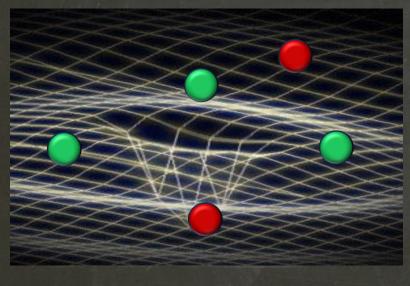
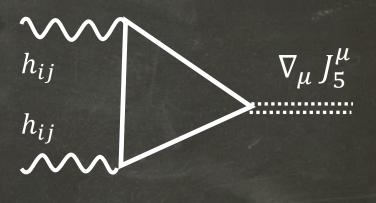
### ROYAL SOCIETY

# Gravitational ABJ Anomaly Stochastic Matter Production & Leptogenesis



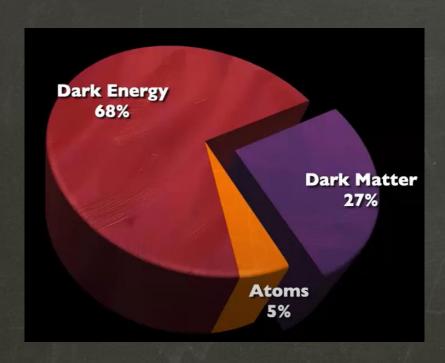
Azadeh Malek-Nejad King's College London



arXiv:2412.09490

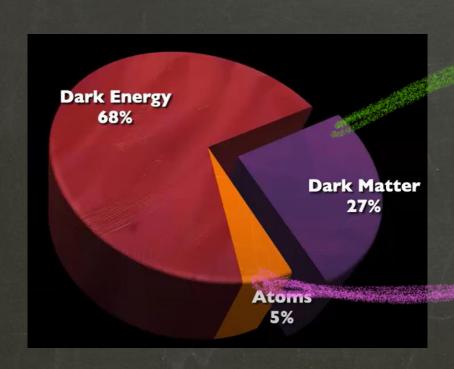
### The Story Behind this Story!

32% of the Universe today is made of Matter, visible & dark



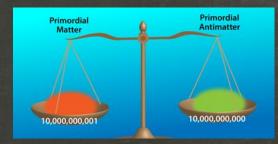
### The Story Behind this Story!

32% of the Universe today is made of Matter, visible & dark



i) Particle Nature of DMii) Its Production mechanism

Requires CP violation beyond the SM Matter – antimatter Asymmetry!

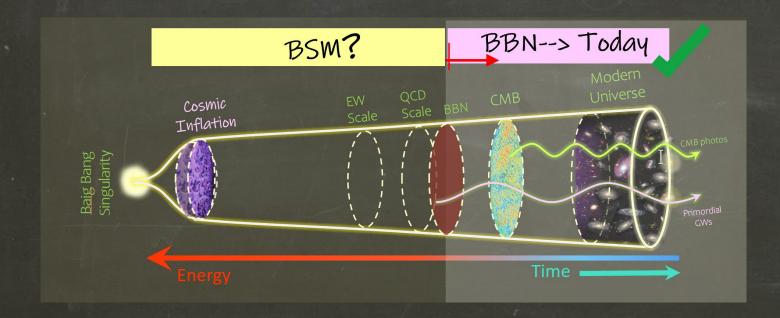


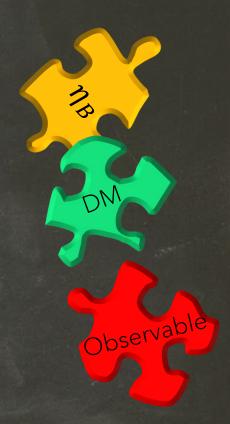
### Common Puzzles of SM & Cosmology

- I) Origin of matter asymmetry
- II) Particle nature of DM
- III) How DM has been produced?

Puzzles Which need Physics Beyond SM

We often assume New Physics Interactions
To Answer this Question!





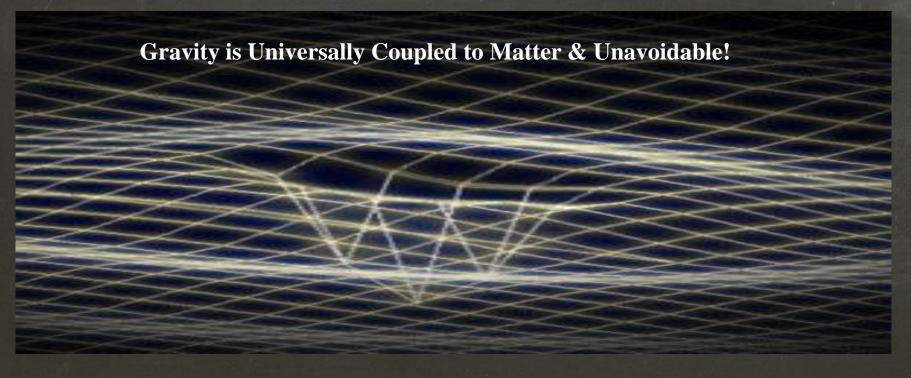
### Common Puzzles of SM & Cosmology

- I) Origin of matter asymmetry
- II) Particle nature of DM
- III) How DM has been produced?

