

# Lattice QCD and BSM: an update

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## Introduction to the Lattice approach

### Updates from QCD

- >  $g - 2$
- > Scattering & Decays

### Updates from BSM

- > Composite Higgs
- > Early Universe Phase transitions

# Introduction

recipe for a lattice calculation

## Ingredients...

- > A euclidean 4-dimensional space-time lattice with spacing  $a$  and volume  $V$ ,
- > Discretized degrees of freedom  $\bar{\psi}(x)$ ,  $\psi(x)$ ,  $U_\mu(x)$ ,
- > An action  $S_E = S_g[U] + S_f[\psi, \bar{\psi}, U]$

## ...and the recipe

- > for any observable  $O(\psi, \bar{\psi}, U)$

$$\langle O \rangle = \frac{1}{Z} \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}U O[\psi, \bar{\psi}, U] e^{-S_E[\psi, \bar{\psi}, U]},$$

with

$$Z = \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}U e^{-S_E[\psi, \bar{\psi}, U]}$$

## FAQ:

- > How do we define  $U_\mu(x)$ ?

$$U_\mu(x) = \exp i \int_x^{x+\hat{\mu}} dx A_\mu(x), \quad A_\mu(x) = A_\mu^a T^a$$

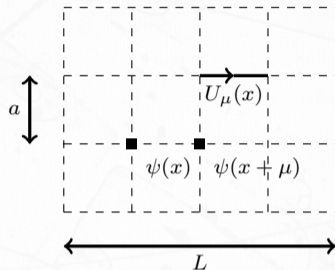


Figure: A two dimensional projection of the hypercubic lattice.

- > How are  $S_g[U]$  and  $S_f[\psi, \bar{\psi}, U]$  chosen?  
Every choice has pros and cons.  
as  $a \rightarrow 0$ , they must tend to the well-known continuum actions.

# Introduction

## Gauge and Fermion actions

> **Wilson gauge action,**

$$S_g[U] = \beta \sum_{x, \mu < \nu} \left( 1 - \frac{1}{N} \text{Re Tr } \mathcal{P}_{\mu\nu}(x) \right),$$

where  $\beta = 2N/g_0^2$ ,  $g_0$  bare coupling at scale  $1/a$ ,

$$\mathcal{P}_{\mu\nu}(x) = U_\mu(x)U_\nu(x + \hat{\mu})U_\mu^\dagger(x + \hat{\nu})U_\nu^\dagger(x)$$

> **Wilson-Dirac fermion action,**

$$S_f[\psi, \bar{\psi}, U] = \sum_x \bar{\psi}(x)(D + m_0)\psi(x)$$

where  $m_0$  is the quark mass parameter, and

$$D = (4/a + m_0)\delta_{yx} - \frac{1}{2a} \sum_{\mu} \left\{ (1 - \gamma_{\mu})U_{\mu}(x)\delta_{y,x+\mu} + (1 + \gamma_{\mu})U_{\mu}(x - \mu)^{\dagger}\delta_{y,x-\mu} \right\}$$

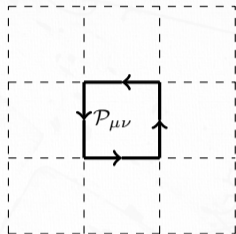


Figure: The elementary plaquette.

### Note:

- > Gauge invariance is safe.
- >  $4/a$  term: doubling problem.
- >  $\bar{\psi}(x)$ ,  $\psi(x)$  Grassmann anti-commuting numbers

# Introduction

Some more details...

## In theory...

$$\mathcal{Z} = \int \mathcal{D}U \det(D + m_0)^{1/2} e^{-S_g[U]},$$

## In practice

- > Amenable to Monte Carlo simulations.
- > Set  $a$  using physical observables.
- > Extrapolate to  $a \rightarrow 0$  using

$$O(a) = O(0) + a^n \delta_n O$$

where  $O(0)$  continuum limit.

- > Similarly for the infinite volume limit  $1/V \rightarrow 0$

## Limitations

- > No local chiral fermions in 4d.
- > Topological Freezing, metastabilities...
- > Sign problem(s):  $\mu, \theta, \dots$
- > Computational power.

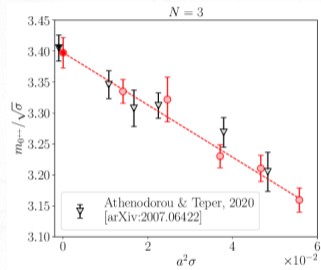


Figure: An example of continuum limit extrapolation, from (Bonanno et al. 2024).

