

Cool quark matter

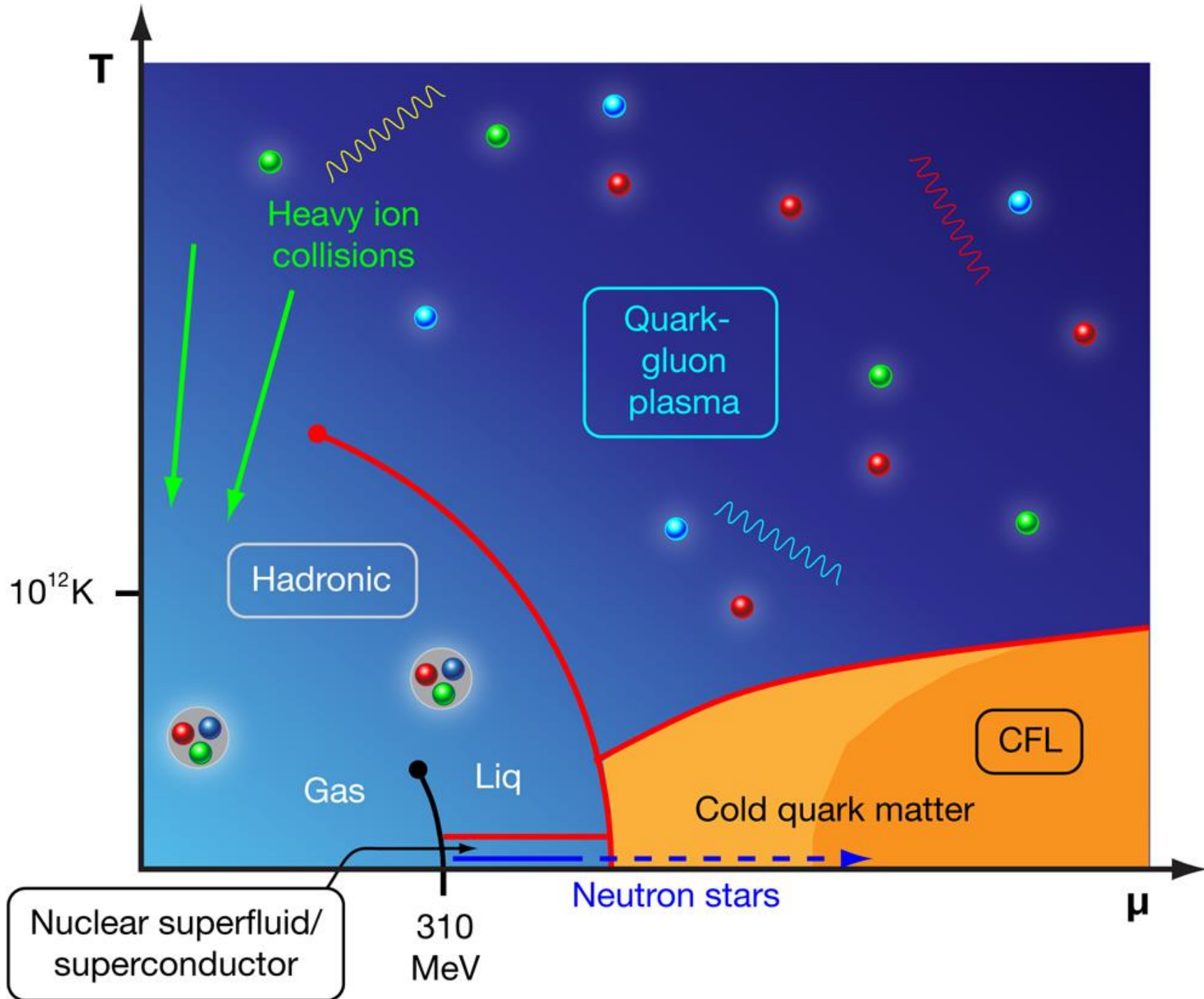
**Alexi Vuorinen
University of Helsinki**

**NewCompStar Conference, Istanbul
28.4.2016**

Based on: Alexi Kurkela, AV, arXiv:1603.00750 [hep-ph]

Ultimate goal in nuclear theory:
quantitative description of systems from
single nuclei to nuclear and quark matter
from first principles, i.e. QCD

Tools: Lattice QCD, Effective Field Theory,
perturbation theory



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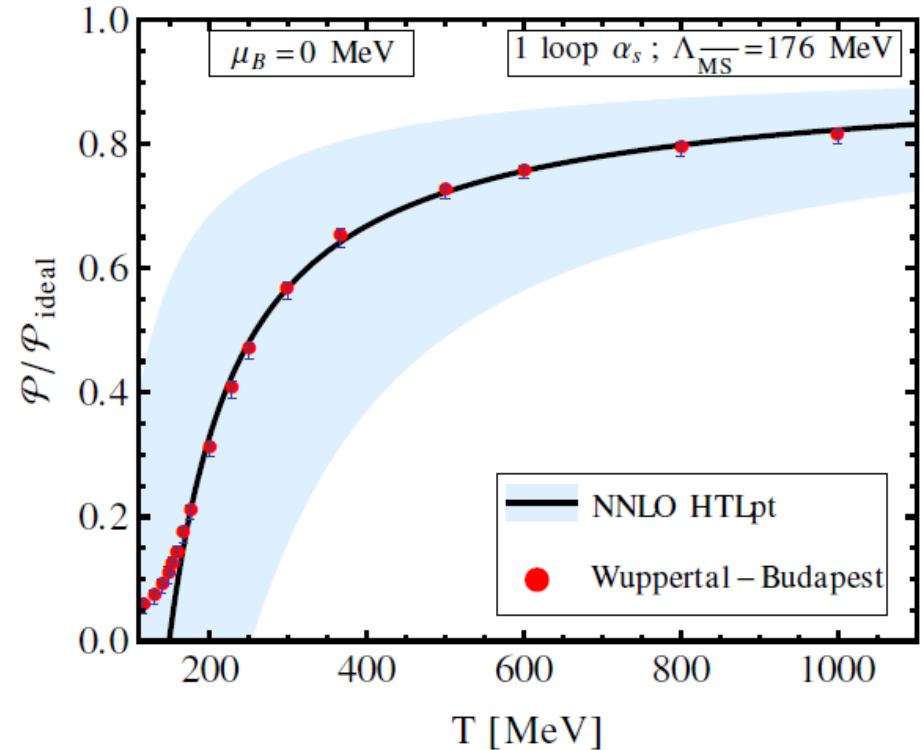
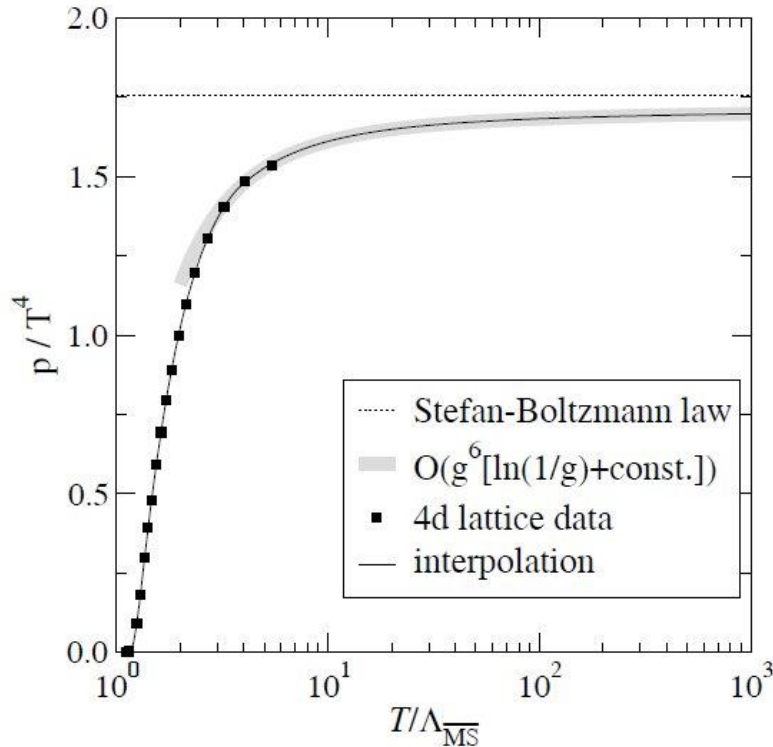
Tools: Lattice QCD, Effective Field Theory,
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This talk: bulk thermodynamics at small but
nonzero temperatures

