

Simple self-consistent prediction methods for the phase space of dark matter: from galactic dynamics to phenomenology

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Phase-space distribution of DM & theoretical uncertainties for direct and indirect searches

Direct searches

$$\frac{dR}{dE} \propto \rho_{\odot} \int_{v_{\min} \leq |\vec{v}| \leq v_{\text{esc}}} \frac{f_{\odot}(\vec{v})}{|\vec{v}|} d^3v$$

Impact at low masses

$$v_{\min} \sim v_{\text{esc}}$$

Speed-dependent annihilation

- $\langle \sigma v \rangle(r) \propto \langle v_r^2 \rangle$ **p-wave**
 $\sigma v = \sigma_1 v_r^2 \Rightarrow \langle \sigma v \rangle \equiv \langle \sigma v \rangle(r)$

$$\langle \sigma v \rangle(r) = \sigma_1 \int d^3v_1 d^3v_2 f_r(\vec{v}_1) f_r(\vec{v}_2) v_r^2$$

- $\langle \sigma v \rangle(r) \propto \langle 1/v_r \rangle$ or $\langle 1/v_r^2 \rangle$

Sommerfeld enhancement

Primordial black holes

- Gravitational microlensing event rates (EROS, MACHO)

$$\frac{d\Gamma}{dt} \propto \rho(r) \int v f(\vec{v}, \vec{r}) d^3v$$

- Merger rates (gravitational waves)

+ DM substructures (test masses), disruption of stellar binaries

Standard approaches 1: "Standard halo model"

Standard halo model (SHM)

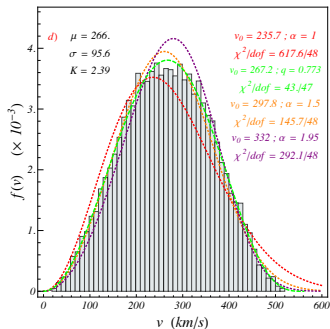
Maxwell-Boltzmann distribution

$$f(\vec{v}) = \frac{1}{v_c^3 \pi^{3/2}} e^{-\left(\frac{\vec{v}}{v_c}\right)^2}$$

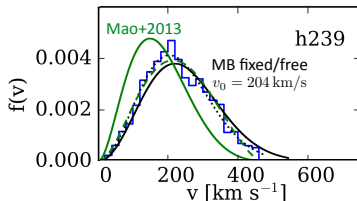
Oversimplification

- Isothermal sphere
- Infinite system
- Ad hoc truncation at v_{esc}

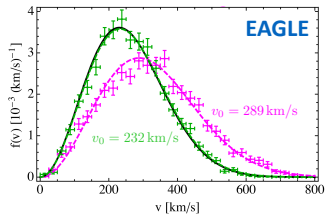
Standard approaches 2: direct fits to simulations



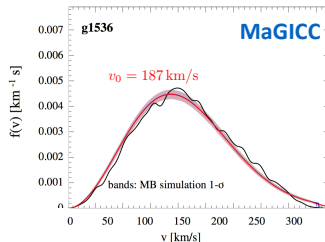
Ling+ 2010, Mollitor+ 2014



Kelso+ 2016



Bozorgnia+ 2016

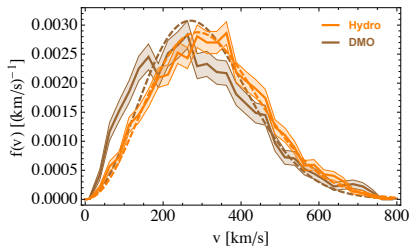


Sloane+ 2016

Standard approaches 2: direct fits to simulations

General insight

Generic features found in simulations (e.g., cusp/cores)



Bozorgnia+ 2017

But insufficient approach

- Extrapolations based on fits at 8 kpc
- **Peak speed free parameter**
⇒ not connected to circular speed
- MW one particular realization
- MW constrained system (e.g., Gaia)
- Subgrid physics

Self-consistent approach required

Eddington-like methods: next-to-minimal approach

Phase space of dark matter from first principles

Phase-space distribution $f(\vec{v}, \vec{r})$: closed system

- Collisionless Boltzmann equation, steady state

$$\{f, H\} = \vec{v} \cdot \frac{\partial f}{\partial \vec{r}} - \frac{\partial \Phi}{\partial \vec{r}} \cdot \frac{\partial f}{\partial \vec{v}} = 0$$

