



The HVP contribution to a_μ from full lattice QCD

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HPQCD collaboration

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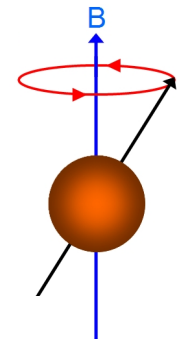
Work with: Bipasha Chakraborty, Jonna Koponen,
Peter Lepage and others ..



Using the Darwin (9600 core) Sandybridge/infiniband
cluster at Cambridge, part of STFC's DiRAC HPC
facility

Muon anomalous magnetic moment

$$\vec{\mu} = g \frac{e}{2m} \vec{S} \quad a_{\mu} = \frac{g - 2}{2}$$



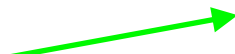
Measure using polarised muons circulating in E and B fields. At a momentum where $\beta \times E$ terms cancel, difference between precession and cyclotron frequencies:

$$\omega_a = -\frac{e}{m} a_{\mu} B$$

BNL result:

$$a_{\mu}^{expt} = 11659208.9(6.3) \times 10^{-10}$$

E989 (FNAL) will
reduce exptl uncty to
1.6, starting 2017

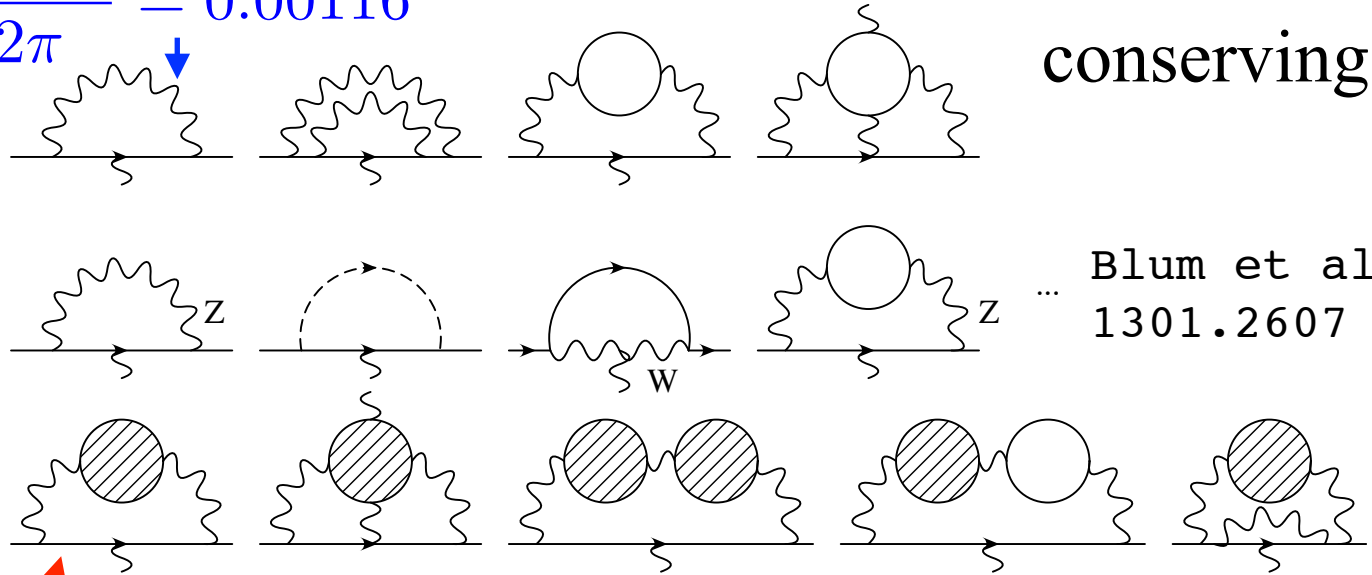


Standard Model theory expectations

Contributions from QED, EW and QCD interactions.

QED dominates.
QCD contriibs start at α_{QED}^2

$$\frac{\alpha_{QED}}{2\pi} = 0.00116$$



flavour and CP conserving

... Blum et al, 1301.2607

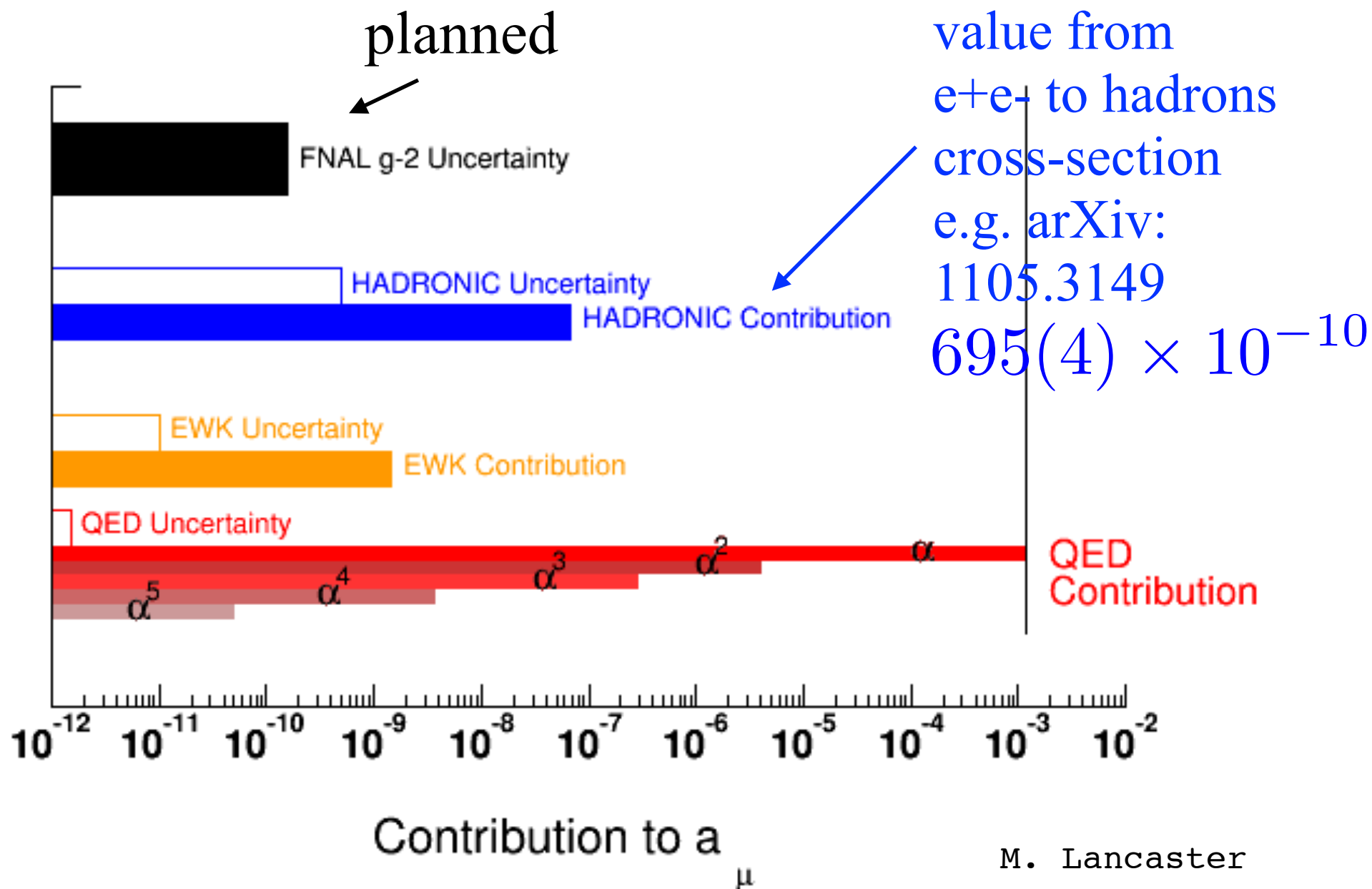
LO Hadronic vacuum polarisation (HVP) dominates uncertainty

$$a_{\mu}^{QED} = 11658471.885(4) \times 10^{-10}$$

$$a_{\mu}^{EW} = 15.4(2) \times 10^{-10}$$

$$a_{\mu}^{E821} = 11659208.9(6.3) \times 10^{-10}$$

Uncertainty dominated by that from hadronic contribns



Hadronic contributions

$$a_{\mu}^{expt} - a_{\mu}^{QED} - a_{\mu}^{EW} = 721.7(6.3) \times 10^{-10}$$
$$= a_{\mu}^{HVP} + a_{\mu}^{HOHVP} + a_{\mu}^{HLbL} + a_{\mu}^{new\ physics}$$

Focus on lowest order hadronic vacuum polarisation,
so assume:

$$a_{\mu}^{HLbL} = 10.5(2.6) \times 10^{-10}$$

$$a_{\mu}^{HOHVP} = -8.85(9) \times 10^{-10} \leftarrow \text{NLO+NNLO}$$

Kurz et al,
1403.6400

$$a_{\mu}^{HVP, no\ new\ physics} = 719.8(6.8) \times 10^{-10}$$

**compare 1105.3149 - discrepancy $24.9(8.0) \times 10^{-10}$

Lattice calculation of HVP

Analytically continue to Euclidean q^2 .

