

Terahertz Acceleration

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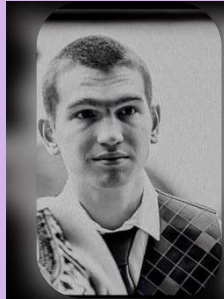
The Cockcroft Institute: Terahertz acceleration group

MANCHESTER
1824

The University of Manchester



Vasileios Georgiadis



Elliott Smith



Dr. Morgan Hibberd



Dr. Darren Graham



Dr. Robert Appleby

Lancaster
University



Oliver Finlay



Dr. Daniel Lake



Dr. Sergey Siaber



Prof. Steven Jamison



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Science & Technology
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Dr. David Walsh



Dr. Edward Snedden



Dr. Conor Mosley



Dr. Oznur Apsimon



Dr. Rosa Letizia



Dr. Osman Akin



Why THz radiation for particle acceleration?

Light sources: xray Free-electron lasers
for femtosecond x-ray pulses

5GeV beams, <10fs bunches sought



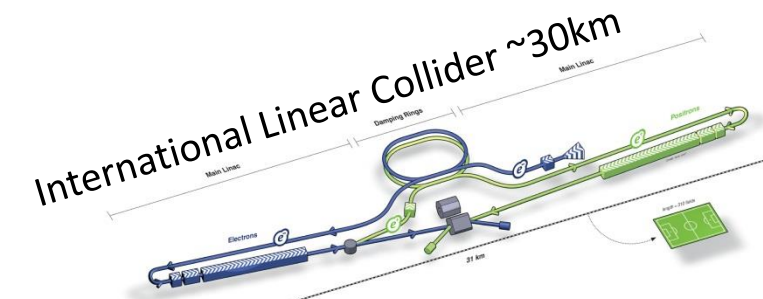
SACLA, Japan



LCLS, USA

Femtosecond manipulation of relativistic
electron bunches using nanosecond period
Radio Frequency sources

Particle physics



'CERN's 'Compact Linear Collider' ~48km

Scale of accelerators set by
 $O(100\text{MV/m})$ gradient limitation
(fundamental limits, atomic lattice
distortion)

THz driven
Acceleration

- Field strengths $O(100\text{MV/cm})$
- Picosecond period, broad-bandwidth,
direct waveform control

Why THz radiation for particle acceleration?

Radio frequencies

- Frequency ≈ 3 GHz
period ≈ 300 ps,
- $\lambda \approx 10$ cm
Cavity aperture ~ 2 cm
- $E_{\text{acc}} \sim 10 - 100$ MV m $^{-1}$
- Pulse length: 5 ms (superconducting)
5 μ s (normal)

THz frequencies

- Frequency ≈ 1 THz
period ≈ 1 ps
- $\lambda \approx 300$ μ m – 1 mm
Cavity aperture ~ 1 cm
- $E_{\text{acc}} \approx 100$ MV m $^{-1}$
- Pulse length 1 ps – 100 fs

Optical frequencies

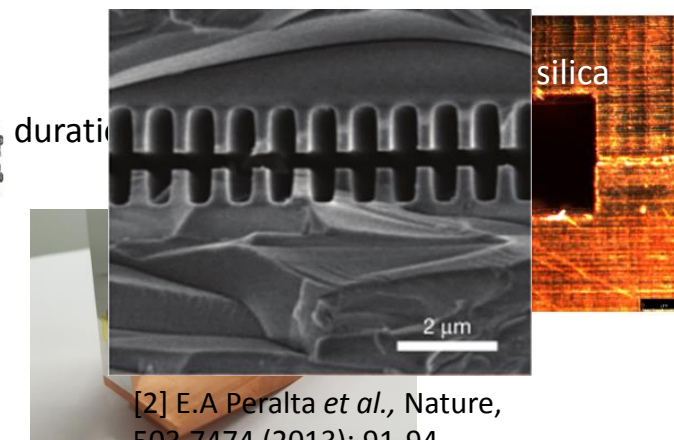
- Frequency ≈ 300 THz
period ≈ 3 fs,
- $\lambda \approx 0.8 - 1$ μ m
Cavity aperture ≈ 1 μ m
- $E_{\text{acc}} \approx 100$ MV m $^{-1}$ to \gg GV/m
- Pulse length 20 fs – 1 ps

THz advantages:



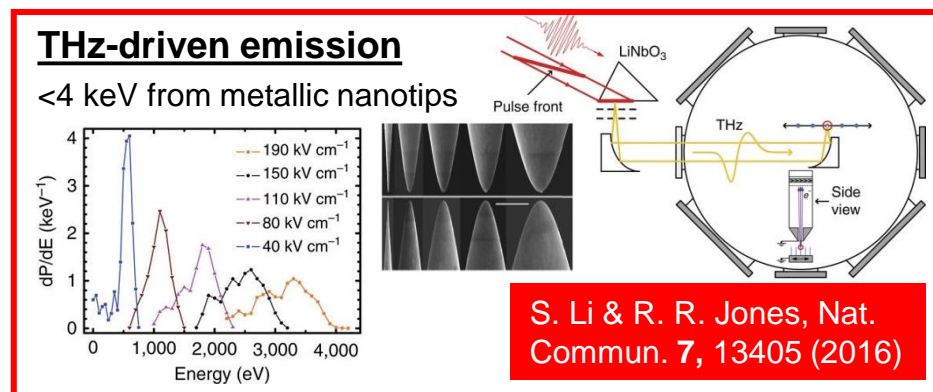
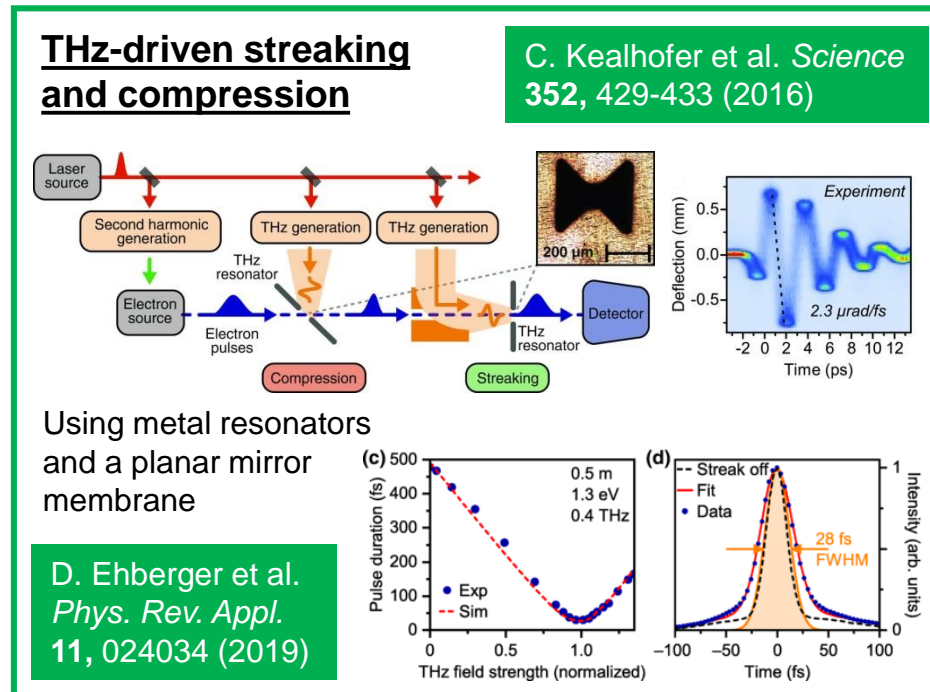
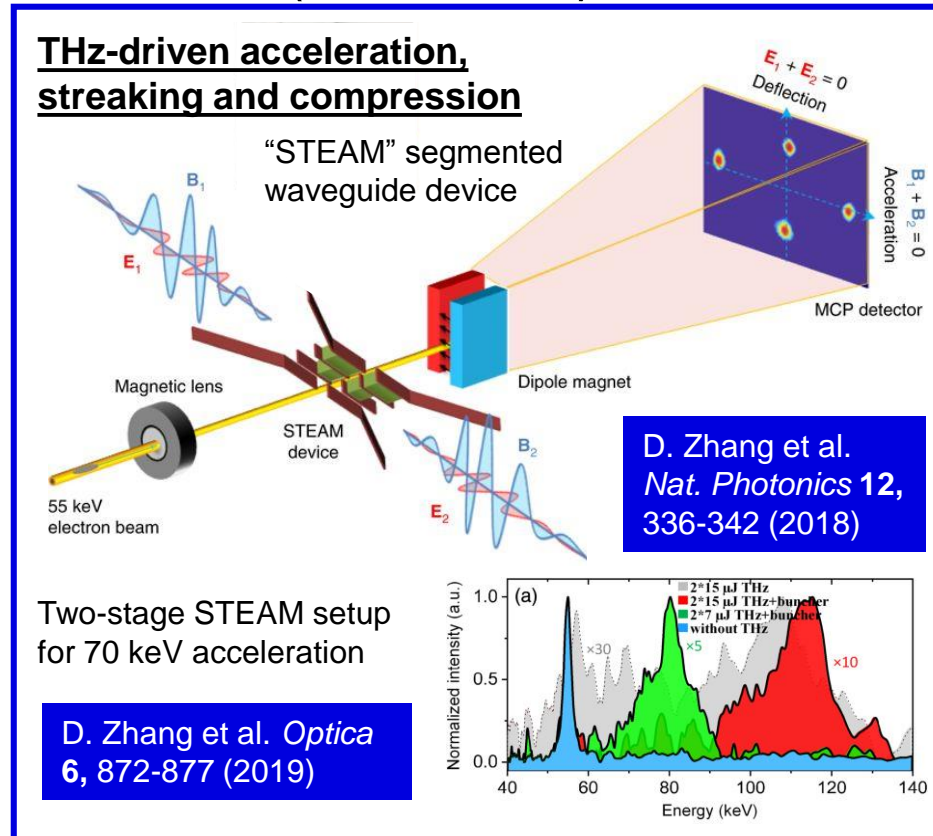
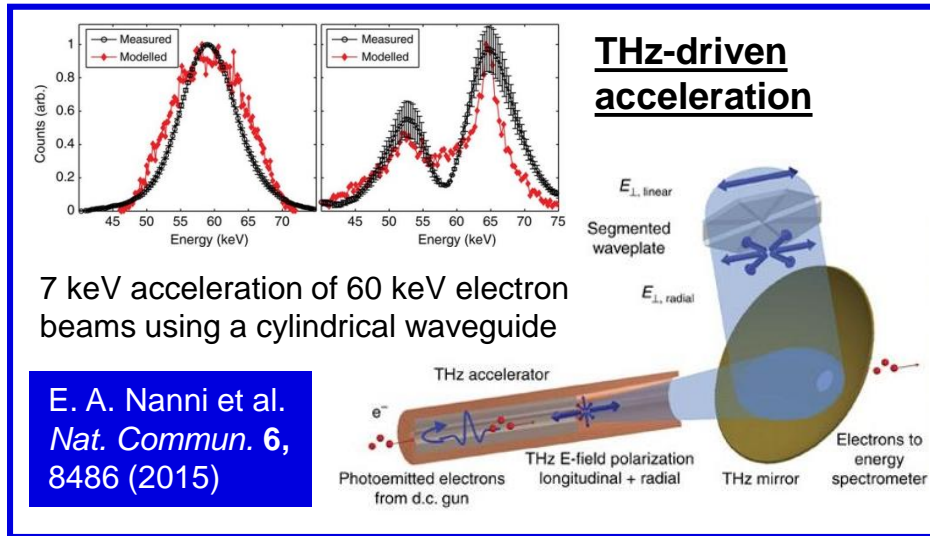
THz challenges:

- Obtaining sub-luminal phase velocities
- Broadband pulses – controlling dispersion
- High field THz sources still challenging

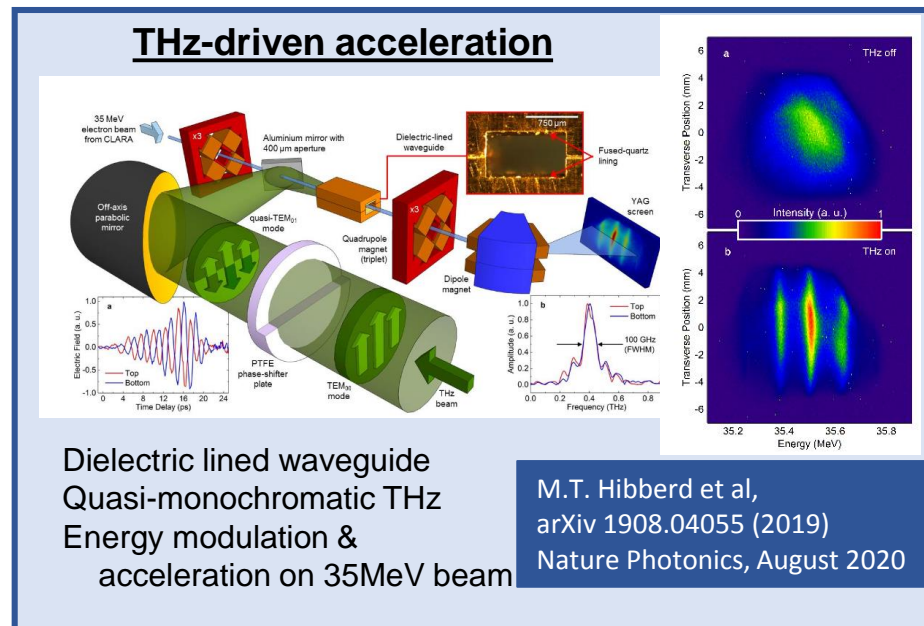
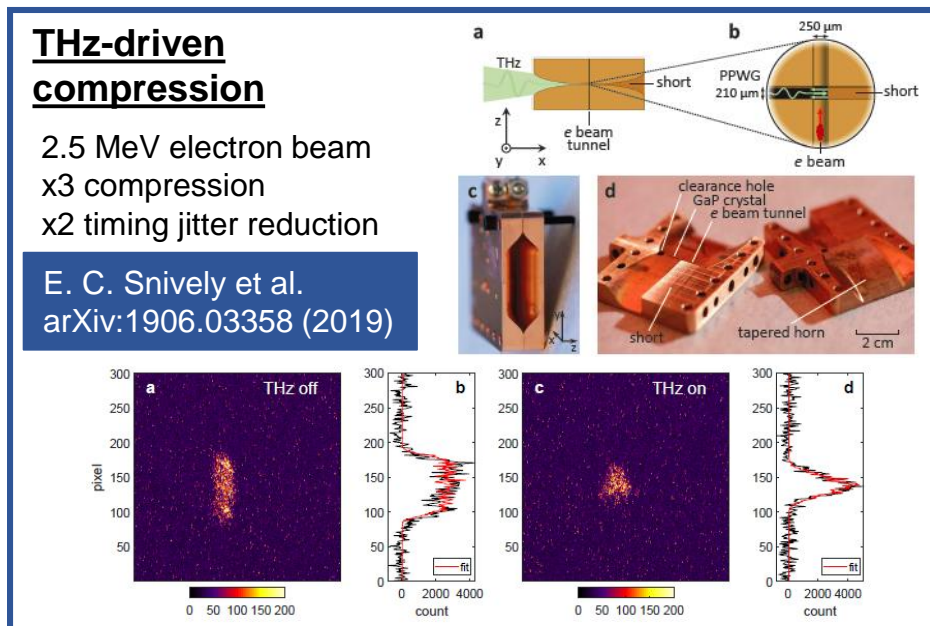
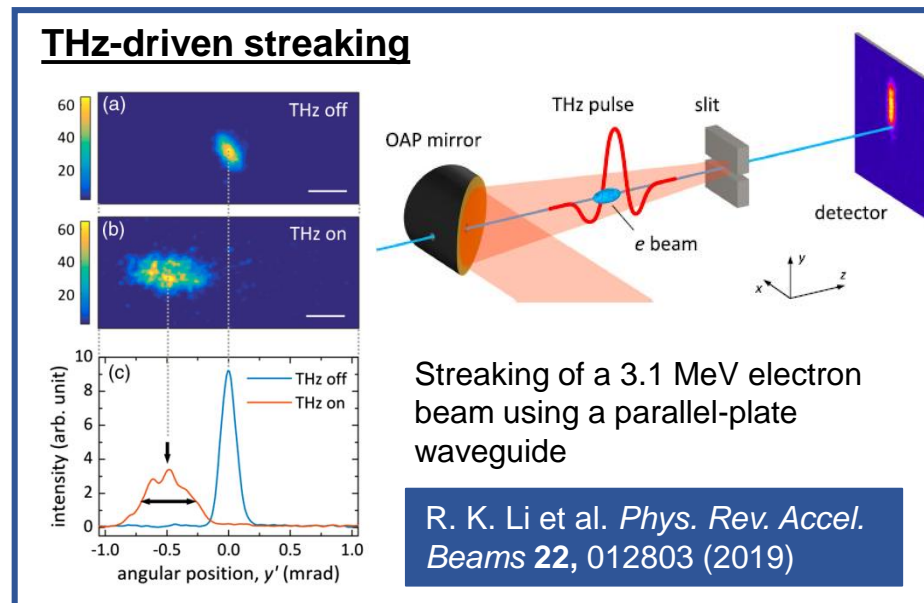
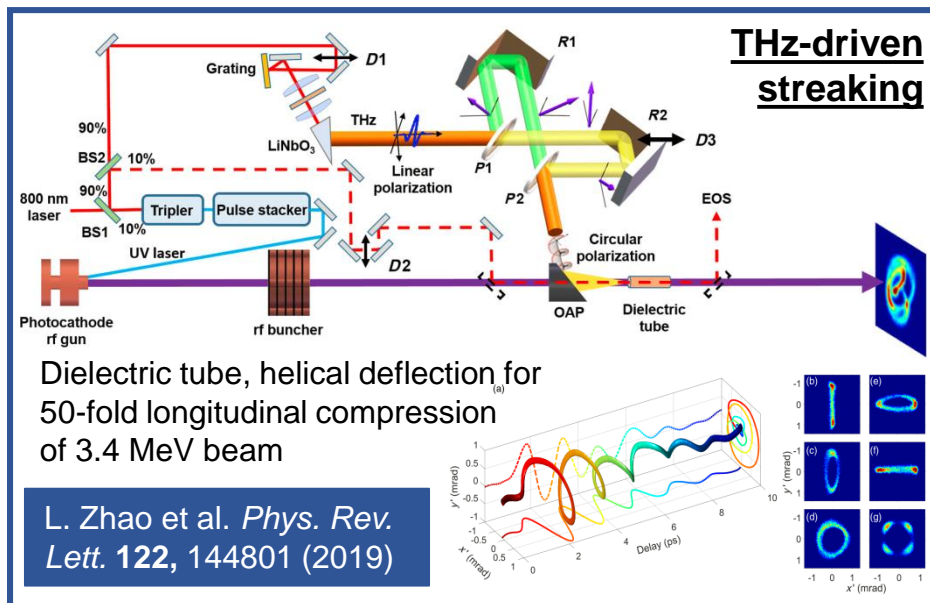


