

Exploring the conformal window and infrared fixed points using lattice field theory

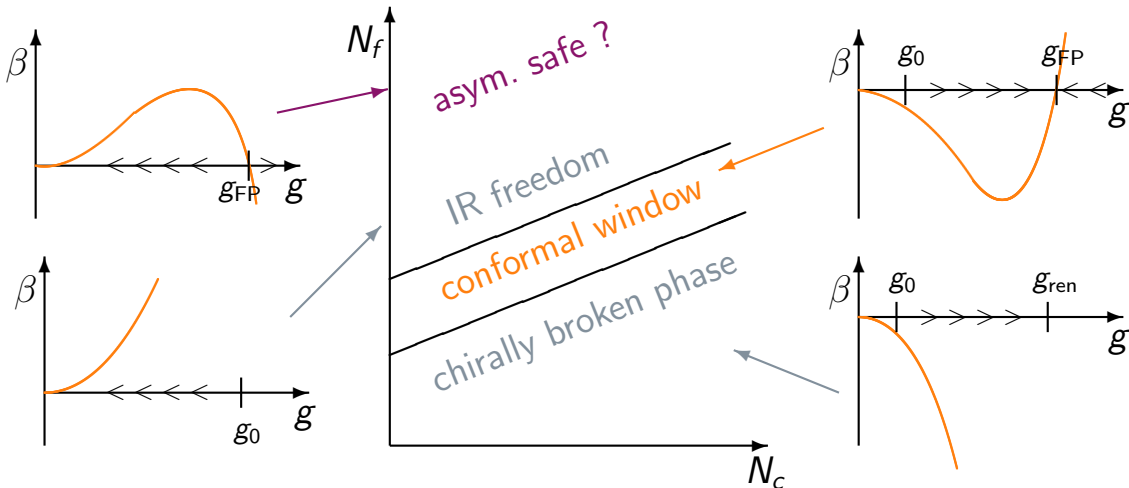
Oliver Witzel



Asymptotic Safety meets Particle Physics & Friends
DESY Hamburg, Germany · December 18, 2023

Gauge-fermion systems

- ▶ Gauge-fermion system with $N_c \geq 2$ colors and N_f flavors in some (fundamental) representation



Renormalization Group β function

$$\beta(g^2) = \mu^2 \frac{dg^2}{d\mu^2}$$

- ▶ Encodes dependence of coupling g^2 on the energy scale μ^2
- ▶ β has no explicit dependence on μ^2 , only implicit through $g^2(\mu)$
- ▶ Known perturbatively up to 5-loop order in the $\overline{\text{MS}}$ scheme (1- and 2-loop are universal)
[Baikov, Chetyrkin, Kühn PRL118(2017)082002] [Ryttov and Shrock PRD94(2016)105015]
- ▶ Known perturbatively at 3-loop order in the gradient flow scheme [Harlander, Neumann JHEP06(2016)161]
- ▶ Perturbative predictions reliable at weak coupling,
nonperturbative methods needed for strong coupling

Lattice calculations

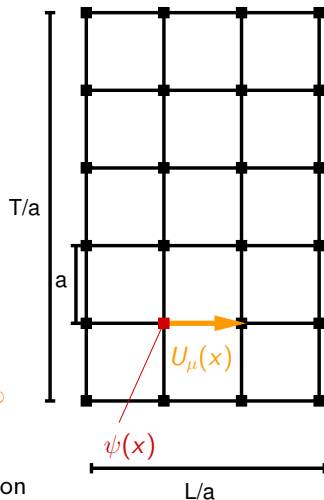
- ▶ Wick-rotate to Euclidean time $t \rightarrow i\tau$
- ▶ Discretize space-time and set up a hypercube of finite extent $(L/a)^3 \times T/a$ and spacing a

- ▶ Use path integral formalism

$$\langle \mathcal{O} \rangle_E = \frac{1}{Z} \int \mathcal{D}[\psi, \bar{\psi}] \mathcal{D}[U] \mathcal{O}[\psi, \bar{\psi}, U] e^{-S_E[\psi, \bar{\psi}, U]}$$

⇒ Large but finite dimensional path integral

- ▶ Finite lattice spacing $a \rightarrow$ UV regulator
- ▶ Finite volume of length $L \rightarrow$ IR regulator
 - Study physics in a finite box of volume $(aL)^3$ plus limit $L \rightarrow \infty$
- ▶ Different discretizations for gauge and fermion actions possible
 - Wilson, Symanzik gauge; Wilson, staggered, domain-wall fermion
 - Discretization effects disappear after taking $a \rightarrow 0$ continuum limit



Gradient flow

[Narayanan and Neuberger JHEP 0603 (2006) 064] [Lüscher CMP 293 (2010) 899][JHEP 1008 (2010) 071]

- ▶ Add flow time coordinate t with dimension [-2] and define gauge field $B_\mu(x, t)$

$$\frac{\partial}{\partial t} B_\mu(t) = \mathcal{D}_\nu(t) G_{\nu\mu}(t) \quad \text{with} \quad B_\mu(t=0) = A_\mu$$

- Ordinary differential equation (ODE) \rightsquigarrow solve numerically using Runge-Kutta
- Covariant derivative defined in terms of the flow field $B_\mu(t)$ and the Yang-Mills action $S_{YM}(B)$
- ▶ Gradient flow is a smoothing/averaging transformation
 - Can be related to RG flow \rightsquigarrow perfect to define RG β function, ...
- ▶ Energy density $\langle E(t) \rangle = -\frac{1}{2} \text{Tr} \langle G_{\mu\nu}(x, t) G_{\mu\nu}(x, t) \rangle$
 - Renormalized coupling, scale setting ($\sqrt{8t_0}$, w_0)
- ▶ Different gradient flows: Wilson flow, Symanzik flow, Zeuthen flow, ...

