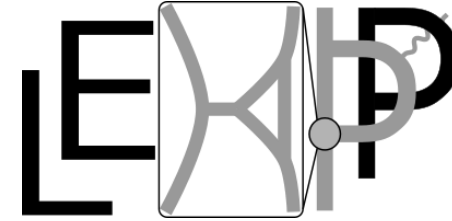
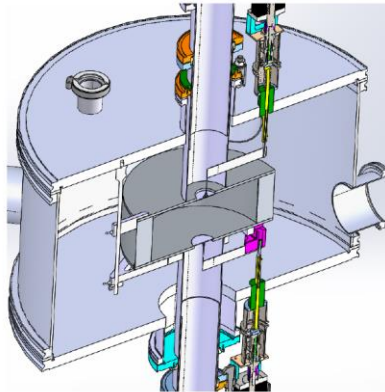


Searching for BSM physics with low-energy bound states: spin-polarized atoms and neutrons



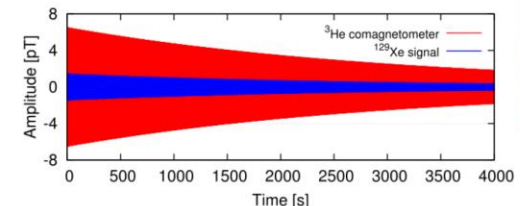
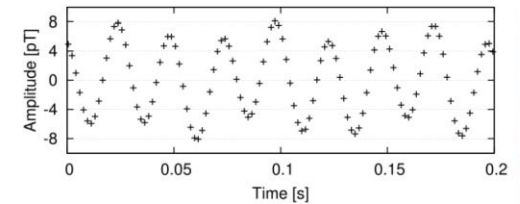
KIT Particle Physics Colloquium, 28 April 2022
Skyler Degenkolb

Case 1: trapped neutron interferometry

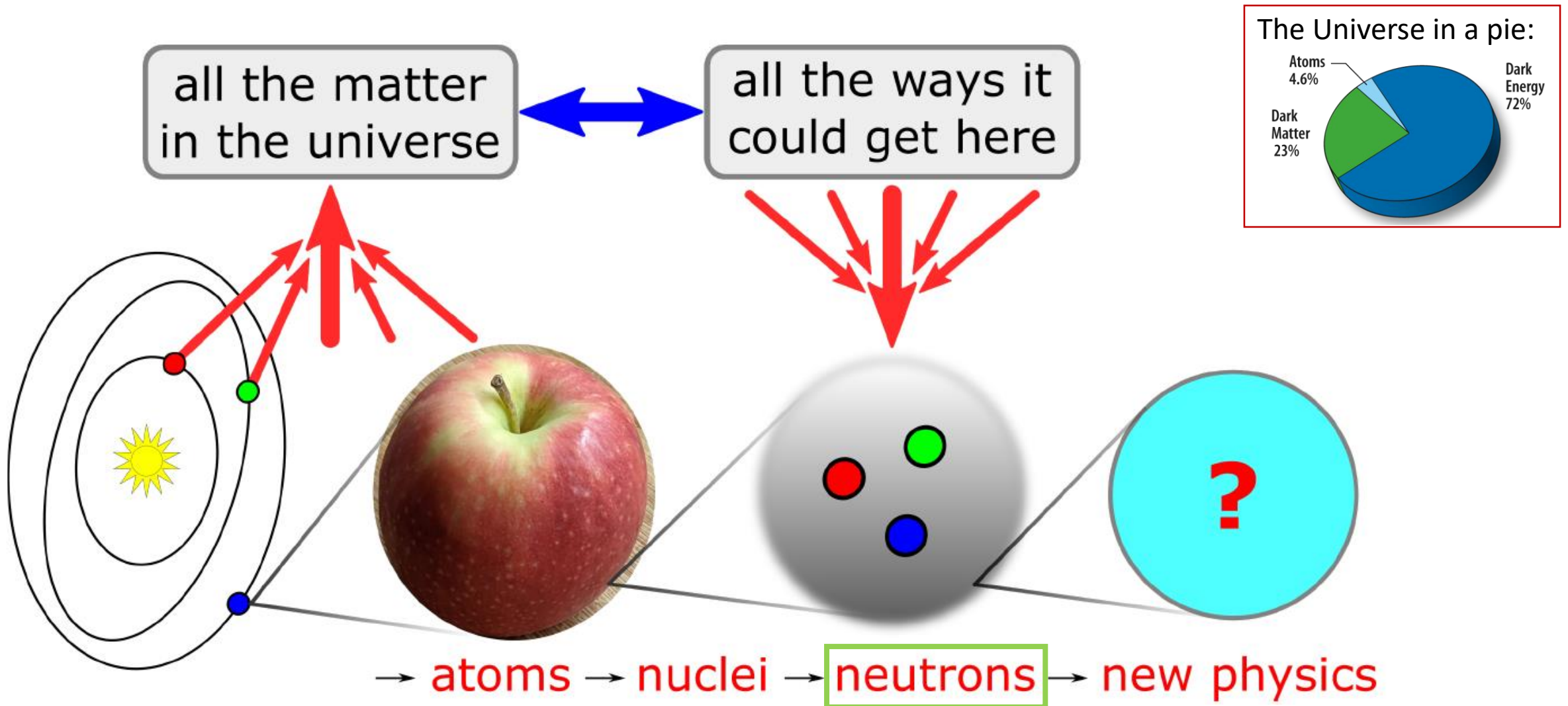


Low-Energy Precision Physics
Physikalisches Institut
Universität Heidelberg

Case 2: nuclear spin gyroscopes



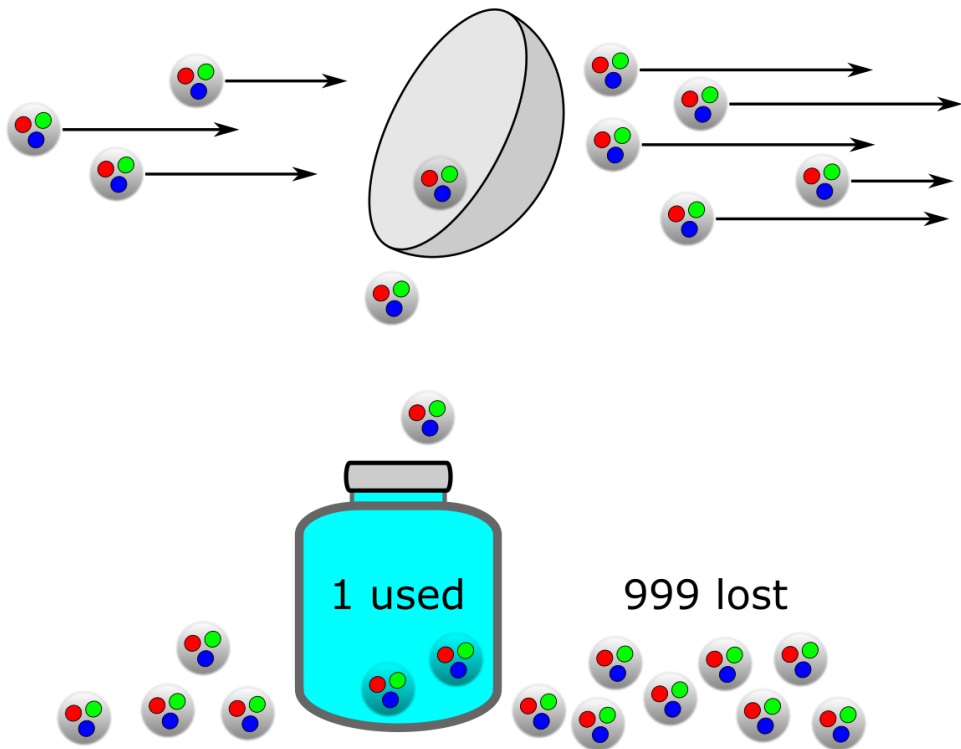
Our motivation, and a sense of scale



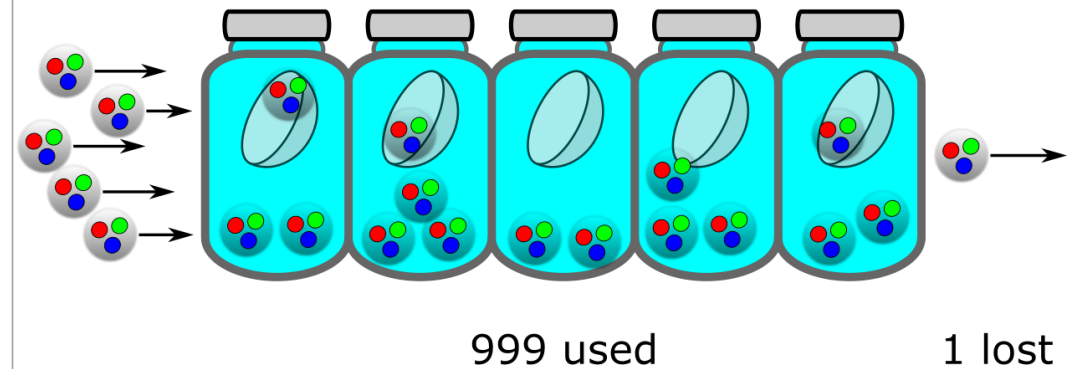
1

Challenge the first: statistics

State of the art: catch/pour
...with 0.1% success



New approach: catch them
all, directly in many bottles

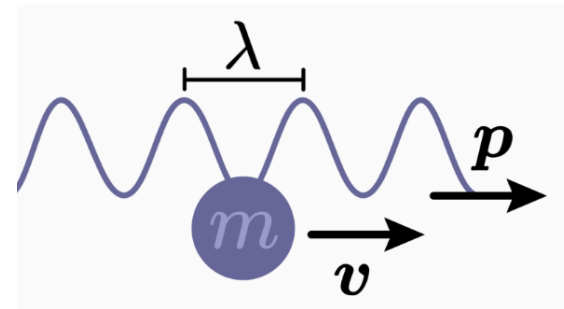


2

Challenge the second: observation time

“Never measure anything but frequency”
–Arthur Schawlow (1981 Physics Nobel Prize)

$$\delta\omega \sim \frac{1}{\delta t}$$

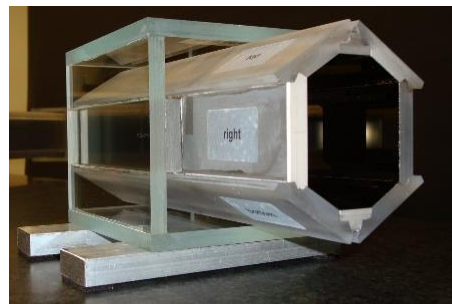


But... how to store or cool ensembles?

Wave optics, with massive particles!

“Cold” beams: O(500 m/s)

particles fly through most experiments in milliseconds



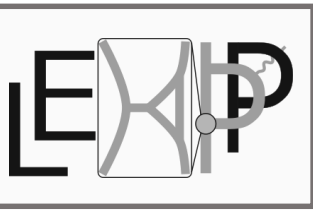
 S-DH

S-DH



“Ultracold” traps: O(5 m/s)

particles stored for minutes (>10⁵ ms)



“Permanent Electric Dipole Moment” = ?

Quantum eigenfrequencies:

$$\hbar\omega_E \propto -d\mathbf{S} \cdot \mathbf{E}$$

$$\hbar\omega_B \propto -\mu\mathbf{S} \cdot \mathbf{B}$$

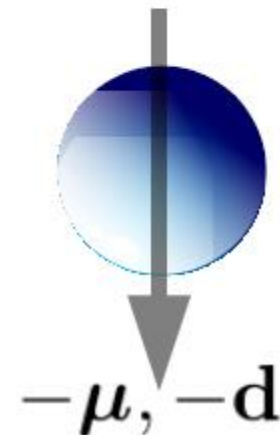
Classical moments:

$$\mathbf{d} = \int \mathbf{r}\rho(\mathbf{r})d\mathbf{r}$$

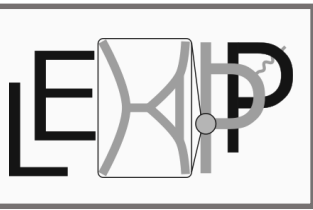
$$\boldsymbol{\mu} = \frac{1}{2} \int \mathbf{r} \times \mathbf{J}(\mathbf{r})d\mathbf{r}$$



$\pm\mathbf{S} \rightarrow \mp\mathbf{S}$



*Beware
of pictures
like this!*



“Permanent Electric Dipole Moment” = ?

Quantum eigenfrequencies:

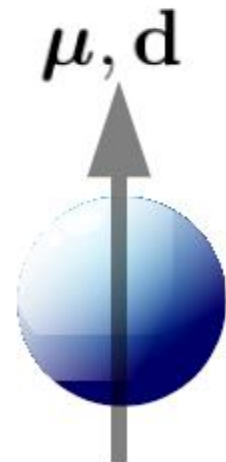
$$\hbar\omega_E \propto -d\mathbf{S} \cdot \mathbf{E}$$

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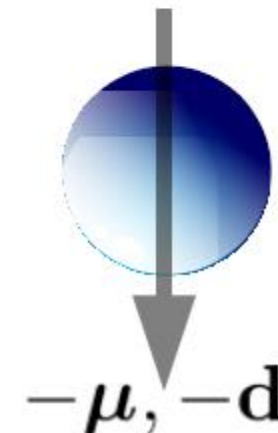
Classical moments:

$$\mathbf{d} = \int \mathbf{r}\rho(\mathbf{r})d\mathbf{r}$$

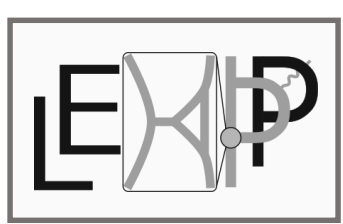
$$\boldsymbol{\mu} = \frac{1}{2} \int \mathbf{r} \times \mathbf{J}(\mathbf{r})d\mathbf{r}$$



$\pm\mathbf{S} \rightarrow \mp\mathbf{S}$

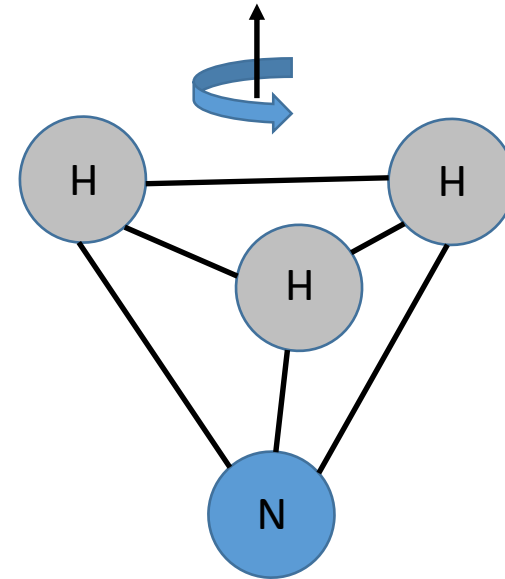
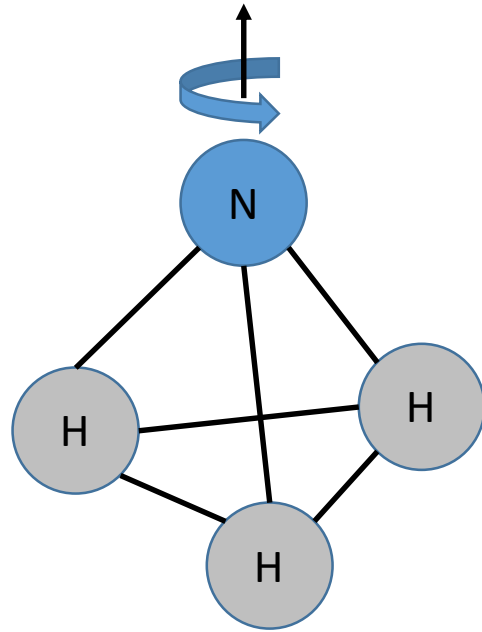


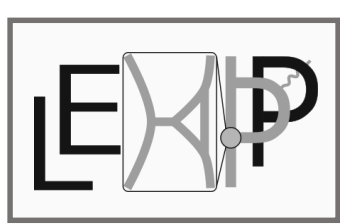
$T\mathbf{r}T^{-1} = \mathbf{r}$
$T\mathbf{p}T^{-1} = -\mathbf{p}$
$T\boldsymbol{\sigma}T^{-1} = -\boldsymbol{\sigma}$



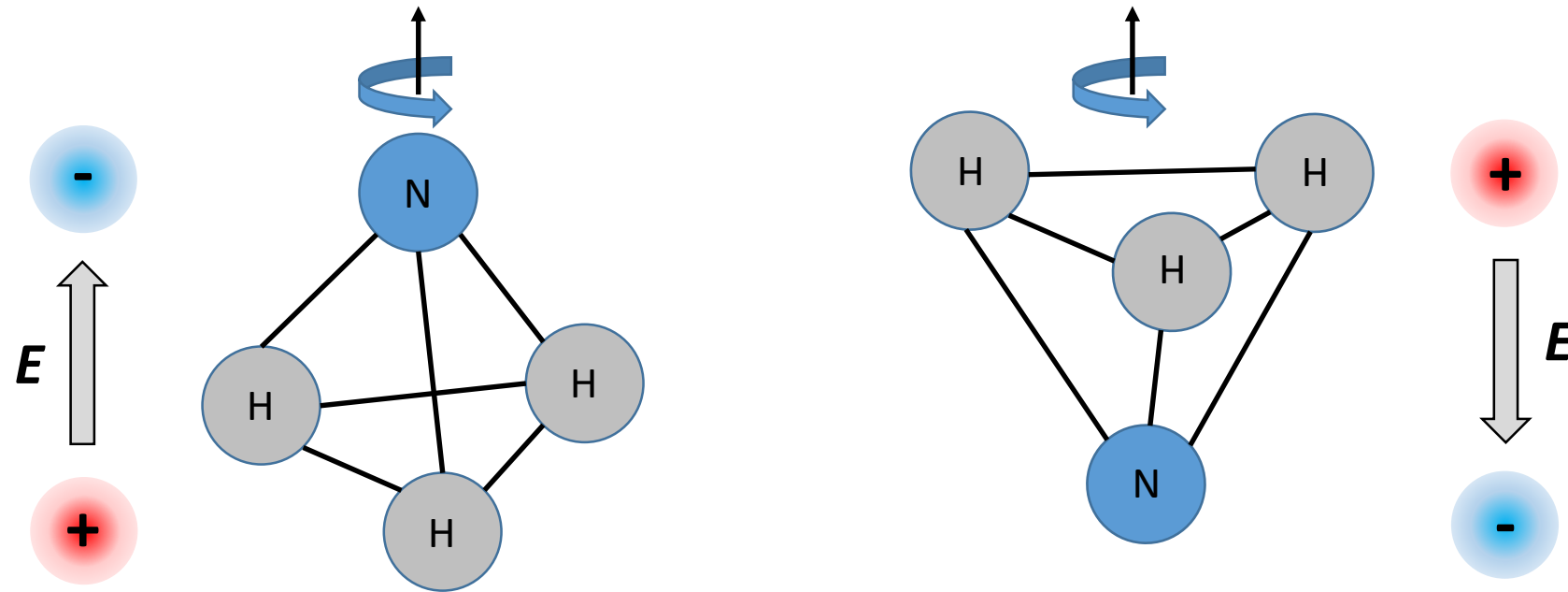
Is it different from a molecular dipole?

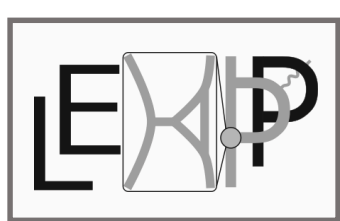
...or, "a warm-up for non-relativistic quantum methods"



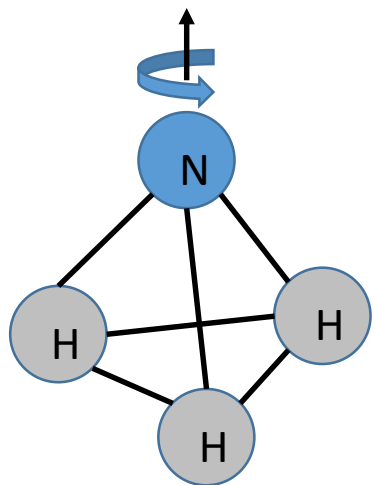


Is it different from a molecular dipole?

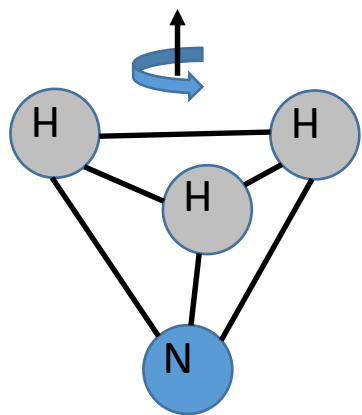




Is it different from a molecular dipole?



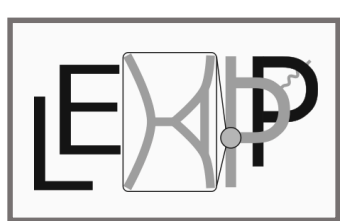
|1⟩



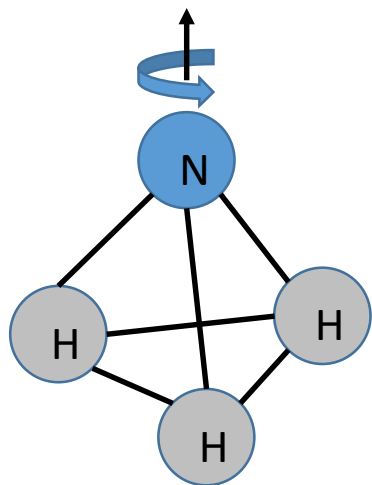
|2⟩

The *energy* eigenstates are:

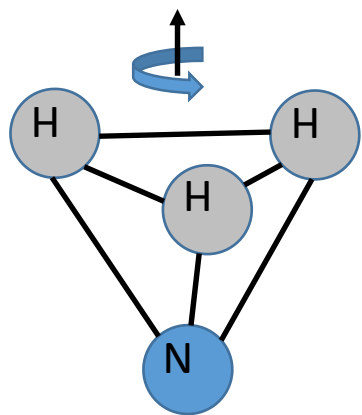
$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$



Is it different from a molecular dipole?



$|1\rangle$

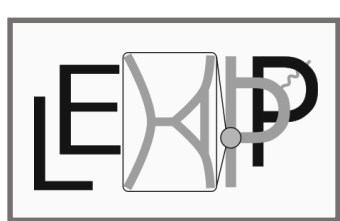


$|2\rangle$

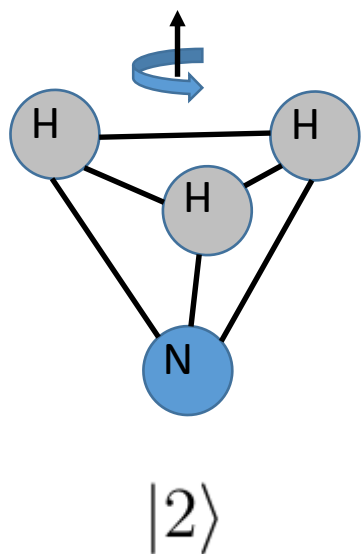
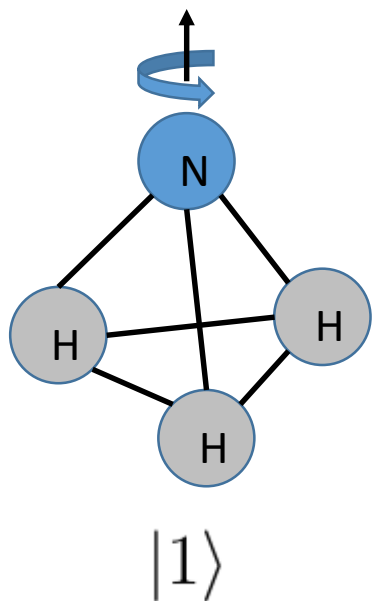
$$E_{\pm} = E_0 + \sqrt{A^2 + d^2 E^2}$$

The *energy* eigenstates are:

$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$



Is it different from a molecular dipole?



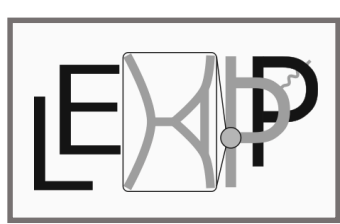
$$E_{\pm} = E_0 + \sqrt{A^2 + d^2 E^2}$$

Two blue arrows point from the square root term to the two resulting energy expressions below.

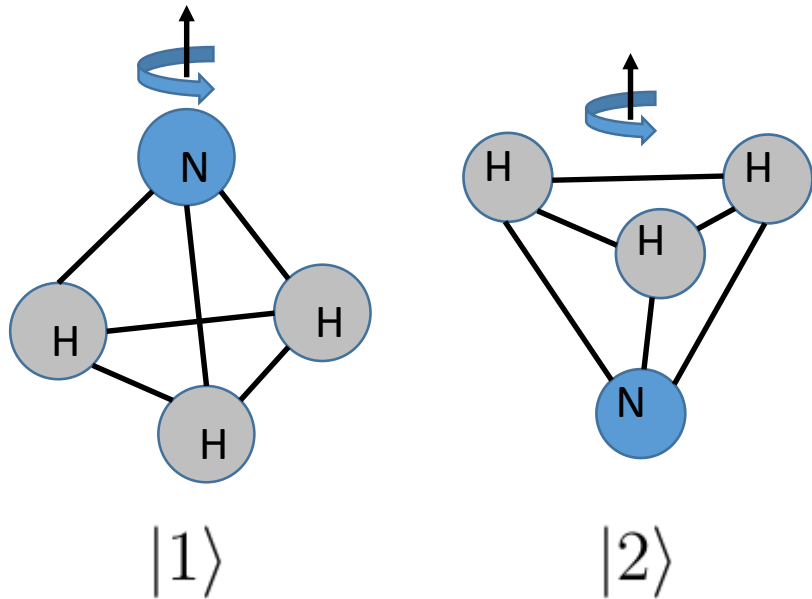
$$E_0 \pm dE \quad (dE \gg A)$$
$$E_0 \pm A \pm \frac{d^2 E^2}{2A} \quad (dE \ll A)$$

The *energy* eigenstates are:

$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$

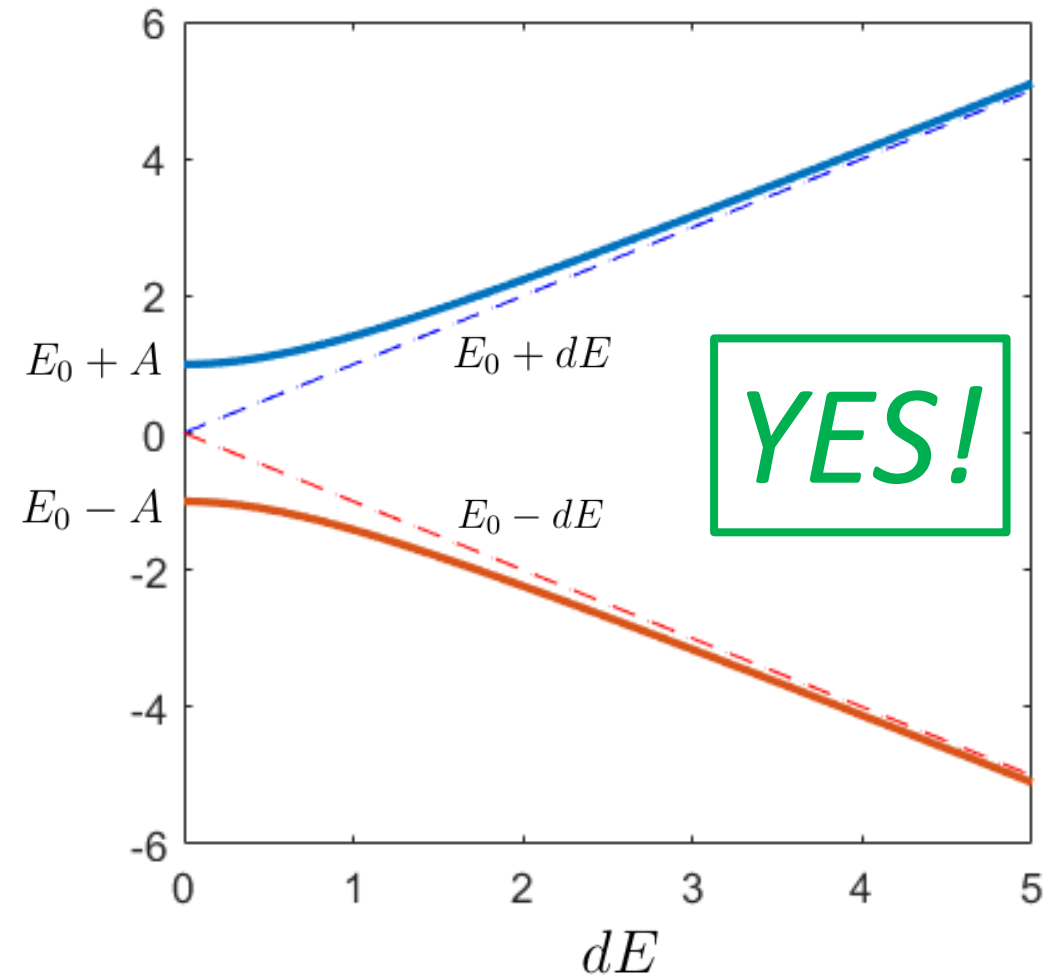


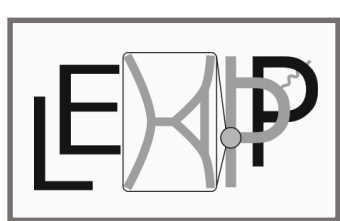
Is it different from a molecular dipole?



The *energy* eigenstates are:

$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$





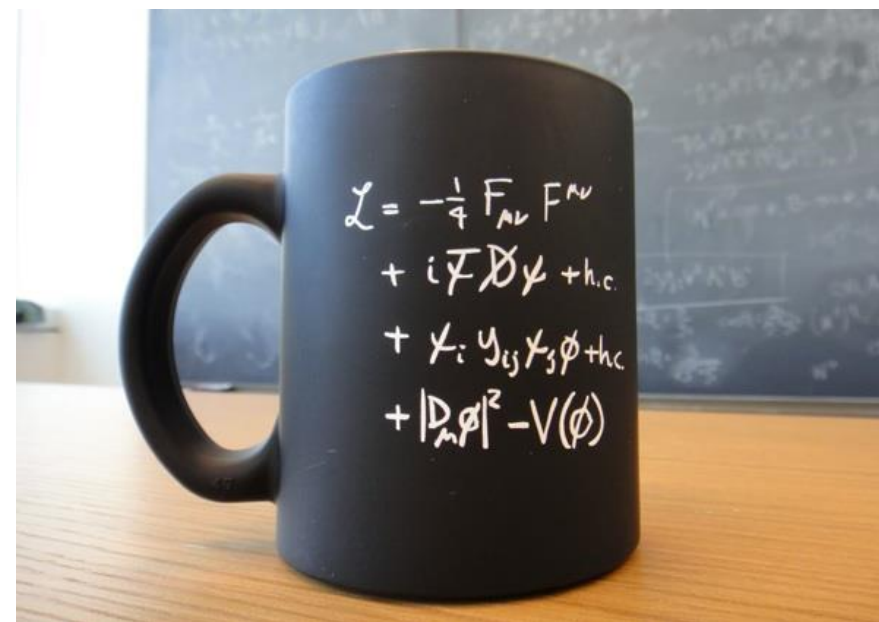
New Physics, in Familiar Terms

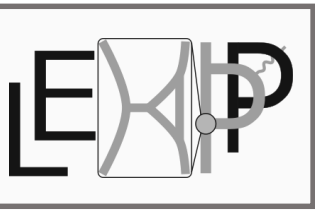
$$\mathcal{L}_{\text{fermion}} = -\frac{\mu}{2}\bar{\psi}\sigma^{\mu\nu}F_{\mu\nu}\psi - i\frac{d}{2}\bar{\psi}\sigma^{\mu\nu}\gamma^5 F_{\mu\nu}\psi$$

↓
MDM

↓
EDM

- Non-conservation of P and T already apparent in EDM term
- Consistency with zero vs. consistency with SM





A Taxonomy of Form Factors*

*which are not just for composite particles!

$$\mathcal{L}_{\text{fermion}} = -\frac{\mu}{2}\bar{\psi}\sigma^{\mu\nu}F_{\mu\nu}\psi - i\frac{d}{2}\bar{\psi}\sigma^{\mu\nu}\gamma^5F_{\mu\nu}\psi$$

↓
MDM

↓
EDM

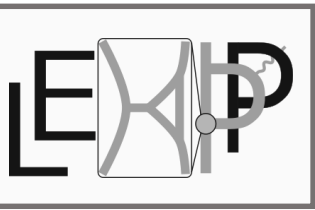
$$\begin{aligned} \langle p_f | j^\mu | p_i \rangle = & \bar{u}(p_f) \left[F_1(q^2)\gamma^\mu \right. \\ & + \frac{i\sigma^{\mu\nu}}{2m}q_\nu F_2(q^2) \\ & + i\epsilon^{\mu\nu\rho\sigma}\sigma_{\rho\sigma}q_\nu F_3(q^2) \\ & \left. + \frac{1}{2m}\left(q^\mu - \frac{q^2}{2m}\gamma^\mu\right)\gamma_5 F_4(q^2) \right] u(p_i) \end{aligned}$$

$$d = -\frac{F_3(0)}{2m}$$

$$Q = F_1(0)$$

$$\mu = \frac{F_1(0) + F_2(0)}{2m}$$

$$a = F_4(0)$$



A Taxonomy of Form Factors

$$\mathcal{L}_{\text{fermion}} = \underbrace{-\frac{\mu}{2} \bar{\psi} \sigma^{\mu\nu} F_{\mu\nu} \psi}_{\text{MDM}} \underbrace{- i \frac{d}{2} \bar{\psi} \sigma^{\mu\nu} \gamma^5 F_{\mu\nu} \psi}_{\text{EDM}}$$

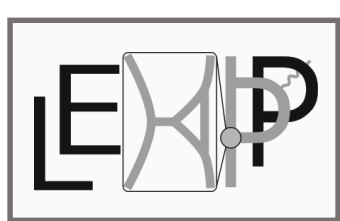
$$\langle p_f | j^\mu | p_i \rangle = \bar{u}(p_f) \left[\begin{aligned} & F_1(q^2) \gamma^\mu \\ & + \frac{i \sigma^{\mu\nu}}{2m} q_\nu F_2(q^2) \\ & + i \epsilon^{\mu\nu\rho\sigma} \sigma_{\rho\sigma} q_\nu F_3(q^2) \\ & + \frac{1}{2m} \left(q^\mu - \frac{q^2}{2m} \gamma^\mu \right) \gamma_5 F_4(q^2) \end{aligned} \right] u(p_i)$$

$$d = -\frac{F_3(0)}{2m}$$

$$\mu = \frac{F_1(0) + F_2(0)}{2m}$$

$$Q = F_1(0)$$

$$a = F_4(0)$$



Summary of Motivation

Problem: 1 extra baryon in every 10^9

- Asymmetry: $\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma}$
- n_γ comes from CMB decoupling
- Actually normalize to entropy density s , since universe expands

Requirements: Sakharov's criteria

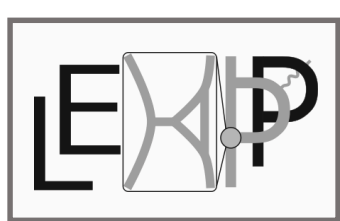
- Baryon-number (B) violation
- C and CP -violation
- Departure from thermal equilibrium

Solution: ???