$$a_\mu^{HVP,i} = \frac{\alpha}{\pi} \int_0^\infty dq^2 f(q^2) (4\pi\alpha e_i^2) \hat{\Pi}_i(q^2)$$

connected contribution for flavour i

$f(q^2)$ divergent function with scale set by m_μ

$$\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$$

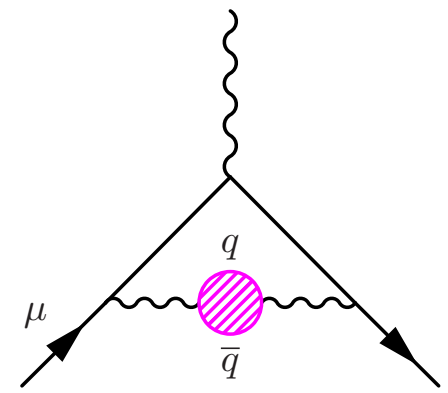
HPQCD method: time-moments of vector JJ correlators give expansion around $q^2=0$

$$G_n \equiv \sum_{t, \vec{x}} t^n Z_V^2 \langle J^j(\vec{x}, t) J^j(0) \rangle$$

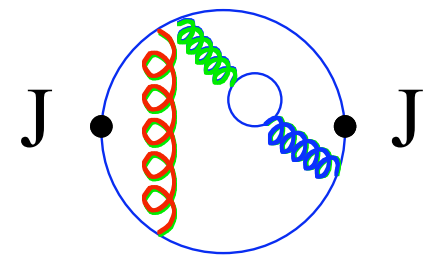
$$\Pi_k = (-1)^{k+1} \frac{G_{2k+2}}{(2k+2)!}$$

$$\hat{\Pi}(q^2) = \sum_{k=1}^{\infty} q^{2k} \Pi_k$$

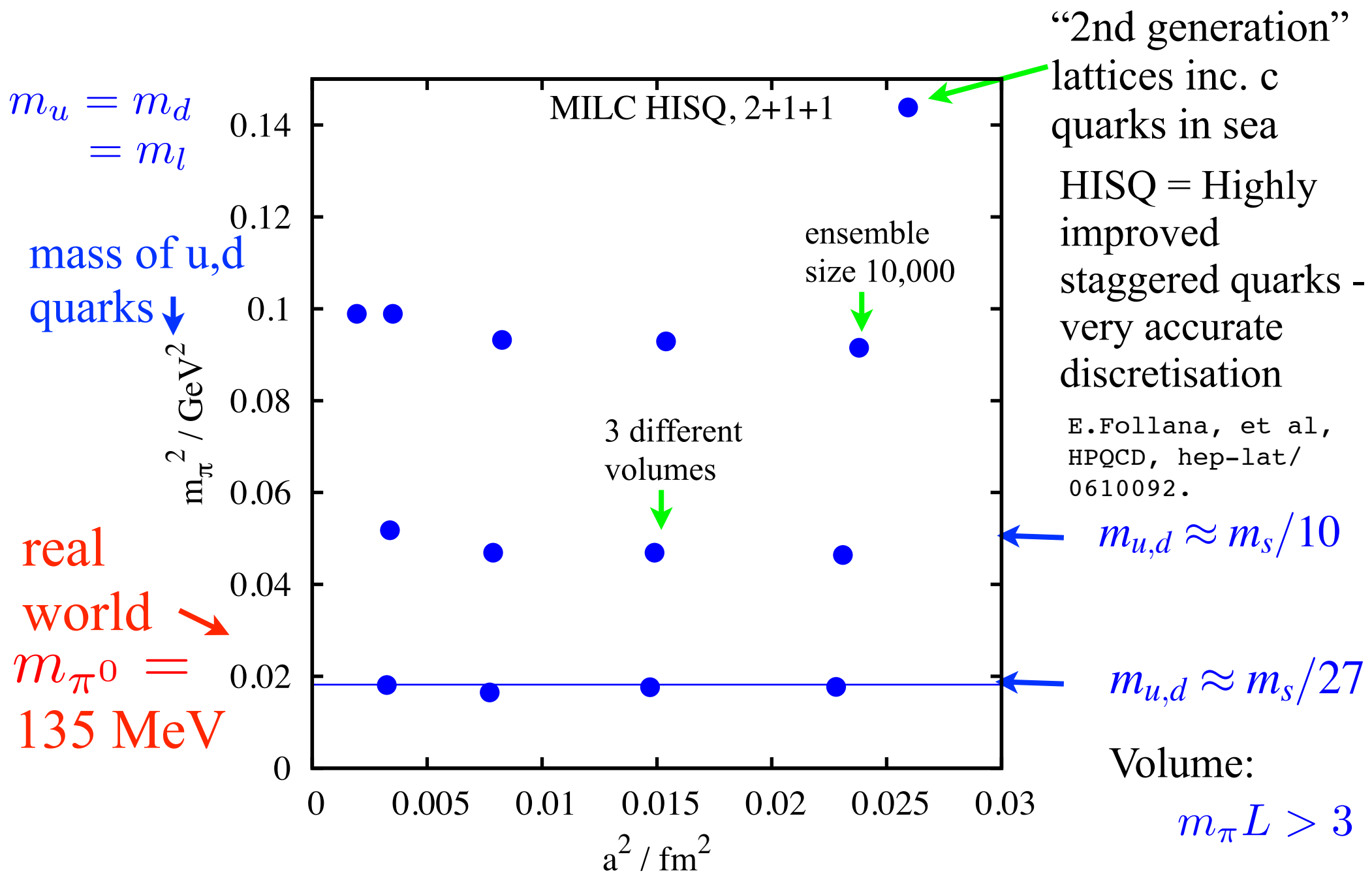
replace with [2,2] Padé



Blum, hep-lat/0212018

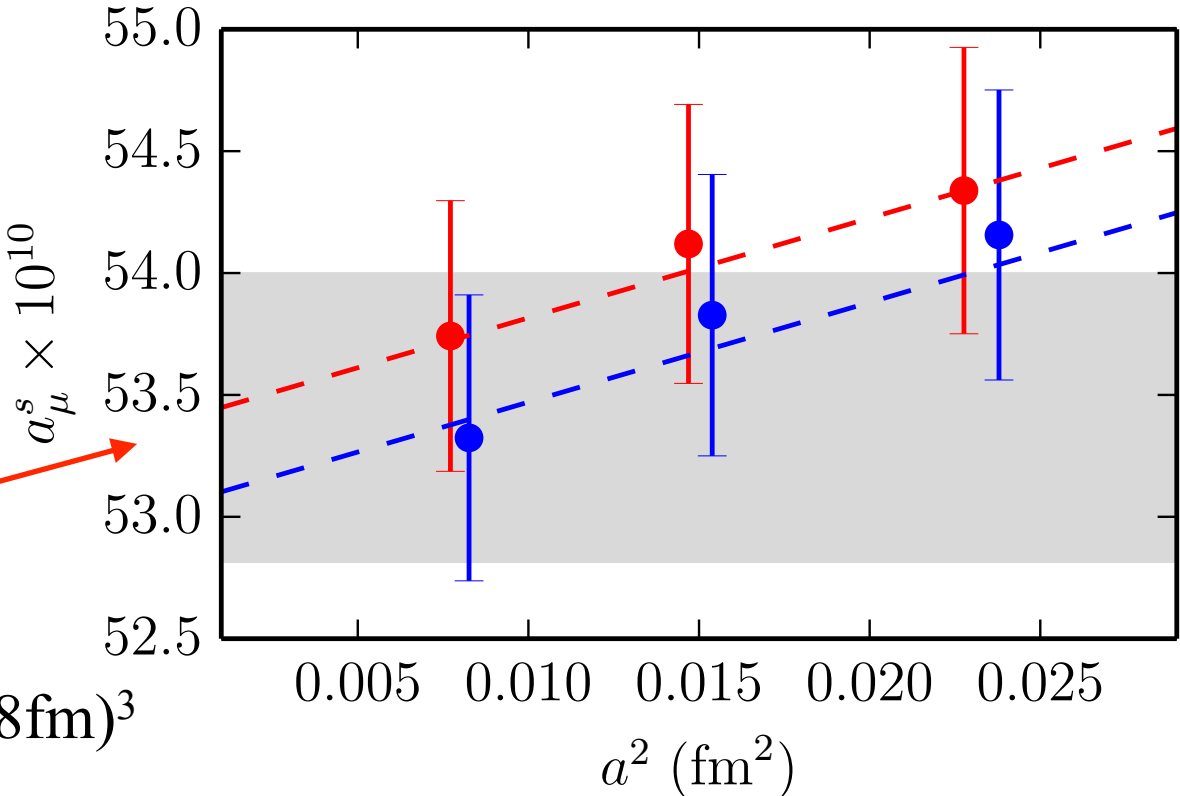


Working with staggered quarks in new ‘2nd generation’ calculations - further improved gluon and quark actions



STRANGE contribution

HISQ valence quarks on MILC 2+1+1 HISQ configs. Local J_V - nonpert. Z_V .
 multiple a (fixed by w_0), m_l (inc. phys.), volumes. Tune s from η_s up to $(5.8\text{fm})^3$

