

An Outsider's perspective on information recovery in de Sitter space

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Based on **2210.12176** with

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and **WIP** with

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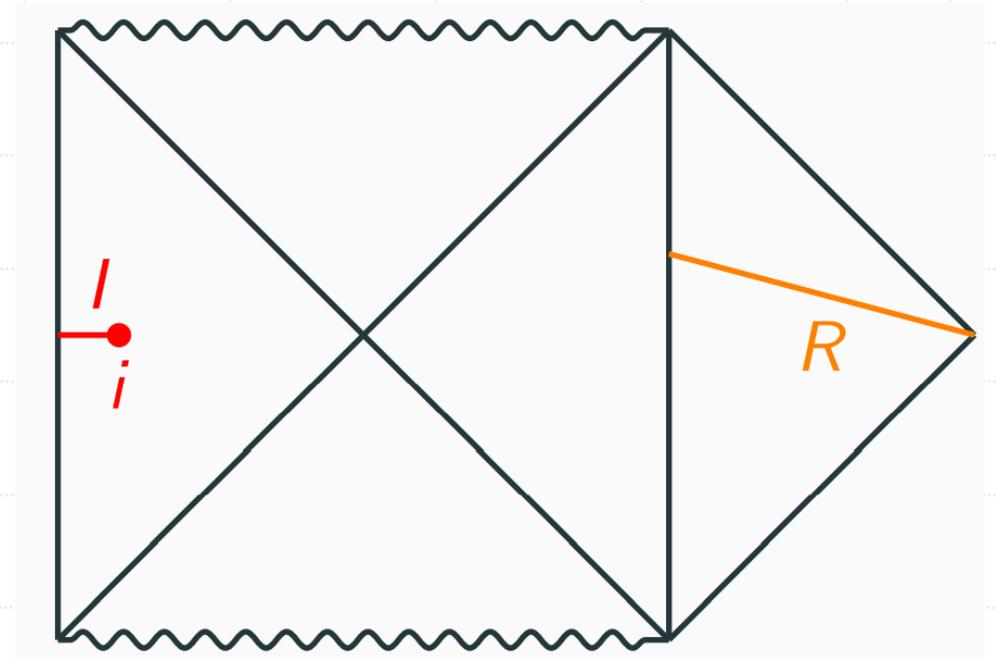
Motivation

- Exciting developments for reproducing the Page curve for BH evaporation.
- **Island rule** [Penington '19], [Almheiri, Engelhardt, Marolf, Maxfield '19]

$$S(R) = \text{Min Ext}_{\partial I} [S_{\text{gen}}(R \cup I)],$$

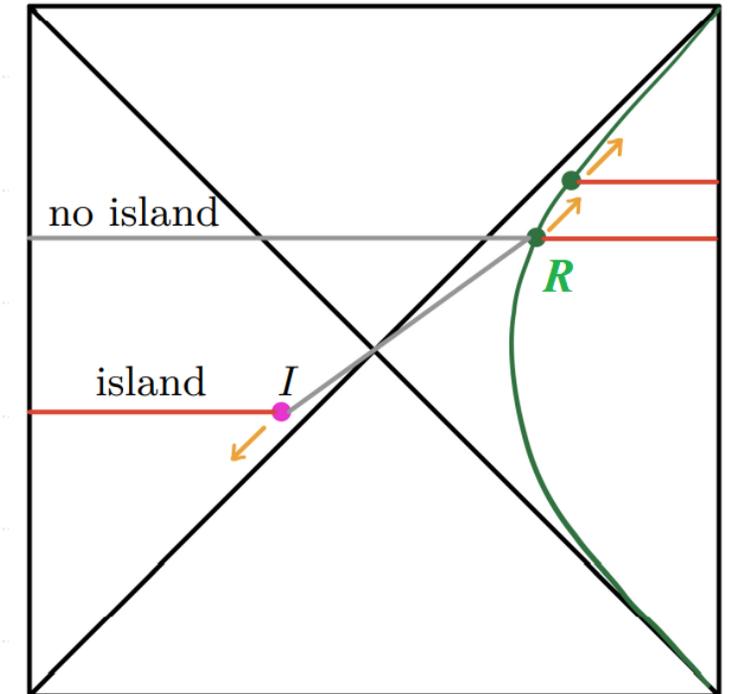
$$S_{\text{gen}}(R \cup I) = \frac{A(\partial I)}{4G_N} + S_{\text{vN}}(R \cup I).$$

- Non-perturbative transitions.