- I. Perturbation theory at finite density:
from quark matter to neutron stars
- II. Challenges at small but nonzero
temperature
- III. New result: perturbative EoS of cool
quark matter

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High temperatures: pQCD standard tool complementing lattice, with excellent agreement above ca. $2T_c$

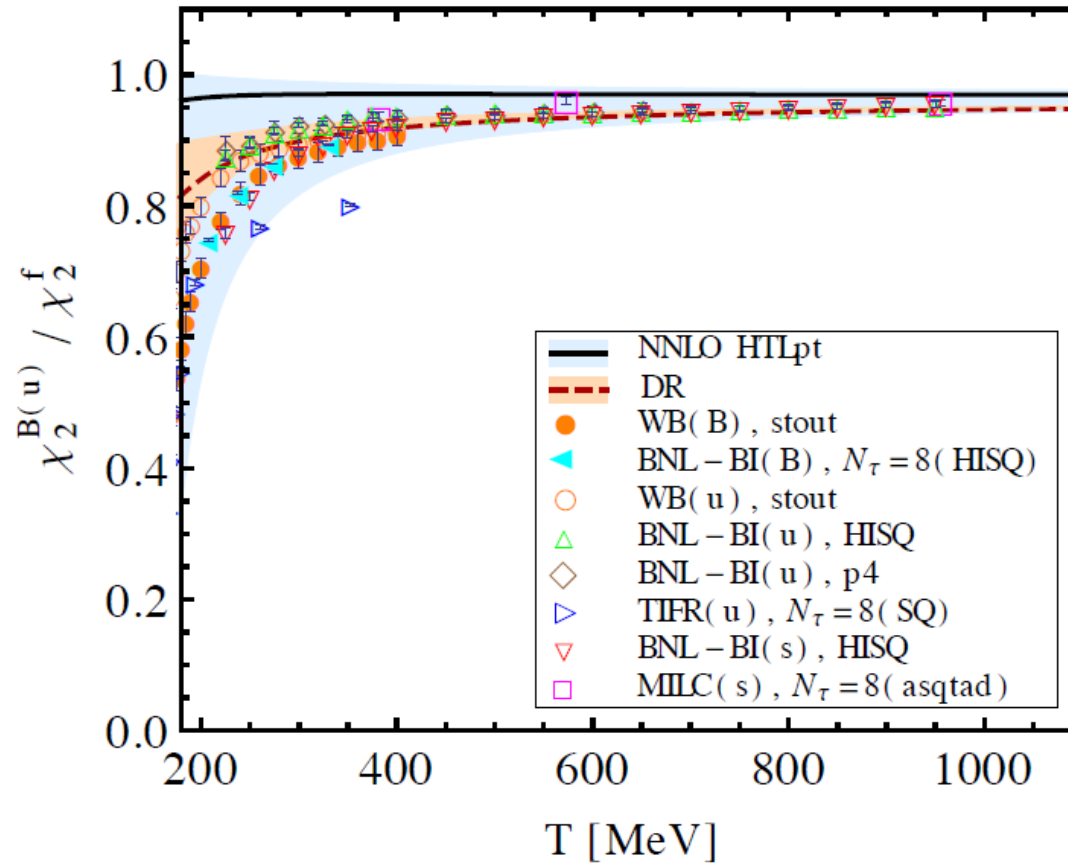


Kajantie, Laine, Rummukainen, Schröder, PRD 67 (2003)

Laine, Schröder, PRD 73 (2006)

Andersen, Leganger, Strickland, Su, JHEP 1108 (2011)

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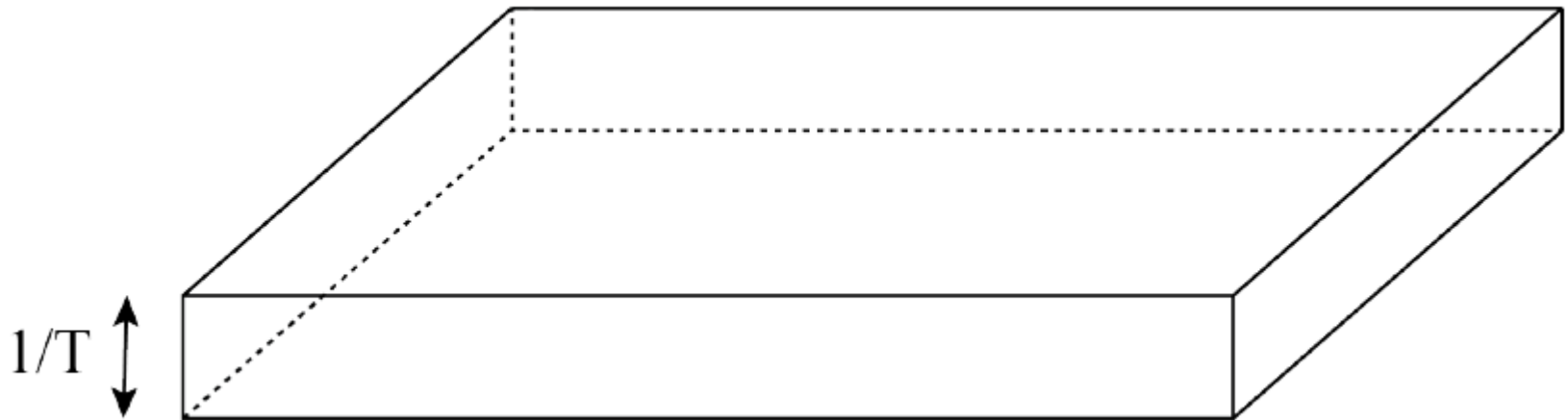