→ Jeans' theorem: $f \equiv f(I_1, \dots, I_N)$ where $\{I_i, H\} = 0$

- Poisson equation

$$\Delta \Phi = 4\pi G \rho \quad \text{with} \quad \rho = \int f(\vec{v}, \vec{r}) d^3v$$

Eddington's inversion formula (Eddington 1916)

Isotropic system + spherical symmetry: $f(\vec{v}, r) \equiv f(\mathcal{E})$

with $\mathcal{E} = \Psi(r) - \frac{v^2}{2}$ and $\Psi(r) = \Phi(R_{\max}) - \Phi(r)$

$$f(\mathcal{E}) = \frac{1}{\sqrt{8\pi^2}} \left[\frac{1}{\sqrt{\mathcal{E}}} \left(\frac{d\rho}{d\Psi} \right)_{\Psi=0} + \int_0^{\mathcal{E}} \frac{d^2\rho}{d\Psi^2} \frac{d\Psi}{\sqrt{\mathcal{E} - \Psi}} \right]$$

+ anisotropic extensions

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Velocity distribution

Central ingredient for observables

$$f_{\vec{v}}(\vec{v}, r) \equiv \frac{f(\mathcal{E}, L)}{\rho_{\text{DM}}(r)}$$

Speed distribution ($v = |\vec{v}|$)

$$f_v(v, r) \equiv v^2 \int d\Omega_v f_{\vec{v}}(\vec{v}, r)$$

Encapsulates most of the dynamical information

For isotropic distribution

$$f_v(v, r) = \frac{4\pi v^2}{\rho_{\text{DM}}(r)} f\left(\Psi(r) - \frac{v^2}{2}\right)$$

Going beyond spherical symmetry

Angle-action coordinates

- More suitable coordinate system if no spherical symmetry
Binney & Tremaine 1987
- Best way to account for complexity revealed by Gaia
- But very cumbersome calculations
- Ansatz for $f(\vec{r}, \vec{v}) \Rightarrow$ theoretical uncertainties
- Level of refinement not necessarily required for DM searches
→ Evaluate astrophysical uncertainties

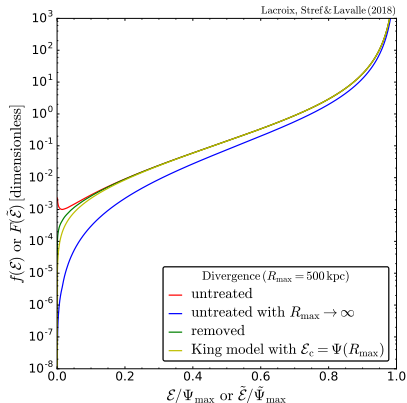
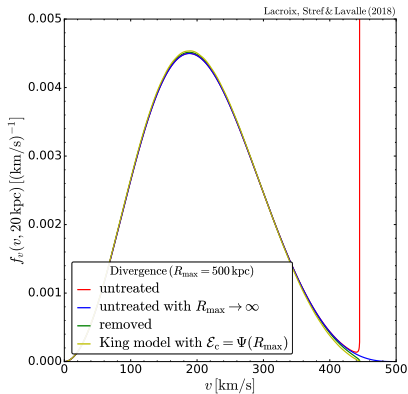
Eddington's formalism

- Lower level of technicalities to account for dynamical constraints
- Method applied blindly to direct searches so far
→ Timely to study validity range in detail

Theoretical consistency and radial boundary

Imposing a radial boundary

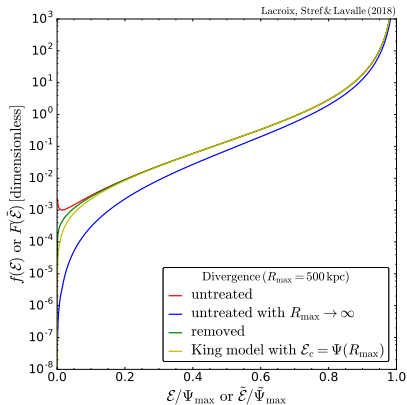
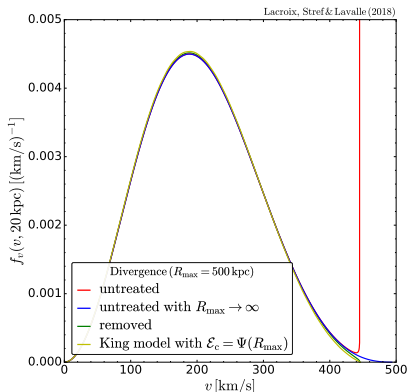
- Finite system (R_{\max}) \Rightarrow divergence of $f(\vec{r}, \vec{v})$ at v_{esc} (from $1/\sqrt{\mathcal{E}}$)
- Phase-space compression
- v_{esc} crucial (direct DM searches at low masses, stellar surveys)



