

# **Absence of $SO(4)$ quantum criticality in Dirac semimetals at two-loop order**

**based on MU, Herbut, Stamou, Scherer – arXiv:2308.12464 & PRB 108, 245130**

**Max Uetrecht**

Asymptotic Safety @ DESY

19 December, 2023

# Outline

- **Critical Phenomena in Dirac Materials**

- Equivalent field-theoretical description of Criticality

*[Liu, Huffmann, Chandrasekharan, Kaul '22, PRL] [Herbut, Scherer '22, PRB]*

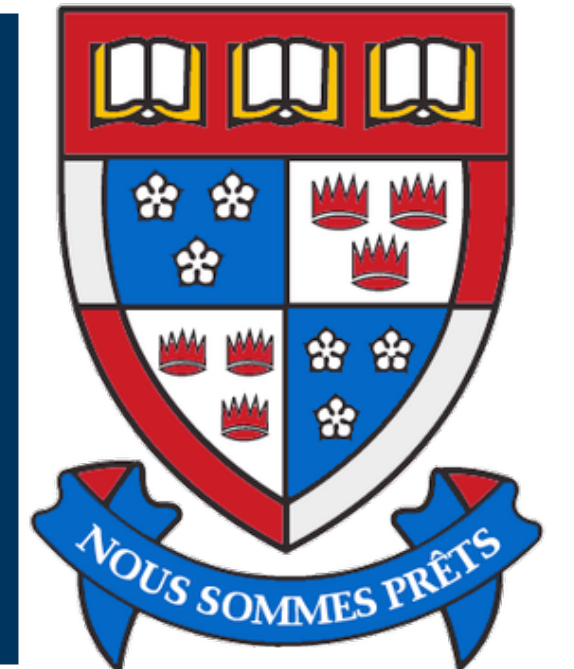
- Recent quantum Monte Carlo simulation  $\longrightarrow$  **Criticality found**

- **Describe QMC Criticality with a Gross–Neveu–Yukawa field theory**

- RG Fixed-Point Analysis at **one-loop** order  $\longrightarrow$  **Criticality lost**  $\longrightarrow$  **two-loop?**

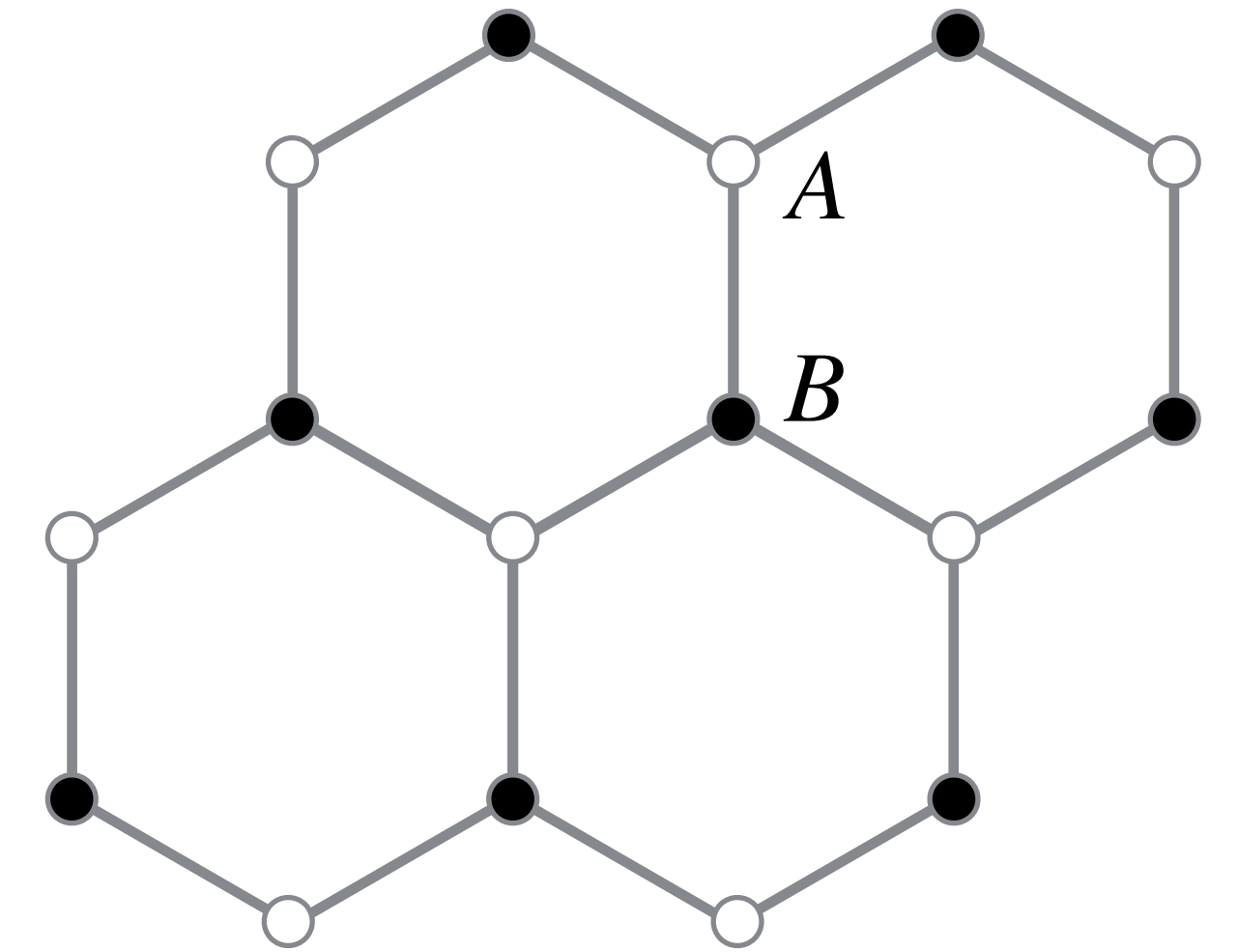
*[Herbut, Scherer '22, PRB]*

**Collaboration with:** Michael M. Scherer (RUB), Igor F. Herbut (SFU), Emmanuel Stamou (TUDO)



# Critical Phenomena in Dirac Materials

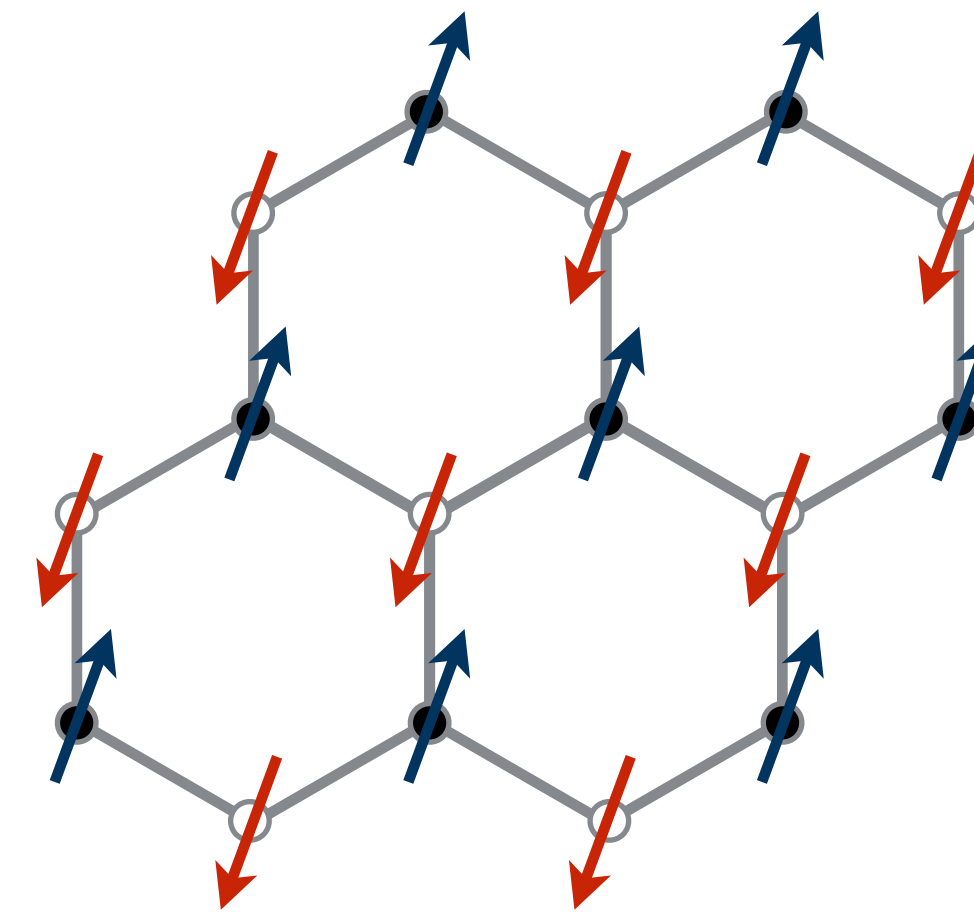
- Electrons in a solid effectively described by the **Dirac equation**
- Examples: graphene, topological insulators, ...
- Strong electron-electron interaction  $\longrightarrow$  massive fermion phase



$$m = 0$$

$$m > 0$$

e.g.:



Unordered  
phase

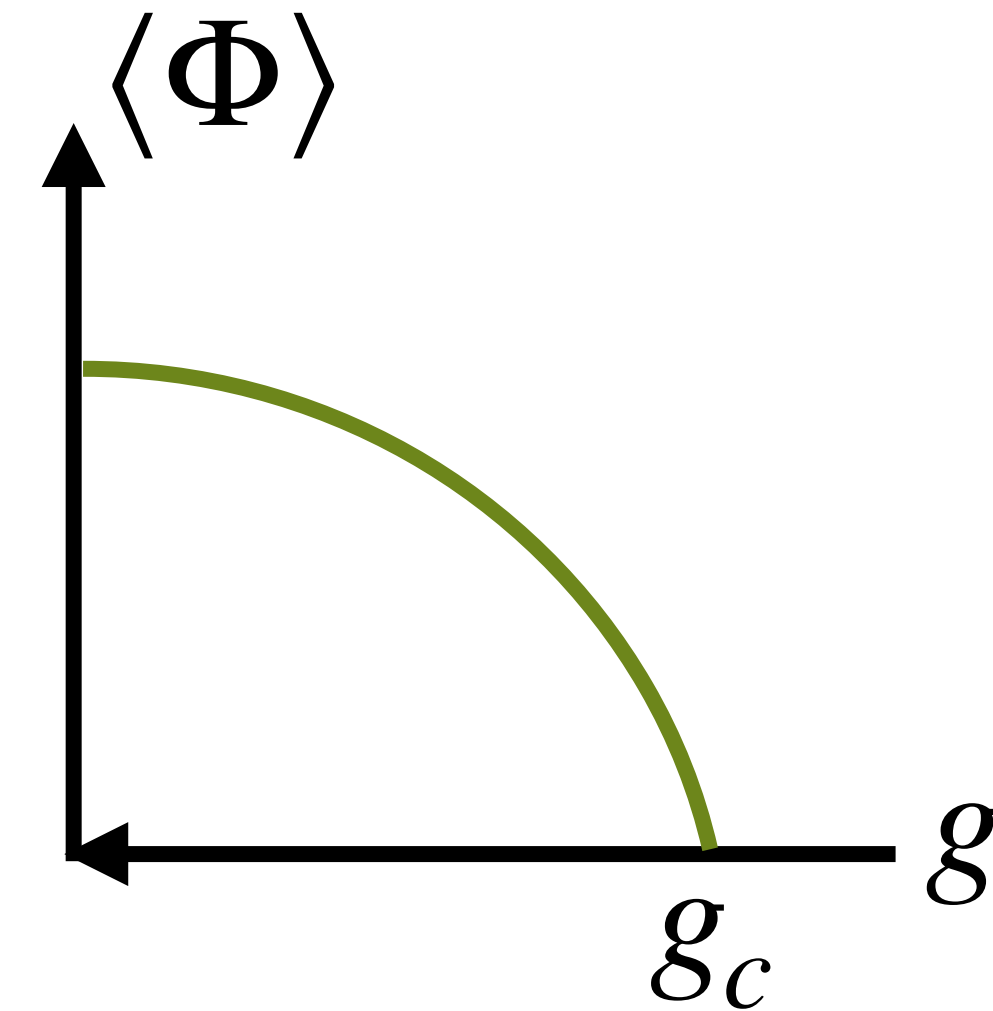
$g_c$

Ordered  
phase

Tuning parameter  
(e.g. coupling  $g$ )

# Critical Phenomena in Dirac Materials

- Phases connected by continuous phase transition  $\longrightarrow$  Critical Phenomena
- Mass generation described by **Spontaneous Symmetry Breaking (SSB)**
- Governed by Order Parameter  $\Phi$  associated to the symmetry



$$\langle \Phi \rangle = 0$$

$$\langle \Phi \rangle \neq 0$$

Unordered  
phase

$g_c$

Ordered  
phase

Tuning parameter  
(e.g. coupling  $g$ )

# Critical Phenomena in Dirac Materials

- At critical point  $g = g_c$  : system is **scale invariant**, i.e.  $\xi \rightarrow \infty$
- **Connection to field theory?**

Renormalize  $\mathcal{L}_{\text{Dirac}}$   $\longrightarrow$  Beta functions  $\beta(g) = \frac{dg}{d \log \mu}$

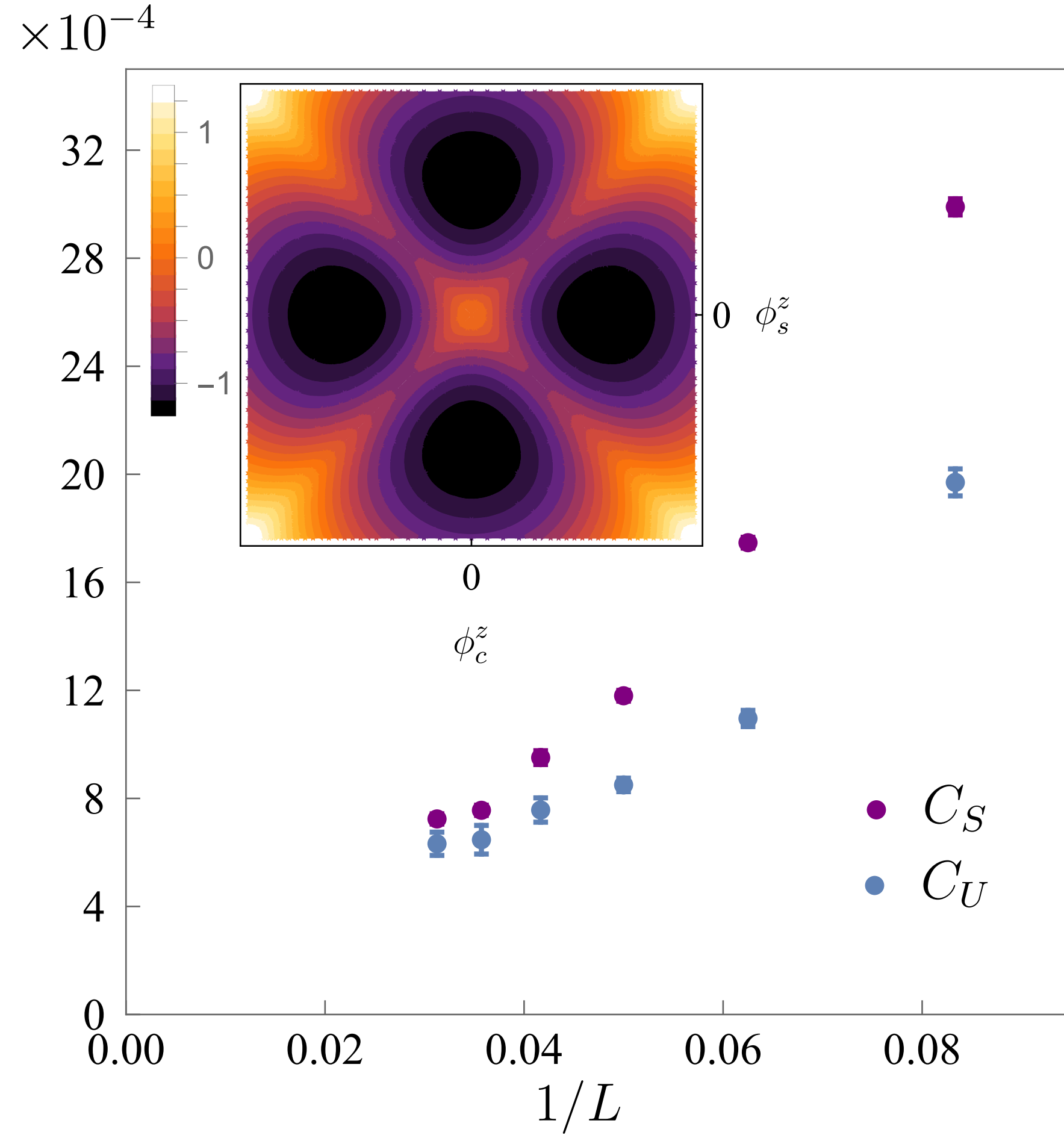
Roots of Beta function = Fixed Points (FPs)  $\longrightarrow$  **scale invariance**

