

Cross sections of light nuclei in the Glauber-Gribov representation

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Abstract

The Glauber-Gribov representation of nucleus-nucleus cross sections is reviewed for light nuclei. The model provides fast and robust calculation of total, inelastic, elastic and production nucleus-nucleus cross sections. Comparisons with experimental data are presented in broad range of nucleus energies and weights.

1 Glauber-Gribov model for nucleus-nucleus cross-sections

The Glauber-Gribov model with Gaussian distributed of point-like nucleons in a nucleus results in the following nucleus-nucleus cross-sections:

$$\sigma_{tot}^{A_p A_t} = 2\pi(R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{tot}^{NN}}{2\pi(R_p^2 + R_t^2)} \right],$$

$$\sigma_{in}^{A_p A_t} = \pi(R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{tot}^{NN}}{\pi(R_p^2 + R_t^2)} \right], \quad \sigma_{prod}^{A_p A_t} = \pi(R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi(R_p^2 + R_t^2)} \right],$$

$$\sigma_{el}^{A_p A_t} = \sigma_{tot}^{A_p A_t} - \sigma_{in}^{A_p A_t}, \quad \sigma_{qel}^{A_p A_t} = \sigma_{in}^{A_p A_t} - \sigma_{prod}^{A_p A_t},$$

Where $\sigma_{tot}^{A_p A_t}$, $\sigma_{in}^{A_p A_t}$, $\sigma_{prod}^{A_p A_t}$, $\sigma_{el}^{A_p A_t}$, and $\sigma_{qel}^{A_p A_t}$, are the total, inelastic (reaction), production, elastic and quasi-elastic cross sections, respectively. The projectile and the target nucleus weights are, $A_p = Z_p + N_p$, and, $A_t = Z_t + N_t$, respectively. The values of Z and N are the corresponding numbers of proton and neutrons in the nuclei, $\sigma_{tot/in}^{NN}$ are the nucleon-nucleon cross-sections.

$$A_p A_t \sigma_{tot}^{NN} = Z_p Z_t \sigma_{tot}^{pp} + N_p N_t \sigma_{tot}^{nn} + (Z_p N_t + N_p Z_t) \sigma_{tot}^{pn}.$$

2 The model parameters

The model is reduced to the selection of $\sigma_{tot/in}^{NN}$ and $R(A)$ values.

The nucleon-nucleon cross-sections are well tabulated and can be found in literature (PDG) or in GEANT4 class [G4HadronNucleonXsc](#).

$R(A)$ is the RMS radius of nucleon distribution inside the nucleus.

For low energies the inelastic and the total cross-sections are corrected for Coulomb barrier:

$$\sigma_{tot}^{A_p A_t} \rightarrow \sigma_{tot}^{A_p A_t} \left[1 - \frac{B_c}{T_{kin}^{cm}} \right], \quad B_c = \frac{\alpha \hbar c Z_p Z_t}{R_{min}},$$

where $R_{min} \sim R_t + R_p$, ($\alpha = e^2/\hbar c$).

3 The light nucleus radii

The RMS radius of nucleon distribution inside a nucleus, $R(A)$ in the medium range of $A \leq 50$ is fitted in GEANT4 by the following expression:

