



Single particle versus collectivity, shapes of exotic nuclei

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Rewriting Nuclear Physics textbooks
30 years with Radioactive Ion Beam Physics
Pisa (Italy), July 20th – 24th, 2015



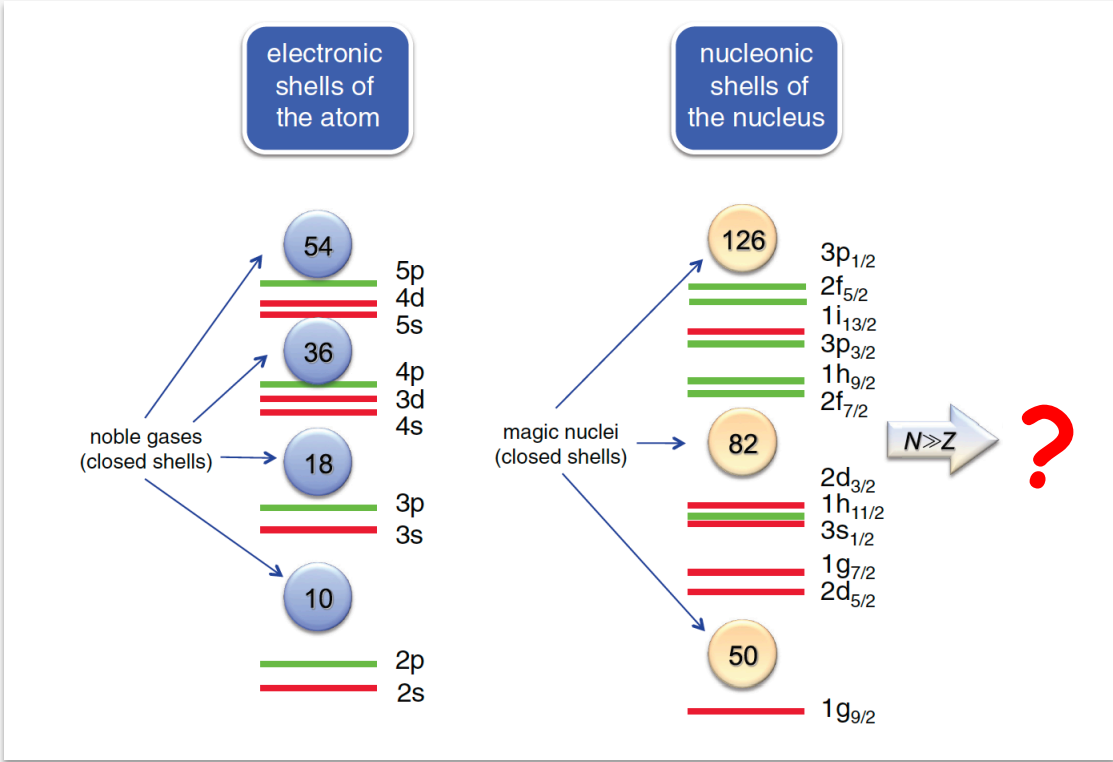
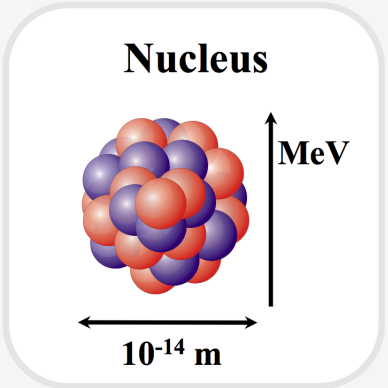
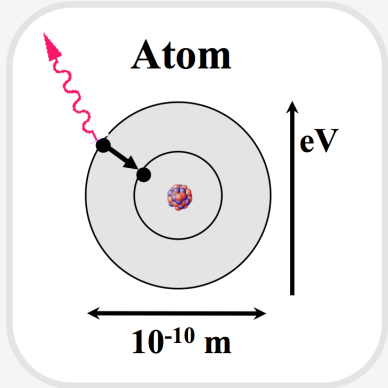
Single particle versus collectivity, shapes of exotic nuclei

- What a title ! A bit of everything ...
- Start with some reminders of “classical” nuclear physics
- What’s new over the last 30 years ?
From the perspective of an observer of the field since 1989 ...
(pure experimentalist, γ -ray spectroscopist interested in $A > 70$ nuclei)

**Rewriting Nuclear Physics textbooks
30 years with Radioactive Ion Beam Physics**

- Radioactive ion beams are an important part, but not the whole story !
- Small selection of examples – only limited time ...
- I will simplify (hopefully without getting things wrong), omit, not always show latest etc.
- Try to avoid topics which others may already have shown

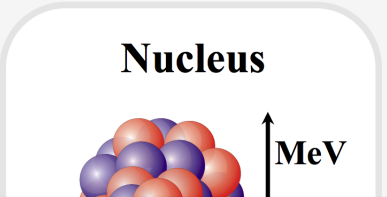
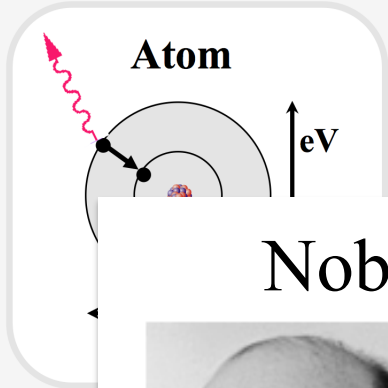
The shell structure of atomic nuclei




Remember:
Two independent parameters
in atomic nuclei, Z and N !

“Shell evolution”:
Next part of my talk ...


The shell structure of atomic nuclei




Nobel Prize in Physics 1963



Eugene Paul Wigner
Prize share: 1/2



Maria Goeppert Mayer
Prize share: 1/4



J. Hans D. Jensen
Prize share: 1/4

noble gases
(closed shells)



ent parameters
ei, Z and N !

**"Shell evolution":
Next part of my talk ...**

The birth of the collective models in 1952

Interpretation of Isomeric Transitions of Electric Quadrupole

AAGE BOHR
Institute for Theoretical Physics, Copenhagen, Denmark

(Received November 19, 1952)

IN the recent classification of isomeric transitions on the basis of the shell model, examples of lifetimes appear more than a factor of a hundred larger than expected on the basis of the shell model. we have here the effect of collective motion.¹

A natural interpretation of the isomeric transitions in terms of collective motion and nuclear surface oscillations in the shell model, the low-lying states of the particle structure of the nucleus, or by an excitation of the particle quantum number states are of the former character. However, that transition is

Nov. 1952

Rotational States in Even-Even Nuclei

AAGE BOHR AND BEN R. MOTTELSON*
Institute for Theoretical Physics, Copenhagen, Denmark

(Received March 24, 1953)

IN a recent note,¹ an interpretation of the short-lived $E2$ isomers has been suggested in terms of rotational states of the deformed nucleus. Empirical evidence is rapidly accumulating on the low energy spectra of even-even nuclei;²⁻⁴ the purpose of the present note is to call attention to the extensive support which exists in these data for the above interpretation, and to suggest its usefulness in the analysis of decay schemes.

In the model describing the nucleus in terms of the coupled particle motion and surface oscillations, low-lying rotational states are associated with the large deformations expected in regions with many particles outside of closed shells. In such regions, the rotational spectrum is expected to be given rather accurately by the simple expression¹

$$E_I = \frac{\hbar^2}{2\mathcal{I}} I(I+1), \quad I=0, 2, 4, 6, \dots \quad (1)$$

even parity

March 1953

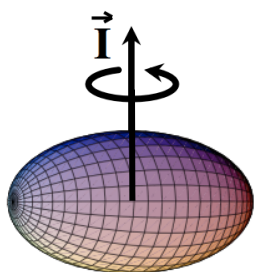
First observation of rotational bands

Three examples from 1953:

Fine structure in α decay

Nobel Prize in Physics 1975

rotations



even-even
nucleus



Aage Niels Bohr

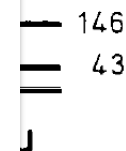
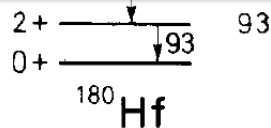


Ben Roy Mottelson

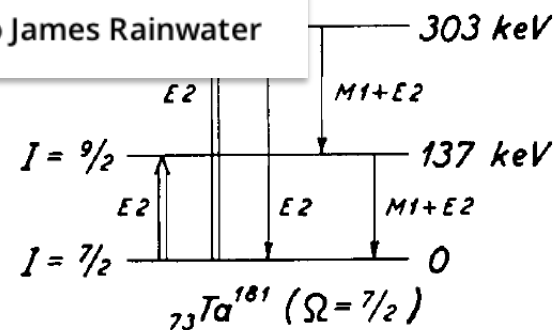


Leo James Rainwater

$$E(4^+)/E(2^+) \sim 3.3$$



experiment



The problem of the "wrong" moment of inertia

Det Kongelige Danske Videnskabernes Selskab

Matematisk-fysiske Meddelelser, bind 30, nr. 1

Dan. Mat. Fys. Medd. 30, no. 1 (1955)

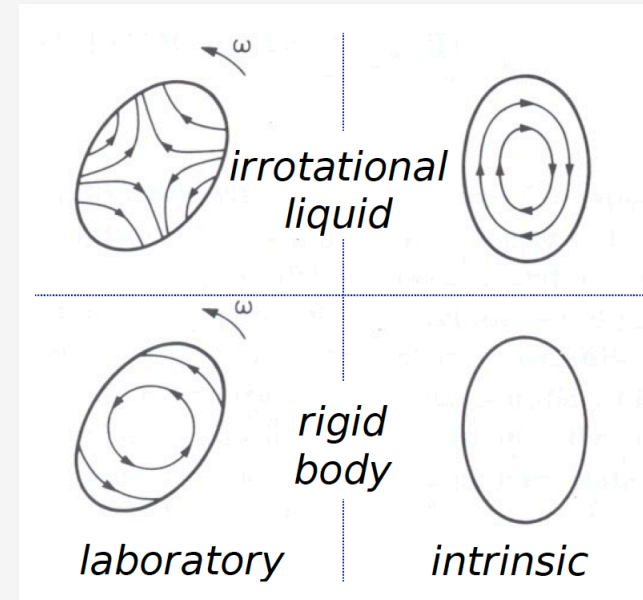
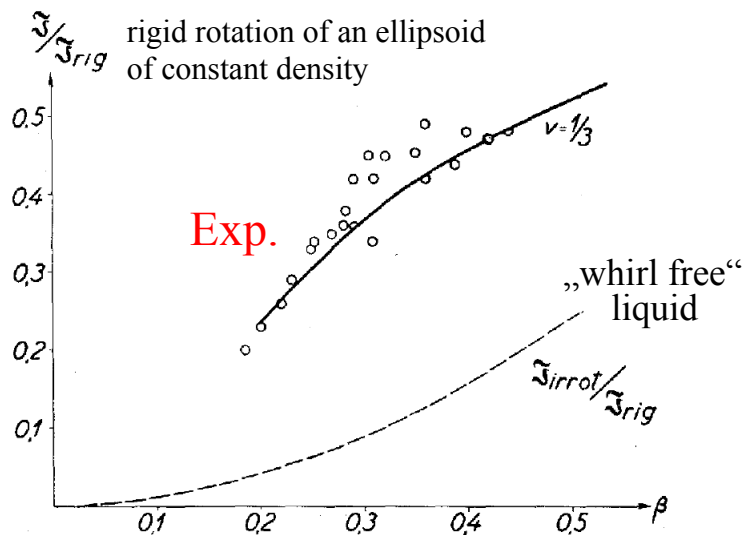
DEDICATED TO PROFESSOR NIELS BOHR ON THE
OCCASION OF HIS 70TH BIRTHDAY

1955

MOMENTS OF INERTIA OF ROTATING NUCLEI

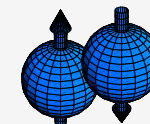
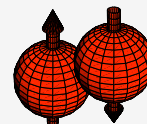
BY

AAGE BOHR AND BEN MOTTELSON

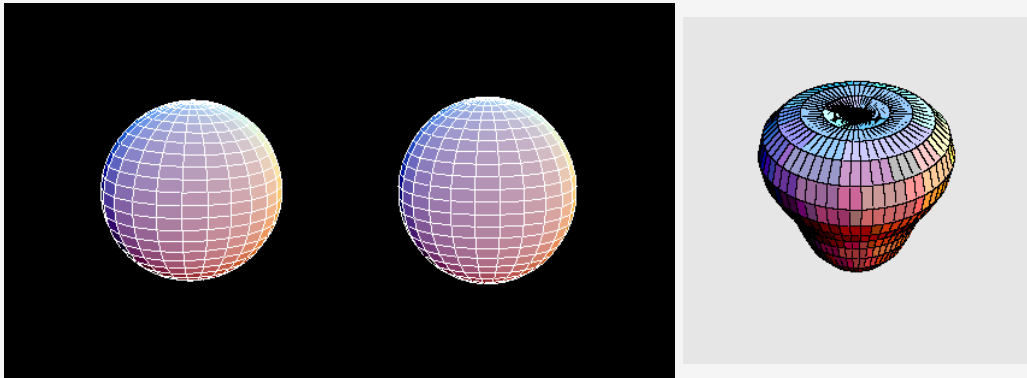
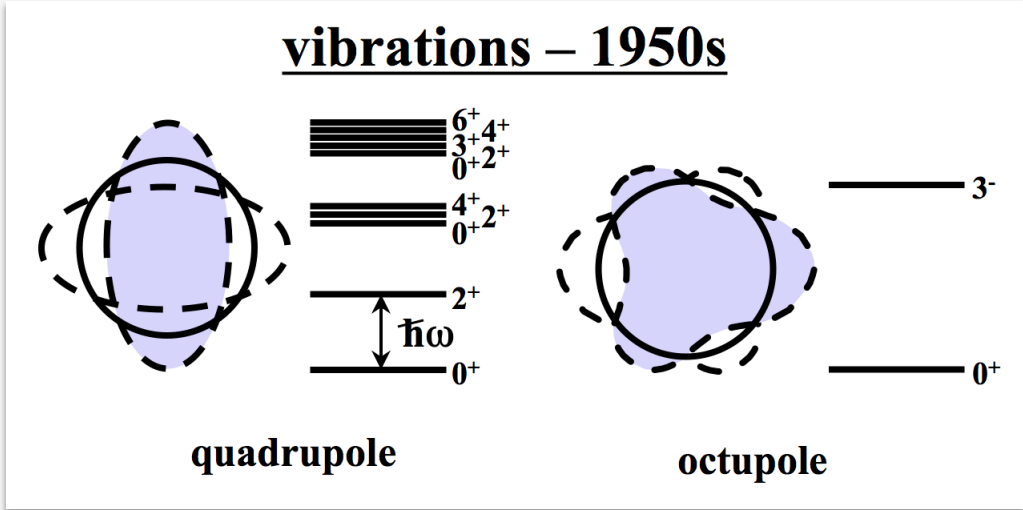
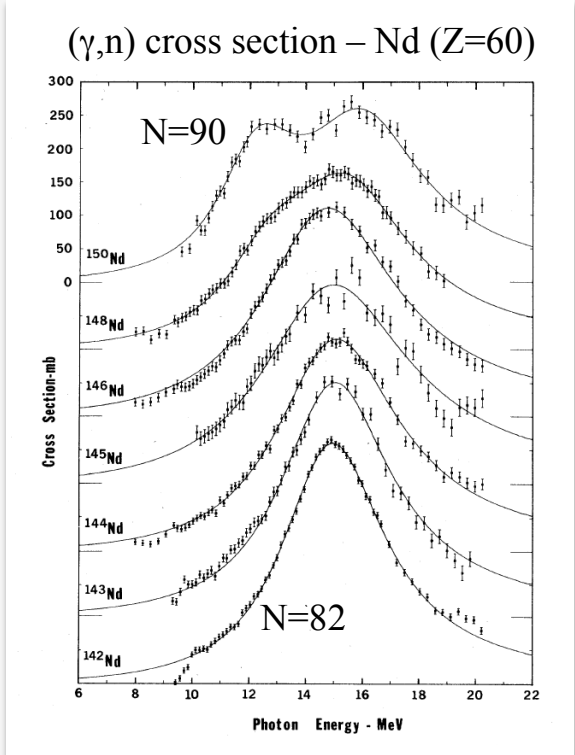


„Nuclei are like egg shells which are filled
with a mixture of a normal and a super-
conducting liquid !“

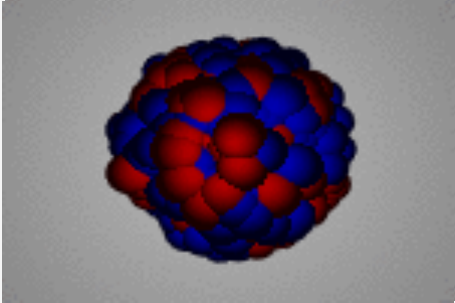
Super conductivity due to **pairing forces**
in analogy to the Cooper pairs (electrons)
in superconductors.



The vibrational degree of freedom



Giant Dipole Resonance



Deformed nuclei in the Nilsson model

Det Kongelige Danske Videnskabernes Selskab

Matematisk-fysiske Meddelelser, bind 29, nr 16

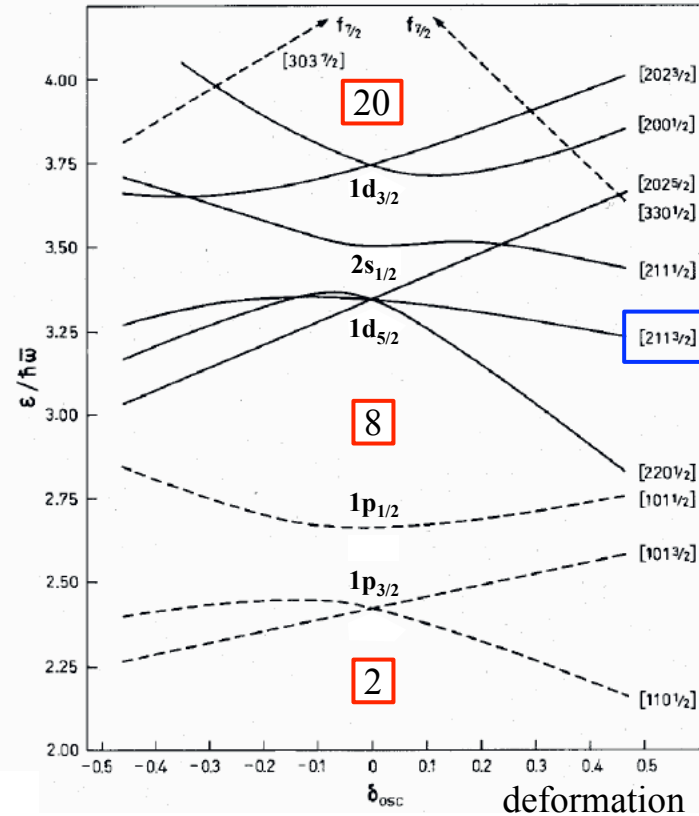
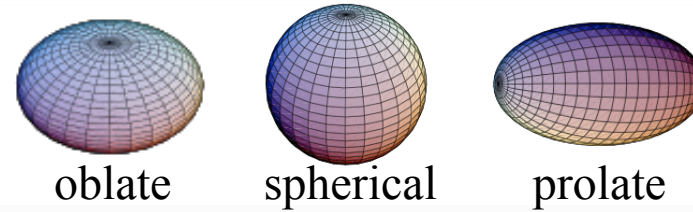
Dan Mat Fys Medd 29, no 16 (1955)

1955

BINDING STATES OF INDIVIDUAL NUCLEONS IN STRONGLY DEFORMED NUCLEI

BY

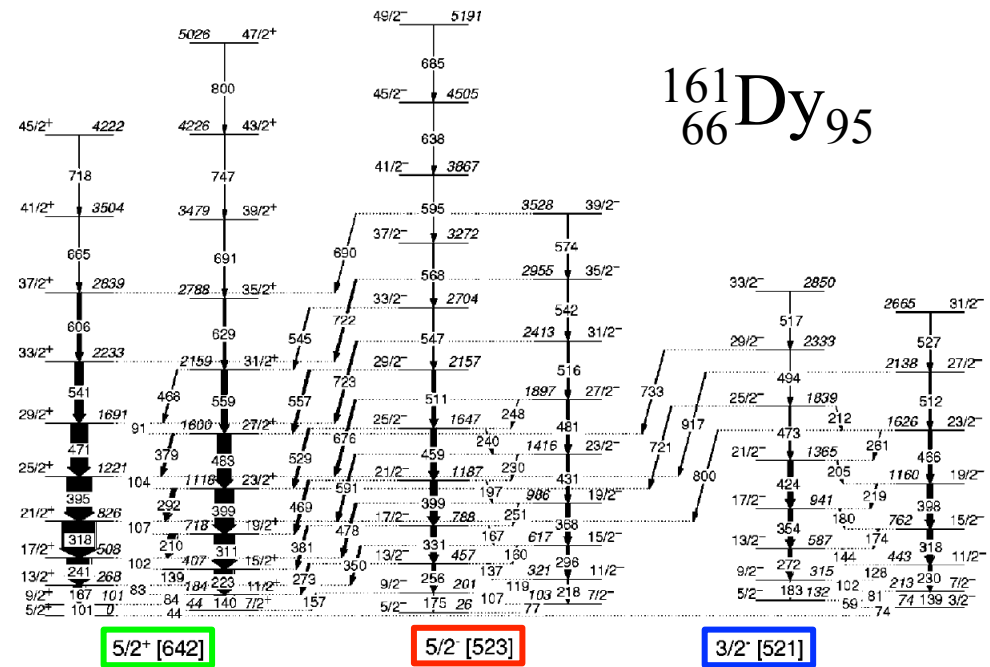
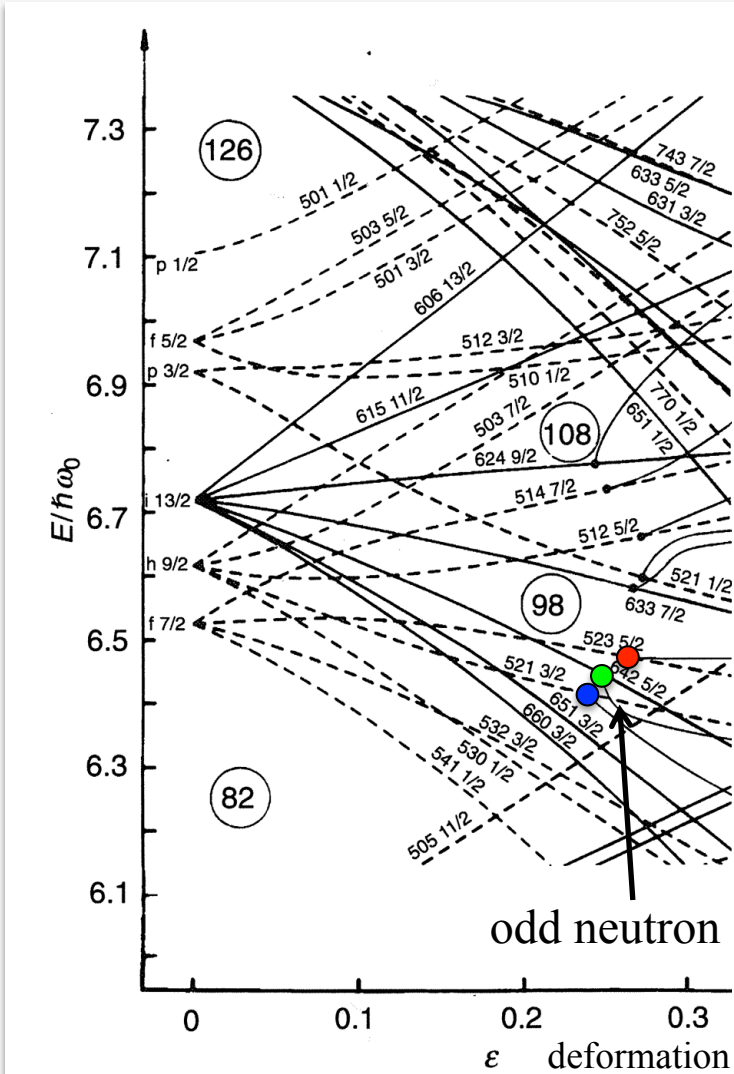
SVEN GÖSTA NILSSON



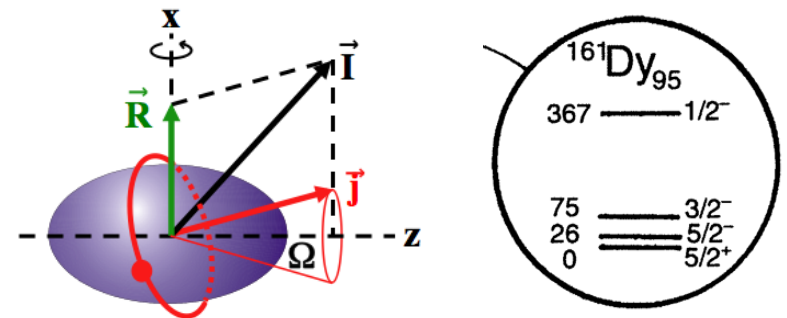
new asymptotic
quantum numbers

$2j+1$ -fold degeneracy is removed,
states with $+m$ and $-m$ still degenerate !