Puzzles Which need Physics Beyond SM

We often assume New Physics Interactions
To Answer this Question!
But

What about Gravitational Fields?





# Setup

- 1) Quantum Fluctuations in Cosmology
- 2) Gravitational Particle Production
- 3) Gravitational ABJ Anomaly

arXiv:2412.09490

4) Outlook

### Quantum Fluctuations in Cosmology

 $\hbar \neq 0$ 

### Quantum Vacuum $\hbar \neq 0$

Due to Uncertainty Principle

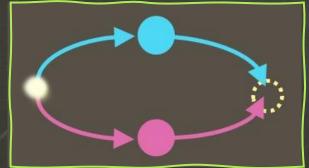
$$\Delta x \, \Delta p \geq \hbar/2$$

quantum vacuum is NOT nothing!

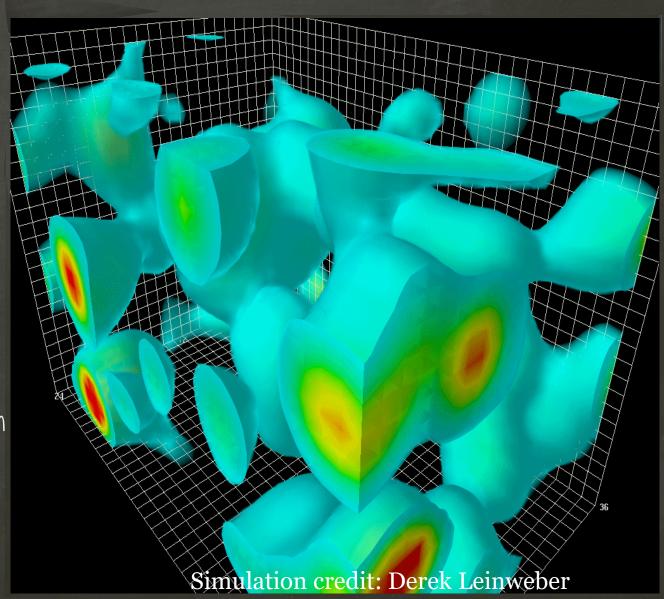
But, a vast ocean made of

#### Virtual particles

Vacuum

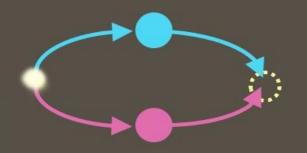


Vacuum



### Particle Production

Virtual particles



background field

Actual particles

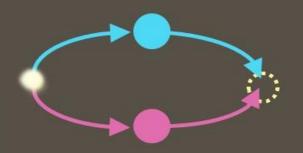
Background field can upgrade them into actual particles!

$$\langle J \rangle = 0$$

$$\langle J \rangle \neq 0$$

### Particle Production

Virtual particles



background field

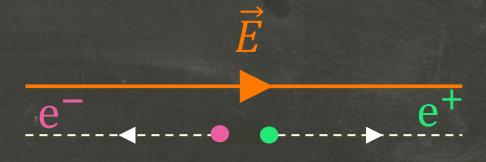
Actual particles

Background field can upgrade them into actual particles!

Examples of such BG fields:

1) Electric Field Schwinger effect

Work of the Lorentz force over Comptom wavelength  $eE \lambda_{comp} = mc^2$ 

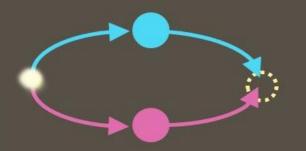


Rest energy of charged particle



#### Particle Production

Virtual particles



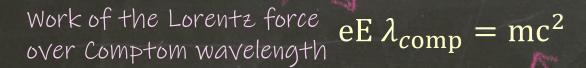
background field

**Actual particles** 

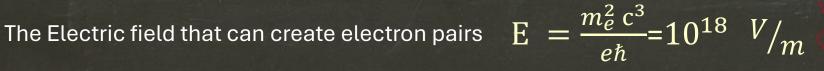
Background field can upgrade them into actual particles!

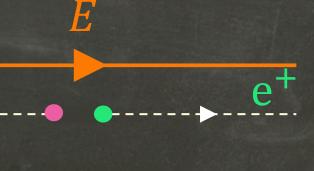
Examples of such BG fields:

1) Electric Field Schwinger effect



Rest energy of charged particle







J. Schwinger (1951)

### What about Schwinger Effect in Early Universe?

Schwinger effect in axion-inflation

#### How about Axion-inflation?!

- i) a natural candidate for the inflaton field
- ii) Naturally coupled to gauge fields



K. Lozanov

E. Komatsu

- K. Lozanov, A. M., E. Komatsu 2018
- A. M., E. Komatsu 2019
- V. Domcke, Y. Ema, K. Mukaida, R. Sato 2019
- L. Mirzagholi, A. M., K. Lozanov 2019



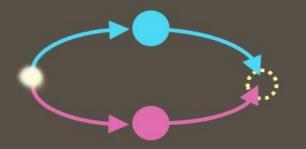
E. Komatsu **2022** nature reviews physics

New physics from the polarized light of the cosmic microwave background



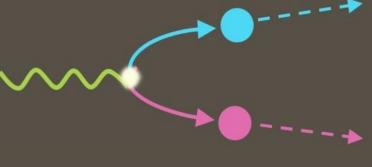
### Particle Production

Virtual particles



background field

Actual particles



Background field can upgrade them into actual particles!