# Typical Computation

## Typical observable

$$C(t_1 - t_2) = \langle O_1(t_1)O_2(t_2) \rangle - \langle O_1 \rangle \langle O_2 \rangle$$

$O_1(t_1)$  and  $O_2(t_2)$  are appropriate operators.

## Example: ground-state energy

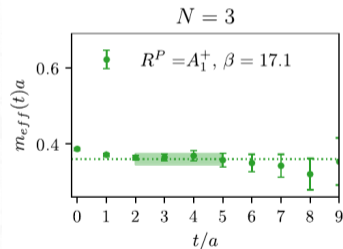
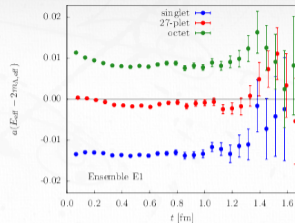
$$C(t) = \sum_n |c_n|^2 e^{-E_n t} \xrightarrow{t \rightarrow \infty} |c_0|^2 e^{-E_0 t}$$

## Effective mass

$$m_{\text{eff}}(t) = -\frac{1}{t} \log \frac{C(t+1)}{C(t)}$$

## Main difficulty

$$\frac{\text{signal}}{\text{noise}} \propto \exp -E_0 t$$



**Figure:** Top: The H dibaryon multiplet masses, taken from (Francis et al. 2019). Bottom: The  $Sp(6)$  glueball mass, taken from (Bennett, Holligan, et al. 2021).

## Take home messages

### **Lattice is...**

- ...first-principle
- ...non-perturbative
- ...not a simplification
- ...amenable to numerical simulations.
- ...systematically improvable.

### **However...**

- ...some systems are still beyond reach,
- ...massive computational needs: sustainability?

# QCD

## QCD

- > Very successful theory of the strong interaction with

$$\beta(\alpha_s) = - \left( \frac{\alpha_s}{4\pi} \right)^2 \left( 11 - \frac{2}{3} n_f \right)$$

## Non-perturbative at low energy

- > Confinement: flux-tubes, static potential, glueballs,...

$$V(R) = -\frac{4}{3} \frac{\alpha_s}{R} + \sigma R + \text{const.}$$

- > Chiral symmetry breaking:  $\pi$  as a pseudo-Nambu-Goldstone boson, EFTs

$$m_\pi^2 \propto (m_u + m_d)$$

**Lattice is the only first-principle method to access this regime**

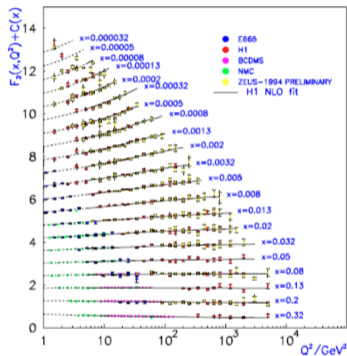


Figure: A compilation of DIS scaling results from (Schlippe)

# Gallery

## Hadron spectrum and Deconfinement

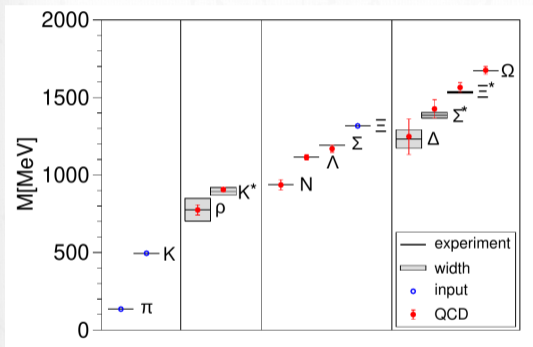


Figure: Determination of the light hadron spectrum, from (Durr et al. 2008).

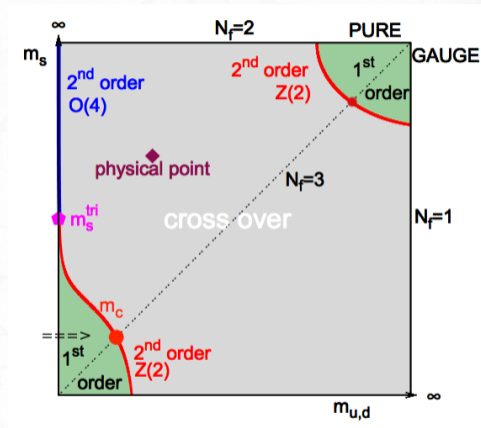


Figure: The so-called Columbia Plot for the deconfinement phase transition, from (Forcrand and D'Elia 2017).