$$\beta(g) = \frac{dg}{d \log \mu} = 0 \implies g = g_c$$

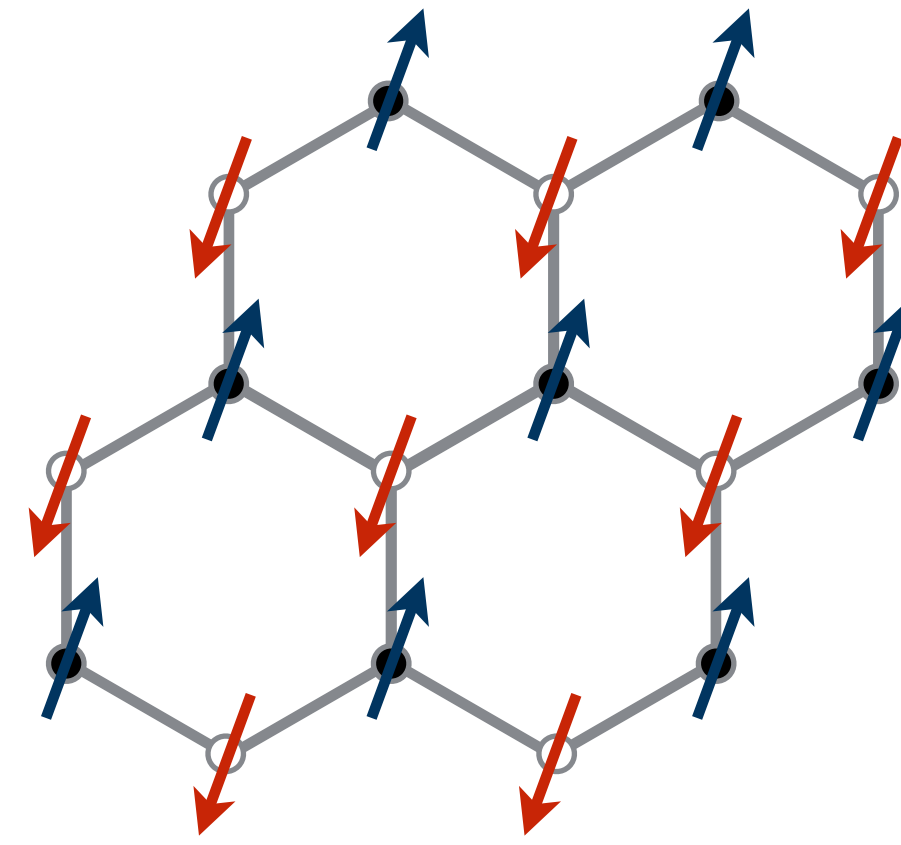
 **FPs of field theory correspond to critical points**

# Critical Phenomena in Dirac Materials: Recent QMC Analysis

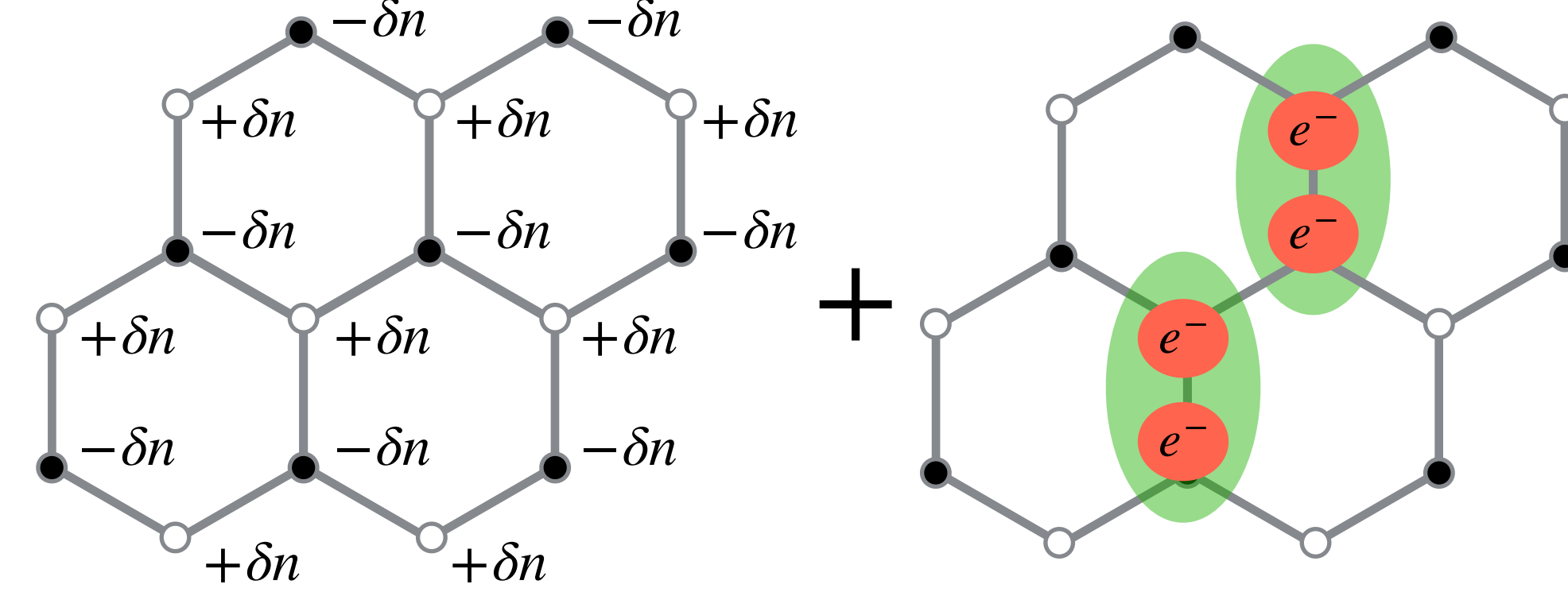
[Liu, Huffmann, Chandrasekharan, Kaul '22, PRL]



- *H. Liu et al. '22*: Quantum Monte Carlo (QMC) analysis of Dirac criticality with **two order parameters (OPs)**



$C_S$  : Néel state

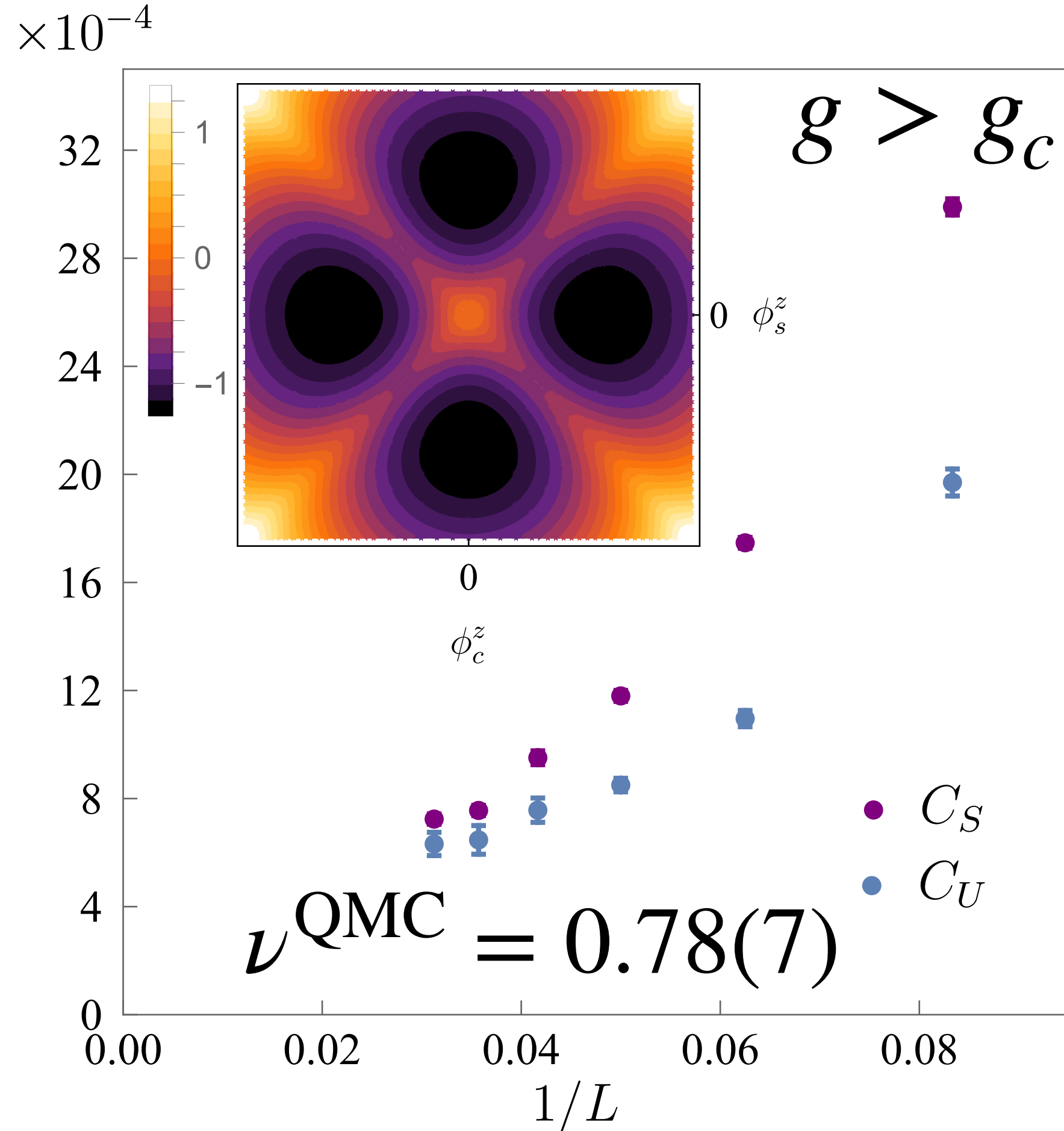


$C_U$  : superconductor-CDW-state



# Critical Phenomena in Dirac Materials: Recent QMC Analysis

[Liu, Huffmann, Chandrasekharan, Kaul '22, PRL]