- Islands are agnostic about the spacetime background.
→ What do they teach us about cosmological horizons?

- **However**, extremal surfaces on the apparent horizon violate basic consistency principles [Shaghoulian '21] and go backwards in time [Sybesma '20].



There is a way around it [Aalsma, Sybesma, '21] at the cost of a large backreaction.

Is there a sense of controlled information recovery in dS space?

Two-Dimensional gravity model

$$I = \frac{1}{2\kappa^2} \int d^d x \sqrt{-g} \Phi \left(\mathbf{R} - \frac{2}{l^2} \right) + (\text{bdy term}) + (\text{matter})$$

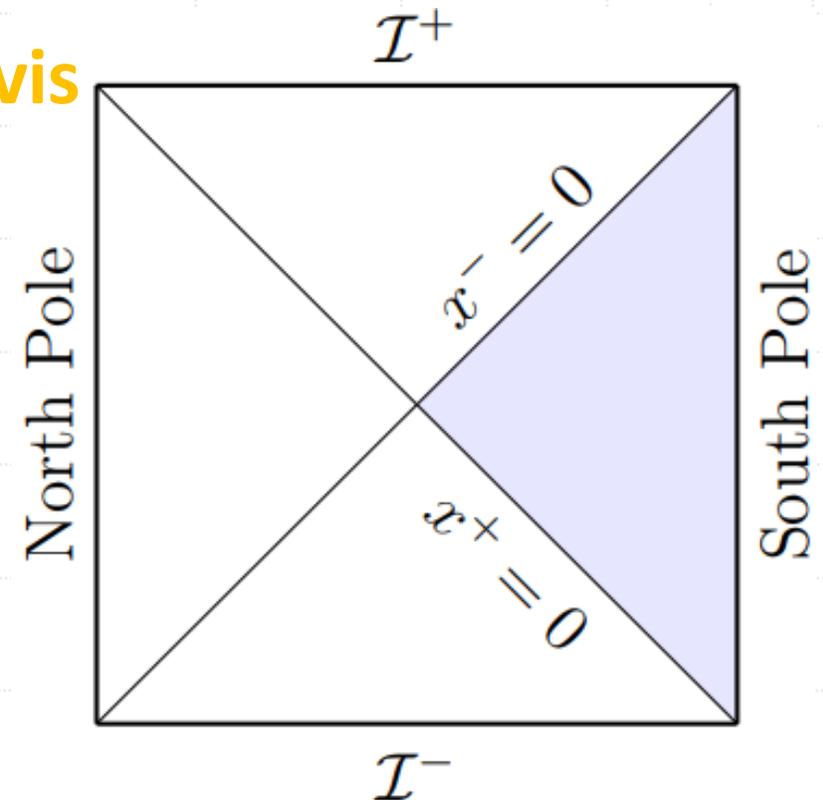
In presence of quantum matter in the **Bunch-Davis** state:

$$ds^2 = -\Omega_{\text{dS}}^{-2}(x^\pm) dx^+ dx^-,$$

$$\langle T_{\pm\pm}(x^\pm) \rangle_{\text{BD}} = 0,$$

$$\langle T_{\pm}(x^\pm) \rangle_{\text{BD}} = \frac{c g_{\pm}}{24\pi l^2},$$

$$\Phi_{\text{dS}}(x^\pm).$$



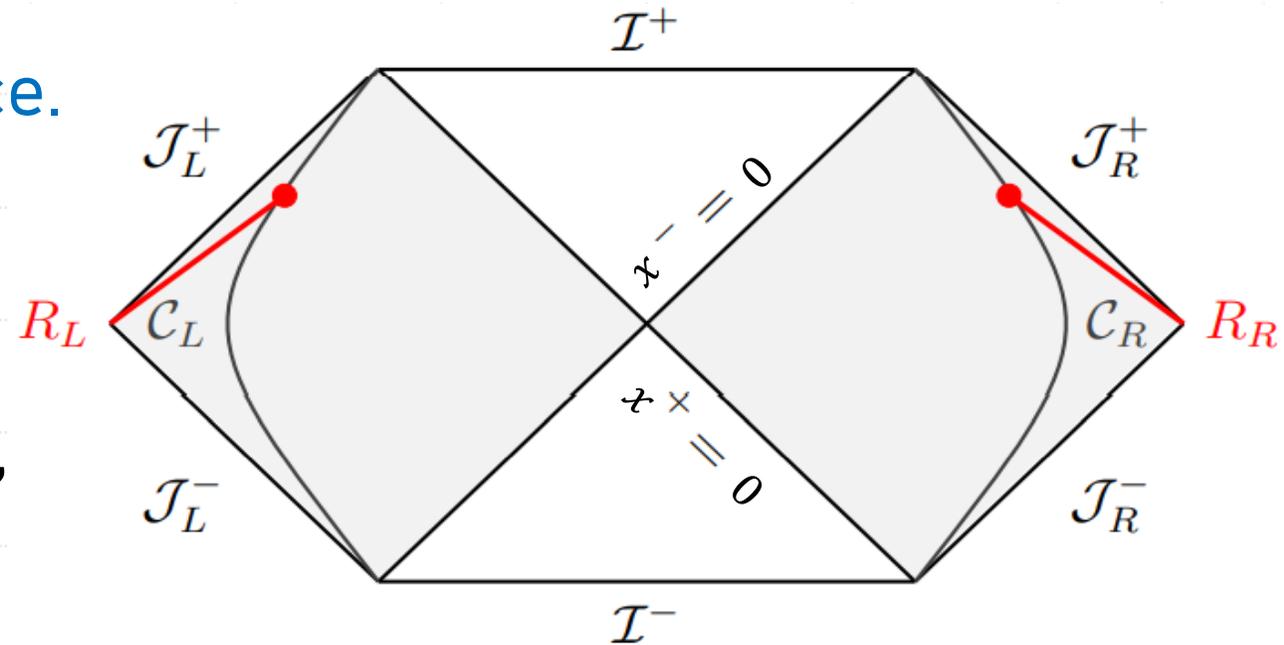
The Outsider's perspective model

- **Problem:** No asymptotic non-gravitational regions in dS space.

- **Proposal:** Glue Rindler wedges.