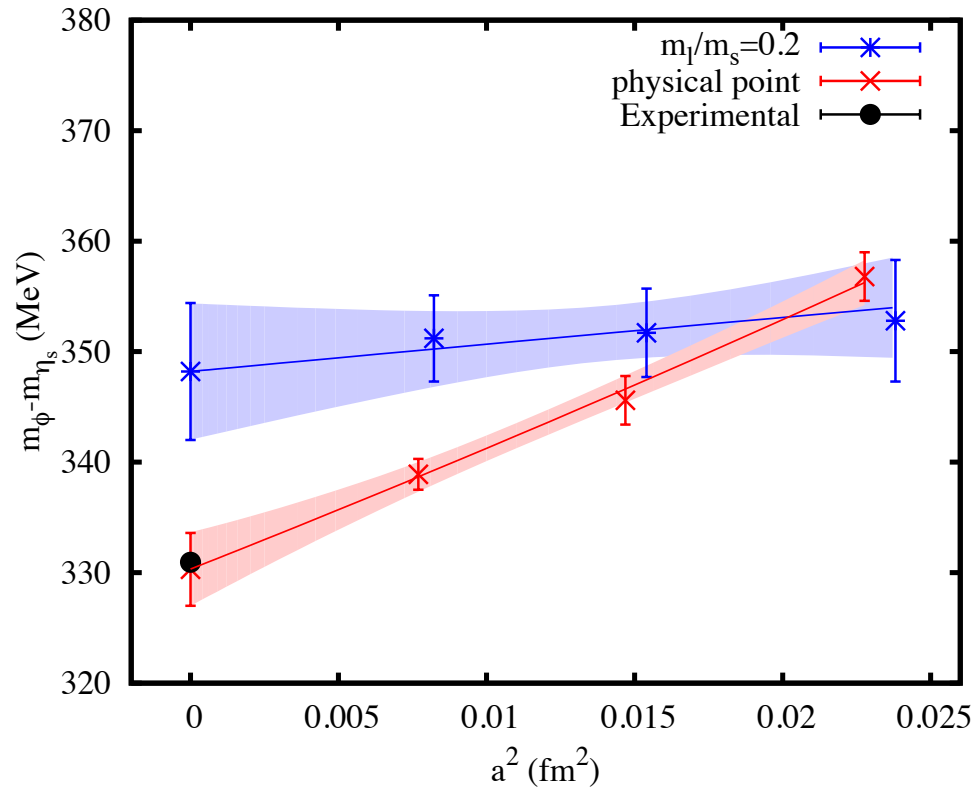


	a_μ^s
Uncertainty in lattice spacing (w_0, r_1):	1.0%
Uncertainty in Z_V :	0.4%
Monte Carlo statistics:	0.1%
$a^2 \rightarrow 0$ extrapolation:	0.1%
QED corrections:	0.1%
Quark mass tuning:	0.0%
Finite lattice volume:	< 0.1%
Padé approximants:	< 0.1%
Total:	1.1%

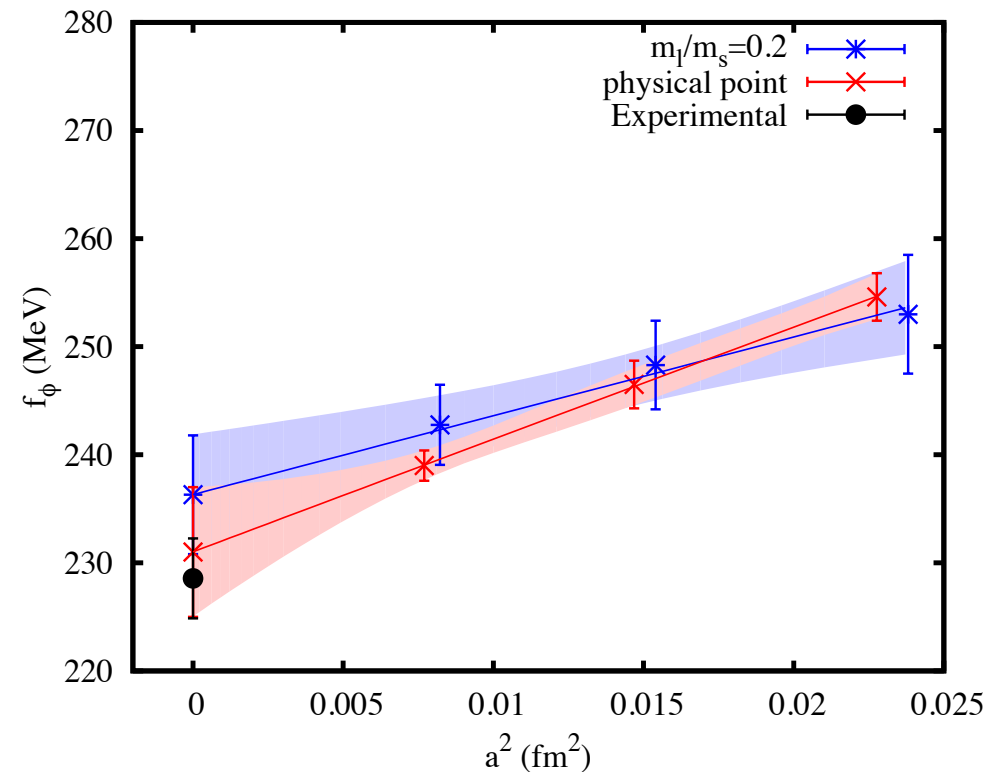
$$a_{\mu, \text{lat}}^s = a_\mu^s \times (1 + c_{a^2}(a\Lambda_{\text{QCD}}/\pi)^2 + c_{\text{sea}}\delta x_{\text{sea}} + c_{\text{val}}\delta x_{\text{val}})$$

$$a_\mu^{\text{HVP},s} = 53.41(59) \times 10^{-10}$$

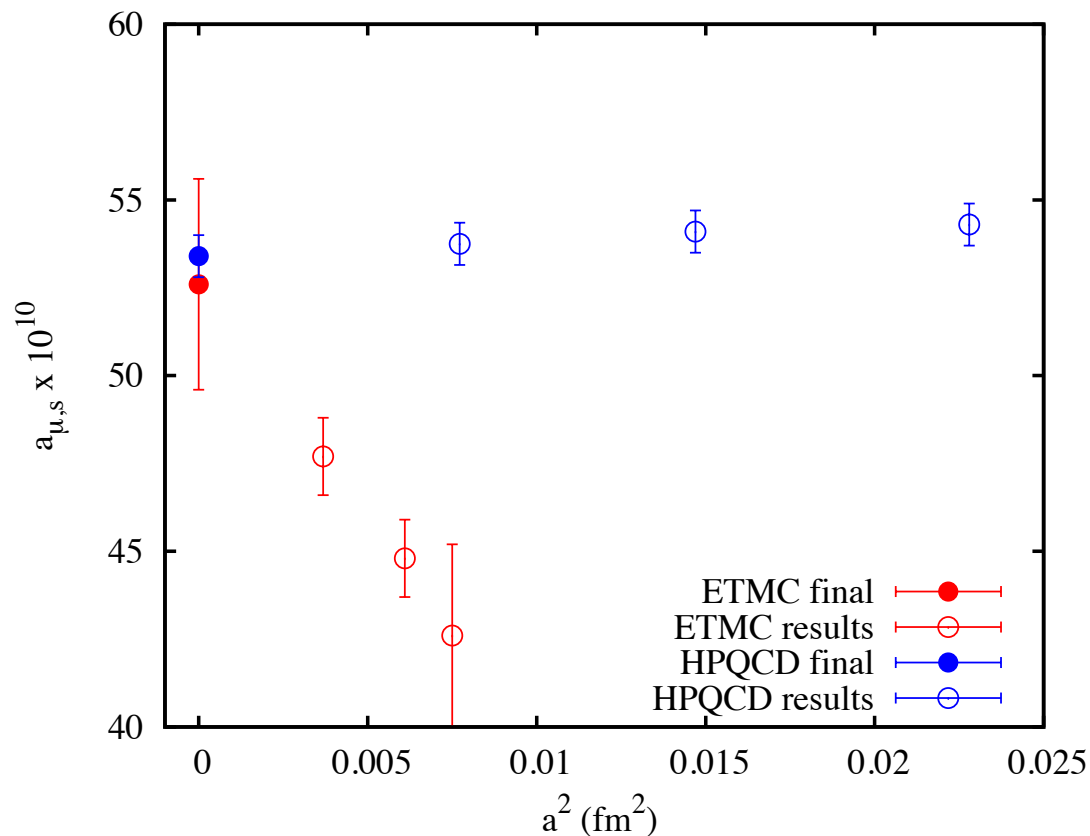
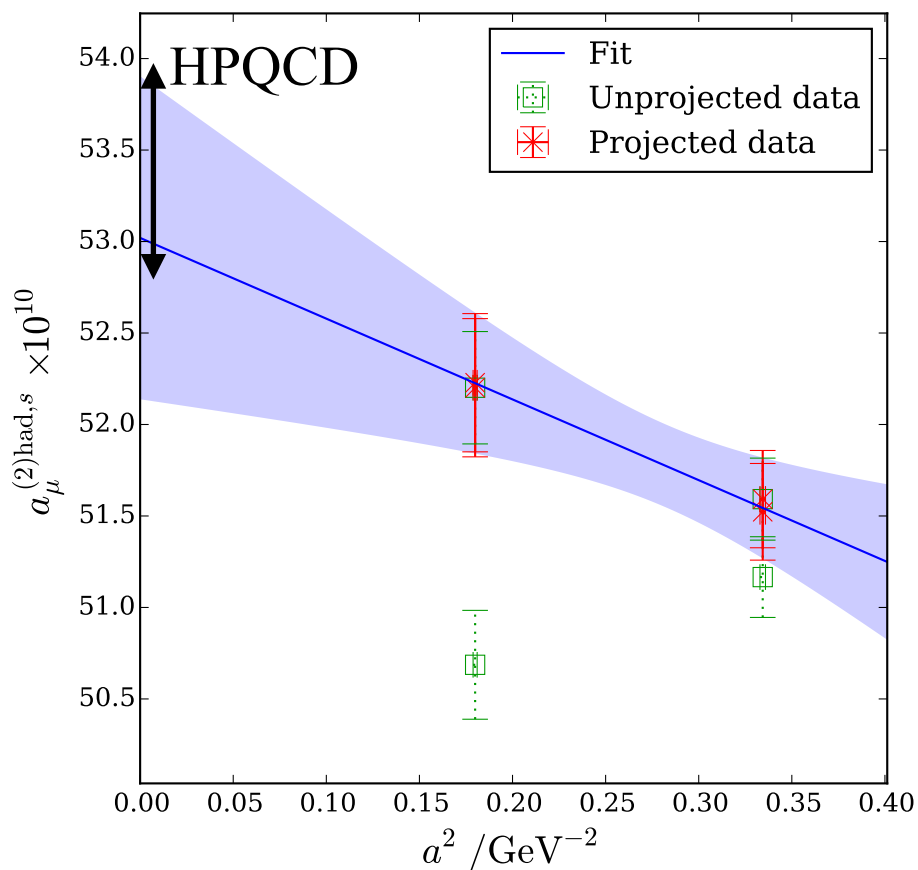
Check mass and decay constant of ϕ from these correlators against expt