- QMC data: **continuous transition** between massless Dirac phase and massive phase
- Both OPs **simultaneously** break their associated symmetry:  $L \rightarrow 0 : C_{S/U} \neq 0!$
- Divergence:  $g \rightarrow g_c : \xi \propto |g - g_c|^{-\nu}$

➔ **Capture criticality with field theory**



**Next:**

**Describe quantum Monte Carlo Criticality with a  
Field Theory**



# Gross–Neveu–Yukawa (GNY) Theory in $d = 2 + 1$

$$\mathcal{L} \simeq \mathcal{L}_{\text{free}} - g_a \bar{\Psi} (\vec{a} \vec{\sigma}) \Psi - g_b \bar{\Psi} (\vec{b} \vec{\sigma}) \Psi - \lambda_a \vec{a}^2 - \lambda_b \vec{b}^2 - \lambda_{ab} \vec{a} \vec{b}$$

- Represent two OPs as **massive scalar fields**  $\vec{a}, \vec{b}$ : vectors of  $\text{SO}(3) \simeq \text{SU}(2)$
- Combine OPs into **global**  $\text{SO}(4) \simeq \text{SU}(2)_A \times \text{SU}(2)_B$  **symmetry; same as QMC lattice**
- Couple OPs to  $\text{SU}(2)$ –bidoublet fermion  $\longrightarrow$  **fermionic “mass terms”**
- Include fluctuations of the OPs  $\longrightarrow$  **quartics**

	$\text{SU}(2)_A$	$\text{SU}(2)_B$
$\Psi$	<b>2</b>	<b>2</b>
$a_i$	<b>3</b>	—
$b_j$	—	<b>3</b>

[Liu, Huffmann, Chandrasekharan, Kaul '22, PRL]

[Herbut, Scherer '22, PRB]

# Gross–Neveu–Yukawa (GNY) Theory in $d = 4 - \epsilon$

$$\mathcal{L} \simeq \mathcal{L}_{\text{free}} - g_a \bar{\Psi} (\vec{a} \vec{\sigma}) \Psi - g_b \bar{\Psi} (\vec{b} \vec{\sigma}) \Psi - \lambda_a \vec{a}^2 - \lambda_b \vec{b}^2 - \lambda_{ab} \vec{a} \vec{b}$$

- $d = 2 + 1$  : couplings are dimensionful  $\longrightarrow$  **no direct perturbative expansion**
- Continue  $\mathcal{L}$  analytically to  $d = 4 - \epsilon$  and expand for **small  $\epsilon$**
- Obtain predictions in the limit  $\epsilon \rightarrow 1$

## RGEs

 **Beta functions & anomalous dimensions at two-loops in  $d = 4 - \epsilon$**

# Gross–Neveu–Yukawa (GNY) Theory in $d = 4 - \epsilon$

$$\mathcal{L} \simeq \mathcal{L}_{\text{free}} - g_a \bar{\Psi} (\vec{a} \vec{\sigma}) \Psi - g_b \bar{\Psi} (\vec{b} \vec{\sigma}) \Psi - \lambda_a \vec{a}^2 - \lambda_b \vec{b}^2 - \lambda_{ab} \vec{a} \vec{b}$$

## Checks of our $\beta$ functions?

- 1-loop results ✓ [Herbut, Scherer '22, PRB], [Liu, Huffmann, Chandrasekharan, Kaul '22, PRL]
- **(Chiral) Heisenberg** subsectors @ 2-loop ✓ [Zerf, Mihaila, Marquard, Herbut, Scherer '17, PRD]
- Independent 2-loop results using **ARGES** ✓ [Litim, Steudtner '21, Comp. Phys. Comm.]

**Finally:**

**Fixed-Point Analysis @ 2-loops for QMC model**

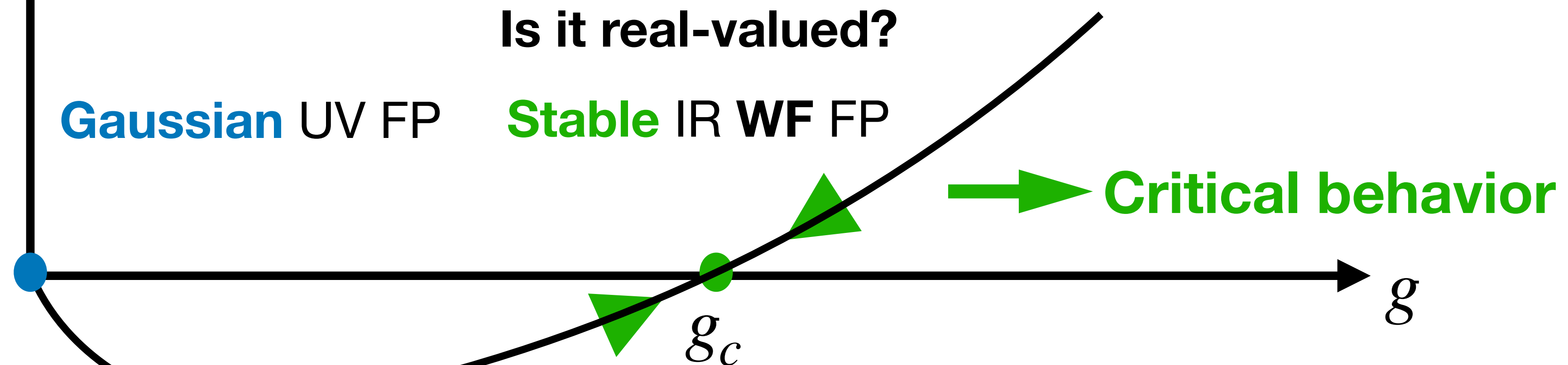
# Fixed-Point Analysis @ 2-loop

- $d = 4 - \epsilon$  theory has **Wilson–Fisher** FPs in the **IR** – **UV**: trivial Gaussian FP
- Linearize  $\beta$  functions around FP  $g_c$   $\longrightarrow$  **stability matrix S**

$$\beta(g) = S(g - g_c) + \mathcal{O}((g - g_c)^2)$$

- **Eigenvalues** determine stability

$$\beta(g) = \frac{dg}{d \log \mu}$$



# Fixed-Point Analysis @ 2-loop

- Study model corresponding to **quantum Monte Carlo**:  $g_a = g_b, \lambda_a = \lambda_b, N_f = 2$

$$\beta_{g^2} = -\epsilon g^2 + \frac{1}{16\pi^2} g^4 (5 + N_f) + \frac{1}{(16\pi^2)^2} g^2 \left[ g^4 \left( -6N_f - \frac{49}{8} \right) - g^2 (5\lambda + 9\lambda_c) + \frac{5}{2} \lambda^2 + \frac{3}{2} \lambda_c^2 \right]$$

$$\beta_\lambda = -\epsilon \lambda + \frac{1}{16\pi^2} \left[ -N_f g^4 + 2N_f g^2 \lambda + 11\lambda^2 + 3\lambda_c^2 \right]$$

$$+ \frac{1}{(16\pi^2)^2} \left[ 8N_f g^6 + N_f g^4 \left( 3\lambda_c - \frac{9}{2} \lambda \right) - N_f g^2 (11\lambda^2 + 3\lambda_c^2) - 3 (23\lambda^3 + 5\lambda\lambda_c^2 + 4\lambda_c^3) \right]$$

$$\beta_{\lambda_c} = -\epsilon \lambda_c + \frac{1}{16\pi^2} \left[ -3N_f g^4 + 2N_f g^2 \lambda_c + 2\lambda_c (5\lambda + 2\lambda_c) \right]$$