$$ds^2 = -\Omega_{\text{Rindler}}^{-2}(x^\pm) dx^+ dx^-,$$

$$\Phi_{\text{Rindler}}(x^\pm).$$



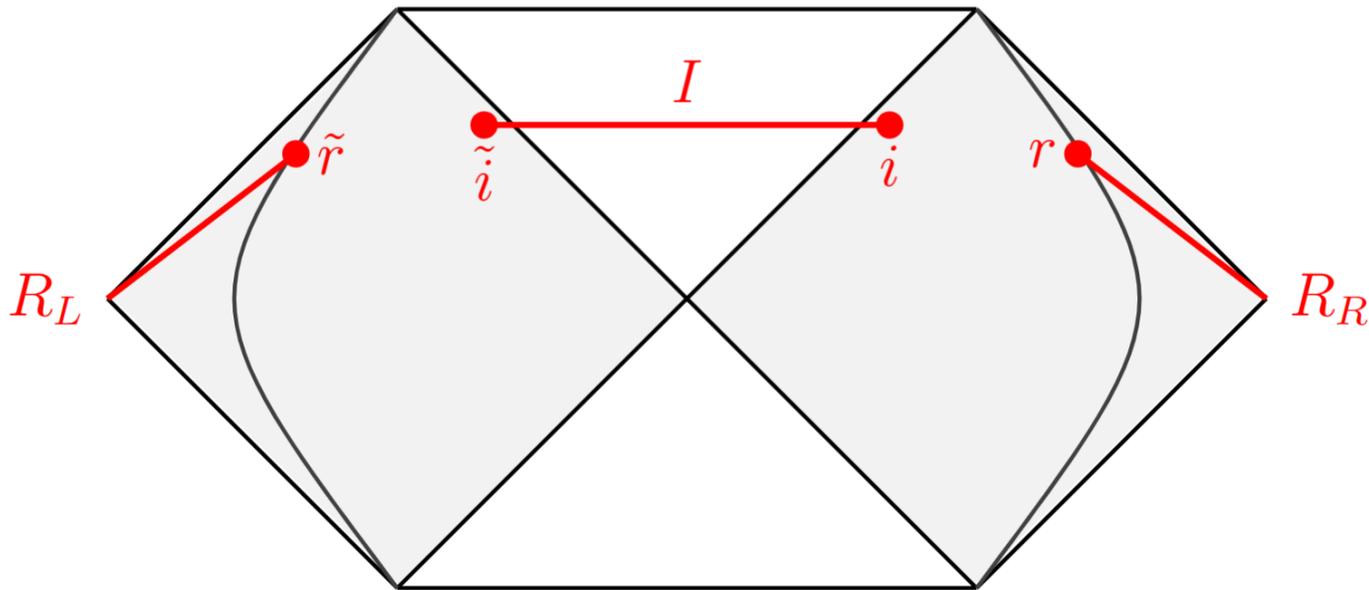
- **Junction conditions:**

$$[\Phi]_{x^-=0} = 0,$$

$$\kappa^2 T_{ab} n^a n^b + [n^a \nabla_a \Phi]_{x^-=0} \delta(x^-) = 0$$

Searching for Islands

$$S_{\text{gen}}(R_L \cup R_R \cup I) = \frac{2\pi}{\kappa^2} \Phi(x_i^\pm) + \frac{c}{6} \log \left[-\frac{(x_i^+ - x_r^+)(x_i^- - x_r^-)}{\epsilon_i \epsilon_r \Omega(x_i^\pm) \Omega(x_r^\pm)} \right] + (i \leftrightarrow \tilde{i}, r \leftrightarrow \tilde{r}).$$



