

# Quantum Gravity Comes of Age

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Windows on Quantum Gravity

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# What Is Quantum Gravity?

Quantum Gravity =

- Gravity +
- Relativity +
- Quantum Mechanics

General relativity = Gravity (G) + Relativity (c)

- Works great IF we ignore fluctuations
- Eg., GPS accounts for slower time on Earth

Quantum Mechanics ( $\hbar$ )

- Each degree of freedom MUST fluctuate a little
- There are a LOT of degrees of freedom!

# The Problem of Quantum Gravity

For frequency  $\omega$  QG theory predicts

$$\left(\frac{\hbar G \omega^2}{c^5}\right)^N \left\{ 1 + \alpha_1 \left(\frac{\hbar G \omega^2}{c^5}\right) + \alpha_2 \left(\frac{\hbar G \omega^2}{c^5}\right)^2 + \dots \right\}$$

Two problems:

1. For  $\omega$ 's we can reach  $\left(\frac{\hbar G \omega^2}{c^5}\right) \ll 1$

Eg. FM radio  $\omega \sim 10^8 \text{ Hz} \rightarrow \left(\frac{\hbar G \omega^2}{c^5}\right) \sim 10^{-71}$

2. Theory predicts divergences in  $\alpha_1, \alpha_2, \dots$

# Mathematical Consistency Isn't Enough

**Gertrude Stein:** “The trouble with Oakland is that, when you get there, there isn't any there there.”

**String Theorist** (on being told the 2011 Nobel Prize-winning result that the universe is accelerating): “I'm sure the data is wrong because string theory predicts a negative cosmological constant. And if it's right, I'll quit doing physics.” (He's still doing string theory.)

**WE NEED DATA!**

## Big questions in M-theory

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Recently there has been a lot of discussion about questions in physics and mathematics that we should confront in the new millenium; see [The top 10 physics problems for the next millenium](#) and [The Clay Mathematics Institute millenium prize problems](#). In this spirit, and to stimulate discussion, I've begun collecting outstanding questions in M-theory. Please submit your favorite question to [giddings@physics.ucsb.edu](mailto:giddings@physics.ucsb.edu); questions added will be credited.

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## Fundamental principles

**What are the fundamental degrees of freedom of M-theory?**

**(Does this question apply, or is it time to abandon the reductionist paradigm?)**

**What are the dynamical laws governing their interaction?**

**What are the observables of M-theory? (*R. Bousso*)**

# Cosmology

Flat space:  $ds^2 = -c^2 dt^2 + d\vec{x} \cdot d\vec{x}$

Cosmology:  $ds^2 = -c^2 dt^2 + a^2(t) d\vec{x} \cdot d\vec{x}$

- $3c^2 H^2 = 8\pi G \rho$  E.g.  $\rho_{mat} \sim \left(\frac{a_0}{a(t)}\right)^3$
- $(-3 + 2\epsilon)c^2 H^2 = 8\pi G p$  E.g.  $p_{rad} = \frac{1}{3} \rho_{rad}$

Two derivatives:

- Hubble parameter:  $H(t) \equiv \frac{\dot{a}}{a}$
- 1<sup>st</sup> slow roll parameter:  $\epsilon(t) \equiv -\frac{\dot{H}}{H^2}$

Current values:

- $H_0 = 67.3 \pm 1.2 \frac{km}{s-Mpc} \sim 2.2 \times 10^{-18} Hz$
- $\epsilon_0 = 0.47 \pm 0.03$

# A Cautionary Tale

## Urgent Mission



WE WERE GOING TO USE THE TIME MACHINE TO PREVENT THE ROBOT APOCALYPSE, BUT THE GUY WHO BUILT IT WAS AN ELECTRICAL ENGINEER.