Step-Scaling β function

- ▶ Discretized β function determined using numerical lattice field theory calculations [Lüscher et al. NPB359(1991)221]
 - Choose symmetric L^4 setup where the size L of the lattice is the **only** scale
 - Determine β function by calculating scale change $L \rightarrow s \cdot L$
 - Conventionally differs by a minus sign compared to the continuum definition

- ▶ Use gradient flow to define a renormalized coupling

$$g_c^2(L) = \frac{128\pi^2}{3(N_c^2 - 1)} \frac{1}{C(c, L)} t^2 \langle E(t) \rangle$$

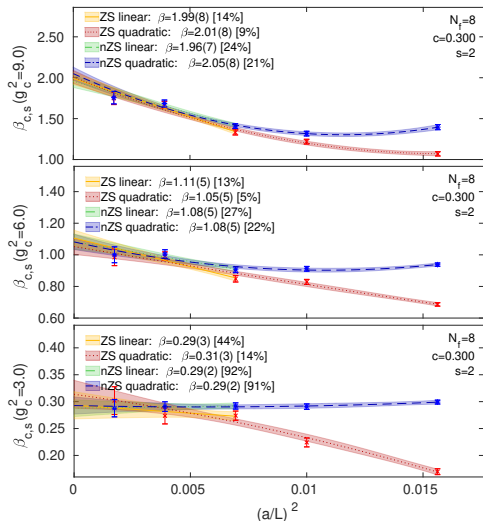
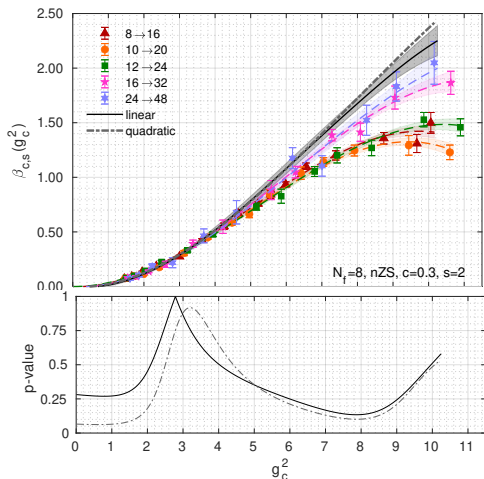
- Relate flow time t to scale L : $\sqrt{8t} = c \cdot L$ [Fodor et al. JHEP11(2012)007][JHEP09(2014)018]

- Calculate scale difference

$$\beta_s^c(g_c^2; L) = \frac{g_c^2(sL) - g_c^2(L)}{\log(s^2)}$$

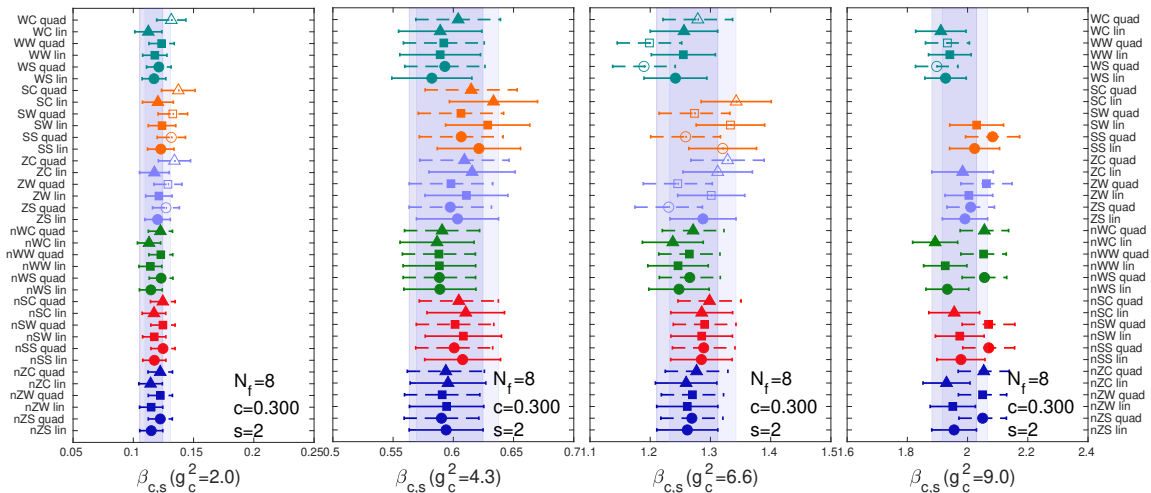
- Extrapolate $L/a \rightarrow \infty$ to remove discretization effects and take combined continuum and infinite volume limit

Analysis SU(3) with $N_f = 8$ fundamental flavors [Hasenfratz, Rebbi, OW PRD 107(2023)114508]