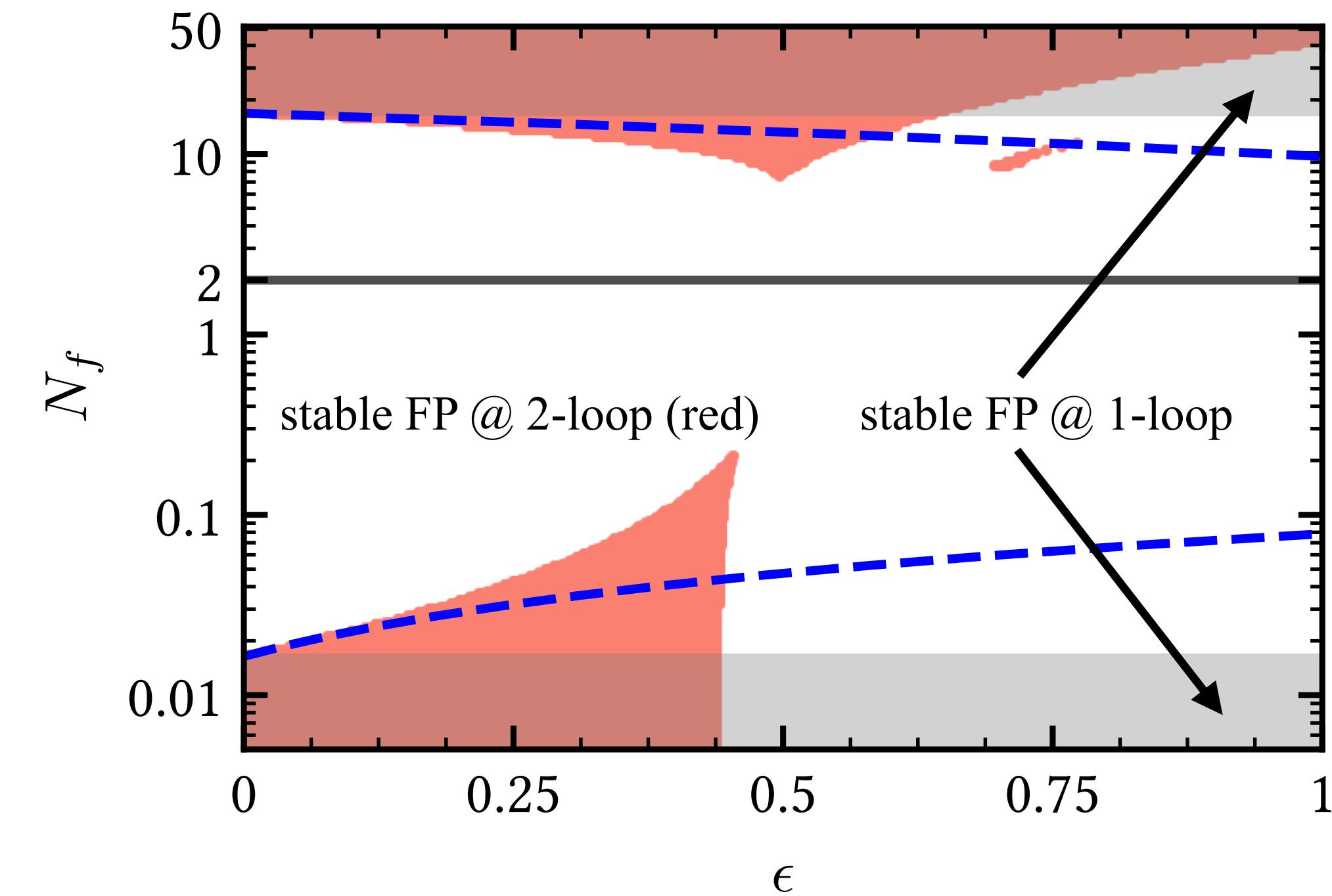
$$+ \frac{1}{(16\pi^2)^2} \left[ 16N_f g^6 + N_f g^4 \left( 5\lambda - \frac{5}{2} \lambda_c \right) - 2N_f g^2 \lambda_c (5\lambda + 2\lambda_c) - \lambda_c (5\lambda + \lambda_c) (5\lambda + 11\lambda_c) \right]$$

[MU, Herbut, Stamou, Scherer '23, PRB]

# Fixed-Point Analysis @ 2-loop

- Study model corresponding to **quantum Monte Carlo**:  $g_a = g_b, \lambda_a = \lambda_b, N_f = 2$

[MU, Herbut, Stamou, Scherer '23, PRB]



**Numerically search for FPs @ FP:  $m_{\text{OPs}} = 0$**

- 1-loop**: no stable FP — FP Annihilation at

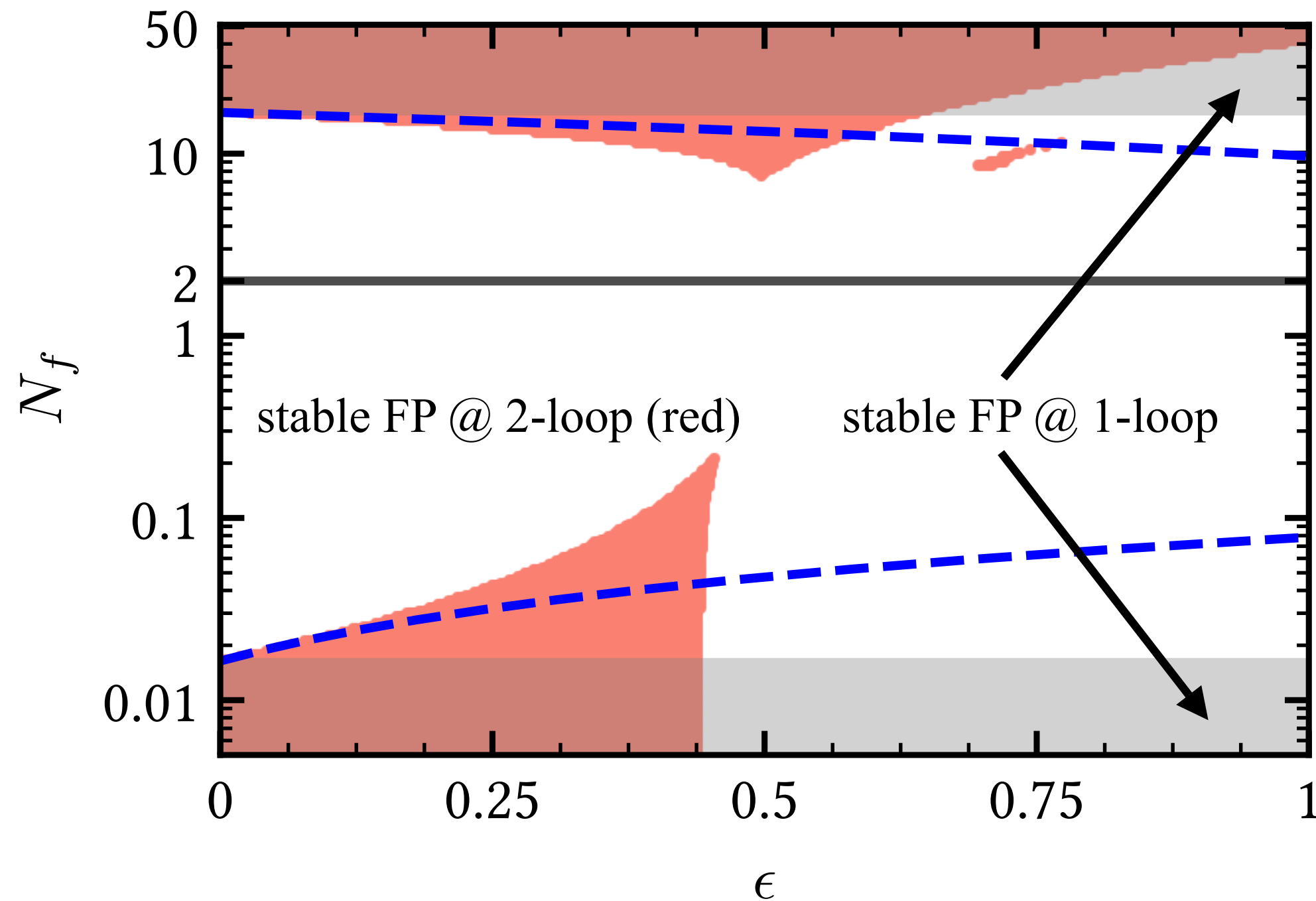
$$N_c^> \approx 16.83$$



# Fixed-Point Analysis @ 2-loop

- Study model corresponding to **quantum Monte Carlo**:  $g_a = g_b, \lambda_a = \lambda_b, N_f = 2$

[MU, Herbut, Stamou, Scherer '23, PRB]



