

Spectral Functions and Conformal Windows

Manuel Reichert

University of Sussex, Brighton, UK

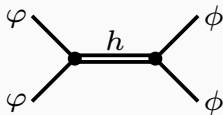
Asymptotic Safety meets Particle Physics & Friends

DESY, Hamburg, 18. Dezember 2023

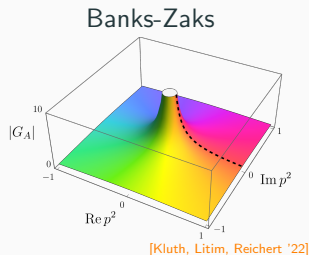
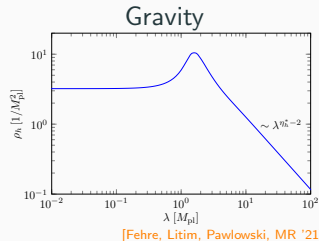


Correlation functions of the quantum effective action

- Want access to full quantum 1PI effective action Γ
- Special importance: propagator $G \sim \Gamma^{(2)^{-1}}$
- Important input for scattering



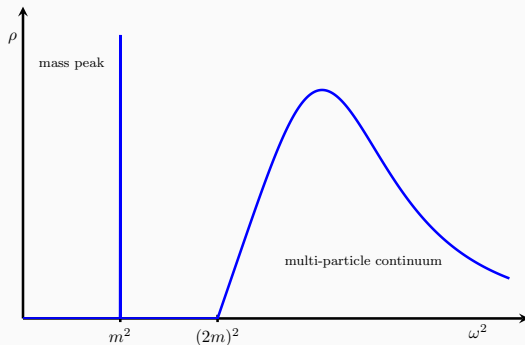
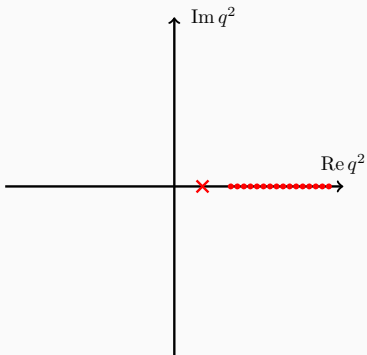
- Test properties of spectral function in conformal windows
- Understand quantum-gravity through comparison with perturbative theories



$$G(q^2) = \int_0^{\infty} \frac{d\lambda^2}{\pi} \frac{\rho(\lambda^2)}{q^2 - \lambda^2}$$

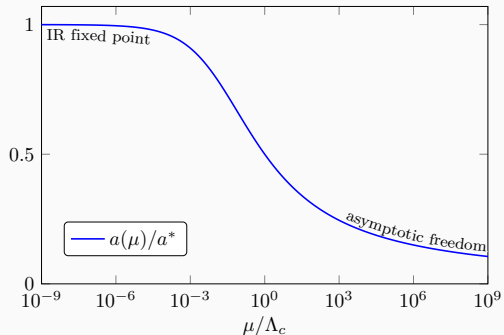
with

$$\rho(\omega^2) = - \lim_{\varepsilon \rightarrow 0} \text{Im } G(\omega^2 + i\varepsilon)$$

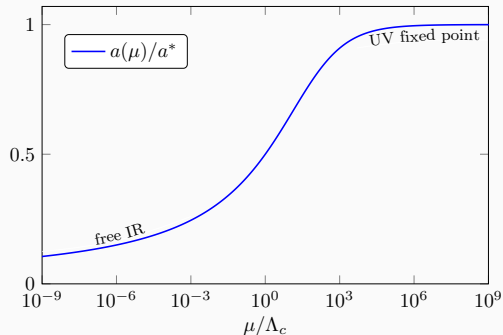


with $\rho(\omega^2) > 0$ and $\int \rho(\lambda^2) d\lambda^2 = 1$

Asymptotic freedom vs safety



- Banks-Zaks with $\frac{N_f}{N_c} = \frac{11}{2} - \varepsilon$
- Wilson-Fisher in $d = 4 - \varepsilon$



- Litim-Sannino with $\frac{N_f}{N_c} = \frac{11}{2} + \varepsilon$
- Quantum Gravity in $d = 2 + \varepsilon$