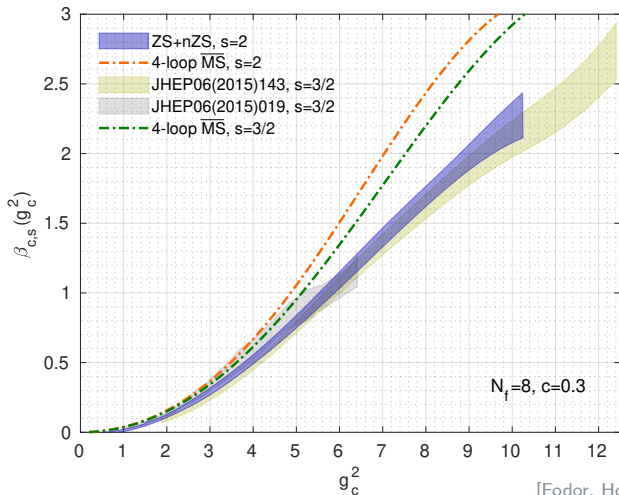


Systematic effects SU(3) with $N_f = 8$ fundamental flavors

[Hasenfratz, Rebbi, OW PRD 107(2023)114508]



Comparison SU(3) with $N_f = 8$ fundamental flavors



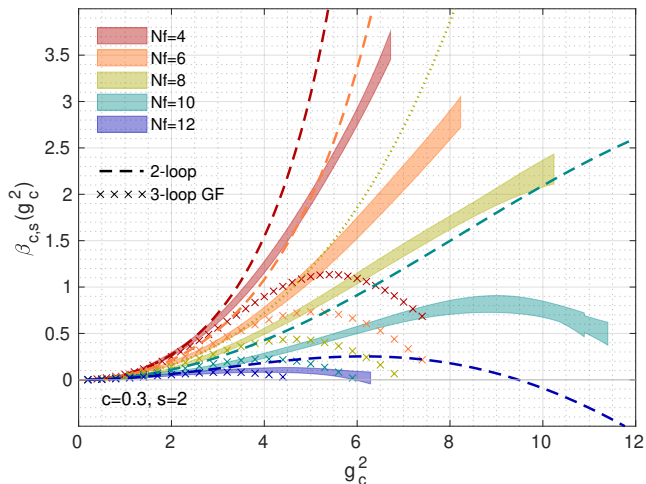
[Hasenfratz, Rebbi, OW PRD 107(2023)114508]

- ▶ Small effect due to different scale change s
- ▶ GF $c = 0.3$ scheme \neq 4-loop $\overline{\text{MS}}$

[Hasenfratz, Schaich, Veernala JHEP06(2015)143]

[Fodor, Holland, Kuti, Mondal, Nogradi, Wong JHEP06(2015)019]

Step-Scaling β function for $N_f = 4, 6, 8, 10, 12$



▶ 3-loop GF

[Harlander, Neumann JHEP06(2016)161]

▶ $N_f = 12$: indication of an IRFP at $g_c^2 \approx 6.3$

▶ $N_f = 10$: hints for an IRFP at $g_c^2 \gtrsim 13$

[Hasenfratz, Rebbi, OW PLB 798(2019)134937]

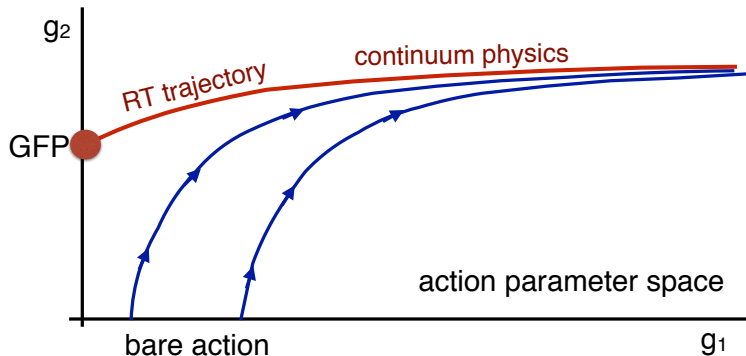
[Hasenfratz, Rebbi, OW PRD 100(2019)114508]

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[Hasenfratz, Rebbi, OW PRD 106(2022)114509]

[Hasenfratz, Rebbi, OW PRD 107(2023)114508]

Gradient flow and real-space renormalization Group (RG) flow



Gradient flow and real-space renormalization Group (RG) flow

- ▶ RG flow: change of (bare) parameters and coarse graining (blocking)
- ▶ Gradient flow is a continuous transformation
 - Define real-space RG blocked quantities by incorporating coarse graining as part of calculating expectation values [Carosso, Hasenfratz, Neil PRL 121 (2018) 201601]
- ▶ Relate GF time t/a^2 to RG scale change $b \propto \sqrt{t/a^2}$
 - Quantities at flow time t/a^2 describe physical quantities at energy scale $\mu \propto 1/\sqrt{t}$
 - Local operator with non-vanishing expectation value can be used to define running coupling
 - ↪ Simplest choice: $t^2 \langle E(t) \rangle$ [Lüscher JHEP 1008 (2010) 071]
- ▶ Continuous RG β function

$$\beta_{GF}(g_{GF}^2) = \mu^2 \frac{dg_{GF}^2}{d\mu^2} = -t \frac{dg_{GF}^2}{dt}$$

Continuous RG β function

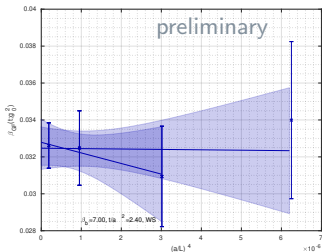
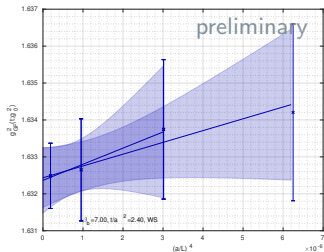
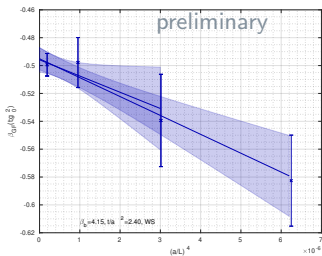
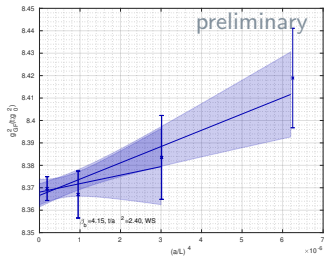
[Fodor et al. EPJ Web Conf. 175 (2018) 08027]