AV, PRD 67, PRD 68 (2003)

Mogliacci, Andersen, Strickland, Su, AV, JHEP 1312 (2013)

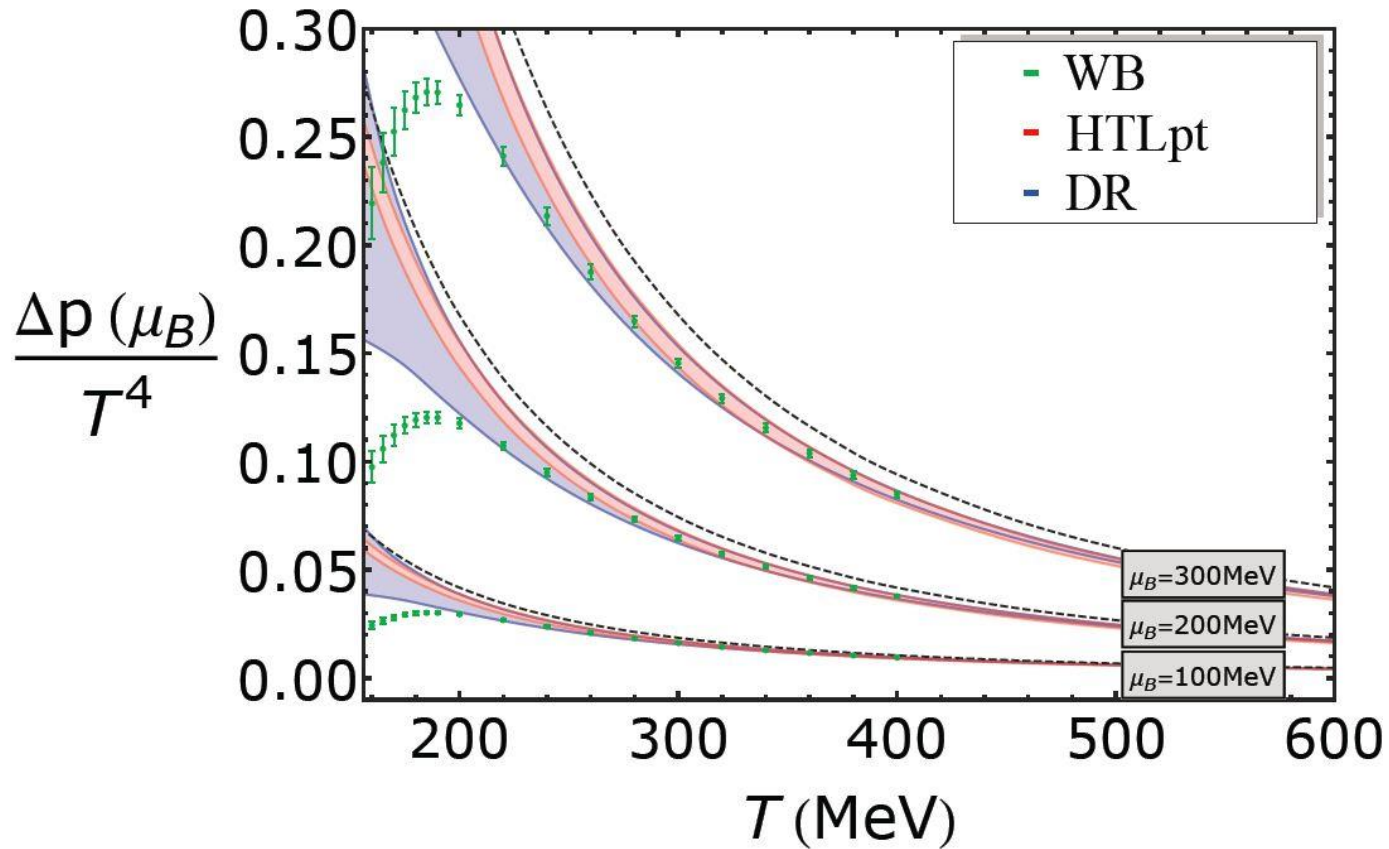
Haque, Andersen, Mustafa, Strickland, Su, PRD 89 (2014)

Key observation behind improvement:
Dimensional reduction and effective theories
HTL and EQCD \rightarrow Resummation of soft modes



$$p_{\text{QCD}}(T, \mu) = p_{\text{E}}(T, \mu) + \frac{T}{V} \ln \int \mathcal{D}A_i^a \mathcal{D}A_0^a \exp \left\{ -S_{\text{E}} \right\}$$

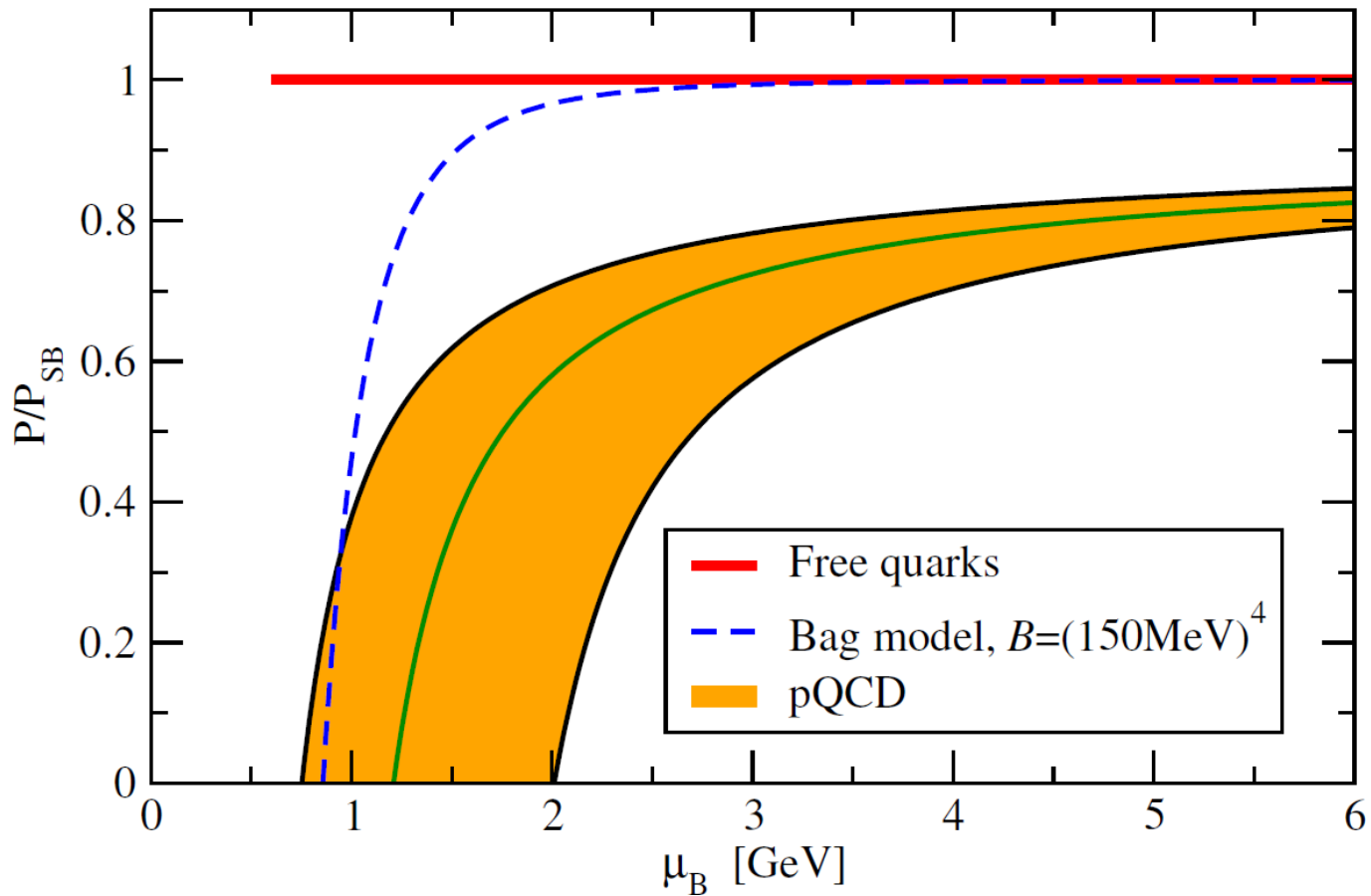
No Sign Problem \rightarrow Perturbation theory works even (better) at finite density!