# Scattering on the Lattice

# Scattering on the Lattice?

## Introduction

### Resonances

- > Most of the known hadrons.
- > Quark model: Exotics, Glueballs,...
- > In experiments, **asymptotic** states

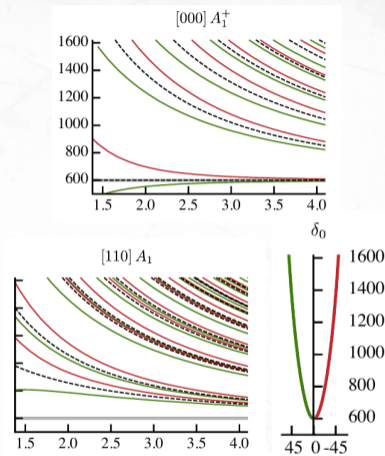
### On the Lattice

- > Euclidean time
- > **Stationary** states in a finite box,
- > Discrete spectrum.

**Lüscher's method:** Infinite-volume scattering phase shifts can be calculated finite-volume spectrum and quantization condition

$$\det [F^{-1}(E, \mathbf{P}, L) - \mathcal{M}(E)] = 0$$

see (Luscher 1986a,b).



**Figure:** Weakly attractive (green), weakly repulsive (red) and non-interacting(dashed) spectra, as a function of size, taken from (Briceno, Dudek, and Young 2018).

# The $\Lambda(1405)$

On the Lattice

## Motivation

- > Difficult to accommodate in the quark model
- > Unitarized  $\chi$ -PT suggests two poles.

## The Lattice setup:

- >  $a = 0.0633(4)(6)$  fm,  $L = 64a$
- >  $m_\pi \sim 200$  MeV,  $m_K \sim 487$  MeV,
- > Improved gauge and fermion actions.

## Results:

- > This analysis suggests the presence of two poles.
- > One is a virtual bound state.
- > One is a resonance.

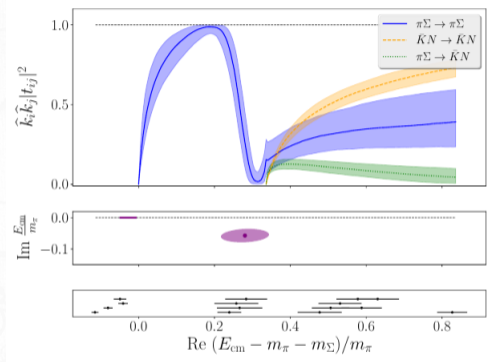


Figure: The isospin  $I = 0$  and strangeness  $S = -1$  coupled channel  $\pi\Sigma - KN^-$  transition amplitudes, and pole position, from (Bulava et al. 2024).

# Muon $g - 2$

# Updates from the $g - 2$ calculation

## Introduction

- > **Electron**  $g - 2$  among the most precise physical observables known to mankind, strong sector practically irrelevant.

$$a_l = \frac{g_l - 2}{2} = \frac{\alpha}{2} + \mathcal{O}(\alpha^2)$$

- > Lepton  $l$  potentially sensitive to new physics at scale  $\Lambda_{\text{NP}}$ ,

$$a_l - a_l^{\text{SM}} \propto \frac{m_l^2}{\Lambda_{\text{NP}}^2}$$

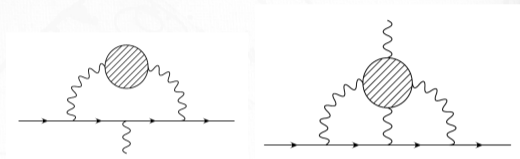
and  $m_\mu/m_e \sim 200$ .