Tests large time behaviour of correlator



New results from other formalisms provide good check



RBC/UKQCD domain wall

$$a_{\mu}^{HVP,s} = 53.1(9) \times 10^{-10}$$

1607.01767

ETMC twisted mass

$$a_{\mu}^{HVP,s} = 53(3) \times 10^{-10}$$

From $R_{e^+e^-}$ we estimate $a_{\mu}^{HVP,s} < \approx 55 \times 10^{-10}$

CHARM contribution

HPQCD 1004.4285,
1208.2855

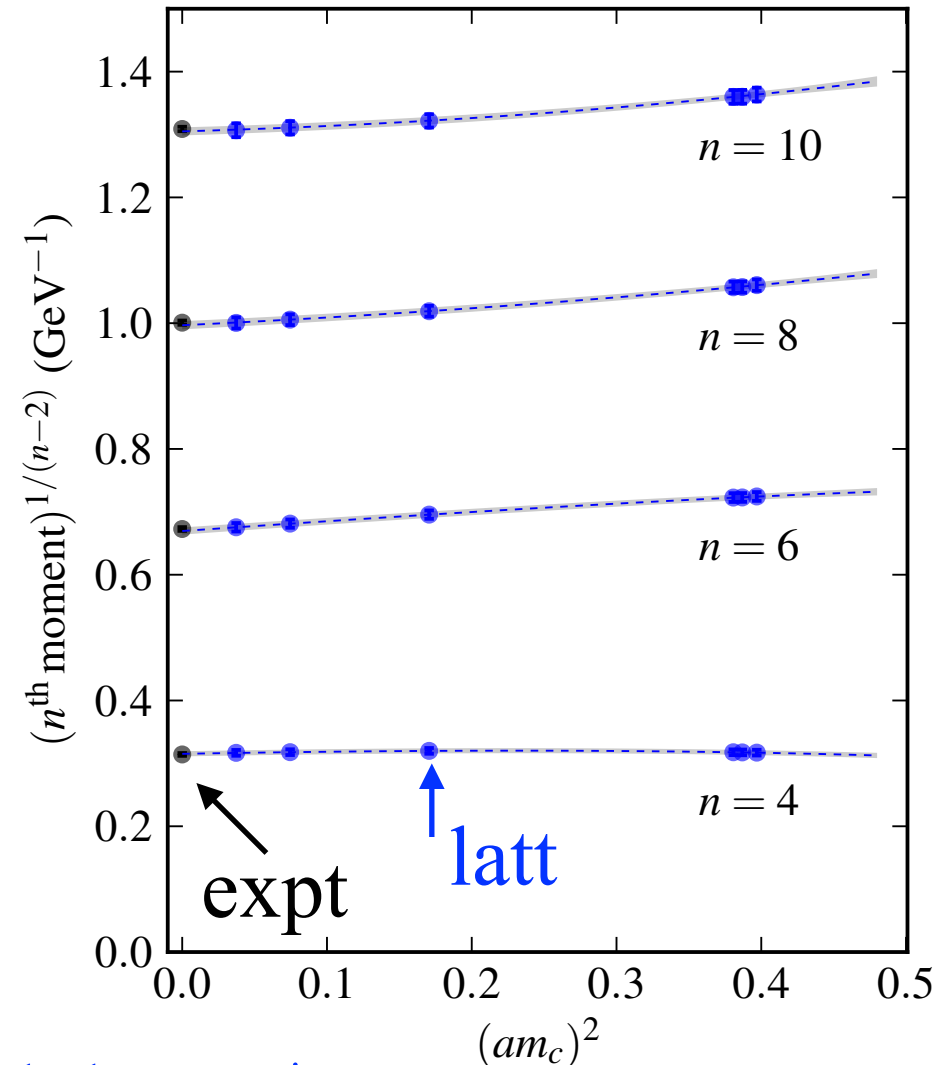
Part of the set of calculations that gave

$$m_c, M(J/\psi) - M(\eta_c), \Gamma(J/\psi \rightarrow e^+e^-), \Gamma(J/\psi \rightarrow \eta_c\gamma)$$

Used HISQ valence quarks on MILC 2+1 asqtad configs. Z_V from contnm QCD pert. th.

Extrapolation to physical point allows us to compare directly to moments from e^+e^- expt. in charm region

$$a_\mu^{HVP,c} = 14.4(4) \times 10^{-10}$$



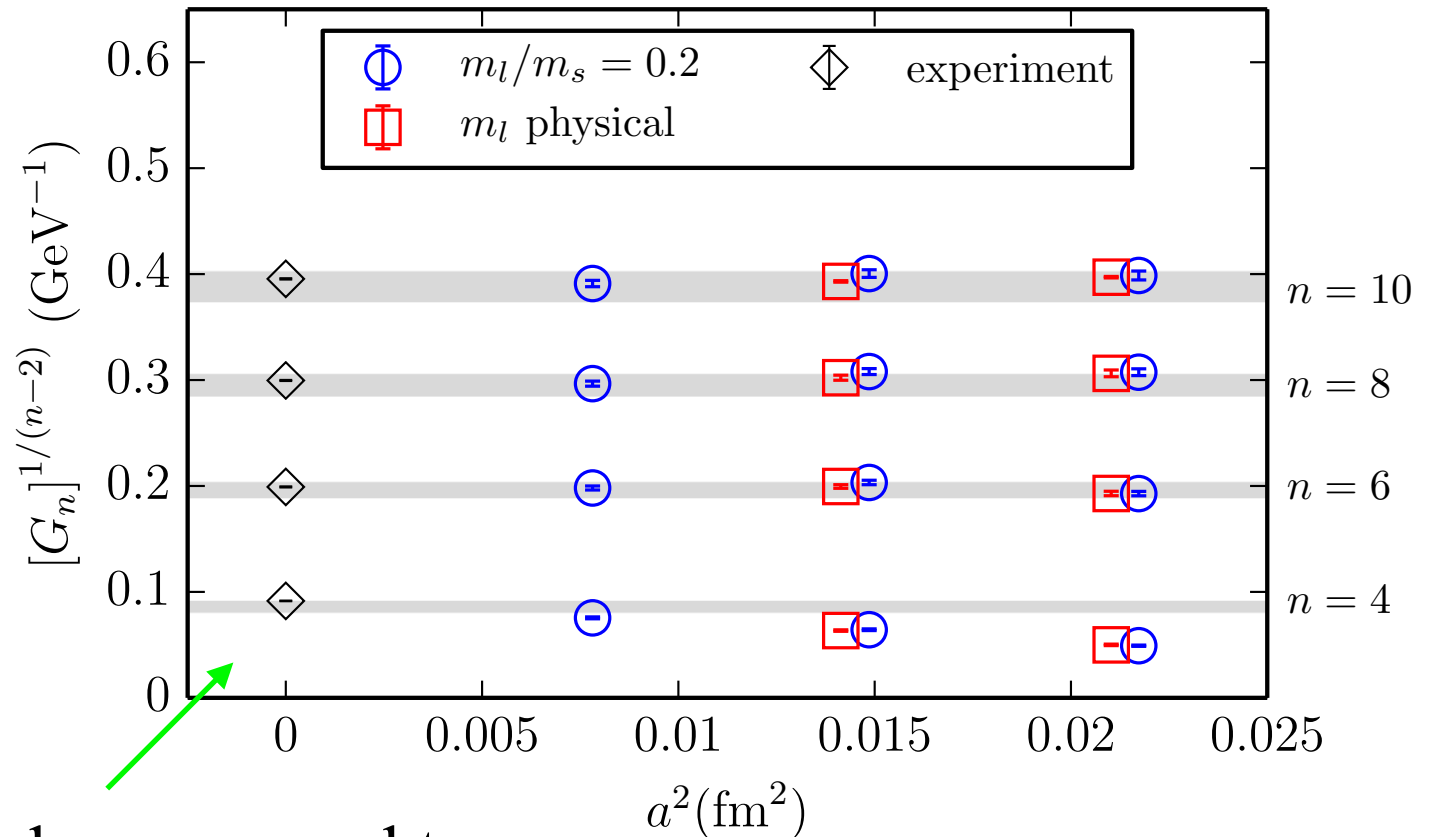
BOTTOM contribution

HPQCD 1110.6887,
1309.5797,
1408.5768

Part of the set of calculations that gave

$$m_b, M(\Upsilon) - M(\eta_b), M(\Upsilon') - M(\eta'_b), \Gamma(\Upsilon \rightarrow e^+e^-), \Gamma(\Upsilon' \rightarrow e^+e^-)$$