AV, PRD 67, PRD 68 (2003)

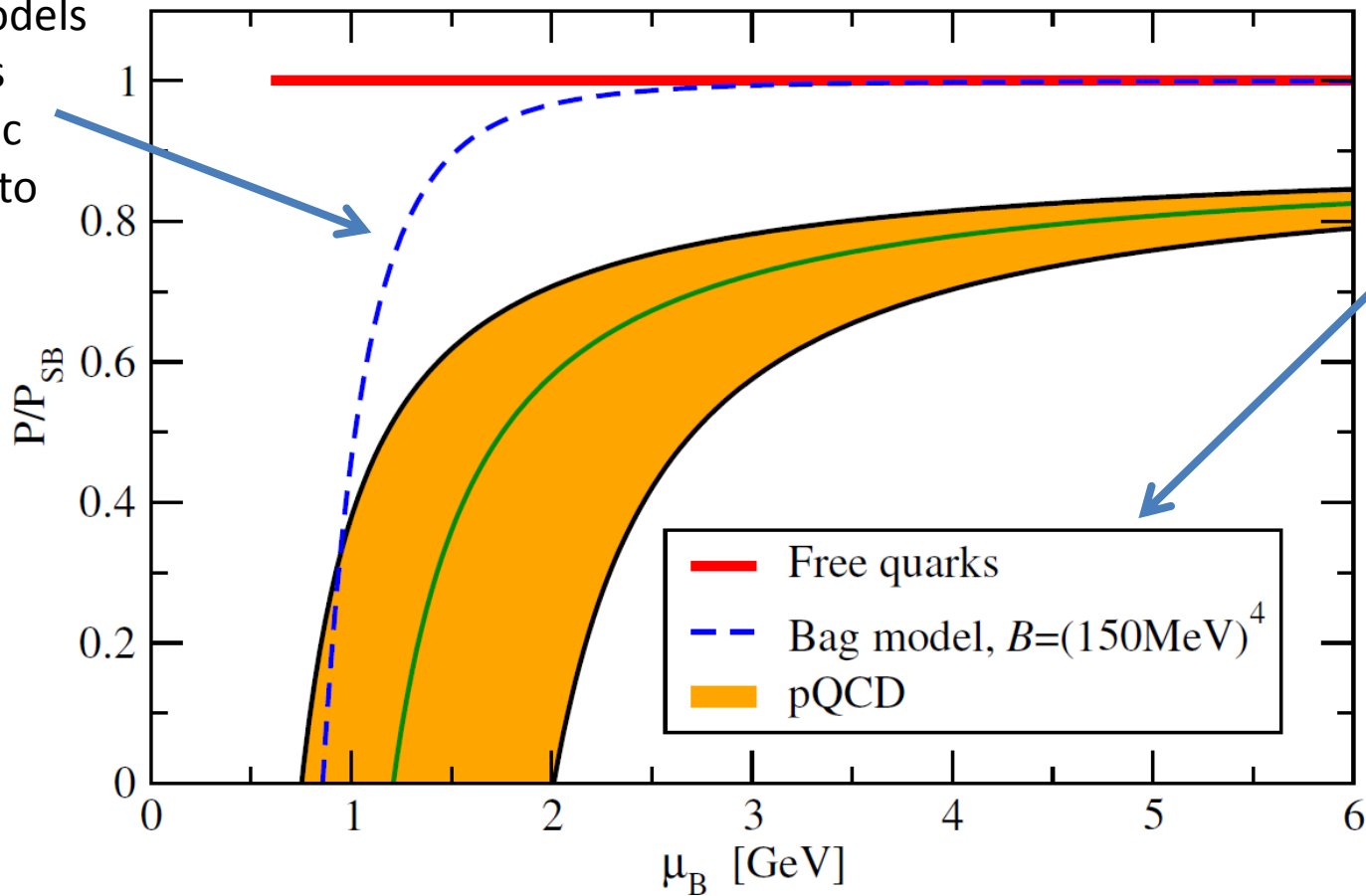
Mogliacci, Andersen, Strickland, Su, AV, JHEP 1312 (2013)