Single-particle motion in a deformed field

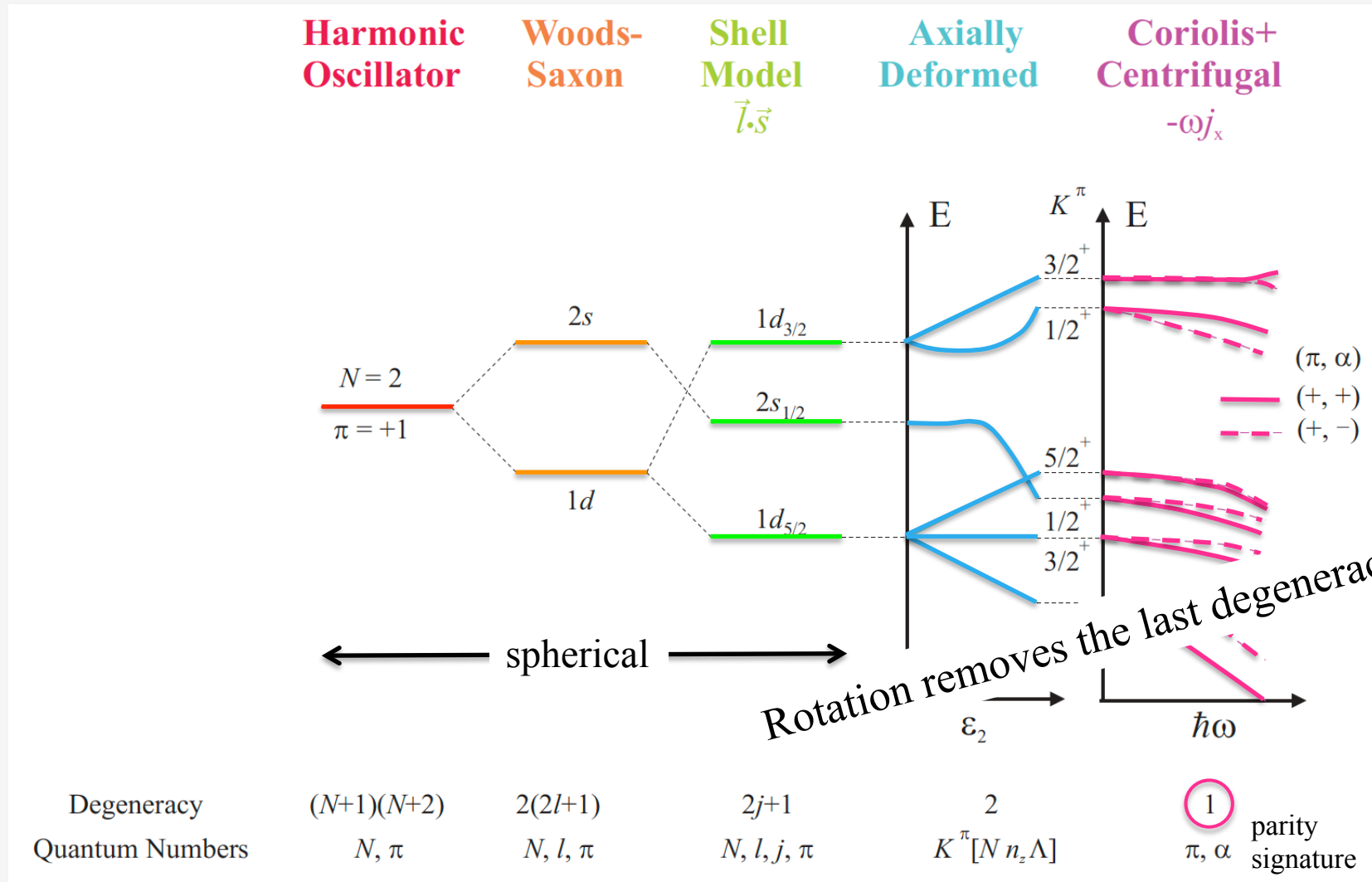


A. Jungclaus et al., Phys. Rev. C67, 034302 (2003)



Three rotational band heads within 75 keV !

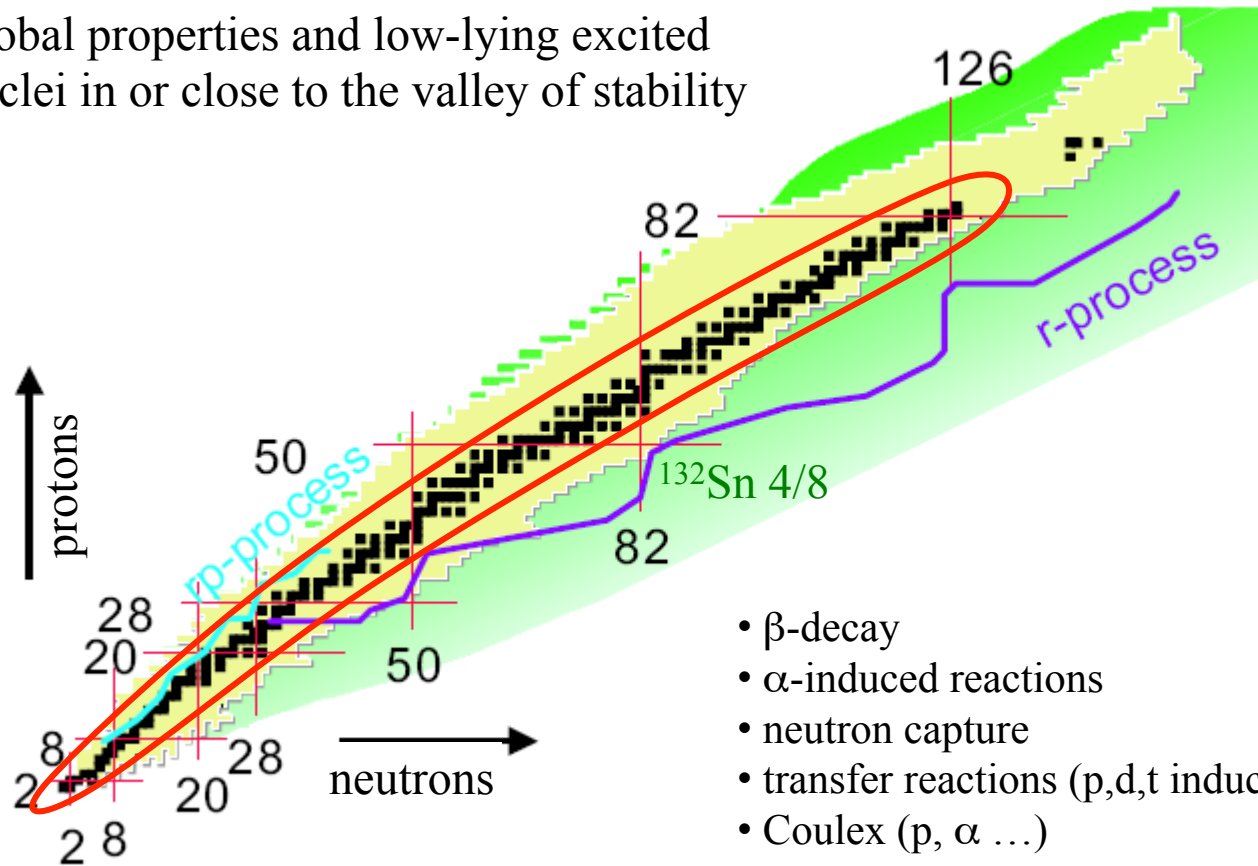
Single-particle motion in a rotating deformed field



Single-particle energies depend on deformation and rotational frequency !

Over the next twenty years ...

Study of global properties and low-lying excited states of nuclei in or close to the valley of stability

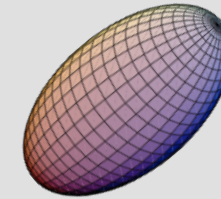
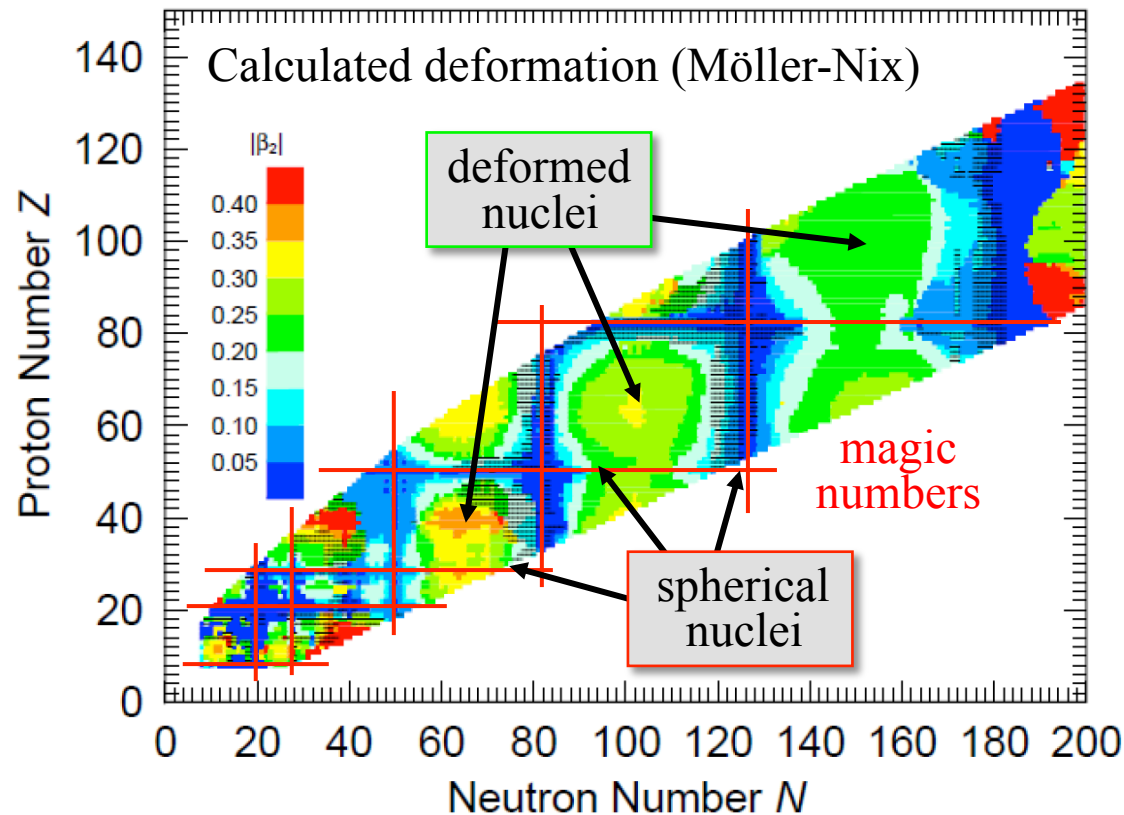


- β -decay
- α -induced reactions
- neutron capture
- transfer reactions (p,d,t induced)
- Coulex (p, α ...)
- γ -induced neutron emission

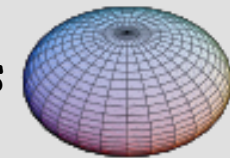
• scattering experiments (e, p ...)

Radii and density distributions !

„Classical“ nuclear physics

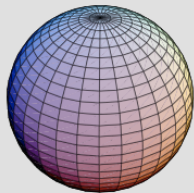


prolate nucleus



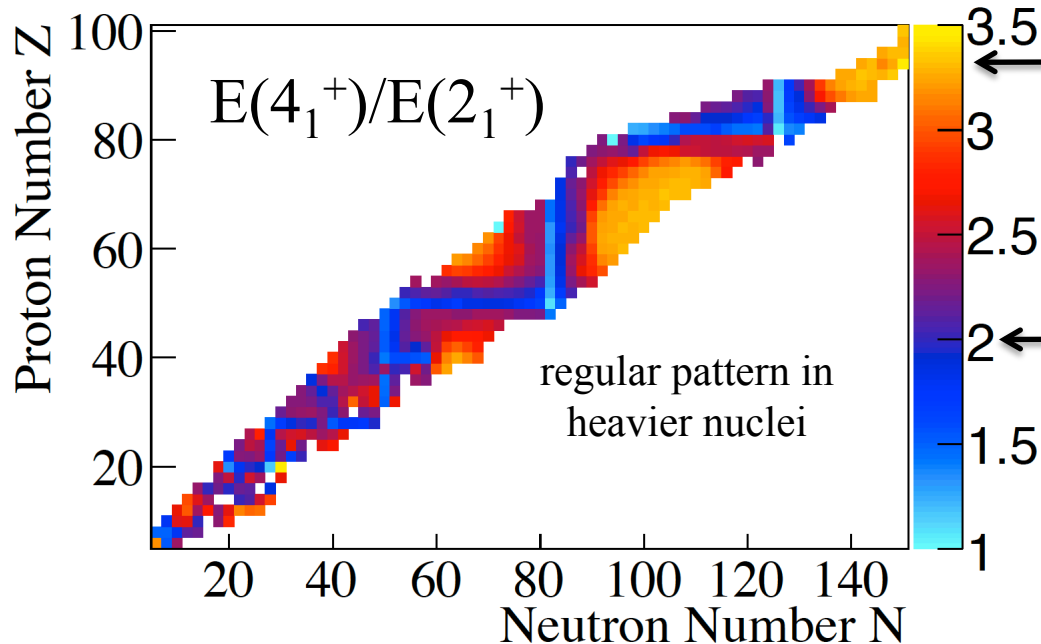
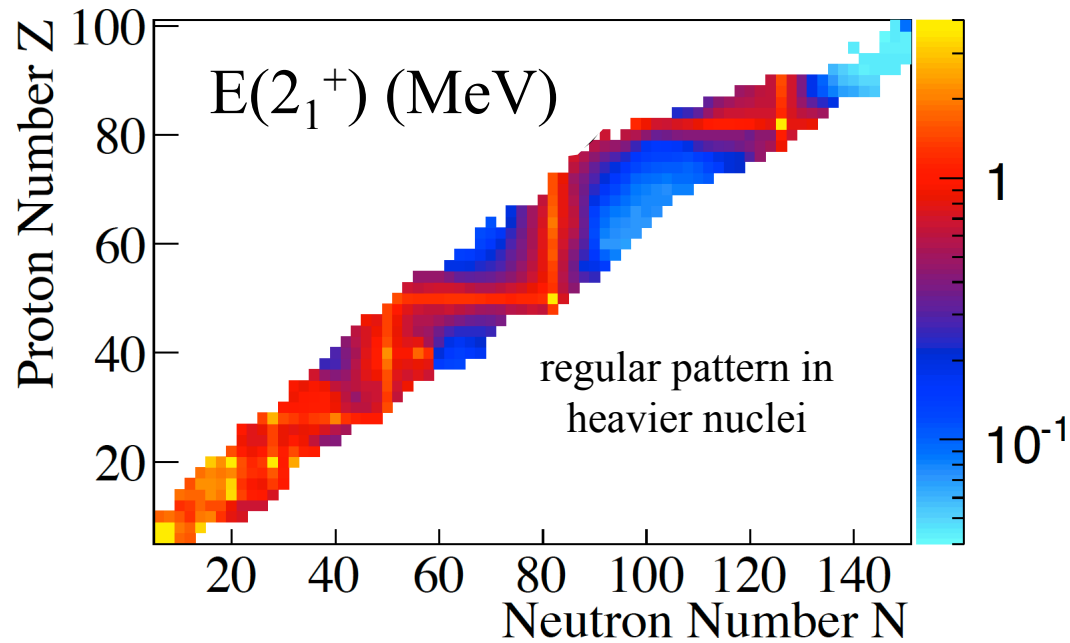
oblate nucleus

rotations and vibrations
→ collective models



spherical nucleus

Single-particle excitations
→ nuclear shell model
vibrations → collective models



$E(2^+)$ and $E(4^+)/E(2^+)$
as global indicators
in even-even nuclei

Rotations

— 8^+

— 6^+

— 4^+

— 2^+

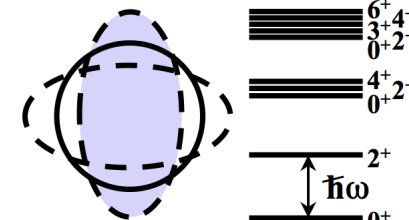
— 0^+

$$E_x = I(I+1) \frac{\hbar^2}{2J}$$

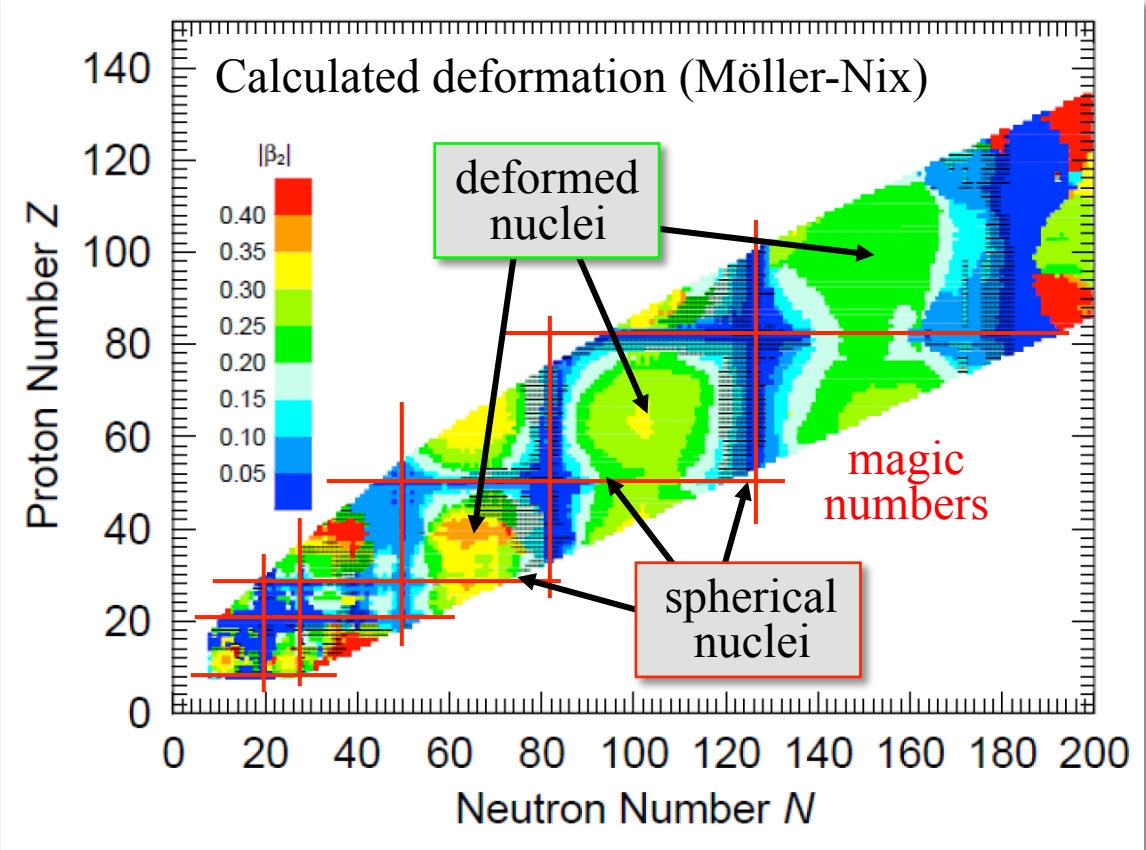
$$\begin{aligned} E(4_1^+)/E(2_1^+) &= (4 \cdot 5)/(2 \cdot 3) \\ &= 20/6 = 3.3 \end{aligned}$$

$$\begin{aligned} E(4_1^+)/E(2_1^+) &= 2\hbar\omega/\hbar\omega = 2 \end{aligned}$$

Quadrupole vibrations



„Classical“ nuclear physics



prolate nucleus

oblate nucleus

rotations and vibrations
→ collective models

spherical nucleus

Single-particle excitations
→ nuclear shell model
vibrations → collective models

Is that already the full story ?

No !!!!!!!!!!!!!!!!