[Hasenfratz, OW PRD 101 (2020) 034514] [Hasenfratz, OW PoS LATTICE2019 (2019) 094]

[Wong et al. PoS LATTICE2022 (2023) 043] [Hasenfratz, Peterson, Van Sickle, OW PRD108 (2023) 014502]

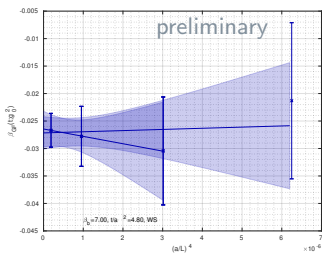
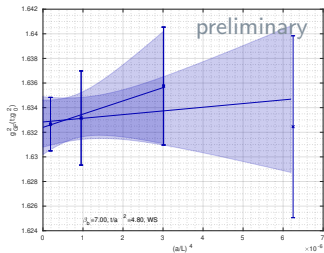
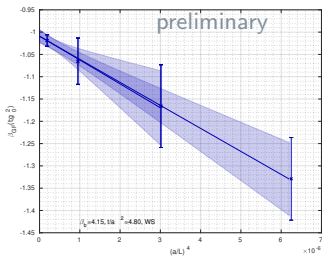
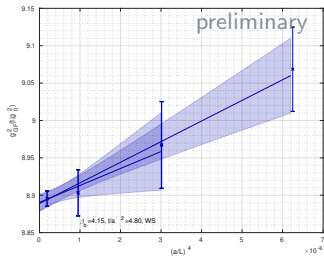
- ▶ Extract $g_{GF}^2(t; \beta_b, L/a)$ its derivative $\beta_{GF}(t; \beta_b, L/a)$ for a range of GF times on each ensembles
→ Different bare coupling β_b on different volumes $(L/a)^4$ or $(L/a)^3 \times T/a$
- ▶ Perform infinite volume extrapolation at fixed bare coupling β_b and GF time t
→ Obtain $g_{GF}^2(t; \beta_b)$ and $\beta_{GF}(t; \beta_b)$
- ▶ Interpolate discrete infinite volume values to get continuous values at fixed flow time
→ $g_{GF}^2(t)$ and $\beta_{GF}(t; g_{GF}^2)$
- ▶ Take continuum limit ($a^2/t \rightarrow 0$) for fixed g_{GF}^2 and obtain $\beta_{GF}(g_{GF}^2)$

⇒ Reanalyze data of the step-scaling calculations

$N_f = 8$: Infinite volume extrapolation

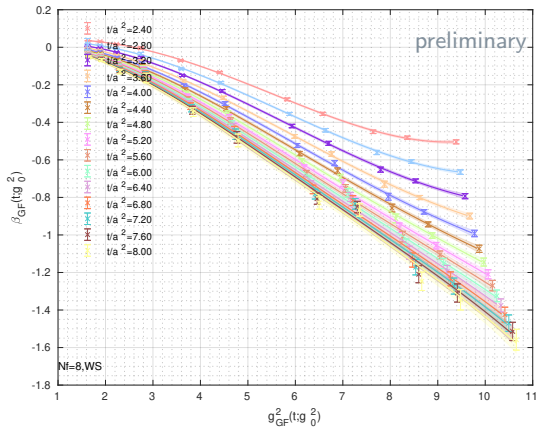
- ▶ For each bare coupling β_b and at all GF times
- ▶ Relevant GF times typically $t/a^2 \in [2.20, 5.00]$
- ▶ Strong ($\beta_b = 4.15$) and weak ($\beta_b = 7.00$) coupling at $t/a^2 = 2.4$

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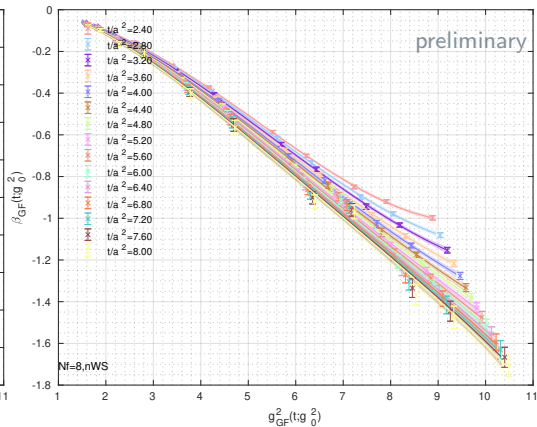


- ▶ For each bare coupling β_b and at all GF times
- ▶ Relevant GF times typically $t/a^2 \in [2.20, 5.00]$
- ▶ Strong ($\beta_b = 4.15$) and weak ($\beta_b = 7.00$) coupling at $t/a^2 = 4.8$