$SU(N_c)$ gauge theory with N_f quarks

$$\beta(a) = \beta_1 a^2 + \beta_2 a^3 + \dots$$

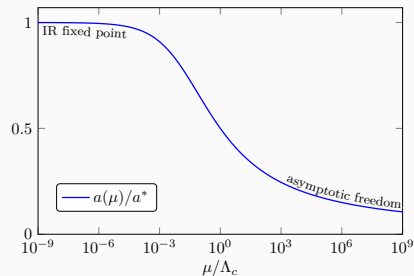
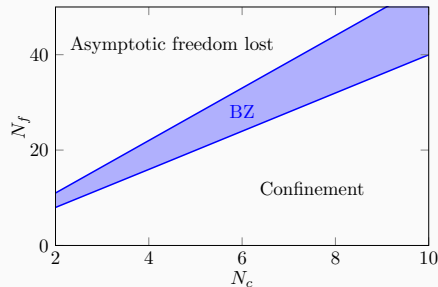
Veneziano limit gives $\beta_1 \sim -\varepsilon$ with

$$\varepsilon = \frac{11}{2} - \frac{N_f}{N_c} > 0$$

$$N_c \ \& \ N_f \rightarrow \infty$$

Perturbative IR fixed point

$$a_* = -\frac{\beta_1}{\beta_2} + \mathcal{O}(\beta_1^2)$$



$SU(N_c)$ with N_f quarks and uncharged $N_f \times N_f$ matrix scalar

$$L \sim L_{\text{YM}} + L_{\text{kin}} - y \text{Tr} \bar{\psi}_L H \psi_R - u \text{Tr} H^\dagger H H^\dagger H - v (\text{Tr} H^\dagger H)^2$$

Veneziano limit N_c & $N_f \rightarrow \infty$ with

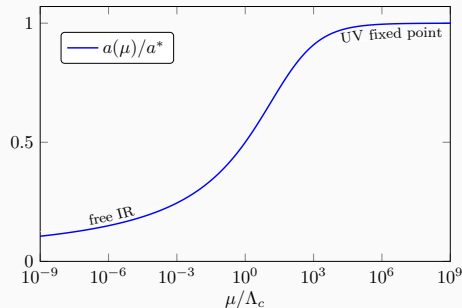
$$\varepsilon = \frac{11}{2} - \frac{N_f}{N_c} < 0$$

Perturbative UV fixed point

$$a_g^*, a_y^*, a_u^*, a_v^* \sim \varepsilon$$

One relevant direction

$$a_y(a_g(\mu)), a_u(a_g(\mu)), a_v(a_g(\mu))$$



Litim-Sannino model – asymptotically free

$SO(2N_c)$ with N_f quarks and uncharged $N_f \times N_f$ matrix scalar

$$L \sim L_{\text{YM}} + L_{\text{kin}} - y \text{Tr} \bar{\psi}_L H \psi_R - u \text{Tr} H^\dagger H H^\dagger H - v (\text{Tr} H^\dagger H)^2$$

Veneziano limit N_c & $N_f \rightarrow \infty$ with

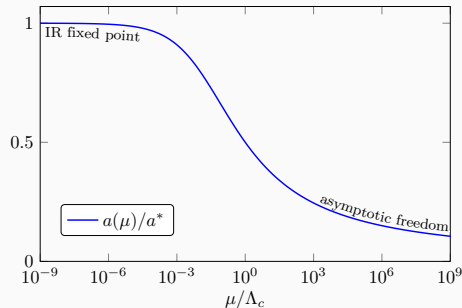
$$\varepsilon = \frac{11}{2} - \frac{N_f}{N_c} > 0$$

Perturbative UV fixed point

$$a_g^*, a_y^*, a_u^*, a_v^* \sim \varepsilon$$

One IR attractive direction

$$a_y(a_g(\mu)), a_u(a_g(\mu)), a_v(a_g(\mu))$$



Propagator with self-energy loop corrections

$$G_\phi(p^2, \mu^2) = \frac{1}{-p^2} \frac{1}{1 + \Pi_\phi(p^2, \mu^2)}$$

Callan-Symanzik equation to resum large logarithms $\log(p^2/\mu^2)$

$$\mu^2 \frac{d}{d\mu^2} (Z_\phi G_\phi) = \left(\mu^2 \frac{\partial}{\partial \mu^2} + \beta(a) \frac{\partial}{\partial a} - \gamma_\phi \right) G_\phi = 0$$

Propagator with self-energy loop corrections

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Resummed Callan-Symanzik propagator with $\bar{a} \equiv \bar{a}(p^2)$ & $a \equiv a(\mu^2)$

$$G_\phi = \frac{1}{p^2} \frac{\mathcal{N}_\phi}{1 + \sum_n \Pi_\phi^{(n)}(p^2 = -\mu^2) \bar{a}(p^2)^n} \left(\frac{a(\mu^2)}{\bar{a}(p^2)} \right)^{\gamma_\phi^{(1)}/\beta_1} \prod_i \left(\frac{a(\mu^2) - a_{*,i}}{\bar{a}(p^2) - a_{*,i}} \right)^{\gamma_\phi(a_{*,i})/\theta_i}$$