pQCD at $T = 0$: $O(g^4)$ with massive quarks

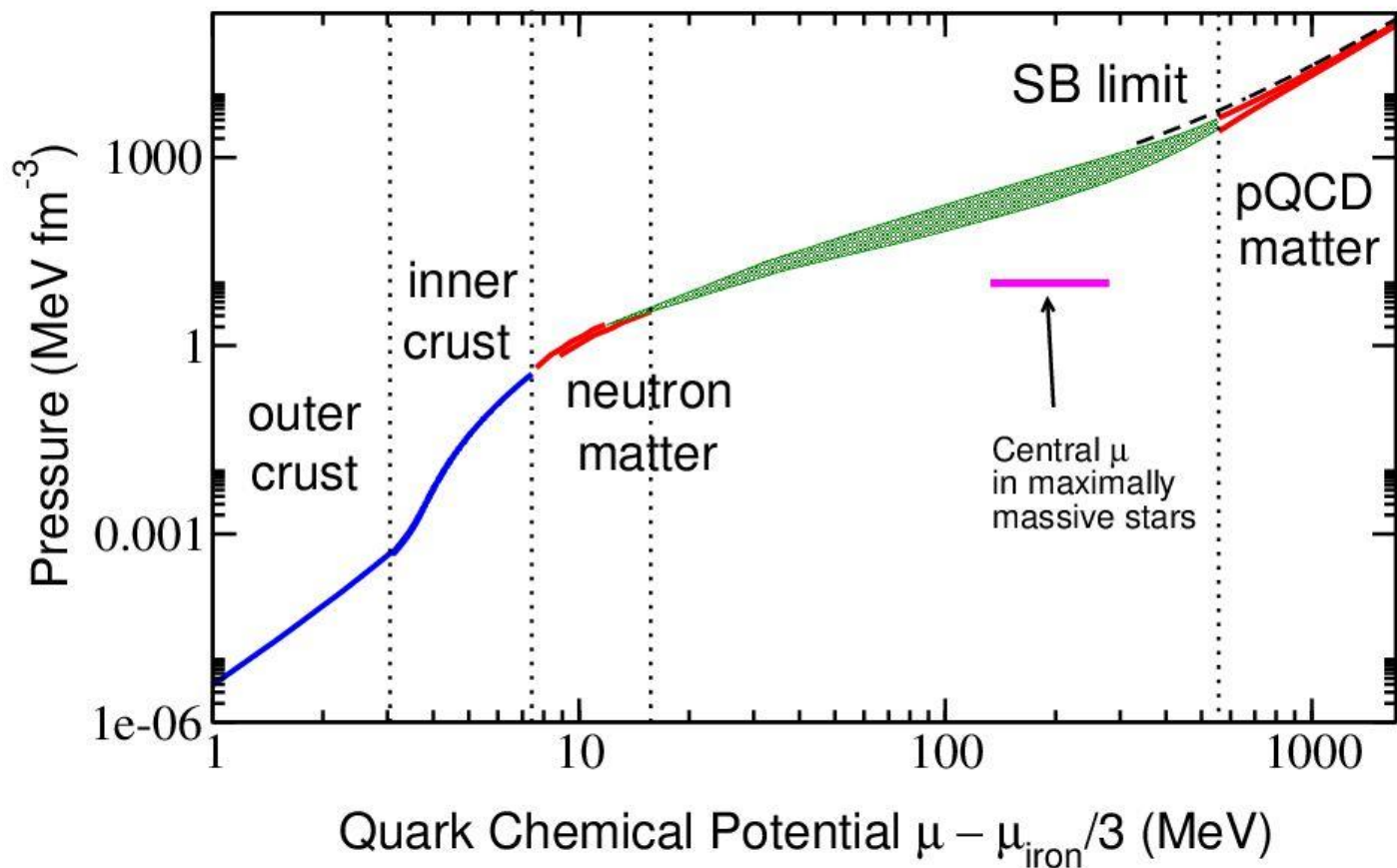
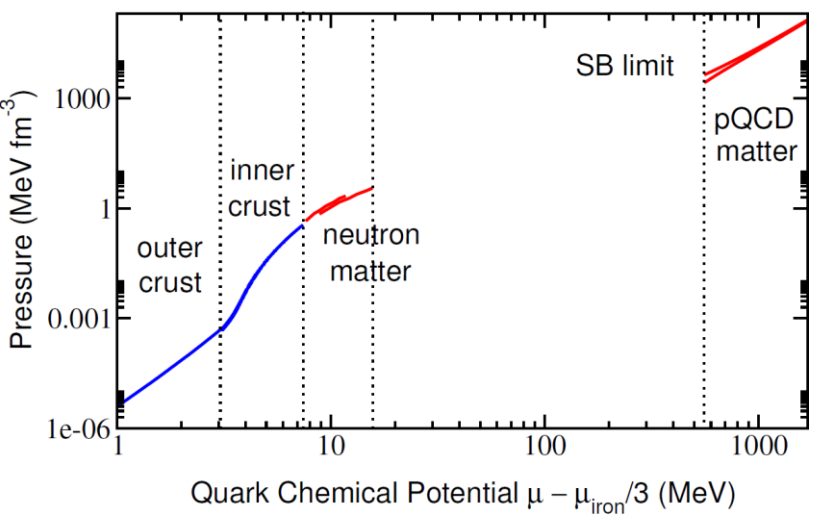


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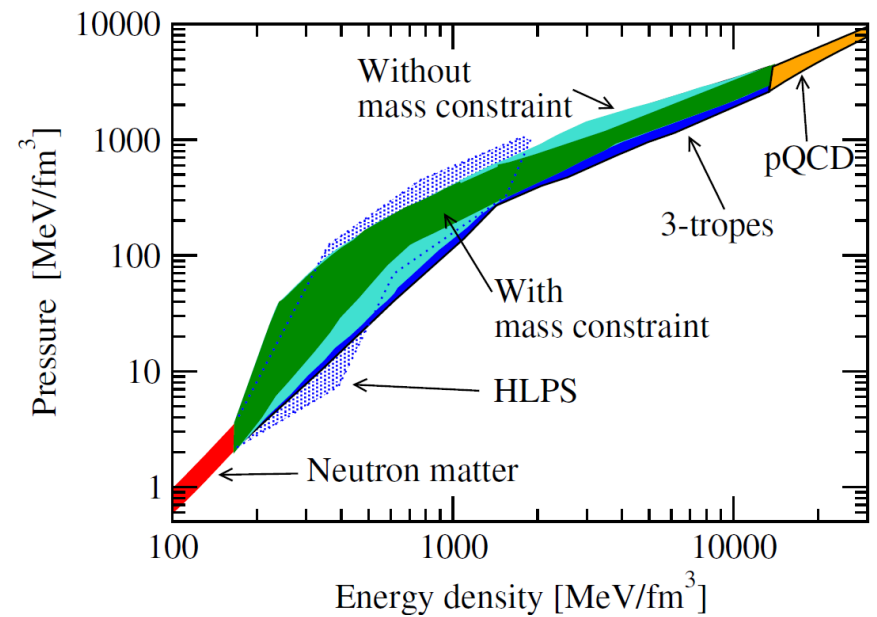
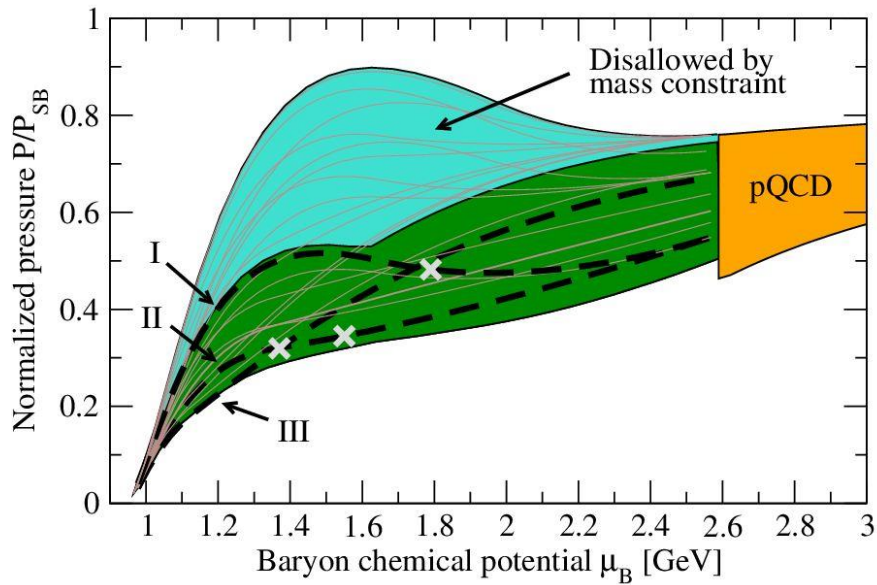
Simple models
often miss
logarithmic
approach to
free limit



No lattice
prediction!