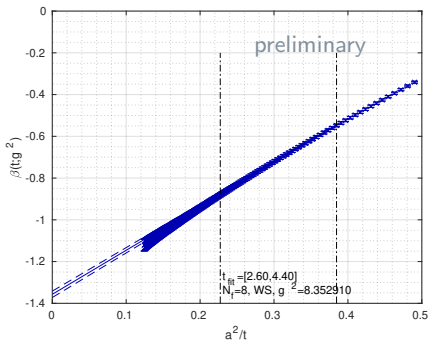
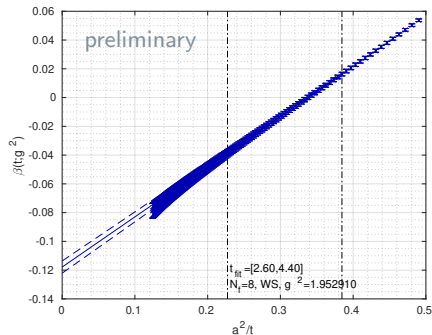
$N_f = 8$: Polynomial interpolation in g_{GF}^2



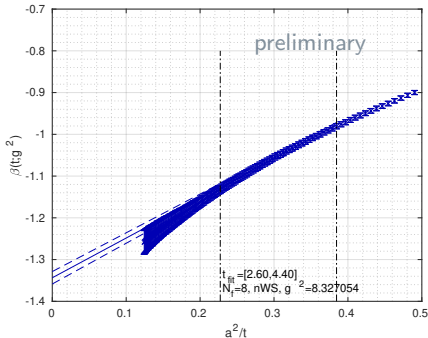
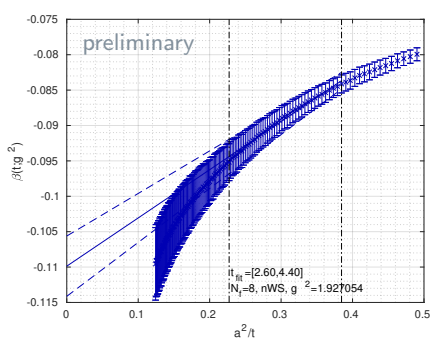
► Without tree-level improvement



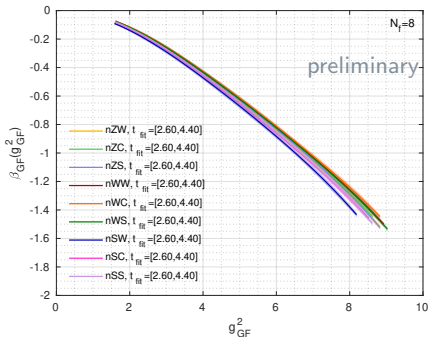
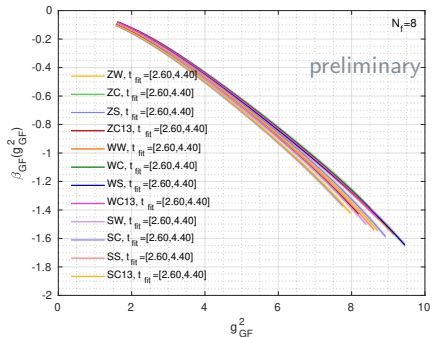
► With tree-level improvement

$N_f = 8$: Continuum limit ($a^2/t \rightarrow 0$)

- ▶ Simple continuum limit for range of flow times
- ▶ Without tree-level improvement

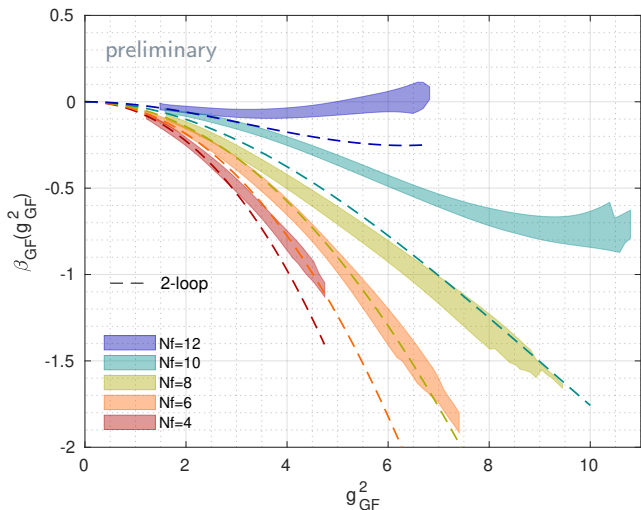
$N_f = 8$: Continuum limit ($a^2/t \rightarrow 0$)

- ▶ Simple continuum limit for range of flow times
- ▶ With tree-level improvement

$N_f = 8$: Continuous β function

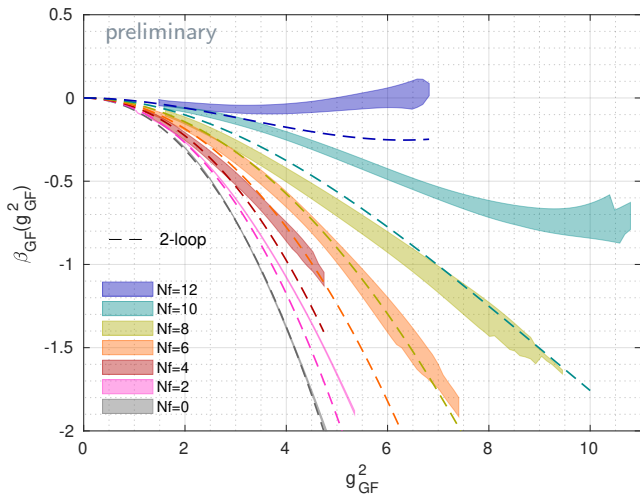
- ▶ Different flow-operator combinations with and without tree-level improvement