[2] E.A Peralta *et al.*, Nature, 503.7474 (2013): 91-94

Terahertz-driven interactions: Sub-relativistic electrons (<100 keV)

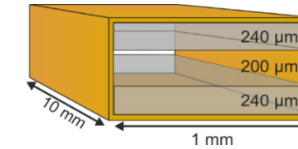


Terahertz-driven interactions: Relativistic electrons

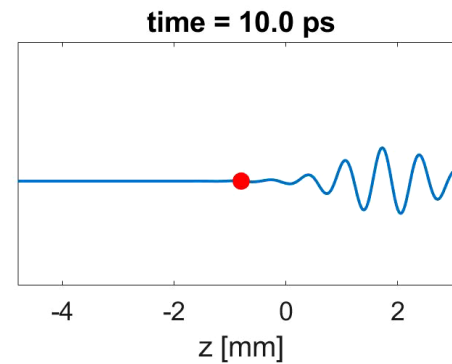


What is needed for THz acceleration?

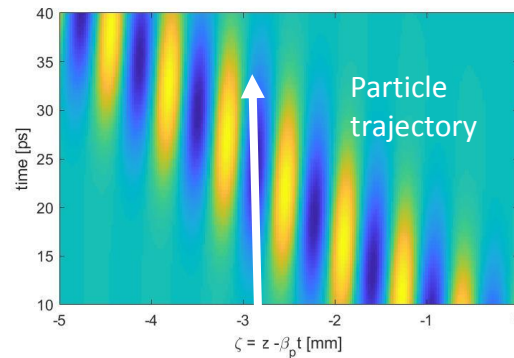
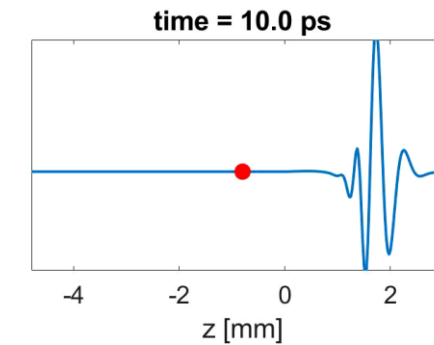
- Strong longitudinal field component
- Phase matched to particle velocity
- Dielectric Lined Waveguides used to achieve both
- Phase velocity matching, $v_\phi = \beta c$
- Group velocity $v_g \ll v_\phi$



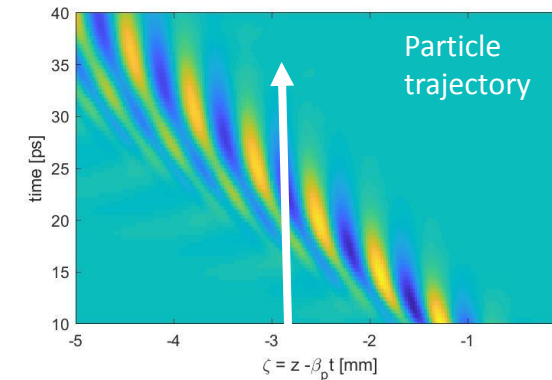
few-cycle THz input
(quasi-monochromatic)