**Numerically search for FPs @ FP:  $m_{\text{OPs}} = 0$**

- 1-loop**: no stable FP — FP Annihilation at

$$N_c^> \approx 16.83$$

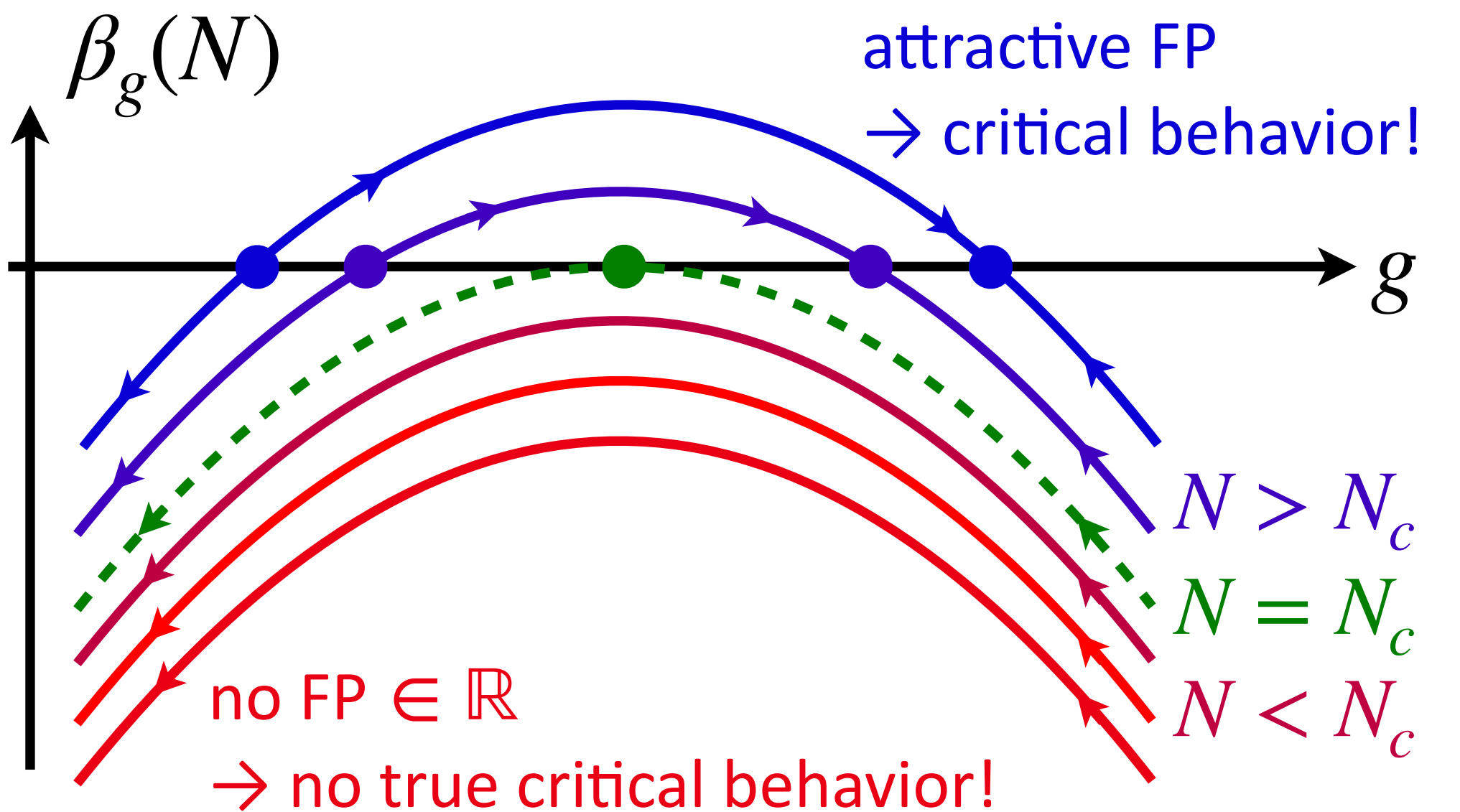
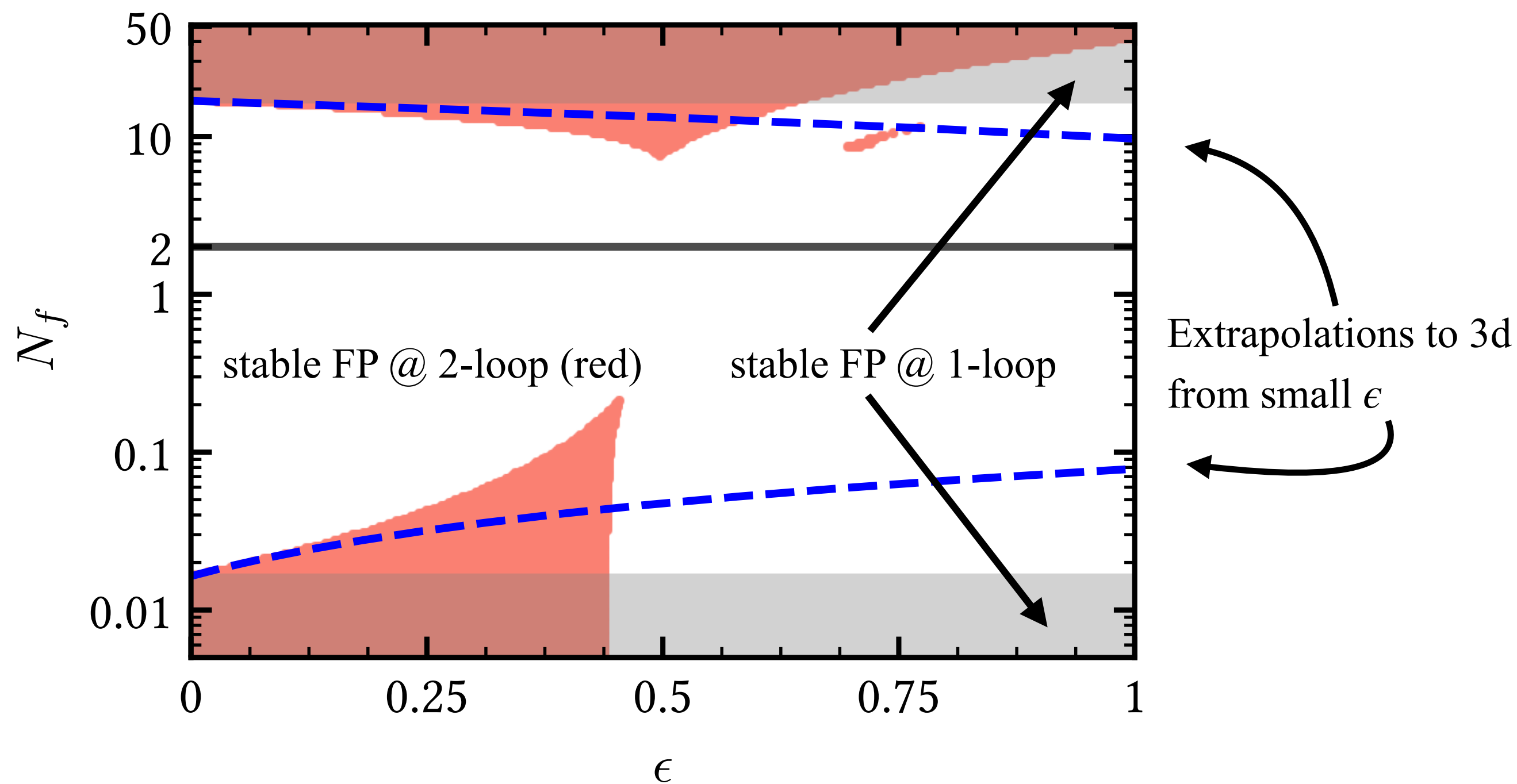


Fig: [Song, Zhao, Janssen, Scherer, Meng '23]

# Fixed-Point Analysis @ 2-loop

- Study model corresponding to **quantum Monte Carlo**:  $g_a = g_b, \lambda_a = \lambda_b, N_f = 2$

[MU, Herbut, Stamou, Scherer '23, PRB]