Used NRQCD
valence
quarks on
MILC 2+1+1
HISQ configs.
 Z_V from
contnm QCD
pert. th.



Again, moments can be compared to those extracted from expt.

$$a_\mu^{HVP,b} = 0.27(4) \times 10^{-10}$$

HISQ valence quarks on MILC 2+1+1 HISQ configs. Use Z_v from s calc.

Multiple a (use w_0), m_1 (inc. phys.), volumes (at $m_l/m_s=0.1$).

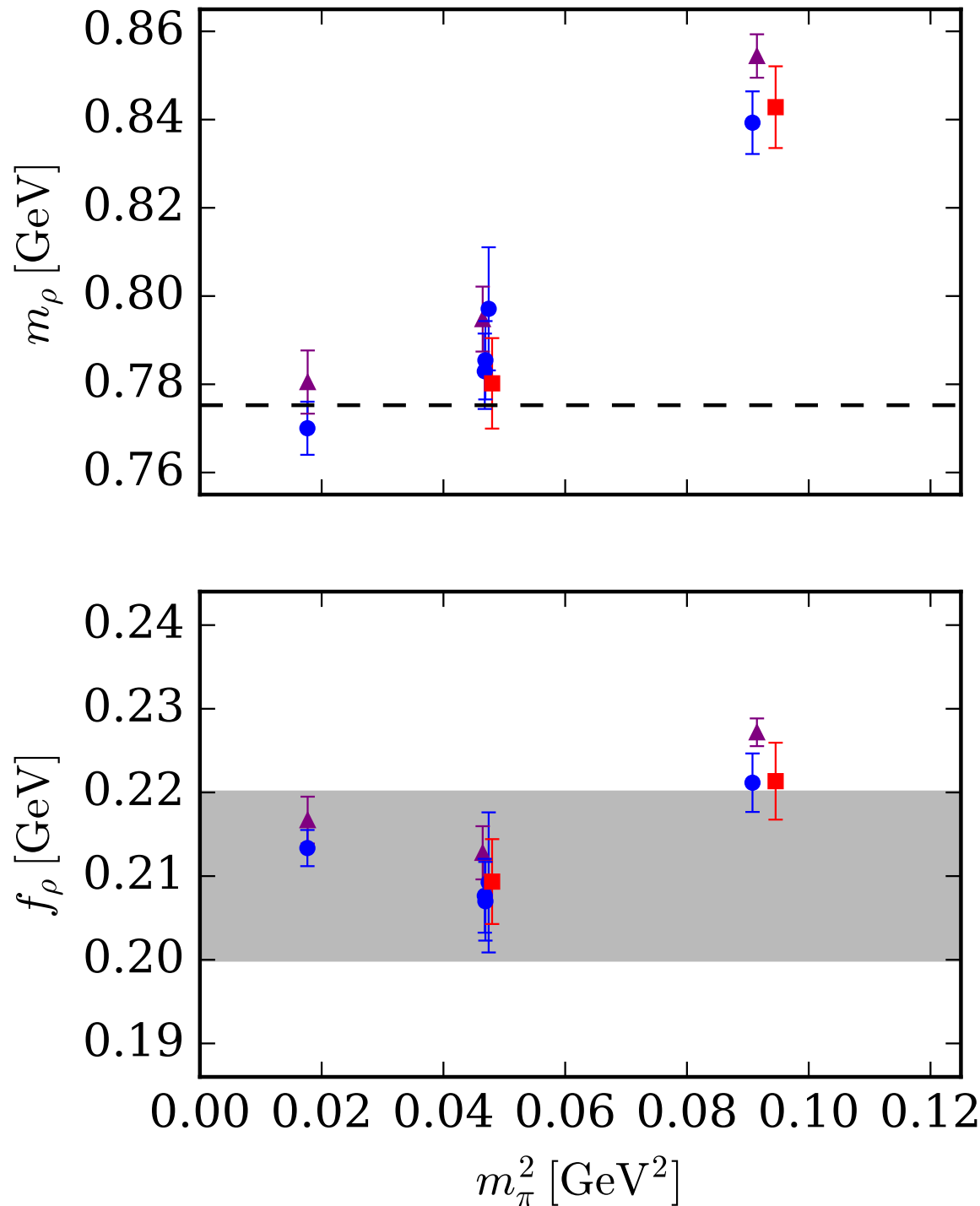
New ingredient since correlators much noisier. Use:

$$G(t) = \begin{cases} G_{\text{data}}(t) & \text{for } t \leq t^* & \longleftarrow \text{from Monte Carlo} \\ G_{\text{fit}}(t) & \text{for } t > t^* & \longleftarrow \text{from multi-exponential fit} \end{cases}$$

$$t^* = 1.5\text{fm} = 6/m_\rho \quad \text{so 70\% of result from } G_{\text{data}}$$

- 80% of result comes from ρ meson pole, so need to understand ρ on lattice
- 10% from $\pi\pi$, sensitive to finite-volume and m_π (so taste-issues for staggered quarks).

Mass and
decay constant of
the ρ from
large time
behaviour of
u/d vector
correlators



Correct key lattice systematics

$$\hat{\Pi}_j^{latt} \rightarrow \left(\hat{\Pi}_j^{latt} - \hat{\Pi}_j^{latt}(\pi\pi) \right) \left[\frac{m_\rho^{2j,latt}}{m_\rho^{2j,expt}} \right] + \hat{\Pi}_j^{cont}(\pi\pi)$$

Remove lattice $\pi\pi$ using effective theory of ρ, π, γ inc. staggered quark effects and finite vol.

Corrections reduce spread of results.

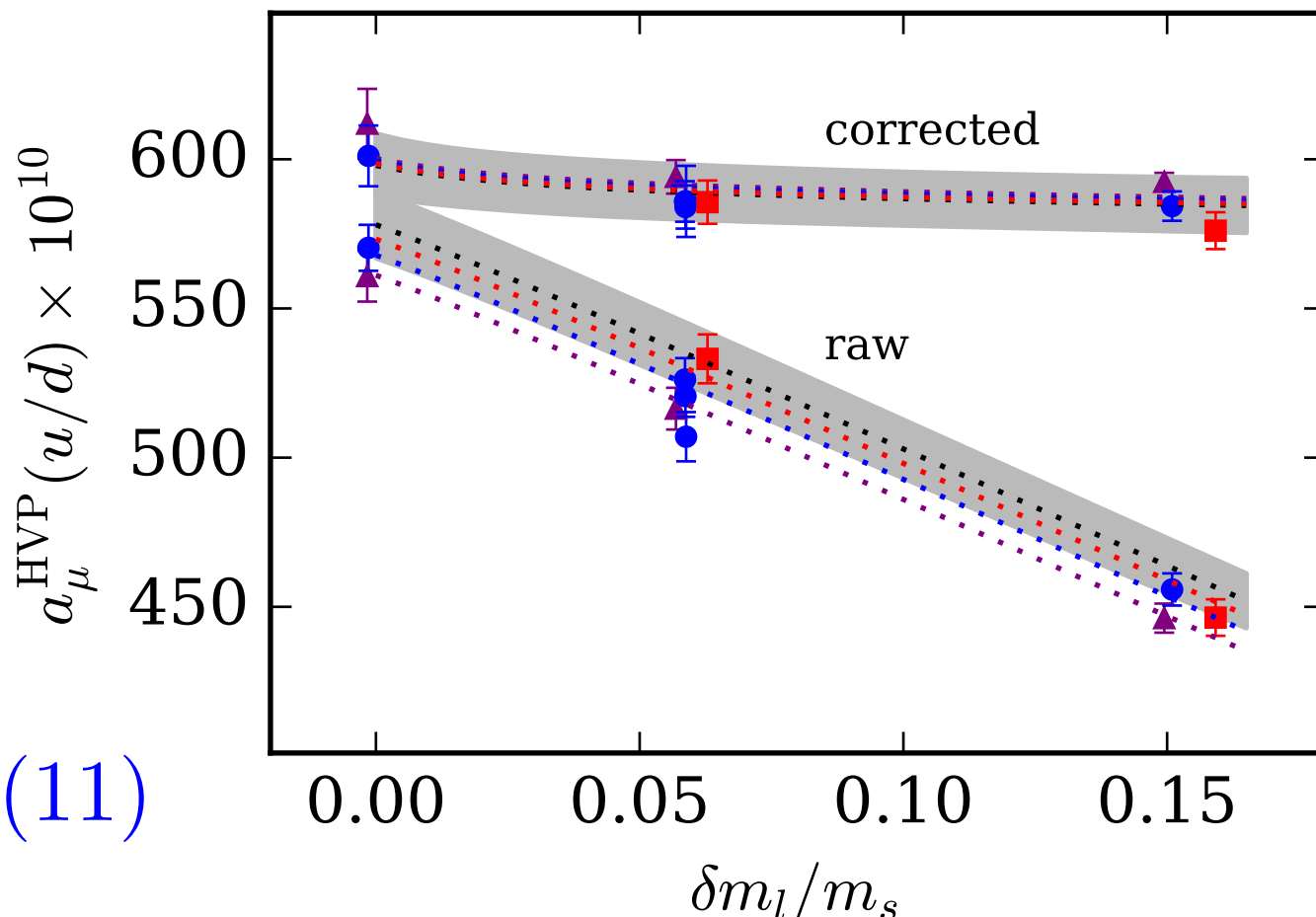
Fit for remaining dependence on a and m_l