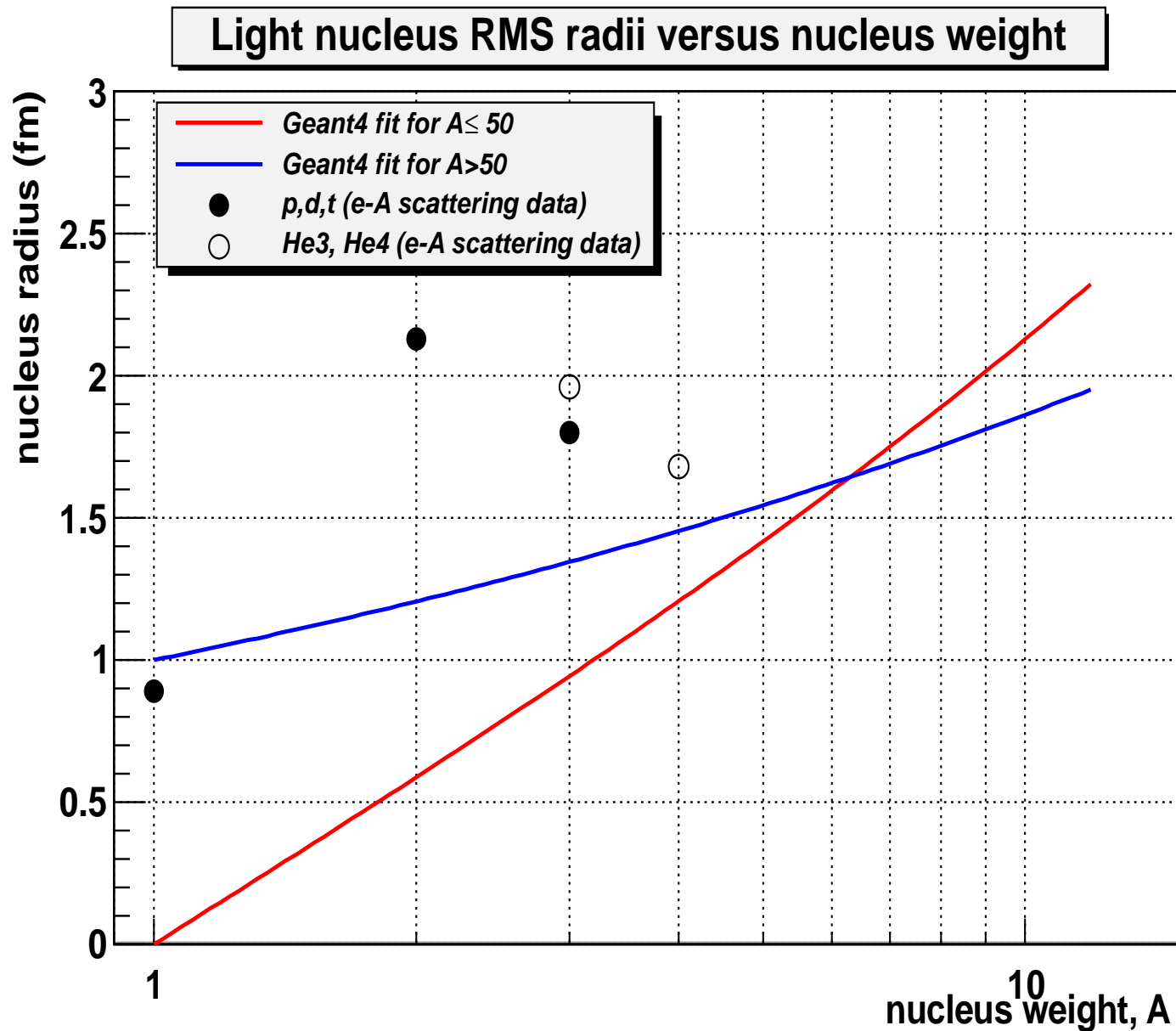
$$R \sim r_o(A)A^{1/3}(1 - A^{-2/3}),$$

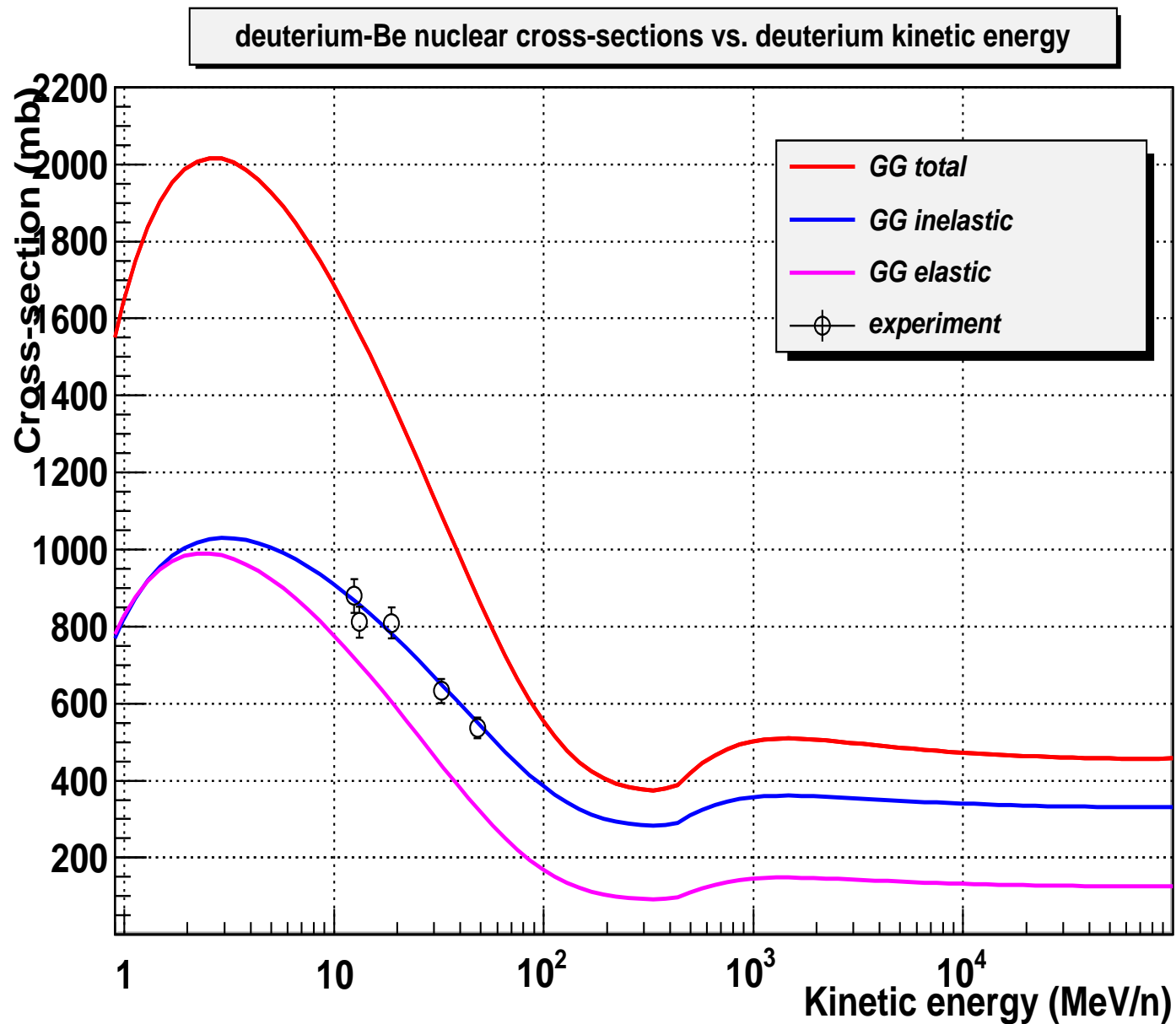
where $r_o(A) \sim 1.1-1.26$ fm is fitted by experimental data. For heavy nuclei $A > 50$ the GEANT4 parametrization reads:

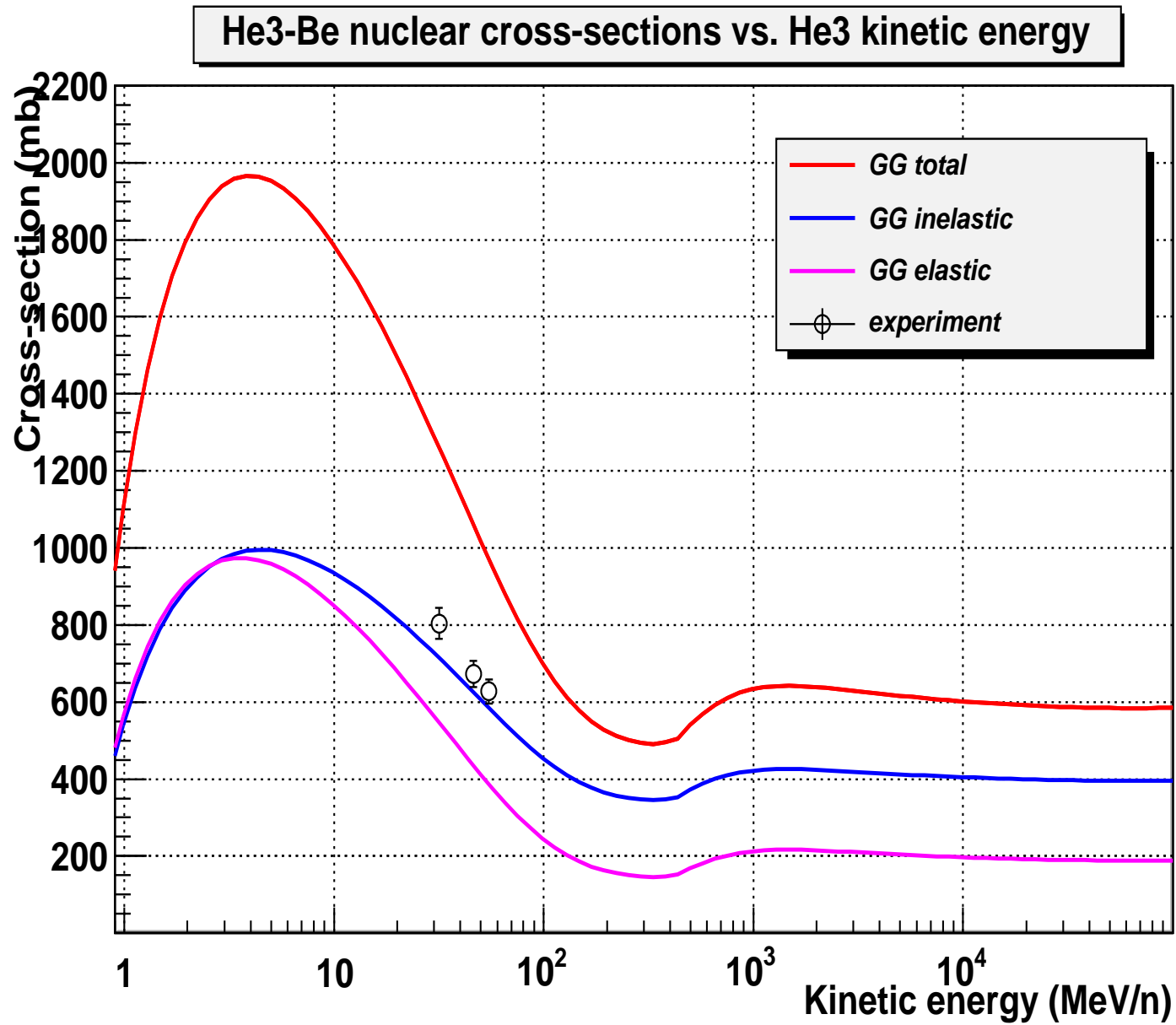
$$R \sim r_o(A)A^{0.27},$$

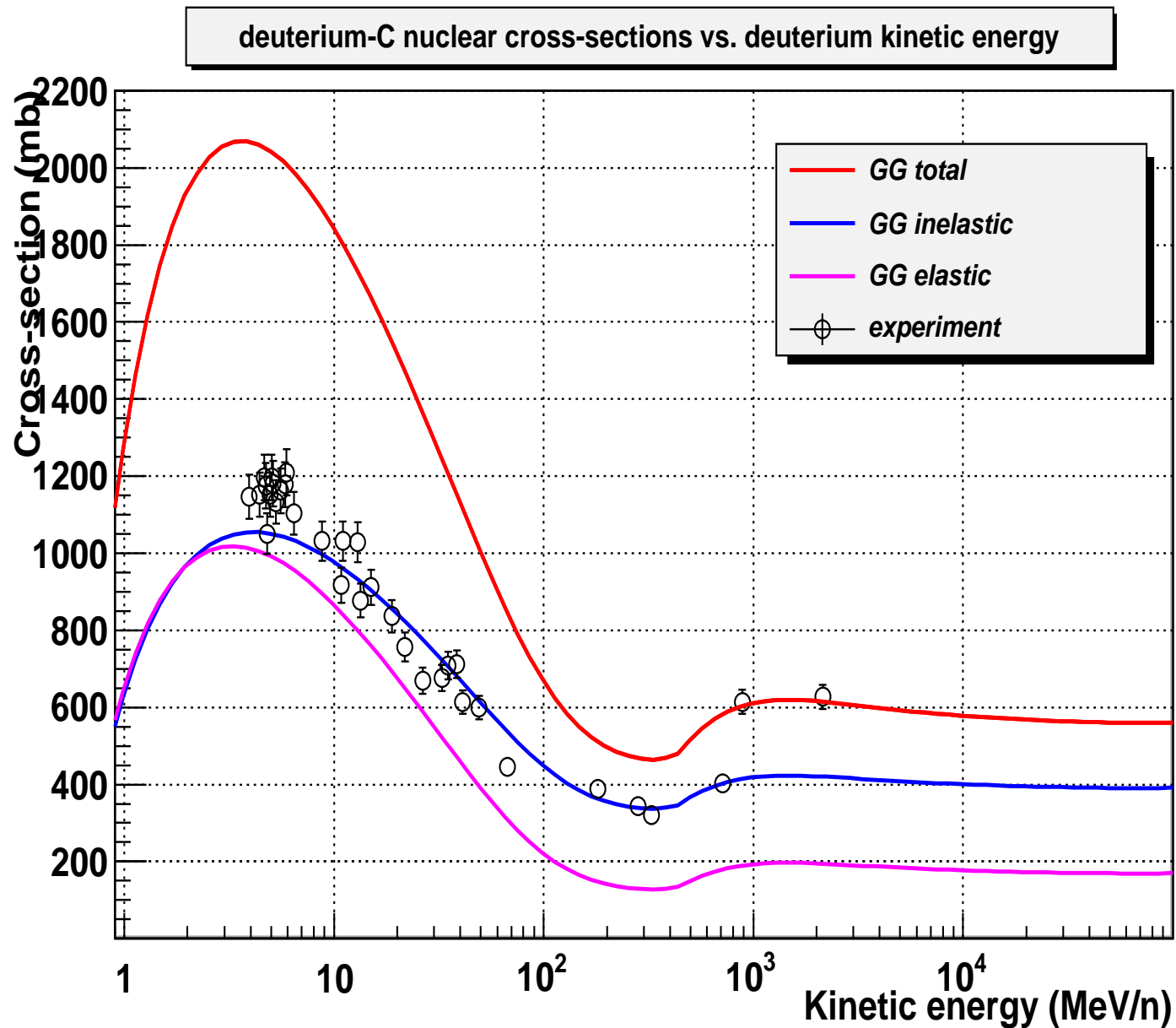
where $r_o(A) \sim 1$ fm is fitted by experimental data.

