





GEANT4 TOOLS FOR THE DESIGN AND ANALYSIS OF PHOTO-FISSION EXPERIMENTS AT ELI-NP

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1. ELI-NP facility

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2. Experimental setup optimization using GEANT4

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3. GEANT4 tool for determining the detector response function

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4. Conclusions & feedback in using GEANT4

1. Extreme Light Infrastructure – Nuclear Physics (ELI-NP)





User facility for nuclear physics experiments

Developed in Romania (near Bucharest)

Co-financed by the Romanian Government

and the European Union

Operational phase starting 2020





High Power Laser System (HPLS)

- 2 arms, 6 outputs
- $\ 2 \ x \ 0.1$ PW, 10 Hz
- 2 x 1 PW, 1 Hz
- 2 x 10 PW, 0.1 Hz (world record!)





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Gamma Beam System (GBS)

- spectral density 0.8-4 \cdot 10⁴ γ /(s \cdot eV)
- narrow bandwidth 0.3-0.5%
- energy range 0.2-19.5 MeV
- linear polarization >99%





Goal: Production and study of exotic **neutron-rich** nuclei

Production method: Photo-fission of actinides targets





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- **thick** because $\sigma(\gamma, f)$ small
- **sliced** in many thin foils: refractory, fast extraction
- tilted: (1) avoid hitting neighboring foils(2) increase γ pathlength





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GEANT4 simulations for IGISOL ion release yield maximization









Optimization of correlated parameters:











One of the important aspects of data analysis: **Unfolding the spectrum from the detector response function**



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3. GEANT4 – tool for determining the detector response function

²³³U+n_{th} experiment

Location: Hungarian Academy of Sciences, Centre for Energy Research (MTA EK – Budapest) – Feb. 2018





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Motivation: Fission fragment deexcitation study, nuclear data, etc.





GEANT4 – simulation of the detector response



Experimental setup construction & detector response simulation



GEANT4 – simulation of the detector response









Detector response matrix construction





GEANT4 – simulation of the detector response







Using matrix inversion method for unfolding

- GEANT4 simulations to determine the response function of the detector in the given setup
- Construct the response matrix
- Using the inverse response matrix to unfold the spectrum from the detector response function

GEANT4 – User feedback



GEANT4 is an important tool in nuclear physics for simulation and design of experiments but also for data analysis.

Necessary additions to GEANT4 implemented by me:

- Photo-fission cross section and data model in 5-25 MeV range (cross section function on energy and compute of secondaries); *P. Constantin et al., NIMB 372 (2016) B. Mei et al., Phys. Rev. C 96 (2017)*

Schiwietz charge state parameterisation for energy loss (default Ziegler):
Schiwietz performs better than Ziegler for slow (E ~ 1 MeV/u) & heavy ions (M > 80 GeV)
P. Constantin et al., NIMB 372 (2016) G. Schiwietz et al., Nucl. Instr. Meth. B 175–177 (2001)

- Personal wish: Have more realistic examples for constructing common detector types & simulating response functions taking into account effects as pile-up and dead time

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ADDITIONAL

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Measurements with the LOHENGRIN spectrometer at ILL Grenoble of ionic charge state distributions in thermal neutron induced fission of ²³⁵U [18] found $\langle Q \rangle = 20-22$ and $\sigma_Q = 2.0-2.4$. Similar values were measured in ²³⁹Pu($n_{th}f$) [19], ²⁴¹Am($2n_{th}f$) [20] and ²⁴¹Pu($n_{th}f$) [21]. These values refer to the regular ionic charge state distributions, unaffected by the presence of nanosecond isomeric states. When such isomers are present, they decay by conversion with electrons from internal shells and the produced vacancies lead to Auger cascades and, consequently, to ionic charge states higher by $\Delta \langle Q \rangle \approx 3$ [18].