Callan-Symanzik equation with multiple couplings

$$\left(\mu^2 \frac{\partial}{\partial \mu^2} + \sum_i \beta_i \frac{\partial}{\partial a_i} - \gamma_\phi \right) G_\phi = 0$$

One monotonic coupling to re-express all others $a_i(\mu^2) = a_i(a(\mu^2))$, leads to same CS equation

$$\mu^2 \frac{d}{d\mu^2} (Z_\phi G_\phi) = \left(\mu^2 \frac{\partial}{\partial \mu^2} + \beta_{\text{eff}} \frac{\partial}{\partial a} - \gamma_\phi \right) G_\phi = 0$$

Callan-Symanzik Resummation – discussion

Resummed Callan-Symanzik propagator with $\bar{a} \equiv \bar{a}(p^2)$ & $a \equiv a(\mu^2)$

$$G_\phi = \frac{1}{p^2} \frac{\mathcal{N}_\phi}{1 + \sum_n \Pi_\phi^{(n)} \bar{a}^n} \left(\frac{a}{\bar{a}}\right)^{\gamma_\phi^{(1)}/\beta_1} \prod_i \left(\frac{a - a_{*,i}}{\bar{a} - a_{*,i}}\right)^{\gamma_\phi(a_{*,i})/\theta_i}$$

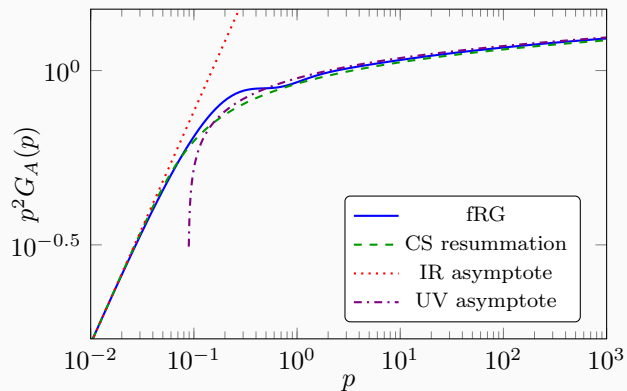
- All fixed points including complex and negative ones contribute
- Bound states can appear through the self-energies
- Self-energies and anomalous dimensions are gauge dependent for charged particles
- Banks-Zaks: 5-loop β and γ -functions & 4-loop self-energies

[Ruijl, Ueda, Vermaseren, Vogt '17; Herzog, Ruijl, Ueda, Vermaseren, Vogt '17; Chetyrkin, Falcioni, Herzog, Vermaseren '17]

Litim-Sannino: 433-loop β and γ -functions & no self-energies

[Bond, Litim, Vazquez, Steudtner '18; Litim, Riyaz, Stamou, Steudtner '23]

Intermezzo – comparison to fRG computation



[Kluth, Litim, Reichert '22]

Very good agreement between fRG and CS resummation

Existence and normalisation of spectral functions

$$G_\phi = \frac{1}{p^2} \frac{\mathcal{N}_\phi}{1 + \sum_n \Pi_\phi^{(n)} \bar{a}^n} \left(\frac{a}{\bar{a}}\right)^{\gamma_\phi^{(1)}/\beta_1} \prod_i \left(\frac{a - a_{*,i}}{\bar{a} - a_{*,i}}\right)^{\gamma_\phi(a_{*,i})/\theta_i}$$

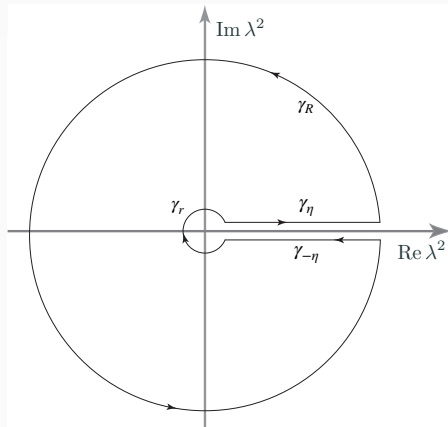
When can we use

$$G(q^2) = \int_0^\infty \frac{d\lambda^2}{\pi} \frac{\rho(\lambda^2)}{q^2 - \lambda^2}$$

with

$$\rho(\omega^2) = - \lim_{\varepsilon \rightarrow 0} \text{Im } G(\omega^2 + i\varepsilon)$$