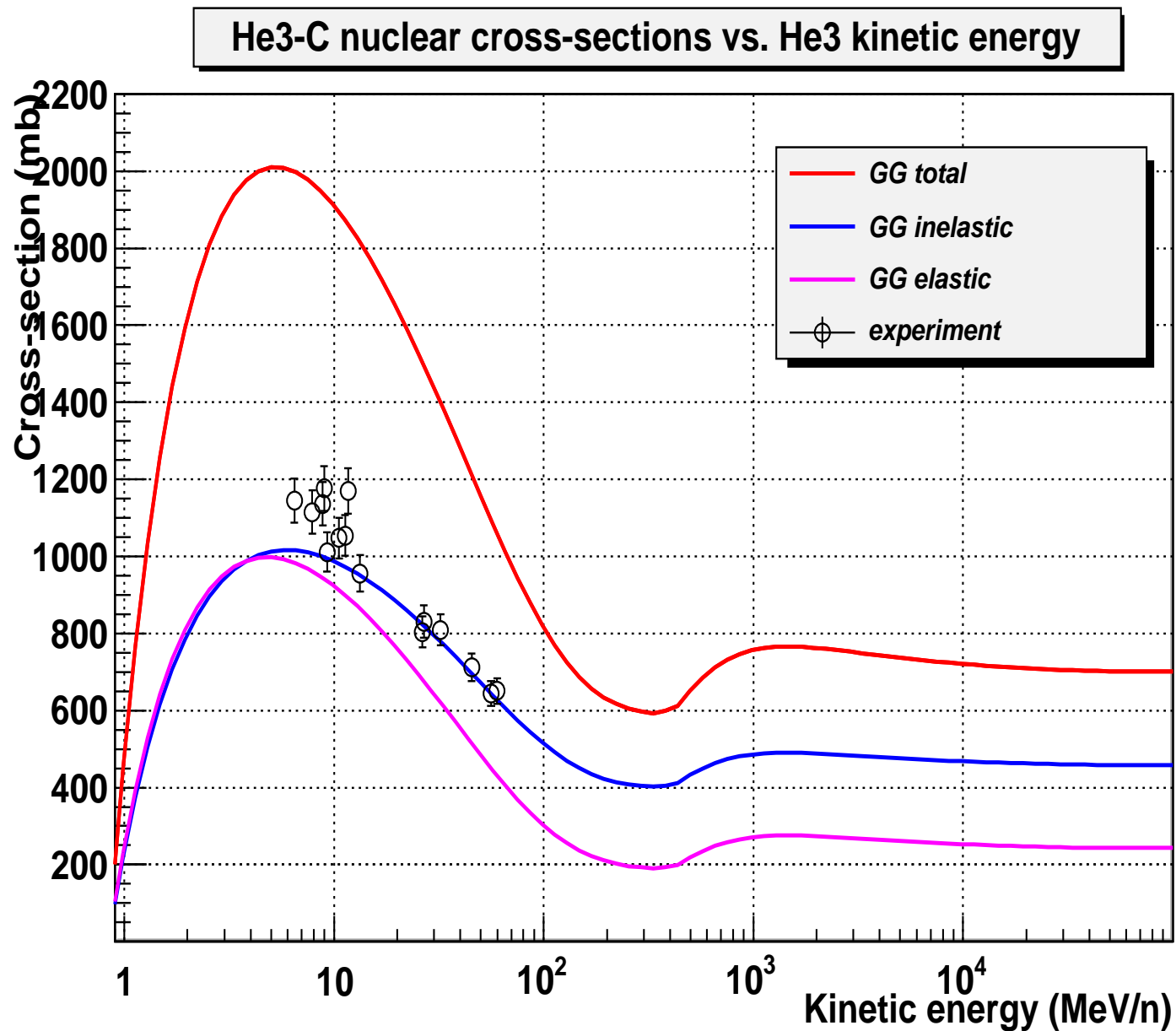
For light nuclei, however, these expressions do not describe experimental data (a compilation of different measurements is shown). The light nucleus-nucleus cross-section experimental data are in agreement with RMS radii extracted directly from electron-nucleus scattering experiments.