Let's play with all degrees of freedom

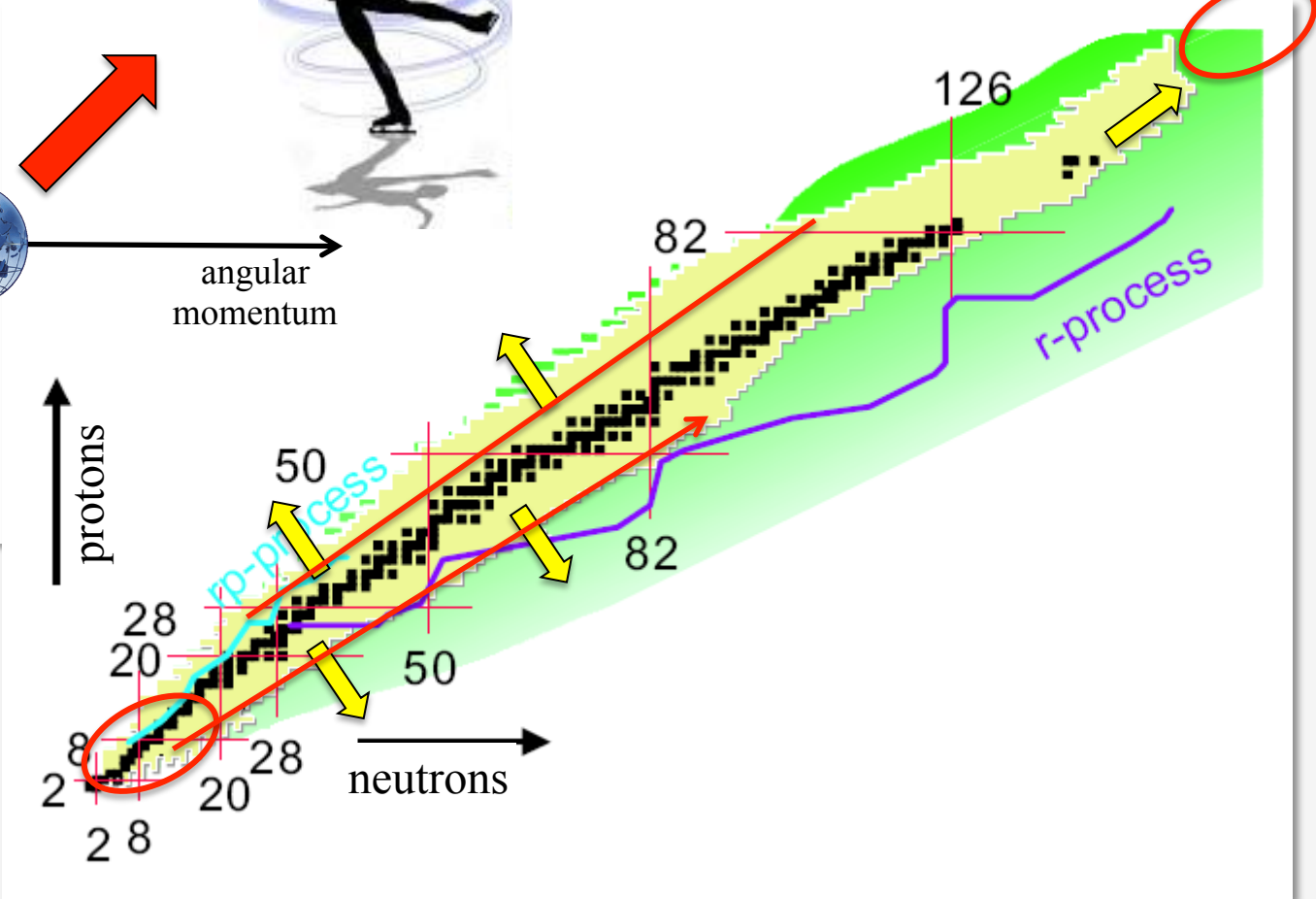


temperature
excitation energy

angular
momentum

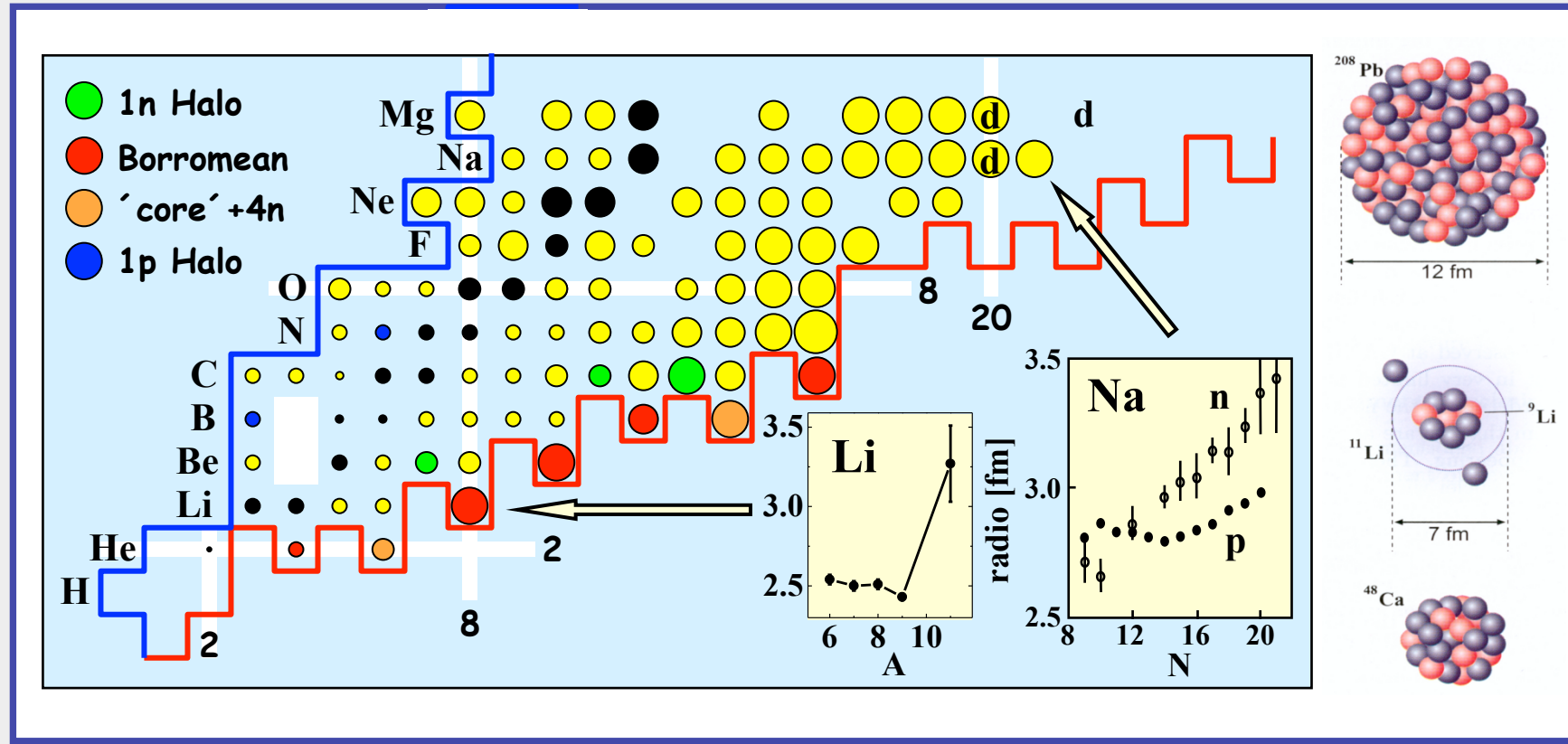
Leave valley
of stability!

isospin, N/Z



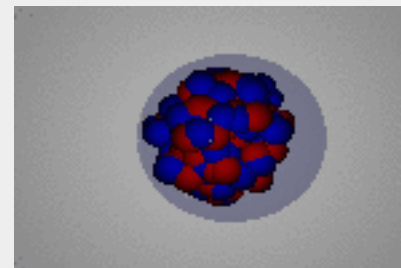
Let's start with the
light nuclei!

Nuclear radii - halos and neutron skins



Remember:

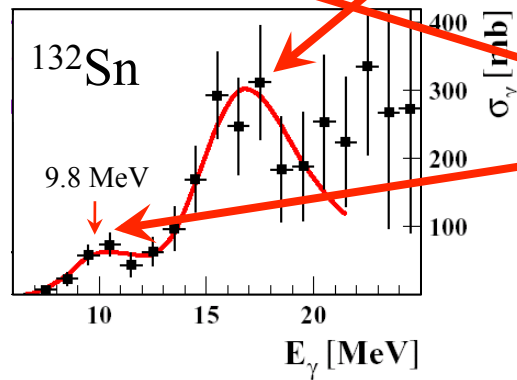
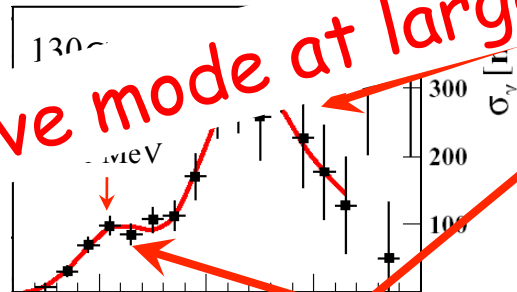
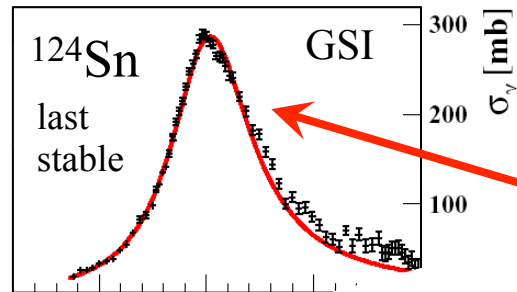
- $R = r_0 \cdot A^{1/3}$ $r_0 = 1.1-1.2$ fm
- The thickness t of the nuclear surface is constant.
- Protons and neutrons are uniformly mixed.



The Pygmy resonance in neutron-rich Sn isotopes

neutron number
 74
 //
 increases
 82
 N ↓

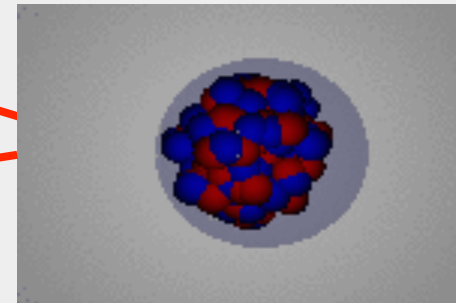
photo-neutron cross section



Giant Dipole Resonance



Pygmy Resonance

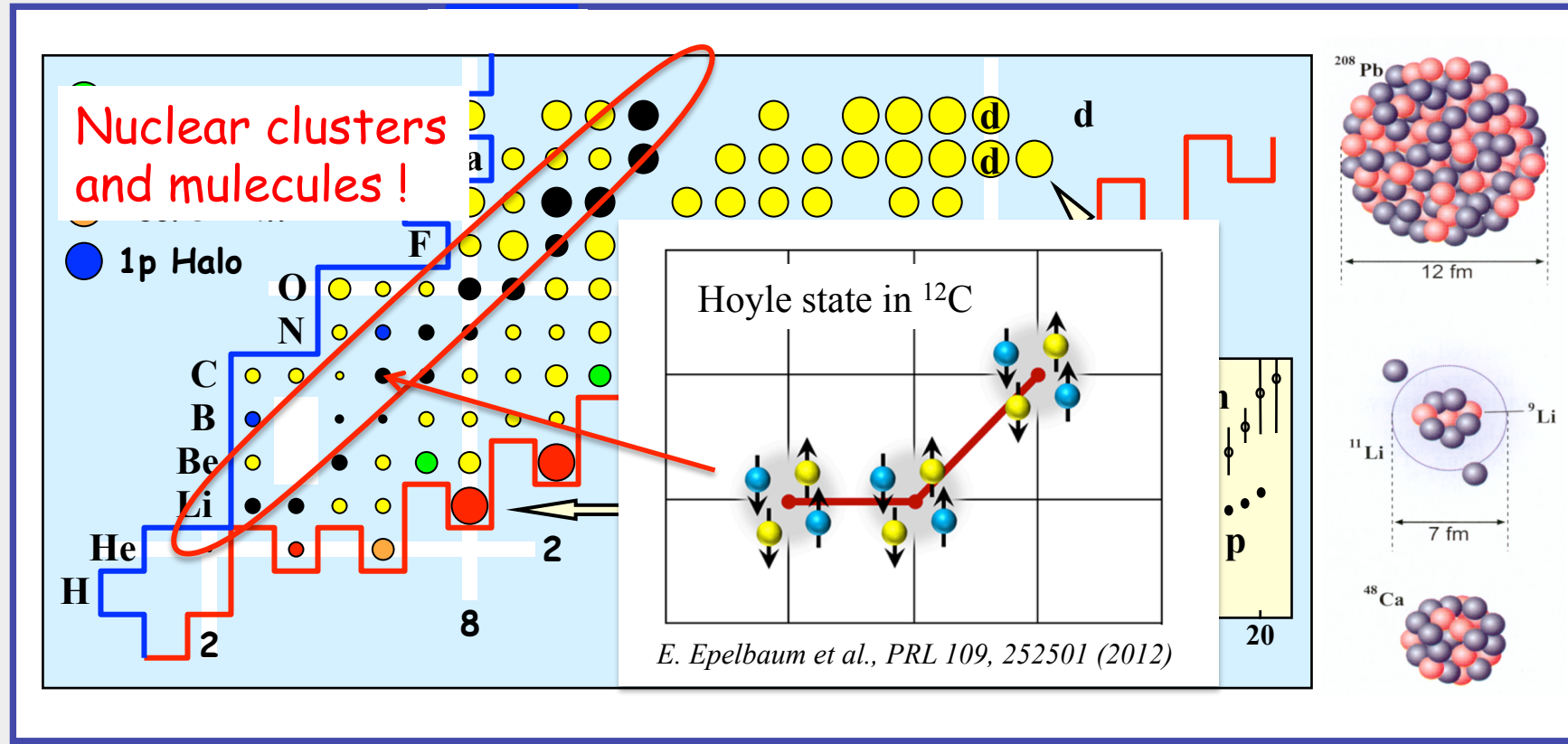


A new collective mode at large isospin!

P. Adrich et al.

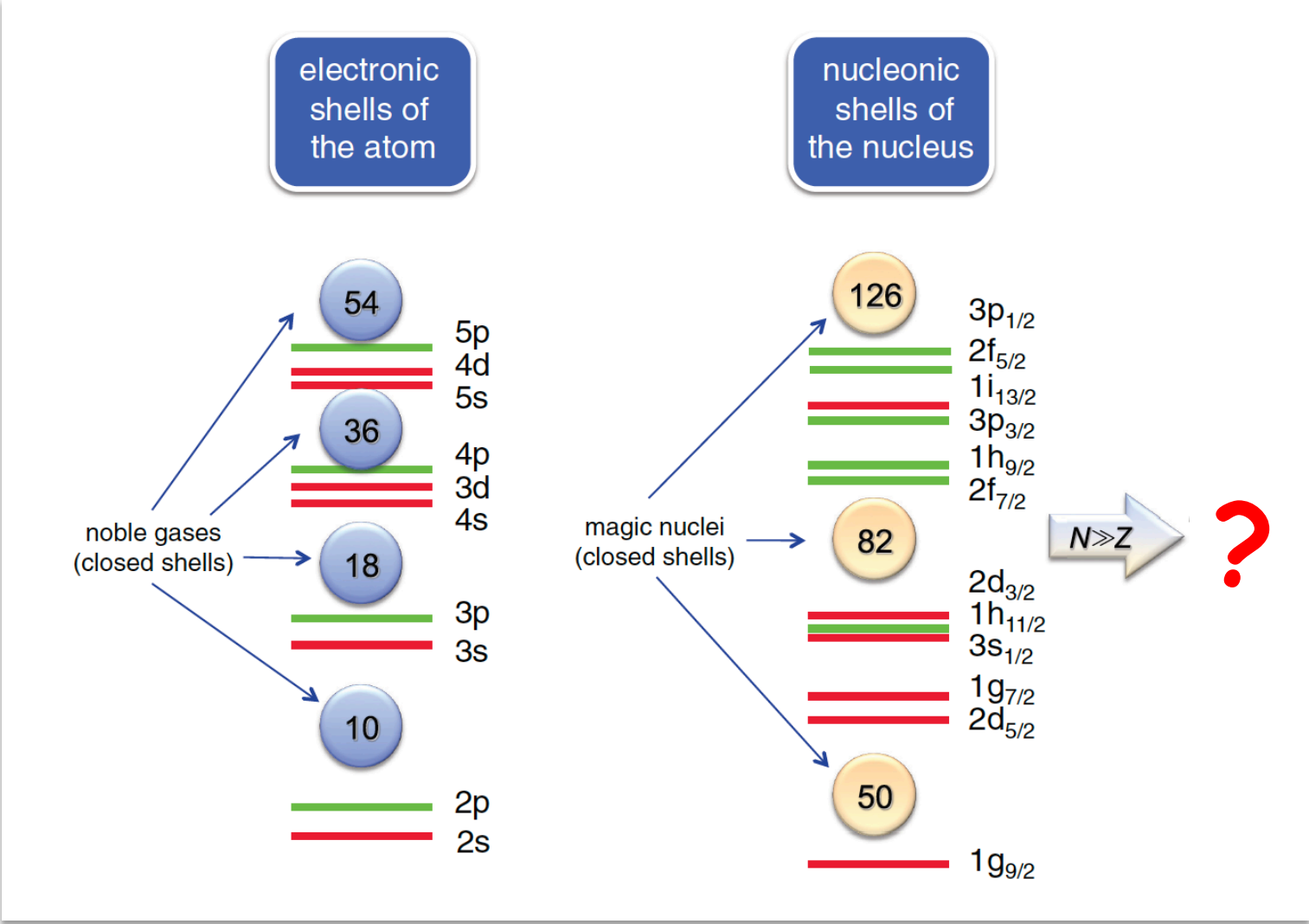
Phys. Rev. Lett. 95 (2005) 132501

Nuclear radii - halos and neutron skins



Light nuclei are indeed a rich playground!

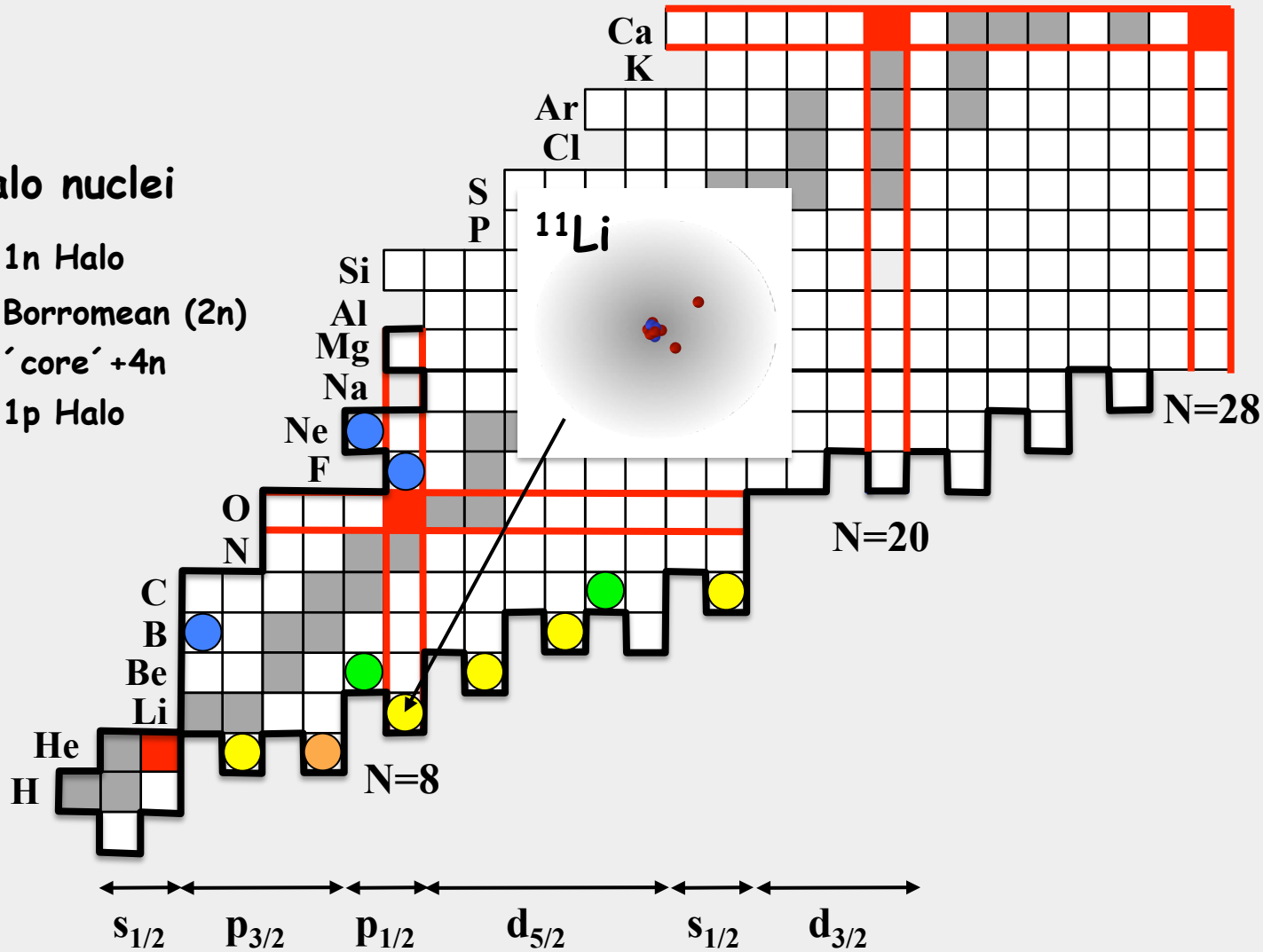
"Shell evolution" on the neutron-rich side