Fig. 6 shows the dependence of the ionic charge Q, at the moment of release from the target, on the ion velocity v for five values of the nuclear charge Z = 36, 38, 40, 56, 58. The lower set of five tightly bound curves, drawn with open circles, is produced with the Ziegler parameterization, while the upper set of five loosely bound curves, drawn with full squares, is produced with the Schiwietz parameterization. Apart from the overall higher Q values, the Schiwietz parameterization has also a more pronounced dependence on the nuclear charge Z than the Ziegler







GIF U238 implementation in GEANT4

Reimplemented two classes

GamaFissionCrossSection (for TCS) & GammaFissionDataModel (for A,Z)

The implemented algorithm is calculating the fission fragments yields in 3 steps:

- 1. The calculation of the Total Cross Section of the Photo-Fission reaction (depending on the incident photon energy)
- 2. After the photo-fission reaction is decided to take place at the current event the mass (A) is calculated according to the mass distribution

3. For the current mass A decided at step 2, the charge (Z) number is being decided according with the charge distribution

Group acknowledgements







$$\begin{split} Y(A) &= Y_{\rm SM}(A) + Y_{\rm ASMI}(A) + Y_{\rm ASMII}(A) \\ &= C_{\rm SM} \exp\left[-\frac{(A - A_{\rm SM})^2}{2\sigma_{\rm SM}^2}\right] \\ &+ C_{\rm ASMI} \exp\left[-\frac{(A - A_{\rm SM} - D_{\rm ASMI})^2}{2\sigma_{\rm ASMI}^2}\right] \\ &+ C_{\rm ASMI} \exp\left[-\frac{(A - A_{\rm SM} + D_{\rm ASMI})^2}{2\sigma_{\rm ASMI}^2}\right] \\ &+ C_{\rm ASMII} \exp\left[-\frac{(A - A_{\rm SM} - D_{\rm ASMII})^2}{2\sigma_{\rm ASMII}^2}\right] \\ &+ C_{\rm ASMII} \exp\left[-\frac{(A - A_{\rm SM} - D_{\rm ASMII})^2}{2\sigma_{\rm ASMII}^2}\right] \\ &+ C_{\rm ASMII} \exp\left[-\frac{(A - A_{\rm SM} + D_{\rm ASMII})^2}{2\sigma_{\rm ASMII}^2}\right] \end{split}$$

The Gaussian parameters $C_{\rm SM}$, $C_{\rm ASMI}$, $C_{\rm ASMII}$ and $\sigma_{\rm SM}$, $\sigma_{\rm ASMI}$, $\sigma_{\rm ASMII}$ are the amplitudes and widths, respectively, of the symmetric (SM) and two asymmetric (ASMI and ASMII) fission modes. $A_{\rm SM}$ is the most probable mass number for the symmetric fission mode. $A_{\rm SM} - D_{\rm ASMI}$ and $A_{\rm SM} + D_{\rm ASMI}$ ($A_{\rm SM} - D_{\rm ASMII}$ and $A_{\rm SM} + D_{\rm ASMII}$) are the most probable mass numbers of the light and heavy fragments, respectively, for the asymmetric mode ASMI (ASMII). These Gaussian parameters can be determined by fitting

Group acknowledgements



$$Z_s = \frac{A + [(a_c A^{2/3})/(2x)]}{(4a_{\rm sym}/x) + [(a_c A^{2/3})/x]},$$

where $x = 2a_{\text{sym}} + [(m_n - m_p)/2]$, and m_n and m_p are the masses of proton and neutron, respectively [27]. The constants $a_c = 0.71$ MeV and $a_{\text{sym}} = 23.21$ MeV are the Coulomb and symmetry energy coefficients, respectively

> $Z_{\text{prob}} = Z_s - 3.8 + \Delta_Z,$ $\Delta_Z = d_1 + d_2A + d_3A^2 + d_4A^3,$

for light fragments with A < 134,

 $\Delta_Z = d_5 + d_6 \exp(A/d_7),$

for heavy fragments with $A \ge 134$.

Group acknowledgements



Space charge = He⁺ cloud created by fragment (>90%) and e^+/e^- (<10%) induced ionization of He gas Above a certain **charge density rate Q**: field saturation, strong e-ion recombination, weak plasma.



Geant4: He, T=70K, p=300mbar (ρ =0.206mg/cm³) >95% of fragments stop in 11.3cm \rightarrow width~24cm





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