## Why Lattice QCD?

- > First-principles, non-perturbative approach.
- > Avoids model-dependent assumptions in data-driven methods.
- > Essential to match future experimental precision

## Largest contribution to error:

- > Hadronic Vacuum polarization (HVP)
- > Hadronic Light by Light scattering (HLbL)



**Figure:** LOHVP and LOHLBL contributions. The blob represents all connected intermediate hadronic states. Taken from (Blum, Hayakawa, and Izubuchi 2012)

# Hadron contributions to $g - 2$

- > **Hadronic Vacuum Polarization (HVP):** main contribution to uncertainty

$$a_{\mu}^{HVP,LO} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^{\infty} dQ^2 f(Q^2) \Pi(Q^2)$$

where

$$\begin{aligned} \Pi_{\mu\nu}(Q) &= \int d^4x e^{iQx} \langle j_{\mu}^{em}(x) j_{\nu}^{em}(0) \rangle \\ &= (Q_{\mu}Q_{\nu} - Q^2) \Pi(Q^2) \end{aligned}$$

and  $j_{\mu}^{em}(x) = \sum_f q_f \bar{\psi}_f \gamma_{\mu} \psi_f$

- > **On the Lattice..**
  - > Evaluate current 2-points function.
- > **..or Data driven method,**

$$a_{\mu}^{HVP,LO} = \frac{1}{4\pi^3} \int_{s_0}^{\infty} ds \sigma_{had,\gamma}(s) K(s)$$

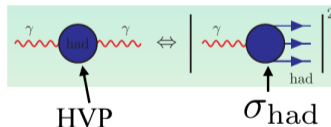
see Figure.

- > **Hadronic Light-by-Light (HLbL):** smaller contribution to uncertainty

$$\begin{aligned} \Pi^{\mu\nu\lambda\sigma}(q_1, q_2, q_3) &= -i \int d^4x d^4y d^4z \\ &e^{iq_1x + iq_2y + iq_3z} \langle j_{\mu}(x) j_{\nu}(y) j_{\lambda}(z) j_{\sigma}(0) \rangle \end{aligned}$$

- > **only Lattice**

- > Direct: evaluate current 4-points function, several methods.
- > Dispersive analysis.



**Figure:** Calculation of HVP with data driven method, taken from (Davis, 2024)

# Updated BMW'23 results

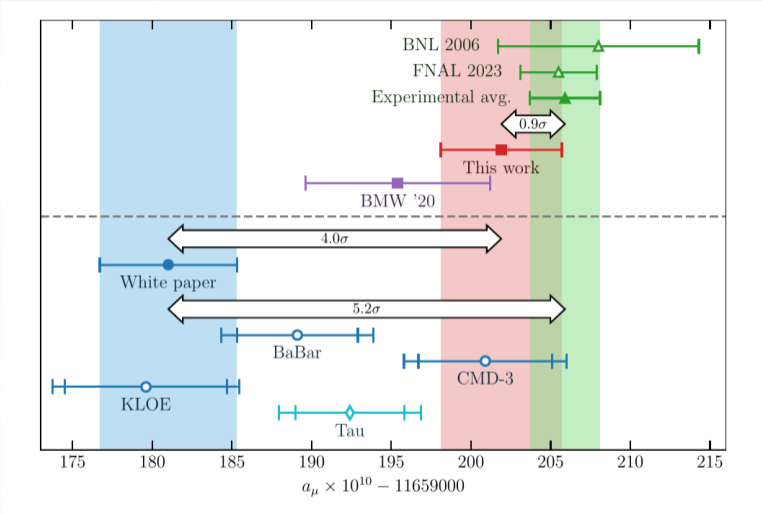


Figure: Blue: data driven methods, Red: Lattice, Green: Experiments. Taken from (Boccaletti et al. 2024).

# BSM

## Shortcomings of the SM..

- > Nature of Dark Matter/Energy?
- > Baryogenesis?
- > Hierarchy Problem, origins of EW sector?
- > No Quantum Gravity!

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## ...and how to address them

- > **New Dark Sectors**, WIMP, **SIMP**,
- > Phase transitions in the early universe,
- > Supersymmetry,
- > **Composite Higgs(CH)**, **Technicolor(TC)**,...
- > Grand Unification and extra dimensions,
- > Strings?

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## The Lattice can...

- ▶ ...help identify deviations from SM:
  - muon  $g - 2$ ,
  - Hadron spectrum,
  - CKM matrix,
  - Rare decays
  - ...
- ▶ ...test BSM scenarios
  - TC and CH Models,
  - SIMP Dark Matter
  - Axions
  - ...

# Composite Higgs

# Naturalness problem

The Higgs boson is...

- ▶ ...what sparks EW symmetry breaking,
- ▶ ...a fundamental scalar,
- ▶ ...phenomenologically extremely successful...

However it is...