**Numerically search for FPs @ FP:  $m_{\text{OPs}} = 0$**

- 2-loop**: no stable FP — FP Annihilation at

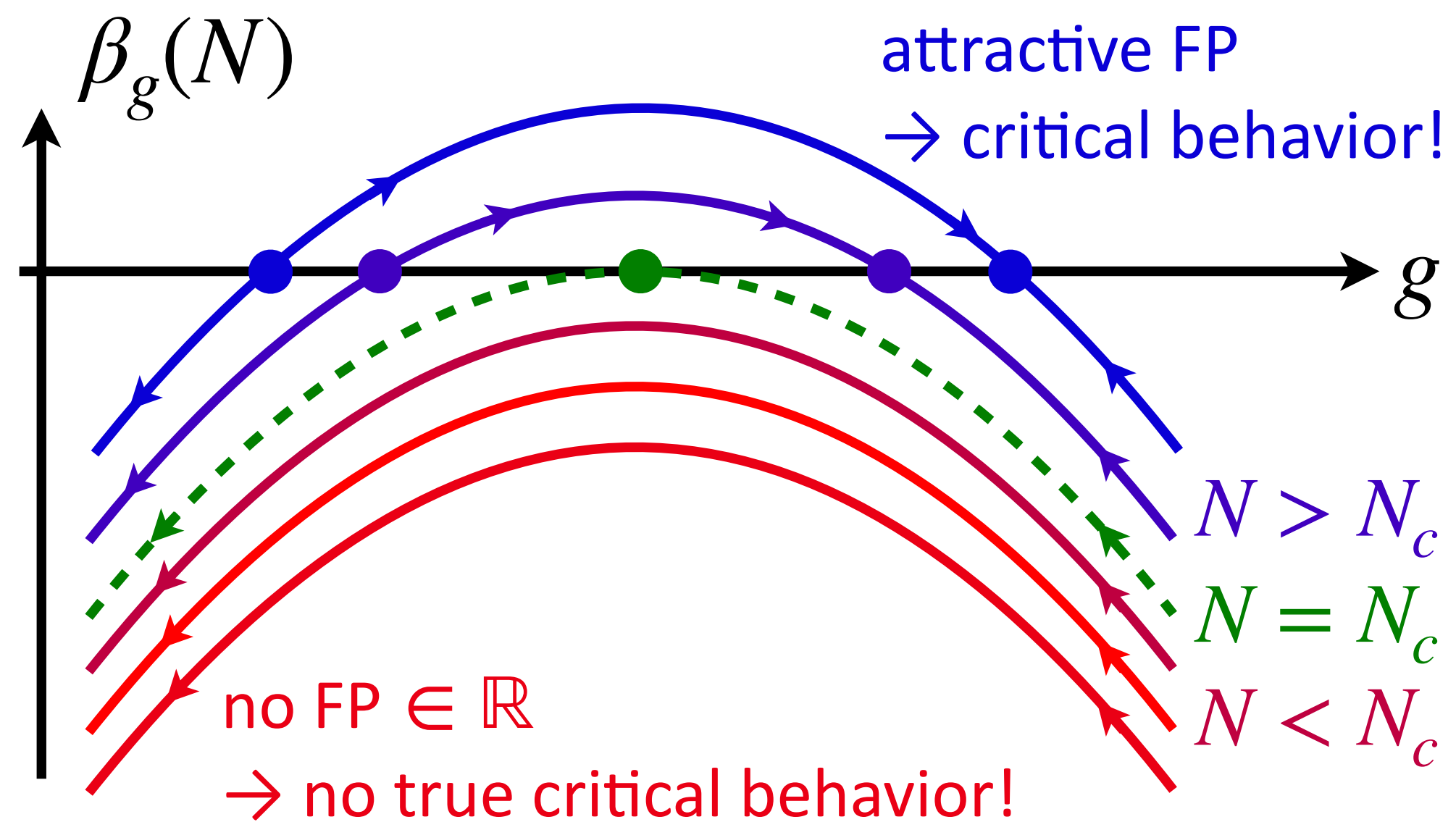
$$N_c^> \approx 16.83 - 7.14\epsilon$$

**➔ how to explain QMC data?**

# Fixed-Point Analysis @ 2-loop

- Study model corresponding to **quantum Monte Carlo**:  $g_a = g_b, \lambda_a = \lambda_b, N_f = 2$

Fig: [Song, Zhao, Janssen, Scherer, Meng '23]



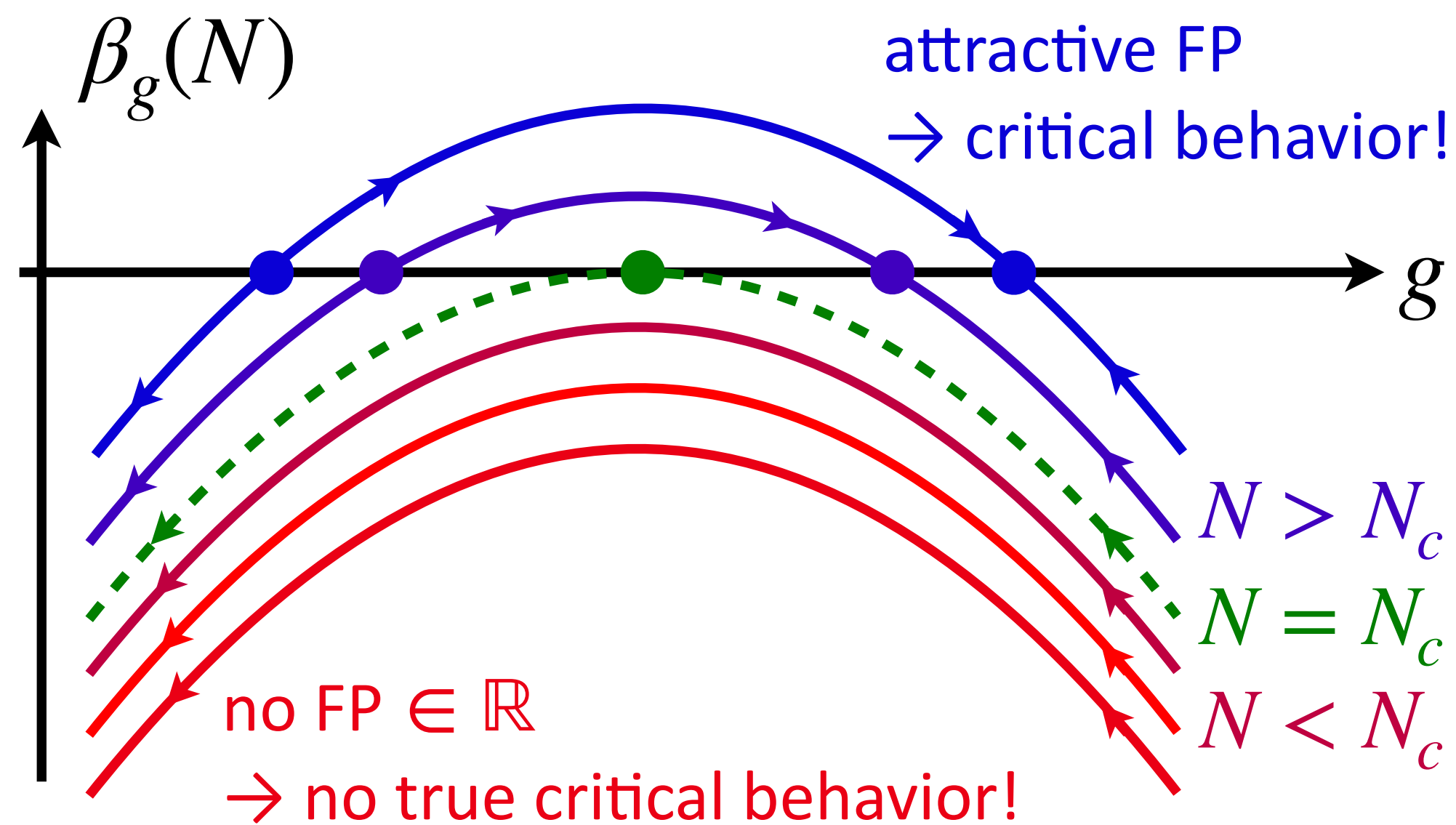
## How to **possibly** explain QMC data?

- For  $N_f < N_c$ : complex FPs
- Close to  $N_c$ : **slowly walking** RG flow
- Corr. length  $\xi$  large but finite

# Fixed-Point Analysis @ 2-loop

- Study model corresponding to **quantum Monte Carlo**:  $g_a = g_b, \lambda_a = \lambda_b, N_f = 2$

Fig: [Song, Zhao, Janssen, Scherer, Meng '23]