Early times → Hawking saddle

$$S_{\text{gen}}(R_L \cup R_R \cup I) = S_{\text{vN}}(R) \sim \frac{c}{3l} t,$$

Late times → Rindler island saddle

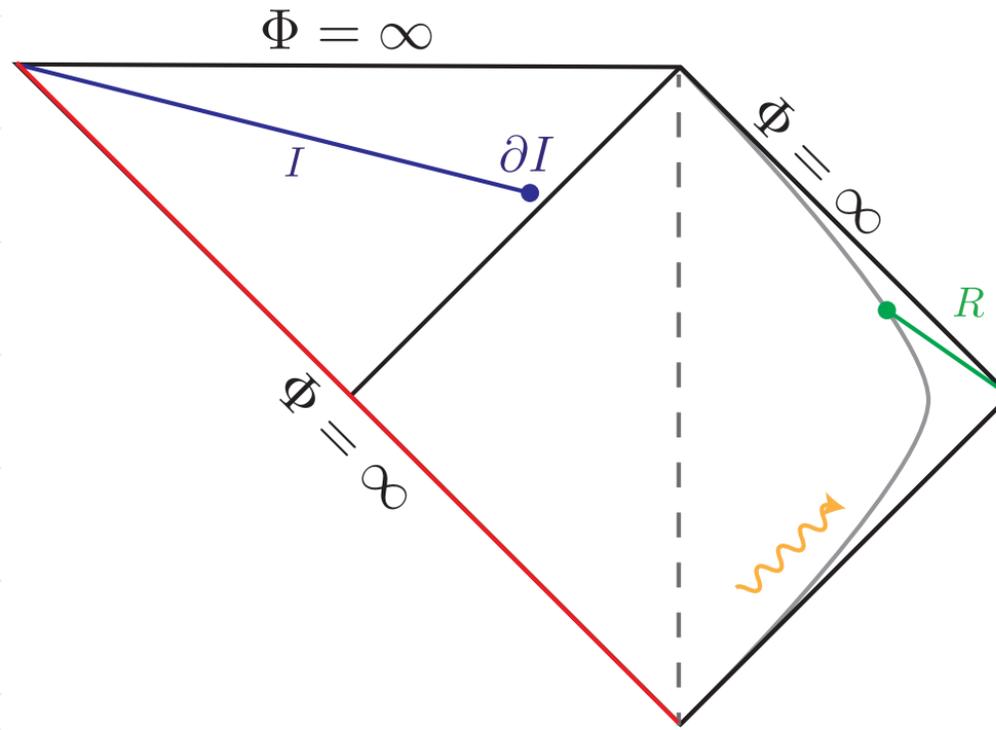
$$S_{\text{gen}}(R_L \cup R_R \cup I) = 2S_{\text{dS}}.$$

Breaking Thermal Equilibrium

- **Unruh-de Sitter state** [Aalsma , Parikh, van der Schaar '19]

$$\langle T_{++}(x^+) \rangle = -\frac{c}{48\pi(x^+)^2}, \quad \langle T_{--}(x^-) \rangle = 0.$$

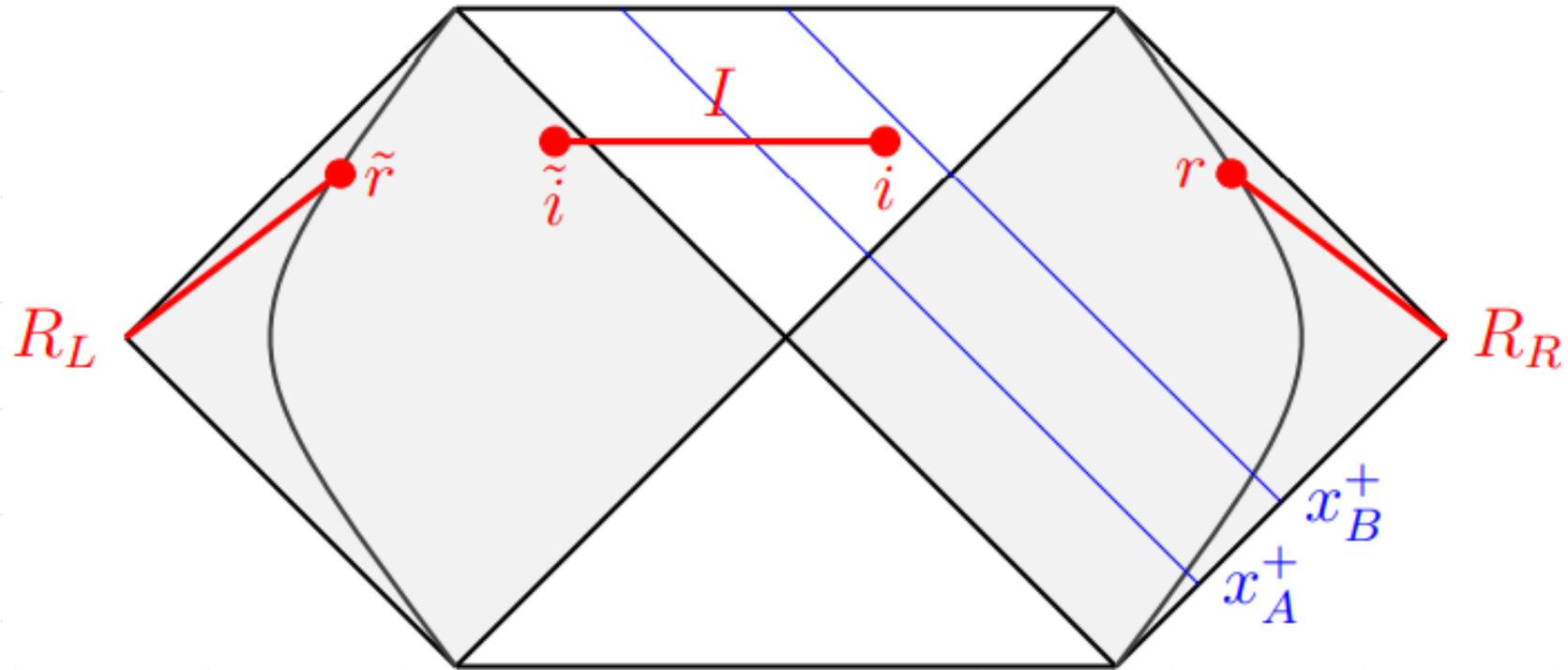
- However, curiosity comes with a **price** [Aalsma, Sybesma '21]