- > trivial?
- > Not an explanation for EW.
- ▶ ...too light? Indeed,

$$m_H^2 = \delta_{\text{SM}} m_H^2 + \delta_{\text{BSM}} m_H^2, \quad \delta_{\text{SM}} m_H^2 \propto \Lambda_{\text{SM}}^2$$

If  $\Lambda_{\text{SM}} \sim M_{\text{GUT}}$

- > SM highly generic at  $\Lambda_{\text{SM}}$
- > Then  $\frac{\delta_{\text{SM}} m_H^2}{m_H^2} \sim 10^{-24}$

If  $\Lambda_{\text{SM}} \sim 1 \text{ TeV}$

- > SM highly constrained at  $\Lambda_{\text{SM}}$ .
- > New physics very close?

No new physics at the TeV scale, so far

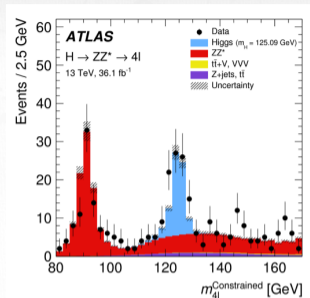


Figure:  $H \rightarrow ZZ \rightarrow 4l$ , @13 TeV, data from the LHC.

# Composite Higgs Models

## Various scenarios...

- > Supersymmetry ?
- > Walking Technicolor?

## Composite Higgs

- > Vacuum Misalignment  
→ **Higgs composite & nearly massless!**
- > Partial Compositeness  
→ **The physical top quark is a superposition of SM+Composite states.**

## On the Lattice

- > Simulate the UV theory.
- > “Measure” observables at low energy.
- > However: many interesting things accessible.

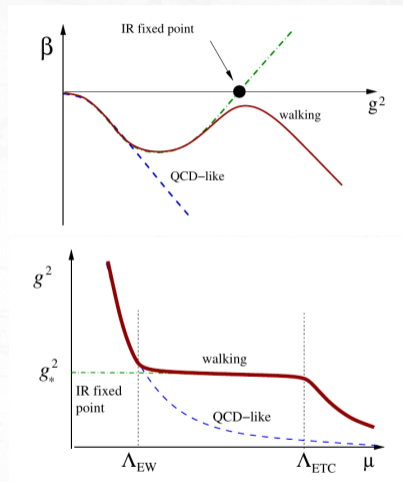
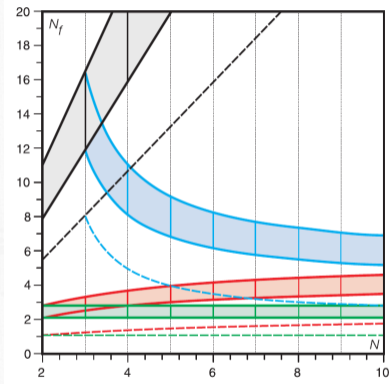


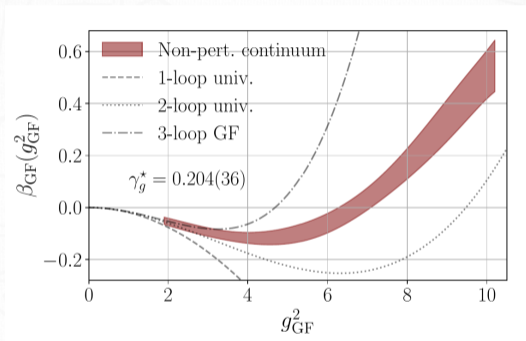
Figure: Sketch of different behaviours for the running coupling, taken from (Bennett, Hong, et al. 2024).

# Conformal window location

For  $SU(N)$



**Figure:** Top: Phase diagram of  $SU(N)$  theories in the  $(N, N_f)$  plane, taken from (Dietrich and Sannino 2007).



**Figure:** Bottom: Non-perturbative prediction for the beta function at  $N_f = 12$ , taken from (Hasenfratz and Peterson 2024).

# $Sp(4)$ Partial Compositeness

## Why $Sp(4)$ ?

- > Landscape of “good” UV completions that can accommodate PC charted. Ferretti and Karateev 2014
- >  $Sp(4)$  minimal model.

## So far

- > Glueball and Meson spectrum for  $Sp(2N_c)$ .
- >  $N_{\text{fund}} = 2$  light mesons spectrum and decay constants.
- >  $N_{\text{as}} = 3$  conformal window?
- > Quenched partial compositeness?