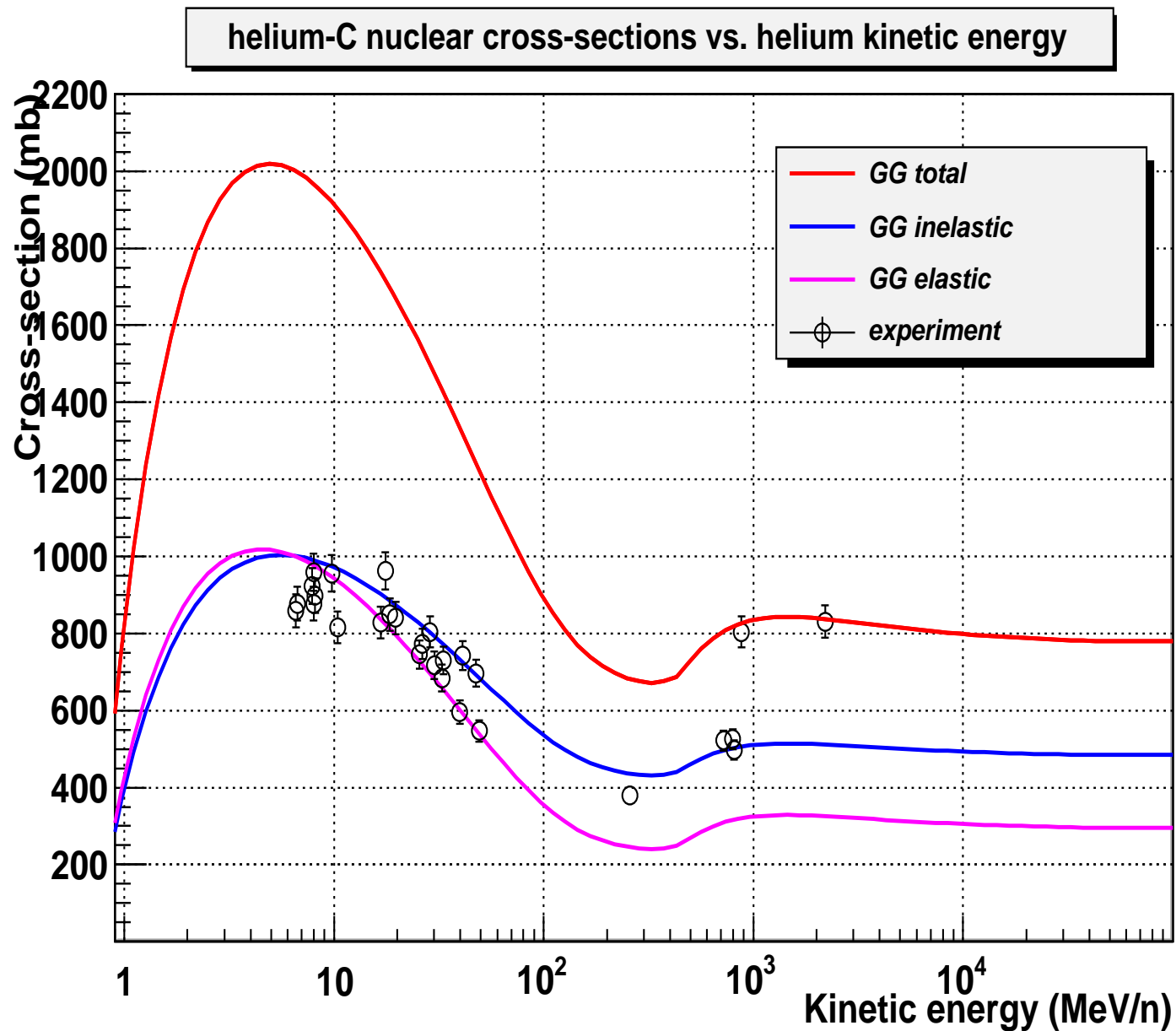


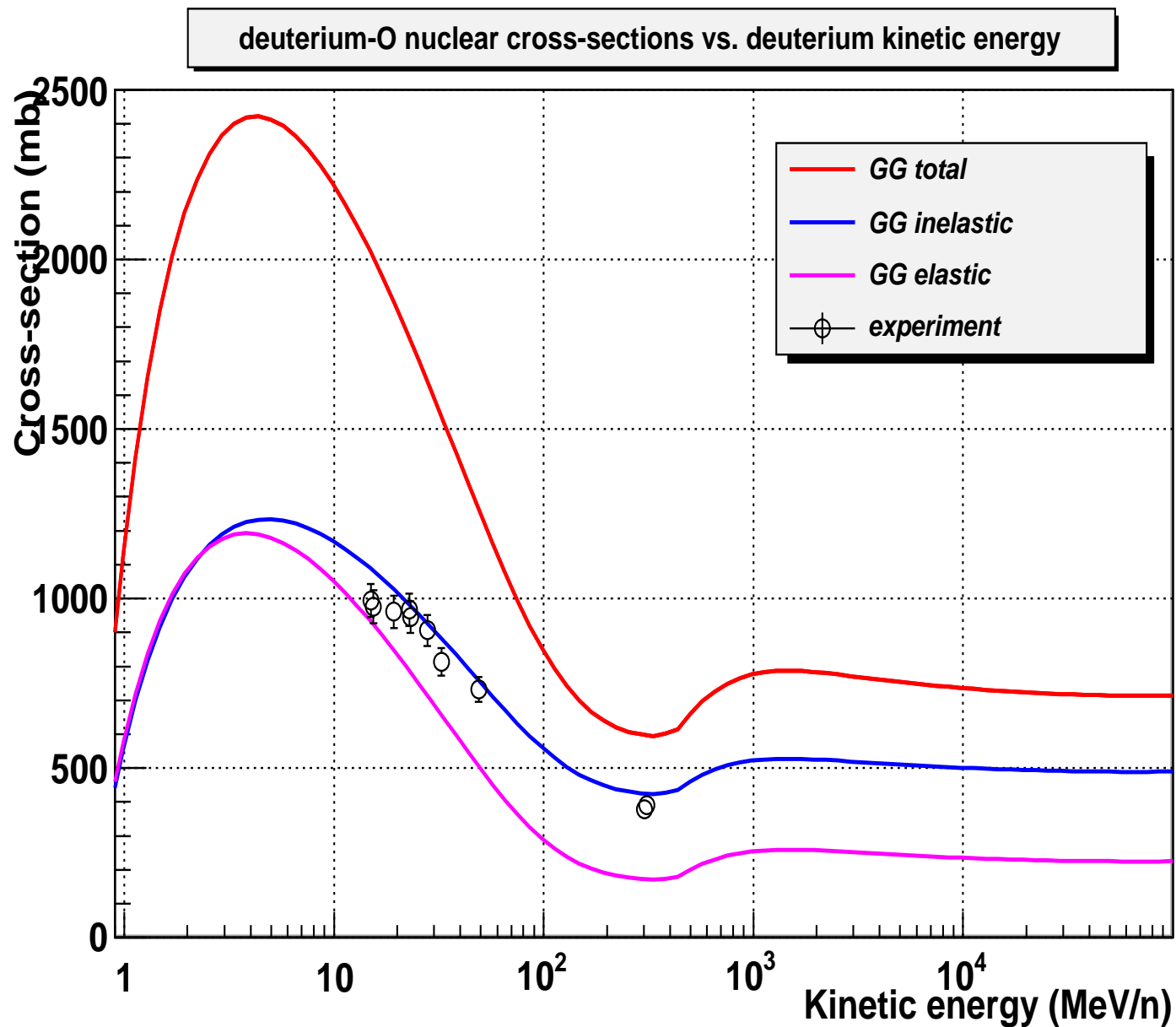


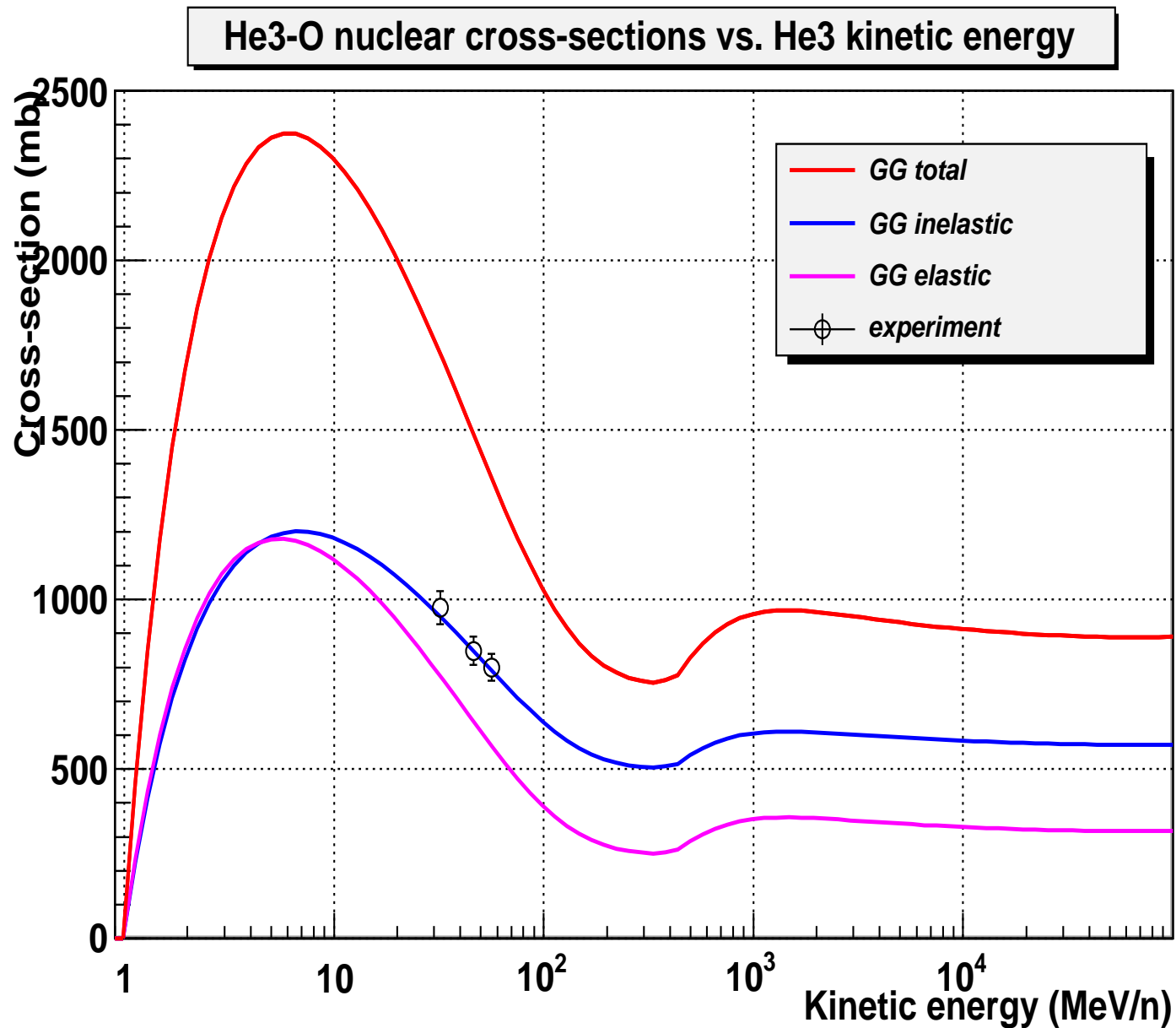


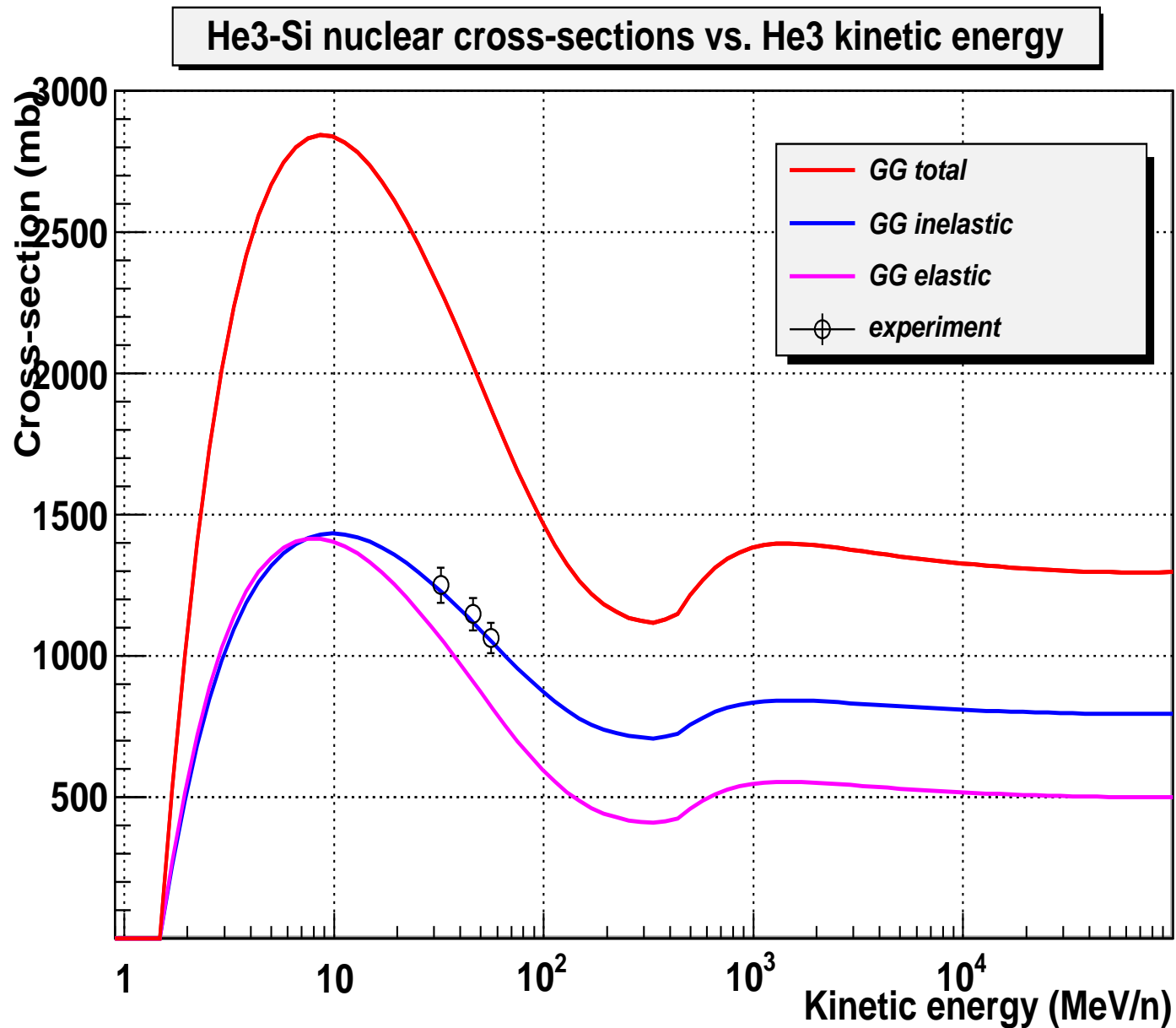


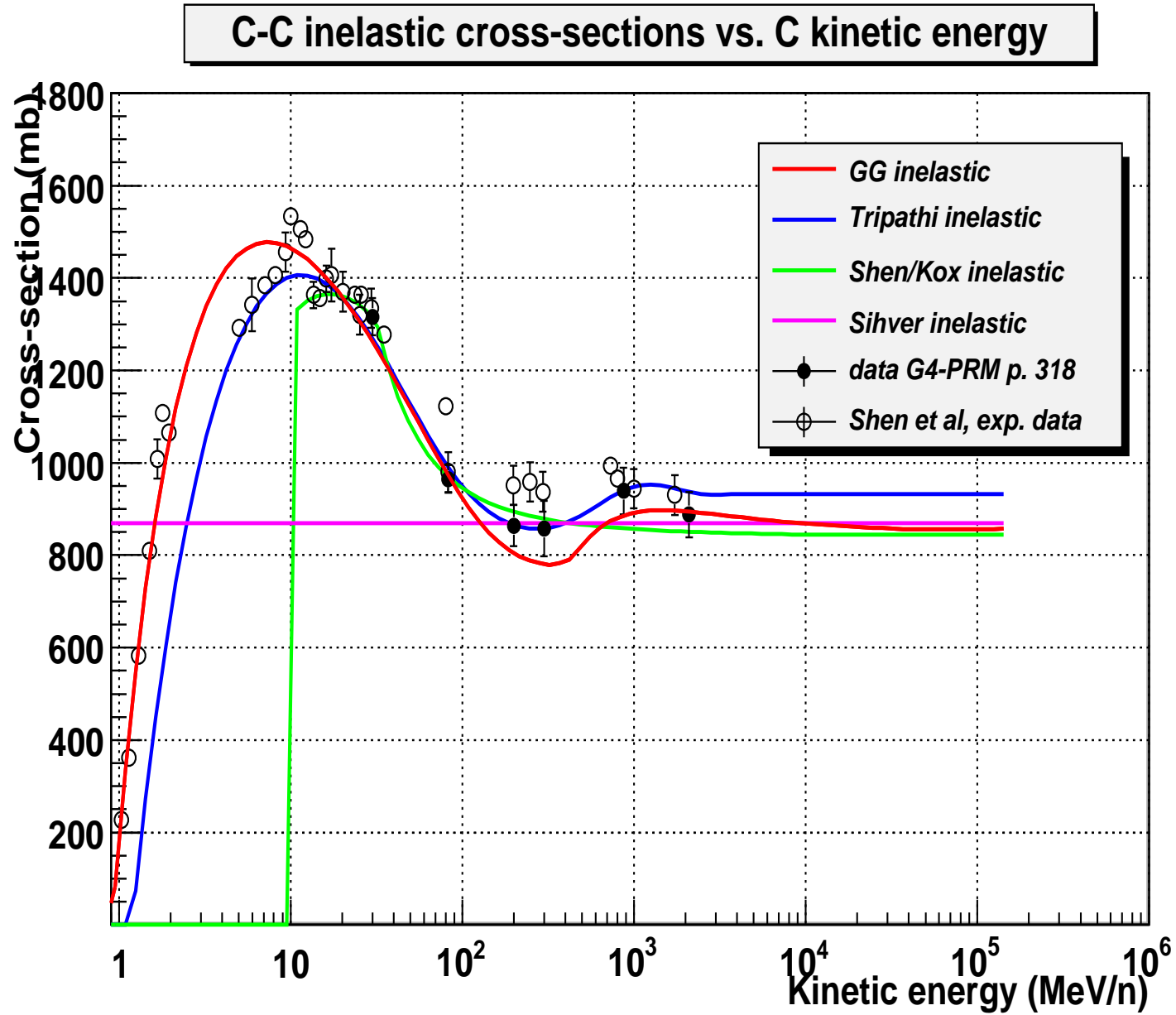