Assume no complex structures for the moment



Existence and normalisation of spectral functions

$$G_\phi = \frac{1}{p^2} \frac{\mathcal{N}_\phi}{1 + \sum_n \Pi_\phi^{(n)} \bar{a}^n} \left(\frac{a}{\bar{a}}\right)^{\gamma_\phi^{(1)}/\beta_1} \prod_i \left(\frac{a - a_{*,i}}{\bar{a} - a_{*,i}}\right)^{\gamma_\phi(a_{*,i})/\theta_i}$$

Existence in the IR

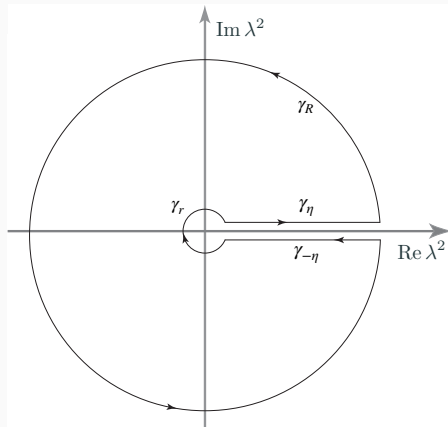
- At most singular with $G_\phi \sim 1/p^2$ for $p \rightarrow 0$
- Exception: $G_\phi \sim 1/p^{2n}$

Existence in the UV

- Decay for $p \rightarrow \infty$: $G_\phi < \text{const}$

Normalisation determined by $p \rightarrow \infty$ behaviour

$$\int \rho_\phi d\lambda^2 = \begin{cases} 0 & G_\phi < 1/p^2 \\ 1 & \text{if } G_\phi = 1/p^2 \\ \infty & G_\phi > 1/p^2 \end{cases}$$



IR safe & UV free

Existence IR:

$$\gamma_\phi^* \geq 0$$

Existence UV:

trivial

Normalisable:

$$\gamma_\phi^{(1)} = 0$$

UV safe & IR free

Existence IR

$$\gamma_\phi^{(1)} / \beta_1 \geq 0$$

Existence UV

$$\gamma_\phi^* < 1$$

Normalisable

$$\gamma_\phi^* = 0$$

IR safe & UV free

Existence IR:

$$\gamma_\phi^* \geq 0$$

Existence UV:

trivial

Normalisable:

$$\gamma_\phi^{(1)} = 0$$

UV safe & IR free

Existence IR

$$\gamma_\phi^{(1)} / \beta_1 \geq 0$$

Existence UV

$$\gamma_\phi^* < 1$$

Normalisable

$$\gamma_\phi^* = 0$$

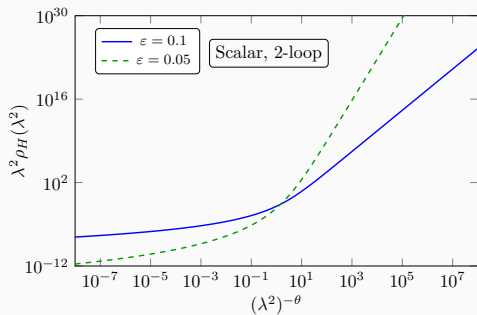
$0 \leq \gamma_\phi < 1$ necessary for existence (gauge dependent if ϕ carries gauge charge)

There is no perturbative fundamental scalar field with a normalisable spectral function

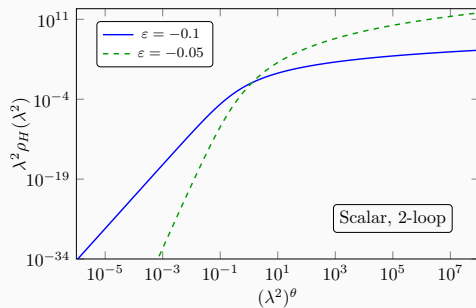
There is no perturbative fundamental scalar field with a normalisable spectral function

Example: Litim-Sannino model at 211 loop order

IR free & UV safe



IR safe & UV free



[Kluth, Litim, Reichert (in prep)]