All hate Franklin's signs  
Clever theorists tried to  
avoid this with  $\ddot{a}(t)$

- Before 1998 we were **CERTAIN**  $\ddot{a}(t_0) < 0$
- Hence "deceleration parameter"  $q(t) \equiv -\frac{a\ddot{a}}{\dot{a}^2}$
- But  $q_0 = -0.53 \pm 0.03$

Moral: don't believe  
before you see the data

# Brief History of $\epsilon(t)$

$N$  = # of e-foldings since inflation

- $a(t) = a_i e^N$
- NB  $e^{60} \cong 10^{26}$
- 6000 Mpc now was 2 m then!

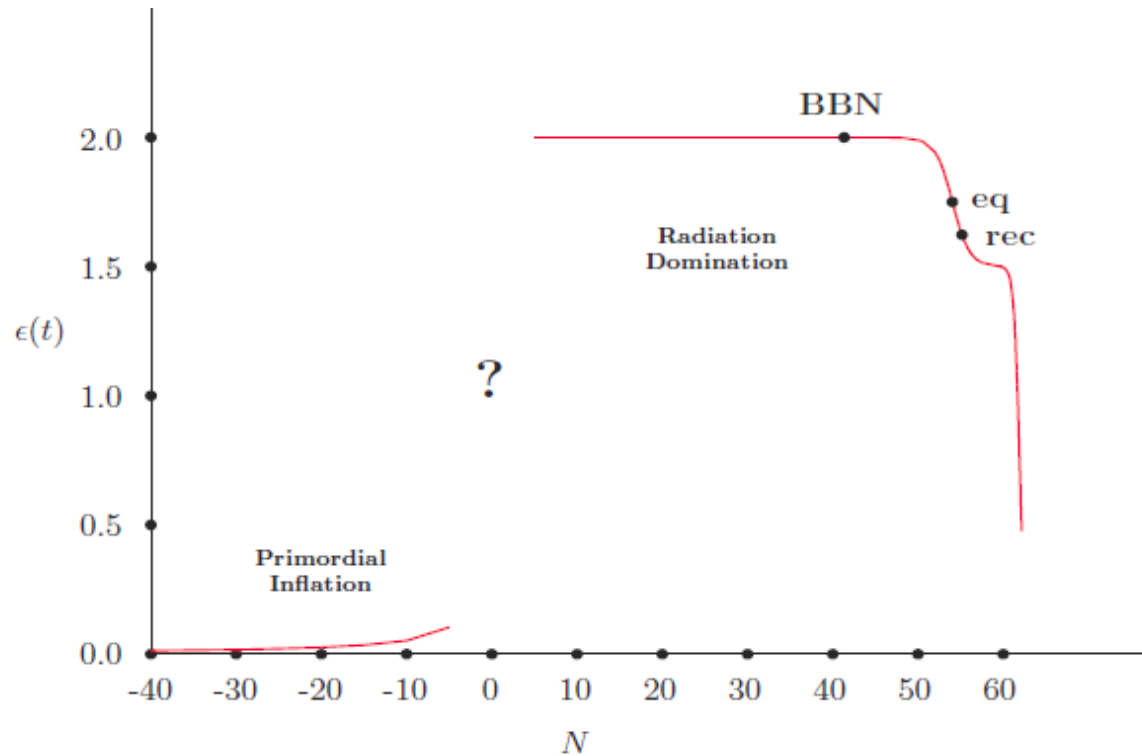
Cosmological Epochs (Roughly)

1. Late acceleration  $59 \leq N \leq 60$
2. Matter domination  $52 \leq N \leq 59$
3. Radiation dom.  $5 \leq N \leq 52$
4. Reheating  $-5 \leq N \leq 5$
5. Primordial Inflation  $N \leq -5$