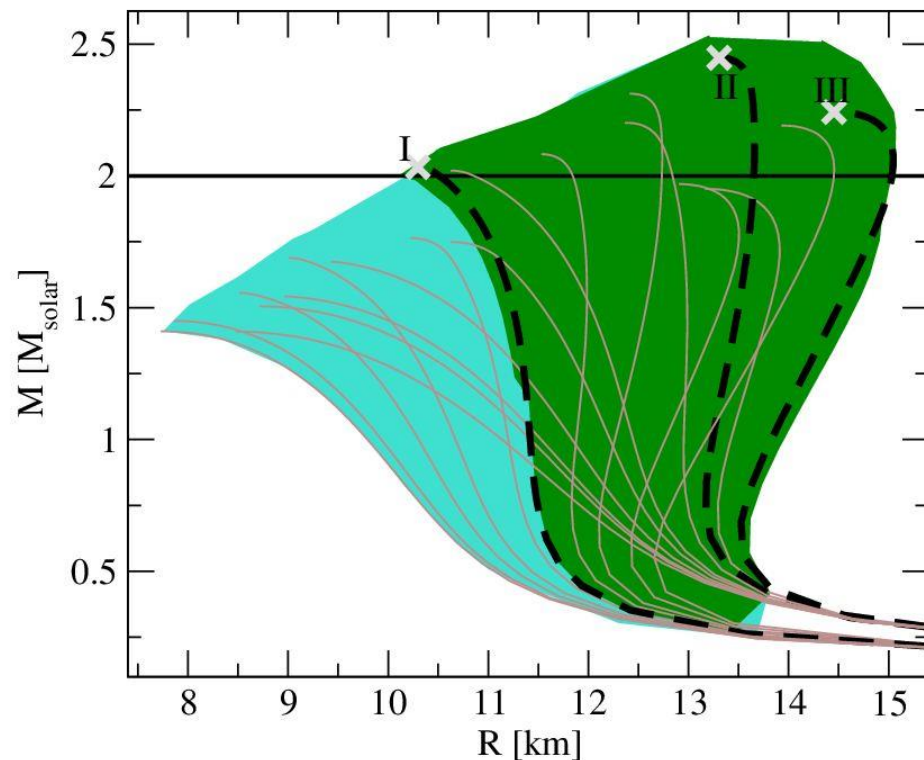


Kurkela, Fraga,
Schaffner-Bielich, AV,
APJ 789 (2014)



Main goal of future work:

- Extend EoS of $T = 0$ QM to full four-loop order \rightarrow dramatically reduced uncert.
- Further constrain neutron star EoS at moderate density
- $O(g^6 \ln g)$ already *almost* complete



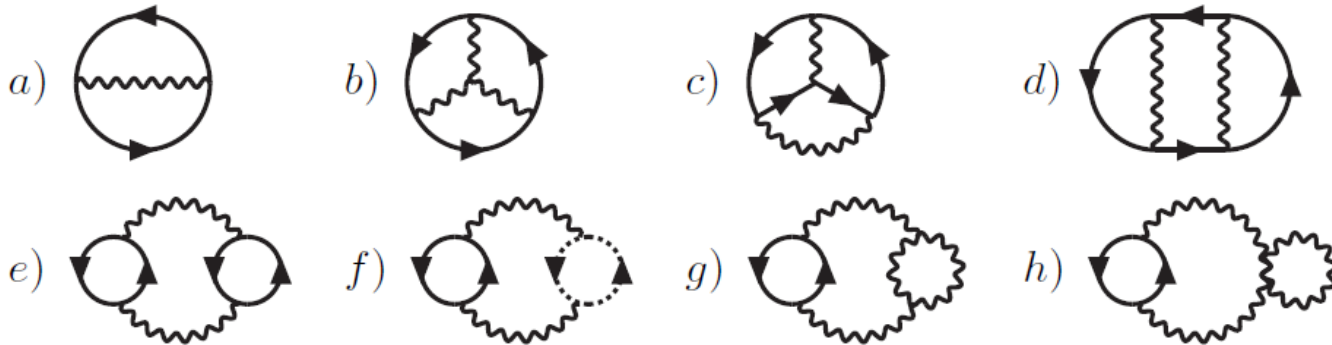
Ghisoiu, Gorda, Kurkela, Romatschke, AV,
In preparation

- I. Perturbation theory at finite density:
from quark matter to neutron stars
- II. Challenges at small but nonzero
temperature**
- III. New result: perturbative EoS of cool
quark matter

$$\Omega(T, \mu_u, \mu_d, \mu_s, m_s) = -T \log \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{D}A_\mu e^{-\int d^3x \int_0^{1/T} d\tau \mathcal{L}_{\text{QCD}}},$$

$$\mathcal{L}_{\text{QCD}} = \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \bar{\psi}_i (\gamma_\mu D_\mu + m_i - \mu_i \gamma_0) \psi_i$$

Perturbation theory: expansion of partition function in powers of gauge coupling $g \rightarrow$ Vacuum or bubble diagrams

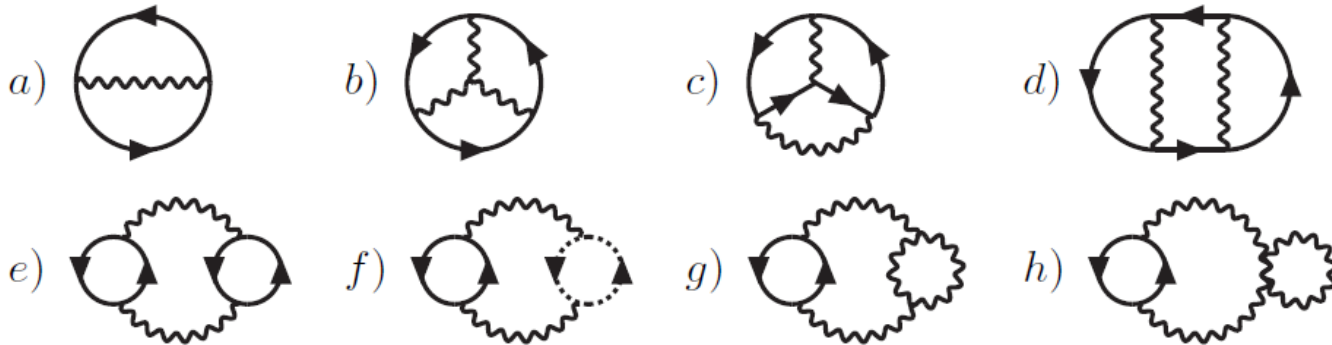


$$\begin{aligned} p_n^{\text{bos}} &= 2\pi n T, \\ p_n^{\text{ferm}} &= (2n + 1)\pi T - i\mu \int \frac{d^{4-2\epsilon} p}{(2\pi)^{4-2\epsilon}} \rightarrow T \sum_{p_n} \int \frac{d^{3-2\epsilon} p}{(2\pi)^{3-2\epsilon}} \end{aligned}$$