Examples of such BG fields:

2) Gravitational

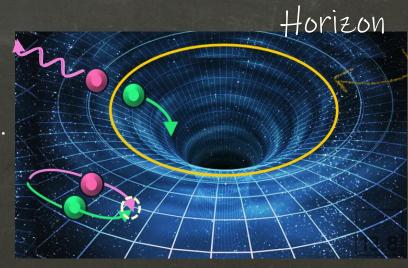
Hawking radiation

one particle fall into the BH, while the other escapes...



Power BH emitted is

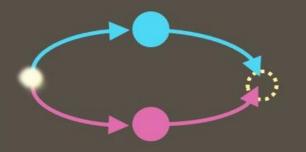
$$P = \frac{\pi c^3 M_{pl}^4}{240} \; \frac{1}{M^2}$$



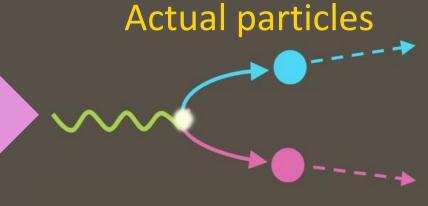
S. Hawking (1974)

### Particle Production

Virtual particles



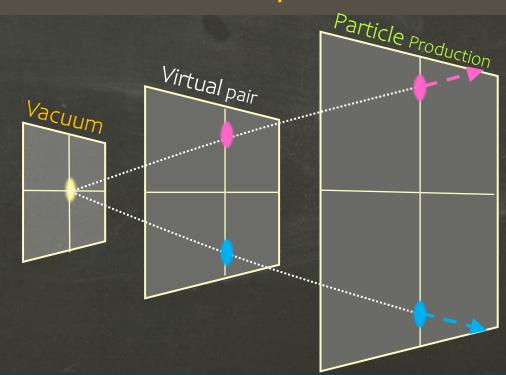
background field



Background field can upgrade them into actual particles!

Examples of such BG fields:

- 1) Electric Field Schwinger effect
- 2) Gravitational
- i) Hawking radiation
- ii) expansion of the Universe!



### Expanding Universe Produces Particles!

Flat Space:

Space

Vacuum Virtual pair Vacuum Time

Vacuum

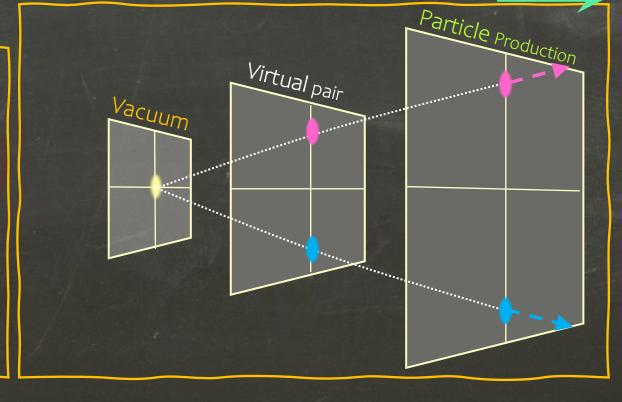
### Expanding space:





E. Schrödinger (1939)

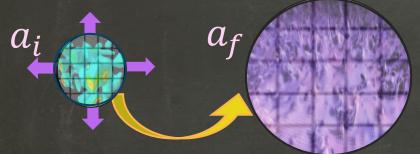
Shocked by his discovery, Schrödinger found it an alarming Phenomenon!



Particle Production

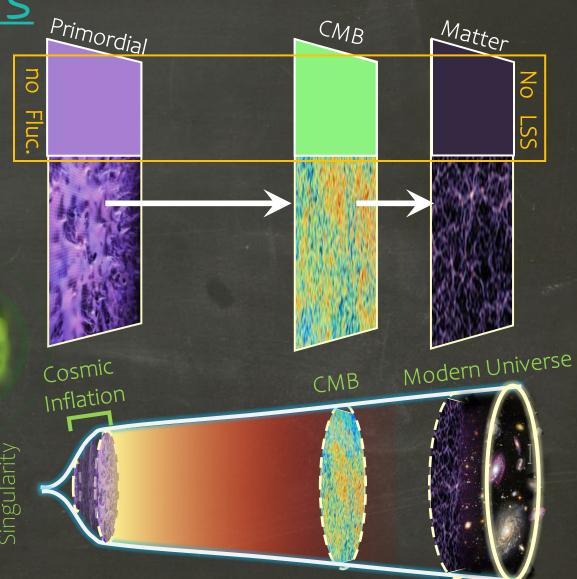
# Cosmic Perturbations Primordial

Exponential expansion turns initial quantum vacuum fluctuations into



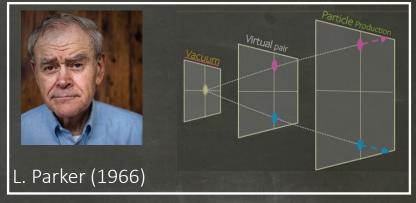
actual cosmic perturbations!

