapses



QCD axion strings T f_a Strings form at PQ phase transition ??? Strings connected by axion domain walls QCD String—wall network collapses



Potential for PQ field

 $\Phi = \rho e^{i\alpha}$



 $V_{\rm PQ}(\Phi)$

Global string solution



Consider the minimal KSVZ axion model with a Higgs portal:

 $\mathcal{V} = V_{\mathrm{PQ}}(|\Phi|) + V_{\mathrm{EW}}(|\mathcal{H}|; 7$

 $V_{\rm PQ}(|\Phi|) =$

$$\begin{split} \Gamma) &+ \kappa \left(|\Phi|^2 - \frac{f_a^2}{2} \right) \left(|\mathcal{H}|^2 - \frac{v^2}{2} \right) \\ &= \eta \left(|\Phi|^2 - \frac{f_a^2}{2} \right)^2 \end{split}$$

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How do axion strings affect electroweak symmetry breaking?





String solutions

• Relevant points in field space:



• Consider first order EWPT with false vacuum B metastable at T = 0

$$V_{\rm EW}(h;T)$$



 $V_{\rm EW}(h;T) = -\frac{1}{2} \left(\mu^2 - c_h T^2\right) h^2 + \frac{\delta}{3} \frac{m_h^2}{v^2} h^3 + \frac{1}{4} \lambda h^4$

• Consider first order EWPT with false vacuum B metastable at T = 0

$$V_{\rm EW}(h;T)$$

Assume too slow hom. nucleation for simplicity



 $V_{\rm EW}(h;T) = -\frac{1}{2} \left(\mu^2 - \frac{1}{2}\right)^2$

$$(-c_h T^2) h^2 + \frac{\delta}{3} \frac{m_h^2}{v^2} h^3 + \frac{1}{4} \lambda h^4$$



FOPT + PQ $\bullet B \rightarrow \bullet A$ $= E_{\rm hom}$

Critical bubble











$(\kappa/\eta \sim \text{crit.})$



 κ/η

Profile of the critical bubble: ★



 $m_h r$



- Percolation as interplay between seeded nucleation rate and density of defects
- Axion—seeded EWPT effectively $\beta/H \sim \xi \sim 10$
- Different velocities parallel or orthogonal to the string?
- Gravitational wave emission before collision (non-spherical bubbles)?



Summary

- The presence of impurities in the early Universe can strongly affect the way a phase transition proceeds
- The xSM with Z_2 symmetry is arguably the simplest (and complete) example for a seeded EWPT
- Other defects can exist at the time of the EWPT: dedicated study of QCD axion strings in KSVZ model with Higgs portal
- Pheno aspects of seeded phase transitions: percolation, slow transitions, expansion of non-spherical bubbles, features in the GW signal?

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- The presence of impurities in the early Universe can strongly affect the way a phase transition proceeds
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Thank you!







Seeded EWPT at LISA



Peak frequency [Hz]