$$a_\mu^{HVP,u/d} = 598(11)$$

Rescale using exptl m_ρ to reduce m_l dependence

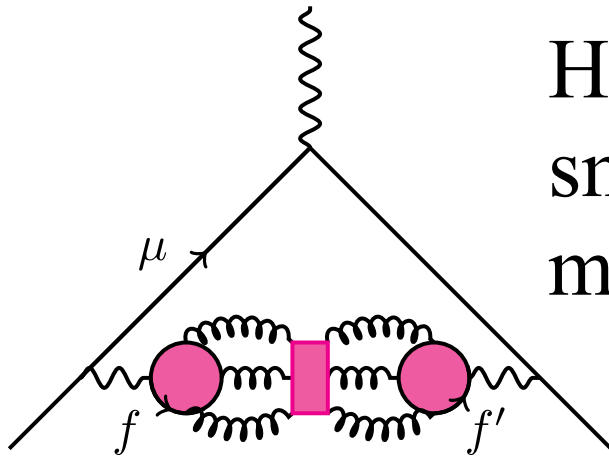
Restore $\pi\pi$ from continuum

HPQCD 1601.03071



Quark-line disconnected contribution

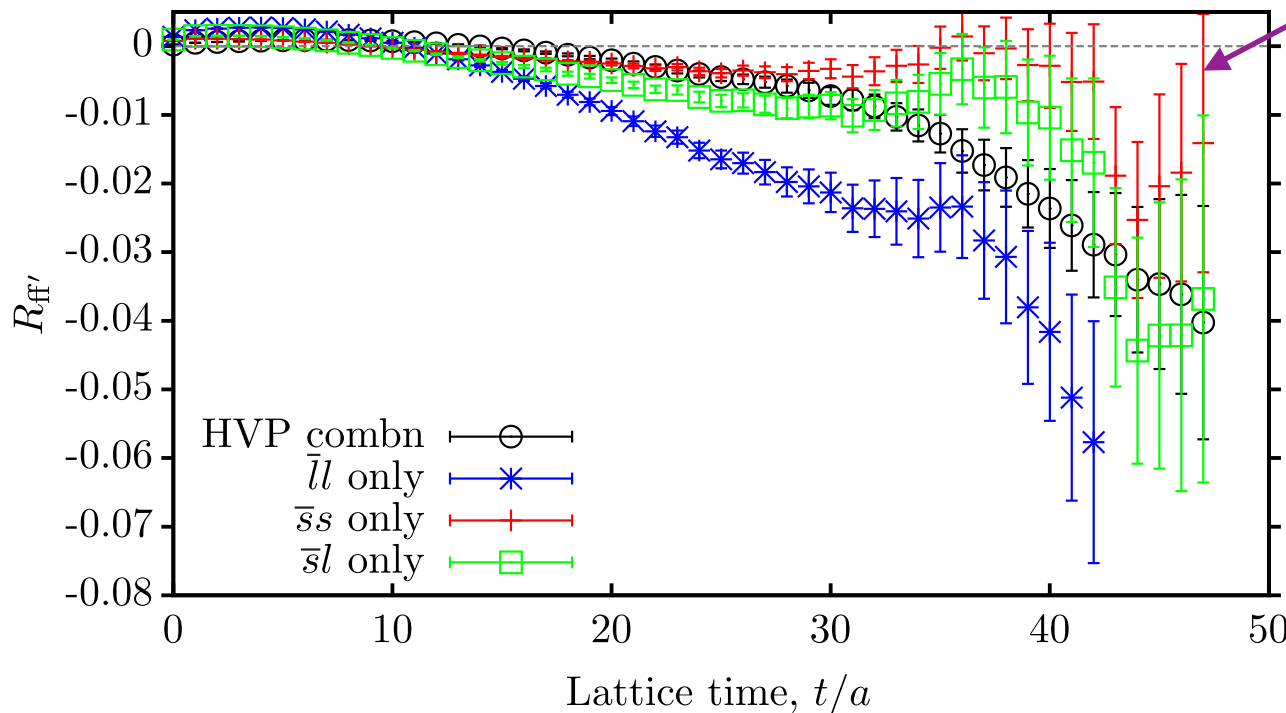
HPQCD/Hadspec 1512.03270



Hard to calculate but small. Suppressed by masses since

$$\sum_{u,d,s} Q_f = 0$$

Simple estimates give ratio of disc. to conn contribution of $-1(1)\%$

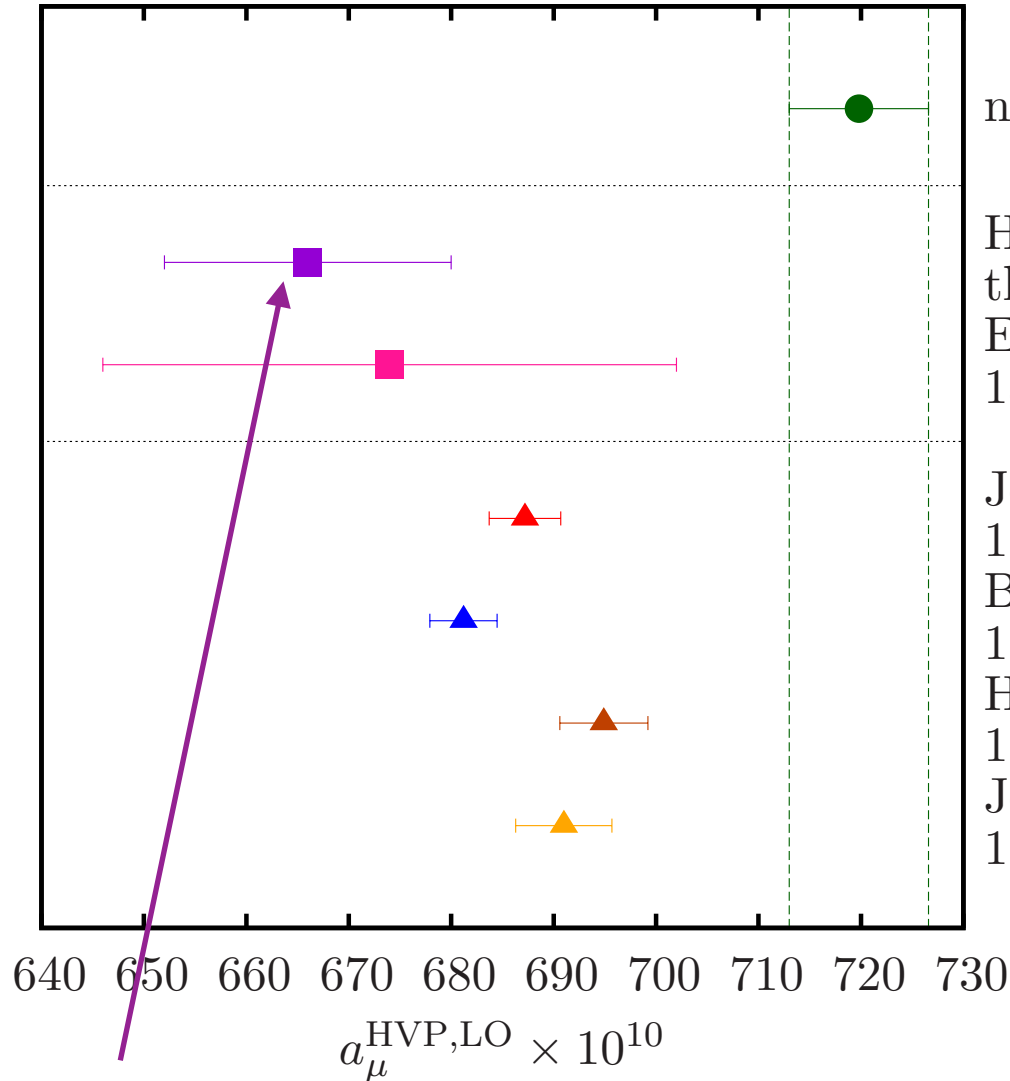


Hadspec ratio of disc. to conn. correlators small and contribution further suppressed (by factor 5 by quark charges)
 Estimate (after fitting):

see also RBC/UKQCD 1512.09054:
 $-9.6(4.0) \times 10^{-10}$

$$a_{\mu}^{HVP,disc} = 0(9) \times 10^{-10}$$

Conclusion: Combining numbers for a total



$$a_{\mu}^{\text{HVP,LO}} \times 10^{-10}$$

$$598(11) \text{ } u/d$$

$$53.4(6) \text{ } s$$

$$14.4(4) \text{ } c$$

$$0.27(4) \text{ } b$$

$$\text{Total } 666(6)(12)$$

add syst from
disc. diags
(1.5%) in quad

3.5σ discrepancy with no new physics

Future focus (with MILC) : improve physical point results and reduce systs from QED, m_u/m_d and disc.