Intermediate summary

- Fully analytic computation of propagator

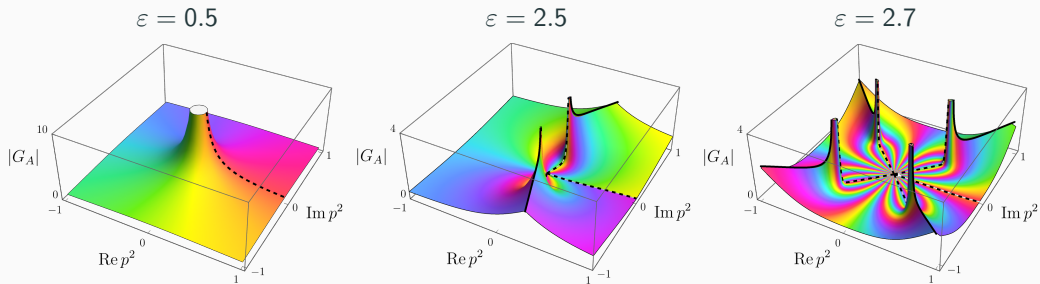
$$G_\phi = \frac{1}{p^2} \frac{\mathcal{N}_\phi}{1 + \sum_n \Pi_\phi^{(n)} \bar{a}^n} \left(\frac{a}{\bar{a}}\right)^{\gamma_\phi^{(1)}/\beta_1} \prod_i \left(\frac{a - a_{*,i}}{\bar{a} - a_{*,i}}\right)^{\gamma_\phi(a_{*,i})/\theta_i}$$

- Momentum-dependent coupling $\bar{a}(p)$ carries all momentum dependences
- $0 \leq \gamma_\phi < 1$ necessary for existence of spectral function
- No normalisable fundamental spectral functions

Can we constrain the size of conformal windows via

- non-analyticities in the complex momentum plane?
- positivity of the spectral function?

Propagator in the complex plane (Banks-Zaks at 2-loop)



[Kluth, Litim, MR '22]

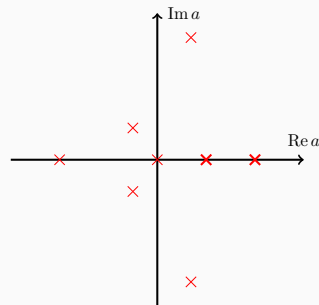
- Branch cuts and violation of KL spectral representation for large ϵ
- Branch cuts are inherited from the running coupling

Perturbative n -loop beta function

$$\beta(a) = d_a a + \beta_1 a^2 + \beta_2 a^3 + \dots + \beta_n a^{n+1}$$

Absence of branch cuts guaranteed if (a_0 initial value)

$$\left| \sum_i^{a_{i,*} > a_0} \frac{1}{\theta_i} \right| > 1$$



- Gauge independent statement due to universal critical exponents $\theta_i = \partial_a \beta(a)|_{a_{i,*}}$
- Only real and positive fixed points contribute
- Absence of branch cuts guaranteed for $\varepsilon \rightarrow 0$ since $\theta_{\text{BZ/LS}} \rightarrow 0$

Size of conformal Banks-Zaks window

$N_c = 3$	$\varepsilon_{\text{branch-cut}}$	ε_{max}
2-loop	2.2723	2.8158
3-loop	2.6798	3.5520
4-loop	2.6817	3.0538
5-loop	–	1.2019
5-loop Padé [1,3]	–	2.2183
5-loop Padé [2,2]	–	1.6993
5-loop Padé [3,1]	–	0.7304

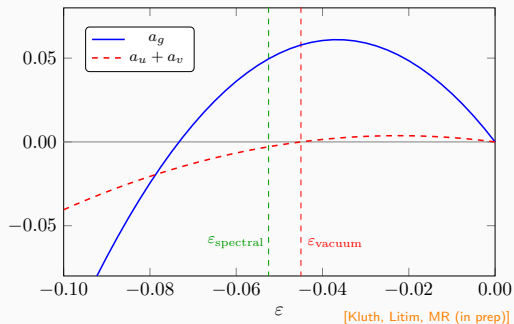
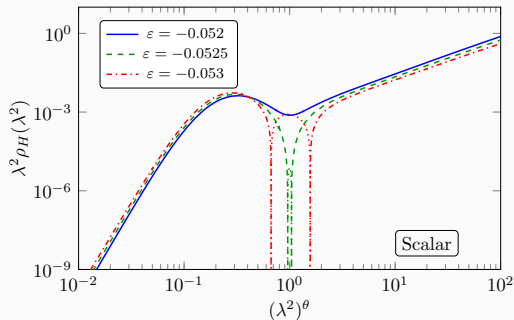
[Kluth, Litim, MR '22]

Existence of branch cuts does not post a tighter constrain on BZ window at 5-loop