- No complex antimatter nuclei
- No annihilation fronts
- No adequate symmetry-breaking

Prediction:

- New CP -violating physics
- Coupling to Standard Model baryons
- Polarization of bound states

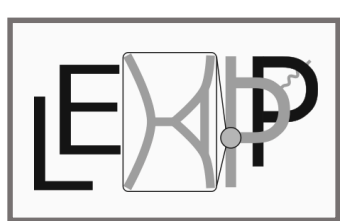


Analysis: Dimensional vs. Global

Naïve estimate for generic new physics:

$$d_n \propto \frac{m_q}{\Lambda^2} \cdot e \cdot \phi_{\text{CPV}}$$

Current experiments: $10^{-26} e \text{ cm} \rightarrow \Lambda \sim 10 - 100 \text{ TeV}$



Analysis: Dimensional vs. Global

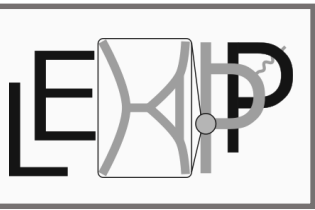
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Current experiments: $10^{-26} e \text{ cm}$ $\rightarrow \Lambda \sim 10 - 100 \text{ TeV}$

Standard Model CKM: $10^{-32} e \text{ cm}$

Standard Model QCD: ??? $\rightarrow d_n \approx (10^{-16} e \text{ cm}) \bar{\theta}$



Analysis: Dimensional vs. Global

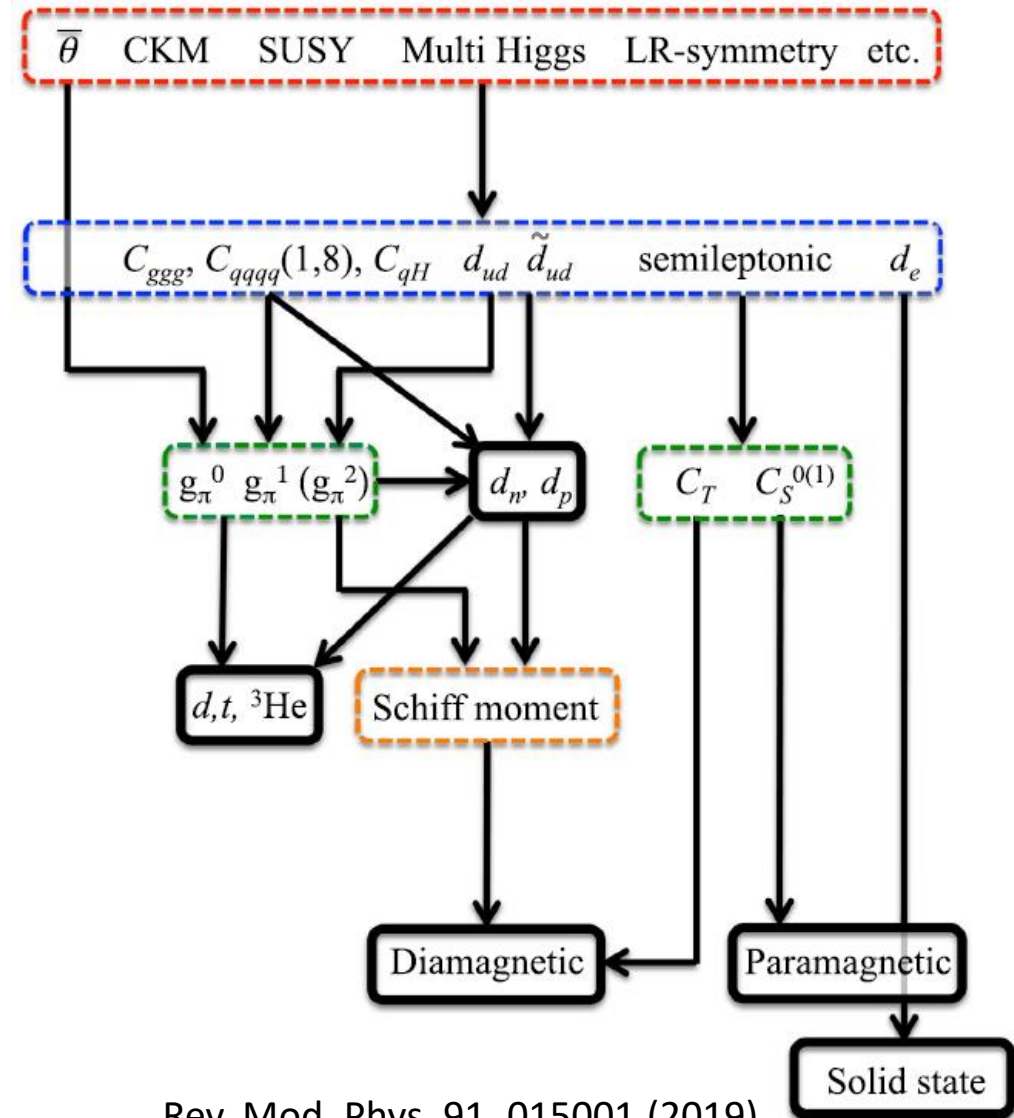
Naïve estimate for generic new physics:

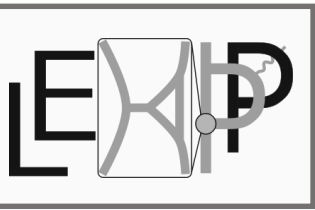
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Analysis: Dimensional vs. Global

Naïve estimate for generic new physics:

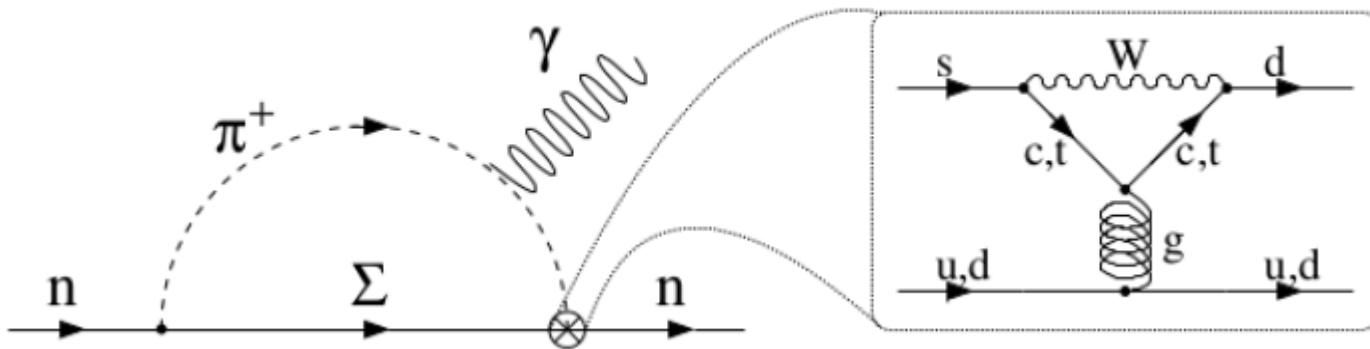
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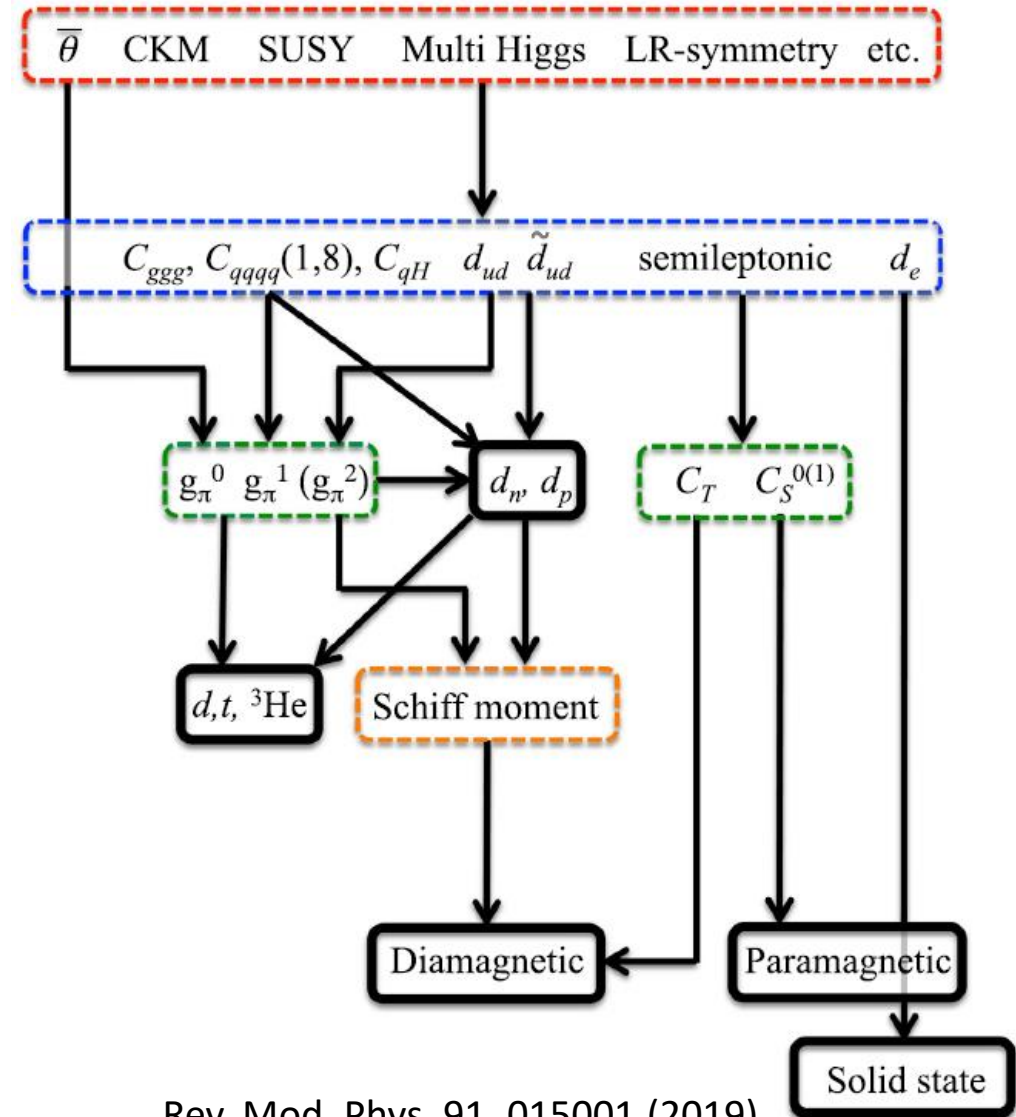
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Standard Model QCD: ???

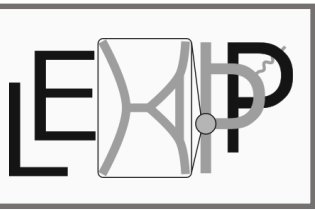
Neutron EDM from CP-violating pion couplings:



Pospelov & Ritz, *Annals of Physics* 318 (2005): 119-169



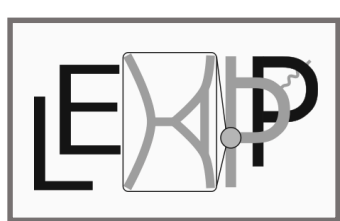
Rev. Mod. Phys. 91, 015001 (2019)



Many Parameters / Many Experiments

Sensitivity: System:	Paramagnetic	Diamagnetic	"Particle"
Trap	Tl, Cs, PbO, HfF ⁺ , Fr, BaF, ...	¹⁹⁹ Hg, ¹²⁹ Xe, ²²⁵ Ra, Rn, Pa, RaO, ...	n (ultra-cold)
Beam	YbF, ThO, WC	TlF	n (cold)
Storage ring	TaO ⁺	?	p, d, ³ He ⁺⁺ , μ, ...

Other: solid state (Gd₃Ga₅O₁₂, Eu_{0.5}Ba_{0.5}TiO₃), colliders (τ, Λ, ν, ...), crystal (n scattering on quartz), ...



Un-natural Units; Orders of Magnitude

$$10^{-26} e \text{ cm} \times \frac{1 \text{ MV}}{m} \times \frac{1}{2\pi\hbar} = 24 \text{ nHz}$$

$$\frac{1}{24 \text{ hours}} = 11.6 \mu\text{Hz}$$

$$\frac{1}{15 \text{ min}} = 1 \text{ mHz}$$

$$\mu_N \times \frac{1\mu\text{T}}{2\pi\hbar} = 8 \text{ Hz}$$

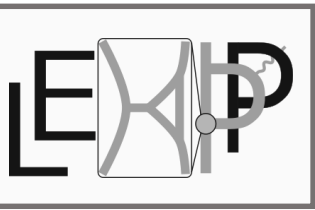
$$\mu_B \times \frac{1\mu\text{T}}{2\pi\hbar} = 14 \text{ kHz}$$

$$1 e \text{ cm} = 10^{13} e \text{ fm}$$

$$1 \text{ neV} = 1 \frac{\text{GeV}}{c^2} \times 1 \text{ cm} \times g$$

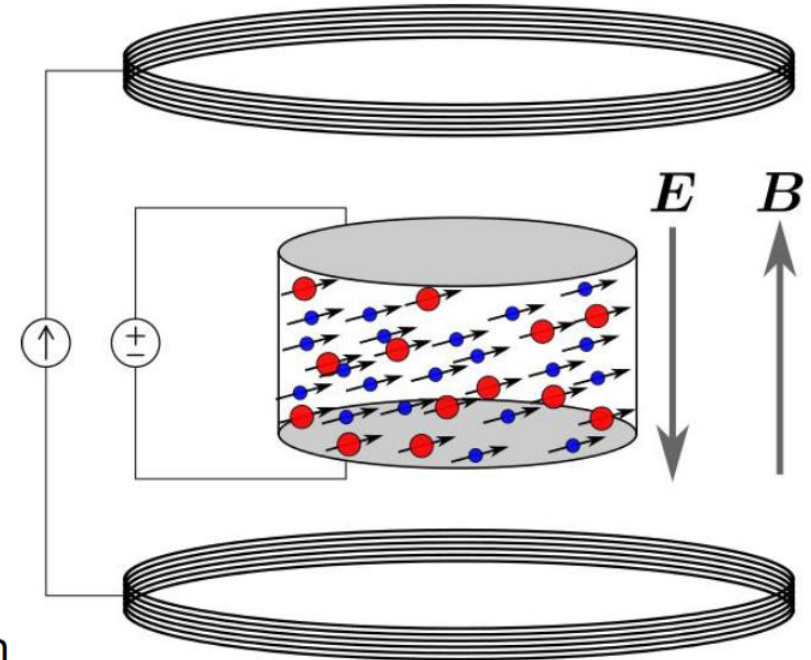
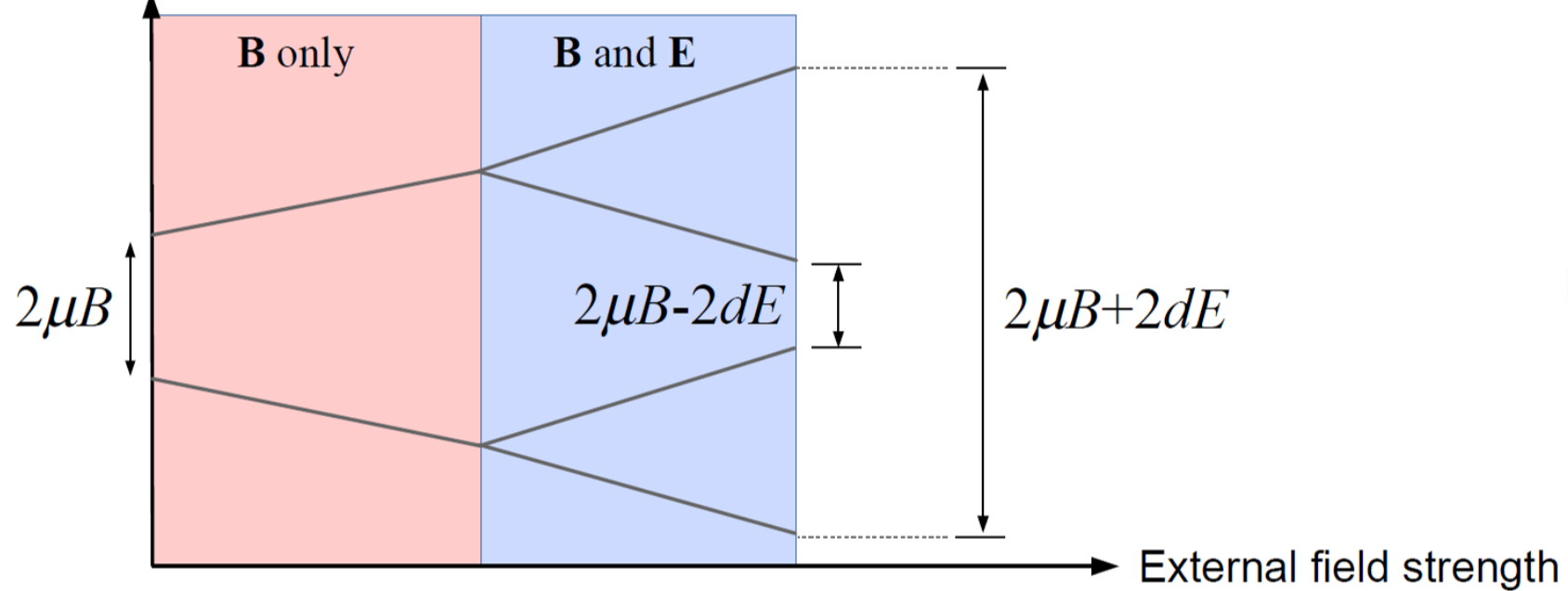
Some recent experimental EDM limits:

$$\begin{aligned} n: |d| &< 1.8 \times 10^{-26} e \text{ cm} \text{ (90\% C.L.)} \\ {}^{129}\text{Xe}: |d| &< 1.4 \times 10^{-27} e \text{ cm} \text{ (95\% C.L.)} \\ \text{ThO}: |d| &< 1.1 \times 10^{-29} e \text{ cm} \text{ (90\% C.L.)} \end{aligned}$$



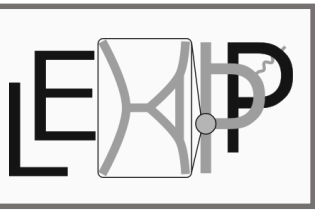
How could you measure an EDM?

$$H_{spin} = -\boldsymbol{\mu} \cdot \mathbf{B} - \mathbf{d} \cdot \mathbf{E}$$



$$\hbar(\omega_+ - \omega_-) = 4dE$$

...up to drift, gradients, etc.



Remember it is "locked" to the spin

Spin-precession based magnetometry:

- $E = -\boldsymbol{\mu} \cdot \mathbf{B}$
- $\boldsymbol{\tau} = \boldsymbol{\mu} \times \mathbf{B}$
- $\boldsymbol{\mu} = \gamma \mathbf{L} \rightarrow \boldsymbol{\omega}_L = -\gamma \mathbf{B}$

Time evolution from Bloch equations:

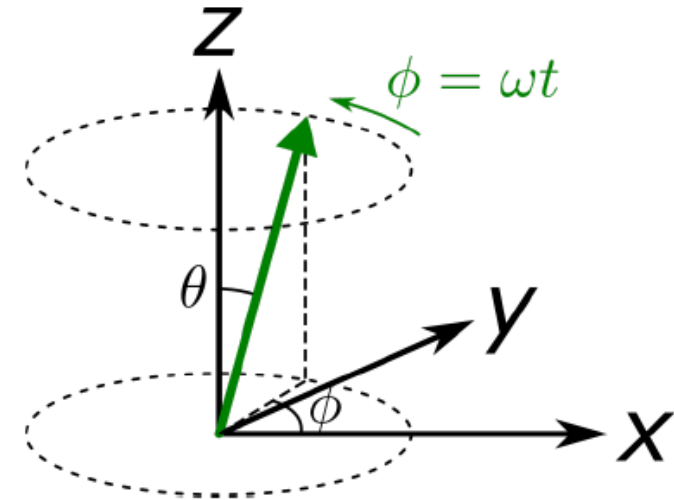
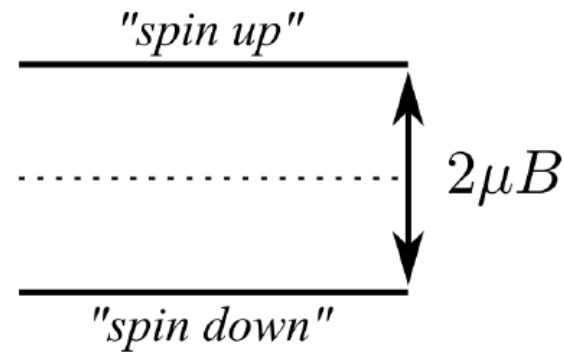
$$\frac{d\boldsymbol{\mu}}{dt} = \gamma \boldsymbol{\mu} \times \mathbf{B} - (\text{relaxation terms})$$

Sensitivity from: $\Delta E \Delta t \geq \hbar/2$

- relaxation limits measurement time
- many particles \rightarrow many measurements

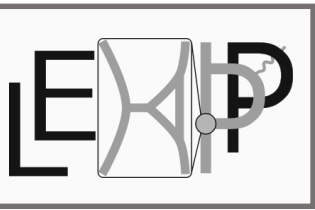
EDM fundamental sensitivity:

$$|\delta\omega| = \frac{|dE|}{\hbar F} \quad (\Delta m_F = 1)$$



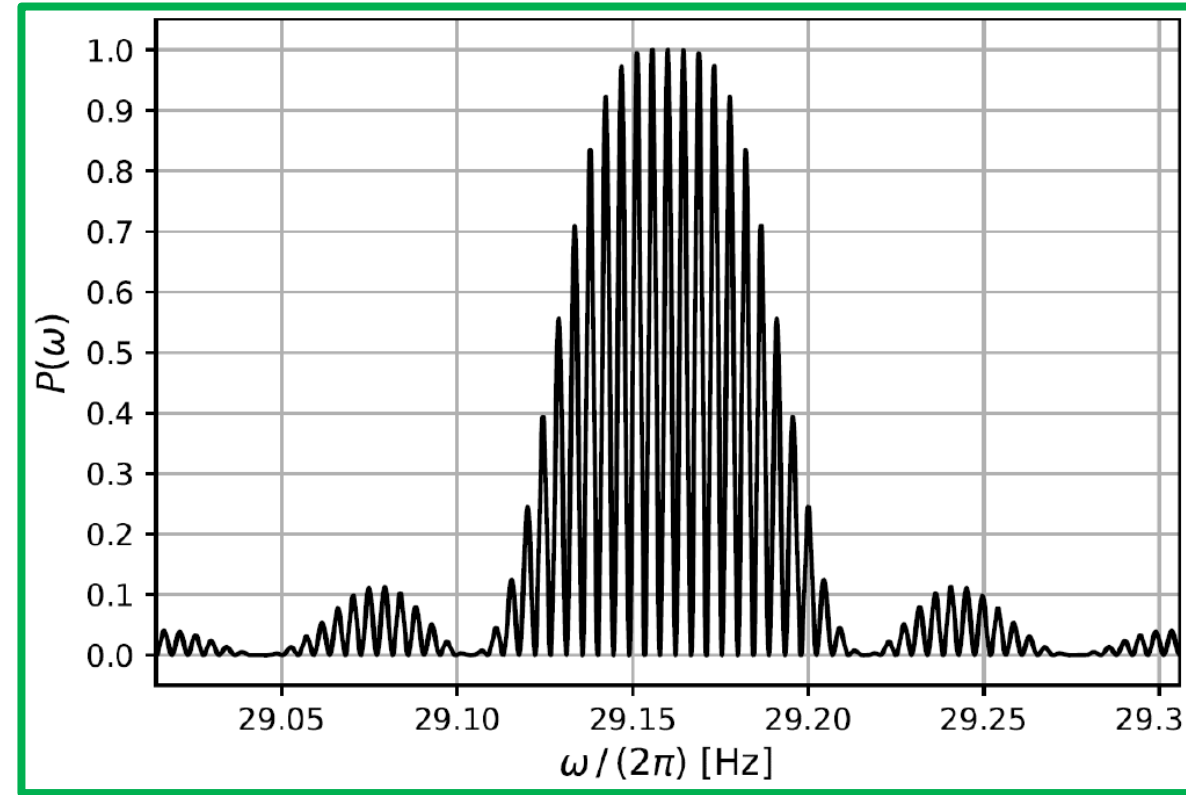
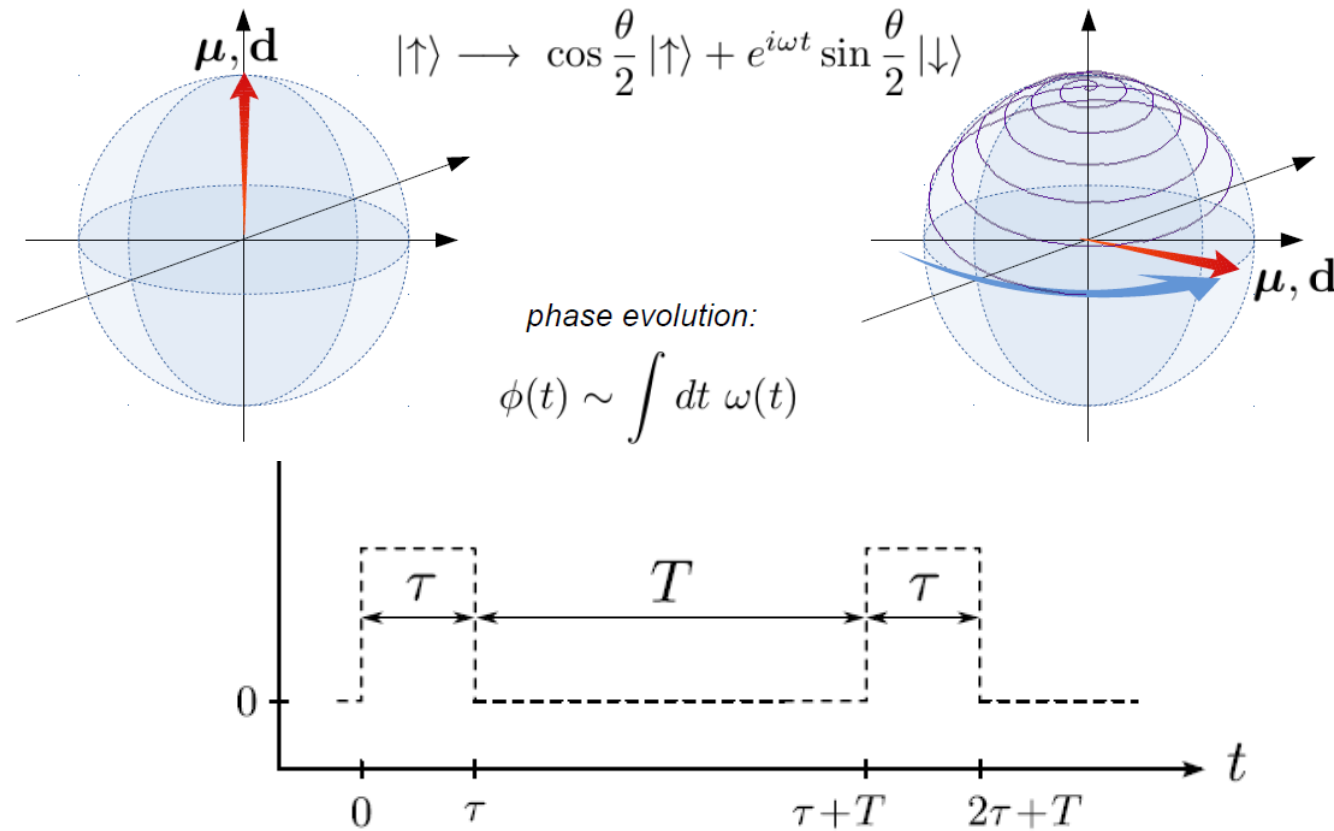
Cornell and Wieman... Nobel 2001, Rev. Mod. Phys. 74, 875 (2002)

vious initial step toward understanding dynamical behavior. Second, in experimental physics a precision measurement is almost always a frequency measurement, and the easiest way to study an effect with precision is to find an observable frequency that is sensitive to that effect. In the case of dilute-gas BEC, the observed fre-

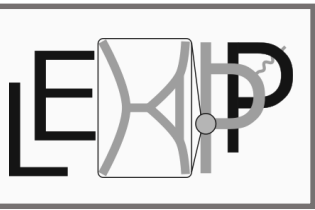


Time-Domain Interferometry

Ramsey's method to measure frequencies*:

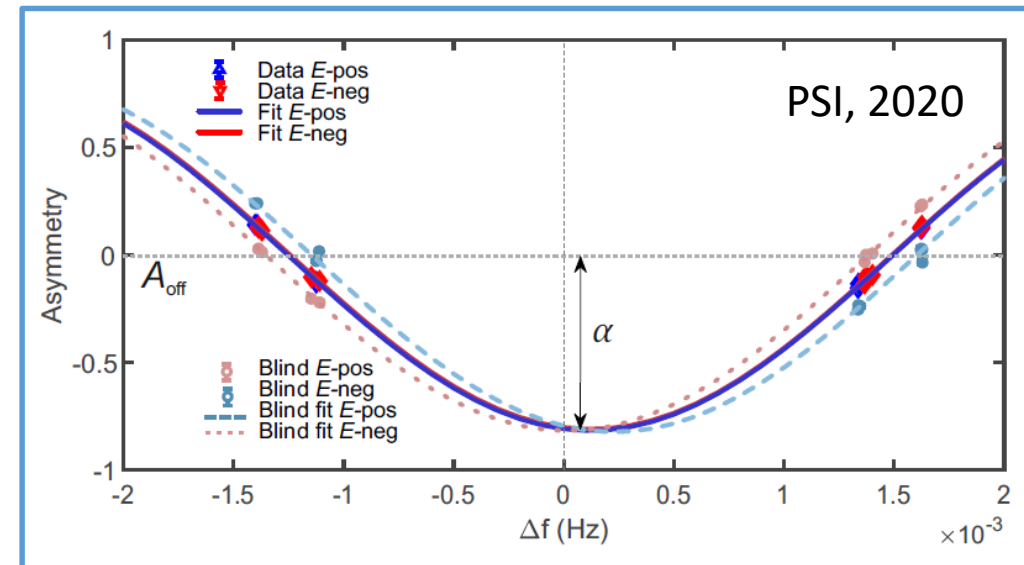
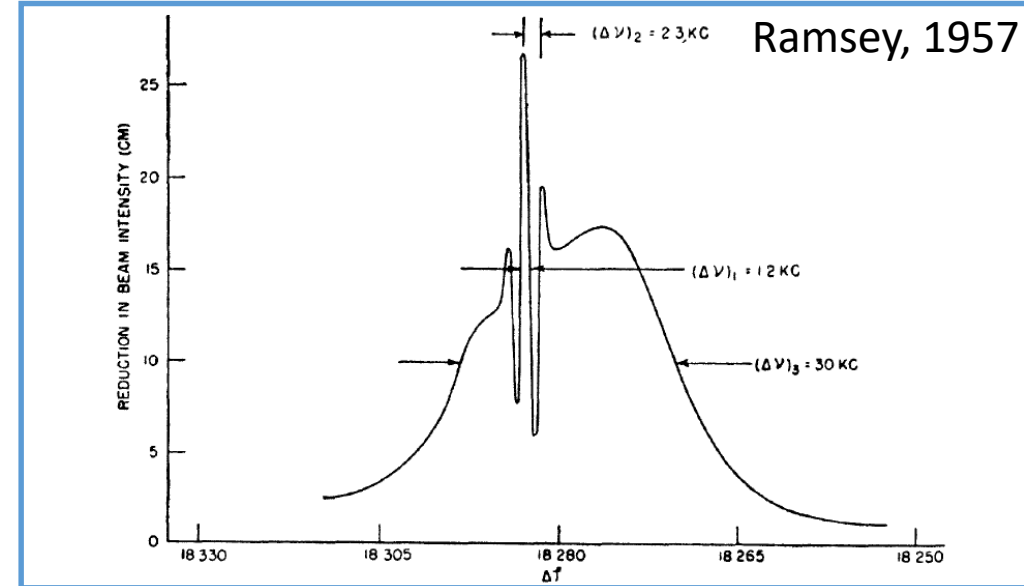
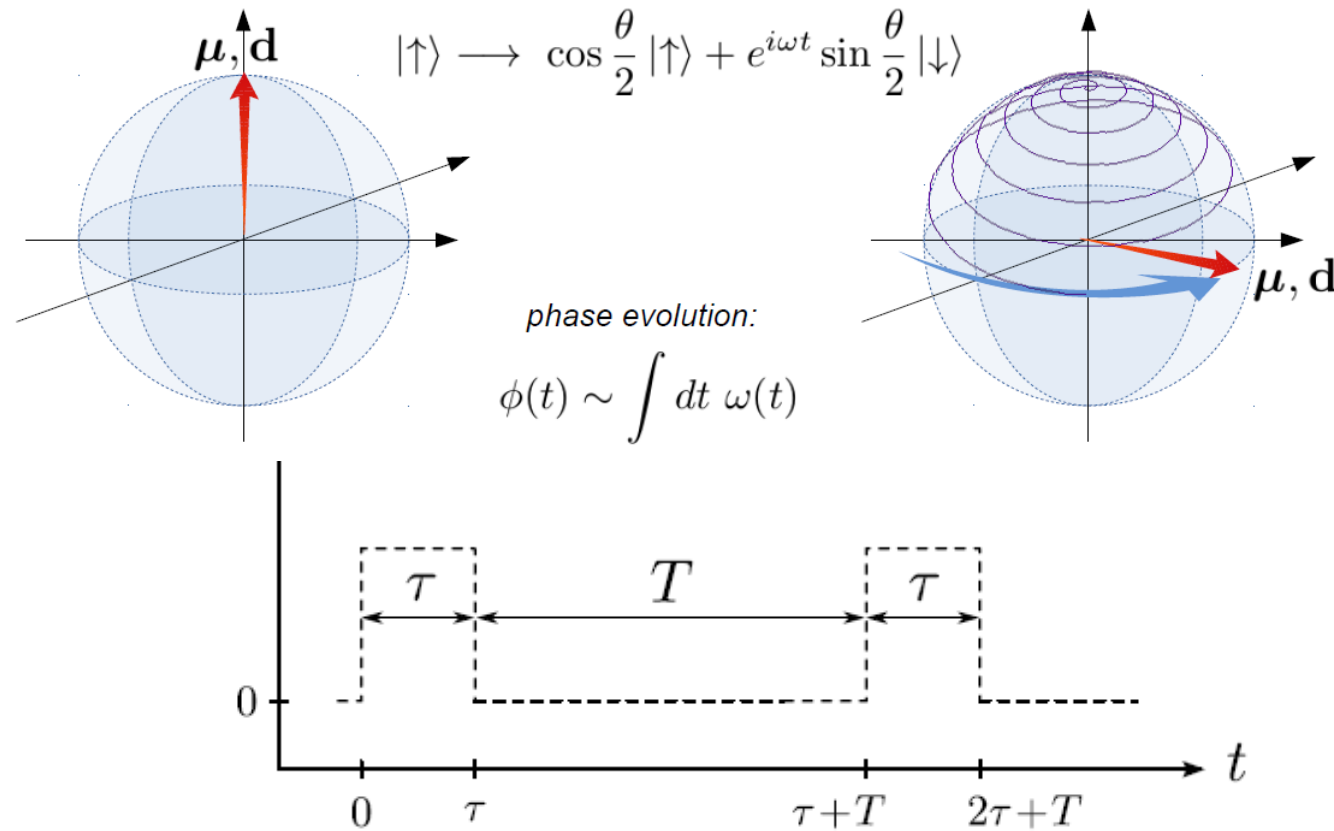


*we'll come back to *frequency vs. phase*

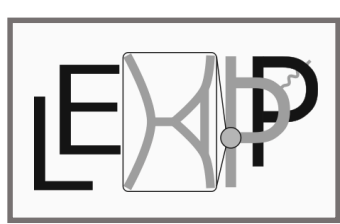


How could you measure an EDM?