Theoretical consistency and radial boundary

Regularization

- Modified profile, flat at R_{\max}
- Energy cutoff (King)



Lacroix+ 2018a

Not possible for radial anisotropy (e.g., Osipkov-Merritt)

Theoretical consistency: instabilities

Validity range of the method

- Standard criterion:

$$f \geq 0$$

- Antonov instabilities for some DM-baryon configurations

- Stable solution if

$$\frac{df}{d\mathcal{E}} > 0 \Leftrightarrow \frac{d^2\rho}{d\Psi^2} > 0$$

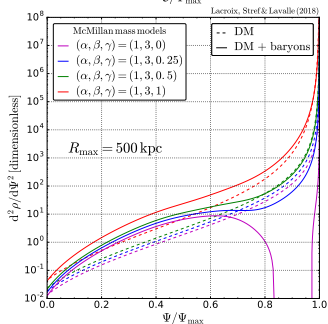
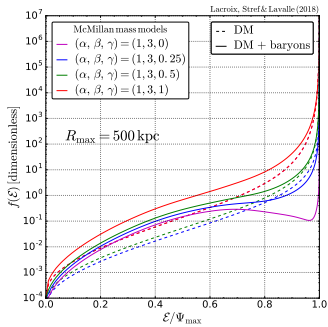
Doremus+ 1971, Kandrup & Sygnet 1985

- Select mass models

Lacroix+ 2018a

For anisotropic systems criteria against radial perturbations only

Doremus+ 1973

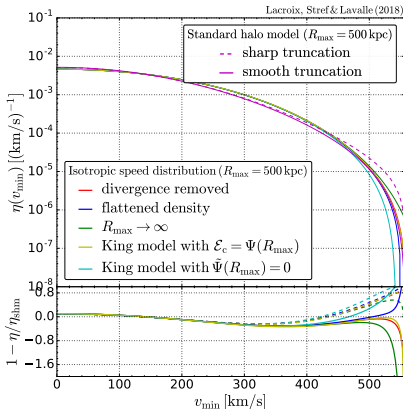


Impact on predictions for direct DM searches

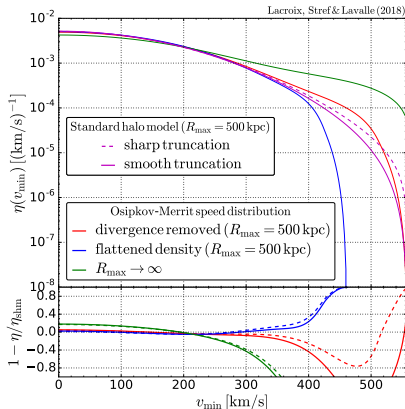
Event rate proportional to

$$\eta(v_{\min}) = \int_{v_{\min} \leq v \leq v_{\oplus} + v_{\text{esc}}} \frac{f_{\vec{v}, \oplus}(\vec{v})}{v} d^3v$$

Isotropic



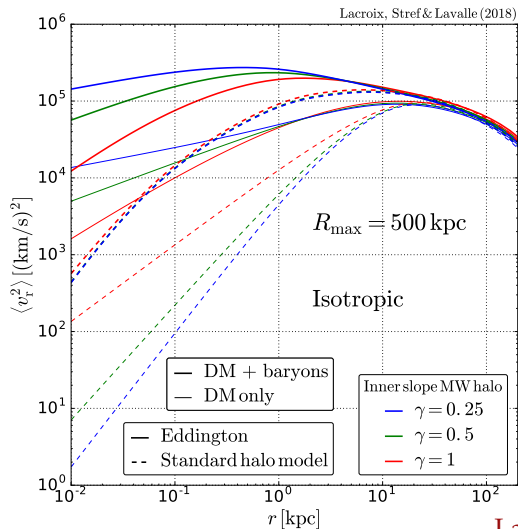
Osipkov-Merritt



Impact on predictions for indirect DM searches

Prototypical case: p-wave annihilation

$$\langle\sigma v\rangle(r) \propto \langle v_r^2\rangle$$



Lacroix+ 2018a

Actual predictivity of Eddington's formalism?