We are the product of quantum fluctuations in the very early universe!



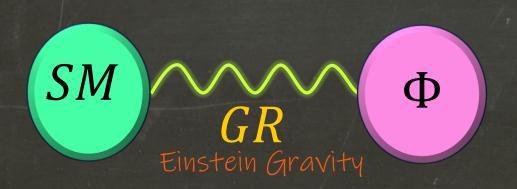
### Cosmological Gravitational Particle Production





The expansion of the Universe creates pair production in FRW geometry.

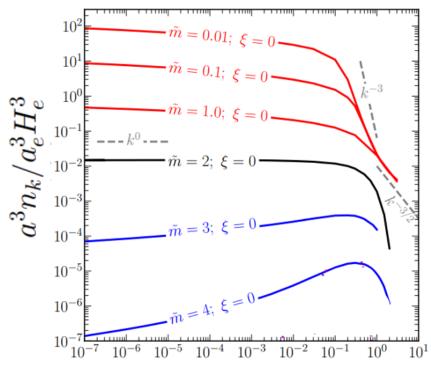
But conformal fields in 4d will not be produced since FRW is conformally flat!



### Scalar Field in Expanding Universe

#### Consider a scalar field

In cosmological background



 $k/a_eH_e$ 

$$\mathcal{L} = \frac{1}{2} g^{\mu\nu} \nabla_{\mu} \Phi \nabla_{\nu} \Phi - \frac{1}{2} m^2 \Phi^2$$

$$g_{\mu\nu} = a^2(\tau) \text{diag}(-1, 1, 1, 1)$$

$$\Phi_k'' + \omega_k^2(\tau) \Phi_k = 0$$

$$\omega_k^2(\tau) = k^2 + a^2(\tau)(m^2 + \frac{R(\tau)}{6})$$

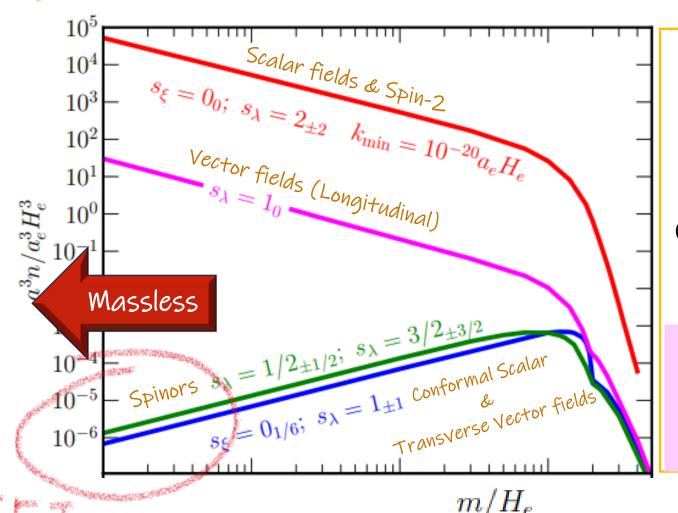
An example of Cosmological Gravitational Particle Production (CGPP)

Plot credit: Kolb & Long Reviews of Modern Physics 2023

 $H_e$  is Hubble scale at the end of inflation



### Spinning Fields in Expanding Universe





L. Parker (1966)

Conformal fields in FRW will not be produced since FRW is conformally flat!

No CGPP for massless fermions! (conformal symmetry)

Plot credit: Kolb & Long Reviews of Modern Physics 2023

### Fermions in Expanding Universe

Consider spin 
$$\frac{1}{2}$$
 massless fermions  $\mathcal{L}_{\psi}=i\psi^{\dagger}~\gamma^{\mu}\mathcal{D}_{\mu}\psi$   
Spinor covariant derivative  $\mathcal{D}_{\mu}^{\dagger}=\nabla_{\mu}-\omega_{\mu}$  Spin connection

In cosmological background

$$g_{\mu\nu} = a^2(\tau) \text{diag}(-1, 1, 1, 1)$$

The field equation of fermion is

$$\left(\gamma^0(\partial_0 + \frac{3}{2}H) + \frac{1}{a}\gamma^i\partial_i\right)\Psi = 0.$$

Effect of gravity

The effect of FRW gravity (conformally flat geometry) can be absorbed as

$$\Psi \equiv a^{3/2} \Psi, \qquad \qquad \left( \gamma^0 \left[ \partial_0 + \frac{1}{a} \gamma^i \partial_i \right] \Psi \right) = 0.$$

canonically renormalized field lives in flat space!

### How to Create Fermions in Expanding Universe?

Breaking the conformal symmetry of Weyl fermions by interactions, e.g.

(dilatation transformation)

Couple your Weyl fermion with
 Inflaton field,
 Standard Model,

Dark sector coupled to thermal bath

o make the fermion massive to produce them gravitationally! (CGPP)

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Is that the best Gravity can do to produce fermions!?

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Breaking the conformal symmetry of Weyl fermions by interactions, e.g.