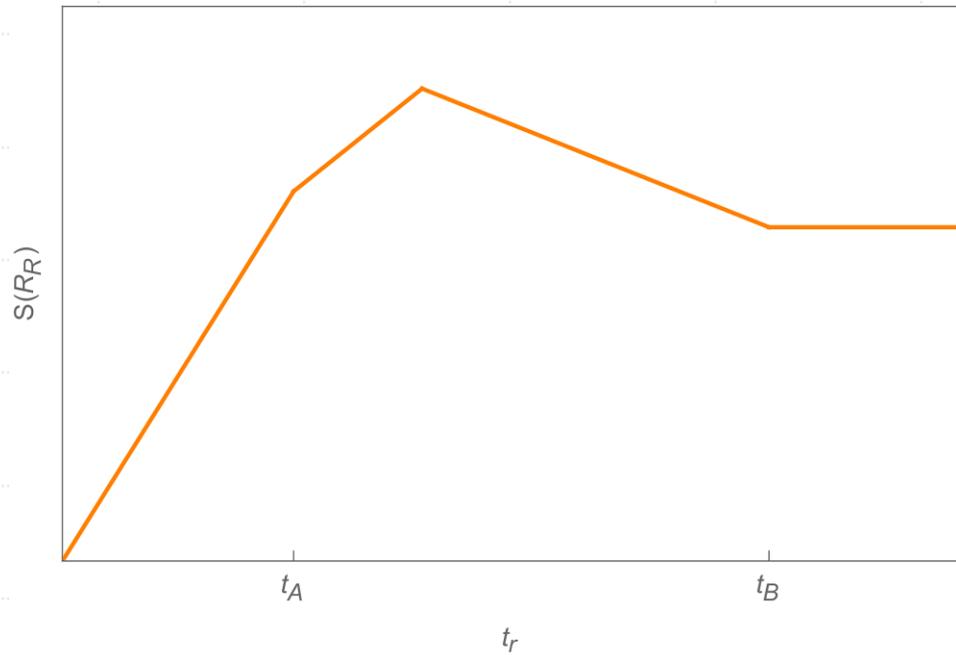
Information recovery

Without destroying the observer.