## Limitations

- >  $\rho$  still stable.
- > Chimera baryons still quenched.

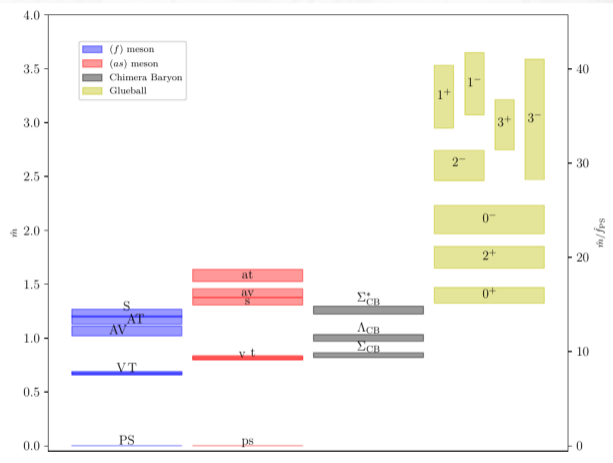


Figure: Multirepresentation quenched spectrum  $Sp(4)$ , taken from (Bennett, Hong, et al. 2024).

# Dark Sectors

# Lattice and GWs

## Baryogenesis

- > Sacharov conditions.
- > SM alone insufficient: CP violation small, EW transition cross-over.

## Dark Matter

- > Many independent clues: structure, lensing, cosmological data,...but no direct detection.
- > Cold DM models in tension with small scale structure.

## Novel Strongly Interacting sectors

- > Phase transitions of the first order:  $SU(N > 2)$   
→ GW signature!
- > Self Interacting DM (**SIDM**) scenarios address some issues of  $\Lambda$ CDM.

## On the Lattice

- > Calculation of transition temperature, and Latent heat.
- > Scattering cross sections
- > Spectrums
- > Axions!

# $Sp(4)$ deconfinement

## Motivation

- > (stochastic) GW signature.
- > Non-perturbative calculation needed.

## Setting

- >  $Sp(4)$  quenched
- > LLR algorithm to access density of states

## Results & Comments

- > Obtained Transition temperature & Latent heat
- > Scaling with volume difficult.

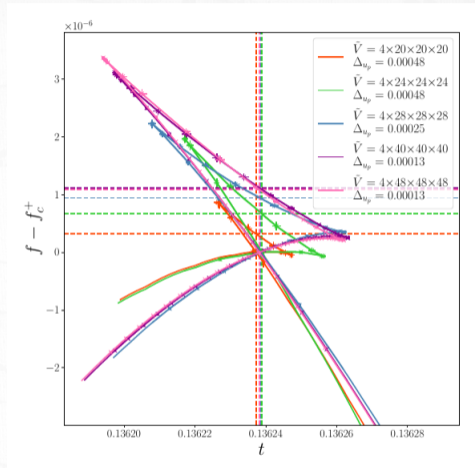


Figure: Behaviour of the free energy of  $Sp(4)$  quenched theory around first-order phase transition, taken from (Bennett, Lucini, et al. 2024).

# Take home messages

- > Lattice first-principle, non-perturbative, systematically improvable approach.
- > Useful for both qualitative and quantitative investigations.
- > Lattice has entered its “Young adult” era: strong, healthy and with many exciting possibilities and challenges ahead.

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# Scattering and Decays

Light and strange vector resonances

## The Lattice setup:

- >  $u$ ,  $d$  and also  $s$  quarks.
- > Domain-Wall fermions, a single lattice spacing.
- > What is computed: spectrum at finite volume and at various momenta.
- > Value of lat. spacing  $a$  set from  $\Omega$  baryon mass.
- > Sources of systematic error: fits ranges, discretization, masses of  $\pi$  and  $K$ ...

## Results:

$$\rho(770) \begin{cases} M_\rho = 796(5)(50)\text{MeV} \\ \Gamma_\rho = 192(10)(31)\text{MeV} \end{cases}$$

$$K^*(892) \begin{cases} M_{K^*} = 893(2)(54)\text{MeV} \\ \Gamma_{K^*} = 51(2)(11)\text{MeV} \end{cases}$$

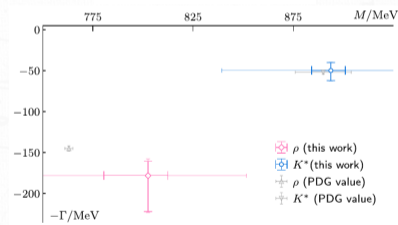
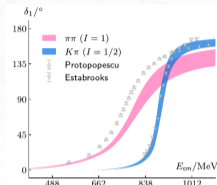


Figure: Figure and results taken from (Boyle et al. 2025a,b).