Summary continuous β function for $N_f = 4, 6, 8, 10, 12$



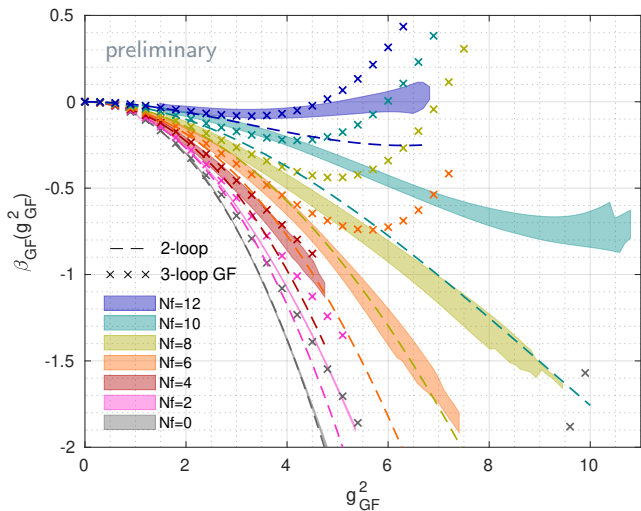
- ▶ Systematic effects for $N_f = 10$ likely underestimated
- ▶ Reach in g^2 limited by 1st order bulk phase transition (lattice artifact)
- ▶ Qualitative behavior captured by 2-loop PT prediction for $N_f \leq 8$
- ▶ IRFP for $N_f = 12$ not (well) resolved

Summary continuous β function for $N_f = 0, 2, 4, 6, 8, 10, 12$

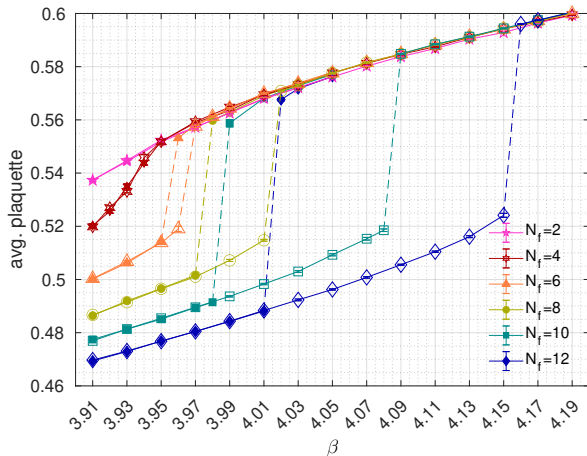


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- ▶ Including $N_f = 2$ and $N_f = 0$
 [Hasenfratz, OW PRD 101(2020)034514]
 [Hasenfratz, Peterson, Van Sickle, OW PRD108(2023)014502]

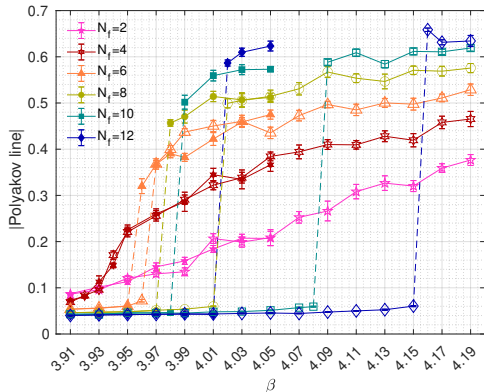
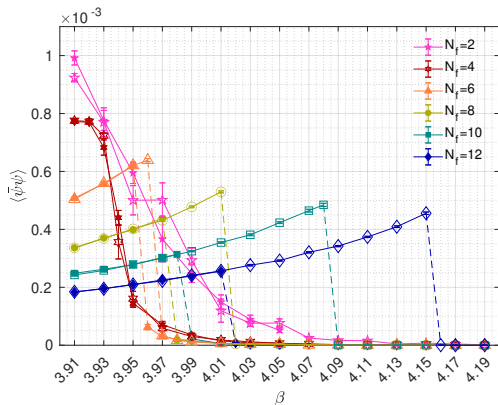
Summary continuous β function for $N_f = 0, 2, 4, 6, 8, 10, 12$



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 [Hasenfratz, Peterson, Van Sickle, OW PRD108(2023)014502]
- ▶ 3-loop GF prediction tracks nonperturbative result longer, but then turns away showing different qualitative behavior

Reach in g_{GF}^2 [Hasenfratz, Rebbi, OW PRD 107(2023)114508]

- ▶ $N_f = 2, 4, 6$: zero mass simulations limited by confinement transition
 - Perform finite mass simulations plus chiral extrapolation
- ▶ $N_f = 8, 10, 12$: 1st order bulk phase transition limits reach in g^2
 - Lattice artifact due to choice of actions ($3\times$ stout-smear MDWF+Symanzik)
 - Wide hysteresis
 - ↪ Artifacts may affect strongest coupling
 - Even larger volumes will not allow to overcome these issues
 - ↪ add e.g. Pauli-Villars field

Reach in g_{GF}^2 [Hasenfratz, Rebbi, OW PRD 107(2023)114508]