Disappearance of fixed point via merger is essential

Positivity of spectral function $SO(2N_c)$ Litim-Sannino

ϵ -expansion at 322 loop order

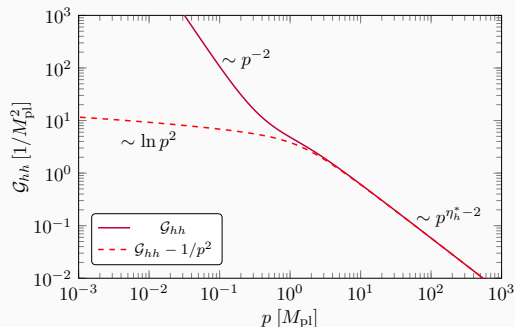
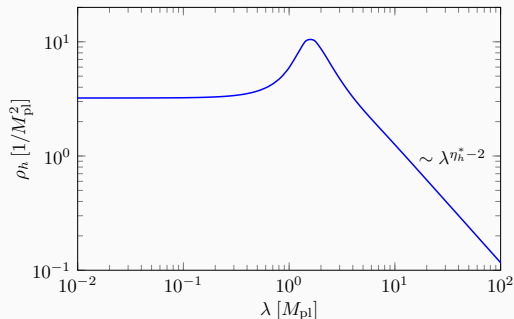


[Kluth, Litim, MR (in prep)]

- Loss of vacuum stability at $\epsilon = -0.0450$
- Loss of positivity of spectral function at $\epsilon = -0.0525$

Comparison to Quantum Gravity

Quantum Gravity results from spectral RG



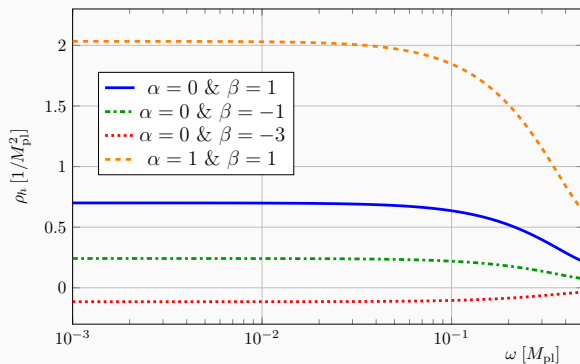
[Fehre, Litim, Pawłowski, MR '21]

- Massless graviton delta-peak with positive multi-graviton continuum
- Non-normalisable spectral function
- No ghosts and no tachyons \rightarrow no indications for unitarity violation
- Good agreement with reconstruction result and EFT

[Bonanno, Denz, Pawłowski, MR '21]

Gauge dependence of graviton spectral function

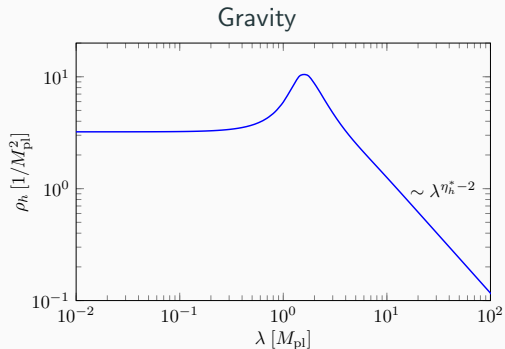
Gauge-fixing $S_{\text{gf}} = \frac{1}{\alpha} \int_x F_\mu^2$ with $F_\mu = \bar{\nabla}^\nu h_{\mu\nu} - \frac{1+\beta}{4} \bar{\nabla}_\mu h^\nu{}_\nu$



[Pawlowski, MR '23]

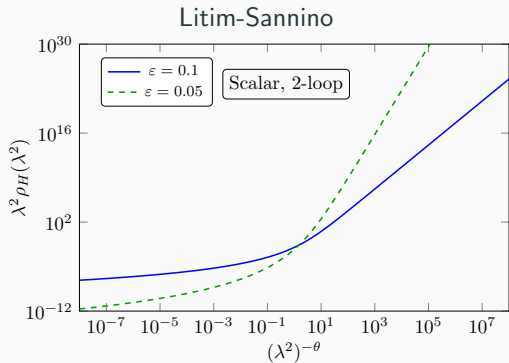
Propagator is gauge-dependent but pole structure is typically not

Spectral functions of massless particles



[Fehre, Litim, Pawłowski, MR '21]

On-shell δ -peak at $\lambda = 0$



[Kluth, Litim, MR (in prep)]