Backup Slides

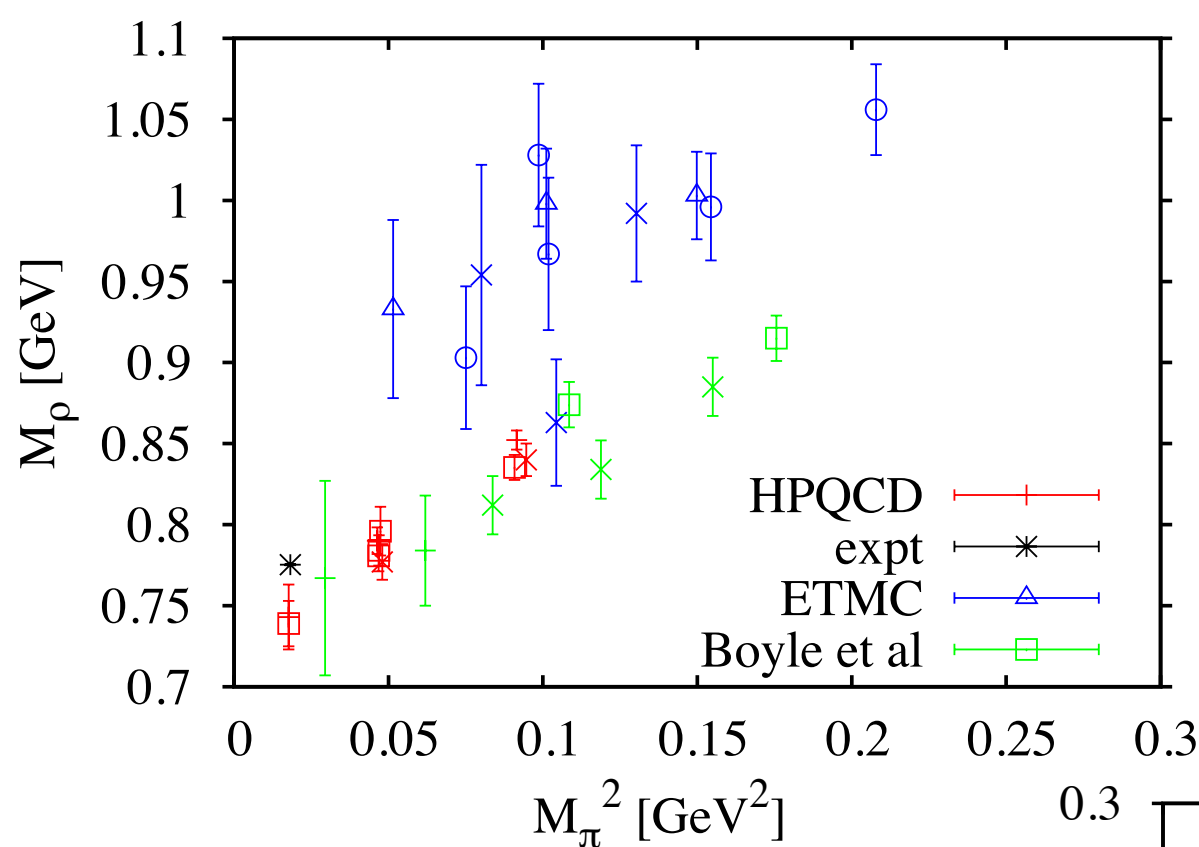
Error budget for u/d HVP

TABLE III: Error budget for the connected contributions to the muon anomaly a_μ from vacuum polarization of u/d quarks.

	$a_\mu^{\text{HVP,LO}}(u/d)$
QED corrections:	1.0 %
Isospin breaking corrections:	1.0 %
Staggered pions, finite volume:	0.7 %
Noise reduction (t^*):	0.5 %
Valence m_ℓ extrapolation:	0.4 %
Monte Carlo statistics:	0.4 %
Padé approximants:	0.4 %
$a^2 \rightarrow 0$ extrapolation:	0.3 %
Z_V uncertainty:	0.4 %
Correlator fits:	0.2 %
Tuning sea-quark masses:	0.2 %
Lattice spacing uncertainty:	< 0.05 %
Total:	1.9 %

Analysis of ρ parameters

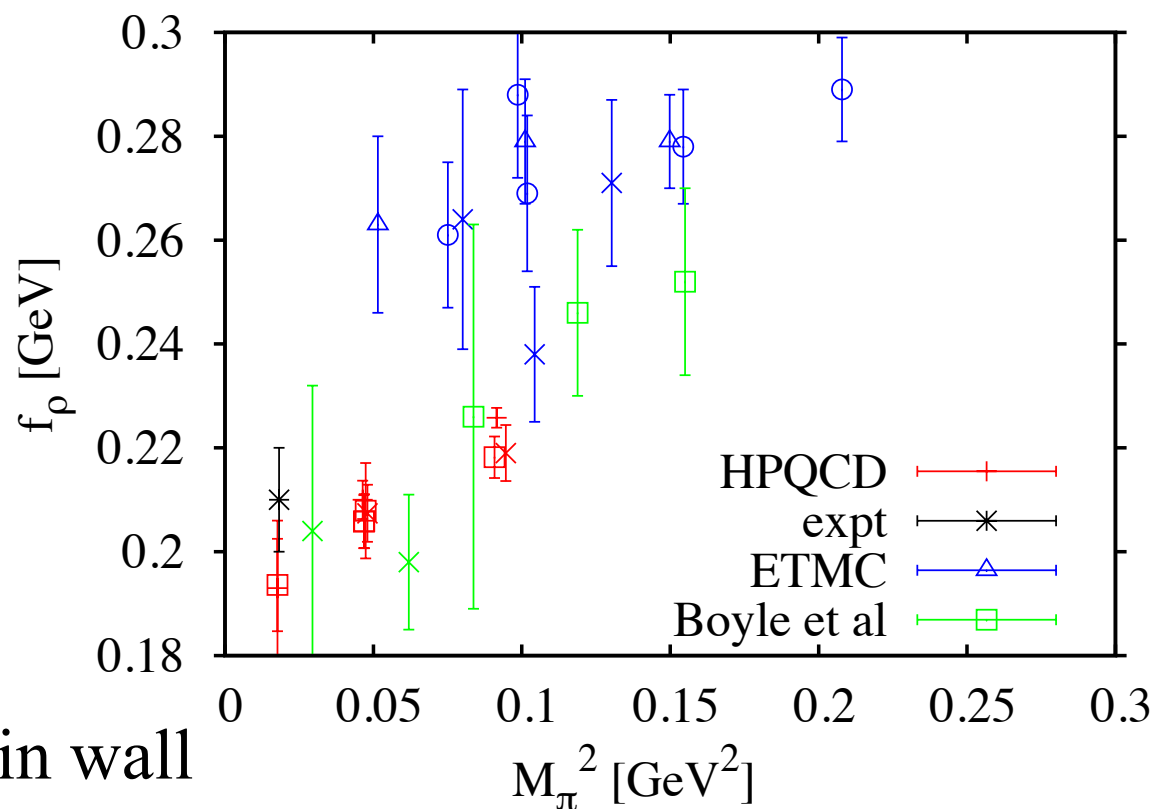
Direct comparison with ETMC (1308.4327) and Boyle et al (1107.1497) possible