We can observe early events

1. Recombination  $t \cong 300,000$  yrs
2. Big Bang Nucleosynthesis  
 $t \cong 1$  s

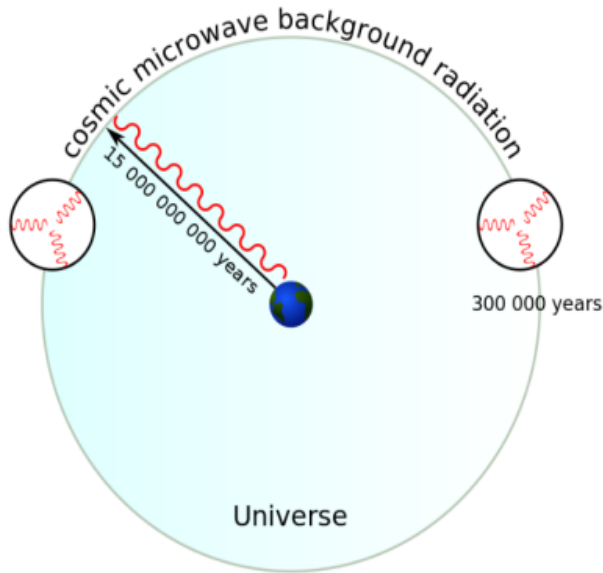
NB  $t_0 \cong 13,800,000,000$  yrs



# Horizon Problem

$$0 = -c^2 dt^2 + a^2(t) dr^2 \rightarrow dr = \frac{cdt}{a(t)}$$

CMB in thermal equilibrium to  
1 part in  $10^5$  (better than this room!)



## Past and Future Light-cones

- $P(t) \equiv \int_t^{t_0} dt' \frac{ca_0}{a(t')}$
- $F(t) \equiv \int_{t_i}^t dt' \frac{ca_0}{a(t')}$

$q > 0 \rightarrow$  upper limit dominates

- ~5000 causally distinct regions in equilibrium!
- $\sim 10^{17}$  regions at BBN!

$q < 0 \rightarrow$  lower limit dominates

- $a(t_i) \rightarrow 0$  makes  $F(t) \rightarrow \infty$



# QG Effects from Inflation observable because:

1.  $\frac{\hbar G H_{inf}^2}{c^5} \sim 10^{-11}$  is small but observable

• Compare  $\frac{\hbar G H_0^2}{c^5} \sim 10^{-122}$

2. Lots of graviton & MMC scalar production

3. Long  $\lambda$  perturbations “fossilize” so they can survive to late times

# Energy-Time Uncertainty Principle in Flat Space

Virtual particles of energy  $E$  exist for  $\Delta t \cong \frac{\hbar}{E}$

Wave number  $k = \frac{2\pi}{\lambda} \rightarrow E = [m^2 c^4 + c^2 \hbar^2 k^2]^{1/2}$

This is why  $m = 0$  particles give  $F(r) \propto \frac{1}{r^2}$  forces

- $m = 0 \rightarrow E = c\hbar k \rightarrow \Delta t \cong \frac{1}{ck}$
- $r = c\Delta t = \frac{1}{k}$
- $F = \frac{\Delta p}{\Delta t} = \frac{\hbar k}{1/ck} = \frac{\hbar c}{r^2}$

Also why  $m \neq 0$  particles give short range forces

- $m \neq 0 \rightarrow E \sim mc^2 \rightarrow \Delta t \cong \frac{\hbar}{mc^2}$
- $r = c\Delta t = \frac{\hbar}{mc}$  fixed

Any inflationary virtual with  $ck < H(t)a(t)$  lives forever

$k = \frac{2\pi}{\lambda}$  fixed but  $\frac{k}{a(t)}$  physical, hence

$$E(t, k) = \left[ m^2 c^4 + \left( \frac{c\hbar k}{a(t)} \right)^2 \right]^{1/2}$$

Virtuals can live for  $\int_t^{t+\Delta t} dt' E(t', k) \cong \hbar$

$m = 0$  virtuals have  $c\hbar k \int_t^{t+\Delta t} \frac{dt'}{a(t')} \cong \hbar$

- We just saw this integral!