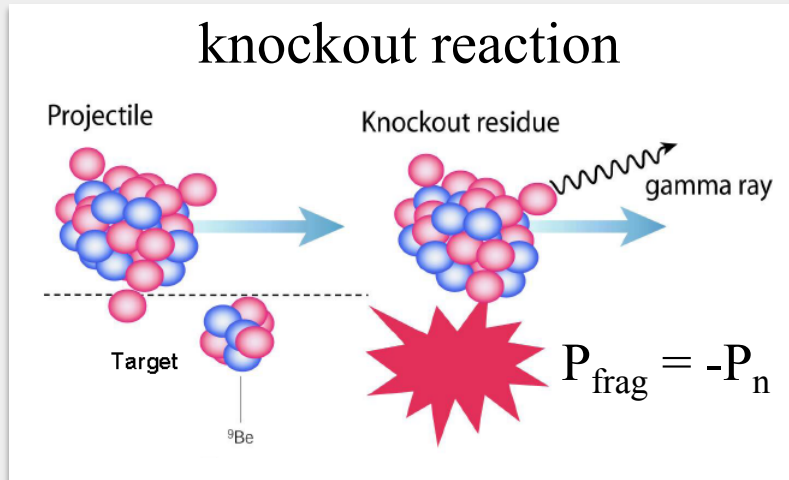
Neutron halos and the N=8 shell closure

Halo nuclei

- 1n Halo
- Borromean (2n)
- 'core' +4n
- 1p Halo

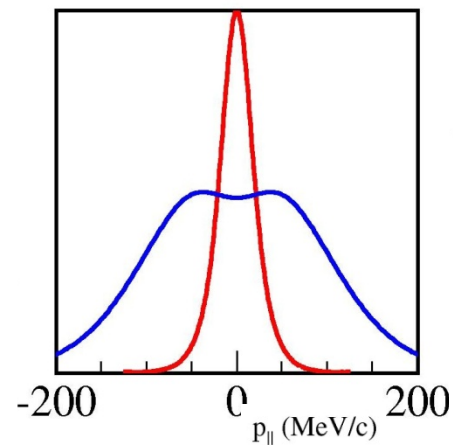


Knockout reactions at relativistic energies



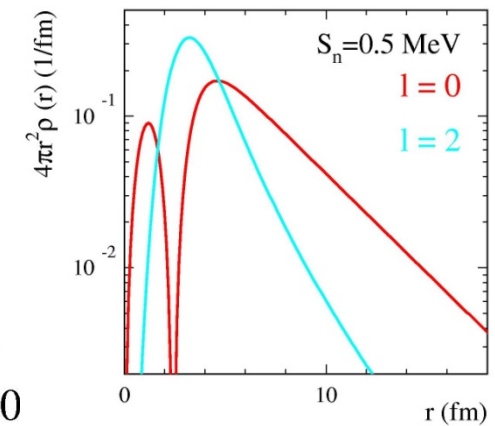
Heisenberg's uncertainty principle

momentum distribution



narrow momentum distribution

density distribution



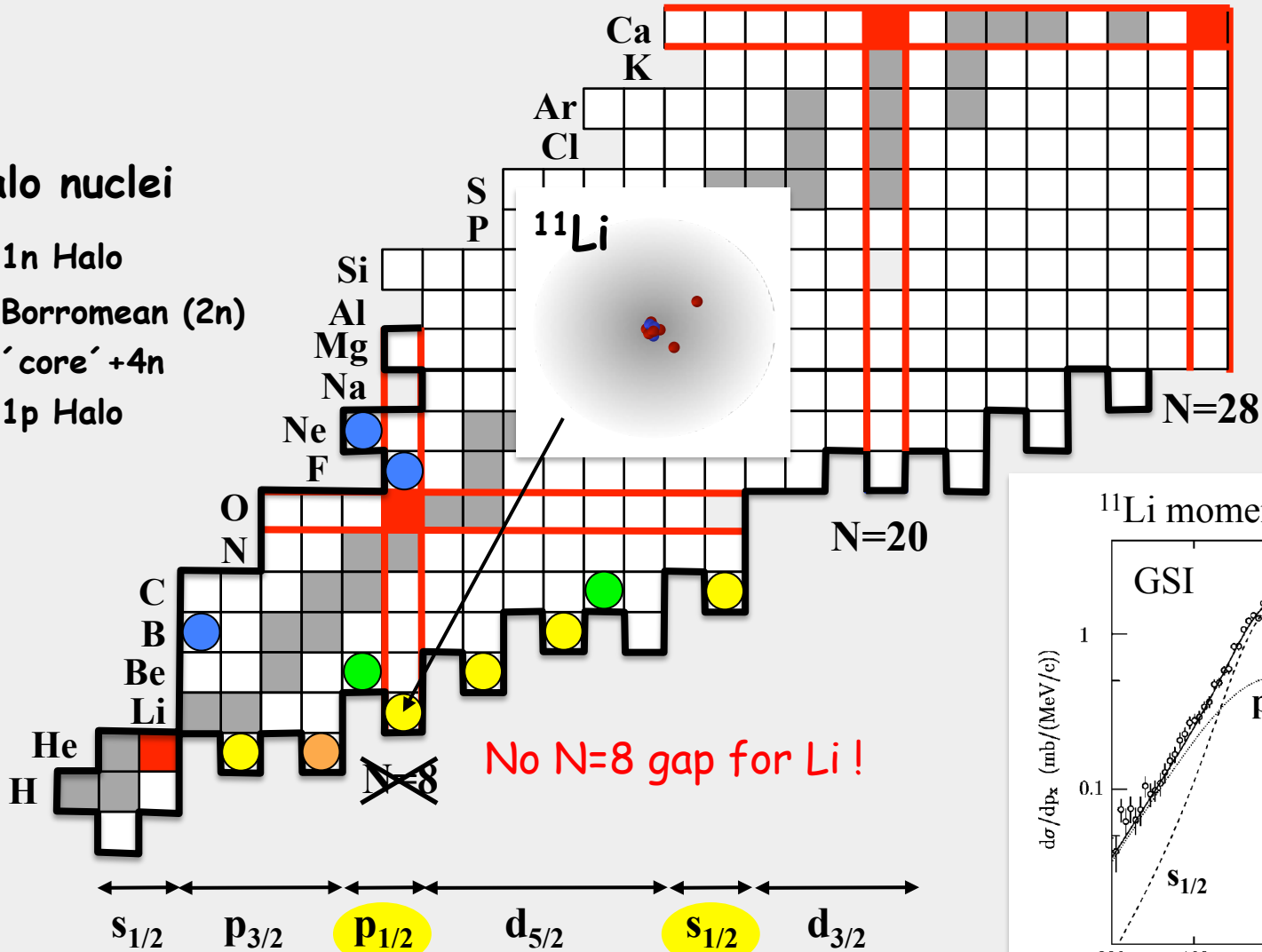
extended wavefunction

→ angular momentum ℓ of the removed nucleon

Neutron halos and the N=8 shell closure

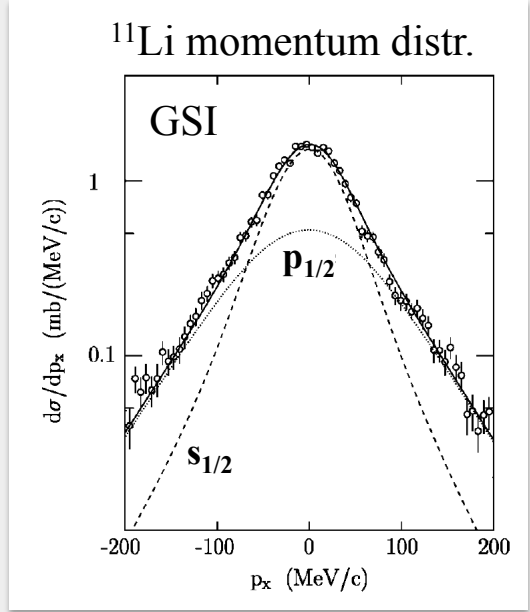
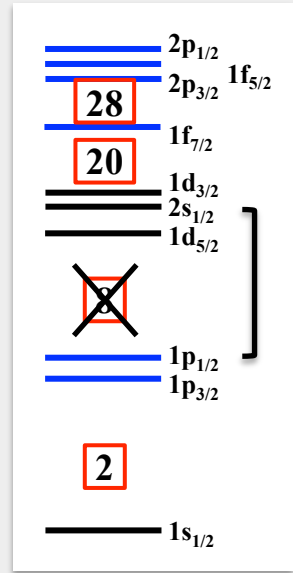
Halo nuclei

- 1n Halo
- Borromean (2n)
- 'core' +4n
- 1p Halo



No N=8 gap for Li!

strong s-wave admixture

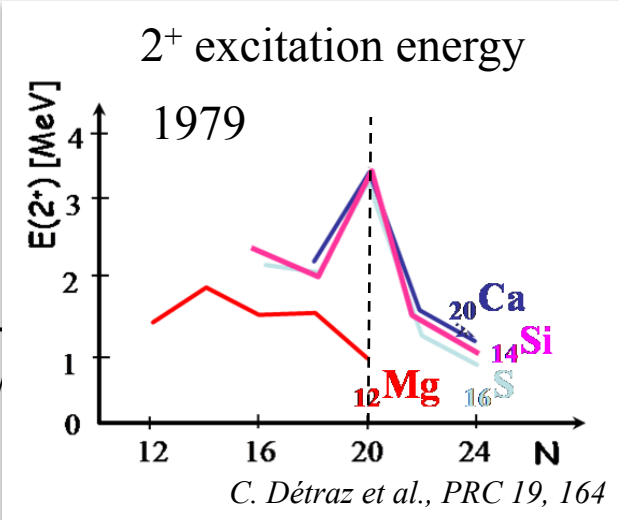
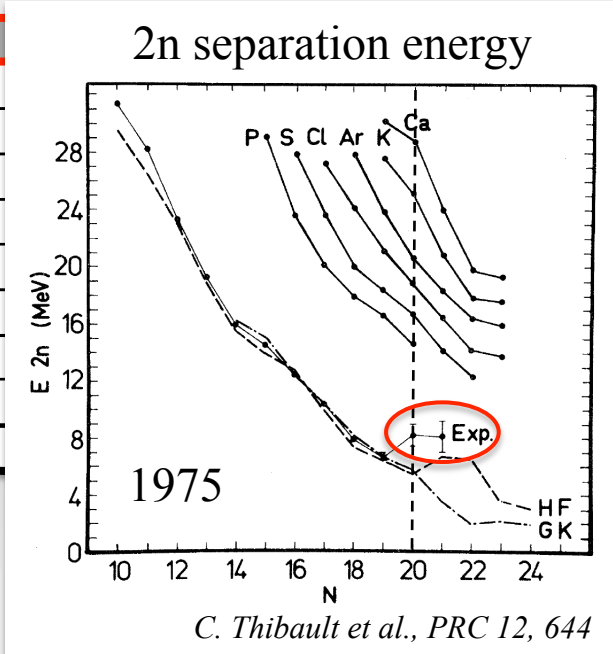
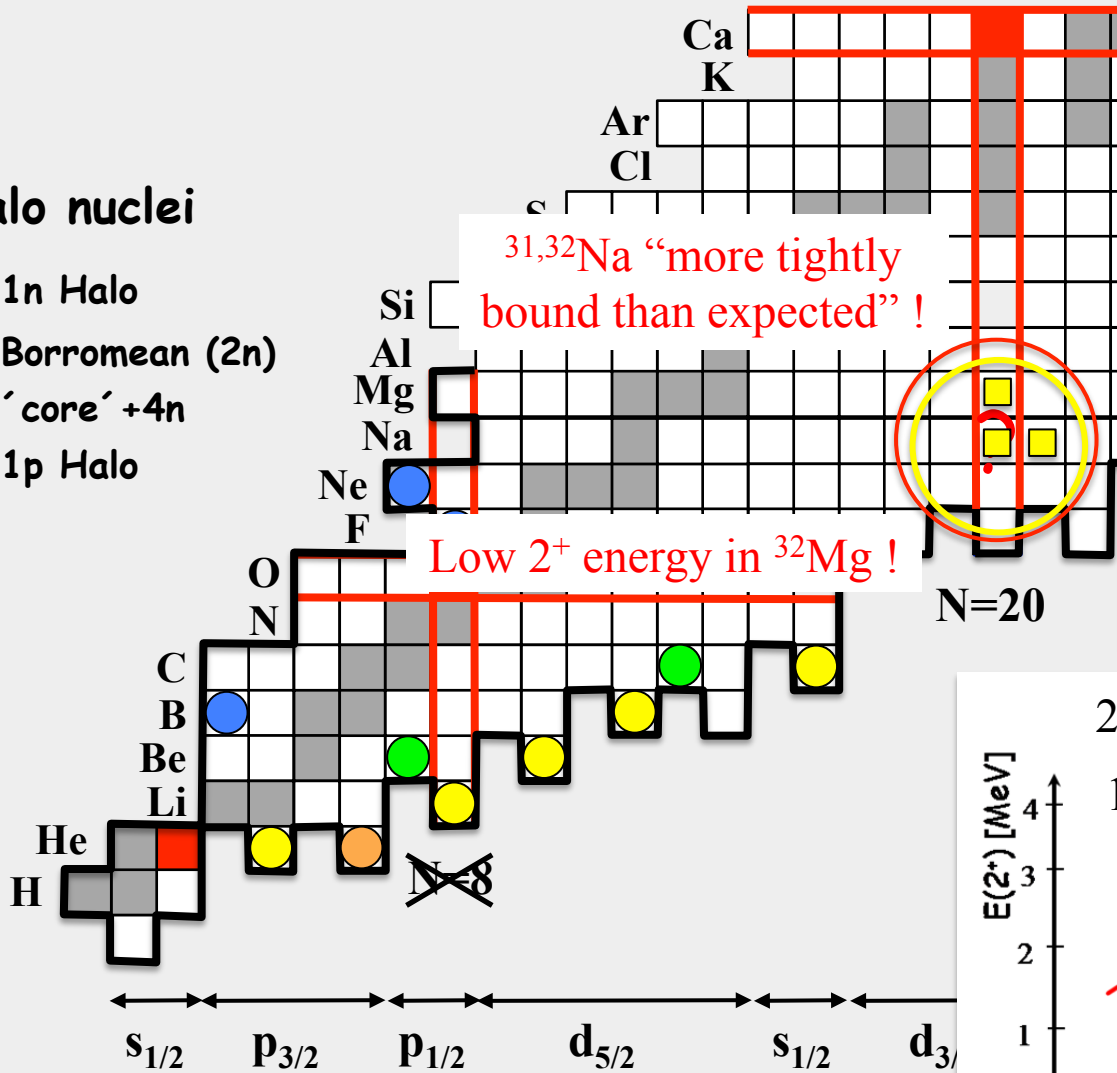


H. Simon et al., Phys. Rev. Lett. 83, 496 (1999)

What about the N=20 shell closure ?

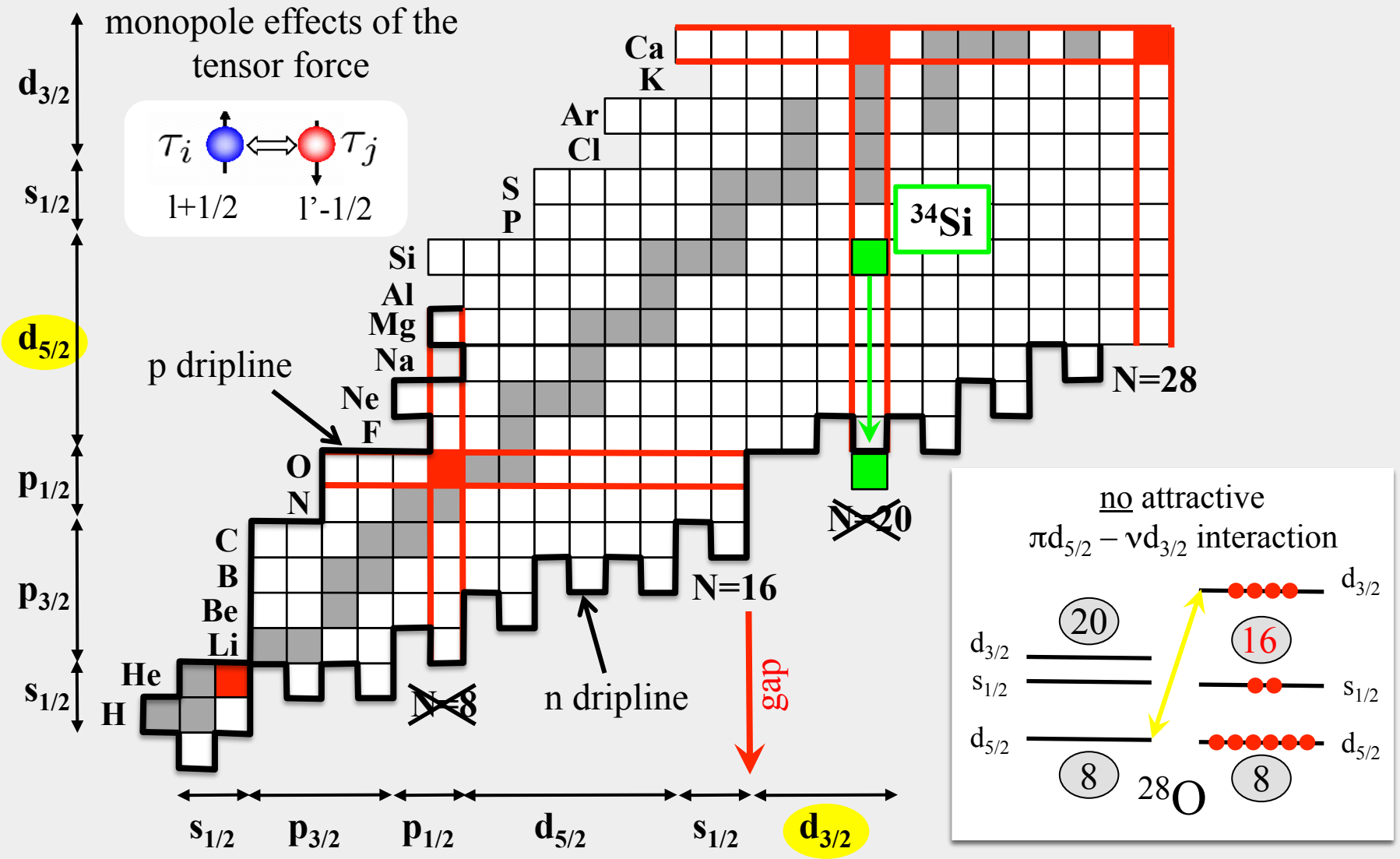
Halo nuclei

- 1n Halo
- Borromean (2n)
- 'core' +4n
- 1p Halo



Since then a wealth of information in this region

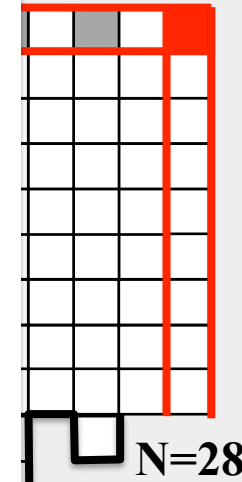
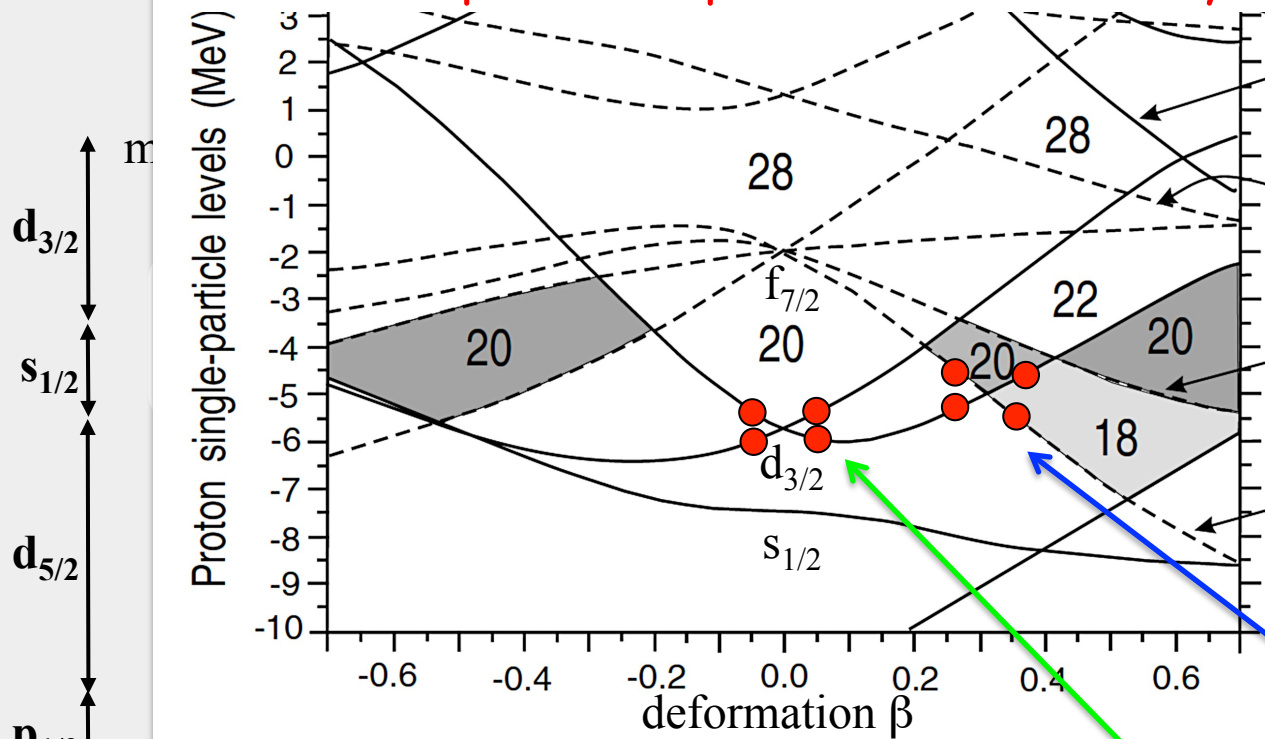
Explanation within the shell model picture



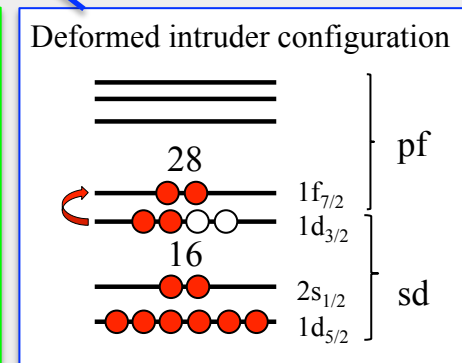
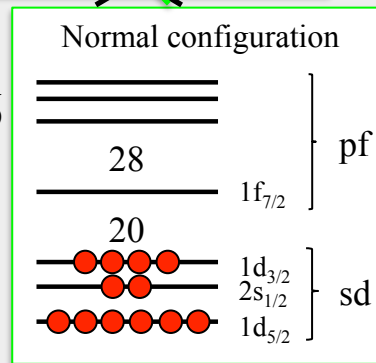
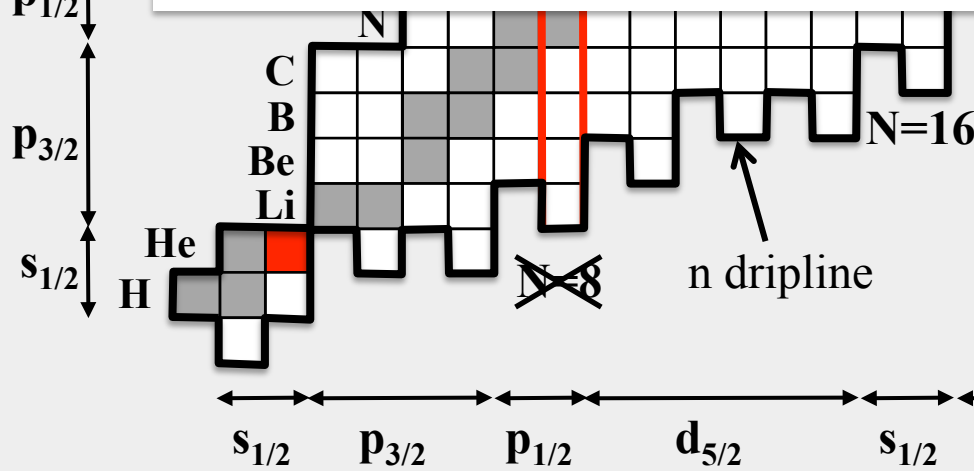
T. Otsuka et al., Phys. Rev. Lett. 87, 082502 (2001); Phys. Rev. Lett. 95, 232502 (2005)

First example of shape coexistence today !

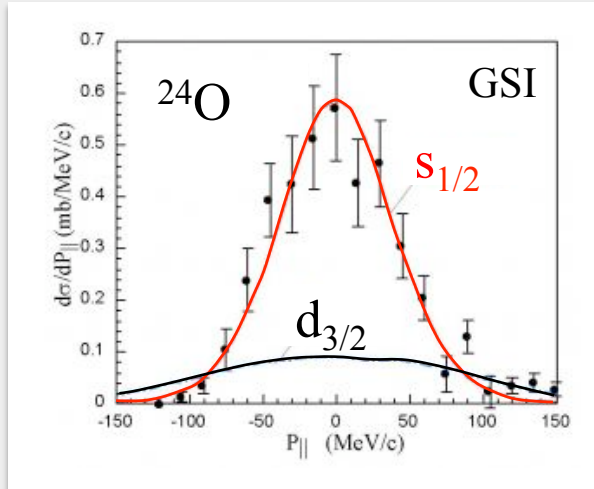
picture



Deformed intruder configuration g.s. in ^{32}Mg !