Tests with cosmological simulations

Description

- 2 sets of simulations

Mollitor+ 2015

$M_{\text{DM}} = 2.3 \times 10^5 M_{\odot}$, Hsml = 150 pc

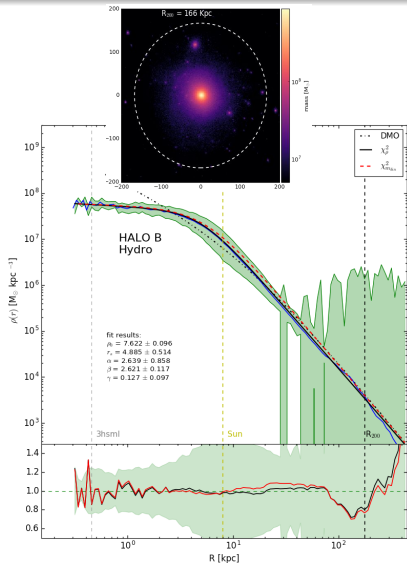
Núñez+ 2018, in prep.

$M_{\text{DM}} = 1.9 \times 10^5 M_{\odot}$, Hsml = 280 pc

- 20 Mpc boxes + zoom-in
- DM-only + hydro

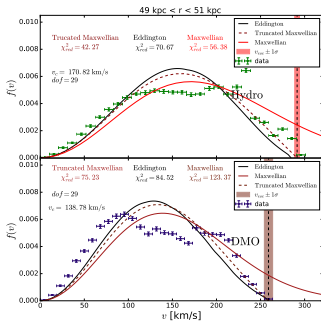
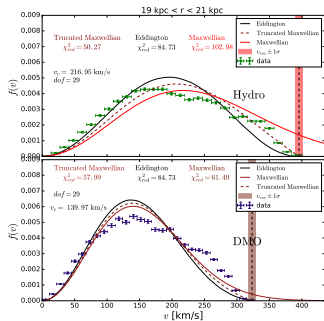
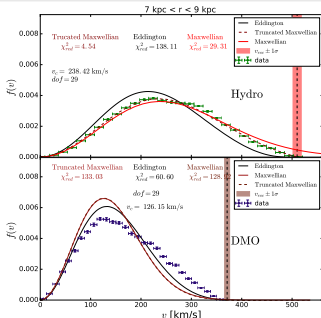
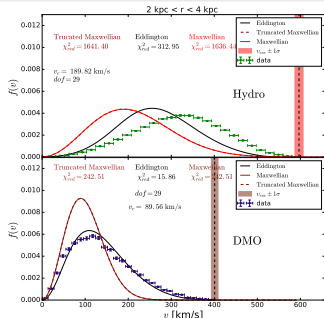
Procedure

- Fit mass model from simulation
 $\Rightarrow \rho_{\text{DM}}, \rho_{\text{B}}, \Psi = \Psi_{\text{DM}} + \Psi_{\text{B}}$
- Input for Eddington's method
- Comparison with simulation outputs