No on-shell δ -peak

Both theories are UV asymptotically safe and free in the IR

Spectral functions of massless particles – QG and QED

One-loop effective action of QG

$$\Gamma_{1\text{-loop,QG}} \sim G_{\text{N}} R \log\left(\frac{\square}{\mu^2}\right) R + \dots$$

Graviton propagator

$$\mathcal{G}_{hh}(p^2) \sim \frac{1}{p^2 + \# G_{\text{N}} \ln(p^2/\mu^2) p^4}$$

Graviton spectral function

$$\rho_h(\lambda^2) \sim \delta(\lambda^2) + \# G_{\text{N}} 2\pi + \dots$$

One-loop effective action of QED

$$\Gamma_{1\text{-loop,QED}} \sim F_{\mu\nu} \log\left(\frac{m_e^2 + \square}{m_e^2}\right) F_{\mu\nu} + \dots$$

Photon spectral function

$$\rho_\gamma(\lambda^2) \sim \delta(\lambda^2) + \# \theta(\lambda^2 - m_e^2) + \dots$$

Mass dimension of G_{N} , electron mass, lack of photon-self interactions preserve on-shell δ peak

Banks-Zaks & Litim-Sannino

- Analytic and perturbatively controlled approach to correlation and spectral functions
- $0 \leq \gamma_\phi < 1$ necessary for existence of spectral function
- Normalisability of fundamental spectral function impossible in perturbation theory
- Positivity and non-analyticities offer independent constraints on conformal windows

Asymptotically safe quantum gravity

- Direct computation of graviton spectral function with spectral fRG
- Well-behaved spectral function without cuts in the complex plane
- Key step towards scattering processes and unitarity

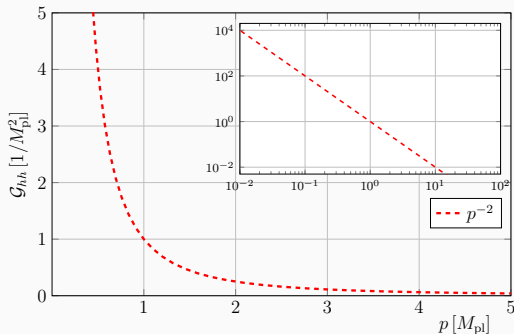
Thank you for your attention!

Back-up slides

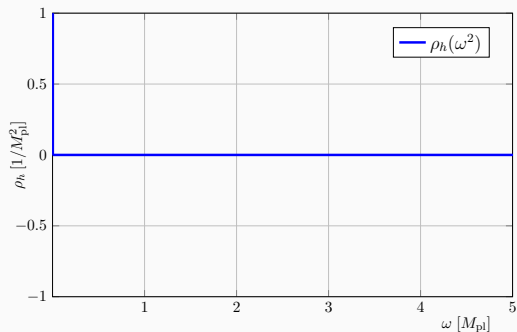
Classical graviton spectral function

$$\text{Einstein-Hilbert action: } S_{\text{EH}} = \frac{1}{16\pi G_{\text{N}}} \int_{\text{x}} \sqrt{g} (2\Lambda - R)$$

$$\text{Flat Minkowski background: } g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$



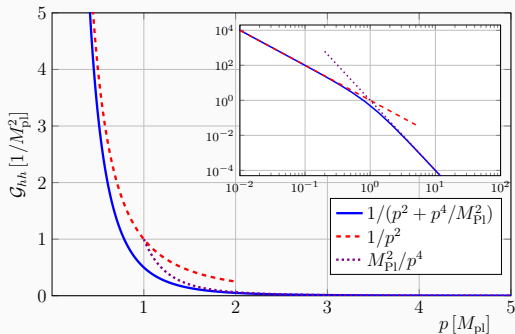
$$\mathcal{G}_{hh}(p^2) \sim \frac{1}{p^2}$$



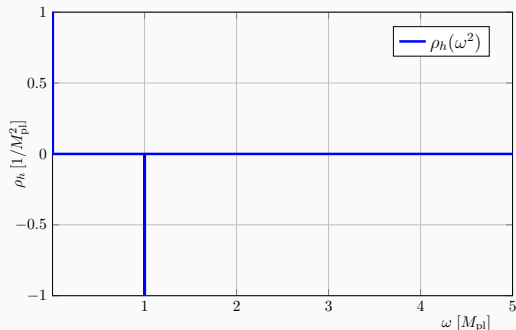
$$\rho_h(\omega^2) \sim \delta(\omega^2)$$