Ramsey's method to measure frequencies*:

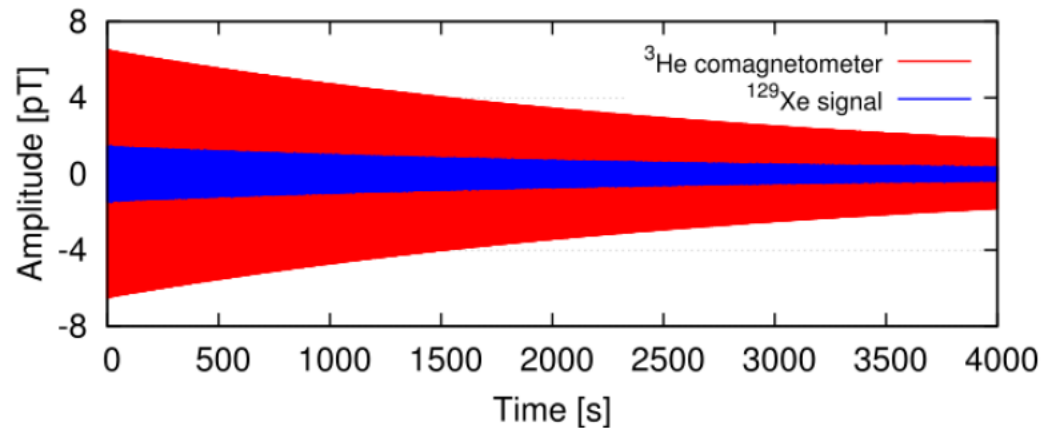
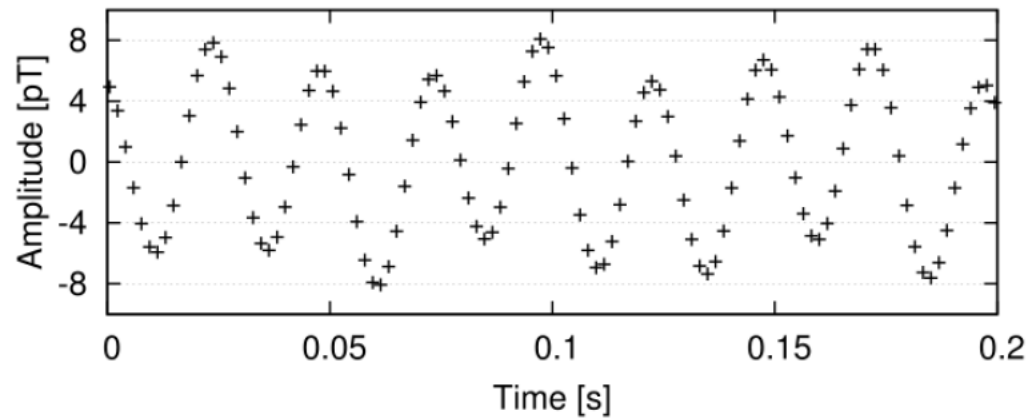


*we'll come back to *frequency vs. phase*



How could you measure an EDM?

What if we could measure continuously?



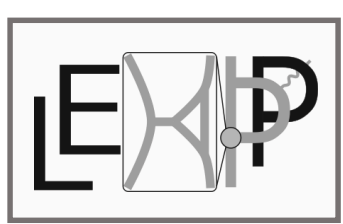
$$\delta d \sim \frac{h}{2ET_2} \frac{1}{S/N}$$

“phase noise” limit

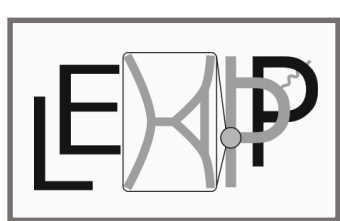
$$\frac{h}{2ET_2} \frac{1}{\sqrt{\phi_n T_2}}$$

$$\frac{h}{2ET_2} \frac{1}{\sqrt{N_n}}$$

“count rate” limit



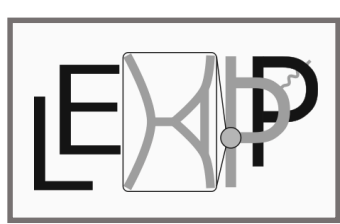
So which system should you measure?



So which system should you measure?

The one where you can discover an EDM, of course!

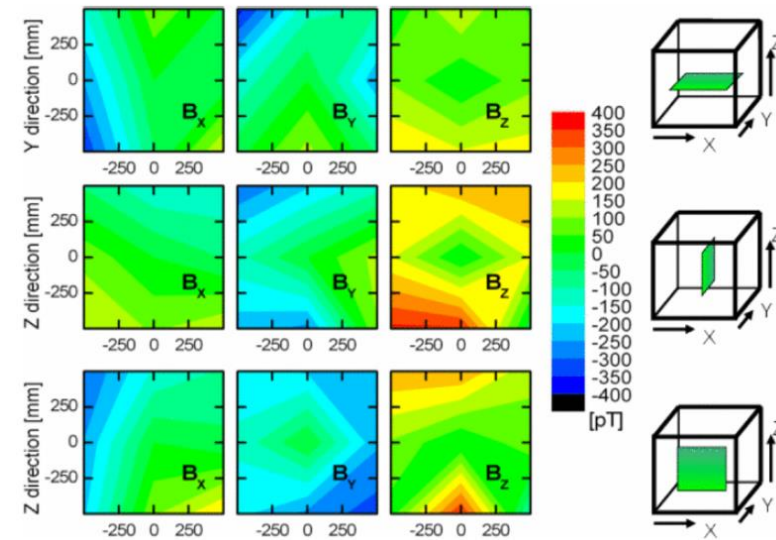
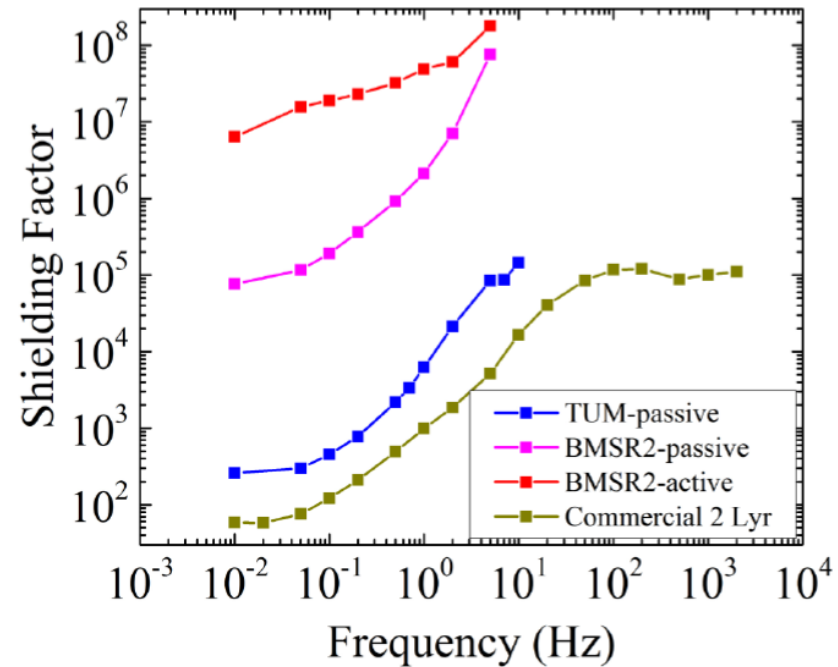
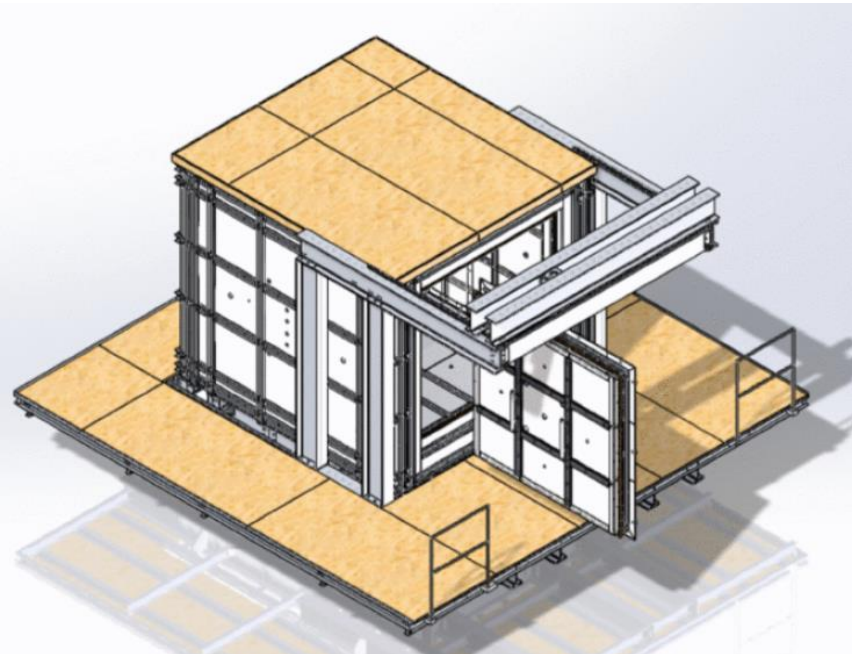
- M. Ramsey-Musolf



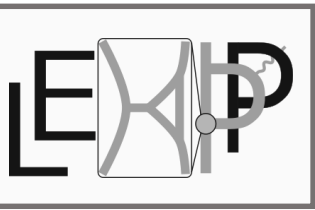
The HeXe Experiment



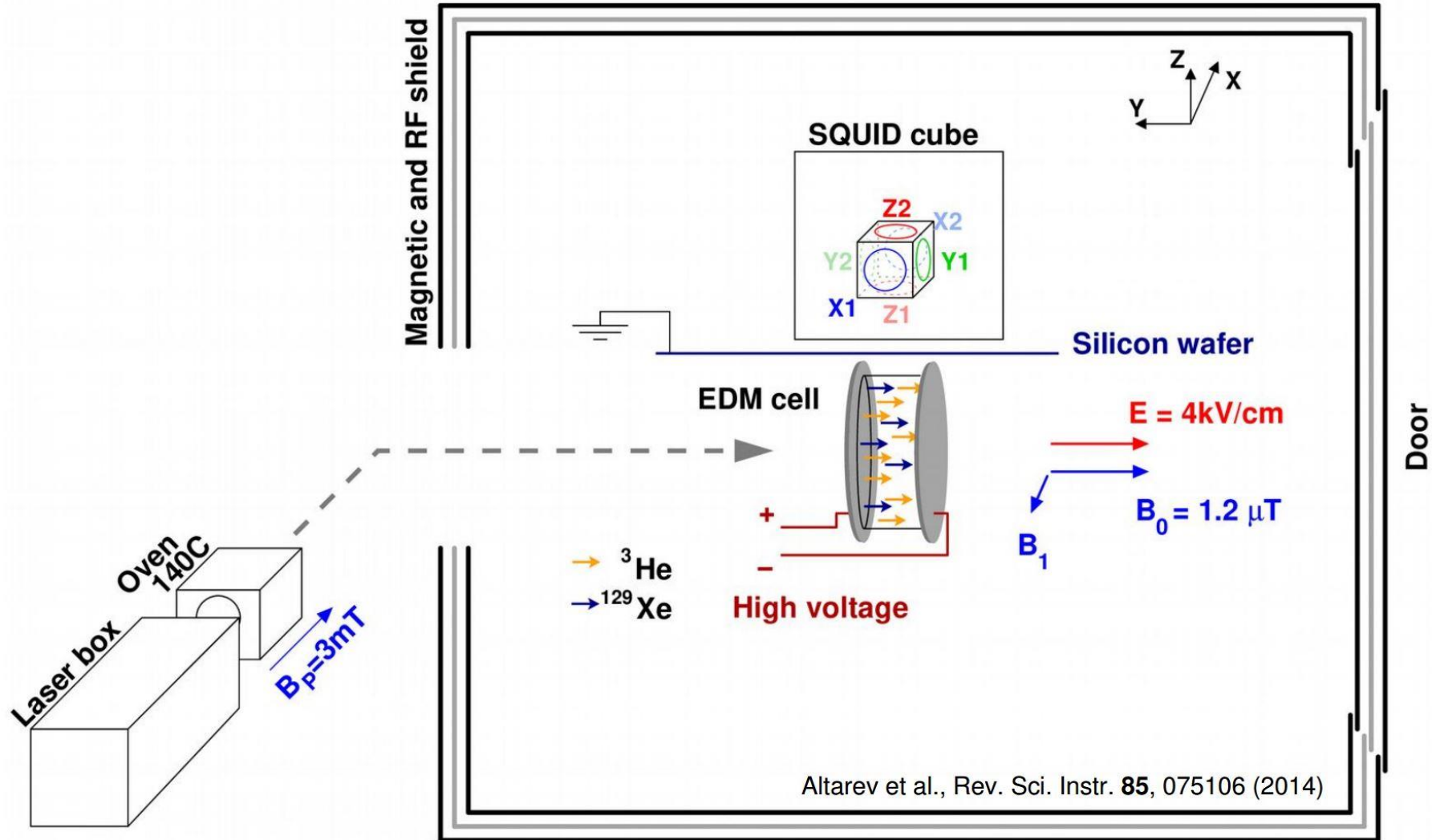
Use the best magnetic shields available (at least to start with...)

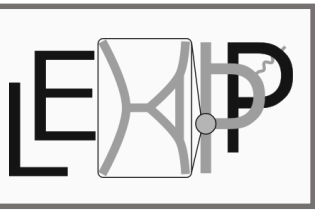


Rev. Sci. Instrum. **85**, 075106 (2014)
J. Appl. Phys. **117**, 183903 (2015)

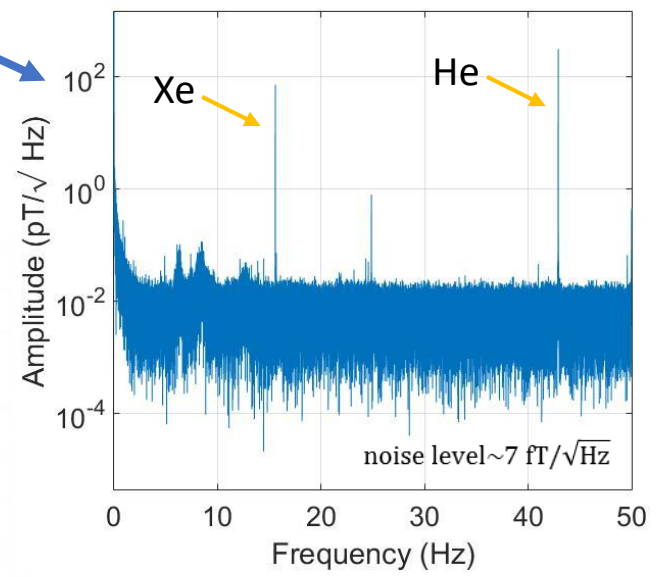
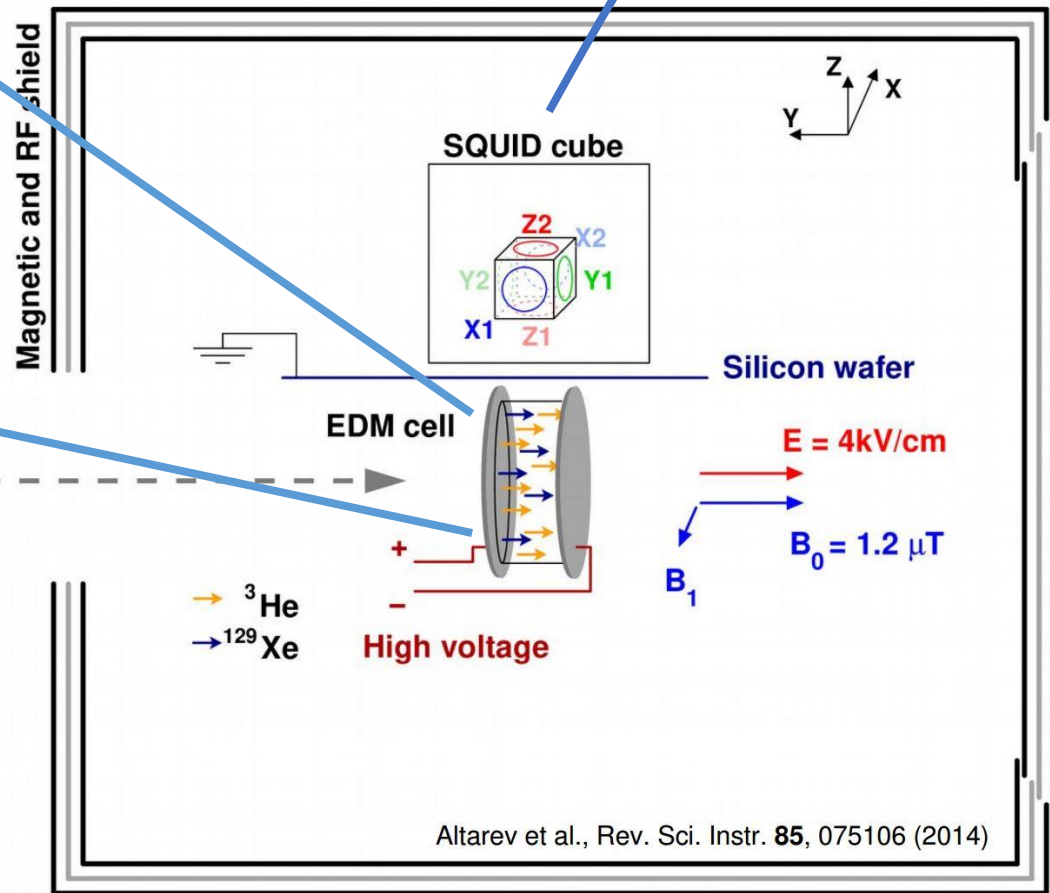
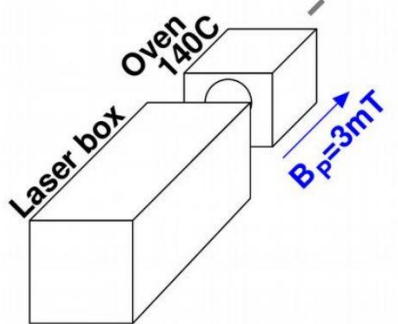
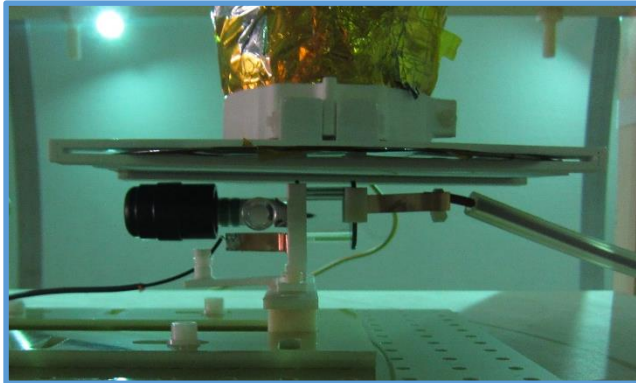


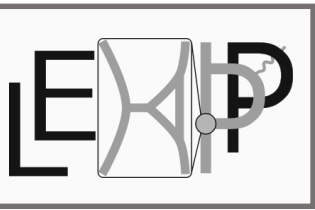
The HeXe Experiment





The HeXe Experiment

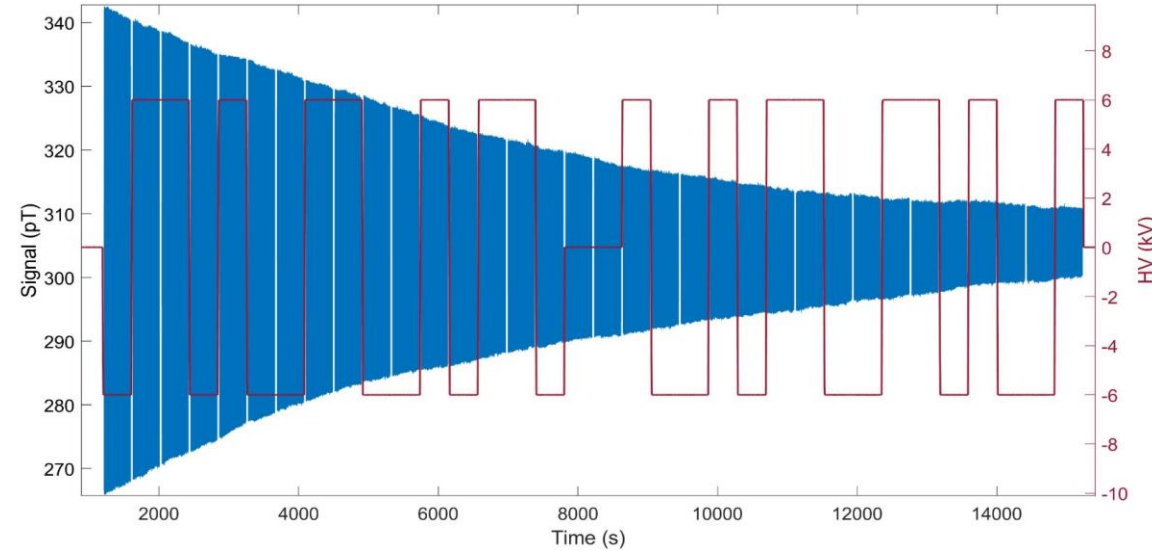
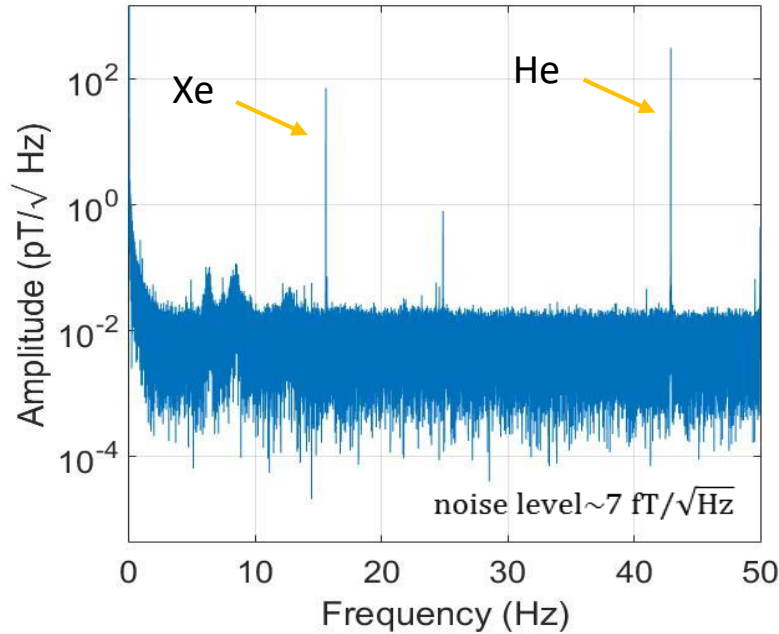




The HeXe Experiment



$E \sim 4 \text{ kV/cm}$
 $\tau \sim 4000 \text{ s}$
 $S \sim 20 \text{ pT}$
 $\epsilon \sim 8 \text{ fT}/\sqrt{\text{Hz}}$

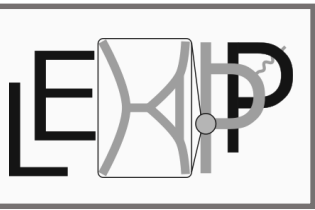


$$\delta\omega = \frac{1}{\tau(S/n)\sqrt{N}} = \frac{\epsilon\sqrt{f_{\text{BW}}}}{\tau S\sqrt{N}}$$

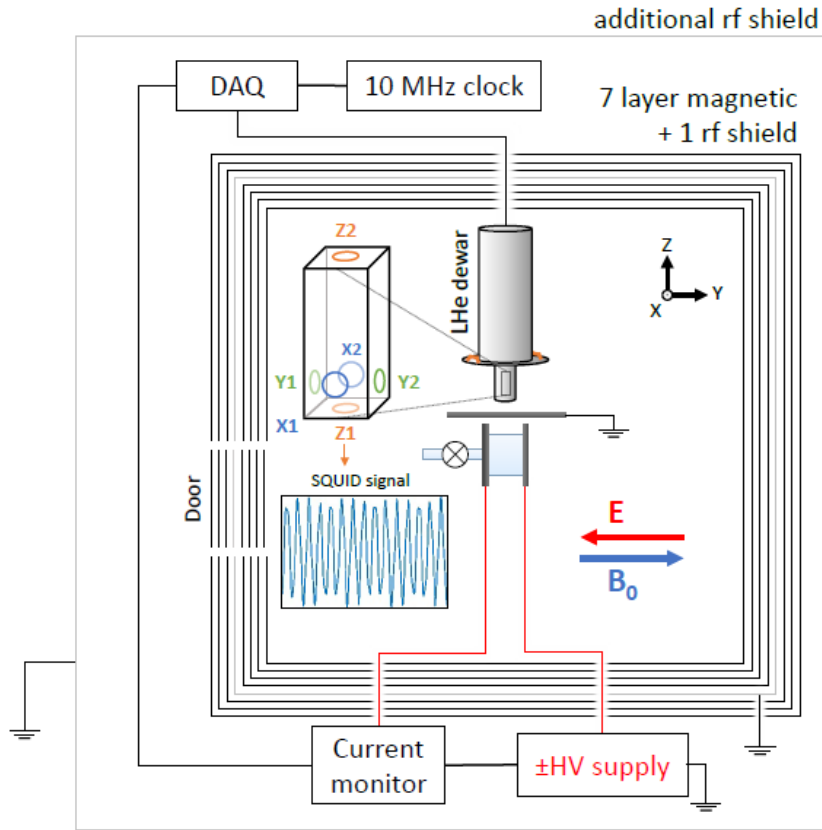
→ require nHz per run

$$\sigma_d = \frac{\hbar}{2E} \frac{\epsilon}{\tau^{3/2} S\sqrt{N}} = \frac{\hbar}{2E} \frac{\epsilon}{\tau S\sqrt{T}}$$

→ few $\times 10^{-27} \text{ e cm}/\sqrt{N}$



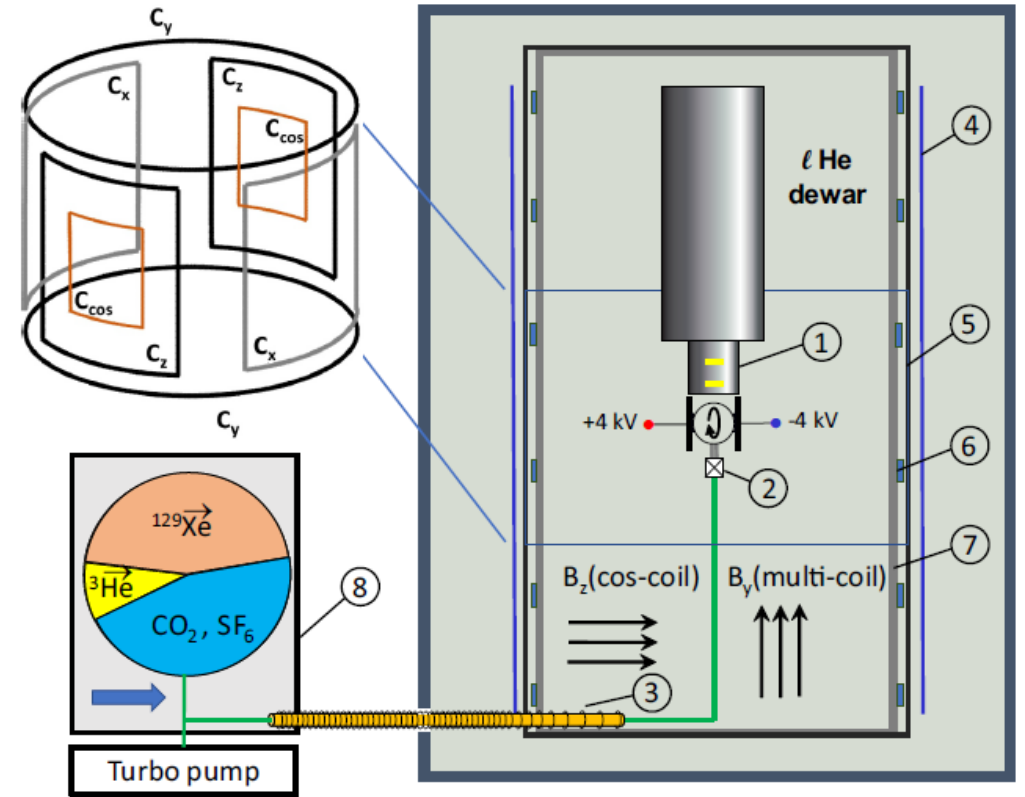
A rapidly-moving field!



Our result from HeXe:

$$d_A(^{129}\text{Xe}) = (1.4 \pm 6.6_{\text{stat}} \pm 2.0_{\text{syst}}) \times 10^{-28} \text{ e cm.}$$

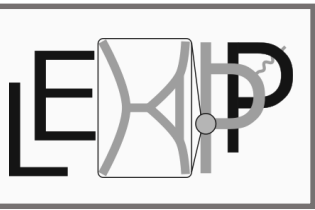
Phys. Rev. Lett. **123**, 143003 (2019)



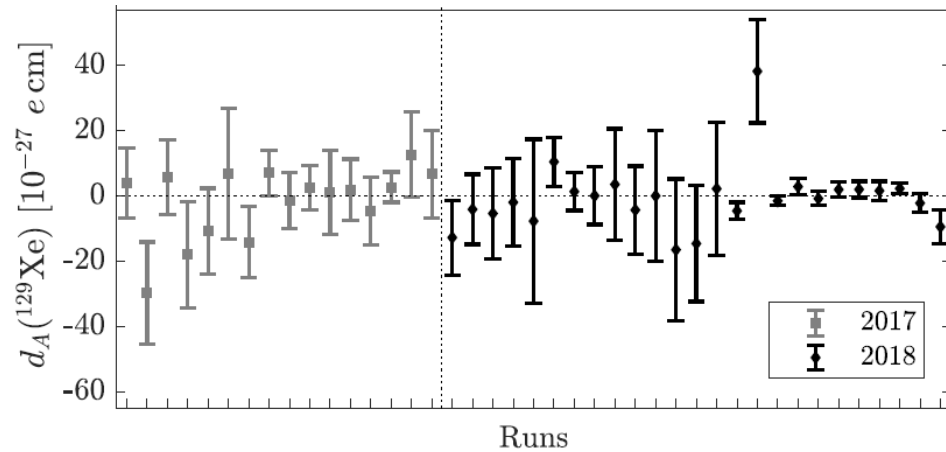
Near-simultaneous from MiXed:

$$d_{\text{Xe}} = (-4.7 \pm 6.4) \cdot 10^{-28} \text{ e cm}$$

Phys. Rev. A **100**, 022505 (2019)



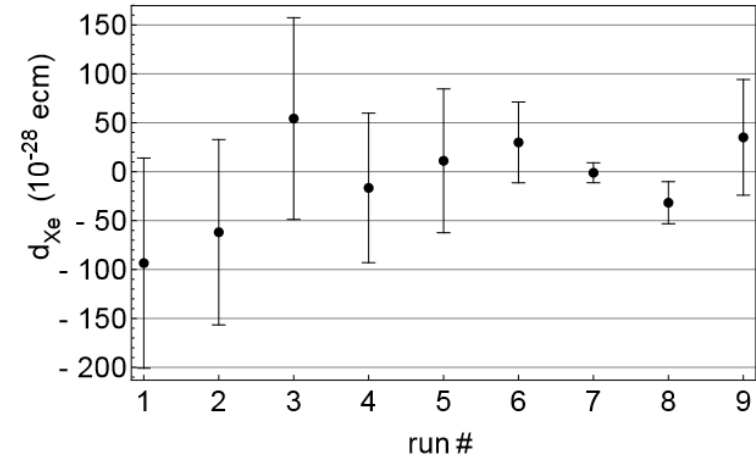
A rapidly-moving field!



Our result from HeXe:

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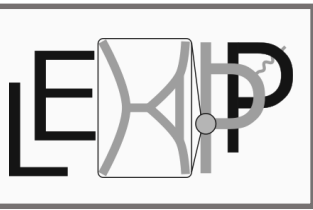
Phys. Rev. Lett. **123**, 143003 (2019)



Near-simultaneous from MiXed:

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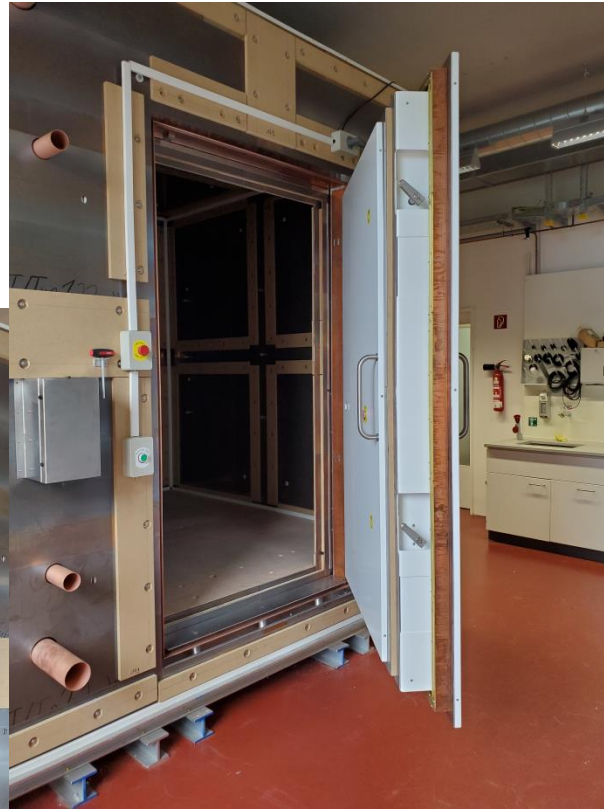
Phys. Rev. A **100**, 022505 (2019)



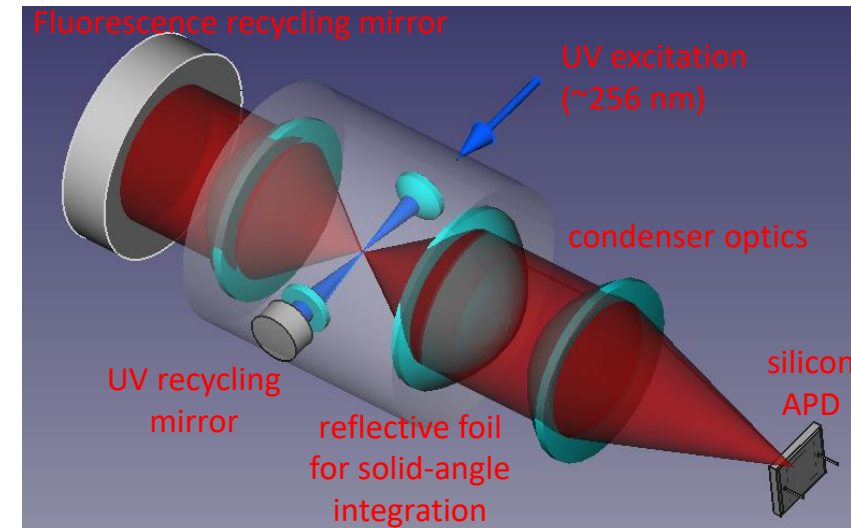
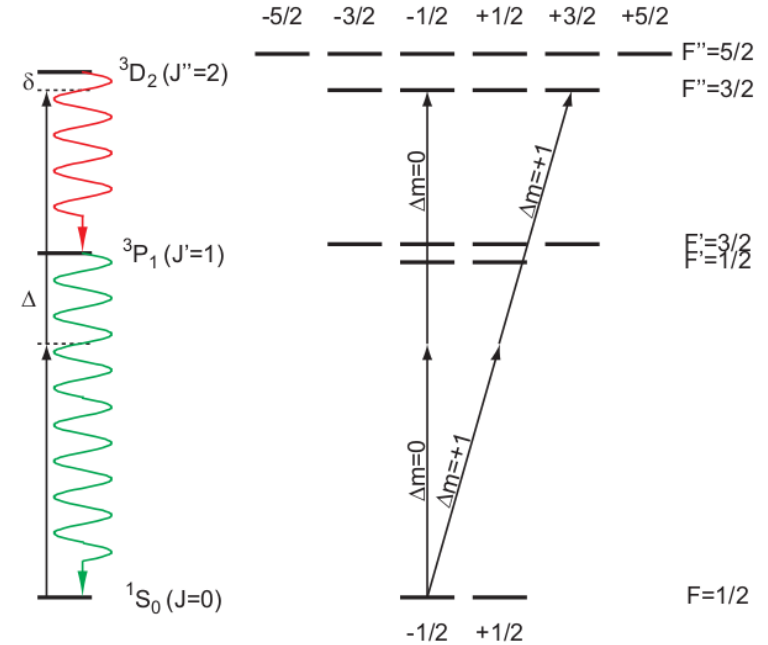
A United Future: collective experience



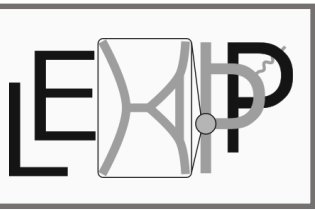
New magnetically shield room @HD!
Dedicated facility...