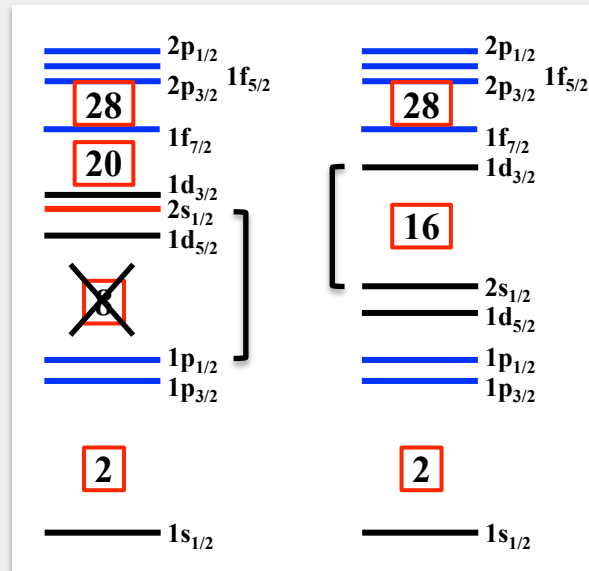


Momentum distributions in $^{23,24}\text{O}$



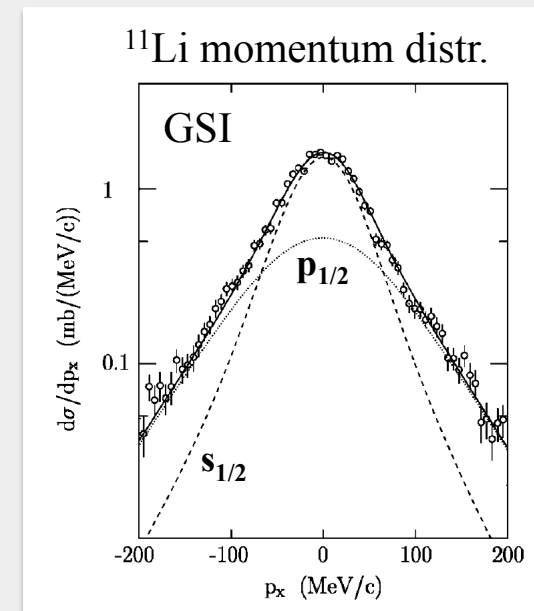
R. Kanungo et al., PRL 102, 152501 (2009)

Last neutron occupies $s_{1/2}$ orbital, no $d_{3/2}$ component !

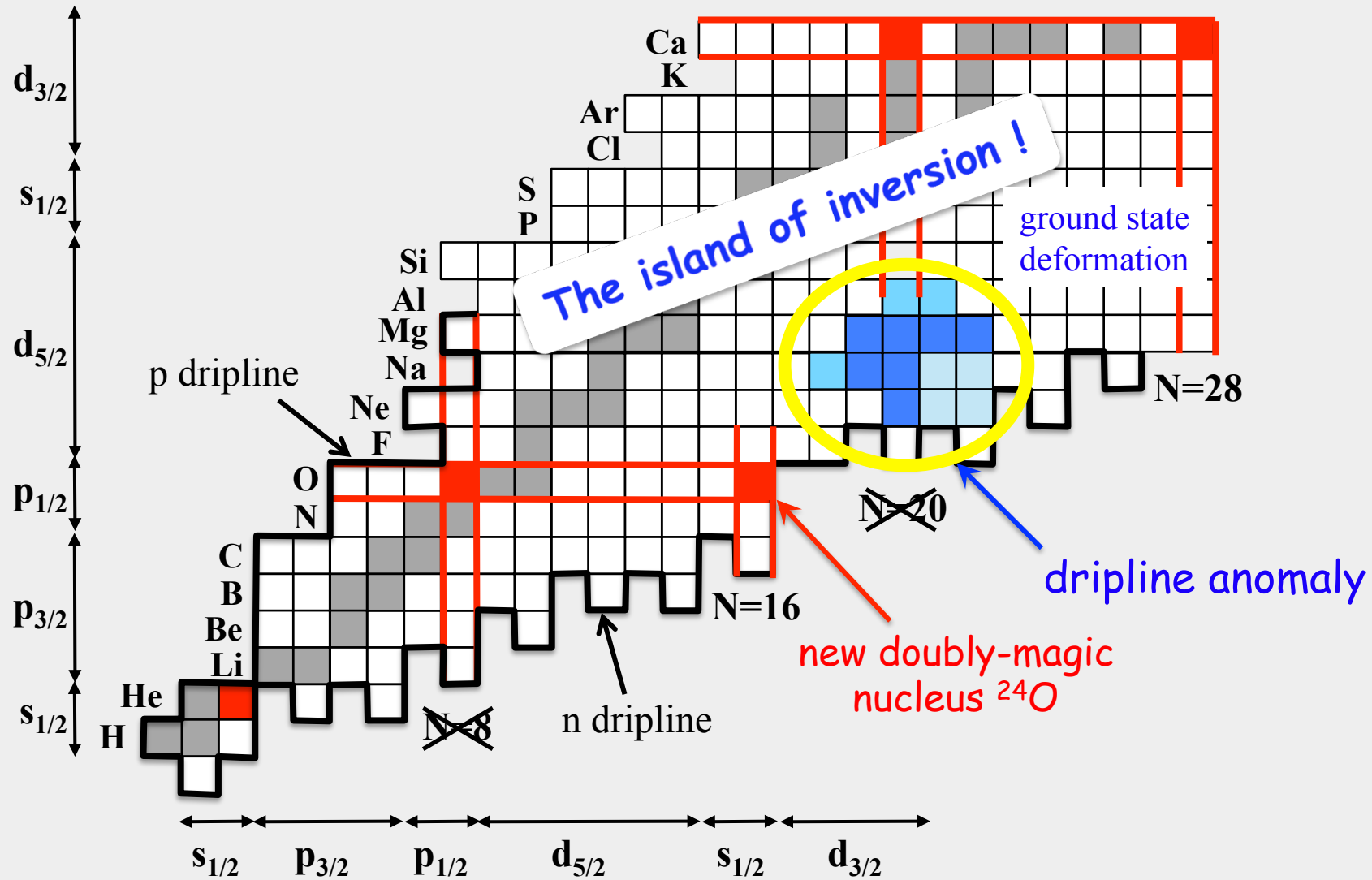


$s_{1/2}$ and $d_{3/2}$ close in energy, mixing to be expected
No mixing \rightarrow N=16 gap !

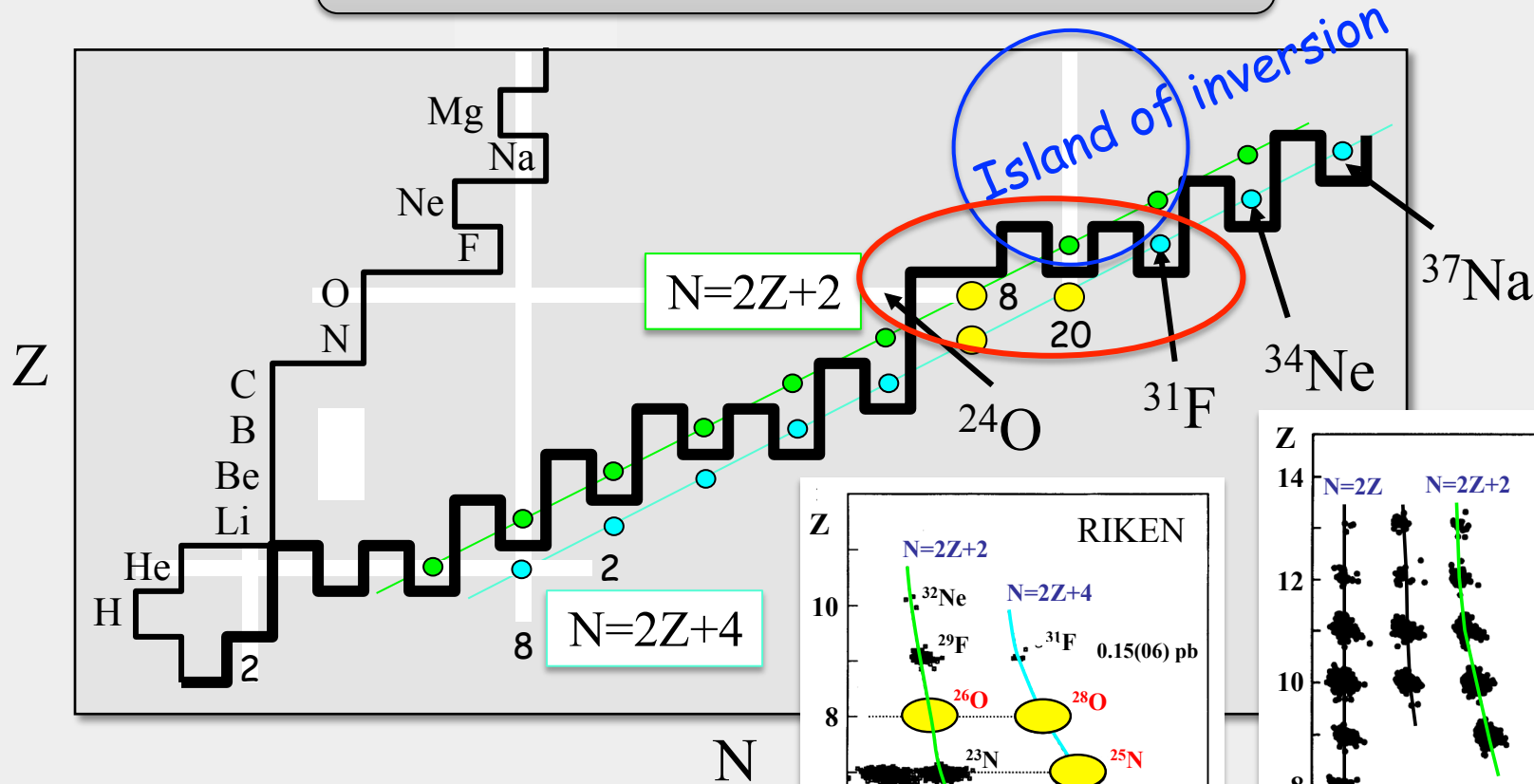
Spherical magic number at N=16,
 ^{24}O doubly magic !



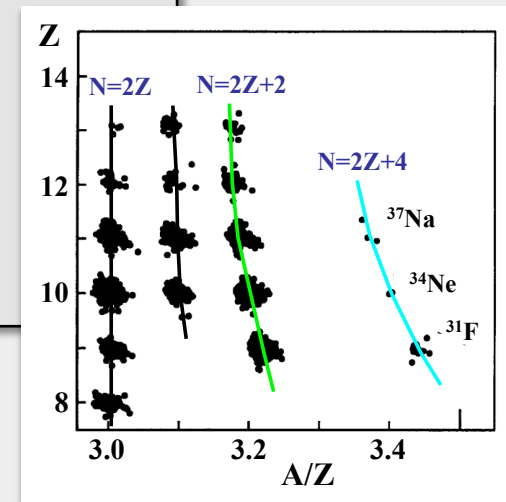
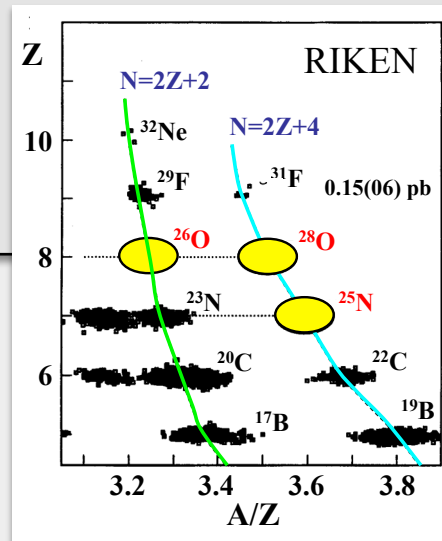
New magic number and the island of inversion



The neutron dripline up to Na



In F ($Z=9$) **six more neutrons** are bound as compared to O ($Z=8$) !

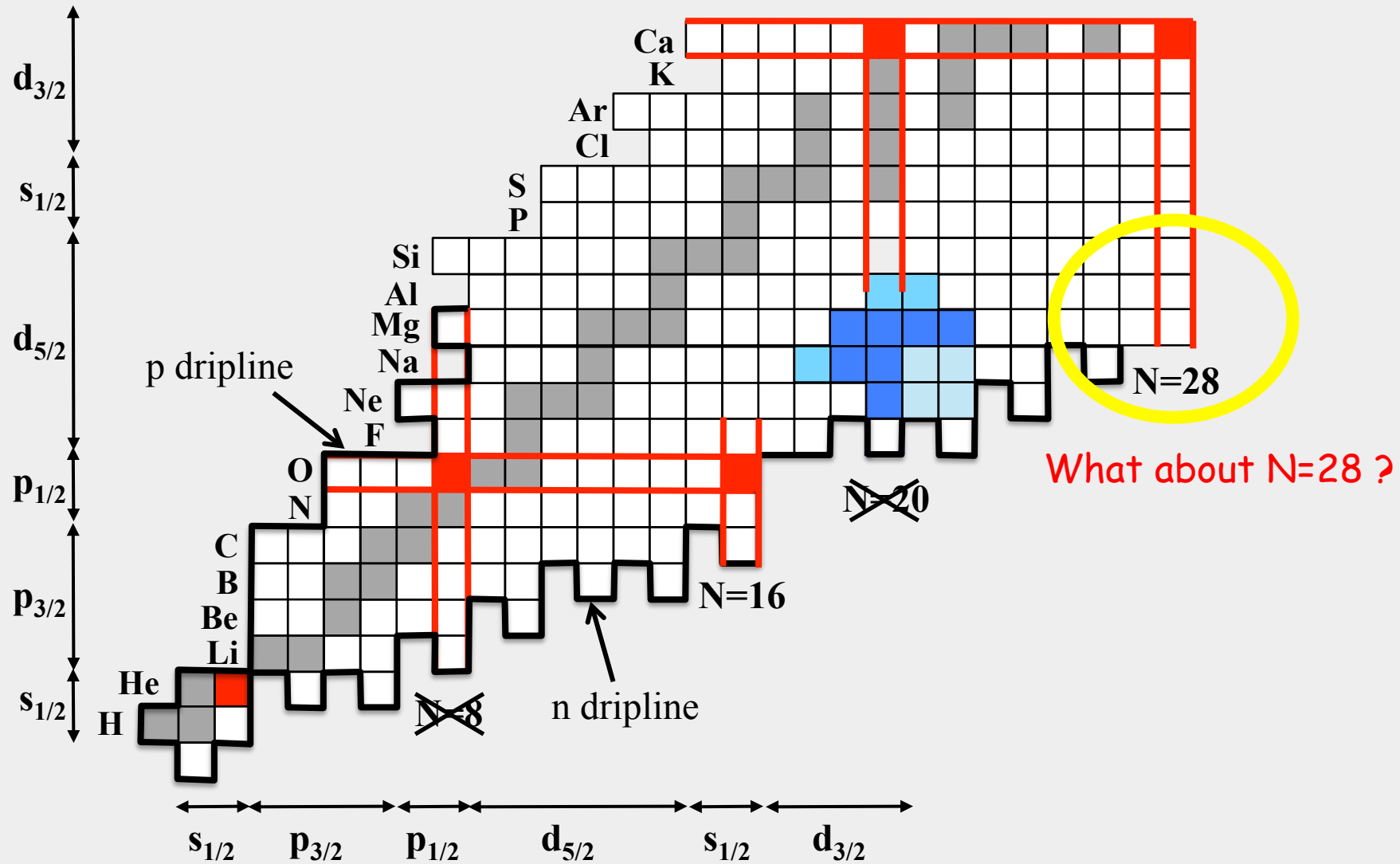


H. Sakurai et al., PLB 448 (1999) 180
M. Notani et al., PLB 542 (2002) 49

Additional binding due to deformation !!!

The mere existence of a nucleus already tells us something ...

New magic number and the island of inversion



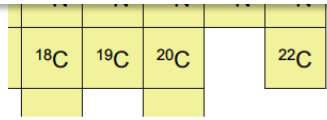
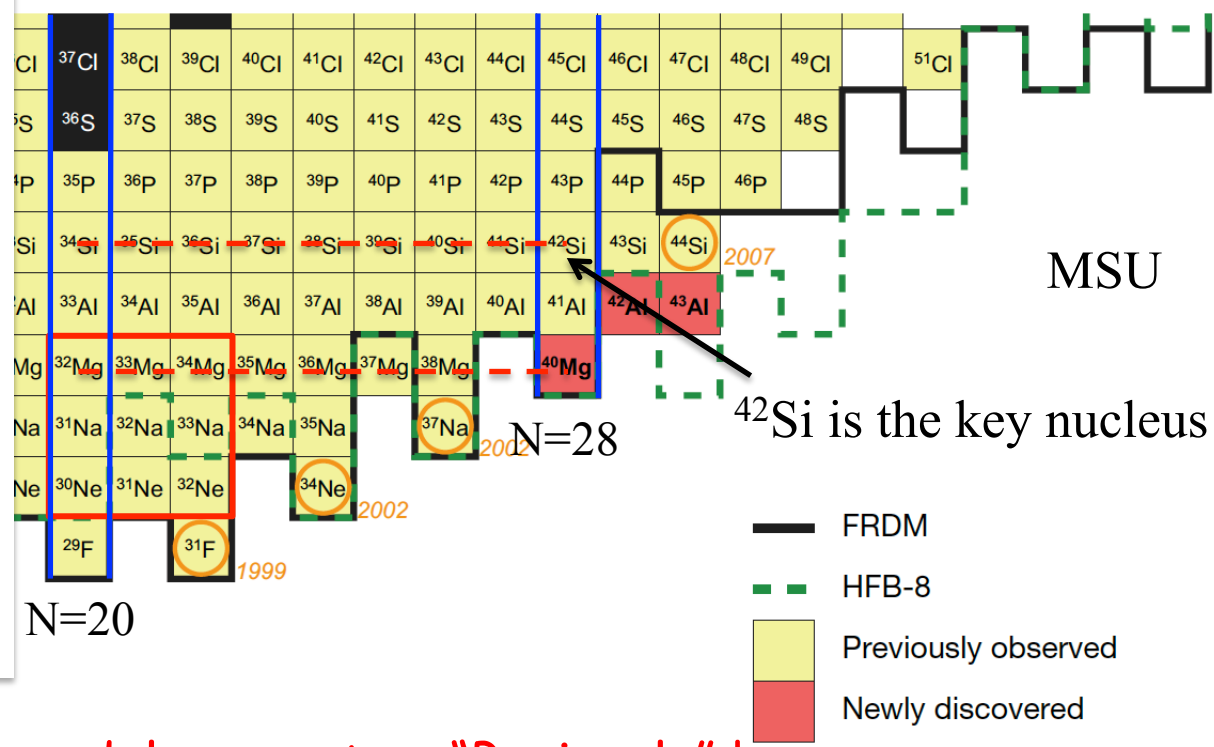
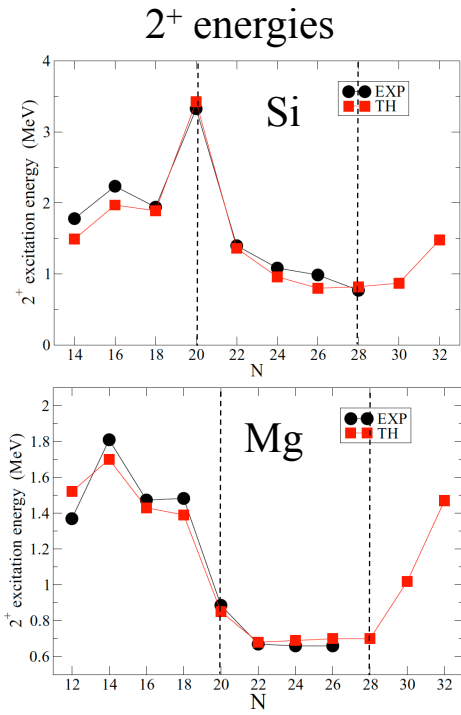
LETTERS

A second "island of inversion" at N=28 ?

Discovery of ^{40}Mg and ^{42}Al suggests neutron drip-line slant towards heavier isotopes

2007

T. Baumann¹, A. M. Amthor^{1,2}, D. Bazin¹, B. A. Brown^{1,2}, C. M. Folden III¹, A. Gade^{1,2}, T. N. Ginter¹, M. Hausmann¹, M. Portillo¹, A. Schiller¹, B. M. Sherrill^{1,2}, A. Stolz¹, O. B. Tarasov^{1,4} & M. Thoennessen^{1,2}



Shell model suggests a "Peninsula" !

What about the N=28 shell closure ?

More complete experimental information needed !