## How to possibly explain QMC data?

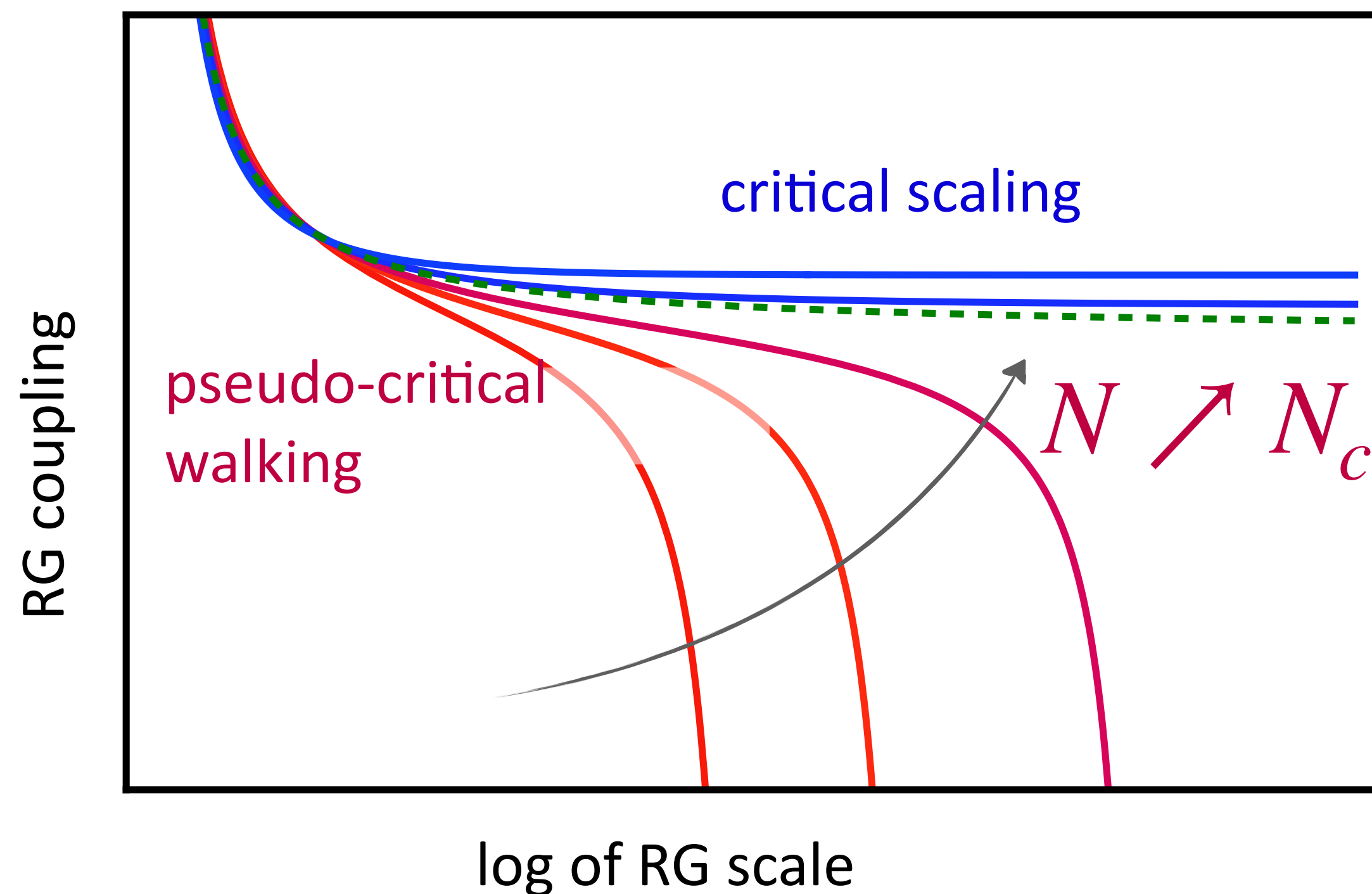
- For  $N_f < N_c$ : complex FPs
- Close to  $N_c$ : **slowly walking** RG flow
  - Corr. length  $\xi$  large but finite
- Higher orders suppress  $N_c$  to  $N_f = 2$**

**$\rightarrow$  Stable FP restored**

# Fixed-Point Analysis @ 2-loop

- Study model corresponding to **quantum Monte Carlo**:  $g_a = g_b, \lambda_a = \lambda_b, N_f = 2$

Fig: [Song, Zhao, Janssen, Scherer, Meng '23]



## How to **possibly** explain QMC data?

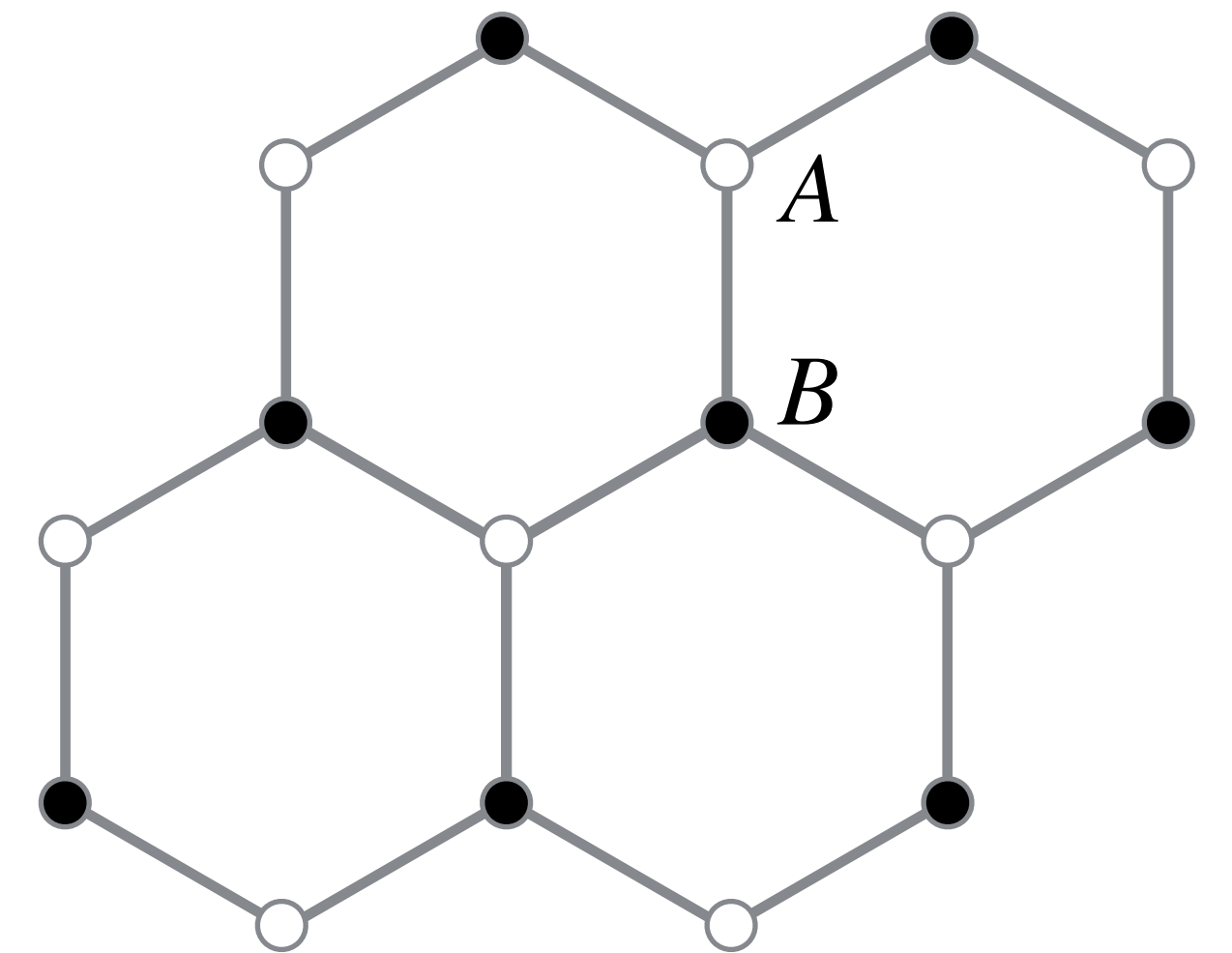
- For  $N_f < N_c$ : complex FPs
- Close to  $N_c$ : **slowly walking** RG flow
- Corr. length  $\xi$  large but finite
- **Drifting critical exponents**  
[Kaplan, Lee, Son, Stephanov '09, PRD]

➔ **Challenging to confirm with QMC**

➔ **Weak First-Order Transition**

# Conclusion & Outlook

- Dirac theory emerges in **Dirac Materials** like graphene
- **Criticality**: Dynamic fermion mass generation through **SSB**
- **Independent benchmark** of QMC / many-body methods [Liu, Huffmann, Chandrasekharan, Kaul '22, PRL]
- RG @ 2-loop in  $d = 4 - \epsilon$ : No stable FP  $\longrightarrow$  **QMC-observed Criticality lost?**  
[MU, Herbut, Stamou, Scherer '23, PRB]
- FP restored at loop orders  $> 2$  ?
- Pseudo-Critical Behavior/Walking ? [Kaplan, Lee, Son '09, PRD]



$$m, \langle \Phi \rangle = 0 \quad | \quad m, \langle \Phi \rangle \neq 0$$

A phase diagram showing a transition at a critical coupling  $g_c$ . The horizontal axis is labeled  $g$ . A vertical tick mark on the axis is labeled  $g_c$ . To the left of  $g_c$ , the mass  $m$  is zero and the order parameter  $\langle \Phi \rangle$  is zero. To the right of  $g_c$ , the mass  $m$  is non-zero and the order parameter  $\langle \Phi \rangle$  is non-zero.