Adding extra Pauli-Villars fields

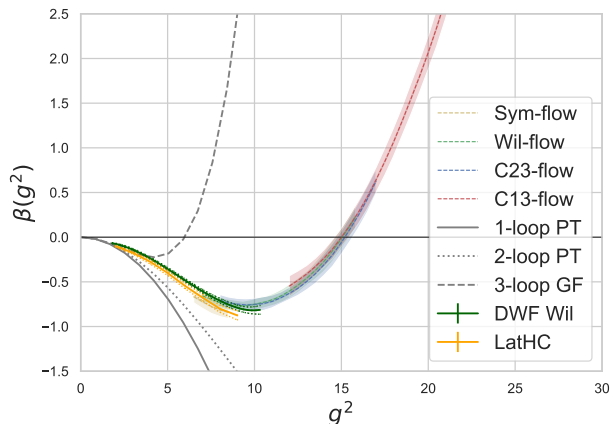
- ▶ Pauli-Villars (PV) are introduced when constructing 5-dim. domain-wall fermions (DWF)
 - Chiral fermions obeying the Ginsparg-Wilson relation
bypassing the Nielsen-Niyomiya no-go theorem
- ▶ Bosonic fields with mass $am_{PV} = 1$
- ▶ Automatically integrated out, when taking the continuum limit

- ▶ Side-effect: PV fields largely absorb effective shift in bare coupling due to (many) fermions
 - At weak coupling avg. plaquette does essentially not vary with N_f for DWF

- ▶ Idea: add PV fields also to Wilson or Staggered fermions for studying many flavor systems
at strong coupling [Hasenfratz, Shamir, Svetitsky PRD 104 (2021) 074509]

$N_f = 10$: Wilson fermion with Pauli-Villars fields

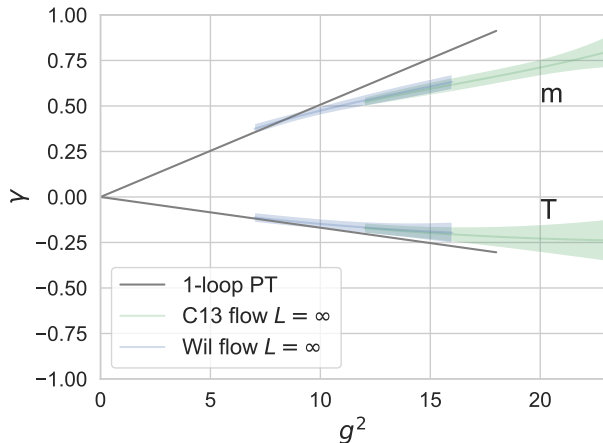
[Hasenfratz, Neil, Shamir, Svetitsky, OW PRD 108 (2023) L071503]



- ▶ Infrared fixed-point around $g^2 \approx 15$
 ⇒ SU(3) with $N_f = 10$ is conformal
- ▶ Consistent result for different gradient flows
- ▶ Weak coupling overlaps with our DWF and LatHC's staggered result
 [Fodor et al. PoS Lattice2018 (2018) 199]
 [Fodor et al. PoS Lattice2019 (2019) 121]
 [Kuti et al. PoS Lattice2021 (2022) 321]
- ▶ Reaching couplings $g^2 > 20$

$N_f = 10$: Anomalous dimension

[Hasenfratz, Neil, Shamir, Svetitsky, OW PRD 108 (2023) L071503]



- ▶ Calculate 2-pt function of flowed mesonic source X' with unflowed source X
 $\langle X(0)X'(t) \rangle \sim t^{-(d+\eta+\gamma)/2}$
 $\rightarrow \gamma$ anomalous dimension of operator
 $\rightarrow \eta$ anomalous dimension of fermion field

- ▶ Cancel η by defining ratio over conserved vector current

$$R(t) = \frac{X(0)X'(t)}{V(0)V'(t)} \sim t^{-\gamma/2}$$

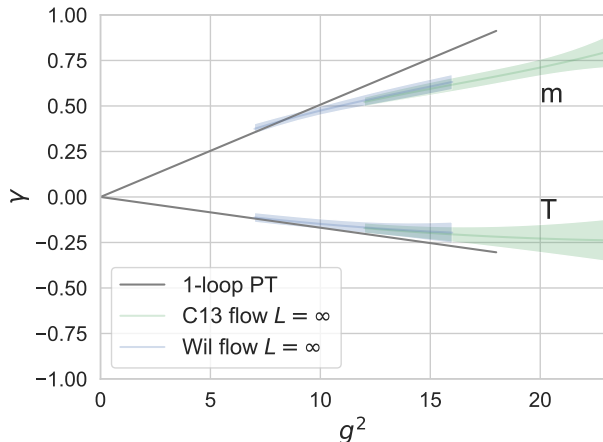
- ▶ Logarithmic derivative

$$\gamma = -2 \frac{t}{R} \frac{\partial R}{\partial t}$$

with $t \gg x_4$ (Euclidean time separation)

$N_f = 10$: Anomalous dimension

[Hasenfratz, Neil, Shamir, Svetitsky, OW PRD 108 (2023) L071503]



▶ PT at 1-loop

→ Mass anomalous dimension (scalar)

$$\gamma_m = \frac{6g^2 C_2}{16\pi^2}$$

$C_2 = \frac{4}{3}$ quadratic Casimir operator

→ Anomalous dimension of the tensor density

$$\gamma_T = -\frac{1}{3}\gamma_m$$

▶ At the IRFP $g^2 \approx 15$

$$\gamma_m \simeq 0.6$$

→ Consistent with [LSD PRD 103 (2021) 014504]