PRL 99, 022503 (2007) PHYSICAL REVIEW LETTERS week ending 13 JULY 2007

2007

Collapse of the $N = 28$ Shell Closure in ^{42}Si

B. Bastin,² S. Grévy,^{1,*} D. Sohler,³ O. Sorlin,¹ D. Baiborodin,⁵ R. Borcea,⁶ C. Bourgeois,⁴ A. Buta,⁶ Z. Elekes,³ S. Franchoo,⁴ S. Iacob,⁶ B. Laurent,² M. N. A. Orr,² Y. Penionzhkevich,¹⁰ Z. M. G. Saint-Laur



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Available online at www.sciencedirect.com



Physics Letters B 649 (2007) 43–48

PHYSICS LETTERS B

www.elsevier.com/locate/physletb

2007

Mass measurements of neutron-rich nuclei near the $N = 20$ and 28 shell closures

PRL 109, 182501 (2012) PHYSICAL REVIEW LETTERS week ending 2 NOVEMBER 2012

2012

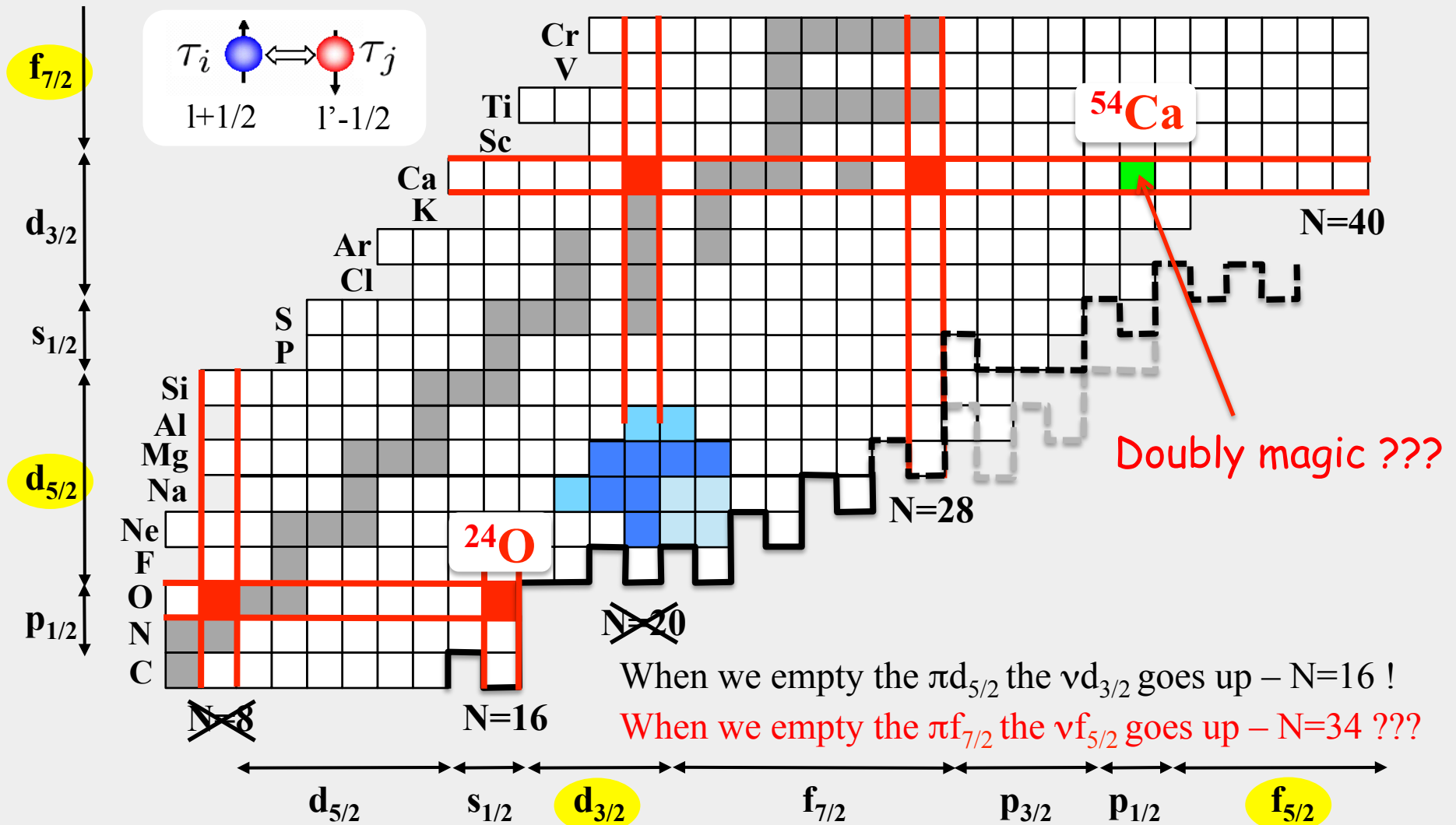
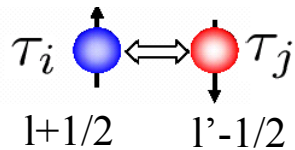
Well Developed Deformation in ^{42}Si

S. Takeuchi,^{1,*} M. Matsushita,^{1,2,†} N. Aoi,^{1,‡} P. Doornenbal,¹ K. Li,^{1,3} T. Motobayashi,¹ H. Scheit,^{1,§} D. Steppenbeck,^{1,†} H. Wang,^{1,3} H. Baba,¹ D. Bazin,⁴ L. C aceres,⁵ H. Crawford,⁶ P. Fallon,⁶ R. Gernh user,⁷ J. Gibelin,⁸ S. Go,⁹ S. Gr vy,⁵ C. Hinke,⁷ C. R. Hoffman,¹⁰ R. Hughes,¹¹ E. Ideguchi,^{9,‡} D. Jenkins,¹² N. Kobayashi,¹³ Y. Kondo,¹³ R. Kr ucken,^{7,||} T. Le Bleis,^{14,15,¶} J. Lee,¹ G. Lee,¹³ A. Matta,¹⁶ S. Michimasa,⁹ T. Nakamura,¹³ S. Ota,⁹ M. Petri,^{6,§} T. Sako,¹³ H. Sakurai,¹ S. Shimoura,⁹ K. Steiger,⁷ K. Takahashi,¹³ M. Takechi,^{1,**} Y. Togano,^{1,**} R. Winkler,^{4,††} and K. Yoneda¹

homaz^a, D. Baiborodin^c,
bert^f, L. Giot^a, A. Khouaja^a,
h^h, S. Pita^a, M. Rousseau^a,

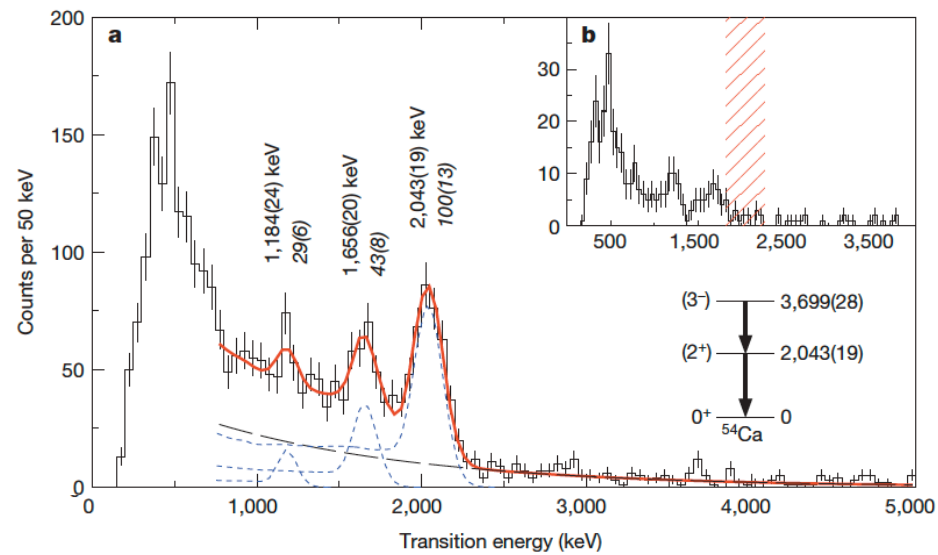
Is N=34 a new magic number far-off stability ?

monopole effects of the tensor force



Evidence for a new nuclear ‘magic number’ from the level structure of ^{54}Ca

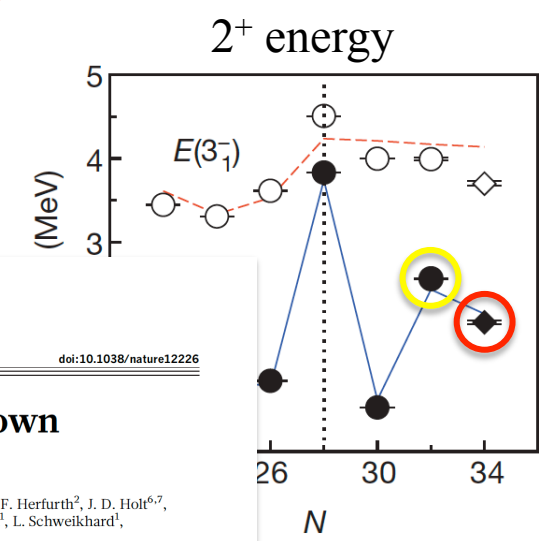
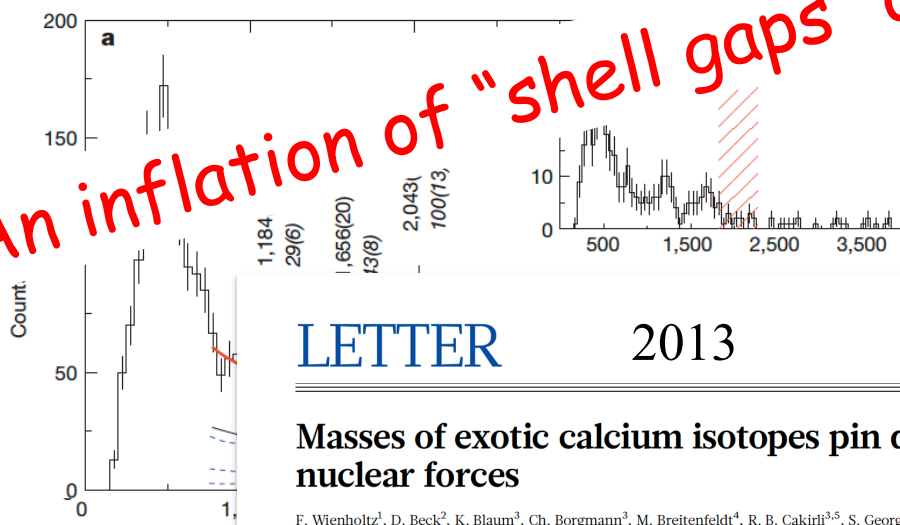
D. Steppenbeck¹, S. Takeuchi², N. Aoi³, P. Doornenbal², M. Matsushita¹, H. Wang², H. Baba², N. Fukuda², S. Go¹, M. Honma⁴, J. Lee², K. Matsui⁵, S. Michimasa¹, T. Motobayashi², D. Nishimura⁶, T. Otsuka^{1,5}, H. Sakurai^{2,5}, Y. Shiga⁷, P.-A. Söderström², T. Sumikama⁸, H. Suzuki², R. Taniuchi⁵, Y. Utsuno⁹, J. J. Valiente-Dobón¹⁰ & K. Yoneda²



Evidence for a new nuclear ‘magic number’ from π^+ level structure of ^{54}Ca

D. Steppenbeck¹, S. Takeuchi², N. Aoi³, P. Doornenbal², M. Matsushita¹, J. Lee², K. Matsui³, S. Michimasa⁴, T. Motobayashi², D. Nishimura⁵, T. Sumikama⁸, H. Suzuki², R. Taniuchi⁵, Y. Utsuno⁹, M. Walterström²

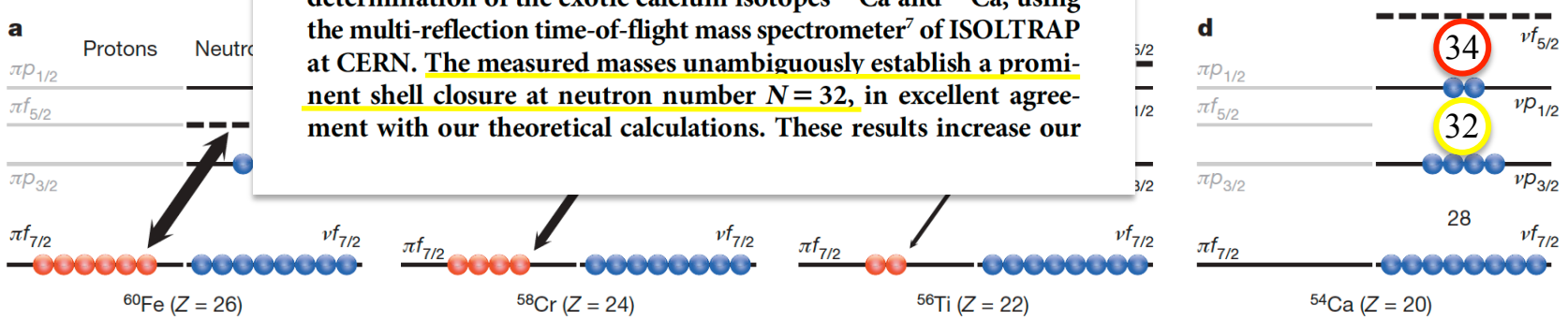
An inflation of "shell gaps" and "magic numbers"!



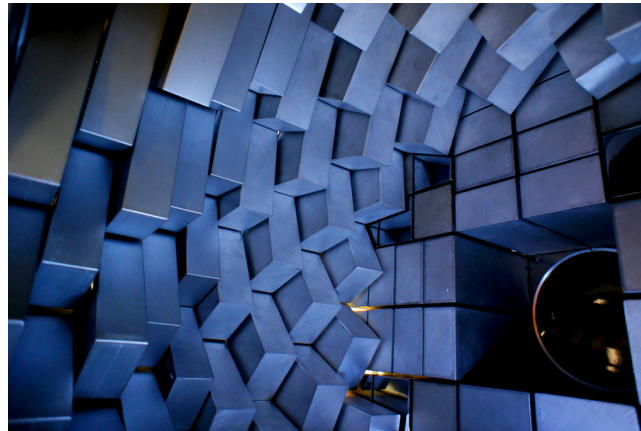
Masses of exotic calcium isotopes pin down nuclear forces

F. Wienholtz¹, D. Beck², K. Blaum³, Ch. Borgmann³, M. Breitenfeldt⁴, R. B. Cakiri^{3,5}, S. George¹, F. Herfurth², J. D. Holt^{6,7}, M. Kowalska⁸, S. Kreim^{3,8}, D. Lunney⁹, V. Manea⁹, J. Menéndez^{6,7}, D. Neidherr², M. Rosenbusch¹, L. Schweikhard¹, A. Schwenk^{7,6}, J. Simonis^{6,7}, J. Stanja¹⁰, R. N. Wolf¹ & K. Zuber¹⁰

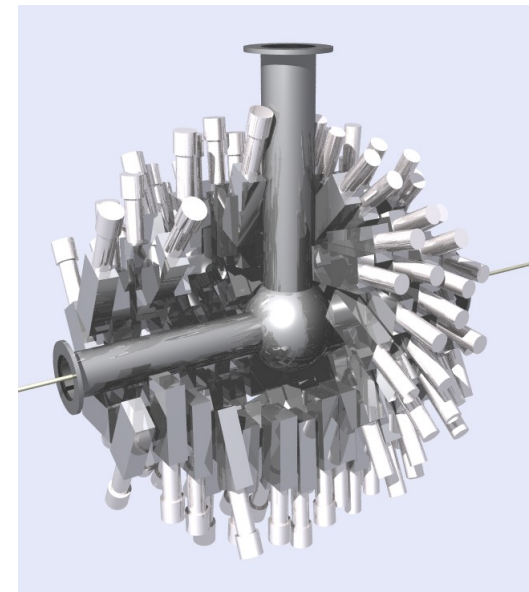
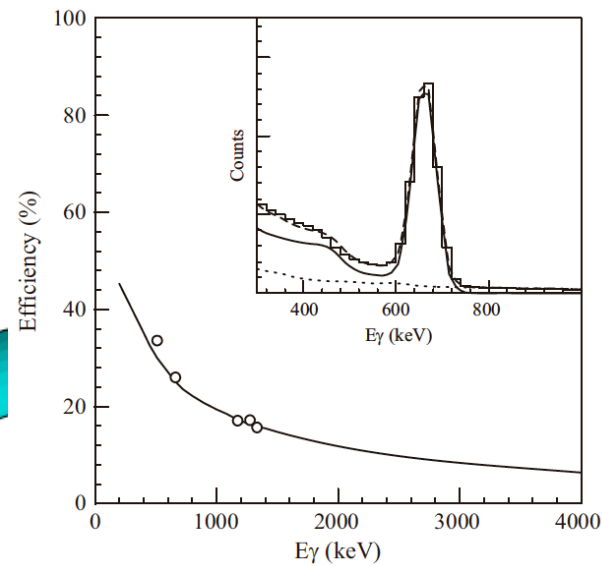
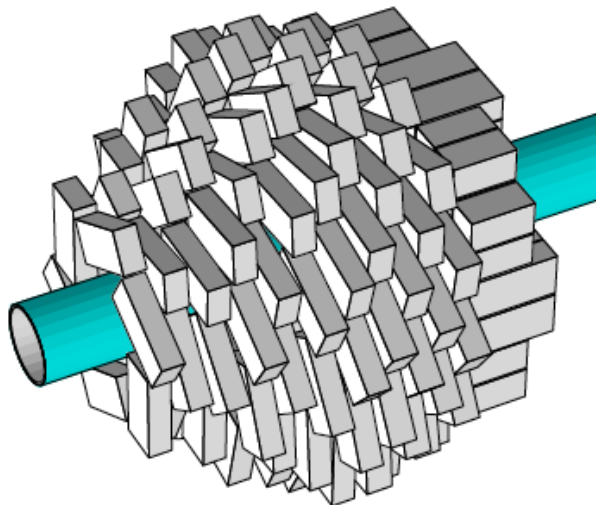
masses evolve for heavier calcium isotopes. Here we report the mass determination of the exotic calcium isotopes ^{53}Ca and ^{54}Ca , using the multi-reflection time-of-flight mass spectrometer⁷ of ISOLTRAP at CERN. The measured masses unambiguously establish a prominent shell closure at neutron number $N = 32$, in excellent agreement with our theoretical calculations. These results increase our



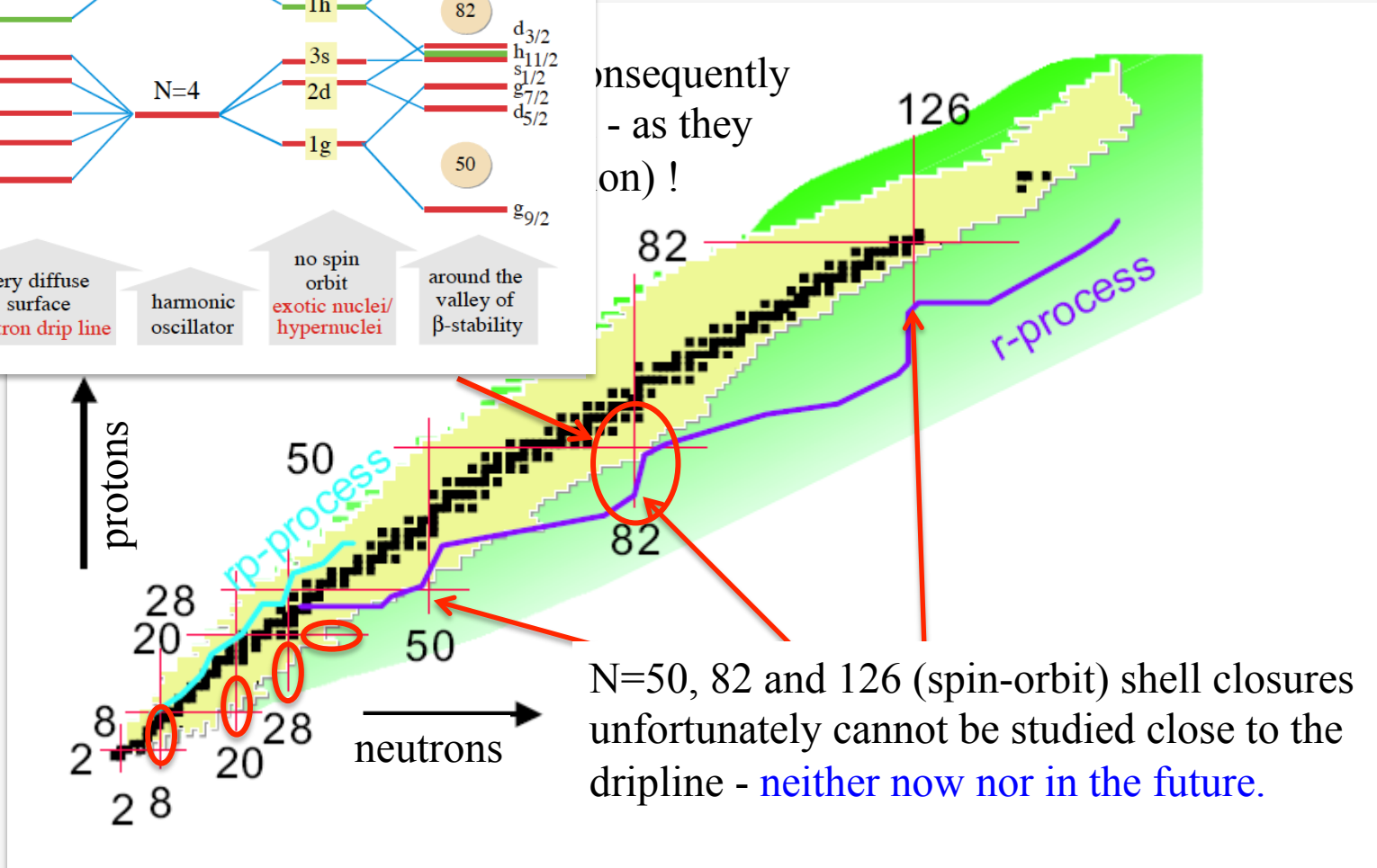
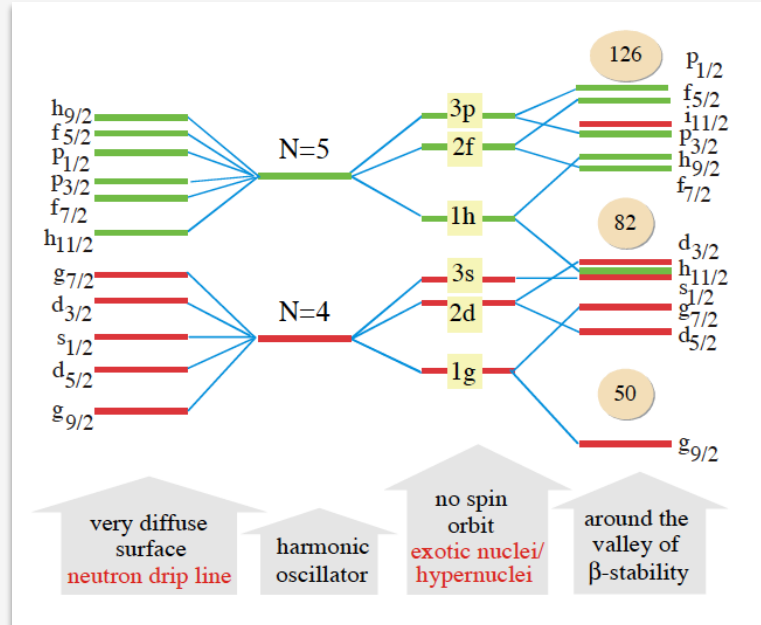
The γ -ray spectrometer DALI2 at RIKEN



$^{186}\text{NaI}(\text{Tl})$
scintillators



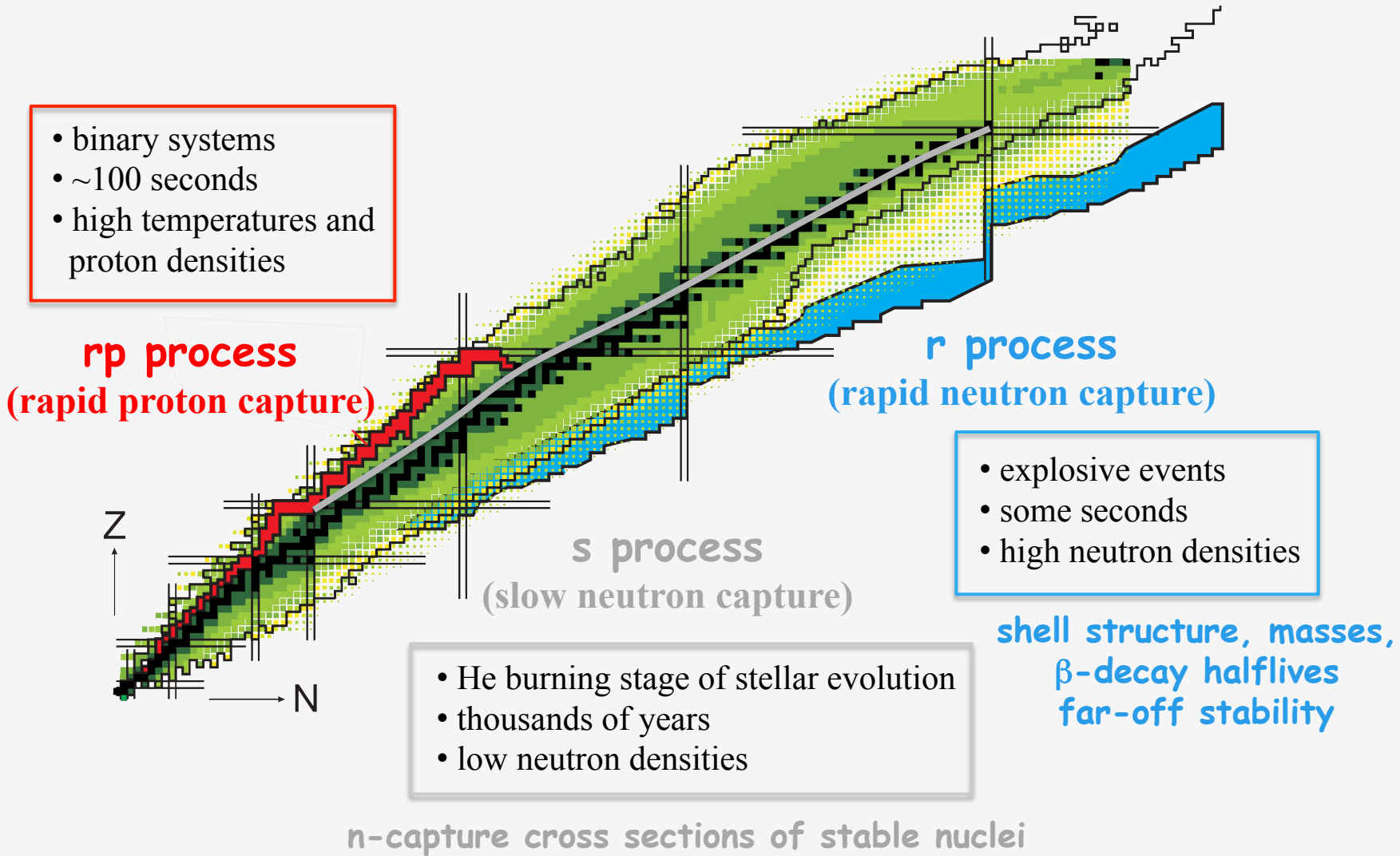
neutron-rich region



consequently
- as they
on) !