Single-cycle THz input

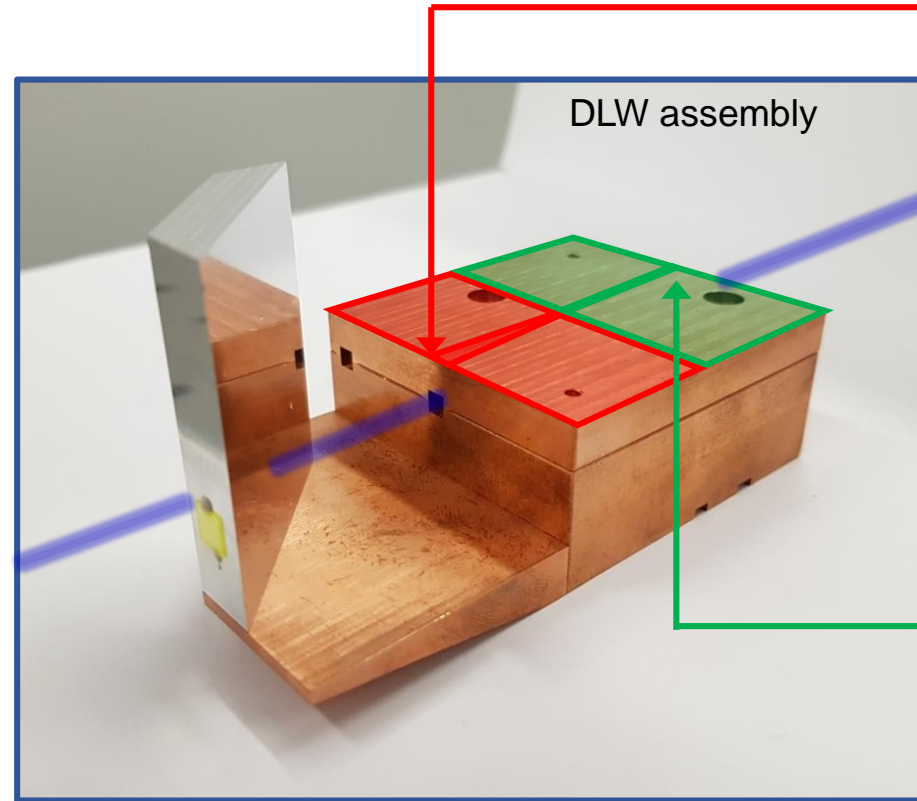
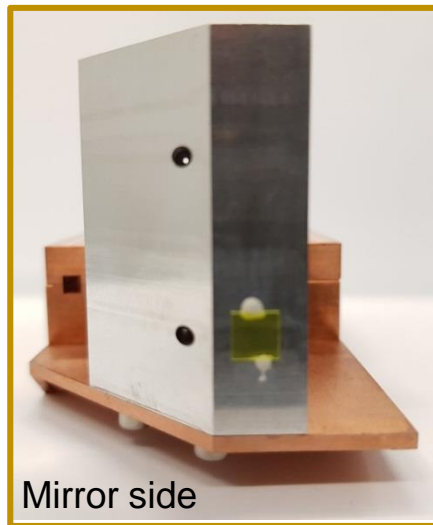
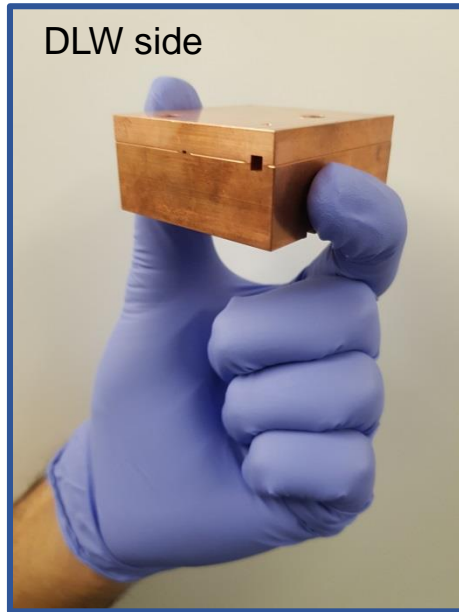


Field in frame of 50 MeV
electron beam ($v = 0.9995 c$)