$$\gamma_m = 0.47(5) \text{ at } g^2 \sim 10$$

$$\gamma_T \simeq -0.2$$

Outlook

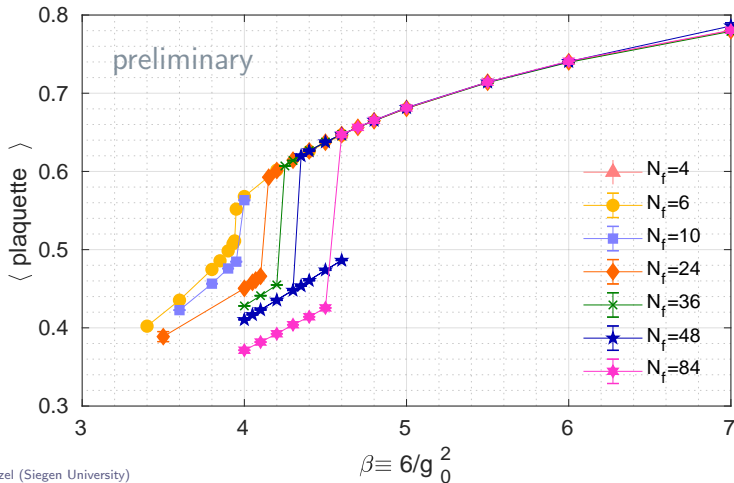
- ▶ SU(3) with $N_f = 12$ is conformal, IRFP around $g^2 \sim 6$
- ▶ SU(3) with $N_f = 10$ is conformal, IRFP around $g^2 \sim 15$

- ▶ What is the nature of $N_f = 8$?
 - Is $N_f = 8$ the onset of the conformal window?
 - Does $N_f = 8$ exhibit a new phase (symmetric mass generation)? [Hasenfratz PRD 106 (2022) 014513]
 - Ongoing and new investigations by the LSD collaboration incl. simulations with PV fields

- ▶ What happens for larger N_f ? Asymptotic safety?

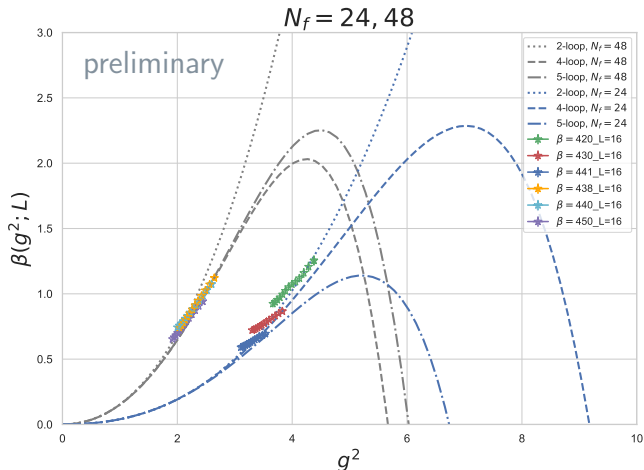
SU(3) with many (more) flavors

- ▶ Exploring bare parameter space for SU(3) with many fundamental flavors (stout-smear MDWF and Symanzik gauge action)



- ▶ Seeking 1st order transition e.g. jump in average plaquette
- ▶ Cheap 8^4 lattices
- ▶ Transition moves to weaker for more flavors
- Verification of bulk phase transition vs. finite T to be completed

SU(3) with many (more) flavors



- ▶ Only $L/a = 16$
- ▶ β function runs backward!
- ▶ Run into bulk phase transition before observing a sign of an UVFP
- ▶ Add four-fermion interaction
→ N-JL type models
- ▶ Repeat with a better action

extra

DWF setup

- ▶ Symanzik gauge action
- ▶ Möbius domain-wall fermions with three levels of stout smearing ($\rho = 0.1$)
- ▶ Input quark mass $am_q = 0$, $L_s = 12$ or 16 such that $am_{res} < 10^{-5}$
- ▶ Fermions with anti-periodic boundary conditions in space and time

$N_f = 4$

7–9 bare couplings

$\beta = 8.50 - 4.60(4.20)$

$N_f = 6$

7–11 bare couplings

$\beta = 8.50 - 4.30(4.05)$

$N_f = 8$

11–16 bare couplings

$\beta = 7.00 - 4.10(4.02)$

- ▶ 5 volume pairs with $s = 2$

$40^4, 32^4, 24^4, 20^4, 16^4, 12^4, 10^4, 8^4$

$48^4, 32^4, 24^4, 20^4, 16^4, 12^4, 10^4, 8^4$

- ▶ Simulations performed using Grid [Boyle et al. PoS Lattice2015 023]
- ▶ Measuring Zeuthen flow, Symanzik flow, and Wilson flow in Qlua [Pochinsky PoS Lattice2008 040]
- ▶ Apply tree-level normalization to reduce cutoff effects [Fodor et al. JHEP09(2014)018]

DWF setup

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$$N_f = 10$$

$$N_f = 12$$

17 bare couplings

17 bare couplings

$$\beta = 7.00 - 4.02$$

$$\beta = 7.00 - 4.10$$

- ▶ 5 volume pairs with $s = 2$

$$32^4, 28^4, 24^4, 20^4, 16^4, 14^4, 12^4, 10^4, 8^4$$

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Wilson setup

- ▶ Wilson gauge action with nHYP dislocation suppression (NDS)
- ▶ nHYP-smearred Wilson-clover fermions with $(\alpha_1, \alpha_2, \alpha_3) = (0.75, 0.6, 0.3)$
 - Mass at zero (κ tuned to κ_{critical})
 - Fermions with periodic (anti-periodic) boundary conditions in space (time)
- ▶ 30 additional Pauli-Villars fields with $am_{PV} = 1$
- ▶ 6 bare couplings $\beta_b = 7.0, 6.7, 6.5, 6.3, 6.2, 6.0$
- ▶ $24^3 \times 48$ and $28^3 \times 56$ volumes
- ▶ Measuring Symanzik flow, Wilson, C13 and C23 flows