ETMC a 0.06-0.08fm
L 2.5- 2.9fm

HPQCD a 0.09-0.15fm
L 2.5-5.8fm

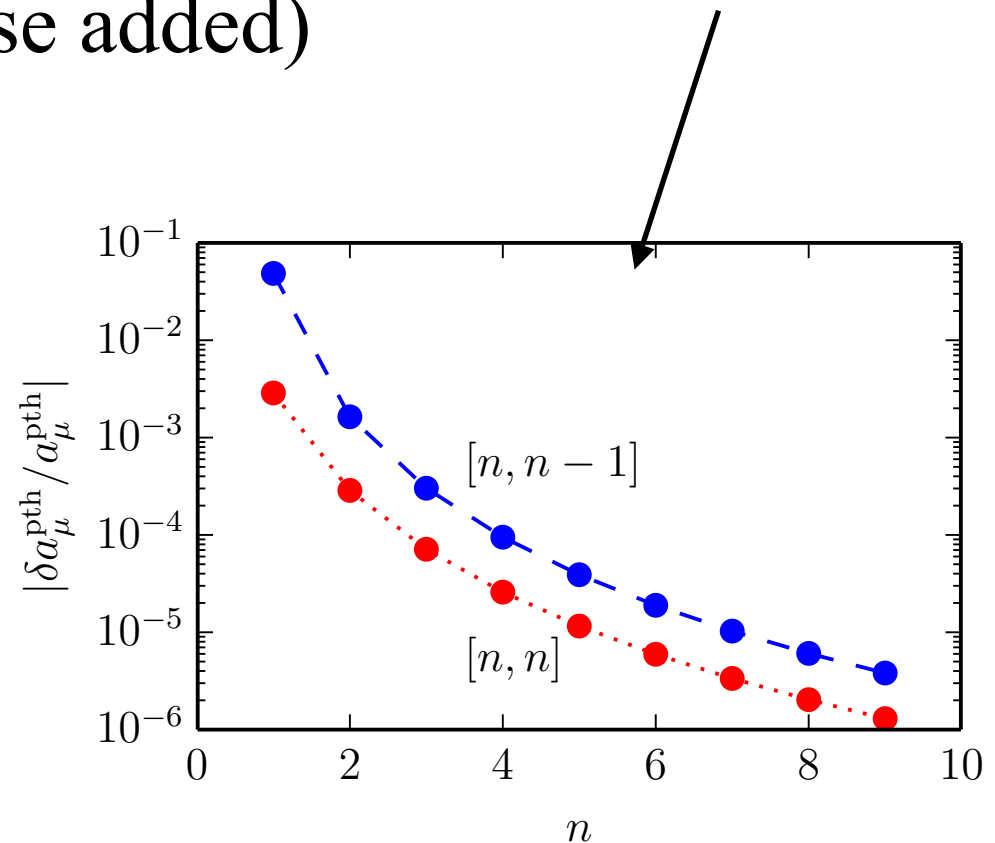
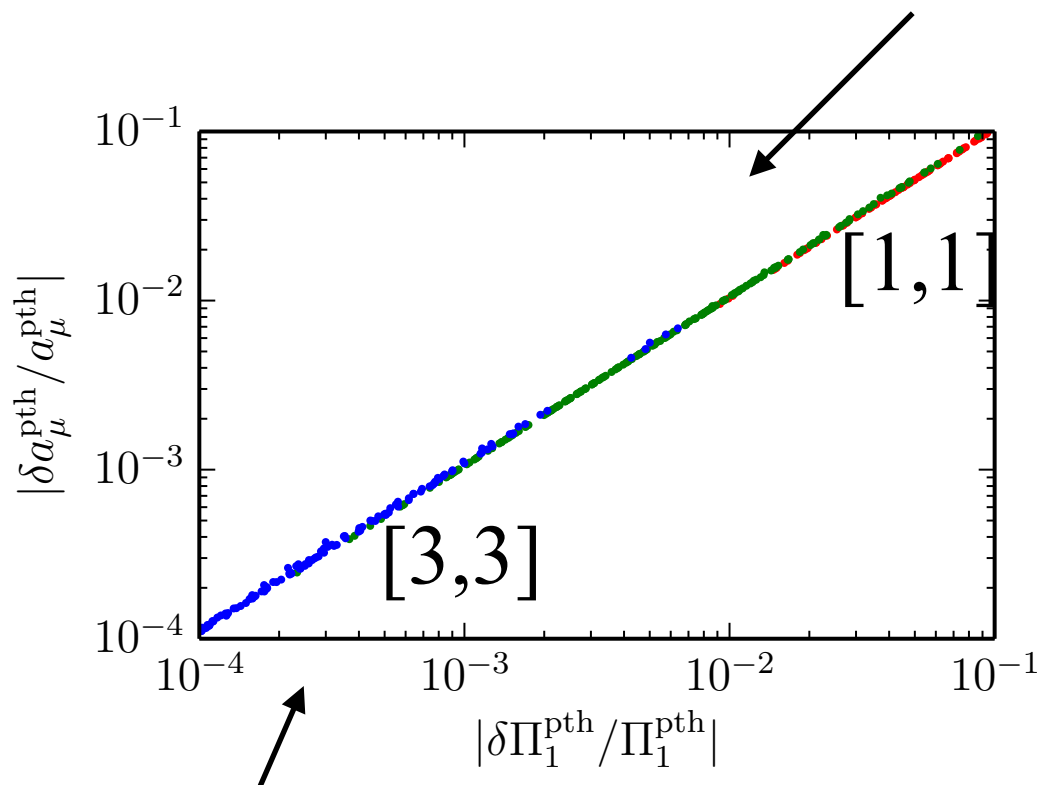
Boyle et al a 0.09-0.14fm
L 2.7-4.6fm ← domain wall



Allows us to reconstruct $\hat{\Pi}(q^2)$ and integrate

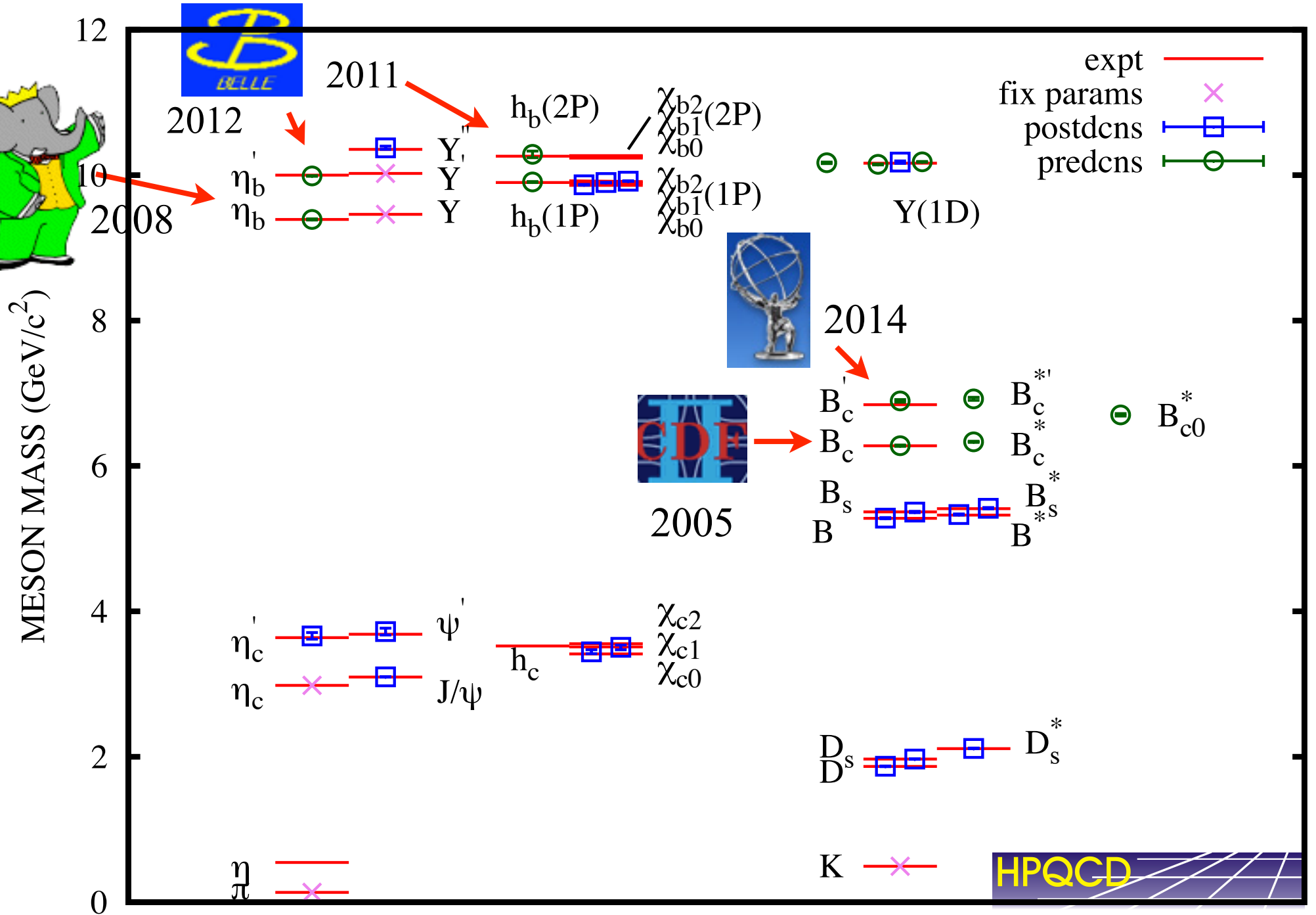
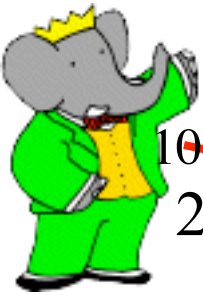
Use Pade approximants (ratio of m/n polynomials) rather than Taylor expansion for better large q^2 behaviour.

Test Pade approximants in similar scenarios (1-loop quark vacuum polarisation, with noise added)



Improved precision allows higher order Pade - we use $[2,2]$

Keep an eye on the 'big' picture whilst doing this



few MeV uncertainties in many cases

Keep an eye on the ‘big’ picture whilst doing this