$$\Omega(T, \mu_u, \mu_d, \mu_s, m_s) = -T \log \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{D}A_\mu e^{-\int d^3x \int_0^{1/T} d\tau \mathcal{L}_{\text{QCD}}},$$

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Problem in pQCD: infrared divergences at three-loop order from long-range (static) gauge fields



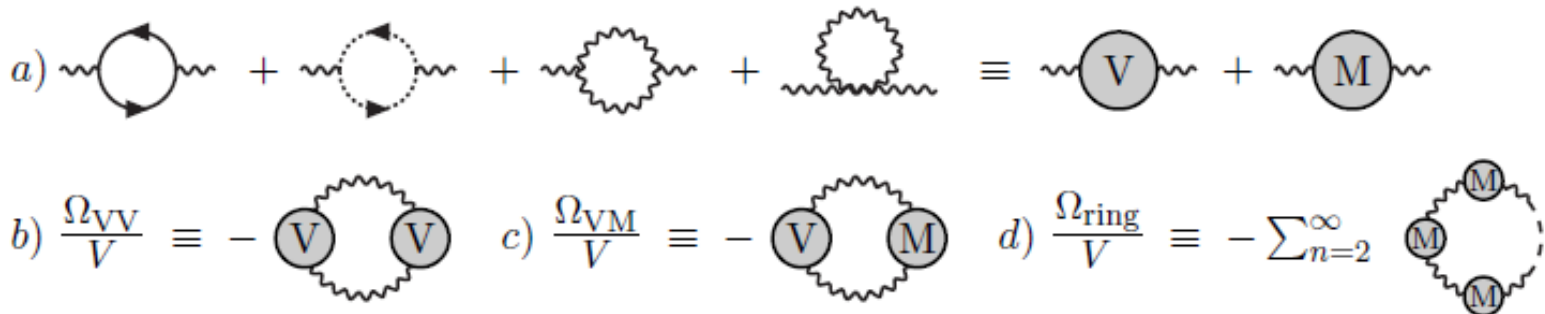
$$-\omega^2 + k^2 \rightarrow -\omega^2 + k^2 + \Pi(\omega, k)$$

$$\Omega(T, \mu_u, \mu_d, \mu_s, m_s) = -T \log \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{D}A_\mu e^{-\int d^3x \int_0^{1/T} d\tau \mathcal{L}_{\text{QCD}}},$$

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Solution: Resummation of IR sensitive contributions to the EoS:
sum (certain) diagrams to infinite order or use EFT



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Solution: Resummation of IR sensitive contributions to the EoS: sum (certain) diagrams to infinite order or **use EFT**

$$p_{\text{QCD}} = p_{\text{QCD}}^{\text{naive}} + p_{\text{DR}}^{\text{res}} - p_{\text{DR}}^{\text{naive}} + p_{\text{HTL}}^{\text{res}} - p_{\text{HTL}}^{\text{naive}}$$

Effective theory for $n = 0$
Matsubara mode. Necessary at
 $T \neq 0$; vanishes when $T \rightarrow 0$.

Effective description for $n \neq 0$
Matsubara mode. Dominates in the
 $T = 0$ limit.

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- III. **New result: perturbative EoS of cool
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Status of weak coupling EoS today:

- $\mu_B = 0$: $O(g^6 \ln g)$ Kajantie, Laine, Rummukainen, Schröder, PRD 67 (2003)
- $0 < \mu_B \leq T/g$: $O(g^6 \ln g)$ AV, PRD 68 (2003)
- $T = 0$: $O(g^4)$ /w finite masses Kurkela, Romatschke, AV, PRD 81 (2010)

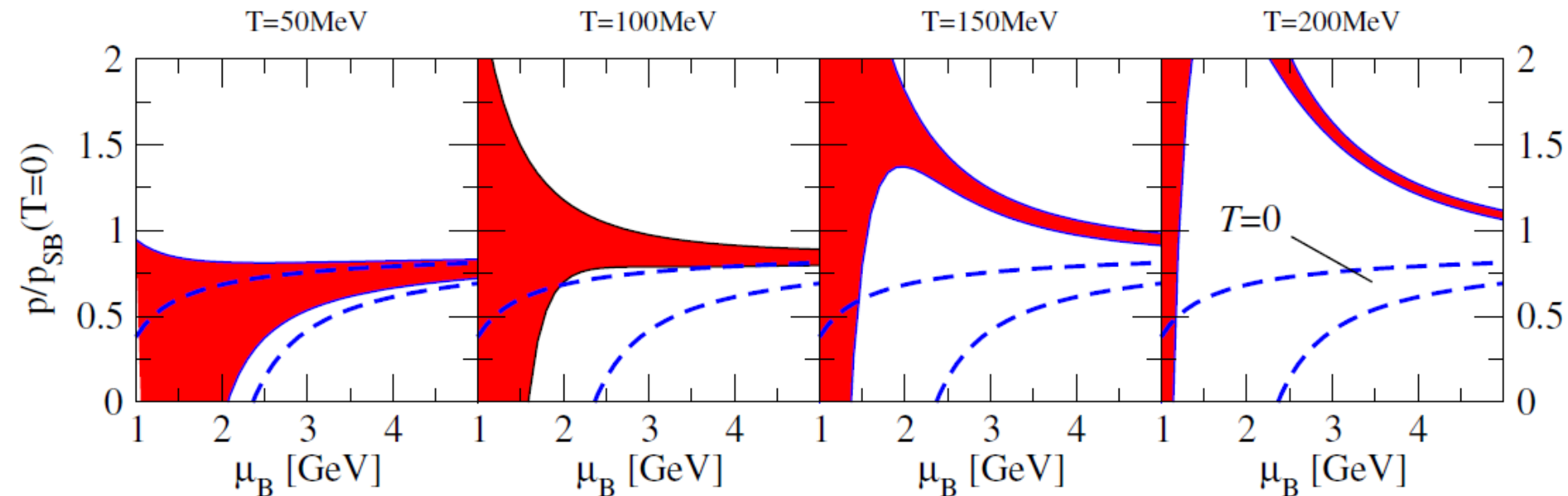
Problematic region: small but nonzero temperatures