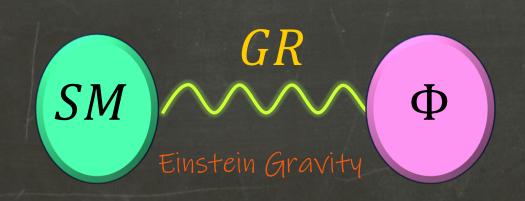
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Is that the best Gravity can do to produce fermions!?  $\sqrt{6}$ 

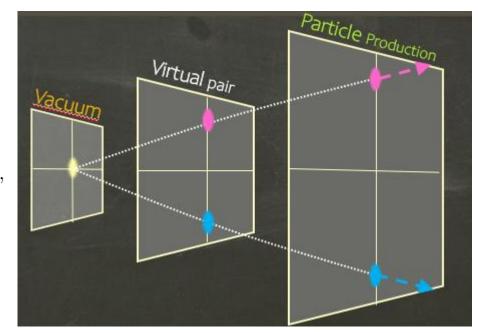


| Production Mechanism       | Underlying Physics | Conditions           |
|----------------------------|--------------------|----------------------|
| Cosmological Gravitational | Cosmic expansion   | super-massive fields |
| Particle Production (CGPP) | )                  | $M > 10^{13} \; GeV$ |

Kolb & Long 2017

#### Relic density by CGPP

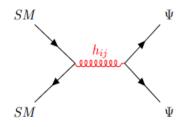
$$\frac{
ho_{\Psi,0}^{\text{CGPP}}}{
ho_{\text{DM},0}} \sim 7 \times \left(\frac{M}{10^{11} \text{GeV}}\right)^2 \left(\frac{T_{\text{reh}}}{10^9 \text{GeV}}\right),$$



|    | Production Mechanism       | Underlying Physics | Conditions           |
|----|----------------------------|--------------------|----------------------|
| D  | Cosmological Gravitational | Cosmic expansion   | super-massive fields |
| 1) | Particle Production (CGPP) |                    | $M > 10^{13} \; GeV$ |
|    |                            |                    |                      |

Kolb & Long 2017

II) Graviton-Mediated
Annihilation (GMA)



Super-massive field High temperature plasma  $T_{reh} > 10^{13} \; GeV$  M. Garny, et al 2016 Bernal, et. al. 2018 Clery et. al. 2022

Relic density by GMA

$$\frac{
ho_{\Psi,0}^{\rm GMA}}{
ho_{{
m DM},0}} \sim 5 \left(\frac{M}{10^{12} {
m GeV}}\right) \left(\frac{T_{
m reh}}{10^{13} {
m GeV}}\right)^3,$$

|      | Production Mechanism                    | Underlying Physics        | Conditions   |  |
|------|---|---------------------------|--|--|
| I)   | Cosmological Gravitational              | Cosmic expansion          | super-massive fields   | Kolb & Long 2017   |
|      | Particle Production (CGPP)              |                           | $M > 10^{13} \; GeV$   |  |
| II)  | Graviton-Mediated<br>Annihilation (GMA) | $SM$ $h_{ij}$ $SM$ $\Psi$ | Super-massive field High temperature plasma $T_{reh} > 10^{13} \; GeV$ | M. Garny, et al 2016<br>Bernal, et. al. 2018<br>Clery et. al. 2022 |
| III) | Gravitational Leptogenesis              | $ abla_{\mu}J_{5}^{\mu}$  | Parity violation $h_L \neq h_R$<br>Chiral GWs<br>Chiral fermions       | Alexander et. al. 2006  A.M. 2014 & 2016  A.M. 2024                |

Chiral fermions

(global) Gravitational anomaly

$$\nabla_{\mu}J_5^{\mu} = \frac{N_L - N_R}{16\pi^2}R\tilde{R}$$

|      | Production Mechanism                 | Underlying Physics        | Conditions   |  |
|------|--------------------------------------|---------------------------|--|--|
| 1)   | Cosmological Gravitational           | Cosmic expansion          | super-massive fields   | Kolb & Long 2017   |
| 1)   | Particle Production (CGPP)           | )                         | $M > 10^{13} \; GeV$   |  |
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**A.M.** 2024

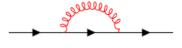
 $h_L \neq h_R$ 

What does Unpolarized Gravitational Waves do!?



|      | Production Mechanism                    | Underlying Physics        | Conditions   |  |
|------|---|---------------------------|--|--|
| I)   | Cosmological Gravitational              | Cosmic expansion          | super-massive fields   | Kolb & Long 2017   |
|      | Particle Production (CGPP)              |                           | $M > 10^{13} \; GeV$   |  |
| II)  | Graviton-Mediated<br>Annihilation (GMA) | $SM$ $h_{ij}$ $SM$ $\Psi$ | Super-massive field High temperature plasma $T_{reh} > 10^{13} \; GeV$ | M. Garny, et al 2016<br>Bernal, et. al. 2018<br>Clery et. al. 2022 |
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IV) GW-Induced Freeze-In

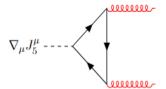


GWs Background

**A.M.** & Kopp 2024

### Matter Asymmetry by Gravitational Anomaly: $\langle R\tilde{R} \rangle \neq 0!$

#### What makes Chiral Gravitational Waves?



Left-handed GW

To generate circularly polarized GWs, we need Parity violation in inflation. Two possible models are