# Scattering of dark pions in $Sp(4)$

## Motivations

- > Scattering information needed.
- > Also important for CH.

## Setting

- > Dirac-Wilson action & HMC algorithm.
- > Effective Range Expansion

$$p^{2l+1} \cos \delta_l = -\frac{1}{a_l^{2l+1}} + \frac{p^2}{2r_l^{2l+1}} + \mathcal{O}(p^4)$$

and

$$\sigma_0(p) = \frac{4\pi a_0^2}{|2 - a_0 r_0 p^2 + 2i p a_0|^2}$$

where  $a_0$  scattering length and  $r_0$  effective range.

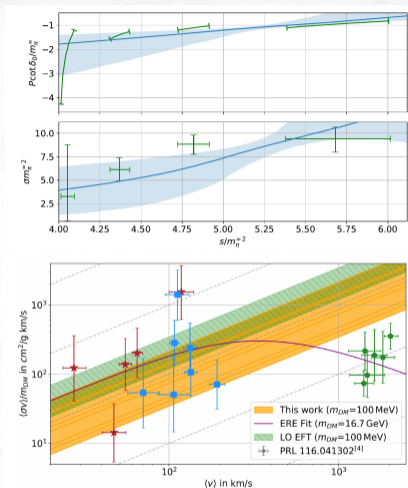


Figure: Top: phase shifts as a function of energy. Bottom: Cross-section as a function of average speed. Taken from (Dengler, Maas, and Zierler 2024).

# Phase transitions in early universe

## Motivation?

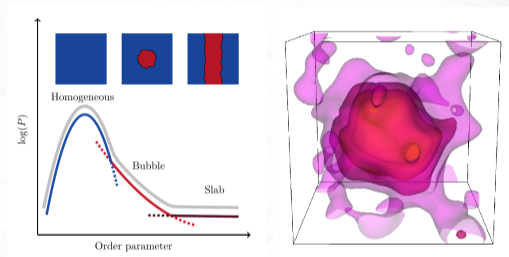
- > (stochastic) GW signature.
- > Non-perturbative calculation needed.

## Setting

- > Transition is strong at  $N(\text{colors}) = 8$ , hence it allows for superheating & supercooling
- > Formulated an EFT and Simulated *that*.
- > Used smeared order parameter.

## Results & Comments

- > Possible to resolve critical bubble with appropriate methods.
- > Larger volumes are necessary.



**Figure:** Left: Free energy behaviour as a function of the order parameter (Hällfors and Rummukainen 2025). Right: Spatial distribution of values of order parameter in two dimensions, taken from (Seppä, Rummukainen, and Weir 2025).

# Charting $Sp(N_c)$ theories

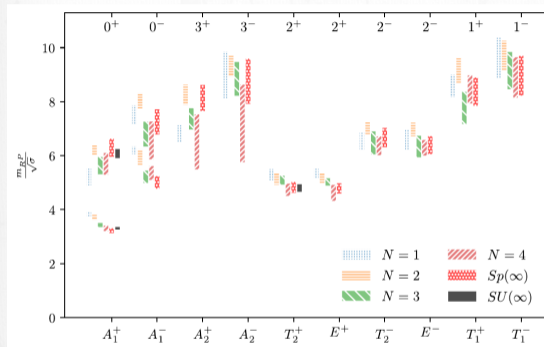


Figure: Quenched glueball spectrum, taken from (Bennett, Holligan, et al. 2021).

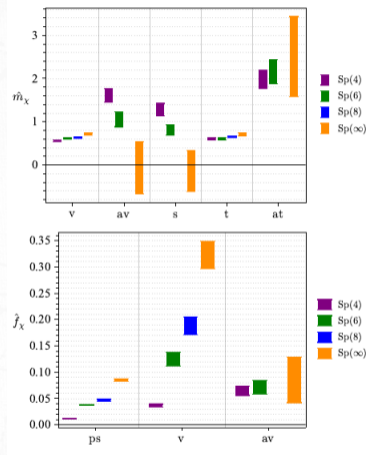


Figure: Quenched meson spectrum, taken from (Bennett, Hong, et al. 2024).