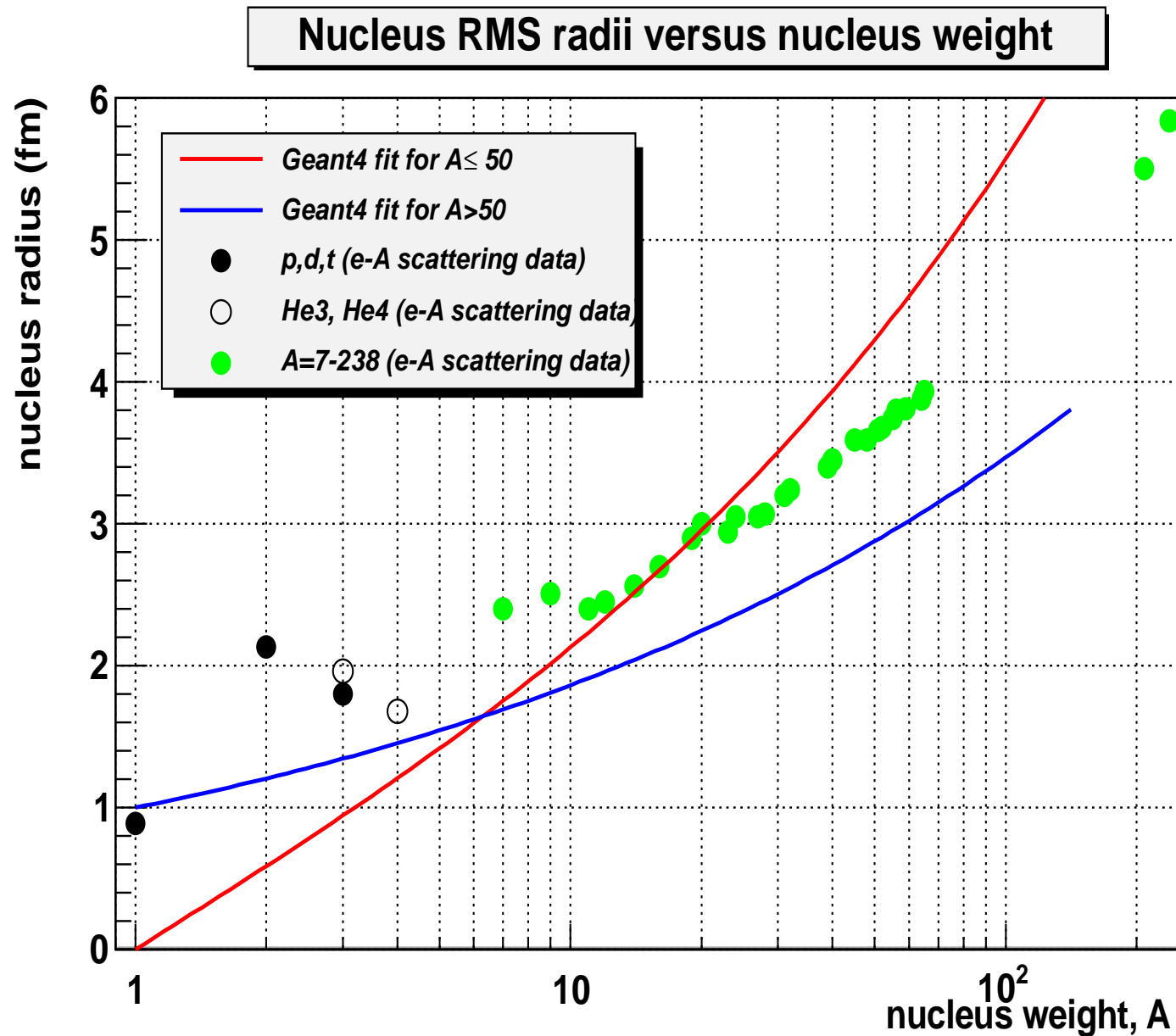












4 Summary

1. Glauber-Gribov approach provides the total, inelastic, and elastic light nucleus-nucleus cross-sections in a wide range of energies and projectile-target combinations.
2. The model algorithm is fast and as robust as $\sigma_{tot/in}^{NN}$ and $R(A)$ parameterizations.
3. The model is in agreement with the experimental data for the inelastic cross-section. The model is applicable starting from ~ 1 MeV/n up to TeV/n range (ALICE, Pb-Pb data). The both limits are defined mostly by the accuracy of the $\sigma_{tot/in}^{NN}$, $R(A)$ parameterizations and the Coulomb barrier phenomenology.