- 1) Chern-Simons Gravity  $\mathcal{L}_{eff} = \frac{1}{\Lambda} \varphi R \tilde{R}$ Alexander, Peskin, Sheikh-Jabbari 2006
- Non-Abelian Gauge fields in axion-inflation
   A.M., Noorbala, Sheikh-Jabbari 2012
   A.M. 2014 & 2016
   Caldwell, Devulder 2017
  - Adshead, Long, Sfakianakis 2017 Alexander, McDonough, Spergel 2018 Kamada, Kume, Yamada, Yokoyama 2019

 $\mathcal{L}_{eff} = \frac{1}{\Lambda} \varphi F \tilde{F}$  (Chiral Gauge Field Chiral GWs)

Axion-inflation is a generic setting for leptogenesis (All the Sakharov conditions are satisfied)

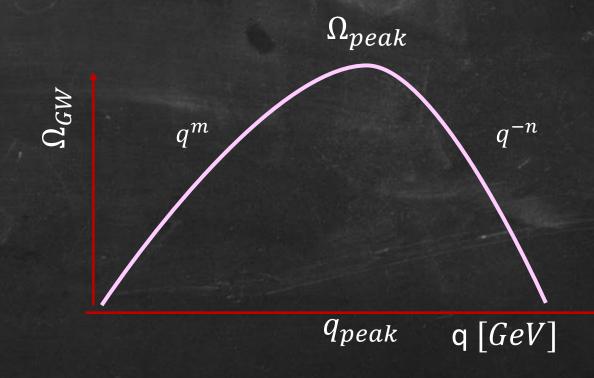
**A.M.** 2014

3) U(1) Gauge fields in axion-inflation

### Phenomenological Model for $\Omega_{GW}$

In radiation era

Broken Power-law Spectrum



### Phenomenological Model for $\Omega_{GW}$

$$\Omega_{\text{gw},0}(q) \approx \begin{cases} \Omega_{\text{p}} \left(\frac{q}{q_{\text{p}}}\right)^{m} & q_{\text{min}} < q < q_{\text{p}}, \\ \Omega_{\text{p}} \left(\frac{q}{q_{\text{p}}}\right)^{-n} & q_{\text{p}} < q < q_{\text{max}}, \end{cases}$$

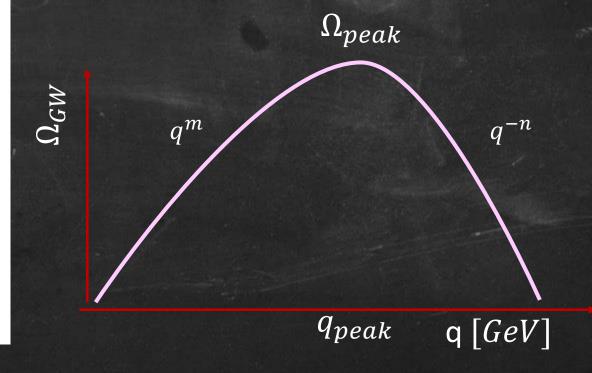
$$\Omega_{\text{gw},0}^{\pm}(q) = \frac{1 \pm \chi(q)}{2} \Omega_{\text{gw},0}(q),$$

$$h_{s,\mathbf{q}}(\tau) = a^{-1}(\tau) \mathcal{T}_s(\tau,q) e^{-iq\tau} h_{s,\mathbf{q},0},$$

$$\mathcal{T}_s(\tau, q) \approx \left(1 - e^{-\pi \beta_s(\tau - \tau_{\rm in})}\right)$$

$$\beta_{\pm} = \beta(1 \pm b_{\chi}),$$

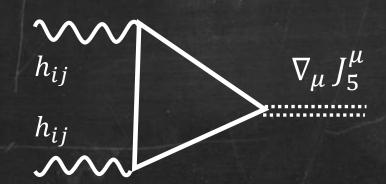
Broken Power-law Spectrum



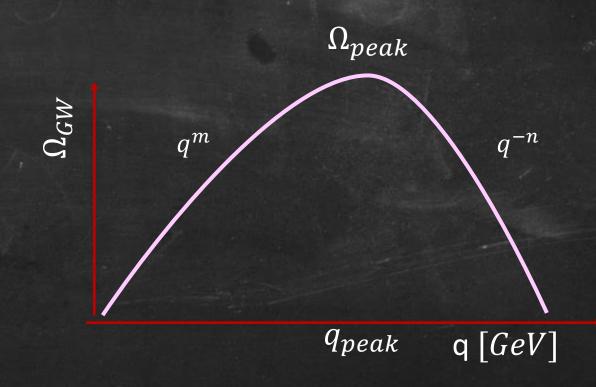
### Phenomenological Model for $\Omega_{GW}$

Gravitational ABJ Anomaly

$$\nabla_{\mu}J_{A}^{\mu} = \frac{N_{\chi}}{24(4\pi)^{2}} \langle R\tilde{R} \rangle,$$