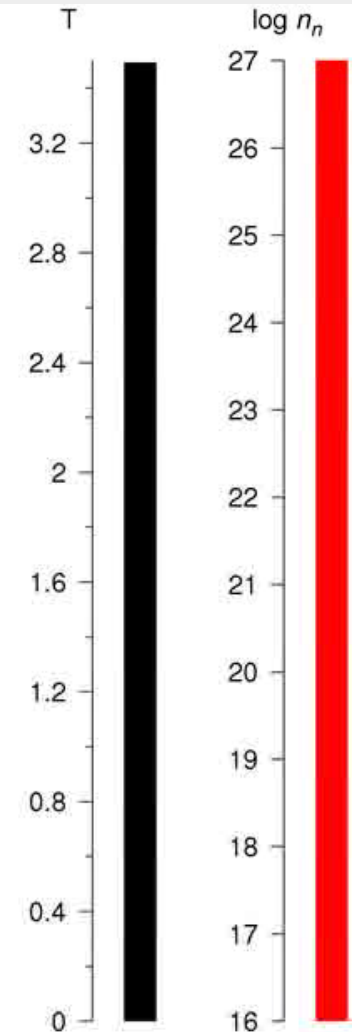
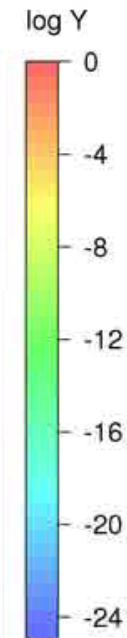
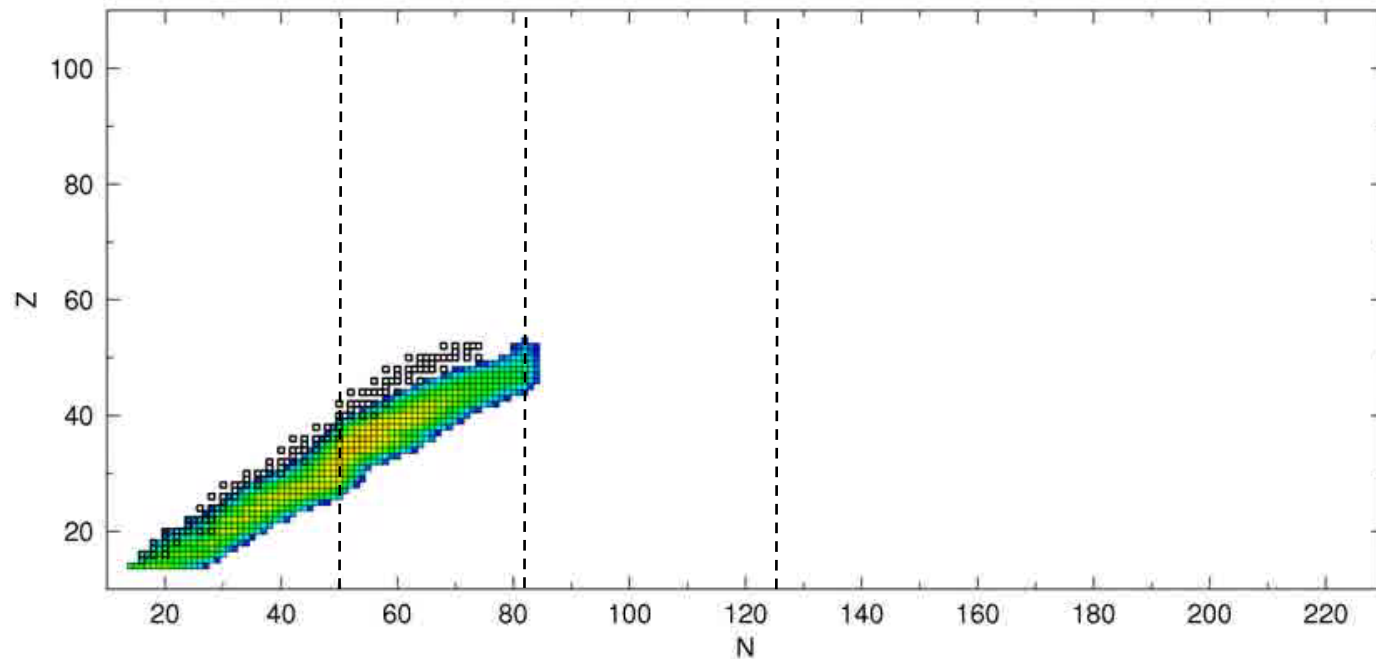
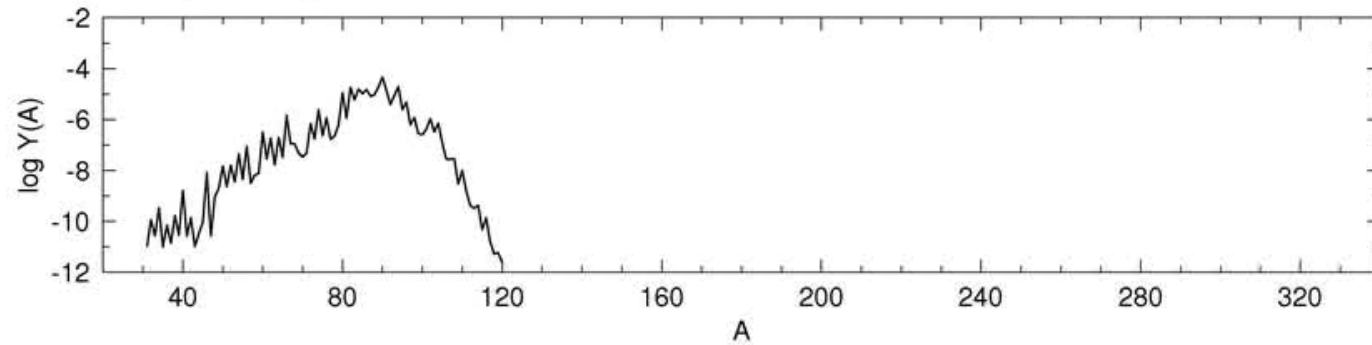
$N=50$, 82 and 126 (spin-orbit) shell closures unfortunately cannot be studied close to the dripline - neither now nor in the future.

Nucleosynthesis of the heaviest elements



Realistic r process calculations (G. Martínez Pinedo)

$T_9=3.49$, $n_r=1.377e+27$, $t=1.471e-14$ s



Most nuclear input data are theoretical estimates !

The importance of the N=82,126 shell closures

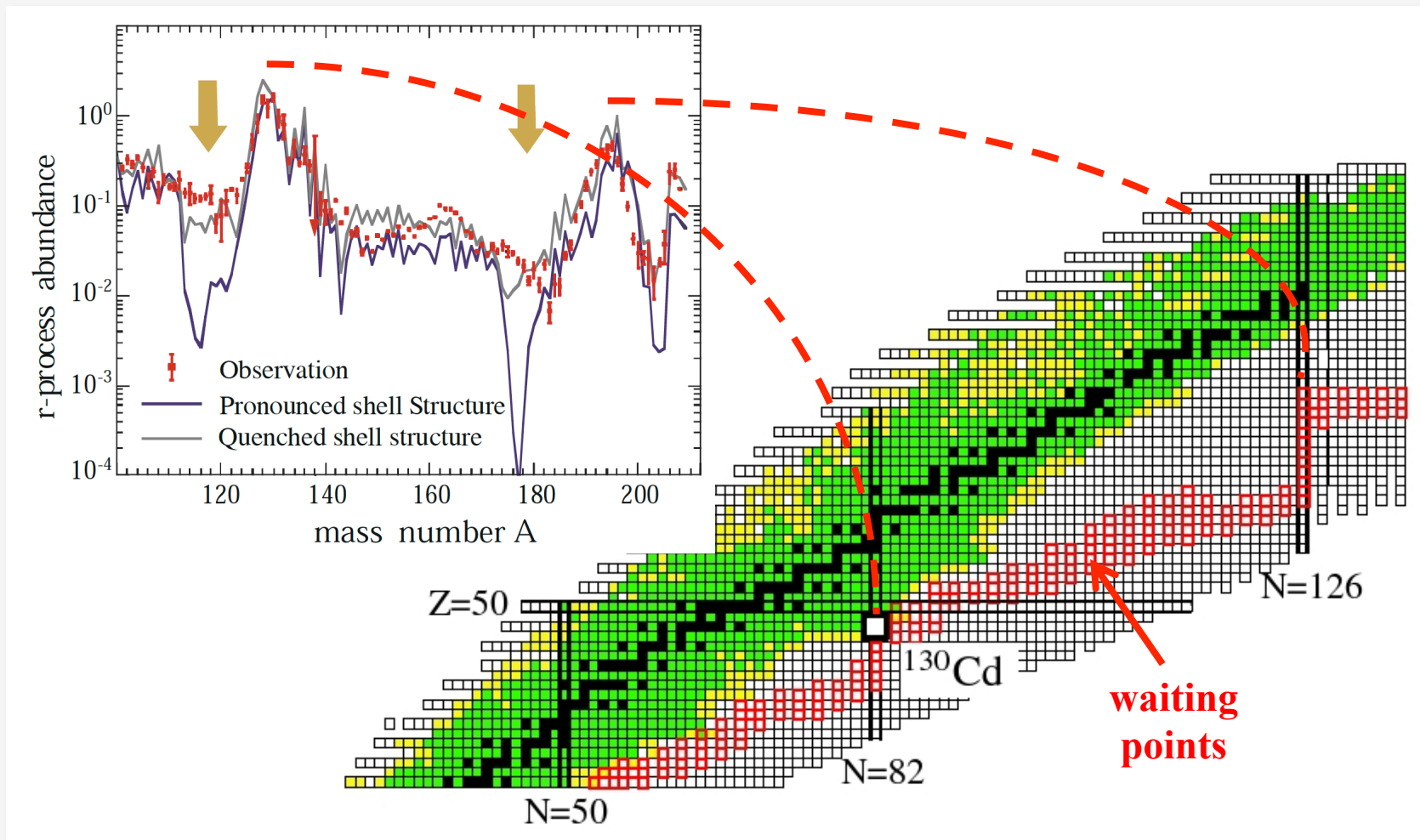


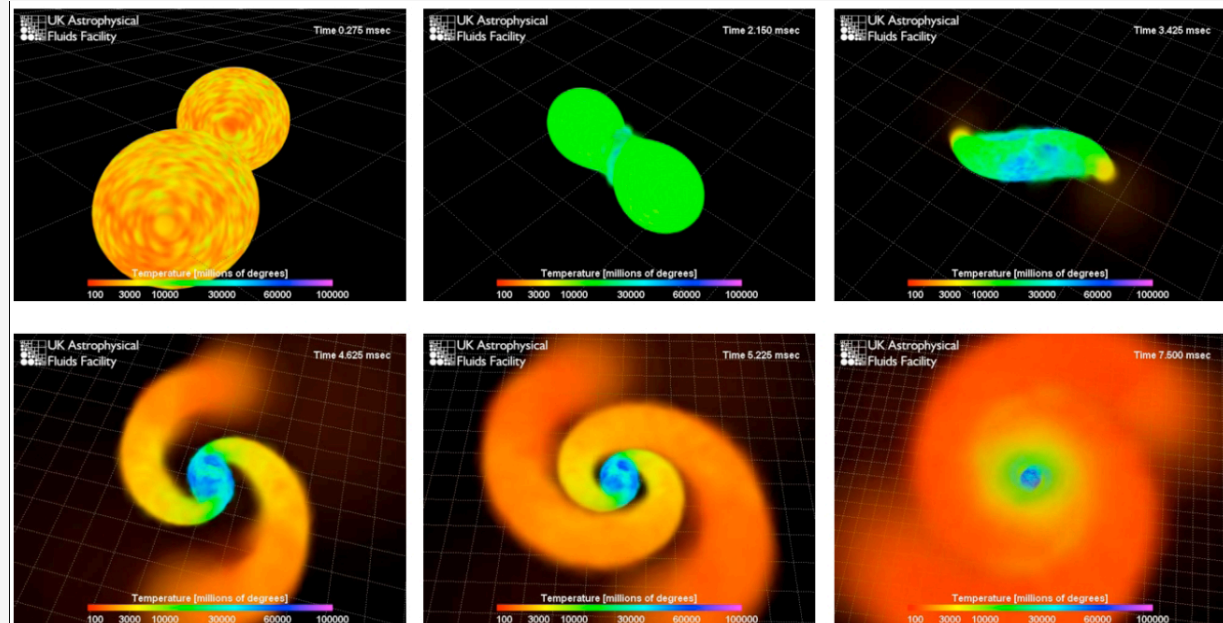
Figure shown in nearly all glossy brochures for future facilities ...

Supernovae remnant Crab nebula

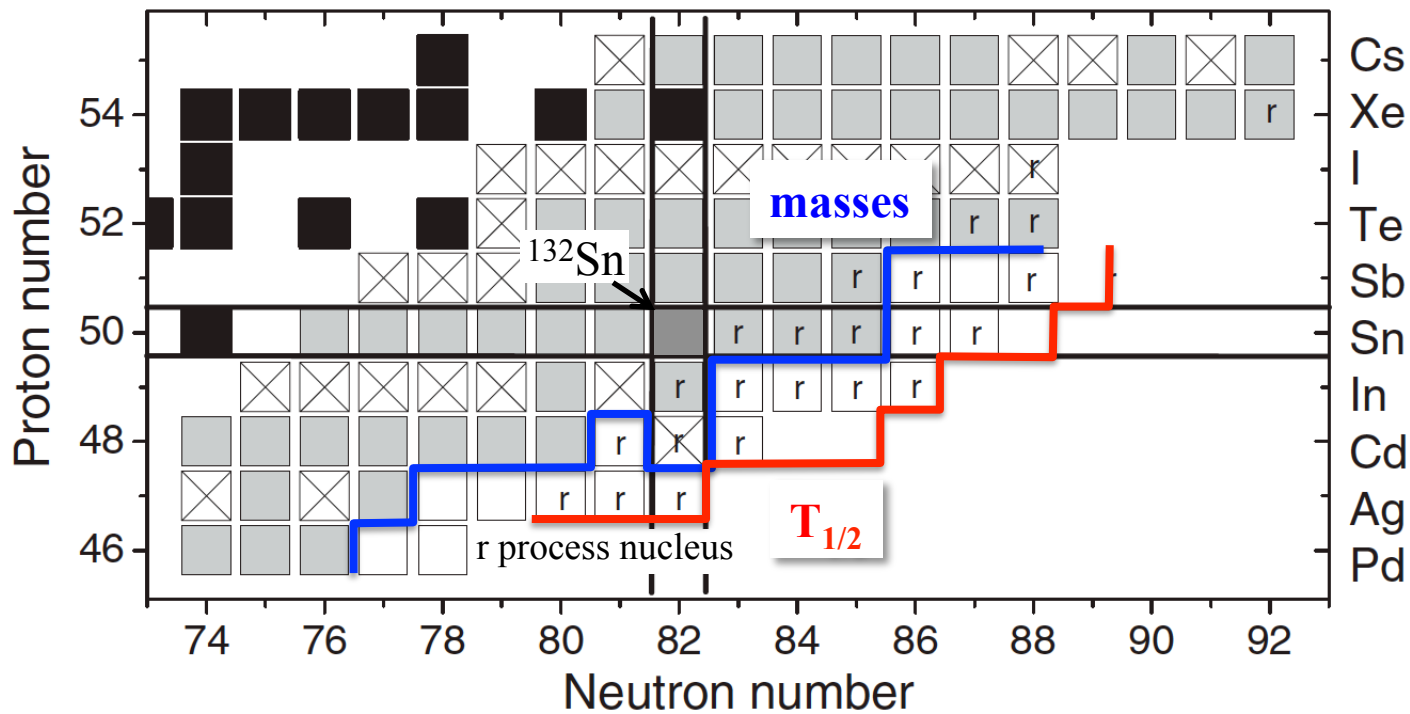


Site of r-process still an open question !

Merging neutron stars



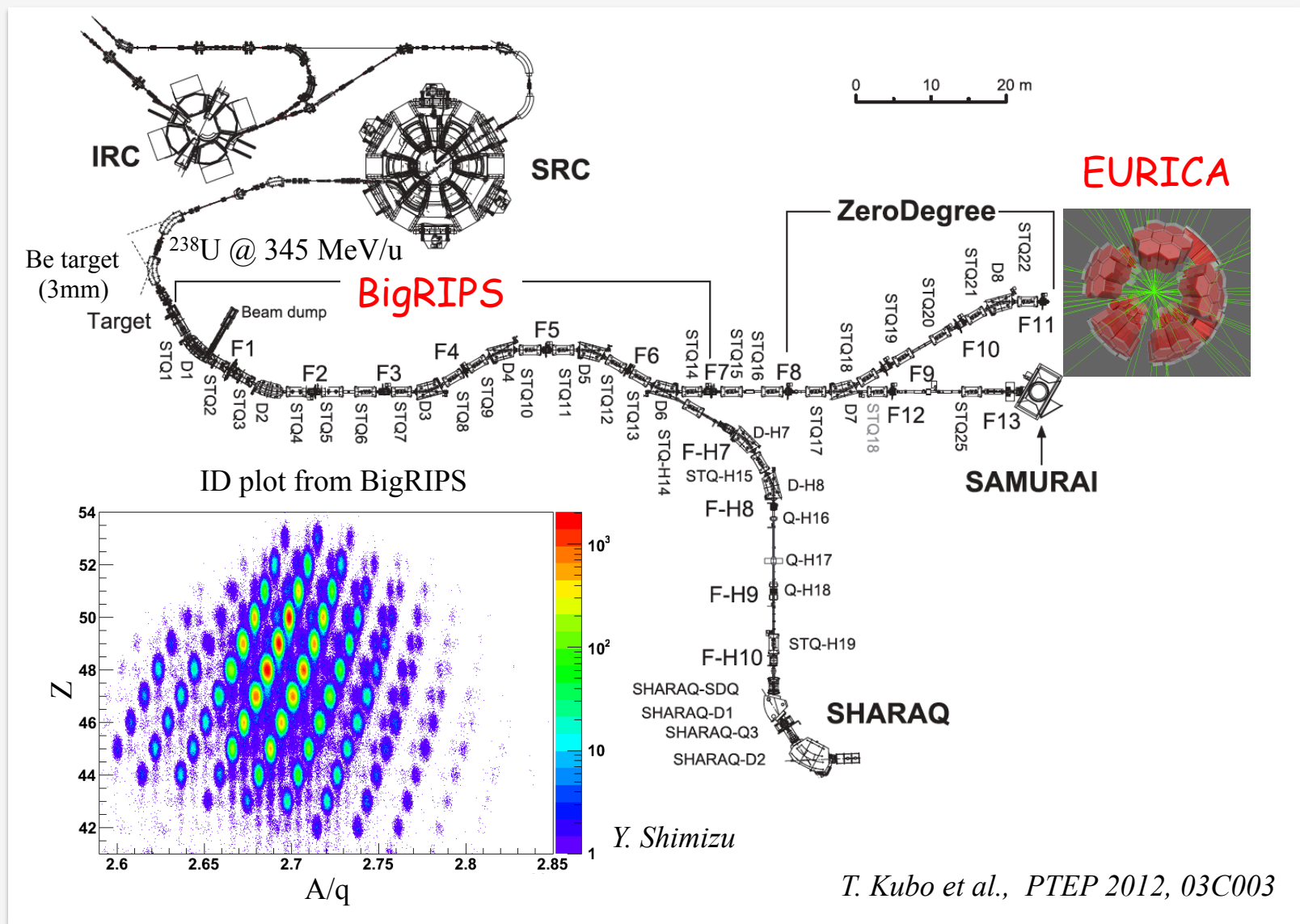
Experimental status of the ^{132}Sn region



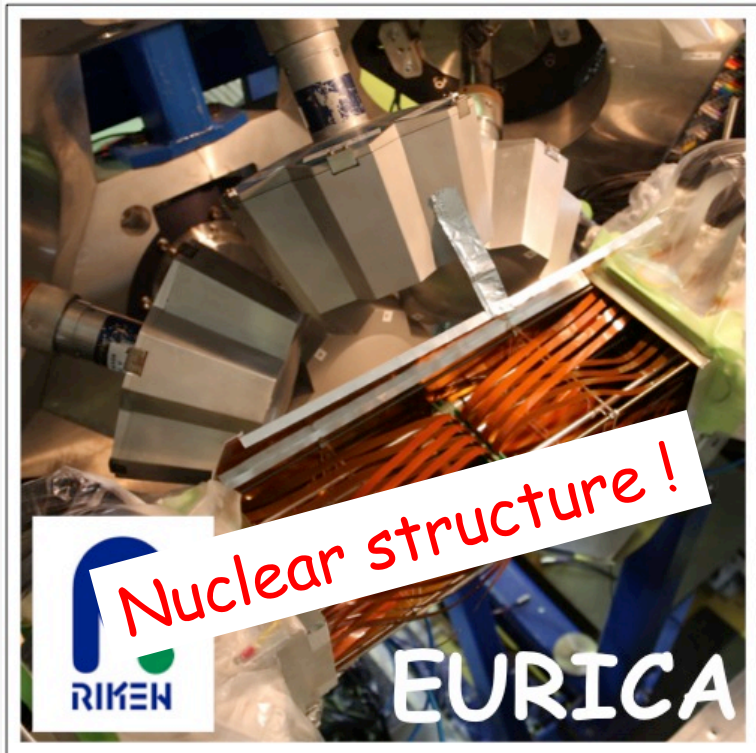
J. Hakala et al., Phys. Rev. Lett. 109 (2012) 032501

Need many more experimental half-lives and masses !