- For inflation can take  $\Delta t \rightarrow \infty \rightarrow \frac{c\hbar k}{H(t)a(t)} \cong \hbar$

# A Killer Symmetry: Conformal Invariance

- $d\eta \equiv \frac{dt}{a(t)} \rightarrow ds^2 = a^2[-d\eta^2 + dx^2] \rightarrow g_{\mu\nu} = a^2\eta_{\mu\nu}$
- Conformal transformation using  $\Omega(\eta, x)$ 
  - $g_{\mu\nu} \rightarrow \Omega^2 g_{\mu\nu}$  ,  $\psi \rightarrow \Omega^{-3/2}\psi$
  - $A_\mu \rightarrow A_\mu$  ,  $\varphi \rightarrow \Omega^{-1}\varphi$
- For example,  $\mathcal{L}_{EM} = -1/4 F_{\mu\nu} F_{\rho\sigma} g^{\mu\rho} g^{\nu\sigma} \sqrt{-g}$
- Conformal invariance  $\rightarrow$  same  $\mathcal{L}$  as flat
- Emission rate  $\frac{dN}{d\eta} = \Gamma_{flat} \rightarrow \frac{dN}{dt} = \frac{\Gamma_{flat}}{a(t)}$
- Any emitted virtual lives forever, but few emitted

# Exceptions: MMC scalars & gravitons

- $L = \frac{m}{2}[\dot{q}^2 - \omega^2 q^2] \rightarrow q(t) = \alpha u(t) + \alpha^\dagger u^*(t)$ 
  - $\frac{d}{dt}(m\dot{u}) + m\omega^2 u = 0 \rightarrow u(t) = \sqrt{\frac{\hbar}{2m\omega}} \text{Exp}\left[-i \int_{t_i}^t dt' \omega\right]$
  - $E(t) = \langle \Omega | \frac{m}{2}[\dot{q}^2(t) + \omega^2 q^2(t)] | \Omega \rangle = \frac{m}{2}[|\dot{u}(t)|^2 + \omega^2 |u(t)|^2] = \frac{1}{2}\hbar\omega$
- $\mathcal{L}_{MMCS} = \frac{a^3}{2c^2} \left[ \dot{\varphi}^2 - \frac{c^2 \|\vec{\nabla}\varphi\|^2}{a^2} \right] \rightarrow$  each  $\vec{k}$  SHO with  $m(t) = \frac{a^3}{c^2}$  &  $\omega(t) = \frac{ck}{a}$ 
  - $u(t, k) = \sqrt{\frac{\hbar c}{2ka^2(t)}} \left[ 1 + \frac{iHa(t)}{ck} \right] \text{Exp}\left[-i \int_{t_i}^t dt' \frac{ck}{a(t')}\right]$  for  $\epsilon = 0$
  - $E(t, k) = \frac{a^3(t)}{2c^2} \left[ |\dot{u}(t, k)|^2 + \frac{c^2 k^2}{a^2(t)} |u(t, k)|^2 \right] = \frac{\hbar ck}{a(t)} \left[ \frac{1}{2} + \left( \frac{Ha(t)}{2ck} \right)^2 \right]$
- Each  $\vec{k}$  has  $N(t, k) = \left[ \frac{Ha(t)}{2ck} \right]^2$  conformal has  $N(t, k) = 0$
- NB  $|u(t, k)|^2 \rightarrow \frac{\hbar H^2}{2ck^3}$  a constant! conformal has  $u(t, k) \rightarrow 0$