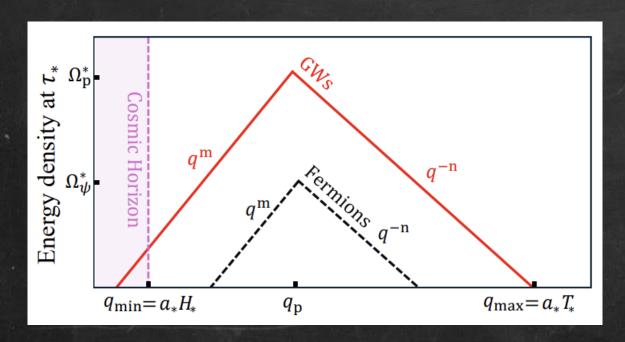


Broken Power-law Spectrum



### Fermion Energy Density

$$\nabla_{\mu}J_{A}^{\mu} = \frac{N_{\chi}}{24(4\pi)^{2}} \langle R\tilde{R} \rangle,$$



### GWs spectral energy density

$$\Omega_{\mathrm{gw},0}(q) \approx \begin{cases} \Omega_{\mathrm{p}} \left(\frac{q}{q_{\mathrm{p}}}\right)^{m} & q_{\mathrm{min}} < q < q_{\mathrm{p}}, \\ \Omega_{\mathrm{p}} \left(\frac{q}{q_{\mathrm{p}}}\right)^{-n} & q_{\mathrm{p}} < q < q_{\mathrm{max}}, \end{cases}$$

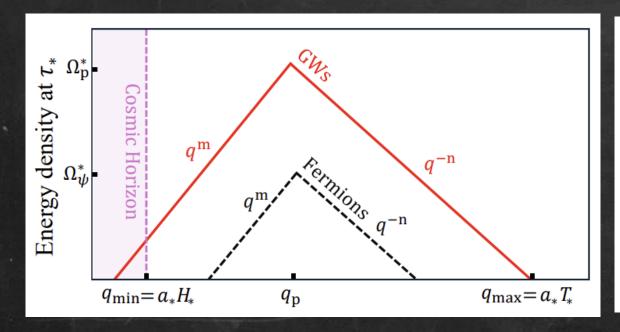
#### Fermions number density

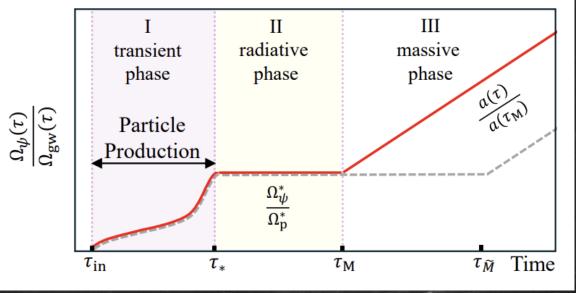
$$n_A(q) \approx \begin{cases} \mathcal{N}_{\mathrm{p}}(q) \, q_{\mathrm{p}}^3 \left(\frac{q}{q_{\mathrm{p}}}\right)^{m-1} & q_{\mathrm{min}} < q < q_{\mathrm{p}}, \\ \mathcal{N}_{\mathrm{p}}(q) \, q_{\mathrm{p}}^3 \left(\frac{q}{q_{\mathrm{p}}}\right)^{-n-1} & q_{\mathrm{p}} < q < q_{\mathrm{max}}, \end{cases}$$

$$\mathcal{N}_{p}(q) = -\frac{\Omega_{p}}{16} \left( \frac{H_{0} \beta}{a_{\text{in}} q_{p}^{2}} \right)^{2} \left[ 2b_{\chi} + (1 + b_{\chi}^{2})\chi(q) \right].$$

### Fermion Energy Density

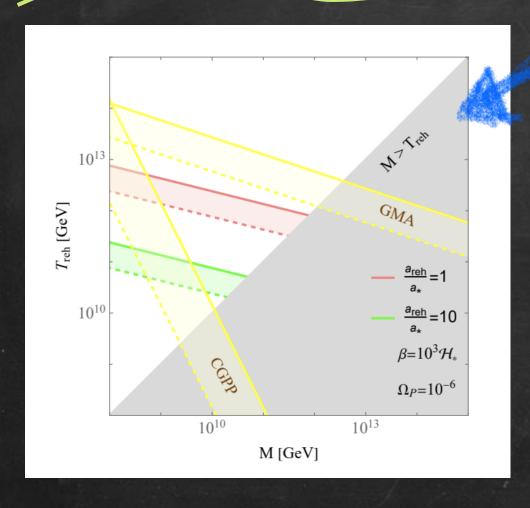
$$\nabla_{\mu}J_{A}^{\mu} = \frac{N_{\chi}}{24(4\pi)^{2}} \langle R\tilde{R} \rangle,$$