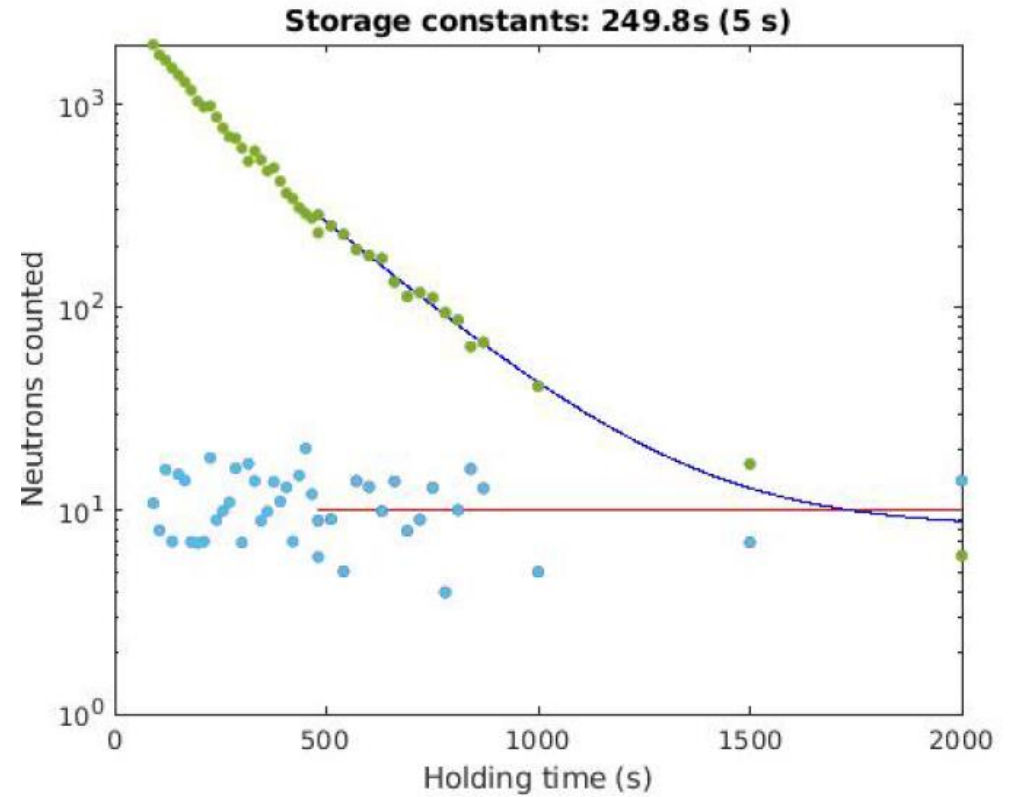
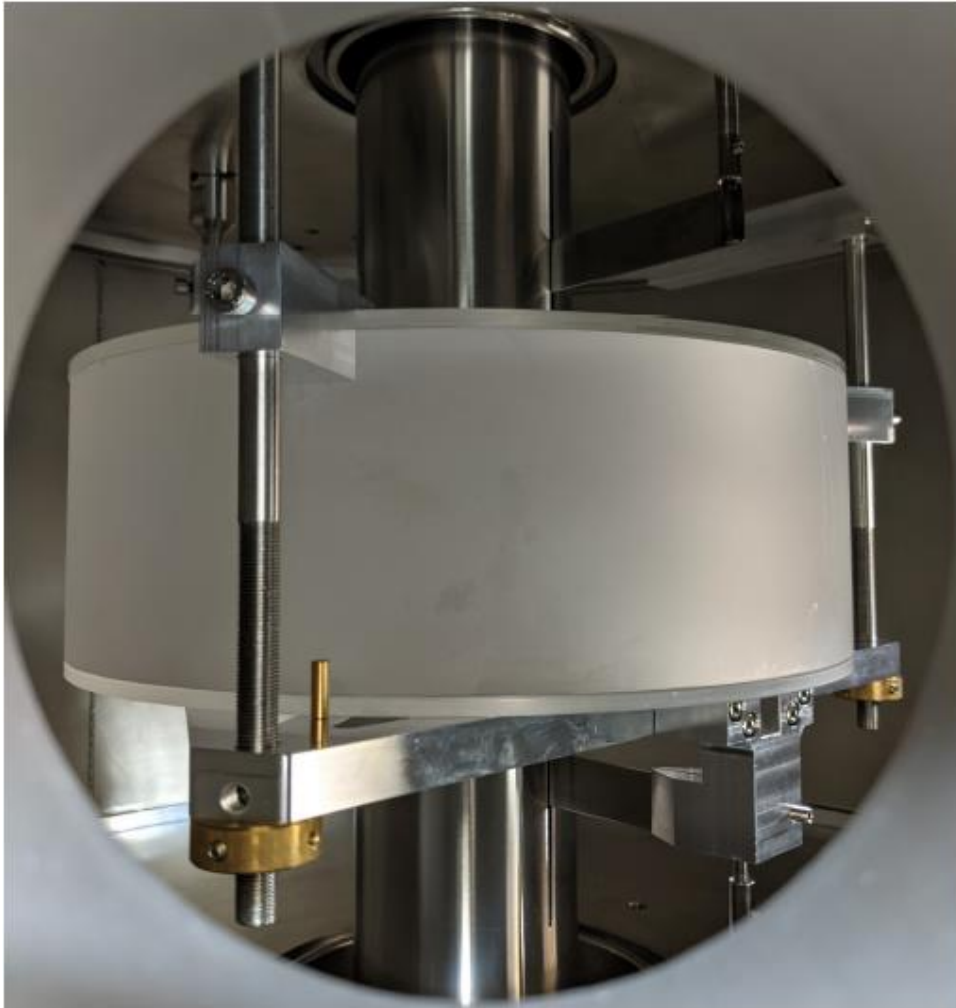
Laser spectroscopy may complement or eventually replace SQUIDS... new tools!



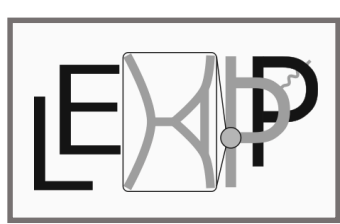
Next order-of-magnitude pursued by refining these now-known methods



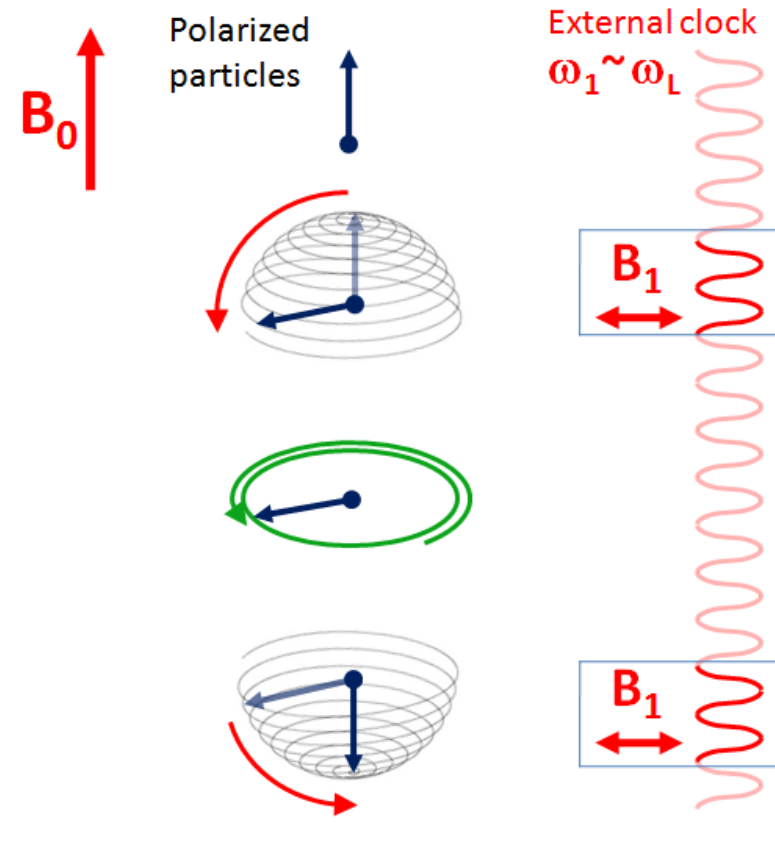
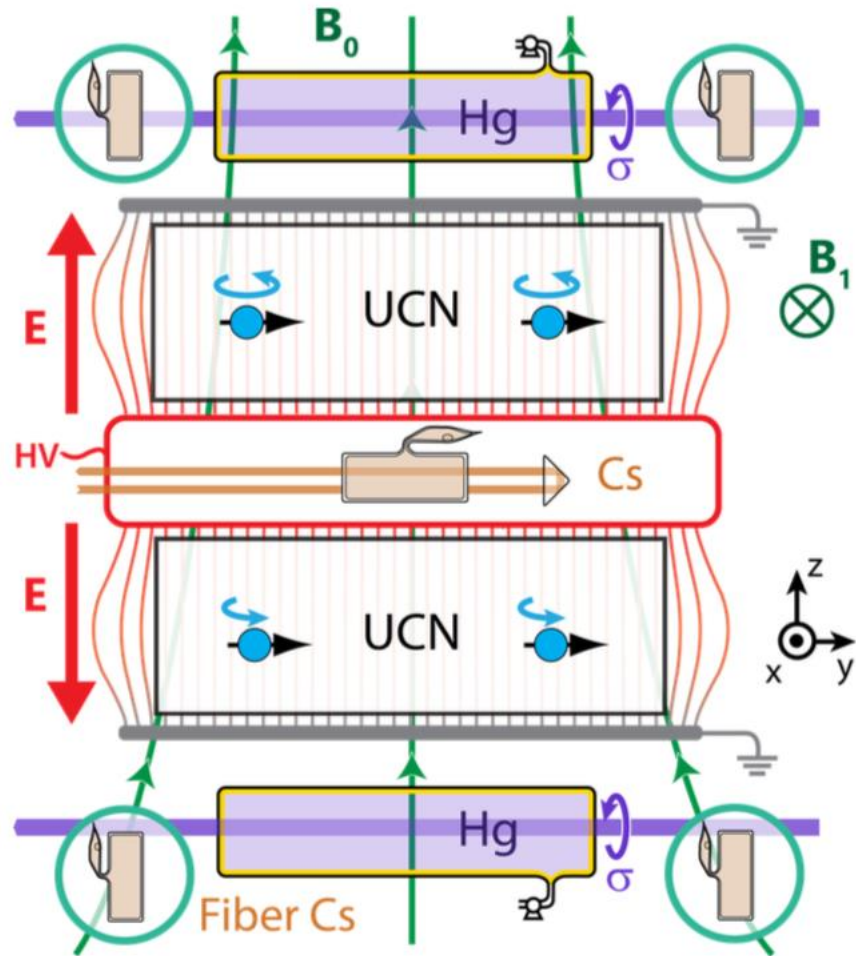
A contrast: neutrons disappear faster!

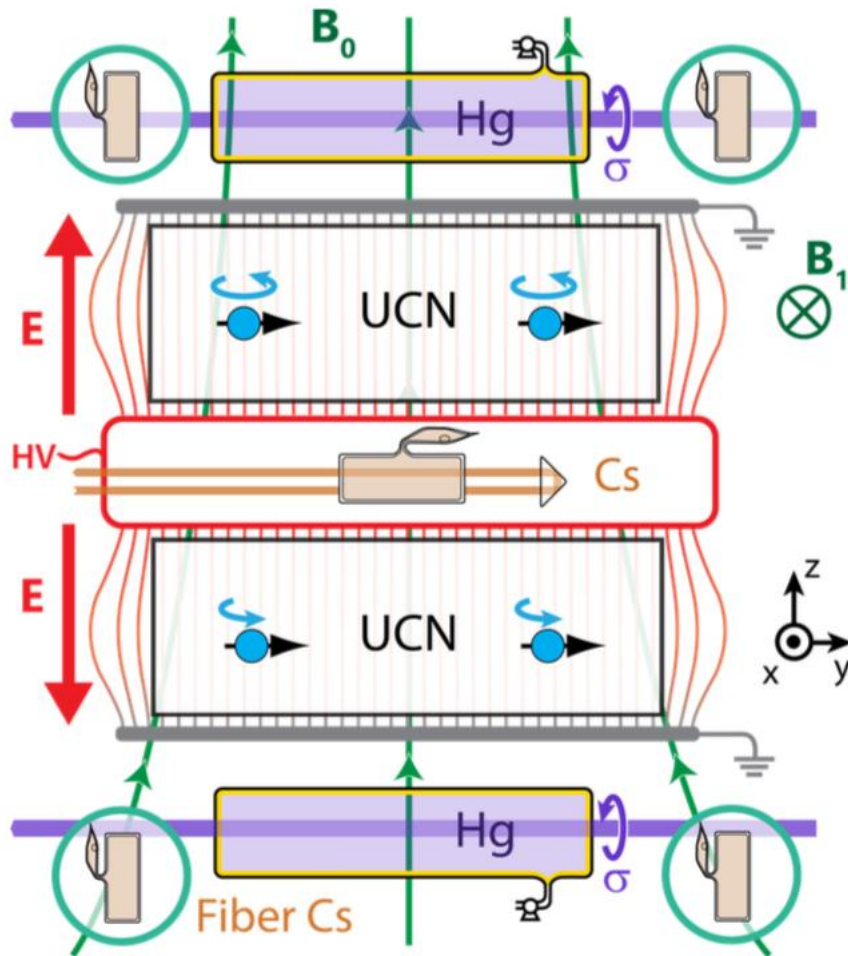


Magnetic fields would be good for loss... bad for systematics

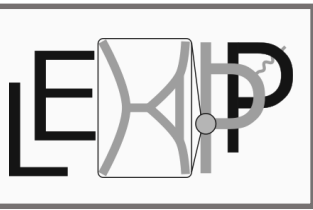


The PanEDM Experiment

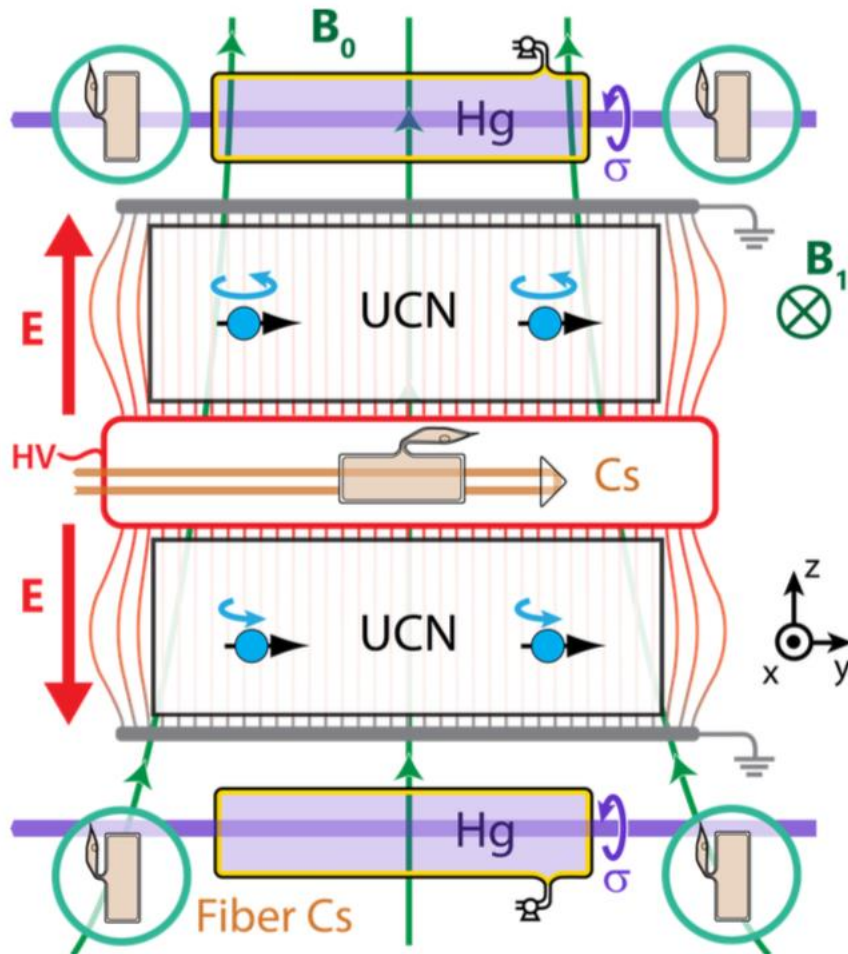




- Double chamber Ramsey interferometer at room temperature (but $T_{UCN} \sim 5\text{mK}$)
- ^{199}Hg magnetometers with few-fT resolution
- Cs magnetometers (also at high voltage)
- Magnetic shielding factor: 6×10^6 at 1 mHz
- Simultaneous spin detection for up/down
- SuperSUN UCN source at ILL in 2 phases:
 - Phase I: unpolarized UCN with 80 neV peak
 - Phase II: polarized UCN, magnetic storage
- Ongoing installation of parts, commissioning with UCN production in 2023-2024



Much lower statistics!



Statistical sensitivity:

$$\sigma(d_n) \gtrsim \frac{\hbar}{2\alpha|\mathbf{E}|T\sqrt{N}}$$

Frequency measurement:

$$|\delta\omega| = \frac{|dE|}{\hbar F}$$

SuperSUN

Phase I

Saturated source

density [cm^{-3}] 330

Diluted density [cm^{-3}] 63

Density in cells [cm^{-3}] 3.9

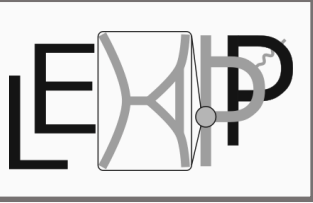
PanEDM Sensitivity [$1\sigma, e\text{ cm}$]

Per run 5.5×10^{-25}

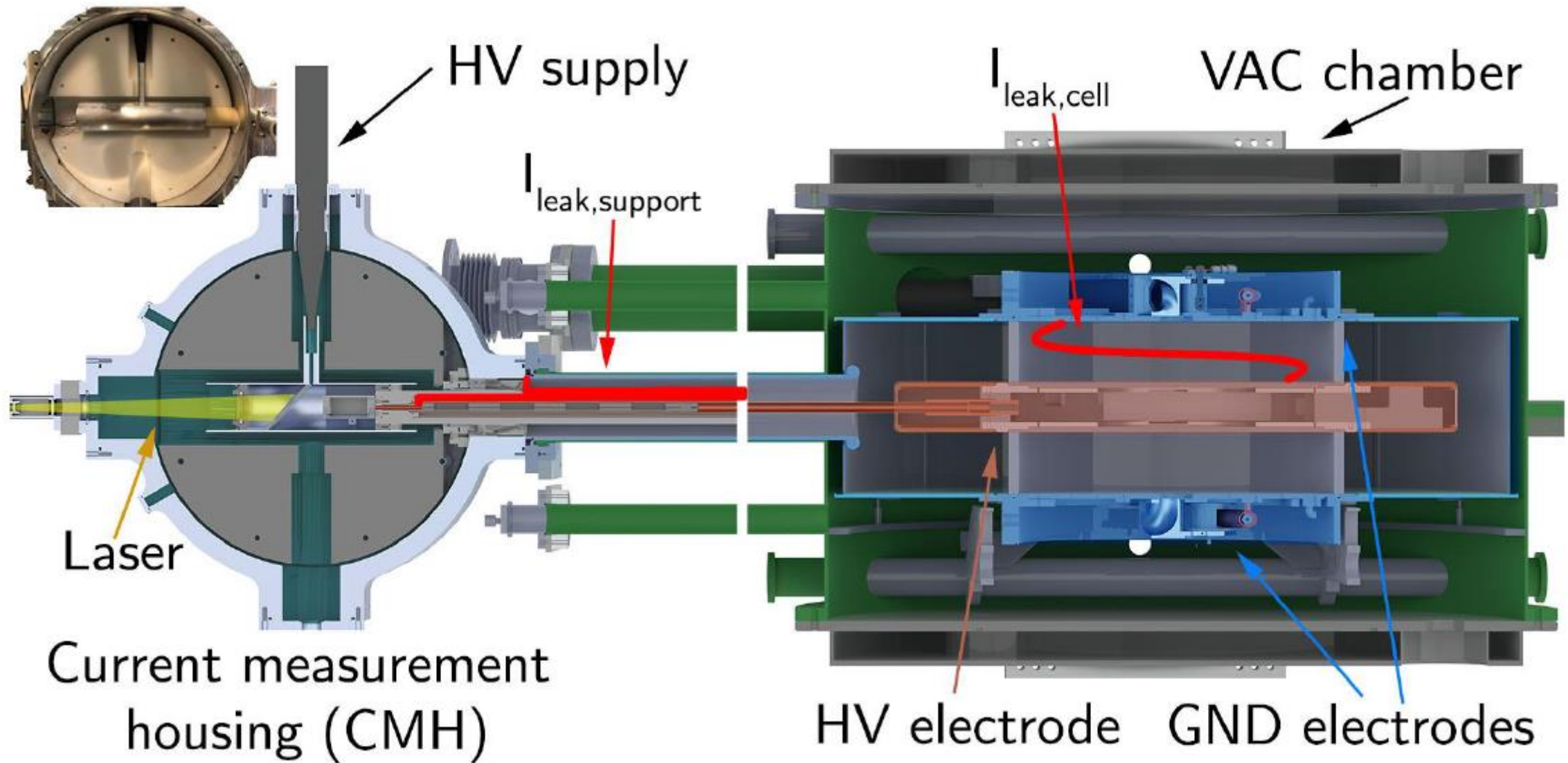
Per day 3.8×10^{-26}

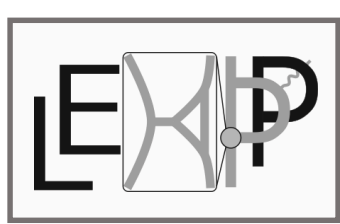
Per 100 days 3.8×10^{-27}

$$\Delta E \Delta t \geq \hbar/2$$

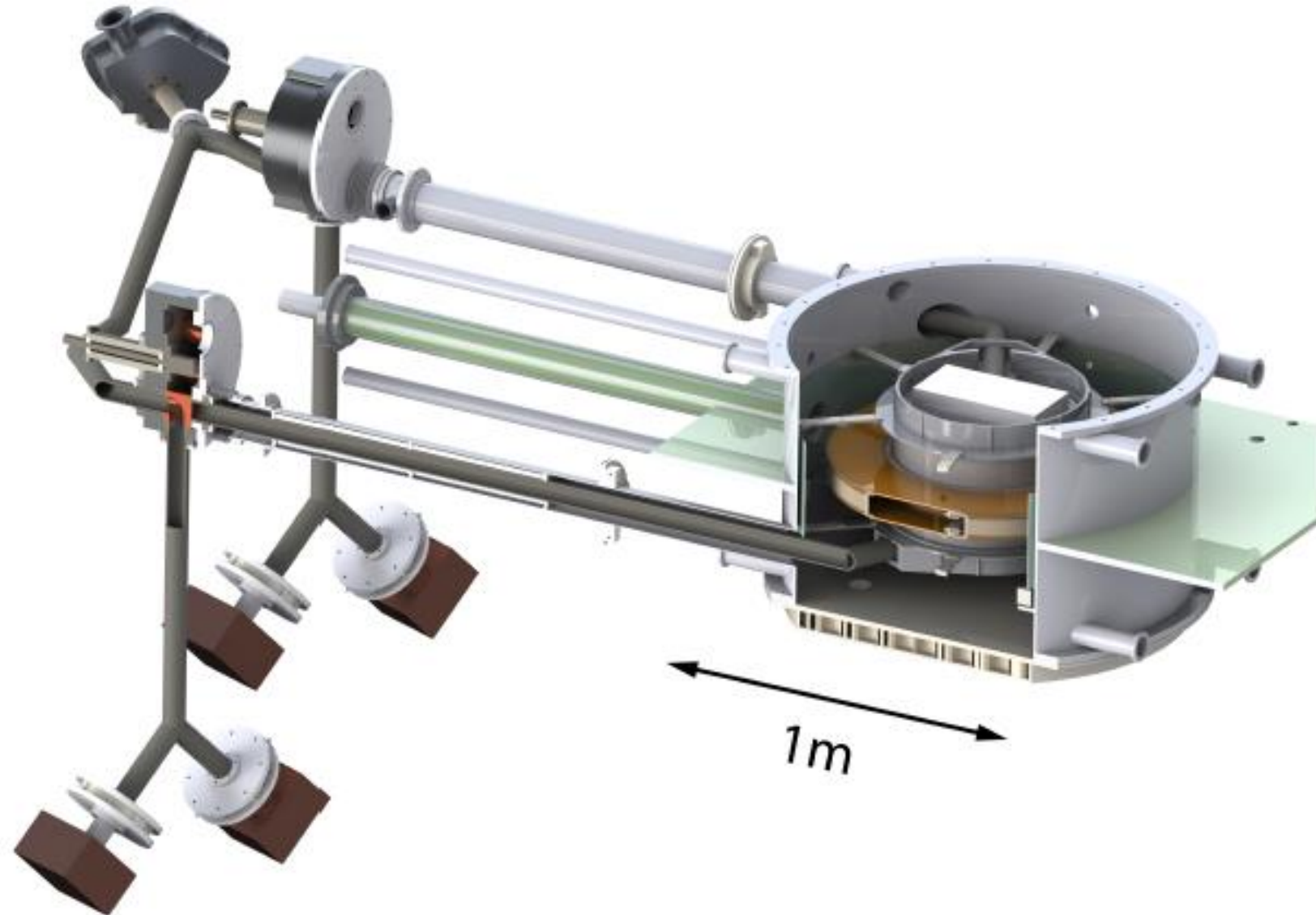


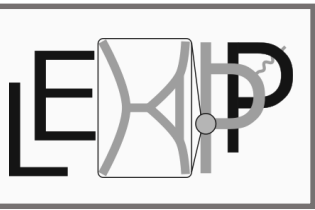
The PanEDM Experiment



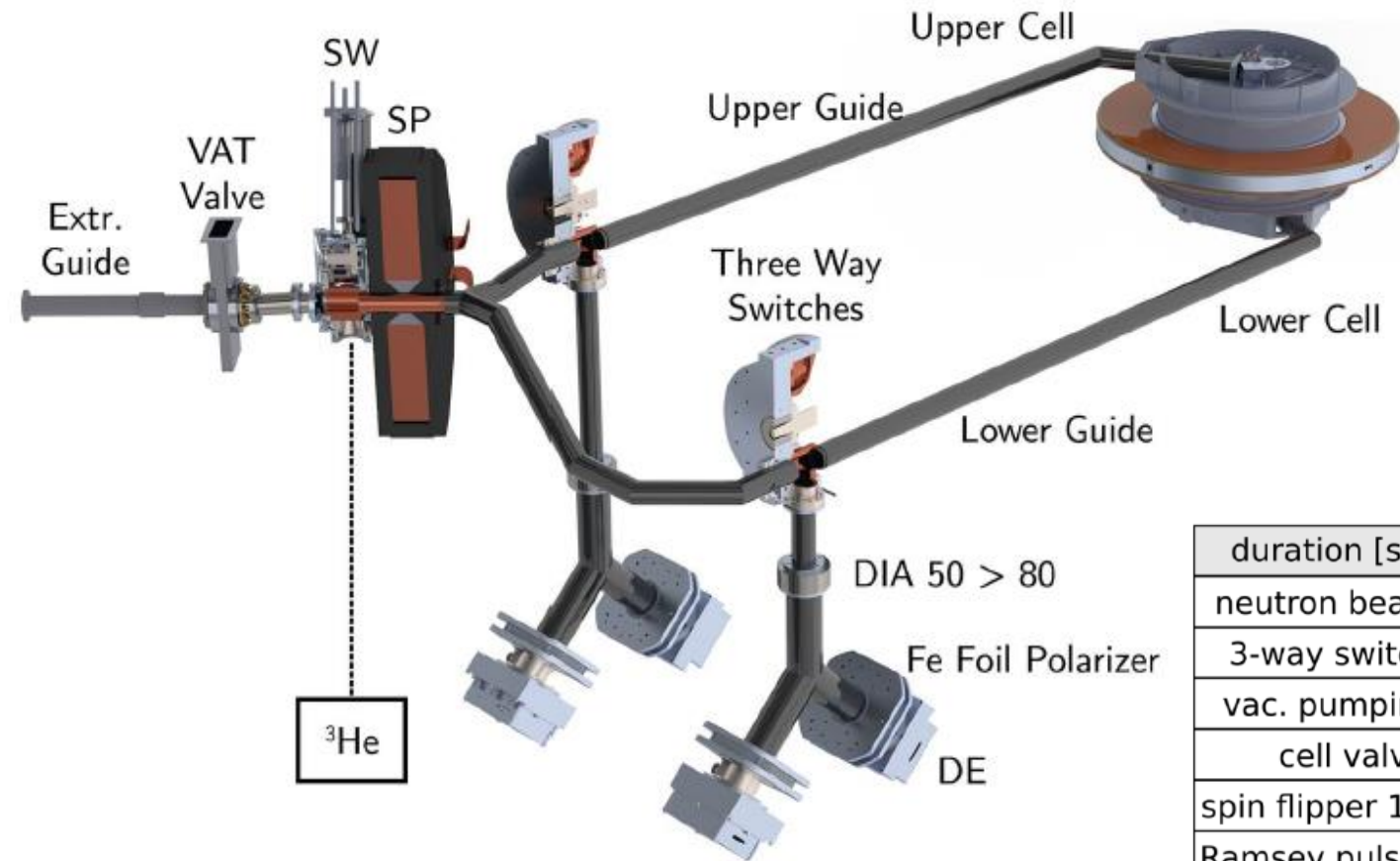


The PanEDM Experiment



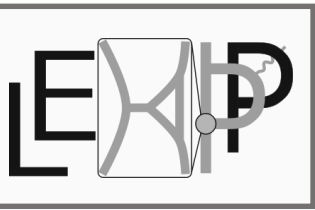


The PanEDM Experiment

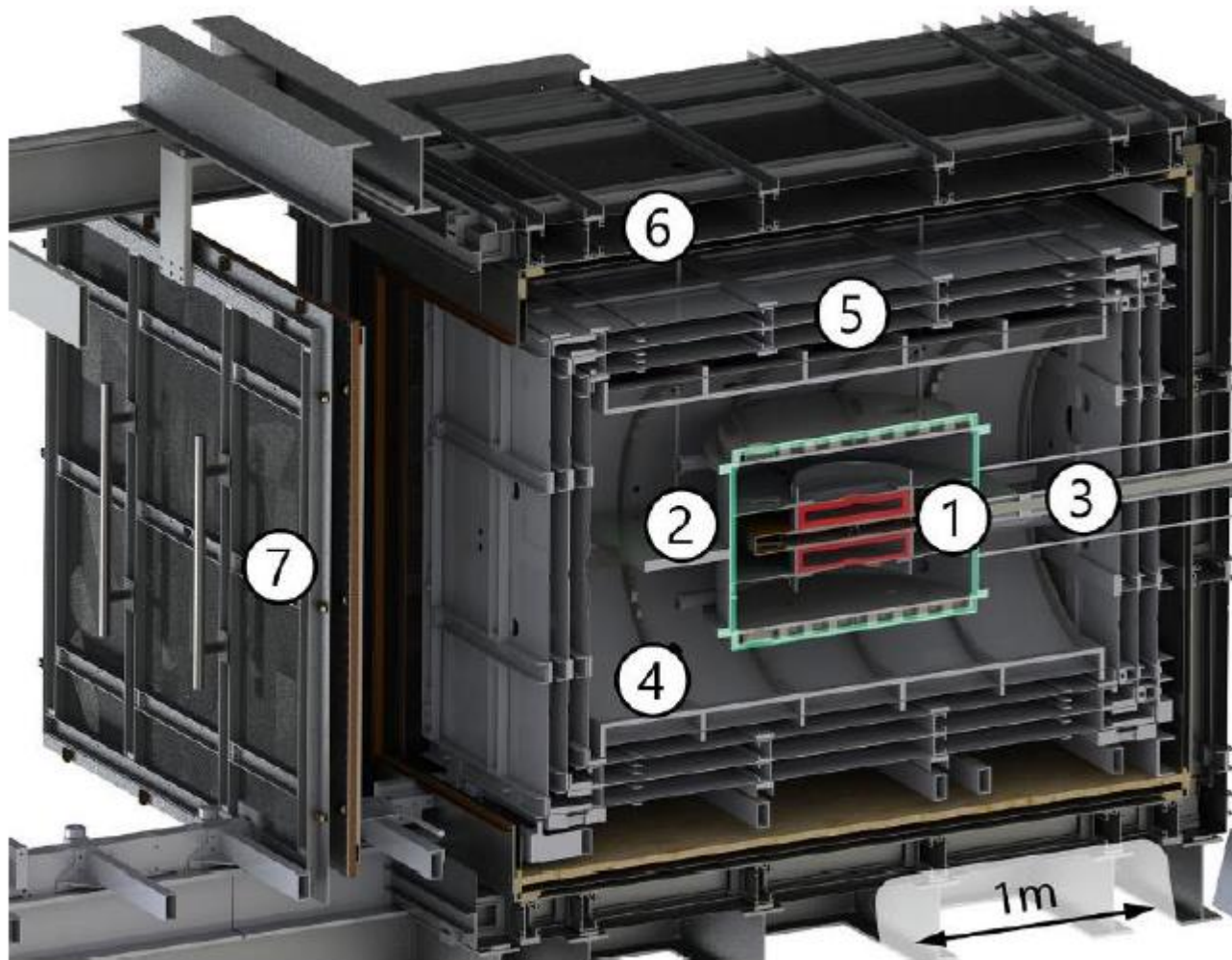


The recipe for an EDM measurement:

	preparation		Ramsey cycle			counting	
duration [s]	30	80	80	110	60	30	50
neutron beam	on	off			on		
3-way switch	vac	fill →	source ↘ detectors		detectors ↙ cells		
vac. pumping	cells	guides					
cell valves	open		closed			open	
spin flipper 1/2			various stability tests			1↑ 2↓	1↓ 2↑
Ramsey pulses			90°	180°	90°		
Hg magnet.		pumping	measure			syst. tests	
UCN detection	background, detector & source - stability					UCN cnt	
B ₀ field	set	measure					
E field	ramp	HV at setpoint					