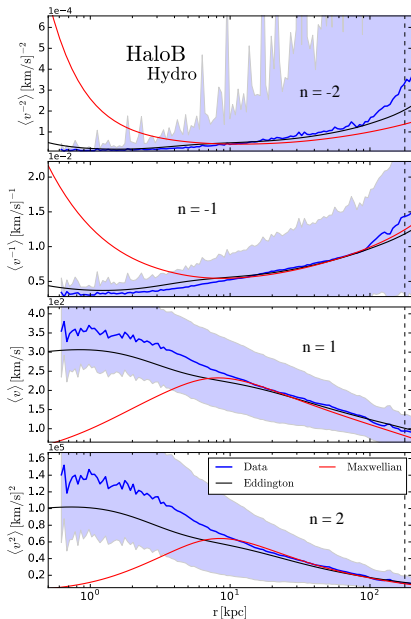
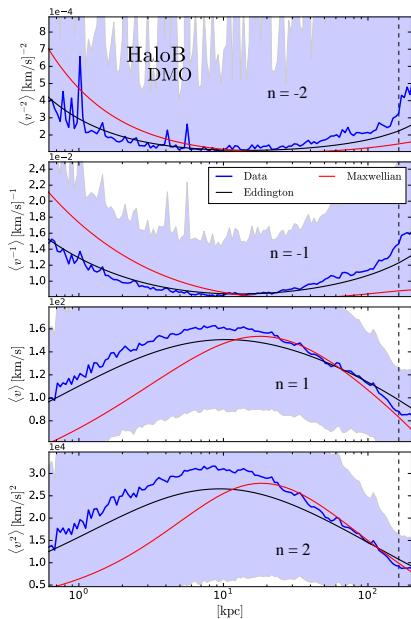
Lacroix+ 2018b, in prep.

Speed distribution $f_v(v, r)$



Lacroix+
2018b, in prep.

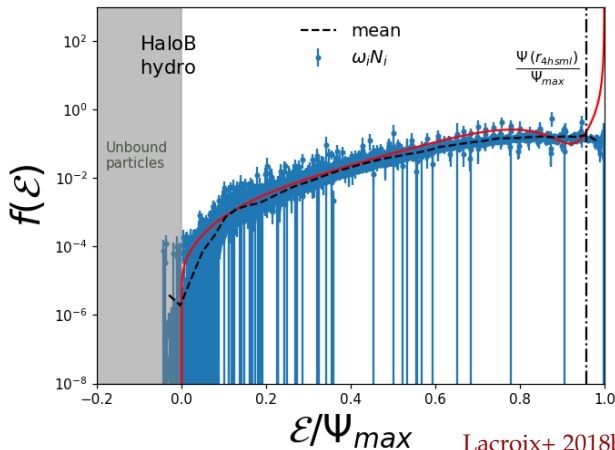
Moments of the speed distribution



Full phase-space distribution

Reconstruction from 2D bins ij in phase space (r_i, v_j)

$$f(\mathcal{E}) = m \frac{d^6 N}{d^3 x d^3 v} \rightarrow f(\mathcal{E})_{ij} = \frac{m}{(4\pi r_i v_j)^2} \frac{N_{ij}}{\Delta r_i \Delta v_j}$$



Lacroix+ 2018b, in prep.

⇒ Very good agreement for isotropic distribution

Application: constraints on sub-GeV DM from cosmic positrons - p -wave annihilation

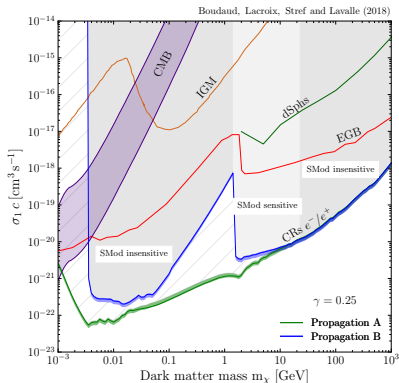
p -wave annihilation

$$\sigma v = \sigma_1 v_{\mathbf{r}}^2 \Rightarrow \langle \sigma v \rangle \equiv \langle \sigma v \rangle(r)$$

$$\langle \sigma v \rangle(r) = \sigma_1 \int d^3 v_1 d^3 v_2 f_{\mathbf{r}}(\vec{v}_1) f_{\mathbf{r}}(\vec{v}_2) v_{\mathbf{r}}^2$$

$$\Rightarrow \psi_e \neq \langle \sigma v \rangle \int \rho^2(r) d^3 r$$

- Very strong e^+ constraints (**Voyager 1**, AMS-02)
- Justifies focusing on Eddington's methods
- Robust w.r.t. uncertainties on anisotropy and propagation



Boudaud+ 2018

Summary

Eddington's inversion method

- A few physical assumptions
- Moderate level of technicalities
- Mass model direct input
- Better control astrophysical uncertainties for DM searches
- Dramatic changes wrt Maxwell-Boltzmann