- High order expansion in $\tau \equiv T/(g\mu_B)$ around $T = 0$ – small radius of converge; “Hard Dense Loops” Rebhan et al. 2002-2005
- Very cumbersome diagrammatic resummation connecting $T = 0$ and $T \sim \mu_B$ Ipp, Kajantie, Rebhan, AV, PRD 74 (2006)

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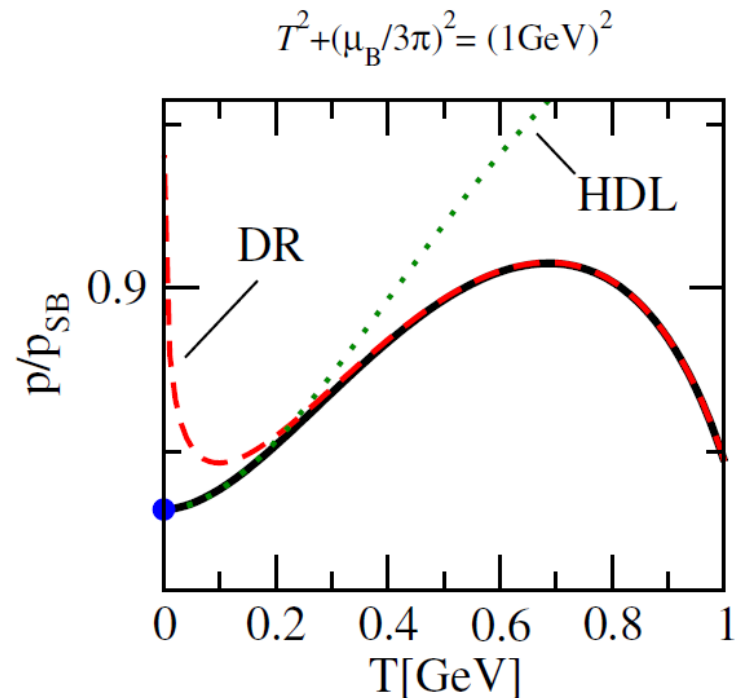
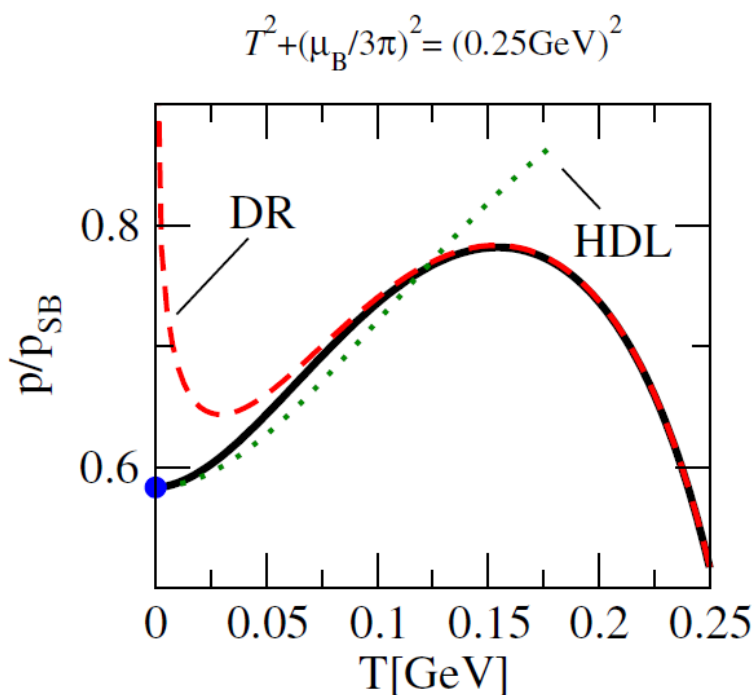
New: analytic result combining DR and HTL resummations →
 Small temperatures under control Kurkela, AV, arXiv:1603.00750 [hep-ph]



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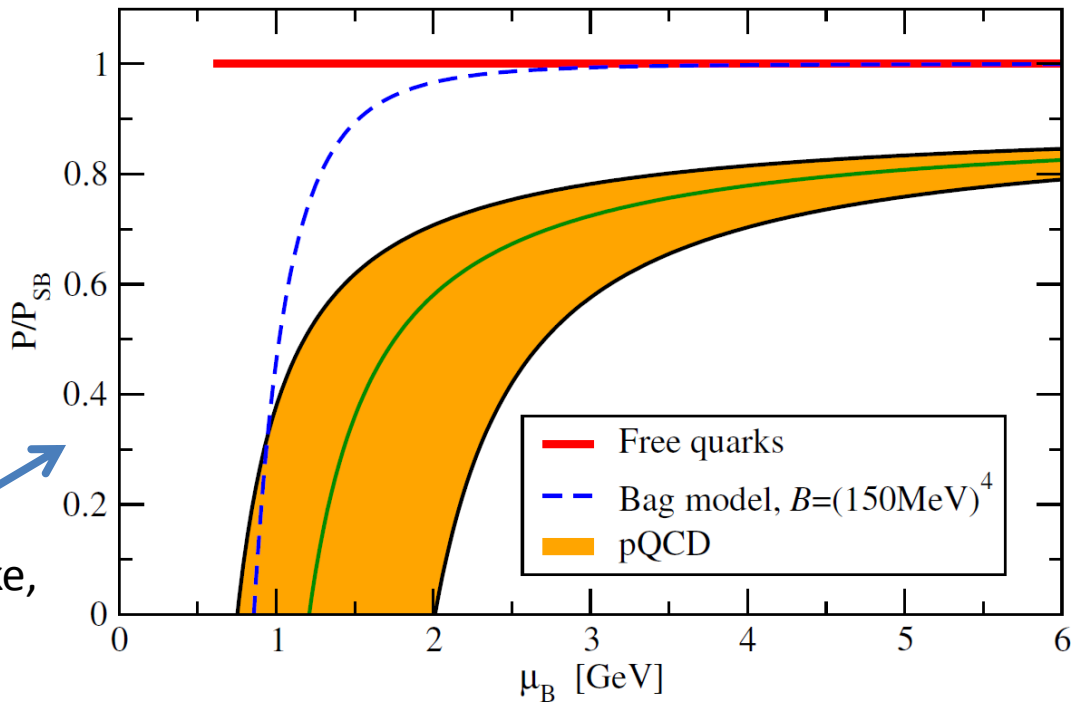
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- Analyticity: result trivially extendible outside beta equilibrium and charge neutrality
 - Leading finite- T correction: $O\left(T^2 \ln \frac{T}{\mu_B}\right)$
- Potential uses in neutron star merger calculations: need finite temperature corrections on a large density interval
 - Plan: Constrain low-density EoSs with new result

Conclusions

Kurkela, Romatschke,
AV, PRD 81 (2010)



Kurkela, AV,
arXiv:1603.00750
[hep-ph]