The PanEDM Experiment



1: EDM cells

3: HV feed

5: Inner shield

7: Outer shield door

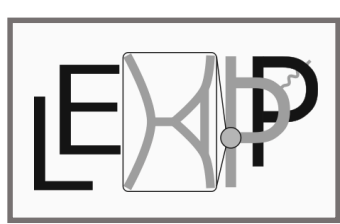
2: Vac. Chamber

4: B_0 & B_1 coil

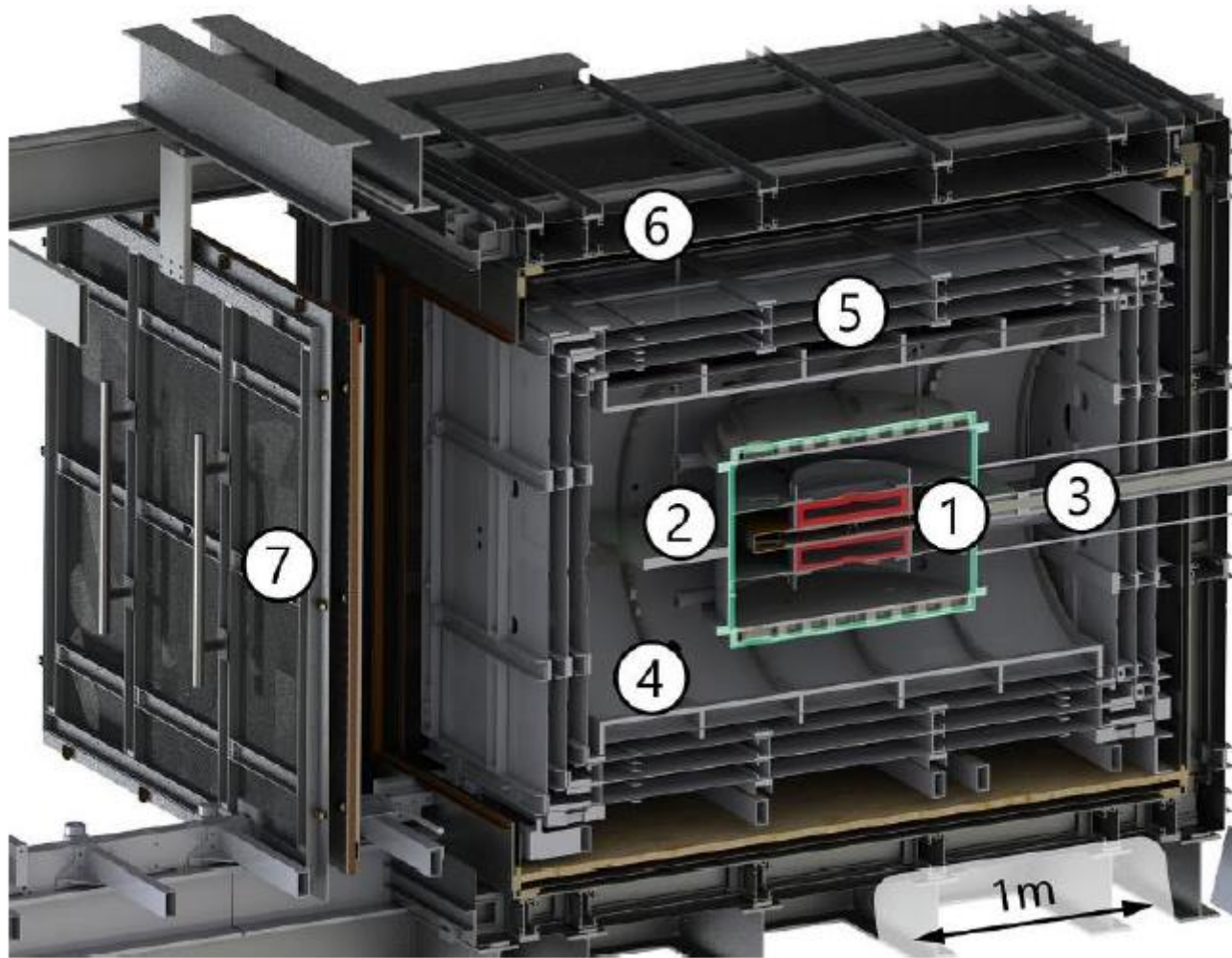
6: Outer shield



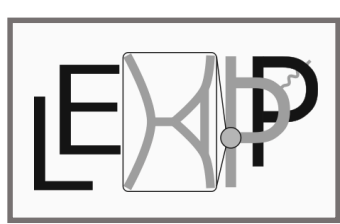
PanEDM @ ILL, 2021



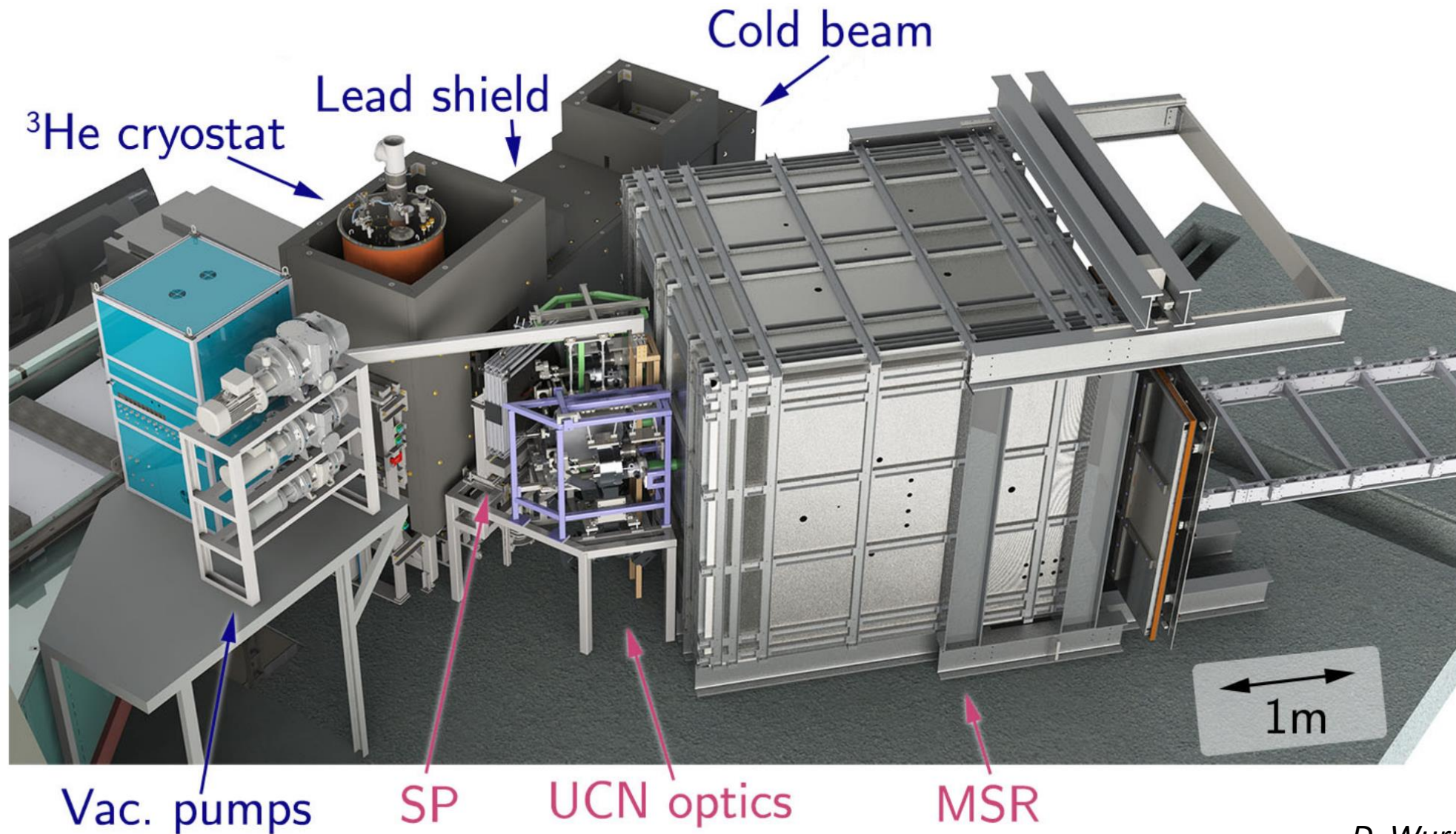
The PanEDM Experiment

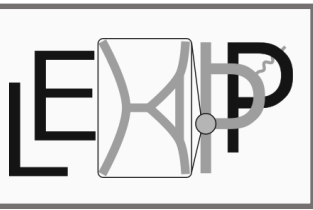


Rev. Sci. Instr. 85(7), 075106 (2014)
J. Appl. Phys. 117(18), 183903 (2015)

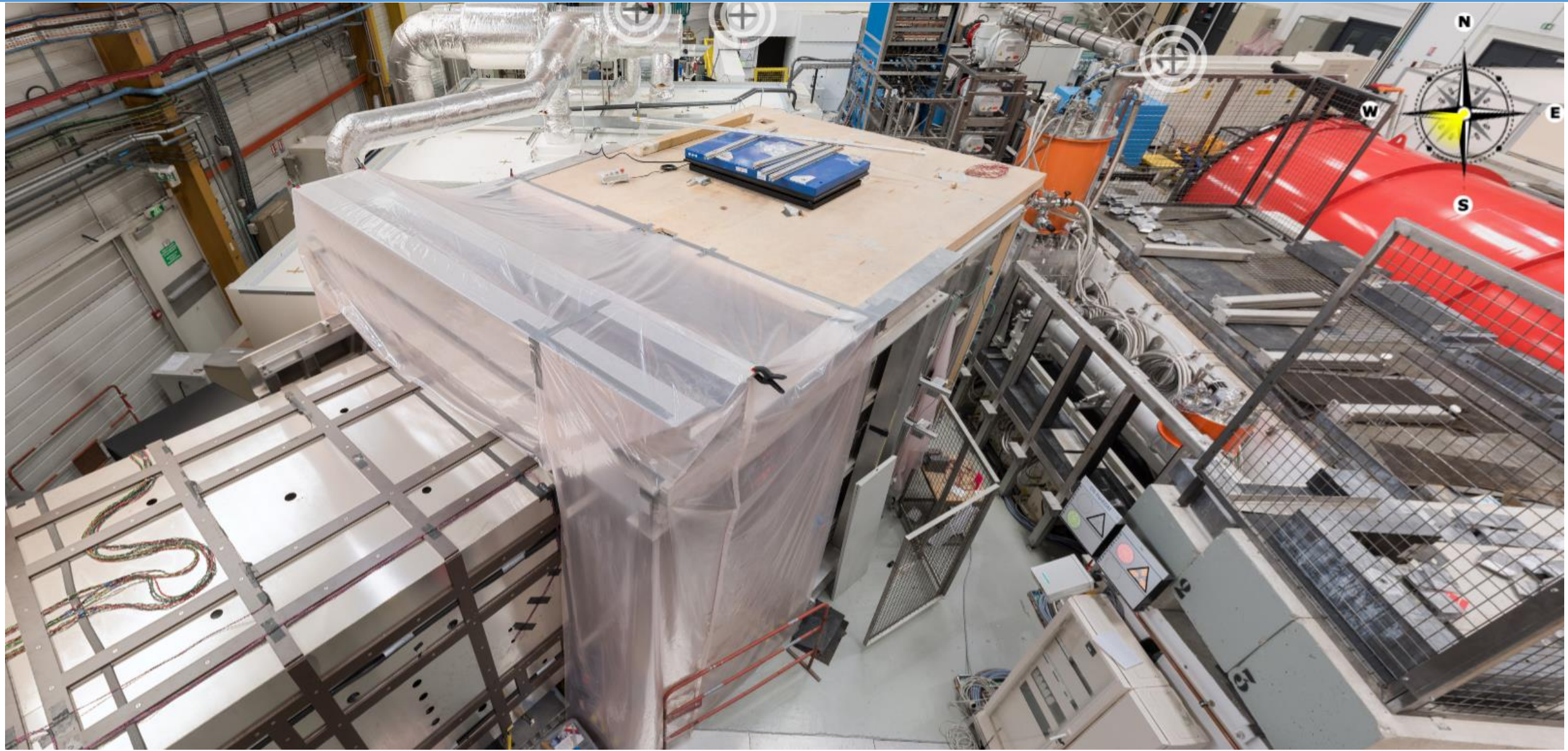


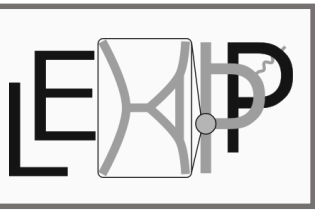
The SuperSUN-PanEDM Installation



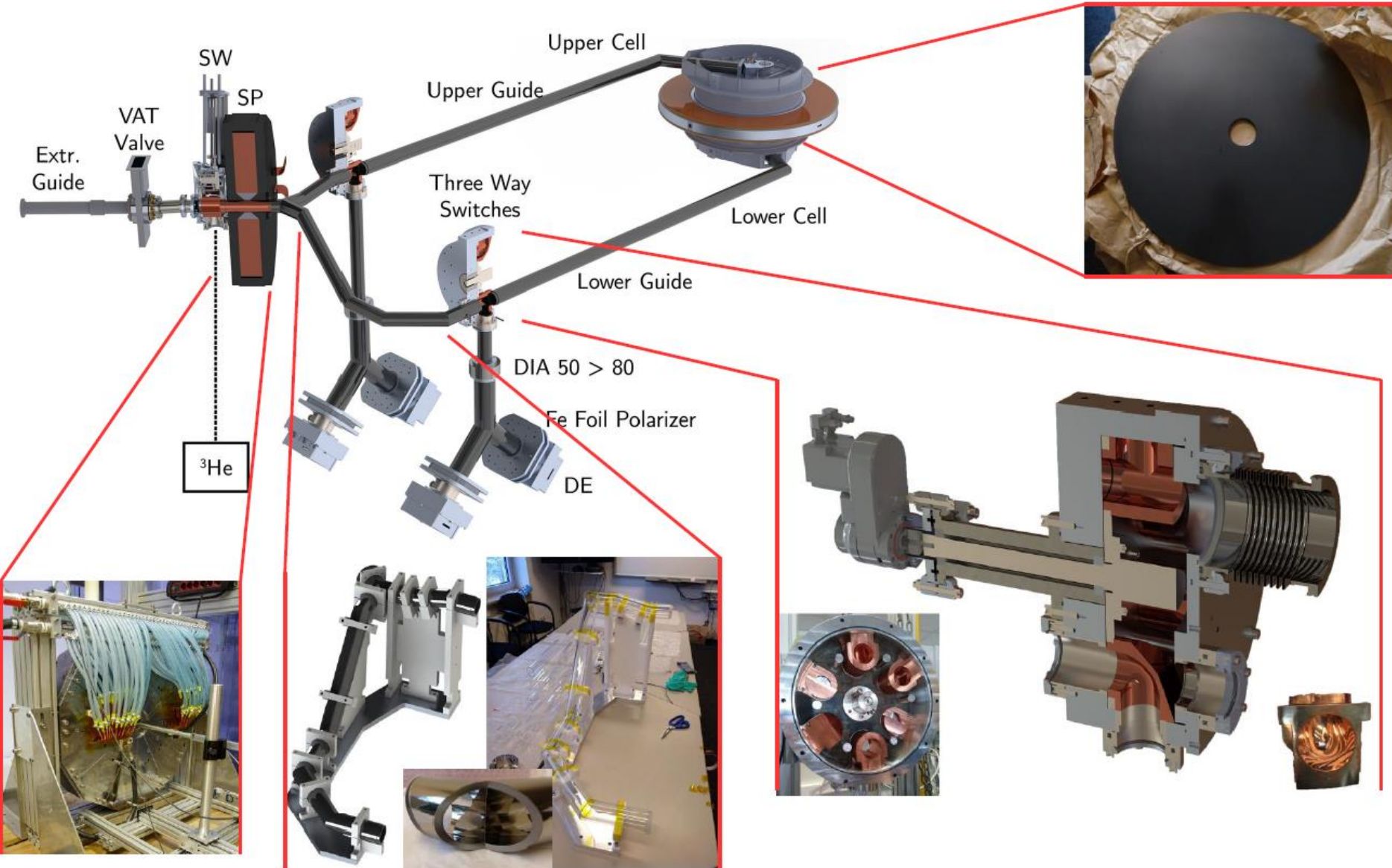


Reality always looks messier!

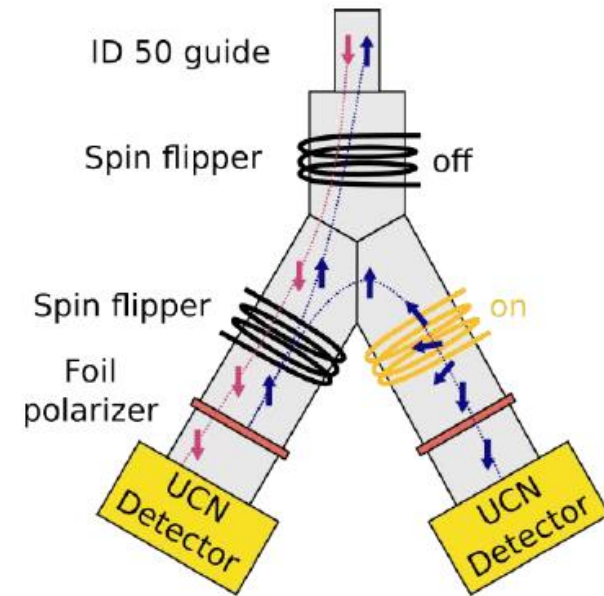


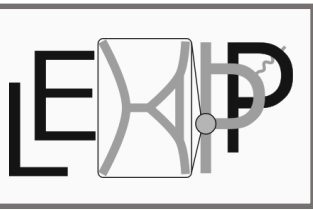


SuperSUN-PanEDM Interface

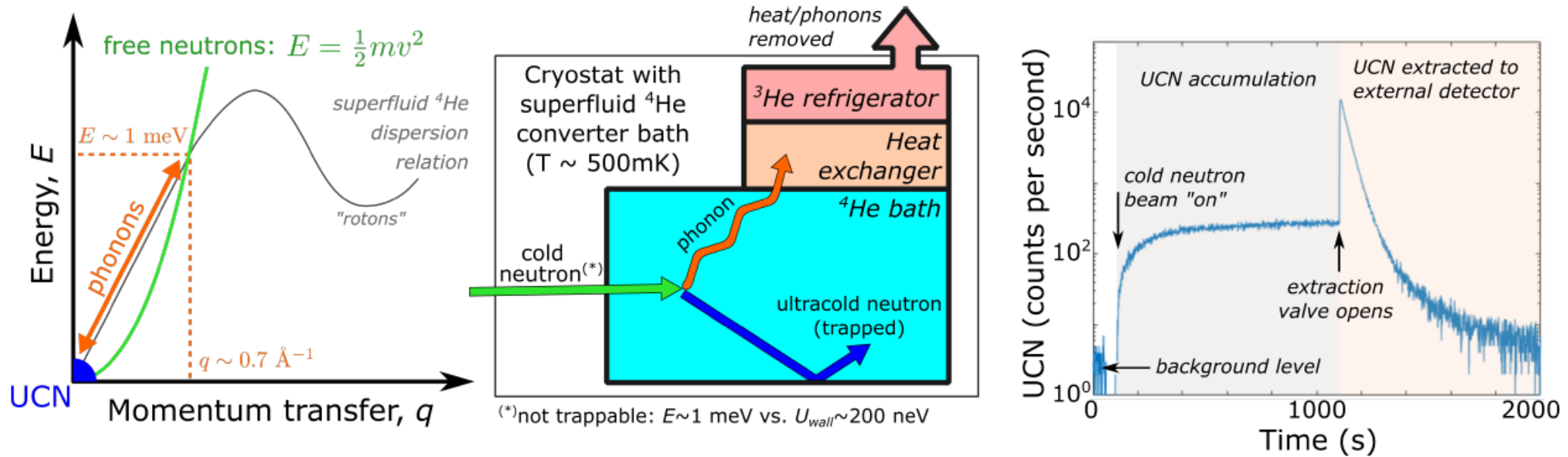


Simultaneous spin detection

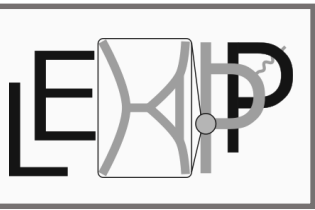




The SuperSUN Neutron Source



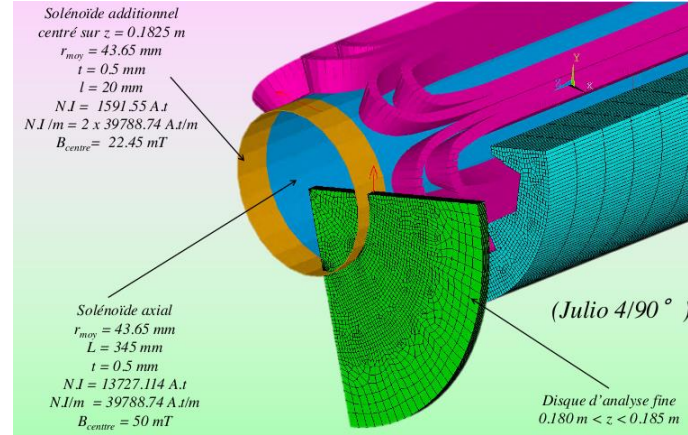
- Figure of merit for production: $\tau \cdot \int d^3\mathbf{k} (1 - e^{-n\sigma l(\mathbf{k})}) \left. \frac{d\Phi}{d\lambda} \right|_{8.9\text{\AA}}$
- Note: partial mean free path $\lambda_{\text{UCN}} \sim 10 \text{ km}$, while $\lambda_{\text{tot}} \sim 10 \text{ m}$
- Loss for a 3m converter: factor 10 (unused CN beam)
- Loss for *ex-situ* storage: factor 100 (UCN extraction/transport/detection)



SuperSUN Neutron Source: Cutaway

^3He pumping

1K pot



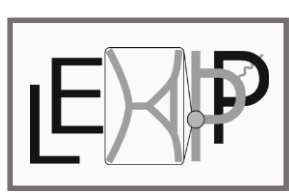
SC Octupole ~2.1T



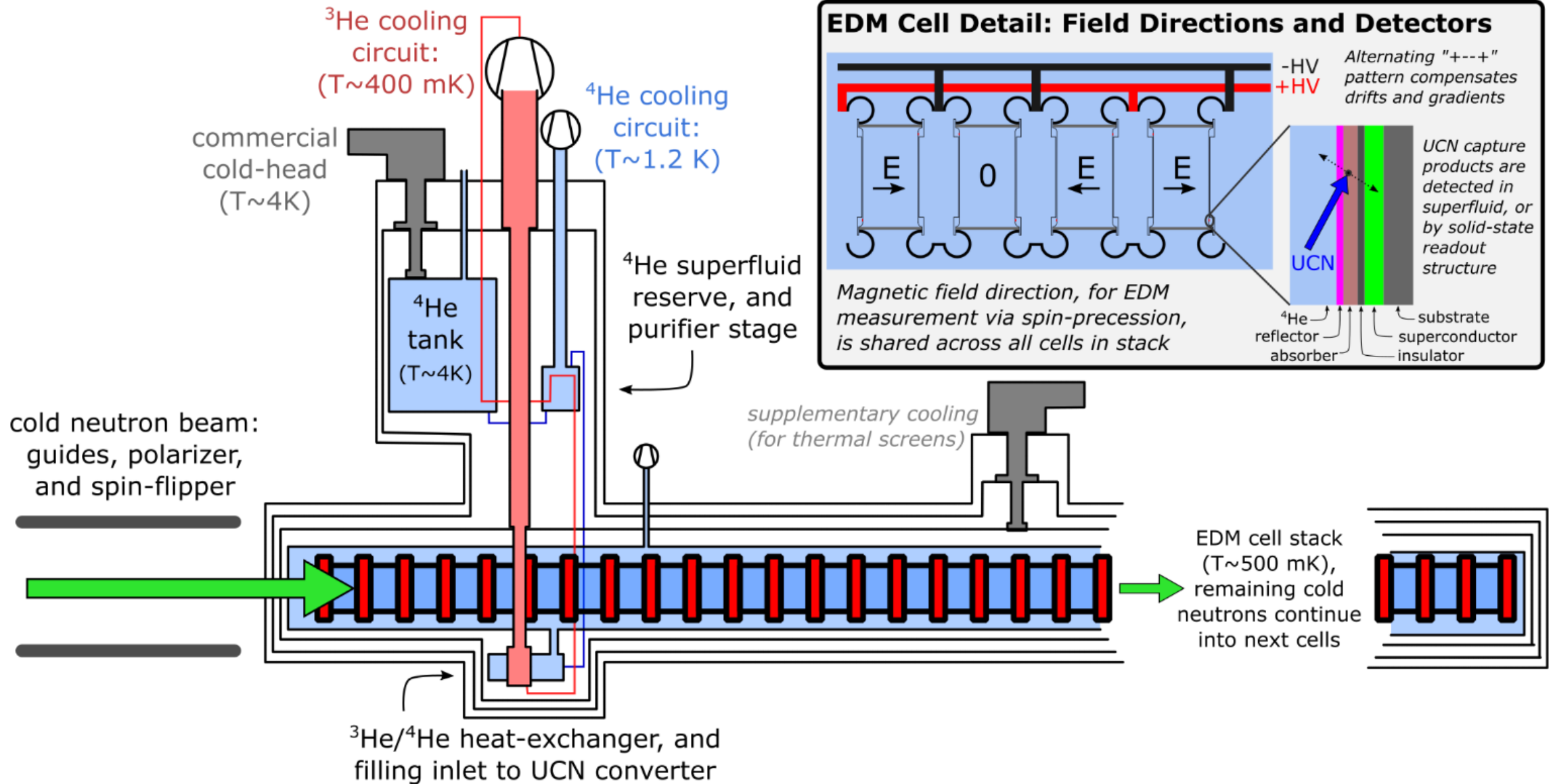
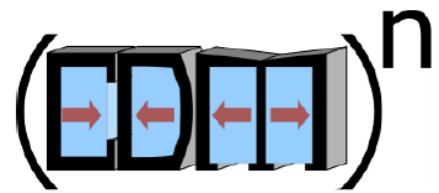
cryogenic CN guide

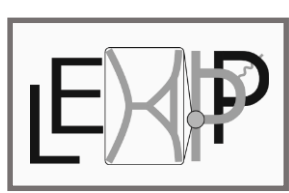
Isotopically pure ^4He

UCN out

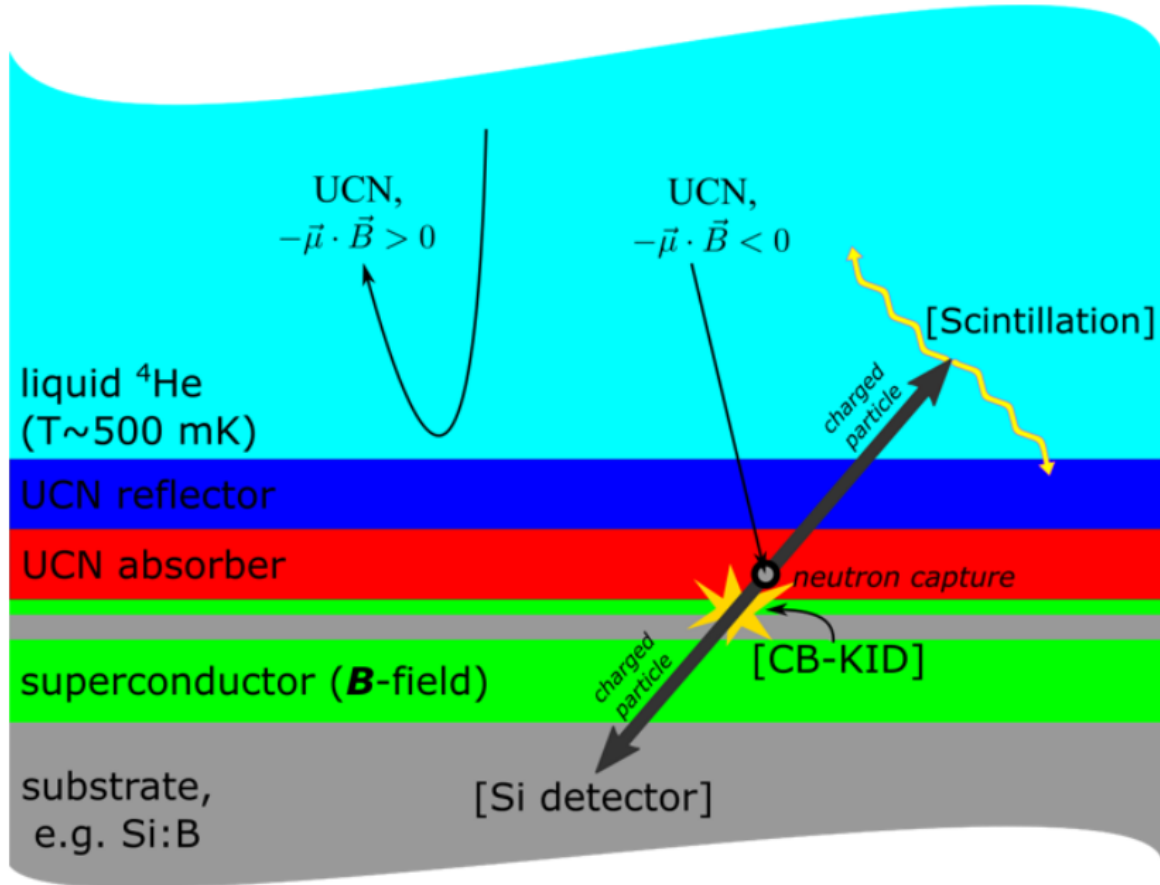
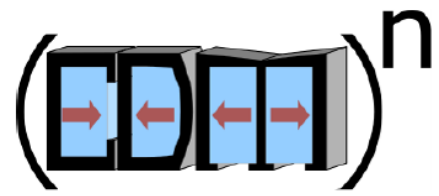


The next generation... scaling up!

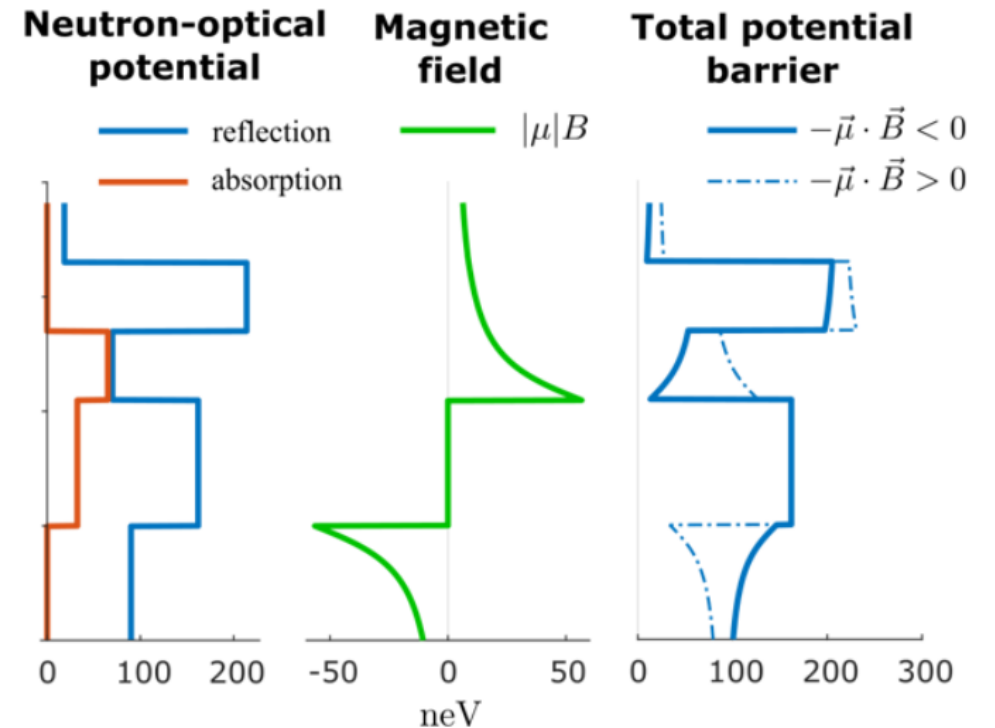


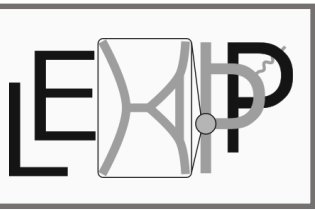


“Quantum Sensing” for Neutrons

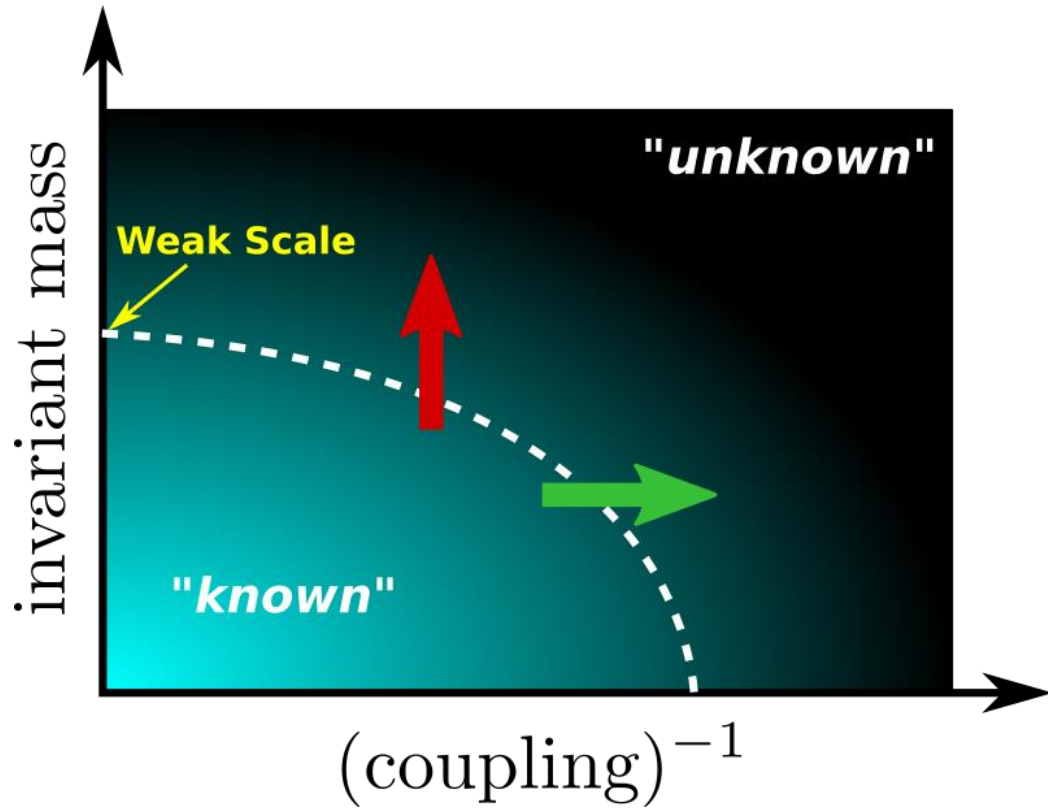


Ultracold neutron (UCN) detection with polarization-sensitivity, via applied magnetic fields partially cancelling the neutron-optical potential for one spin state. Various readout mechanisms to be explored.