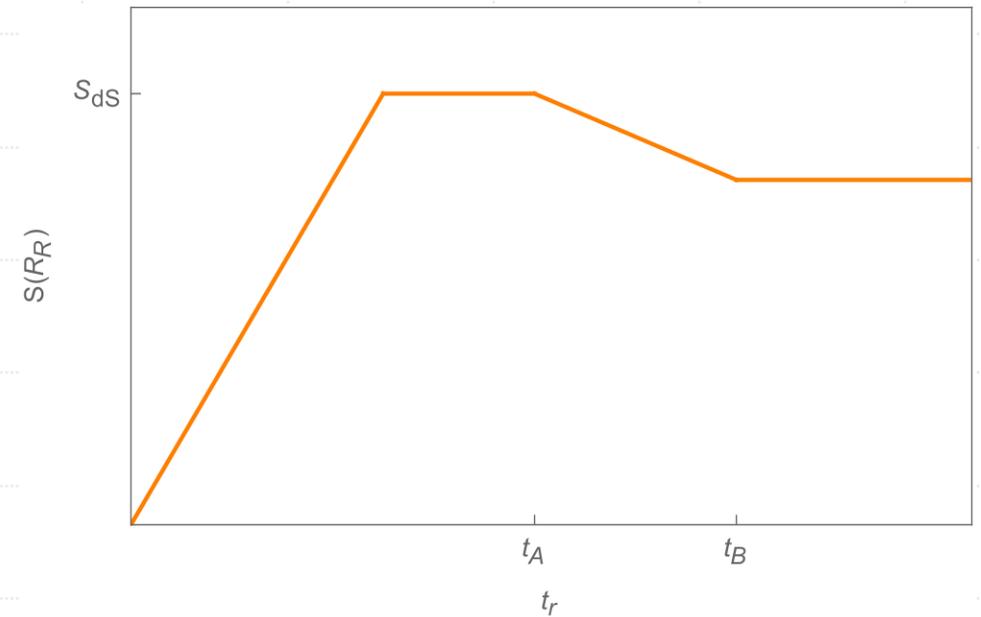
[Aalsma, SEAG, Sybesma '22]