Classical graviton spectral function

$$\text{Higher-derivative action: } S_{\text{HD}} = S_{\text{EH}} + \int_x \sqrt{g} (aR^2 + bC_{\mu\nu\rho\sigma}^2)$$



$$\mathcal{G}_{hh}(p^2) \sim \frac{1}{p^2} - \frac{1}{M_{\text{Pl}}^2 + p^2}$$



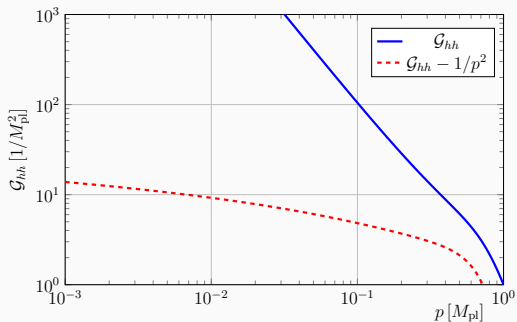
$$\rho_h(\omega^2) \sim \delta(\omega^2) - \delta(\omega^2 - M_{\text{Pl}}^2)$$

EFT graviton spectral function

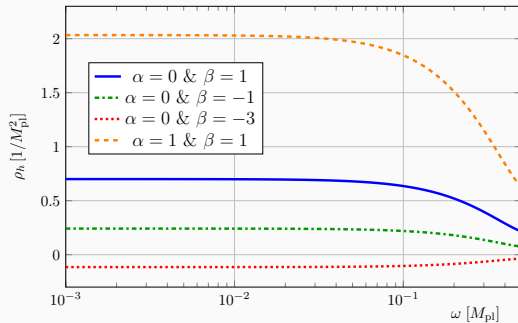
One-loop effective action:

$$\Gamma_{1\text{-loop}} = S_{\text{EH}} + \int_x \sqrt{g} (c_1 R \ln(\square) R + c_2 C_{\mu\nu\rho\sigma} \ln(\square) C^{\mu\nu\rho\sigma}) + \dots$$

Gauge-fixing $S_{\text{gf}} = \frac{1}{\alpha} \int_x F_\mu^2$ with $F_\mu = \bar{\nabla}^\nu h_{\mu\nu} - \frac{1+\beta}{4} \bar{\nabla}_\mu h^\nu{}_\nu$



$$\mathcal{G}_{hh}(p^2) \sim \frac{1}{p^2 + \ln(p^2)p^4}$$

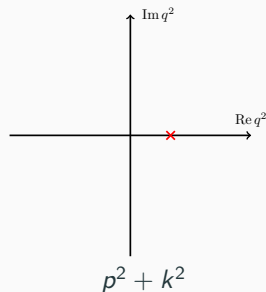
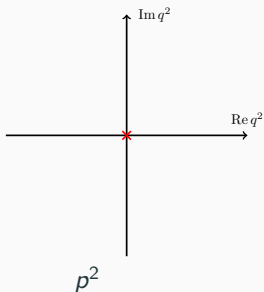


$$\rho_h(\omega^2) \sim \delta(\omega^2) + 2\pi + 4\pi\omega^2 \ln(\omega^2) + \dots$$

- Callan-Symanzik cutoff $R_k \sim k^2$ allows use of spectral representation
- Dimensional regularisation of UV divergences in $d = 4 - \varepsilon$

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \mathcal{G}_k \partial_t R_k - \partial_t \mathcal{S}_{\text{ct},k}$$

[Braun, Chen, Fu, Geißel, Horak, Huang, Ihssen, Pawłowski, MR, Rennecke, Tan, Töpfel, Wessely, Wink '22]



Lorentzian setup

- Einstein-Hilbert action expanded about flat Minkowski background
- Directly compute flow of spectral function

$$\partial_t \rho_h = -2 \operatorname{Im} \mathcal{G}_{hh}^2 \left(\partial_t \Gamma_{\text{TT}}^{(hh)} + \partial_t R_k \right)$$

- Flow of transverse-traceless graviton two point function

$$\partial_t \Gamma_k^{(hh)} = -\frac{1}{2} \left[\text{Diagram 1} + \text{Diagram 2} - 2 \text{Diagram 3} \right] - \partial_t \mathcal{S}_{\text{ct},k}^{(hh)}$$

- Schematically

$$\text{Diagram 1} = \prod_{i=1}^3 \int_0^\infty \frac{d\lambda_i}{\pi} \lambda_i \rho_h(\lambda_i) \int \frac{d^d q}{(2\pi)^d} \frac{V_{\text{3-point}}(p, q)}{(q^2 + \lambda_1^2)(q^2 + \lambda_2^2)((p+q)^2 + \lambda_3^2)}$$