Thematic Recap



1

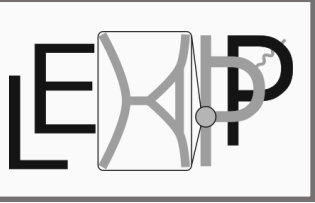
Statistics are the first key!

2

Observation time is the second

3

... and yes, it finally makes sense to follow the green arrow!



Questions?

EXPERIMENT

OUR NEW ~~TELESCOPE~~ WILL
ANSWER TWO KEY QUESTIONS:

- 1) WHY IS THERE ALL THIS MATTER?
- 2) CAN WE DO ANYTHING ABOUT IT?



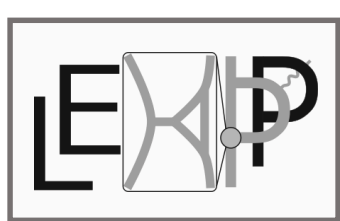
what-if.xkcd.com

Special thanks to:

U. Schmidt, Heidelberg
Heidelberg mechanical workshop
Heidelberg technical design office

Institut Laue-Langevin, NPP division
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THERE ARE
GHOSTS HERE.

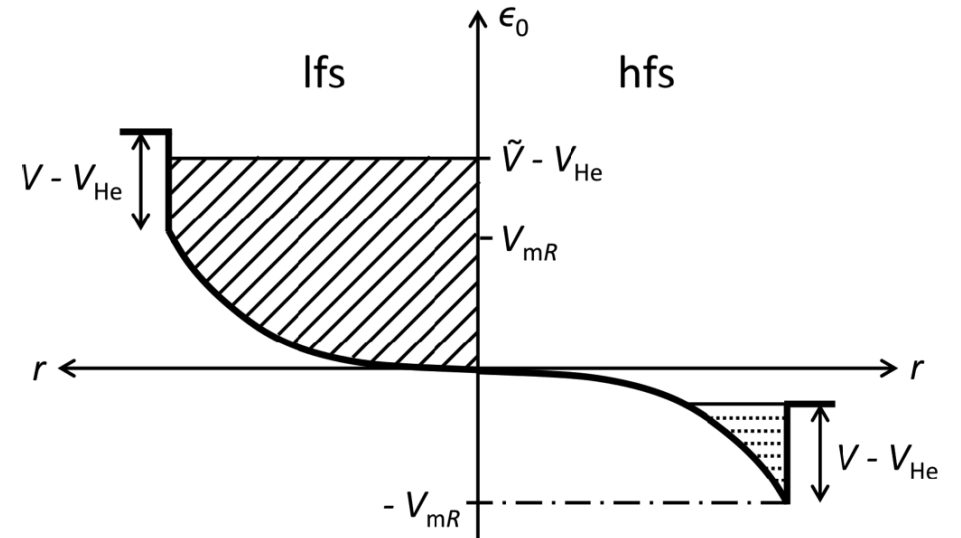
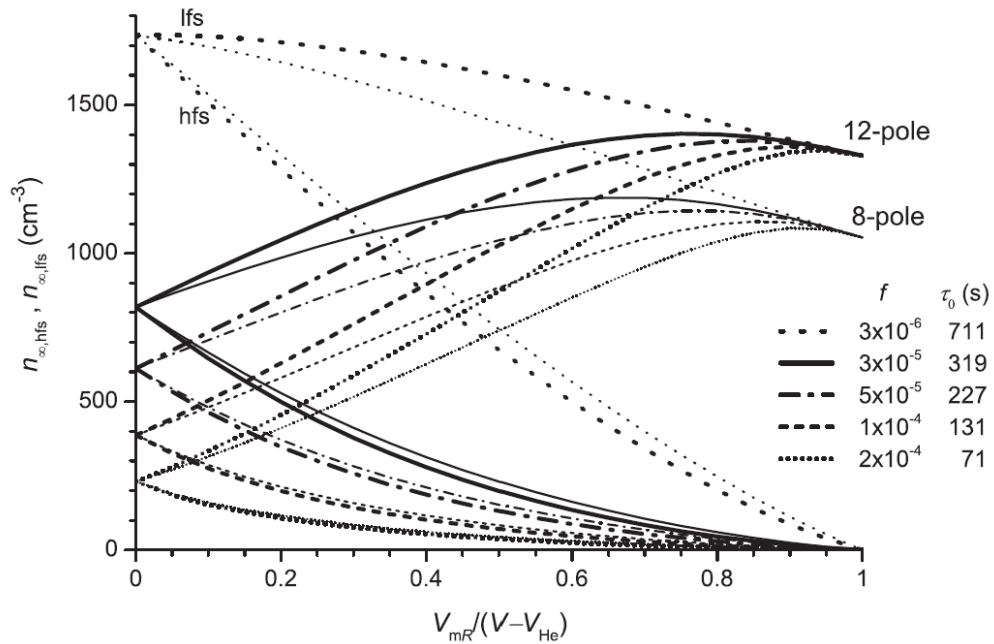
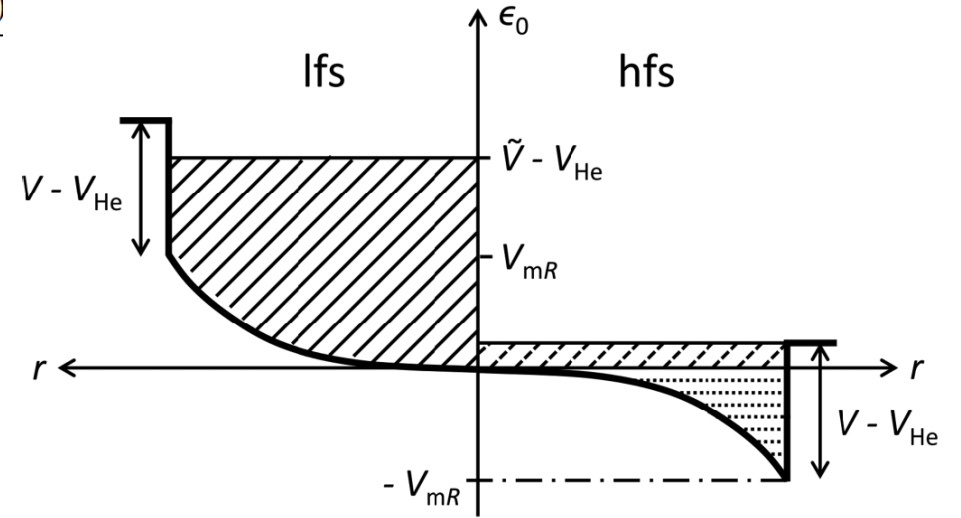
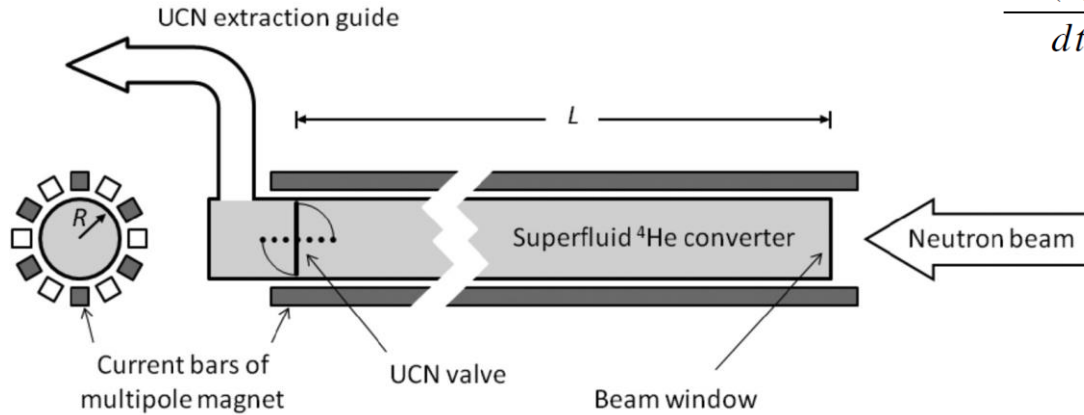


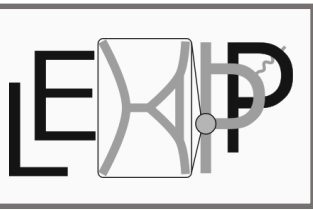
xkcd.com

Minimizing UCN Storage losses

PHYSICAL REVIEW C **92**, 015501 (2015)

$$\frac{dn(\epsilon_0, t)}{dt} = p(\epsilon_0) - \frac{n(\epsilon_0, t)}{\tau(\epsilon_0)}$$





Neutron Delivery to SuperSUN

In SuperSUN's converter vessel:

- $R \sim 15 \text{ UCN}/(\text{cm}^3 \text{ s})$ expected

End of guide H523:

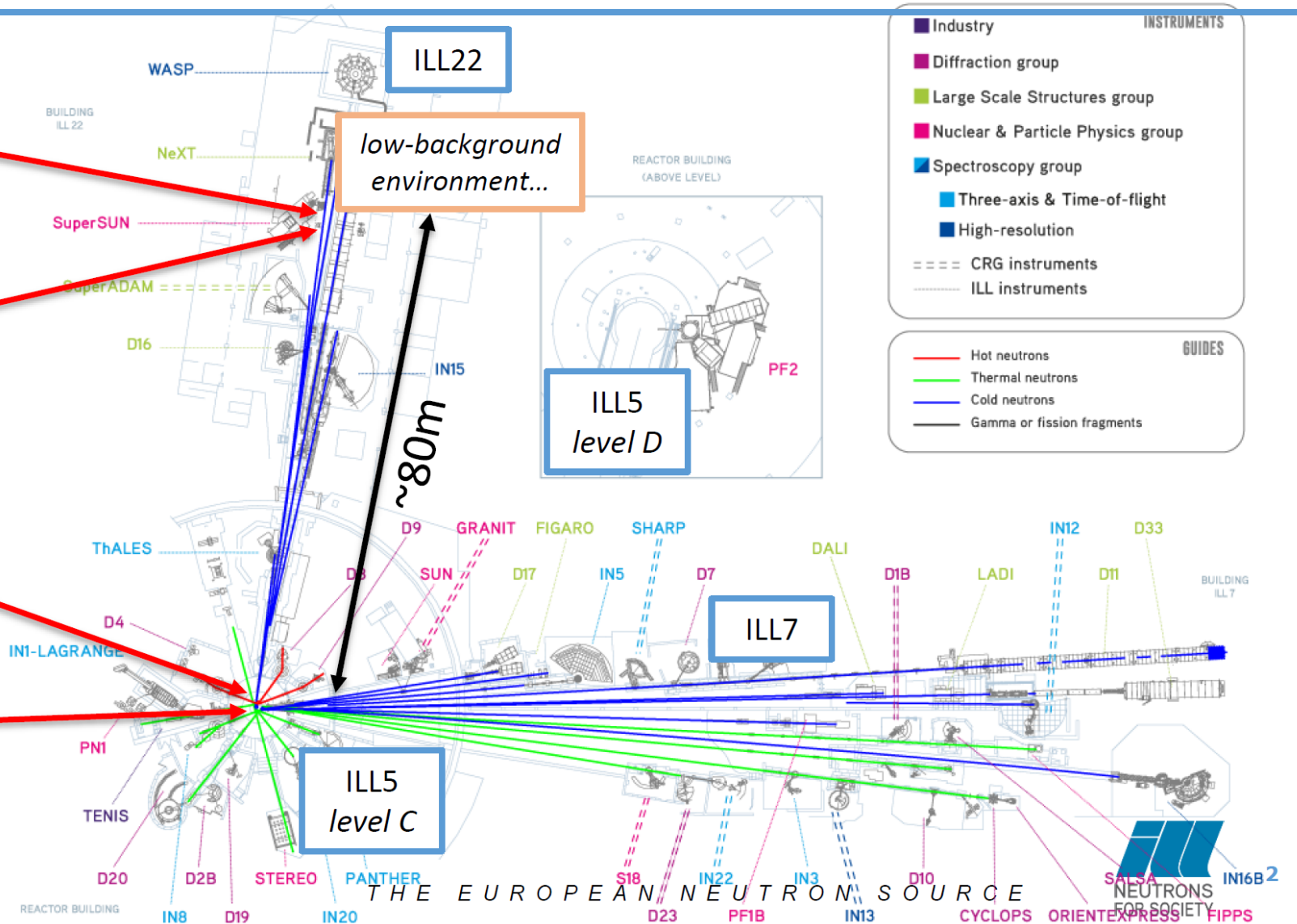
- $\Phi \sim 2 \times 10^{10} \text{ n}/(\text{cm}^2 \text{ s})$

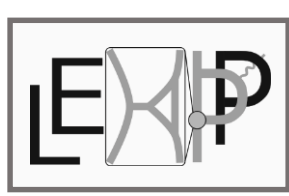
Horizontal Cold Source:

- $\Phi \sim 10^{14} \text{ n}/(\text{cm}^2 \text{ s})$

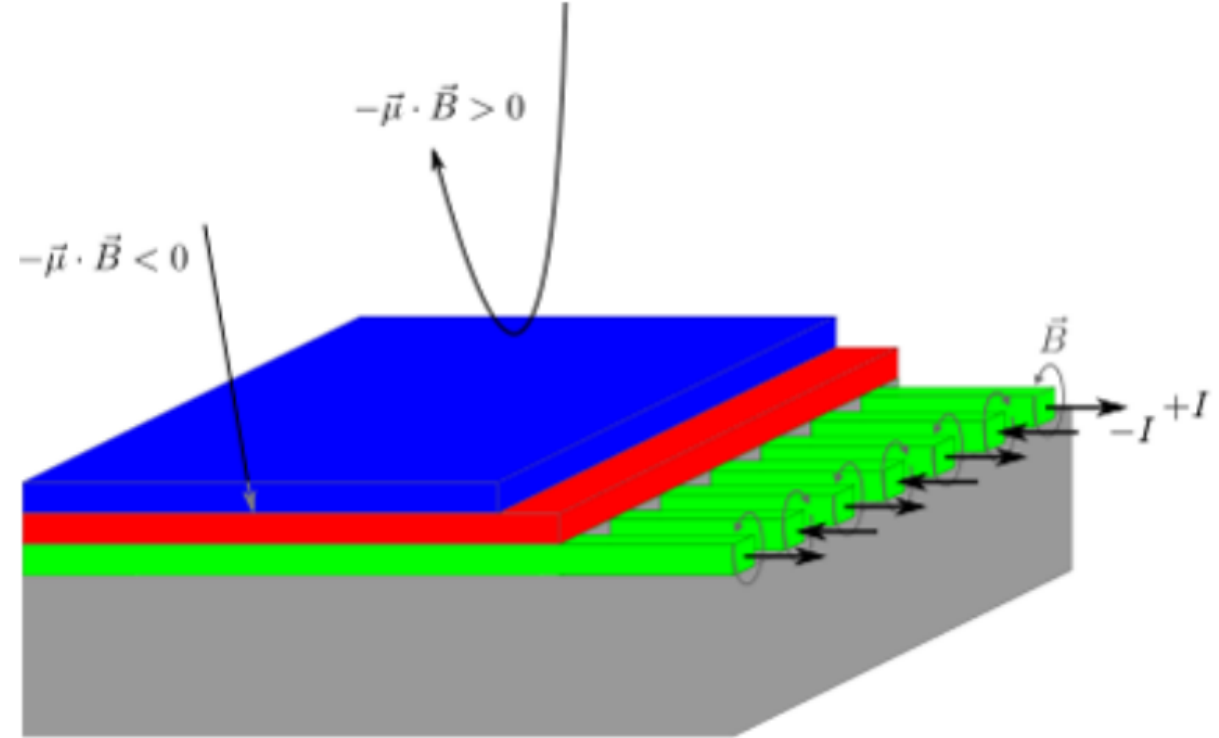
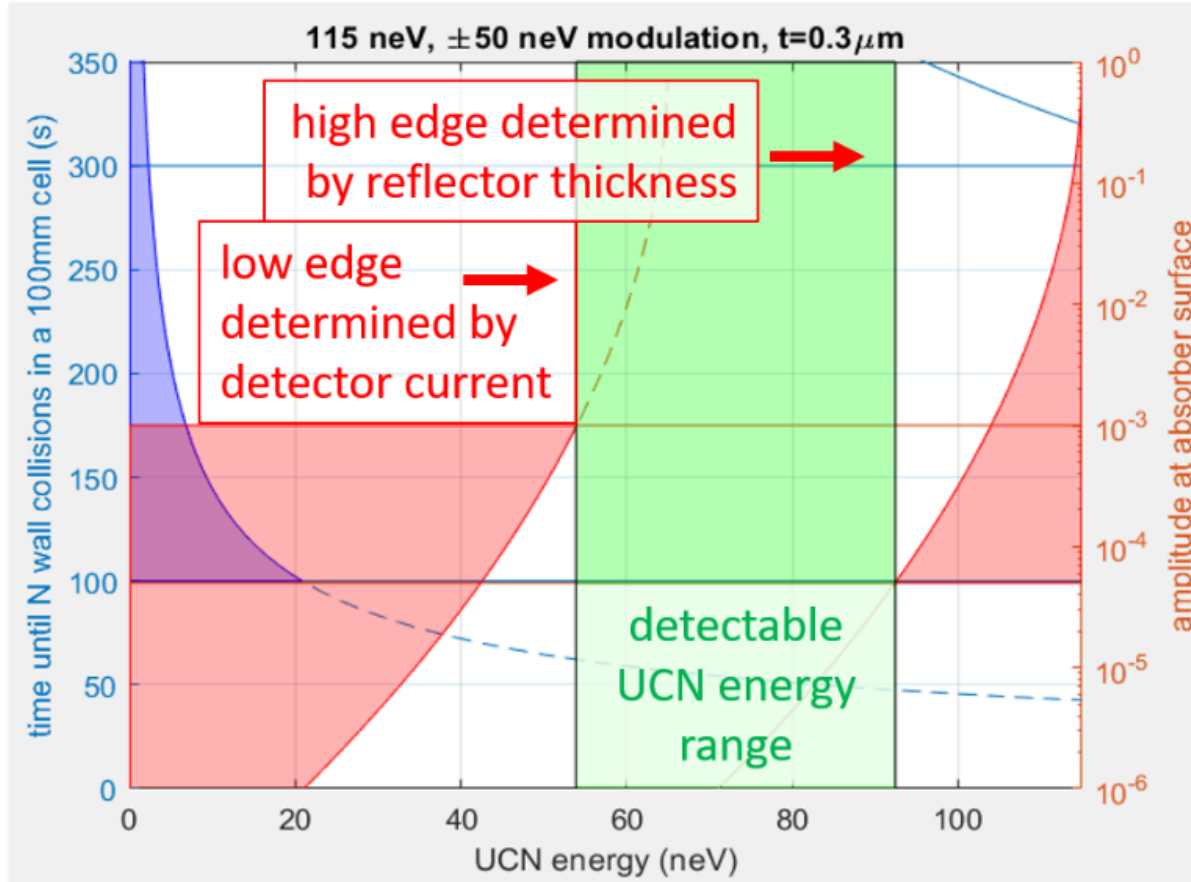
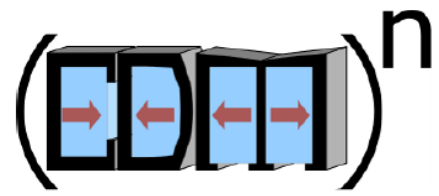
In pile:

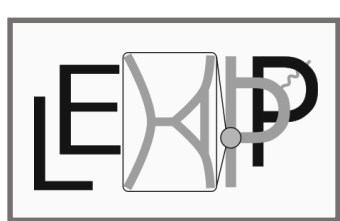
- $\Phi \sim 1.5 \times 10^{15} \text{ n}/(\text{cm}^2 \text{ s})$





“Quantum Sensing”: Spin and Energy





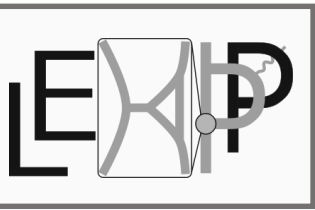
A digression on childrens' toys...

*The SuperSUN Tapered Transition Guide
with Irregular Octagonal Cross-Section,*

or,

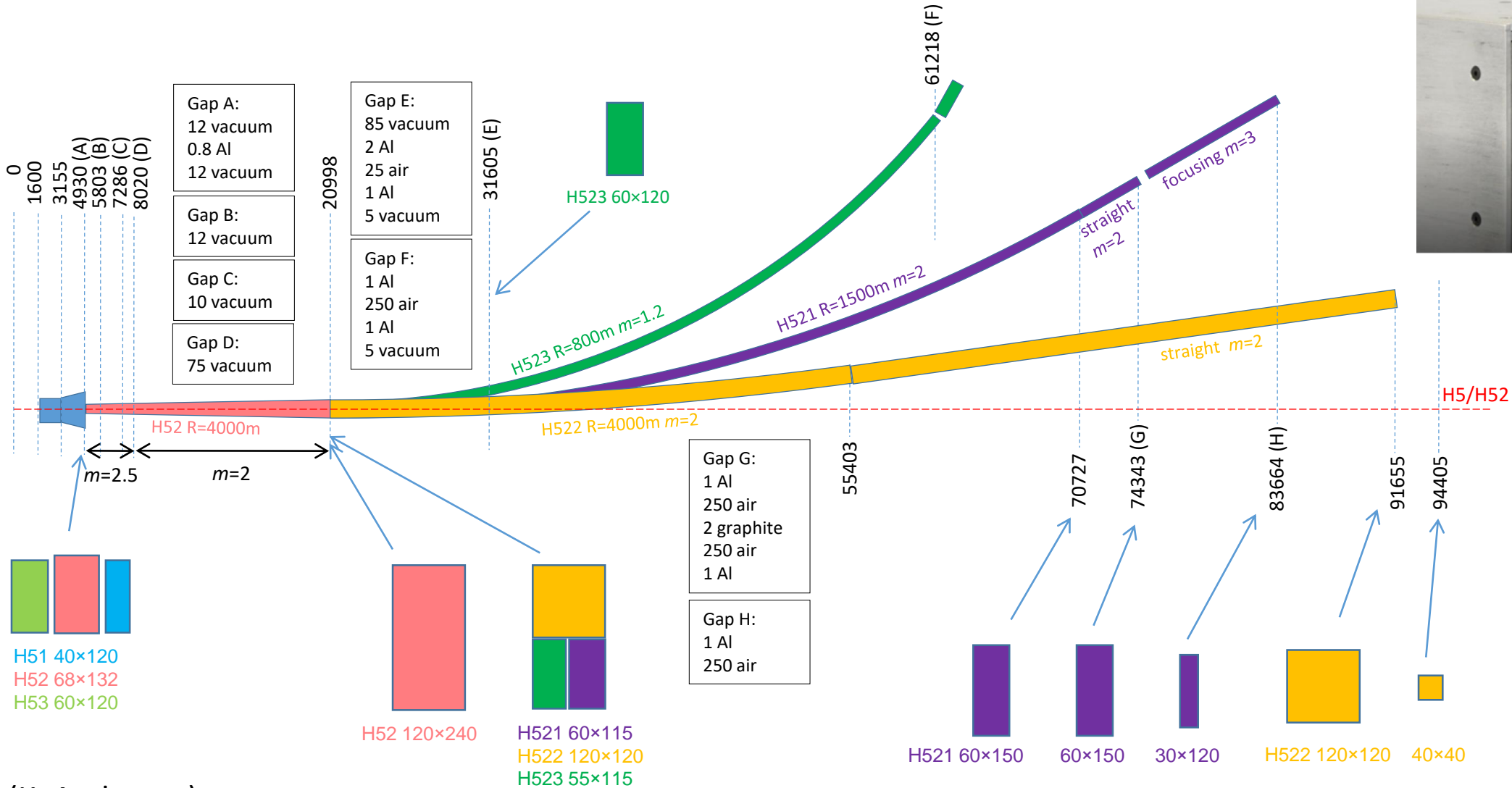
**How to Fit
Square Pegs
into Round Holes**



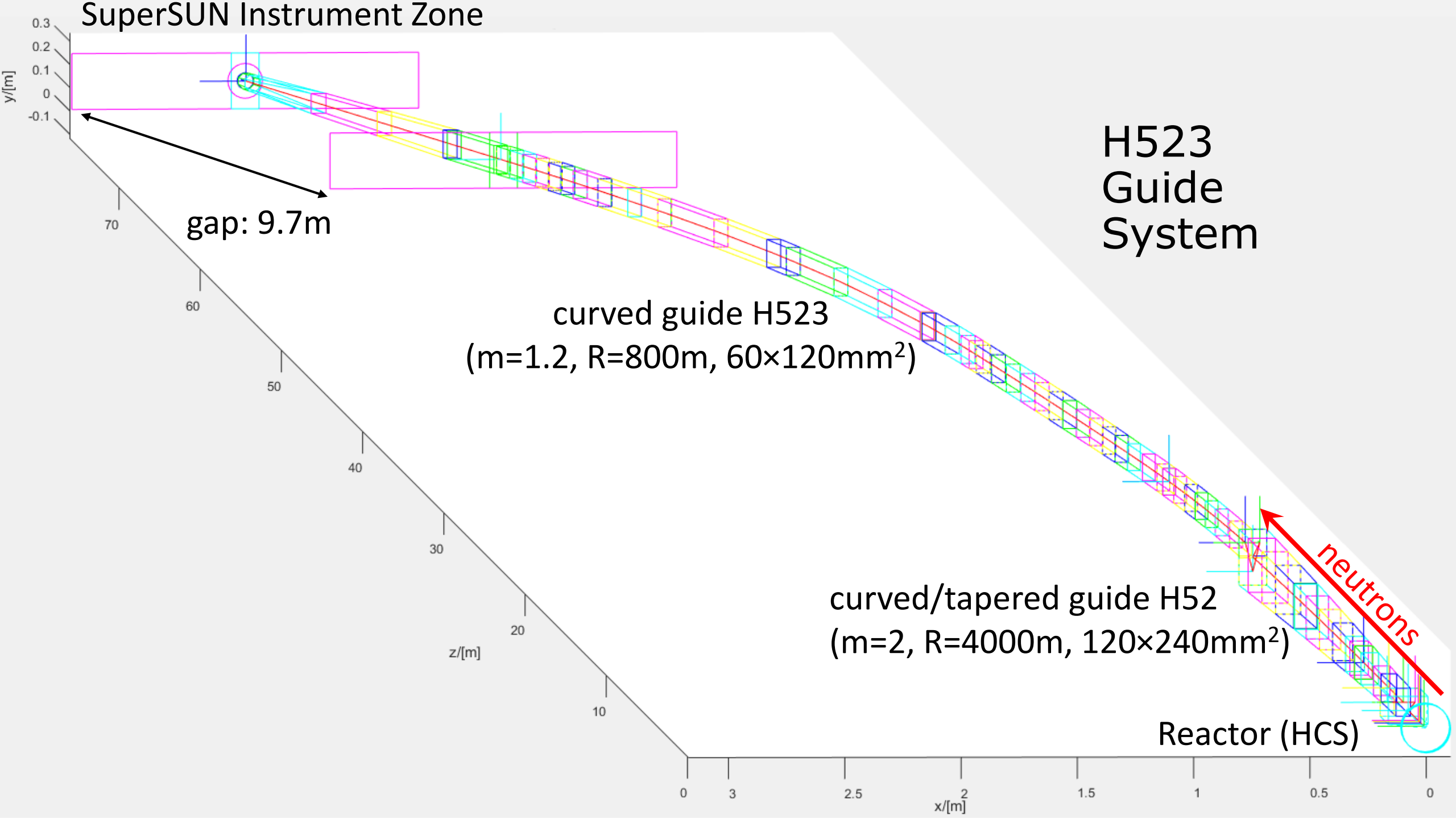


"Peg"

...or the ILL22 cold neutron guides



swissneutronics.ch



SuperSUN Instrument Zone

H523
Guide
System

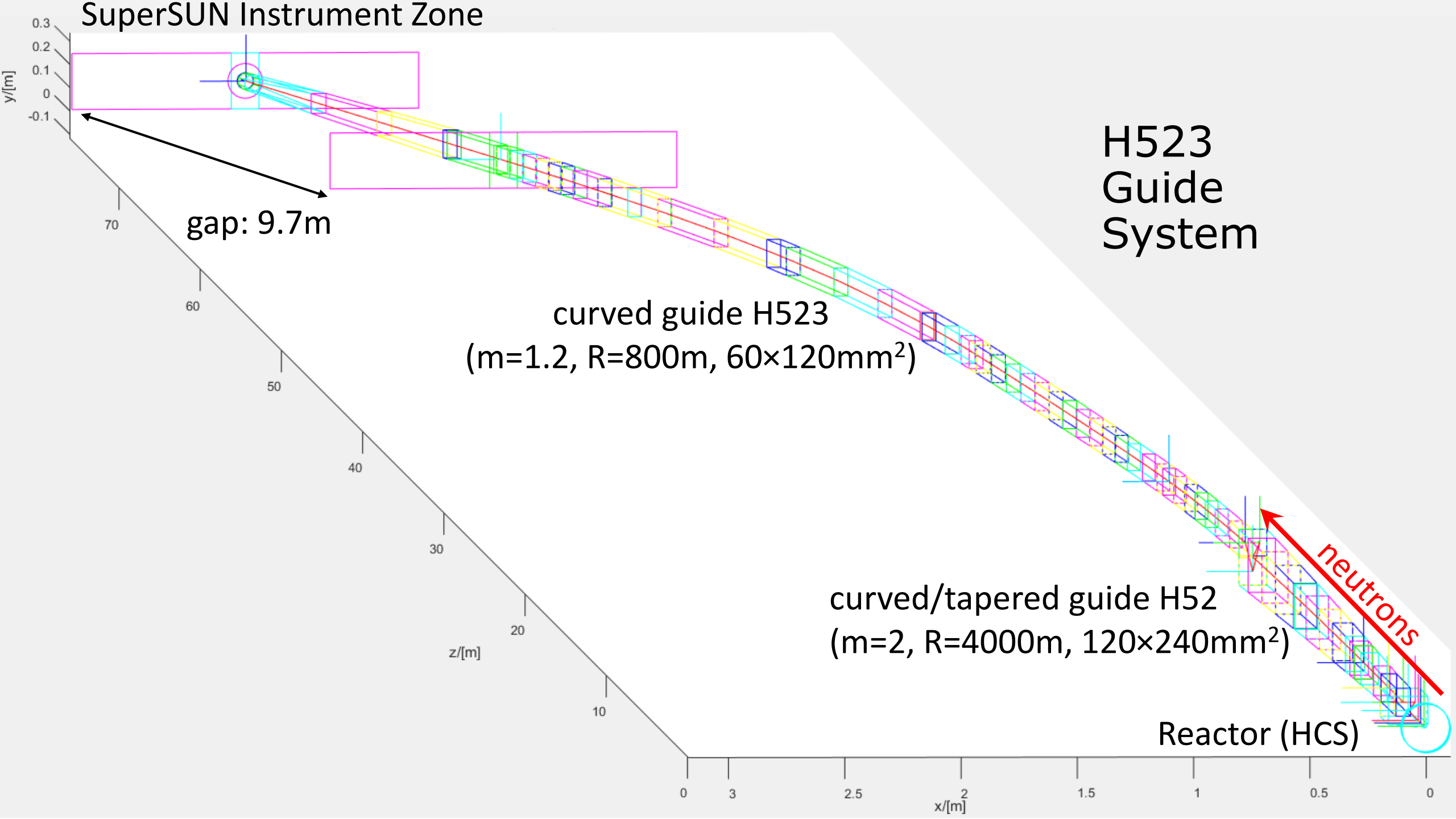
curved guide H523
($m=1.2, R=800\text{m}, 60\times 120\text{mm}^2$)

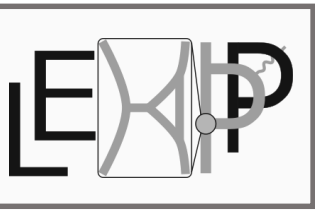
curved/tapered guide H52
($m=2, R=4000\text{m}, 120\times 240\text{mm}^2$)

Reactor (HCS)

gap: 9.7m

neutrons



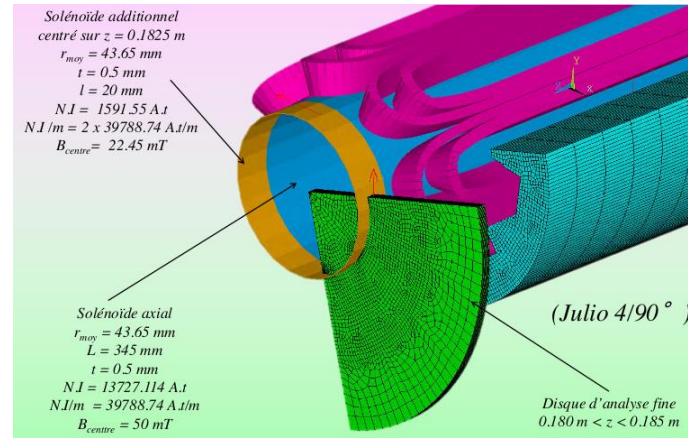


"Hole"