# We Have Data!

- Scalar Power spectrum

$$\Delta_{\mathcal{R}}^2(k) = A_s \left( \frac{k}{k_0} \right)^{n_s - 1}$$

$$A_s = 2.196^{+0.051}_{-0.060} \times 10^{-9}$$

$$n_s = 0.9603 \pm 0.0073$$

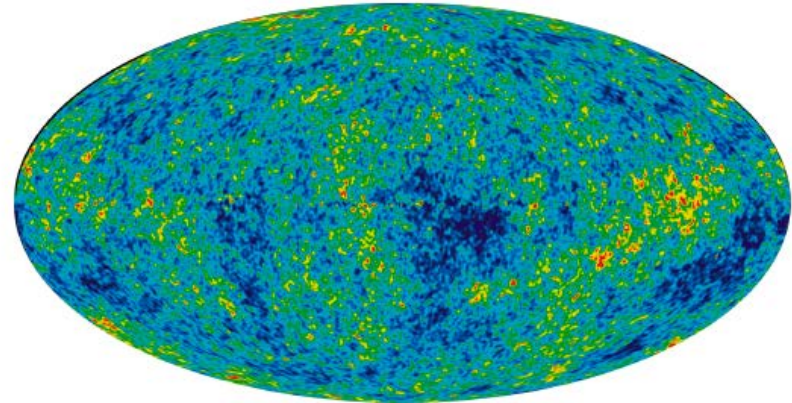
$$k_0 = 0.050 \text{ Mpc}^{-1}$$

- Tensor-to-Scalar ratio

$$r(k) \equiv \frac{\Delta_h^2(k)}{\Delta_{\mathcal{R}}^2(k)}$$

$$r(0.002 \text{ Mpc}^{-1}) < 0.09$$

- Many new polarization probes
- Already  $10^7$  data points
- Factor of  $10^{14}$  increase in 21 cm radiation from early structures



$$\Delta_{\mathcal{R}}^2(k) \approx \frac{\hbar G H^2(t_k)}{\pi c^5 \epsilon(t_k)}$$

$$\Delta_h^2(k) \approx \frac{16 \hbar G H^2(t_k)}{\pi c^5}$$

# This is quantum gravitational data!

## The first ever collected!

1. “It’s not QG if it’s only tree order”
  - $\Delta_{\mathcal{R}}^2$  &  $\Delta_h^2 \sim \hbar G \rightarrow$  it’s quantum gravity
  - Famous tree effects in QFT:
    - $\rightarrow$  Bhabha scattering (QED) & beta decay (weak int)
2. “It’s not QG if it doesn’t involve gravitons”
  - Solar system tests response to classical matter
  - $\Delta_{\mathcal{R}}^2(k)$  tests response to quantum matter
  - More interesting for QG than gravitons
3. “It’s not QG if we can’t predict it uniquely”
  - Rotation curves & Hubble plots measure gravity
  - Even though we don’t understand what drives them

# Impact on Quantum Gravity

## 1. Discretization to solve UV no longer viable

- Minimum  $\Delta x$  cuts off UV divergences

But QG only small if  $\Delta x^2 \gg \frac{\hbar G}{c^3} \sim (10^{-35} m)^2$

- QG effects small at  $t_k$  with  $\frac{a_0}{a(t_k)} \sim e^{110} \sim 10^{48}$

$a(t_k)\Delta x \sim 10^{-35} m \rightarrow a_0\Delta x \sim 10^{13} m$

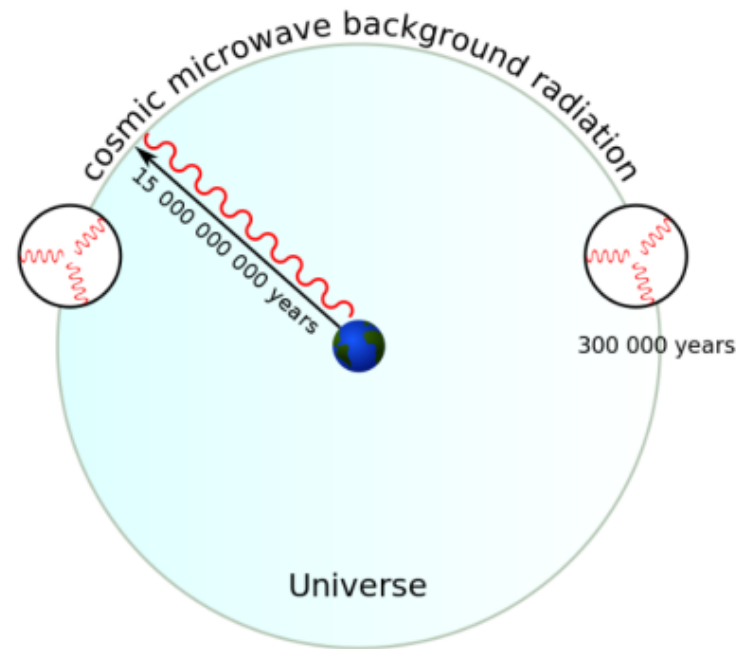
## 2. Failing to quantize gravity no longer viable