### DM and Leptogenesis

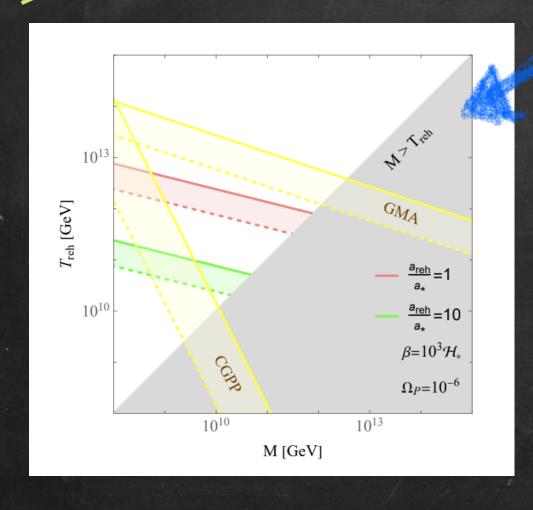
$$\nabla_{\mu}J_{A}^{\mu} = \frac{N_{\chi}}{24(4\pi)^{2}} \langle R\tilde{R} \rangle,$$



Gravitational ABJ Dark Matter

### DM and Leptogenesis

$$\nabla_{\mu}J_{A}^{\mu} = \frac{N_{\chi}}{24(4\pi)^{2}} \langle R\tilde{R} \rangle,$$



Gravitational ABJ Dark Matter

Gravitational ABJ Leptogenesis

$$T_{\rm reh} \sim 10^{14} {\rm GeV}$$

$$\Omega_{\rm p} \sim 10^{-6}$$

## Summary

Gravity and Quantum Effects in Cosmology can still surprise us: We discussed an effect that is zero at tree level and non-zero at 1-loop in cosmic perturbations!

Cosmic Perturbations (like GWs) naturally break the conformal symmetry of Weyl Fermions in Cosmology

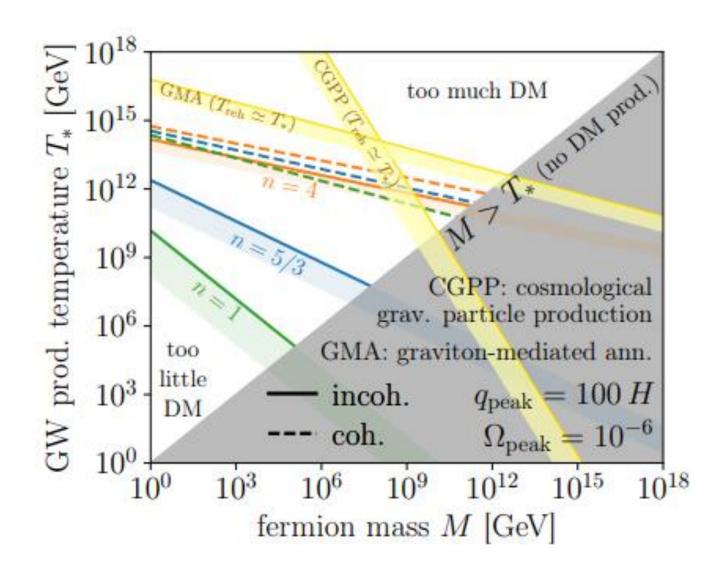
It leads to a new mechanism for dark matter production and baryogenesis in early universe.

# Questions?!



### Parameter space of GW-induced freeze-in of fermion

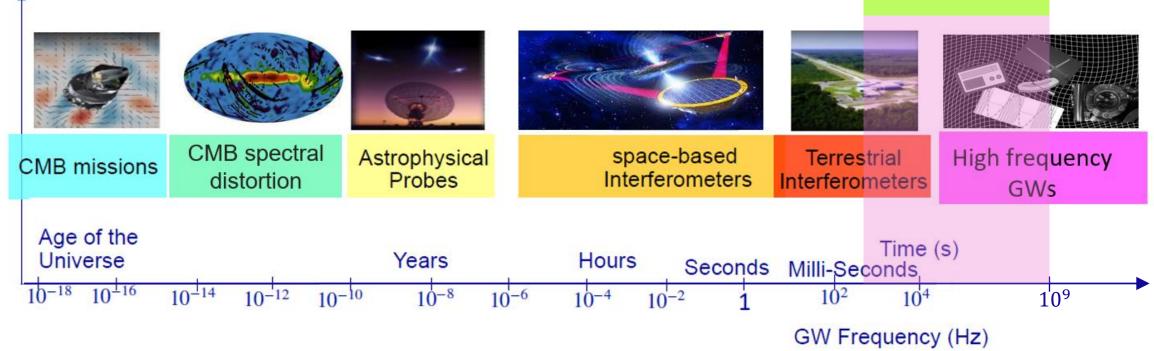
**A.M.** & Kopp 2024



### Gravitational Waves Spectrum

GW-induced freeze-in mechanism requires a GWs spectrum with peak frequency

 $f_{peak} \in$  (kHz-GHz)



Age of the Universe = Billions of Years

### Scalar number density for minimal coupling

$$\frac{a^3 n_k}{a_e^3 H_e^3} \approx \begin{cases} \frac{1}{8\pi^2} \tilde{m}^{-1} \tilde{k}^0 & 0 < \tilde{k} < \tilde{m}^{1/3} \\ \frac{1}{8\pi^2} \tilde{k}^{-3} & \tilde{m}^{1/3} < \tilde{k} < \frac{m_\varphi \kappa}{H_e} \\ C\tilde{k}^{-3/2} & \frac{m_\varphi \kappa}{H_e} < \tilde{k} < \tilde{a}_{\mathrm{RH}} \frac{m_\varphi \kappa}{H_e} \end{cases},$$