...or admittance matching

^3He pumping

1K pot

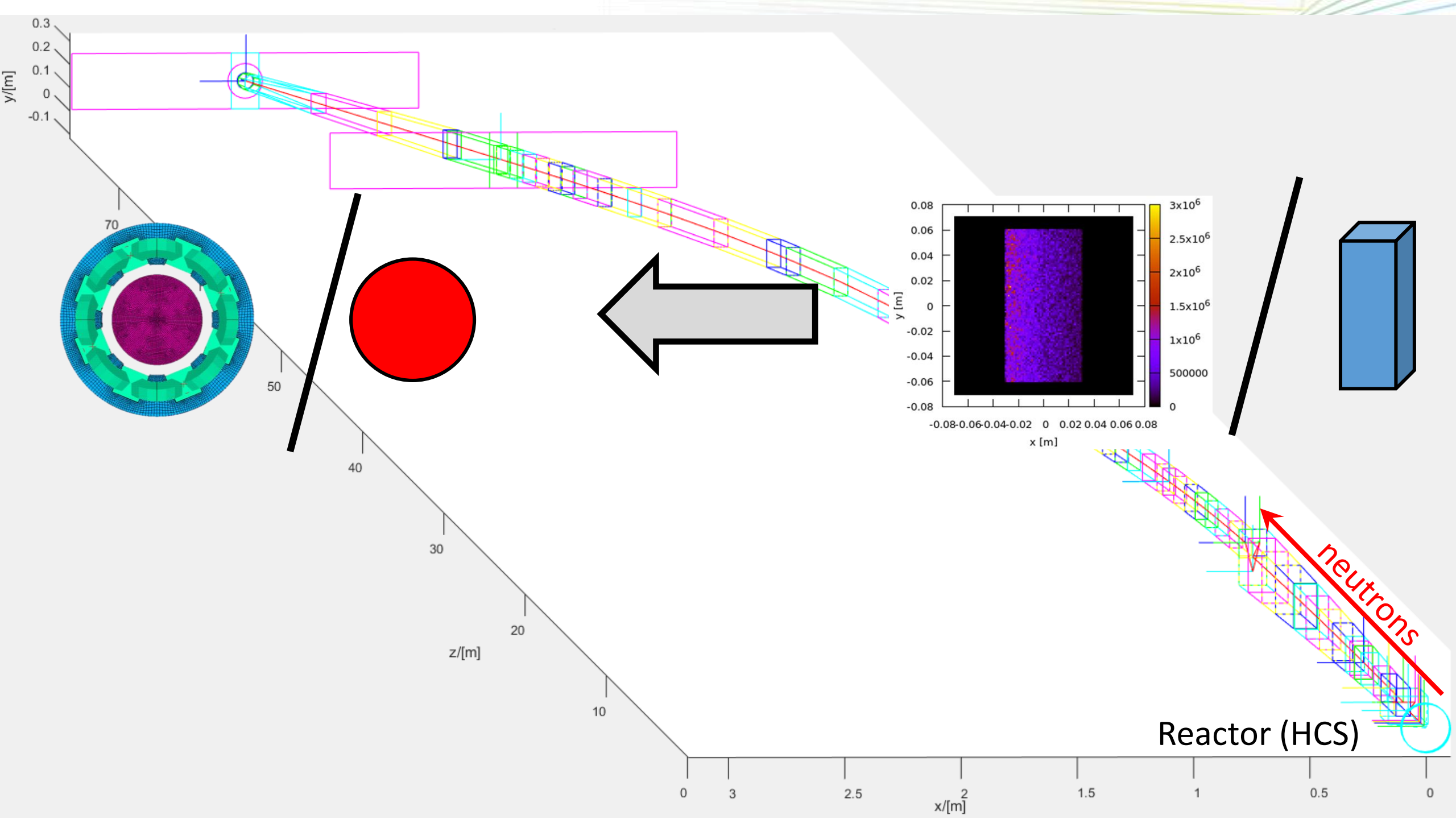


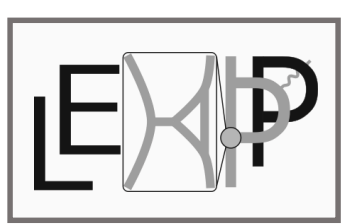
SC Octupole ~2.1T

cryogenic CN guide

UCN out

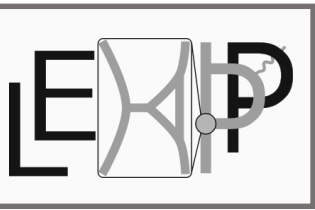
Isotopically pure ^4He





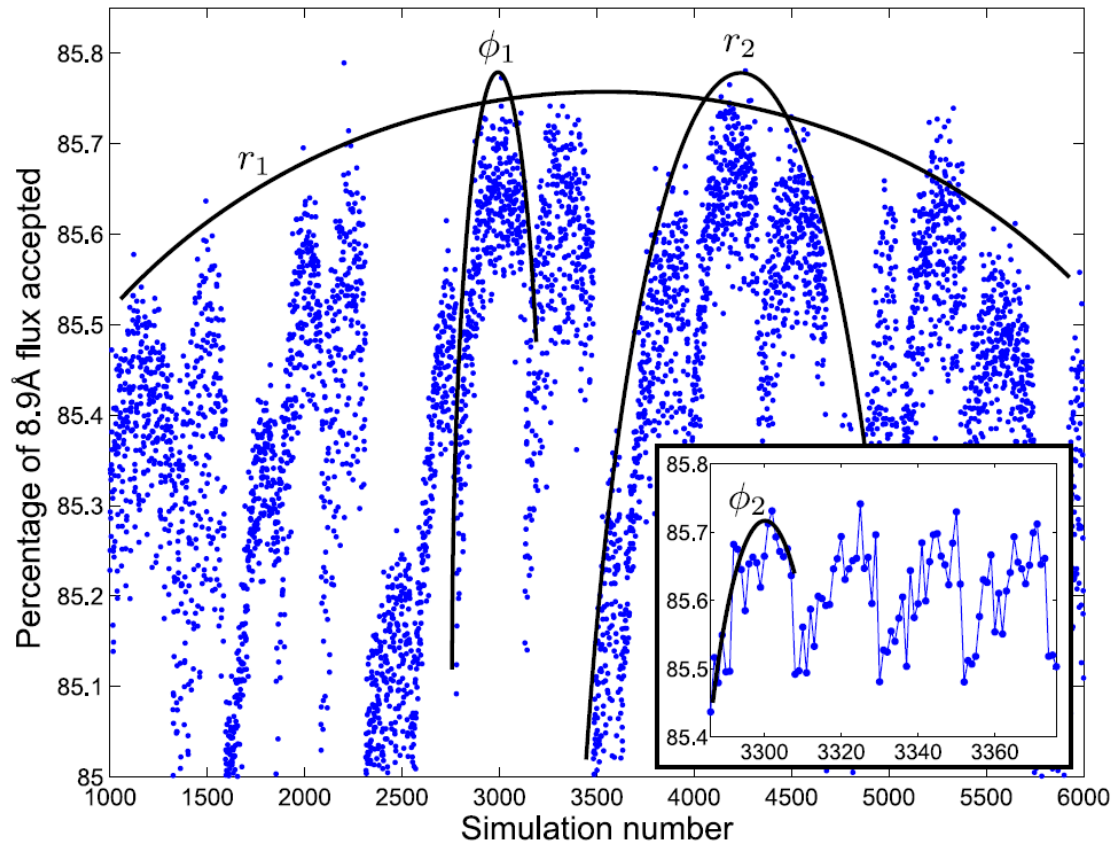
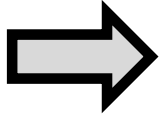
We are still missing a tool...



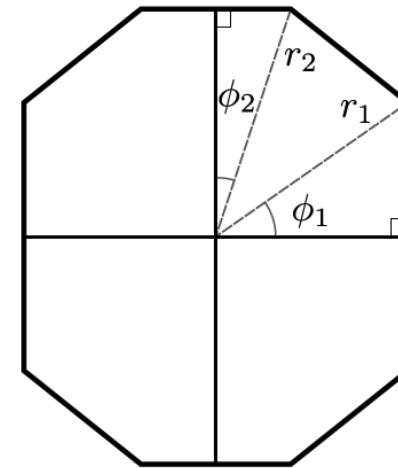


“Hammer”

...or brute force solution

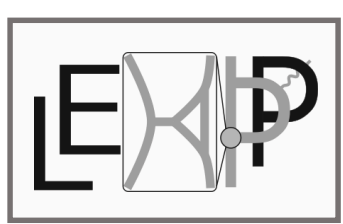


Each point:
~ 10^7 particles
~5min



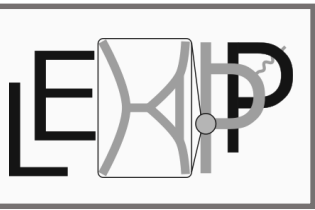
“general” octagon
parameterization

Journal of Neutron Research
20(4), 117-122 (2018)



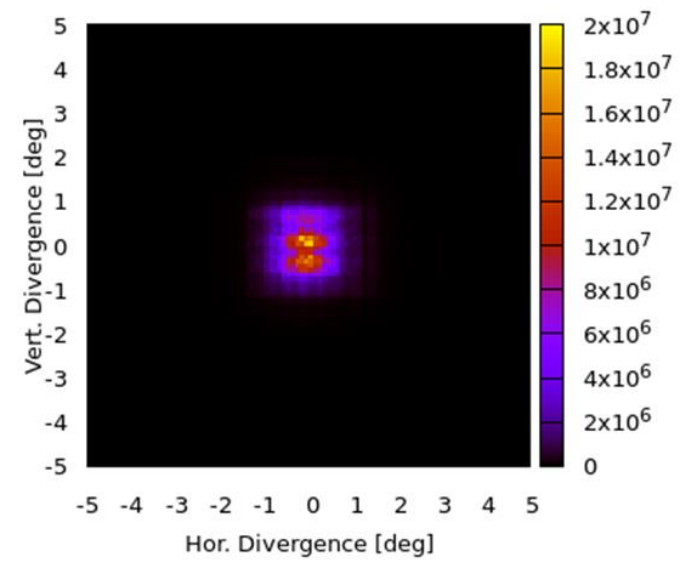
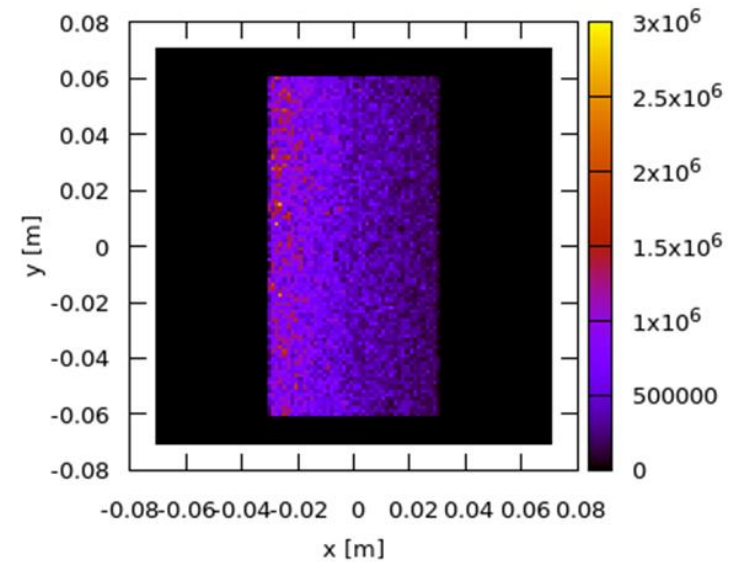
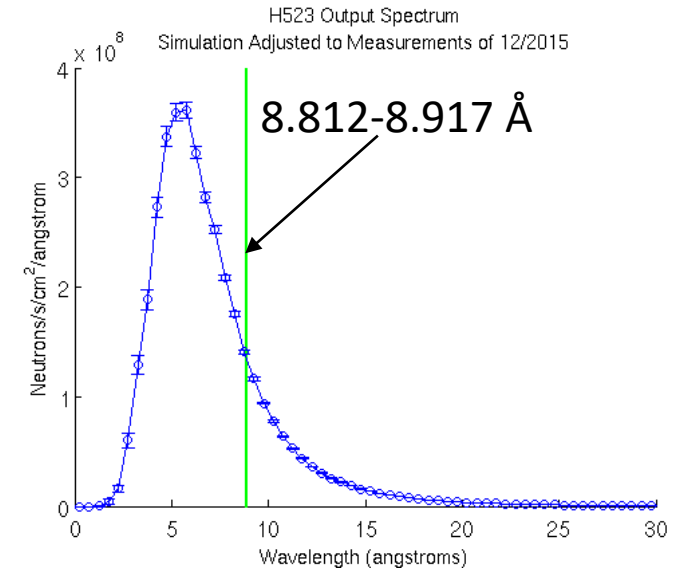
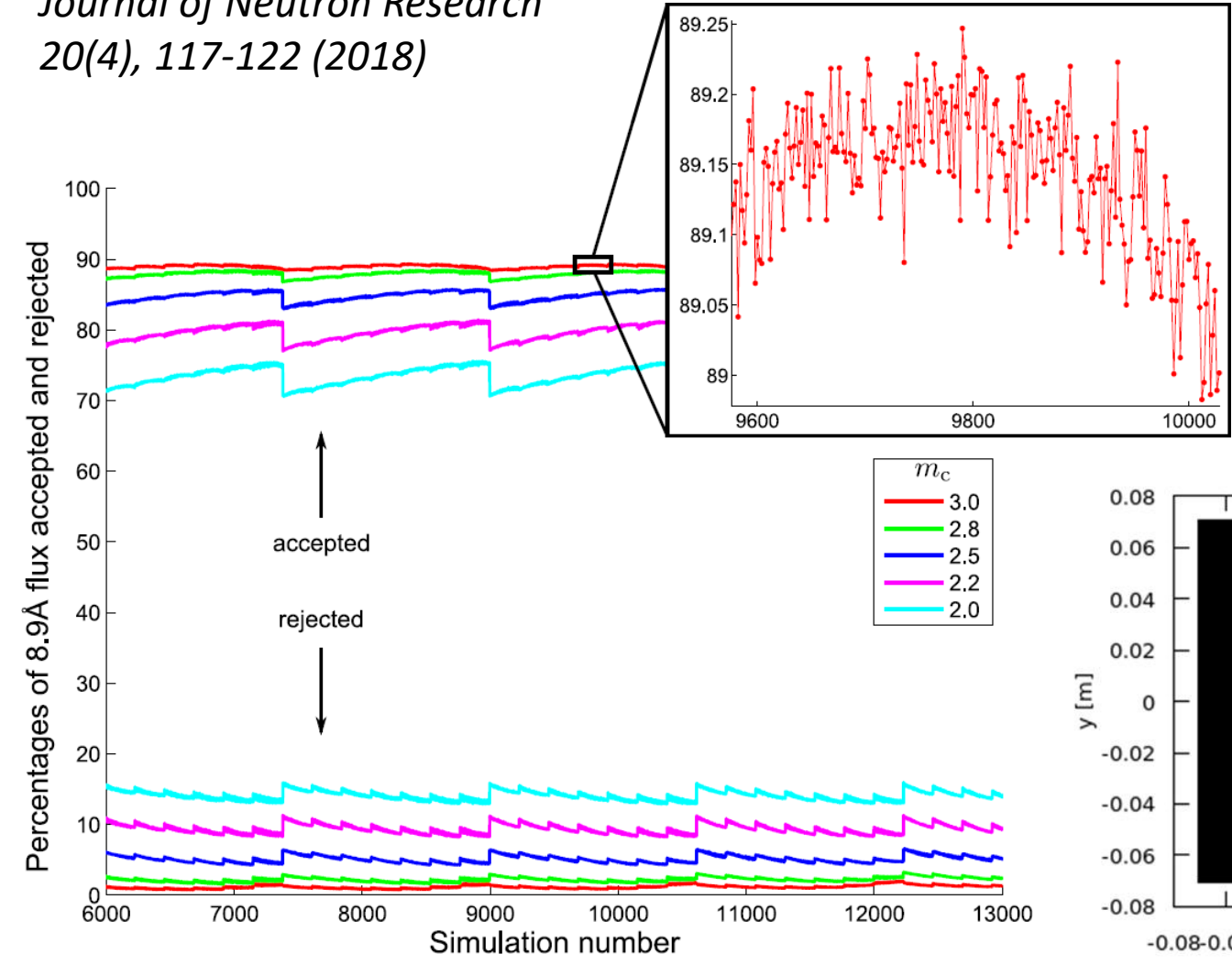
The complete toolset

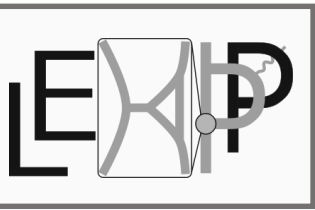




Optimize the remaining parameters...

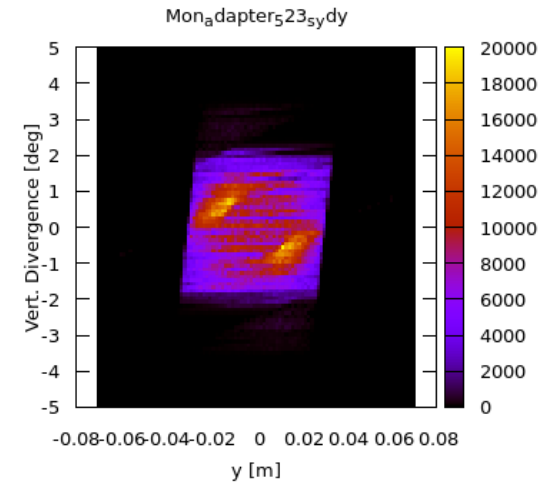
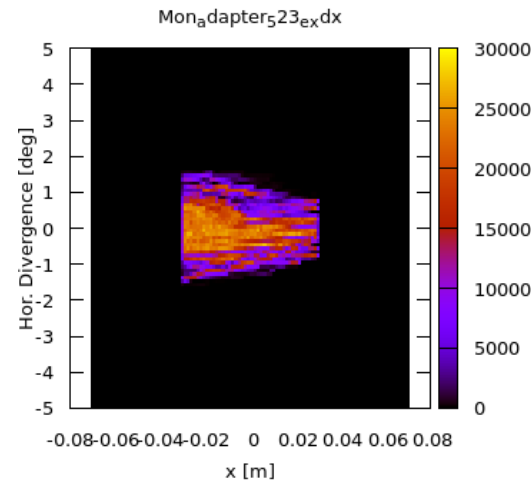
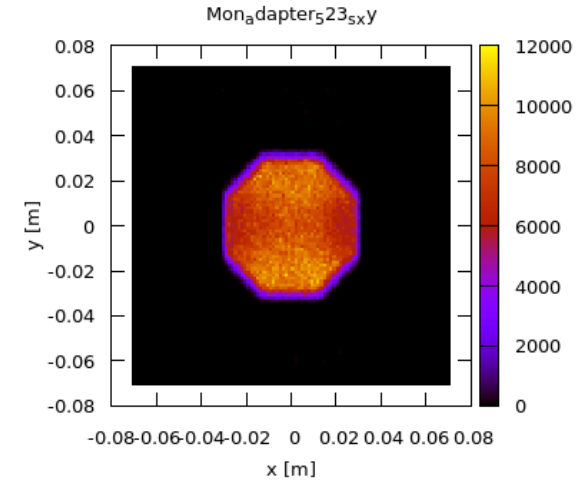
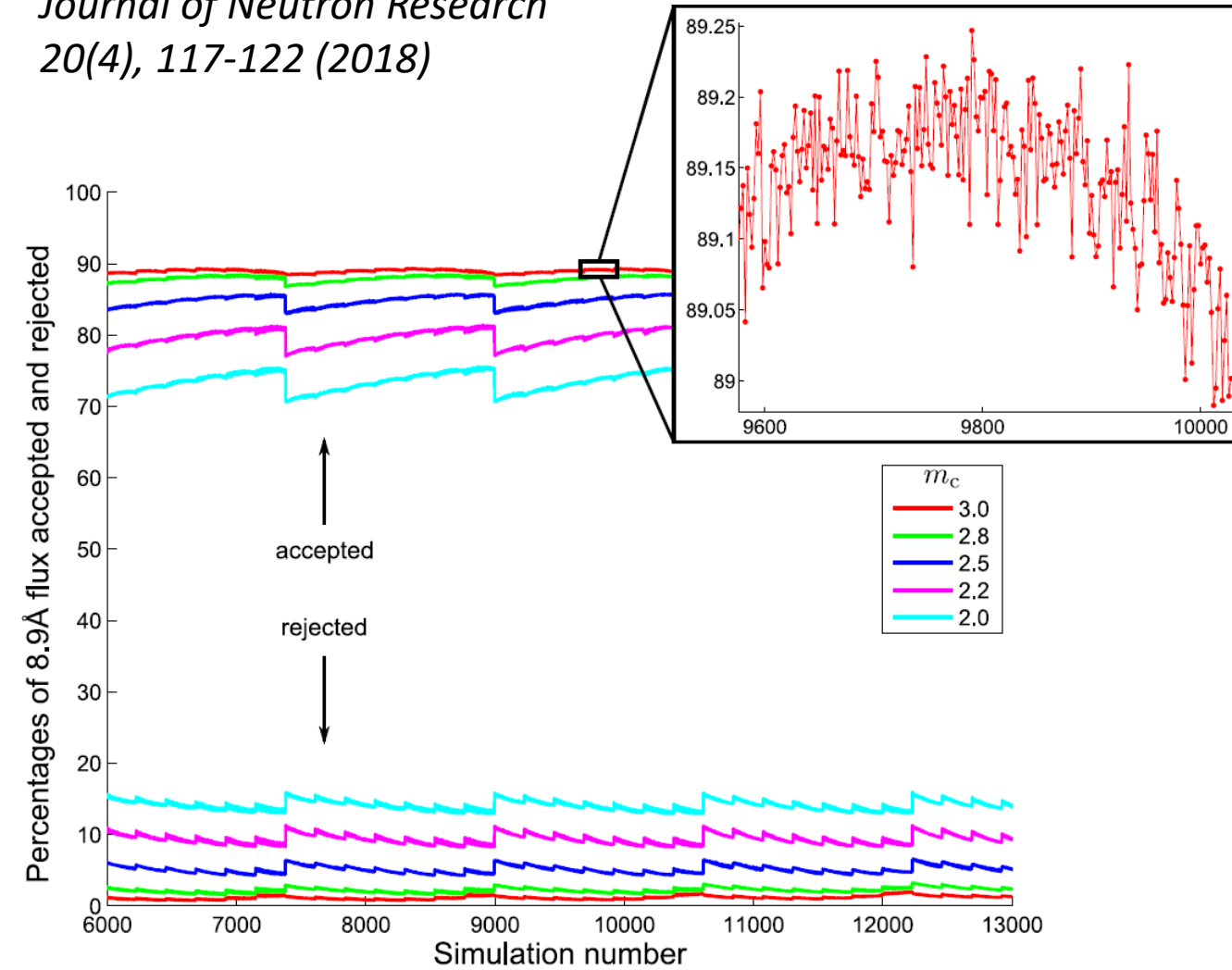
Journal of Neutron Research
20(4), 117-122 (2018)

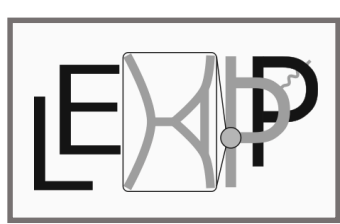




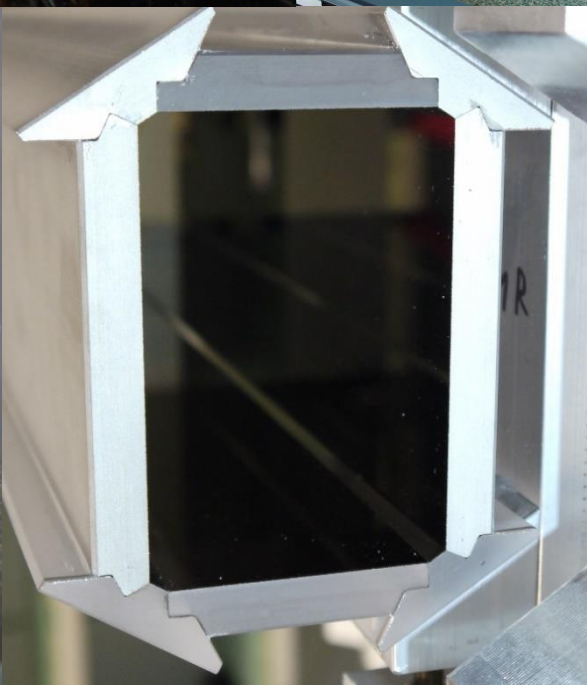
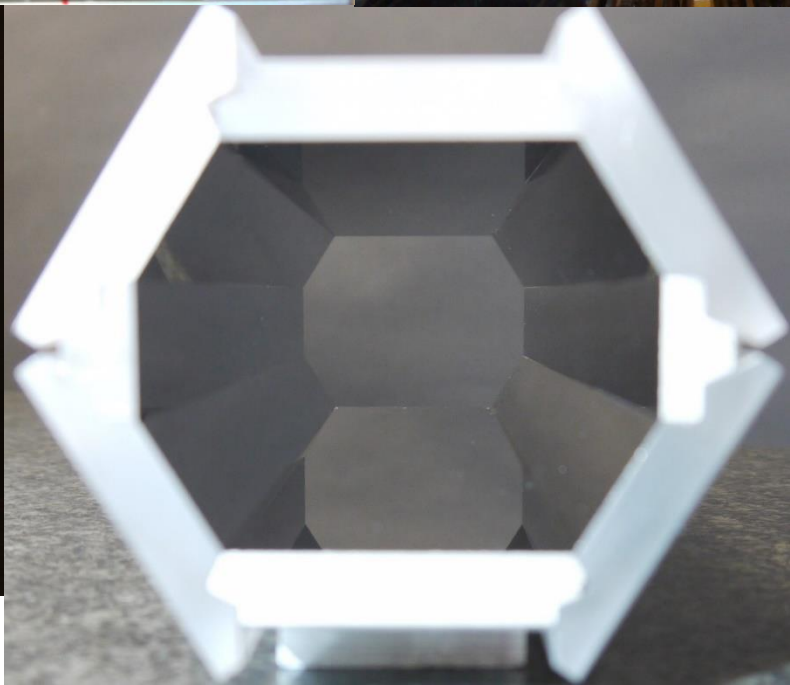
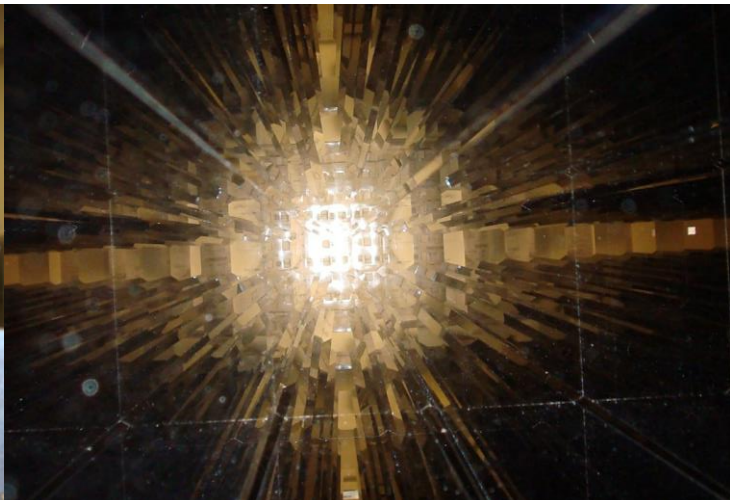
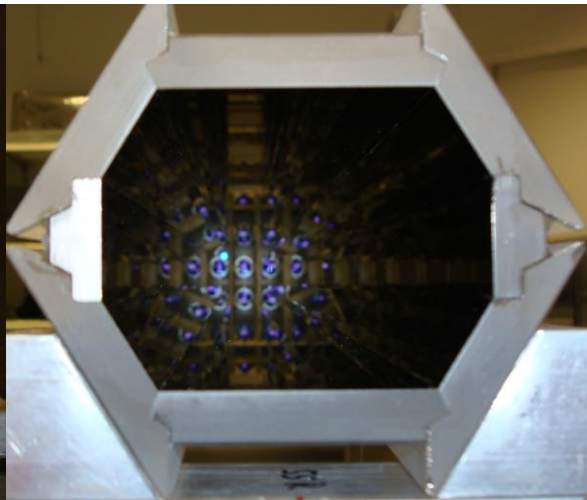
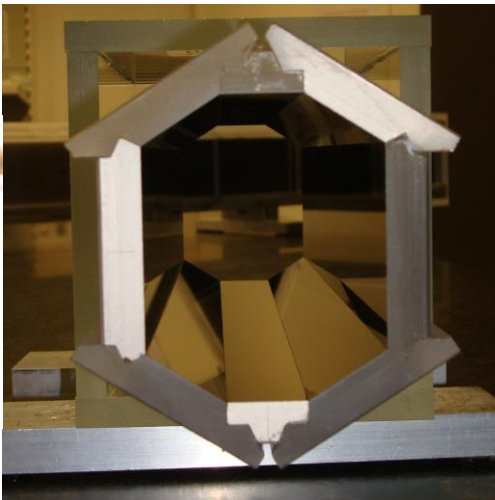
Optimize the remaining parameters...

Journal of Neutron Research
20(4), 117-122 (2018)

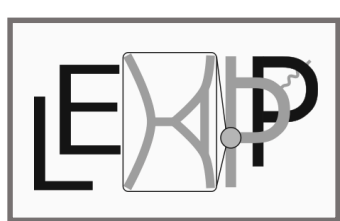




...and implement the solution



S-DH



What would a finite neutron EDM mean?

- CP violation from three sources (ignoring neutrinos):

$$\mathcal{L}_{\text{CPV}} = \mathcal{L}_{\text{CKM}} + \mathcal{L}_{\bar{\theta}} + \mathcal{L}_{\text{BSM}}$$

- CKM CP-violation (Standard Model):

$$\mathcal{L}_{\text{CKM}} = -\frac{ig_2}{\sqrt{2}} \sum_{p,q} V^{pq} \bar{U}_L^p W^+ D_L^q + \text{H.c.}$$

- Strong CP-violation (Standard Model):

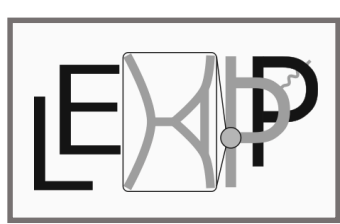
$$\mathcal{L}_{\bar{\theta}} = -\frac{\alpha_S}{16\pi^2} \bar{\theta} \text{Tr}(G^{\mu\nu} \tilde{G}_{\mu\nu})$$

details:

Rev. Mod. Phys. **91**, 015001 (2019)

Phys. Rev. C **91**, 035502 (2015)

Prog. Part. Nucl. Phys. **71**, 21 (2013)



Effective Field Theory

General Effective Lagrangian:

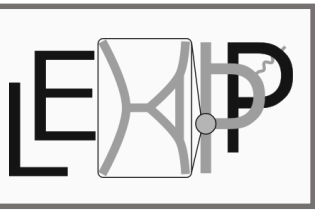
$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{C^{(5)}}{\Lambda} O^{(5)} + \sum_i \frac{C_i^{(6)}}{\Lambda^2} O_i^{(6)} + \dots$$

Dimension-Six terms for the neutron:

$$\begin{aligned} \mathcal{L}_{\text{eff}}^{(6)} = & -\frac{i}{2} \sum_{l,q} d_q \bar{q} \sigma_{\mu\nu} \gamma^5 F^{\mu\nu} q \\ & -\frac{i}{2} \sum_q \tilde{d}_q g_s \bar{q} \sigma_{\mu\nu} \gamma^5 G^{\mu\nu} q \\ & + d_W \frac{g_s}{6} G \tilde{G} G + \sum_i C_i^{(4f)} O_i^{(4f)} \end{aligned}$$

Prog. Part. Nucl. Phys. **71**, 21 (2013)

Wilson coefficient	Operator (dimension)	Number
$\bar{\theta}$	Theta term (4)	1
δ_e	Electron EDM (6)	1
$\text{Im } C_{\ell e q u}^{(1,3)}, \text{Im } C_{\ell e q d}$	Semi-leptonic (6)	3
δ_q	Quark EDM (6)	2
$\tilde{\delta}_q$	Quark chromo EDM (6)	2
$C_{\tilde{G}}$	Three-gluon (6)	1
$\text{Im } C_{quqd}^{(1,8)}$	Four-quark (6)	2
$\text{Im } C_{\varphi ud}$	Induced four-quark (6)	1
Total		13



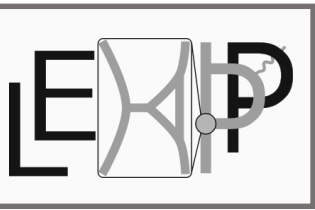
“Global analysis”

Define a matrix α according to $d_i = \hat{a}_i a_{ij} C_j$,

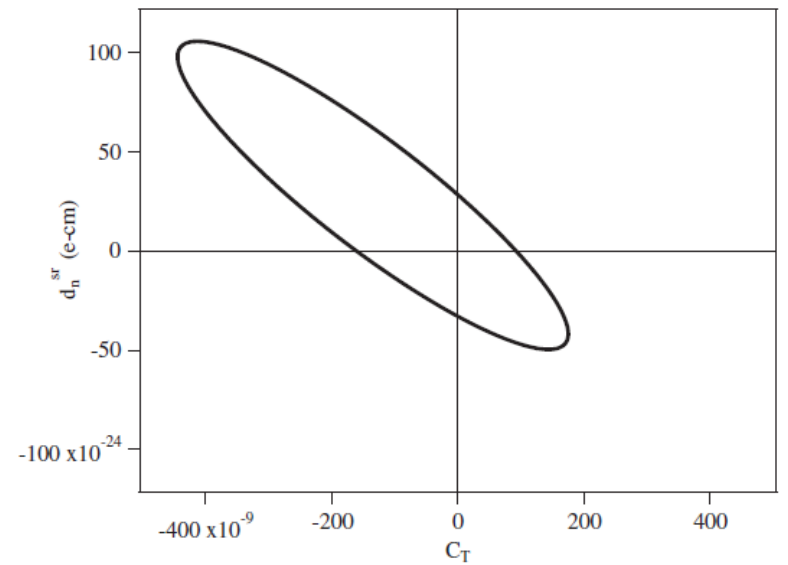
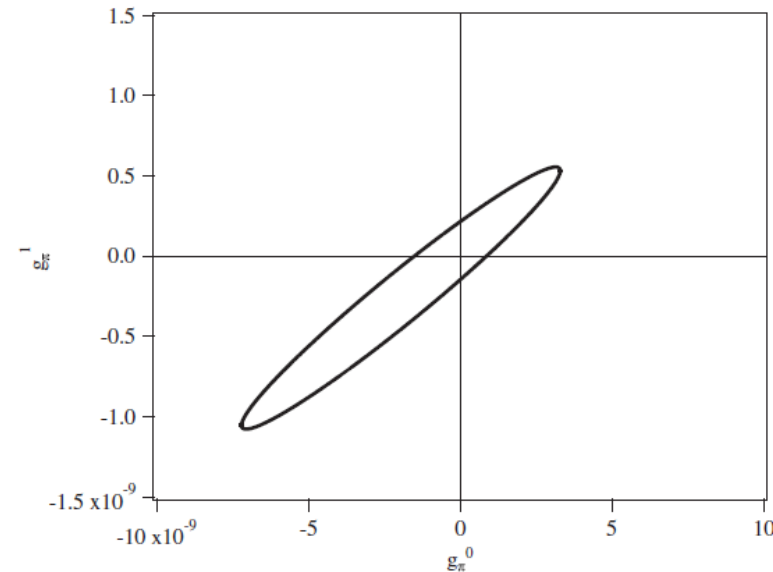
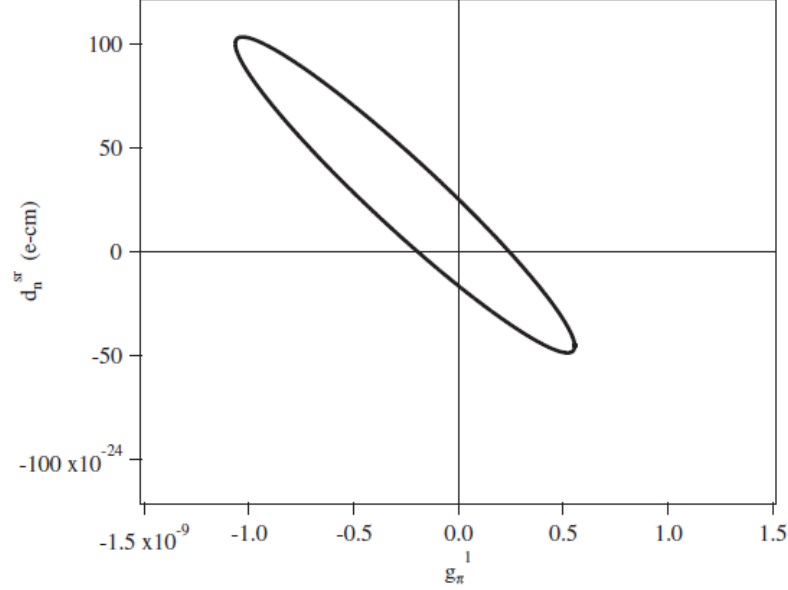
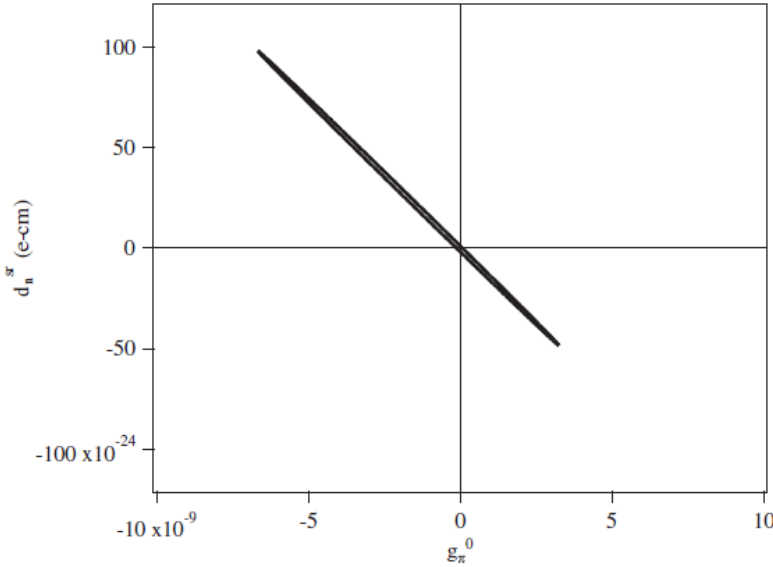
	d_{Hg}	d_{Xe}	d_{TIF}	d_n	
$a_{ij} =$	$-2.0 \cdot 10^{-20}$	$-3.8 \cdot 10^{-18}$	0	0	C_T
\hat{e}	$4.0 \cdot 10^{-21}$	$-2.9 \cdot 10^{-19}$	$-2.2 \cdot 10^{-19}$	0	\tilde{g}_π^0
\hat{e}	$1.1 \cdot 10^{-16}$	$1.2 \cdot 10^{-14}$	$-1.6 \cdot 10^{-13}$	0	\tilde{g}_π^1
\hat{e}	0	$1.5 \cdot 10^{-14}$	$1.4 \cdot 10^{-16}$	1	d_n^{sr}

...and invert it:

$$\begin{bmatrix} C_T \\ \tilde{g}_\pi^0 \\ \tilde{g}_\pi^1 \\ d_n^{sr} \end{bmatrix} = \begin{bmatrix} -1.48 \times 10^{19} & 1.83 \times 10^{20} & -2.52 \times 10^{14} & 0 \\ -1.85 \times 10^{17} & -9.64 \times 10^{17} & 1.32 \times 10^{12} & 0 \\ -2.41 \times 10^{16} & 5.36 \times 10^{16} & -6.32 \times 10^{12} & 0 \\ 2.78 \times 10^3 & 1.44 \times 10^4 & -1.90 \times 10^{-2} & 1 \end{bmatrix} \times \begin{bmatrix} d_{\text{Hg}} \\ d_{\text{Xe}} \\ d_{\text{TIF}} \\ d_n \end{bmatrix}$$

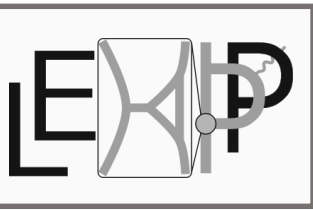


So what is the situation today?