- $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4} \langle \Psi | T_{\mu\nu} | \Psi \rangle$
- But  $\langle \Psi | T_{\mu\nu} | \Psi \rangle$  has no fluctuations!



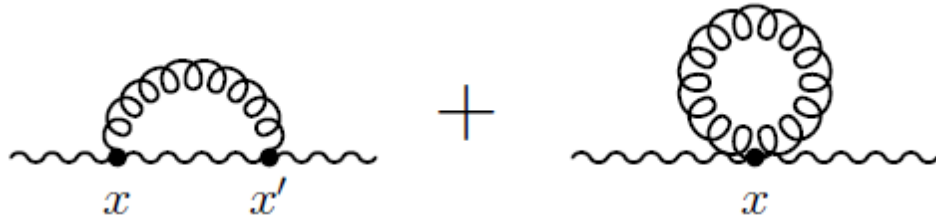
# A Tomograph is better than an X-ray!

- CMB from entire path
  - $N \sim \ell_{max}^2$
  - $\ell_{max} \sim 3000$
- 21 cm redshift fixes WHERE it came from
  - $N \sim \ell_{max}^3$
  - $\ell_{max} \sim 10^7$
  - Data to resolve one loop
- Not easy, or quick
  1. Lunar radio telescope
  2. Superb astrophysics
  3. Unique inflation model



# Scalars & gravitons alter kinematics of other quanta

- Vacuum polarization  $i[\mu\Pi^\nu](x; x')$



- Quantum-corrected Maxwell Equation

$$\partial_\nu[\sqrt{-g}g^{\nu\rho}g^{\mu\sigma}F_{\rho\sigma}(x)] + \int d^4x'[\mu\Pi^\nu](x; x')A_\nu(x') = J^\mu(x)$$

- $J^\mu(x) = 0 \rightarrow$  dynamical photons
- $J^\mu(x) \neq 0 \rightarrow$  electromagnetic forces

# Known Secondary Effects

## MMC Scalar Effects

1. On themselves
  - Brunier, Kahya, Onemli
2. On fermions
  - Garbrecht, Prokopec
3. On photons
  - Degueldre, Prokopec, Tornkvist
4. On gravitons
  - Leonard, Park, Prokopec

## Graviton Effects

1. On scalars
  - Kahya, Park
2. On fermions
  - Miao
3. On photons
  - Glavan, Leonard, Miao, Prokopec
4. On gravitons
  - Mora, Tsamis

# Conclusions → Dawning “Golden Age” of QG

- QG from primordial inflation small but observable

1.  $H_{inf} \sim 10^{56} \times H_0 \rightarrow \frac{\hbar G H_{inf}^2}{c^5} \sim 10^{-11}$

2.  $\epsilon_{inf} < 0.0056 \rightarrow$  gravitons & MMC scalars have  $N(t, \vec{k}) \sim \left[ \frac{H a}{2 c k} \right]^2$

3. Fossilized perturbations survive to late times

- Primary effects

- $\Delta_{\mathcal{R}}^2(k) \sim \frac{\hbar G H^2(t_k)}{\pi c^5 \epsilon(t_k)}$  (resolved)       $\Delta_h^2(k) \sim \frac{16 \hbar G H^2(t_k)}{\pi c^5}$  (not yet)

- Secondary effects → loops & kinematics of other quanta

- Enough data to eventually resolve

- Crucial questions

1. What is causing late time acceleration?
2. What drove primordial inflation?
3. What are the observables?