Page curve



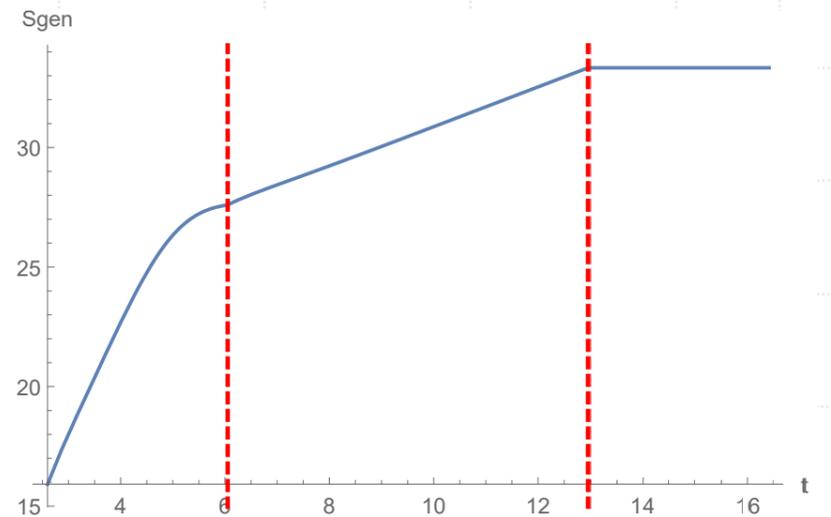
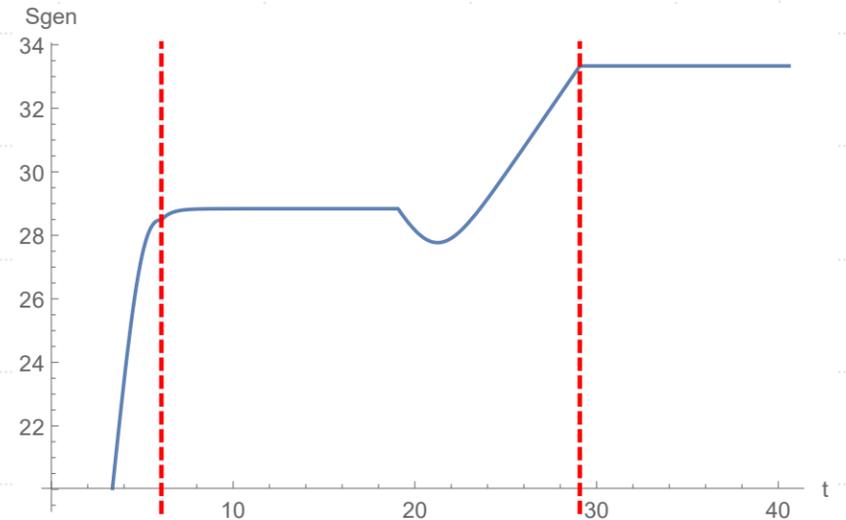
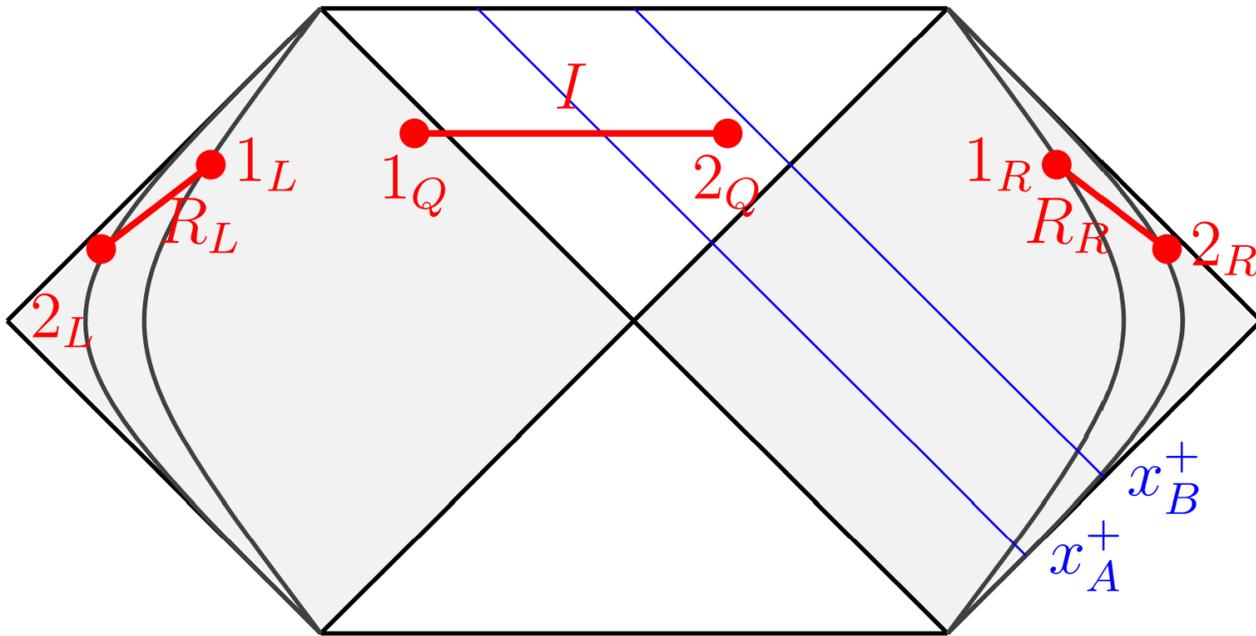
Before saturation time



After saturation time

Collecting radiation in finite segments

[SEAG, Hernandez, Khramtsov, Knysh **WIP**]



Summary

- The entanglement wedge of Hawking radiation includes regions behind the horizon, **independent on the background** [Bousso, Penington '22].
- The cost for recovering information in dS space is large backreaction, which can be **tamed** with our protocol.
- Entropy is sensitive to a notion of non-locality present in quantum gravity in the form of islands.
- Are there other low-energy observables that have this property?