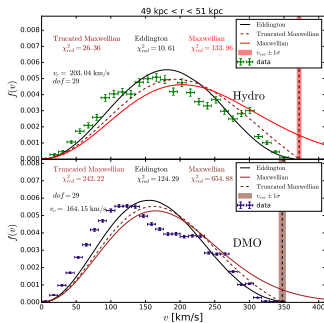
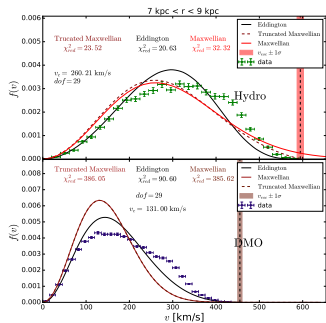
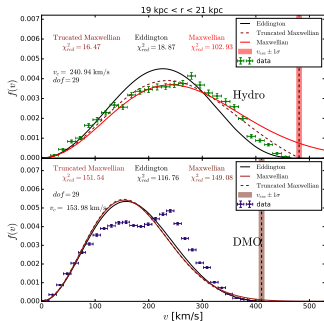
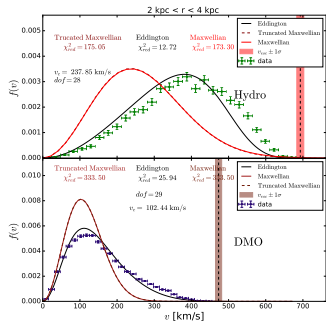
Self-consistency: theoretical validity range

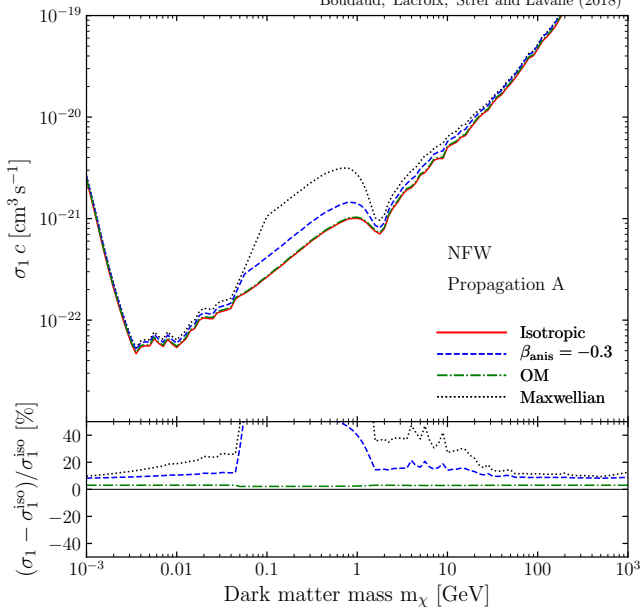
- Radial boundary (direct searches)
- Positive DF + stability

Actual predictivity?

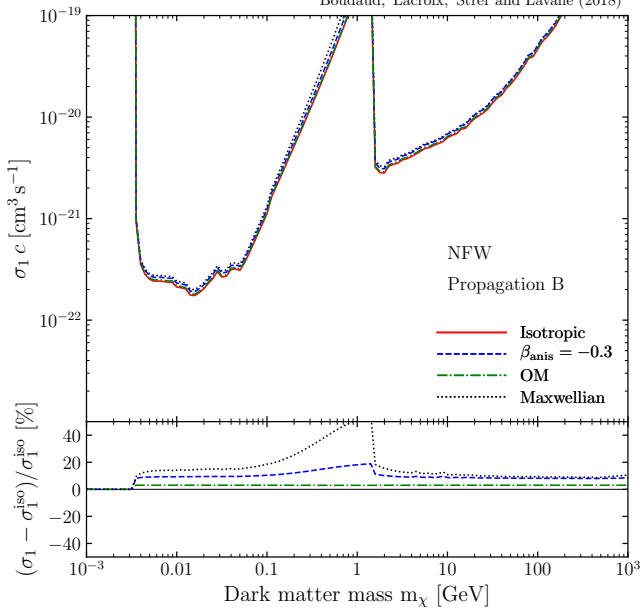
- Testing the method against cosmological simulations
- Not direct fits!!!
- Preliminary results: Eddington method globally performs much better than SHM [Lacroix+ 2018b, in prep.](#)

Thank you for your attention!





Boudaud+ 2018



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