The EURICA project at RIKEN - since 2012



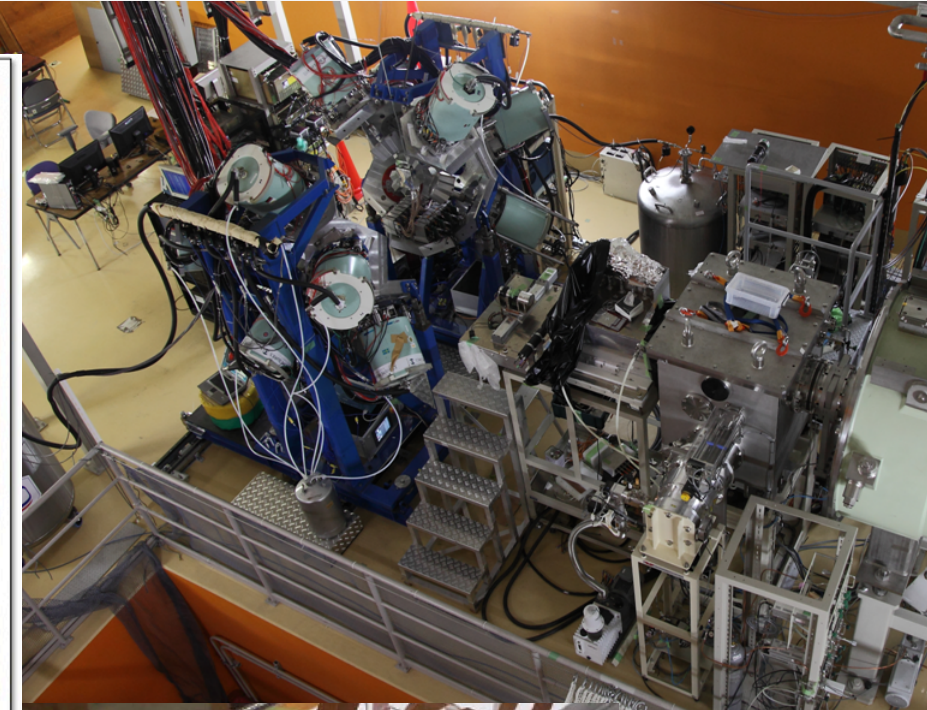
The EURICA project at RIKEN - since 2012



EURICA

12 EUROBALL Cluster detectors, i.e. 84 Ge crystals

γ decay of isomeric states
or after β decay

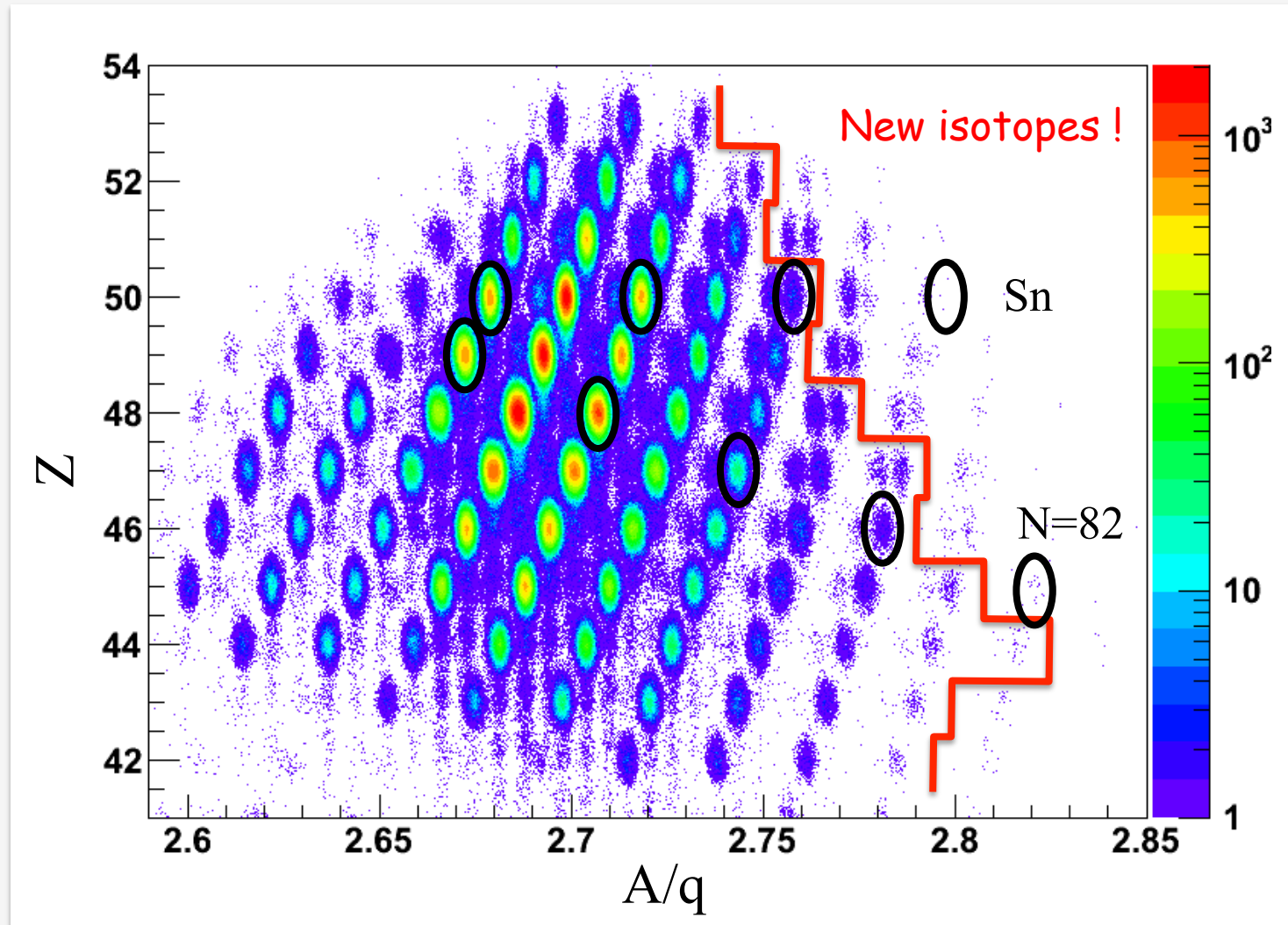


WAS3ABi

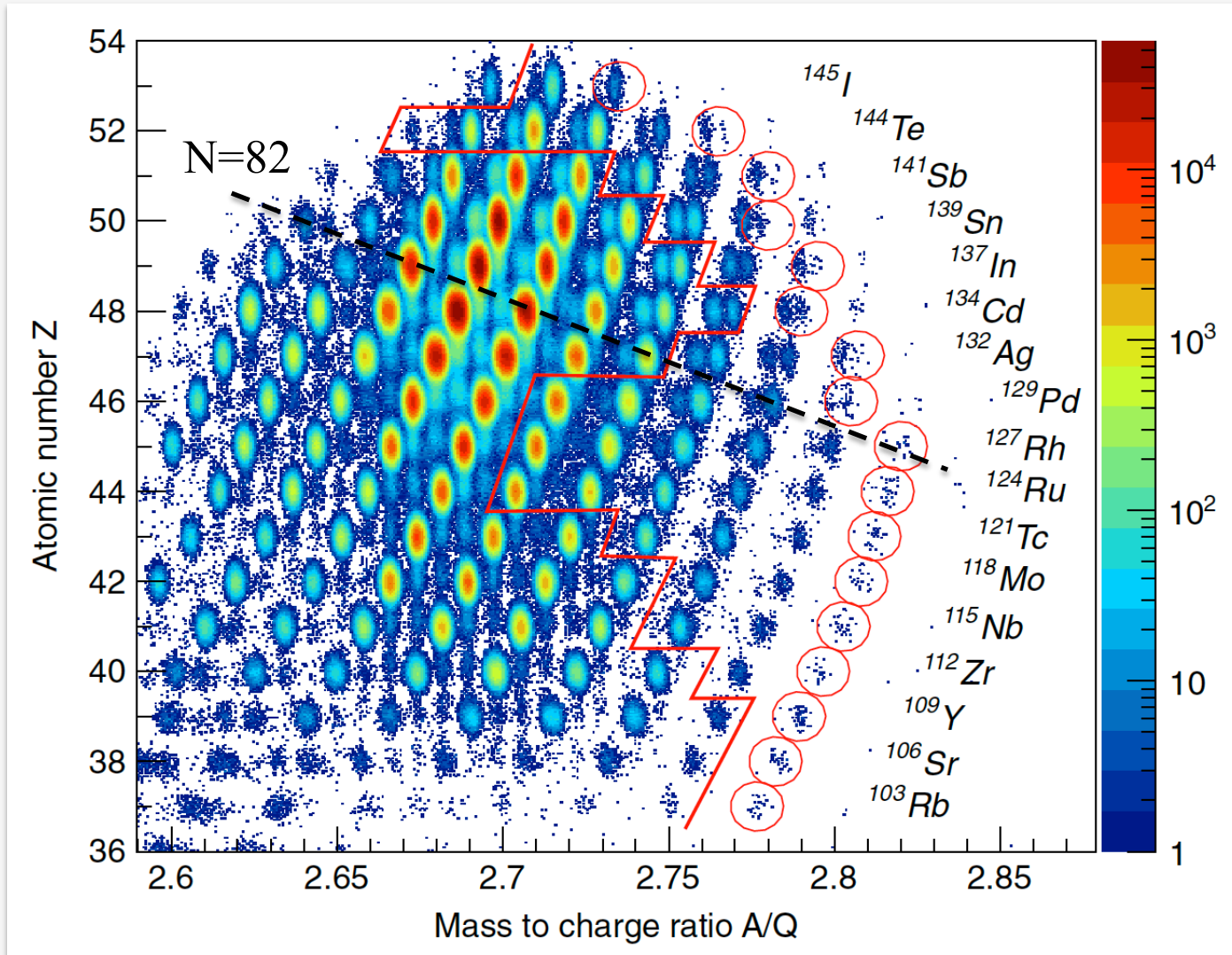
Stack of segmented
Si detectors

ion implantation
 β decay

BigRIPS ID plot: Sum of two settings

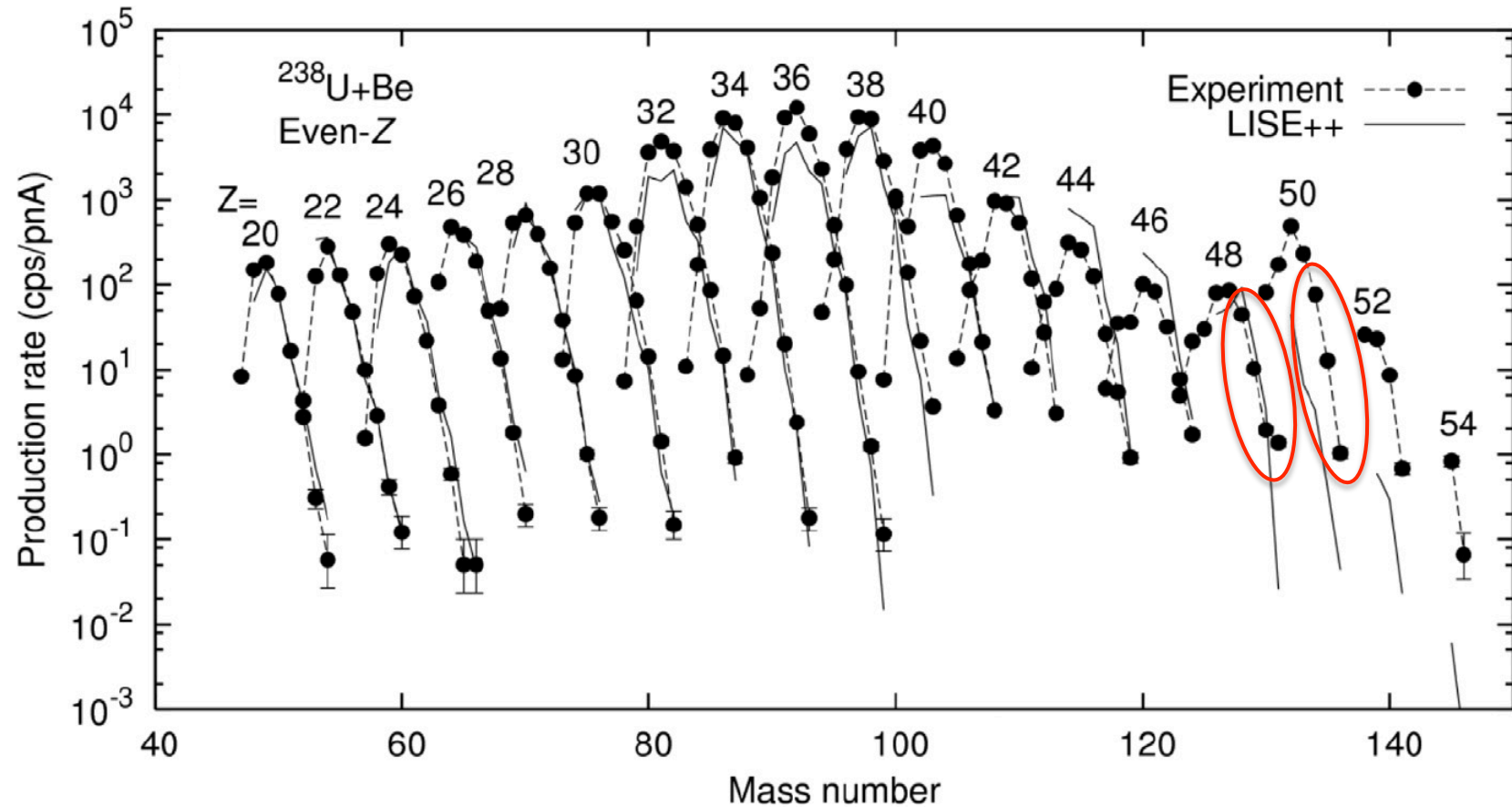


Systematic half-life measurement in the ^{132}Sn region



110 half-lives measured, 40 for the first time !

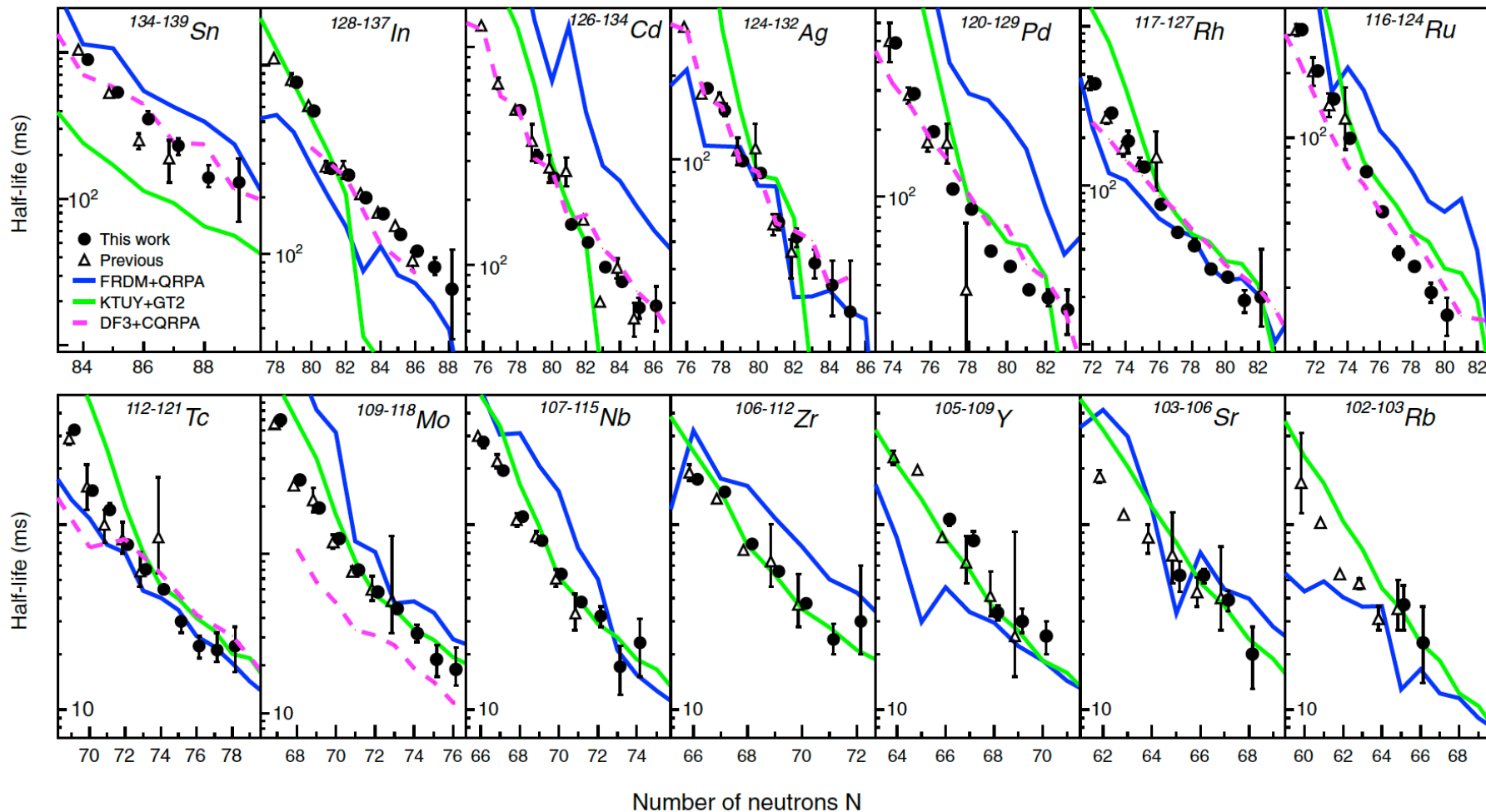
Production cross sections for $^{238}\text{U}+\text{Be}$ @ 345 MeV



H. Suzuki et al./Nuclear Instruments and Methods in Physics Research B 317 (2013) 756–768

One order of magnitude in beam intensity means 1-2 neutrons further out !

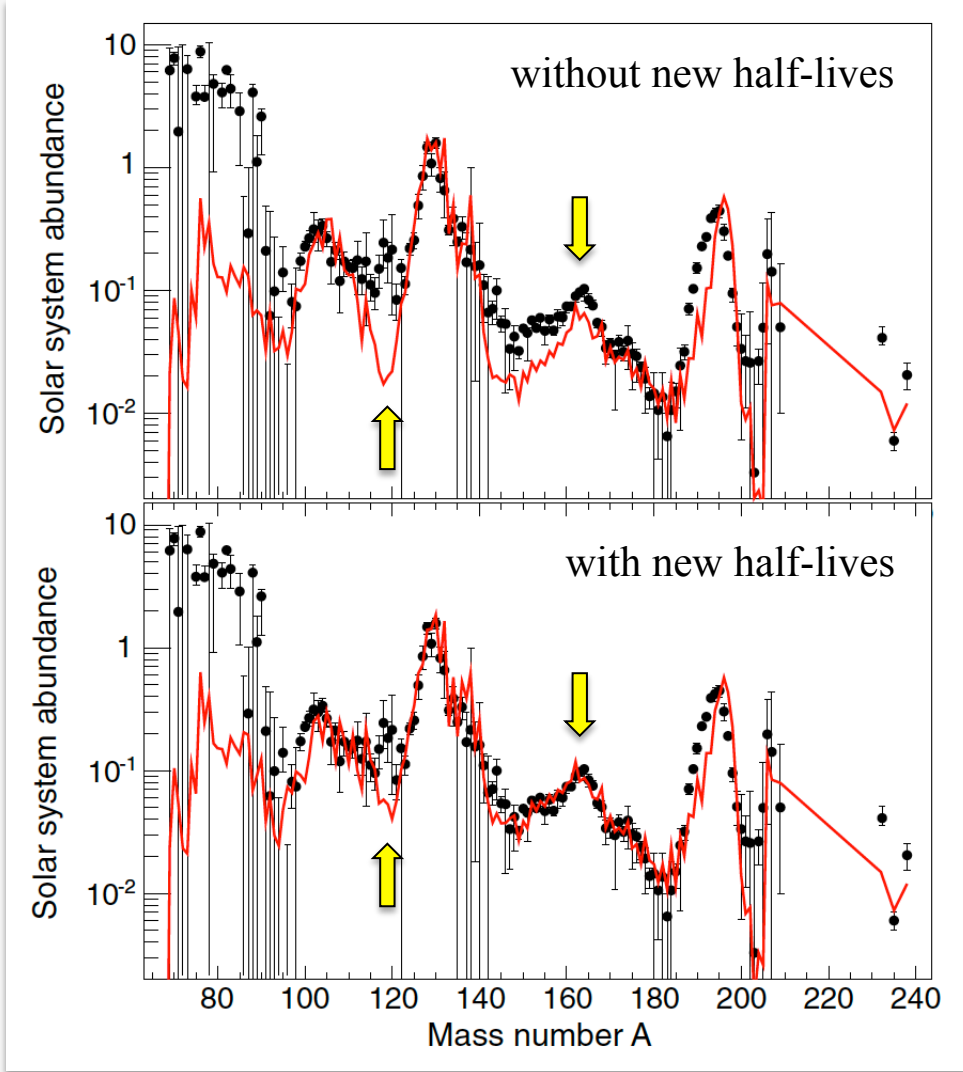
Systematic half-life measurement in the ^{132}Sn region



G. Lorusso et al., Phys. Rev. Lett. 114, 192501 (2015)

Systematic measurements crucial to test theoretical model predictions for inaccessible regions of the nuclear chart !

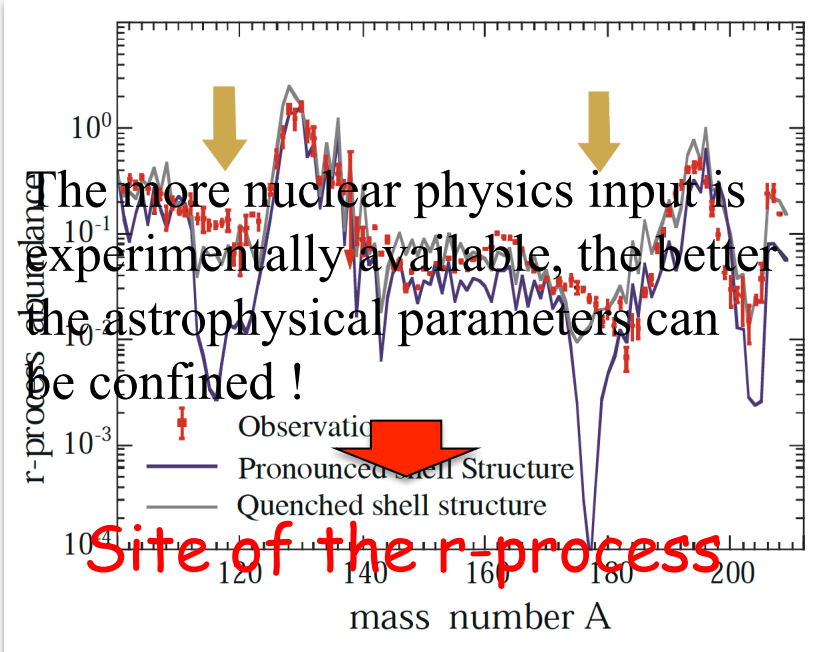
The r-process solar system abundance pattern



G. Lorusso et al., Phys. Rev. Lett. 114, 192501 (2015)

New experimental half-lives

- alleviate the underproduction just below the $A \sim 130$ peak and
- greatly improve the description of the rare earth element abundances.



The more nuclear physics input is experimentally available, the better the astrophysical parameters can be confined !

Site of the r-process

No need for fancy explications !