“Sole source” limits:

LE parameter	System	95% u.l.
d_e	ThO	$9.2 \times 10^{-29} e \text{ cm}$
C_S	ThO	8.6×10^{-9}
C_T	^{199}Hg	3.6×10^{-10}
$\bar{g}_\pi^{(0)}$	^{199}Hg	3.8×10^{-12}
$\bar{g}_\pi^{(1)}$	^{199}Hg	3.8×10^{-13}
$\bar{g}_\pi^{(2)}$	^{199}Hg	2.6×10^{-11}
\bar{d}_n^{sr}	Neutron	$3.3 \times 10^{-26} e \text{ cm}$
\bar{d}_p^{sr}	TIF	$8.7 \times 10^{-23} e \text{ cm}$
\bar{d}_p^{sr}	^{199}Hg	$2.0 \times 10^{-25} e \text{ cm}$
Other parameters		
d_d	$\approx 3/4 d_n$	$2.5 \times 10^{-26} e \text{ cm}$
$\bar{\theta}$	$\approx \bar{g}_\pi^{(0)} / (0.015)$	2.5×10^{-10}
$\tilde{d}_d - \tilde{d}_u$	$5 \times 10^{-15} \bar{g}_\pi^{(1)} e \text{ cm}$	$2 \times 10^{-27} e \text{ cm}$



So what is the situation today?

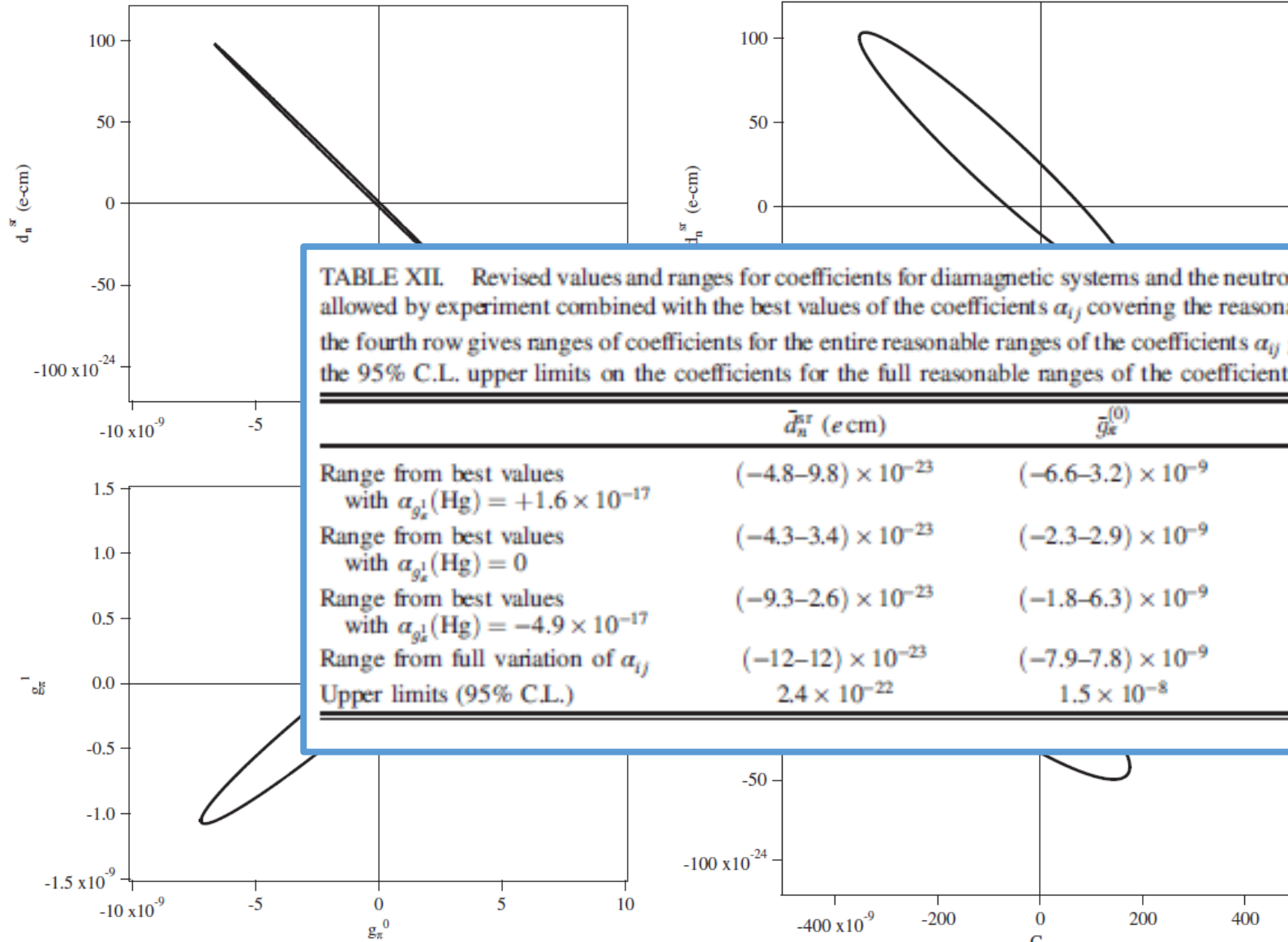
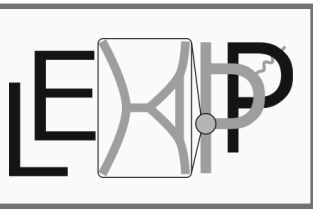
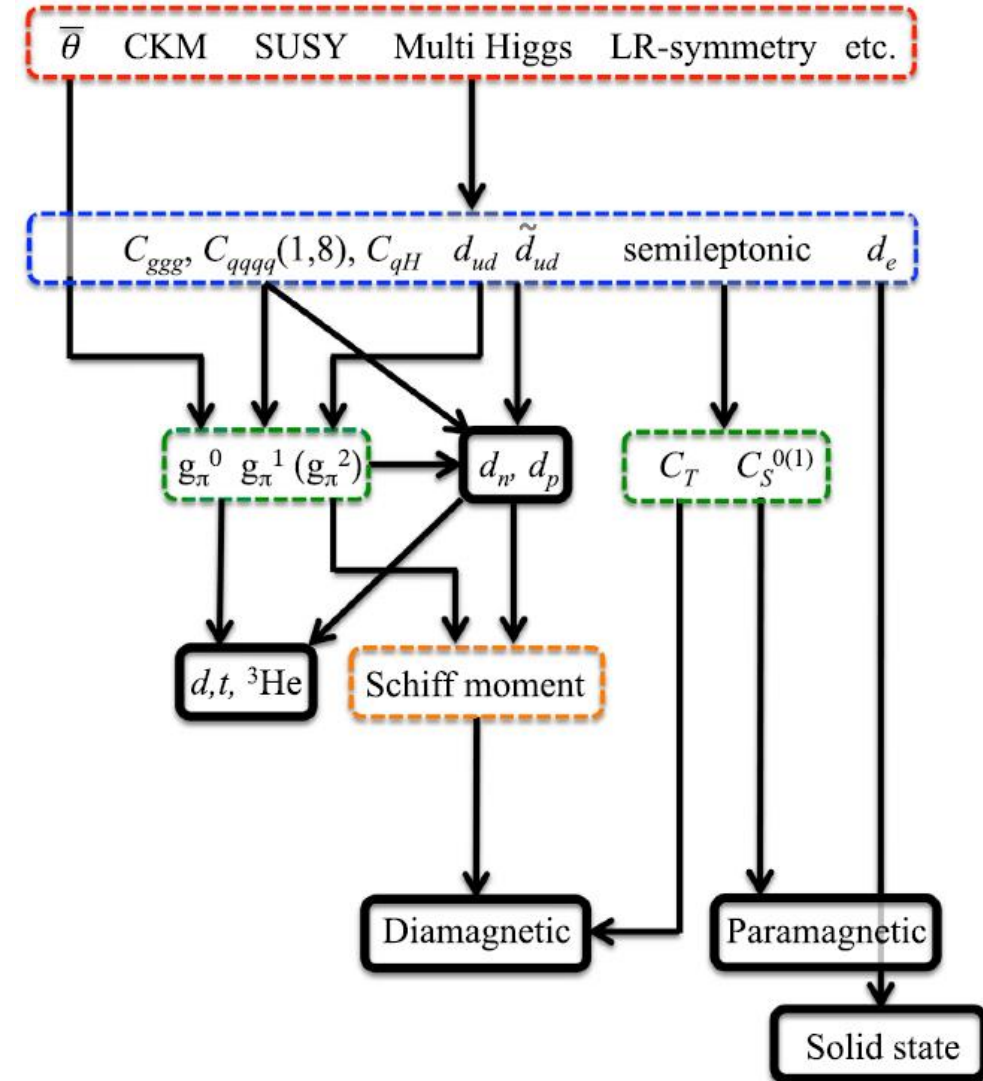
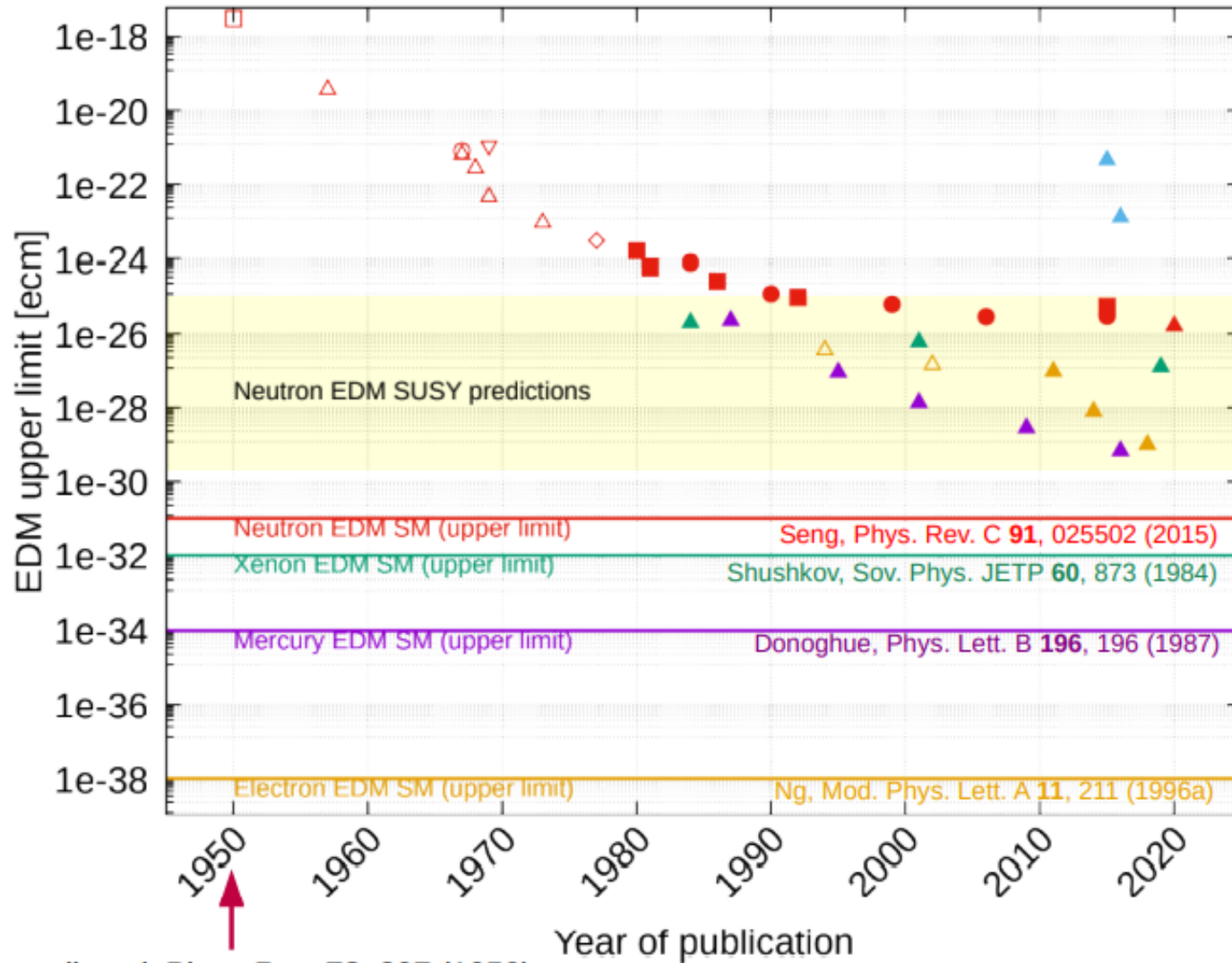


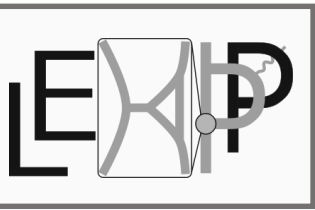
TABLE XII. Revised values and ranges for coefficients for diamagnetic systems and the neutron. The first three rows give the 68% C.L. range allowed by experiment combined with the best values of the coefficients α_{ij} covering the reasonable range of $\alpha_{\text{Hg}, \bar{g}_x^{(1)}}$ with $\alpha_{\text{Ra}, \bar{g}_x} = -8 \times 10^{-4}$; the fourth row gives ranges of coefficients for the entire reasonable ranges of the coefficients α_{ij} given in Table IV, and the bottom row presents the 95% C.L. upper limits on the coefficients for the full reasonable ranges of the coefficients.

	$\bar{d}_n^{\text{sr}} (e \text{ cm})$	$\bar{g}_x^{(0)}$	$\bar{g}_x^{(1)}$	$C_T^{(0)}$
Range from best values with $\alpha_{g_x^1}(\text{Hg}) = +1.6 \times 10^{-17}$	$(-4.8-9.8) \times 10^{-23}$	$(-6.6-3.2) \times 10^{-9}$	$(-1.0-0.5) \times 10^{-9}$	$(-3.5-1.6) \times 10^{-7}$
Range from best values with $\alpha_{g_x^1}(\text{Hg}) = 0$	$(-4.3-3.4) \times 10^{-23}$	$(-2.3-2.9) \times 10^{-9}$	$(-0.6-1.3) \times 10^{-9}$	$(-3.2-4.0) \times 10^{-7}$
Range from best values with $\alpha_{g_x^1}(\text{Hg}) = -4.9 \times 10^{-17}$	$(-9.3-2.6) \times 10^{-23}$	$(-1.8-6.3) \times 10^{-9}$	$(-1.2-0.4) \times 10^{-9}$	$(-11-3.8) \times 10^{-7}$
Range from full variation of α_{ij}	$(-12-12) \times 10^{-23}$	$(-7.9-7.8) \times 10^{-9}$	$(-1.3-1.1) \times 10^{-9}$	$(-6.6-4.6) \times 10^{-7}$
Upper limits (95% C.L.)	2.4×10^{-22}	1.5×10^{-8}	2.4×10^{-9}	1.1×10^{-6}



What does it mean, if we see it?





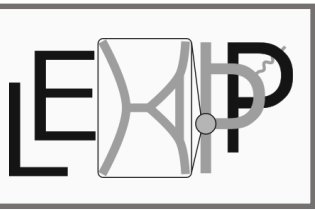
It's not so simple after all...

- **Schiff's theorem:** the field due to an EDM induces a displacement of the bound charges, which exactly cancels it*

$$H_0 = \sum \frac{p^2}{2m} + U(\mathbf{r})$$

Hamiltonian of the charge-system (no EDM)

*Schiff: *Phys. Rev.* **132**, 2194 (1963)
J. Engel: elegant formulation used here



It's not so simple after all...

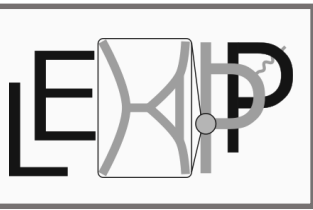
- **Schiff's theorem:** the field due to an EDM induces a displacement of the bound charges, which exactly cancels it

$$H_0 = \sum \frac{p^2}{2m} + U(\mathbf{r})$$

*Add constituent EDMs
As a perturbation...*

$$\mathbf{d}_{\text{tot}} = \sum_i \mathbf{d}_i$$

(sum over constituents)



It's not so simple after all...

- **Schiff's theorem:** the field due to an EDM induces a displacement of the bound charges, which exactly cancels it

$$H_0 = \sum \frac{p^2}{2m} + U(\mathbf{r})$$

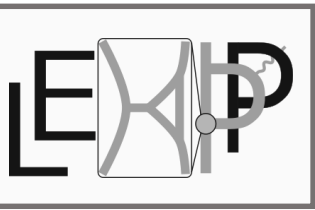
*Add constituent EDMs
As a perturbation...*

$$\mathbf{d}_{\text{tot}} = \sum_i \mathbf{d}_i$$

(sum over constituents)

$$\begin{aligned} H &= H_0 - \sum \mathbf{d} \cdot \mathbf{E} \\ &= H_0 + \sum \mathbf{d} \cdot \frac{\nabla U(\mathbf{r})}{q} \\ &= H_0 + \sum \frac{i}{q} [\mathbf{d} \cdot \mathbf{p}, H_0] \end{aligned}$$

Now see what effect this has...



It's not so simple after all...

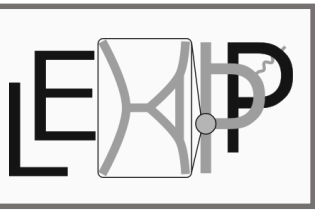
- **Schiff's theorem:** the field due to an EDM induces a displacement of the bound charges, which exactly cancels it

$$H_0 = \sum \frac{p^2}{2m} + U(\mathbf{r})$$

$$\begin{aligned} H &= H_0 - \sum \mathbf{d} \cdot \mathbf{E} \\ &= H_0 + \sum \mathbf{d} \cdot \frac{\nabla U(\mathbf{r})}{q} \\ &= H_0 + \sum \frac{i}{q} [\mathbf{d} \cdot \mathbf{p}, H_0] \end{aligned}$$

Eigenstates receive an energy shift due to the perturbation:

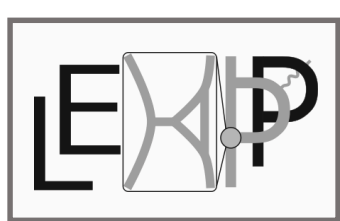
$$\begin{aligned} |0\rangle \rightarrow |\tilde{0}\rangle &= |0\rangle + \sum_n \frac{|n\rangle \langle n| \sum \frac{i}{q} [\mathbf{d} \cdot \mathbf{p}, H_0] |0\rangle}{E_0 - E_n} \\ &= \left(1 + \sum \frac{i}{q} \mathbf{d} \cdot \mathbf{p} \right) |0\rangle \end{aligned}$$



It's not so simple after all...

- What is the total, observable, dipole moment after this shift?

$$\begin{aligned}\tilde{\mathbf{d}} &= \sum \mathbf{d} + \langle \tilde{0} | \sum q\mathbf{r} | \tilde{0} \rangle \\ &= \sum \mathbf{d} + \langle \tilde{0} | \left(1 - \sum \frac{i}{q} \mathbf{d} \cdot \mathbf{p} \right) \sum q\mathbf{r} \left(1 + \sum \frac{i}{q} \mathbf{d} \cdot \mathbf{p} \right) | \tilde{0} \rangle \\ &= \sum \mathbf{d} + i \langle 0 | \left[\sum q\mathbf{r}, \sum \frac{1}{q} \mathbf{d} \cdot \mathbf{p} \right] | 0 \rangle \\ &= \sum \mathbf{d} - \sum \mathbf{d} \\ &= 0\end{aligned}$$



But some details can save us!

- Schiff's theorem assumes:

- pointlike particles → *incorrect for nuclei*

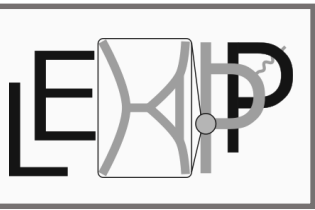
$$\mathbf{S} = \frac{1}{10} \langle r^2 \mathbf{d} \rangle - \frac{1}{6Z} \langle r^2 \rangle \langle \mathbf{d} \rangle$$

...see Prog. Part. Nucl. Phys. **71**, 21 (2013)

- non-relativistic treatment → *incorrect for atomic electrons*

$$U_{\text{lab}} = -\mathbf{d}_{\text{lab}} \cdot \mathbf{E} = -\mathbf{d}_{\text{rest}} \cdot \mathbf{E} + \frac{\gamma}{1 + \gamma} (\boldsymbol{\beta} \cdot \mathbf{d})(\boldsymbol{\beta} \cdot \mathbf{E})$$

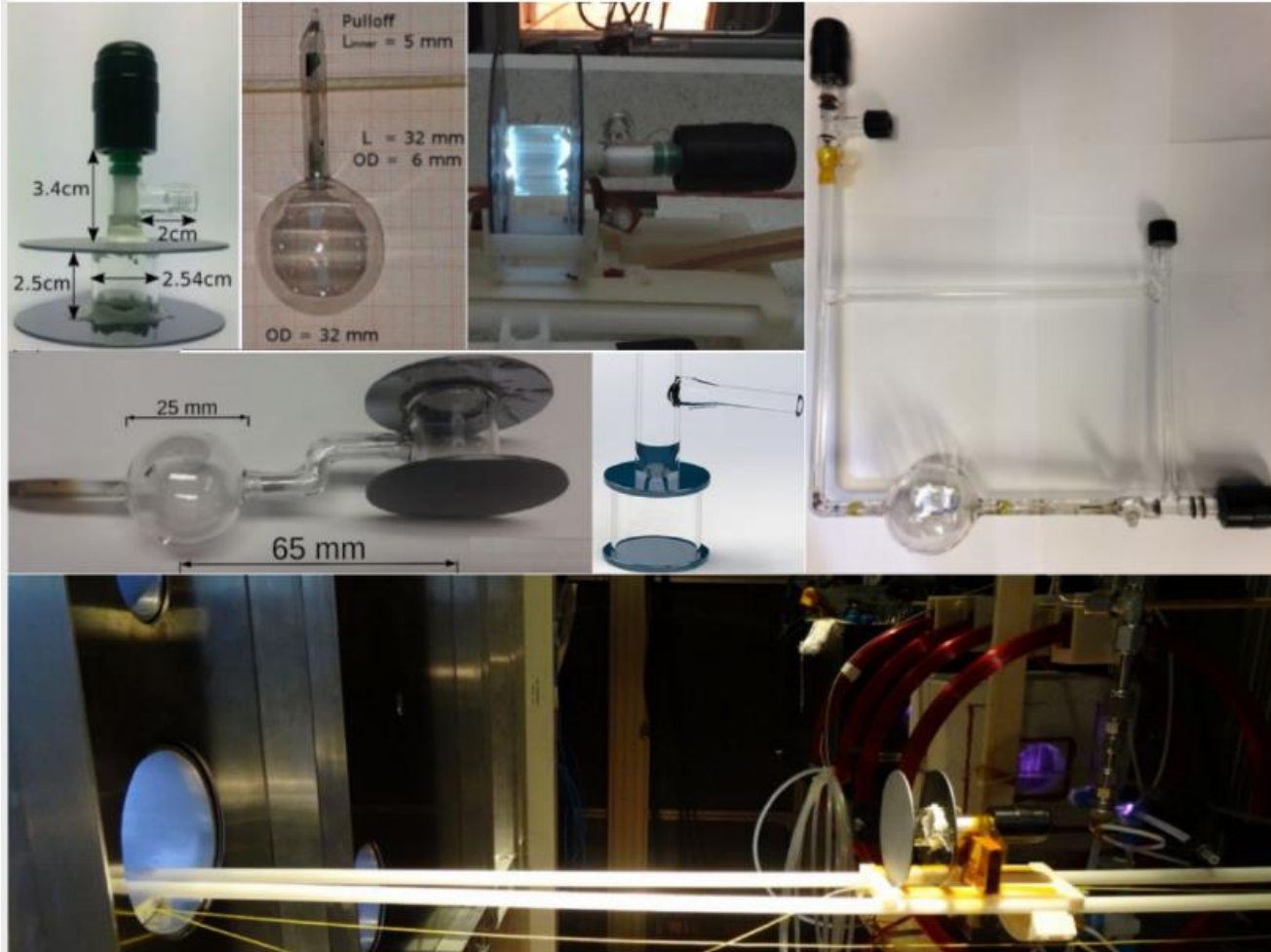
...see American Journal of Physics **75**, 532 (2007)



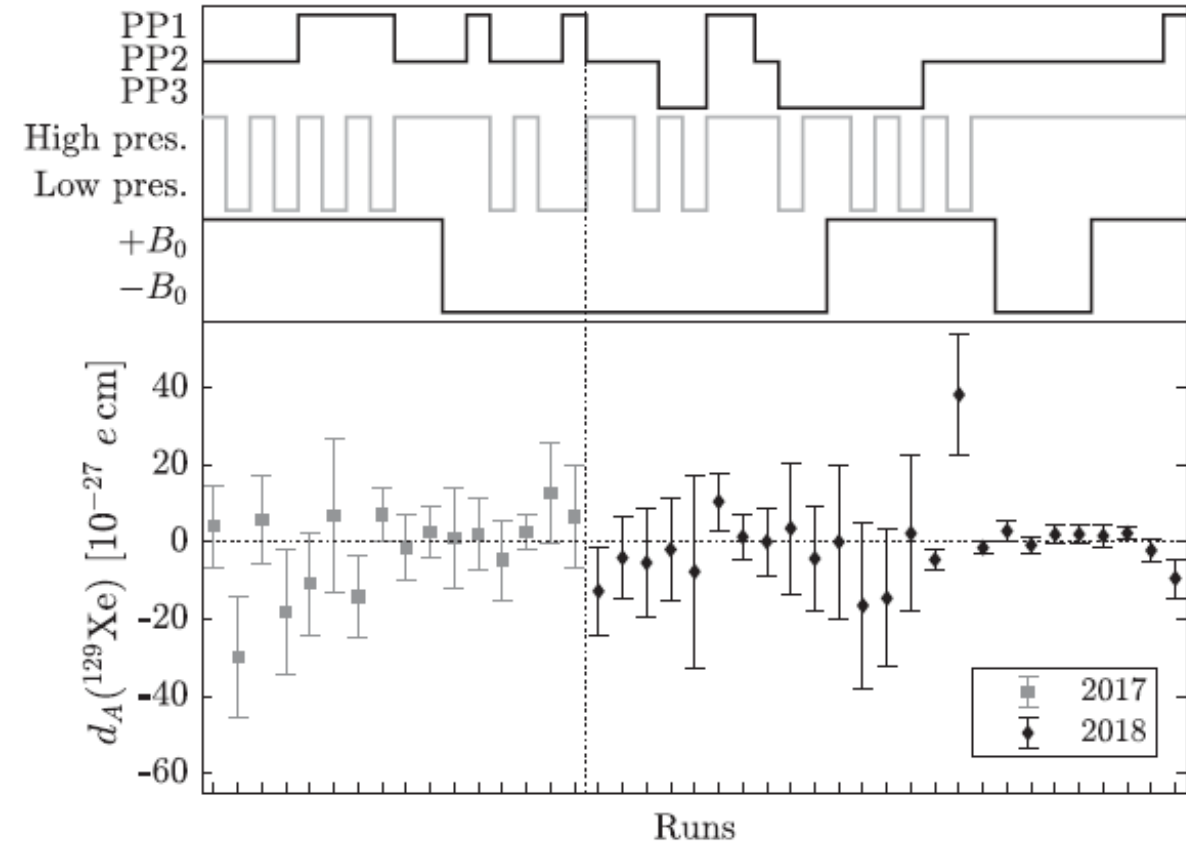
More on HeXe and “complementarity”

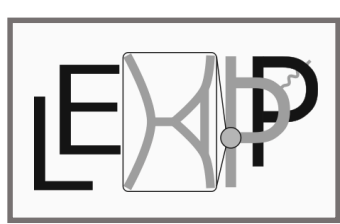


Cells, spin-polarization, co-magnetometry



Analysis methods





HeXe EDM @TUM/PTB...

...laid groundwork for SuperSUN/PanEDM @ILL

- Gave a reference point for what it takes to go from concept to result in a new type of EDM experiment... obviously less complex than nEDM, but still with significant equipment and technique overlap

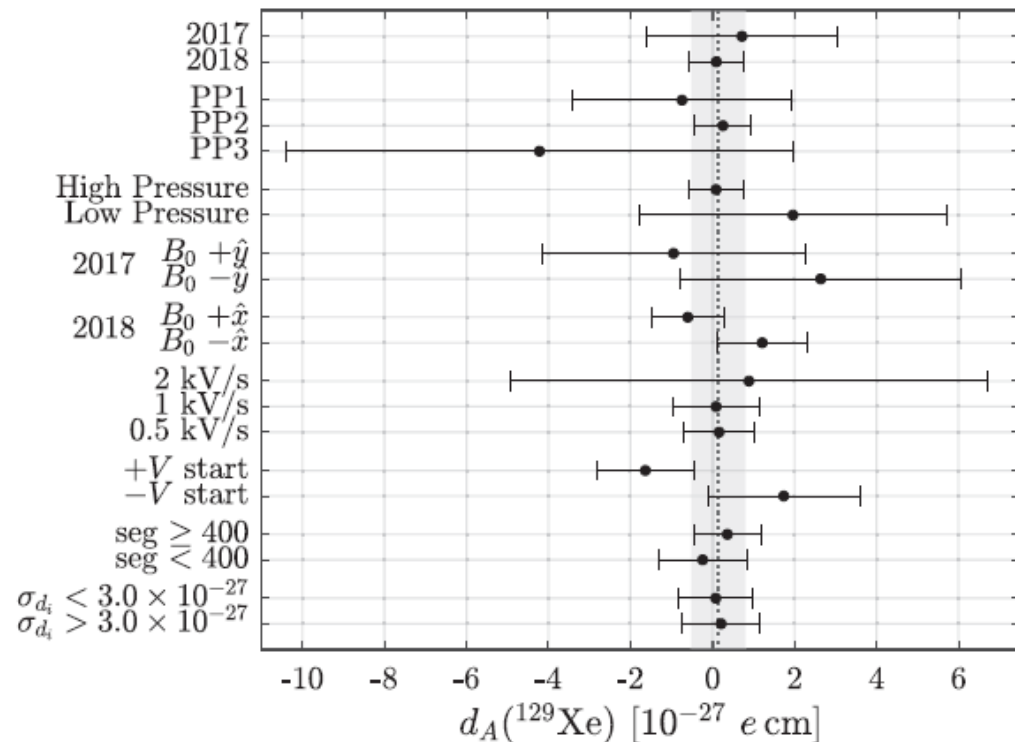
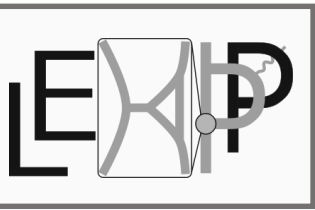


TABLE I. Summary of EDM results and systematic effects discussed in the text.

	2017 ($e\text{ cm}$)	2018 ($e\text{ cm}$)
<i>EDM</i>	7.2×10^{-28}	0.9×10^{-28}
<i>Statistical error</i>	23.5×10^{-28}	6.8×10^{-28}
<i>Systematic Source</i>		
Leakage current	1.2×10^{-28}	4.5×10^{-31}
Charging currents	1.7×10^{-29}	1.2×10^{-29}
Cell motion (rotation)	4.2×10^{-29}	4.0×10^{-29}
Cell motion (translation)	2.6×10^{-28}	1.9×10^{-28}
Comagnetometer drift	2.6×10^{-28}	4.0×10^{-29}
$ \vec{E} ^2$ effects	1.2×10^{-29}	2.2×10^{-30}
$ \vec{E} $ uncertainty	2.6×10^{-29}	9.4×10^{-30}
Geometric phase	$\leq 2 \times 10^{-31}$	$\leq 2 \times 10^{-31}$
Total Systematic Error	3.9×10^{-28}	2.0×10^{-28}



Measuring the ^{129}Xe EDM at Heidelberg PI