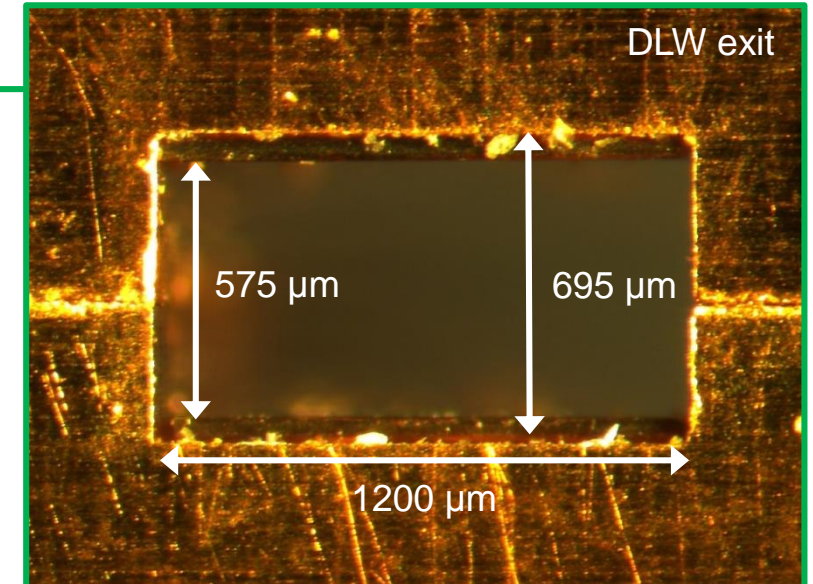
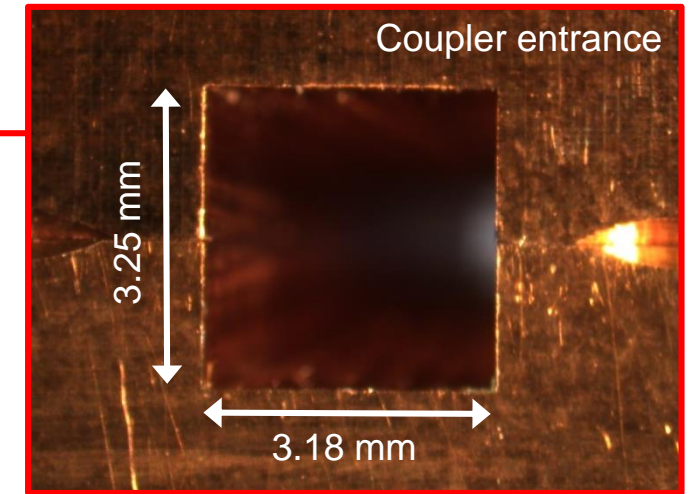
Particle see constant electric field phase (accelerator or decelerating)

Our approach: Rectangular dielectric-lined waveguide



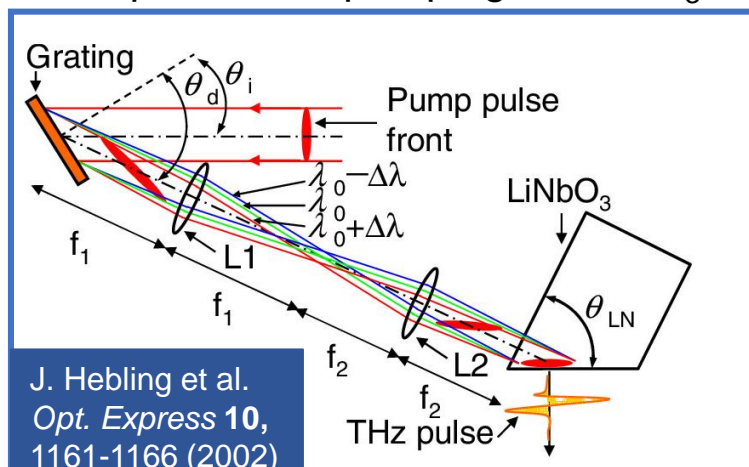
Dielectric lining results in:
Longitudinal Section Magnetic modes

- LSM₁₁ for acceleration
- Designed for velocity-matching to the 35 MeV electron beam



Our approach: Narrowband terahertz source

Tilted-pulse-front pumping in LiNbO₃

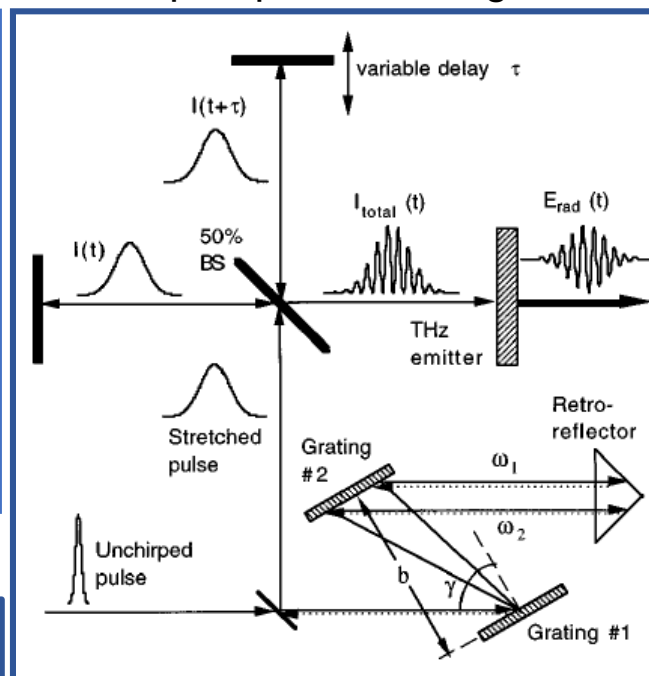


J. Hebling et al.
Opt. Express **10**,
1161-1166 (2002)

Diagram from:

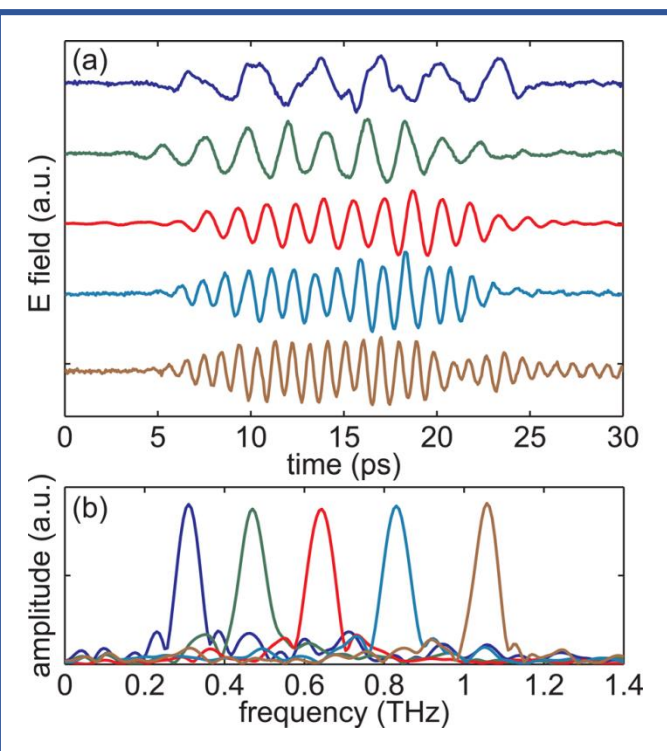
H. Hirori et al. *Appl. Phys. Lett.* **98**, 091106 (2011)

Chirped-pulse beating for narrowband THz generation



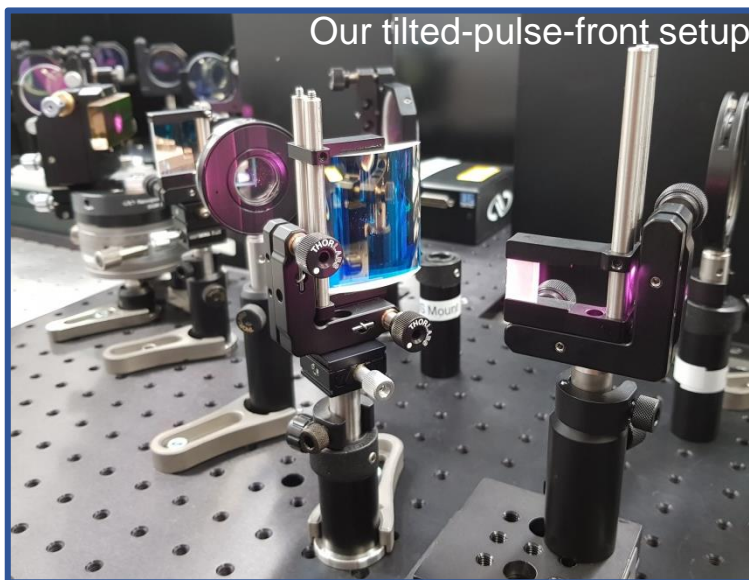
First demonstrated
using a Michelson
interferometer with a
photoconductive
antenna

A.S. Weling & D.H.
Auston, *J. Opt. Soc.
Am. B* **13**, 2783 (1996)



Also shown using an
etalon with a LiNbO₃
crystal for a high-power,
tunable multi-cycle THz
source

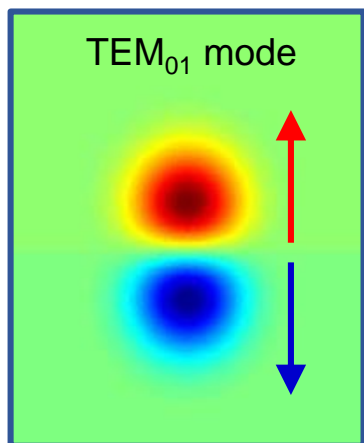
Z. Chen et al. *Appl.
Phys. Lett.* **99**, 071102
(2011)



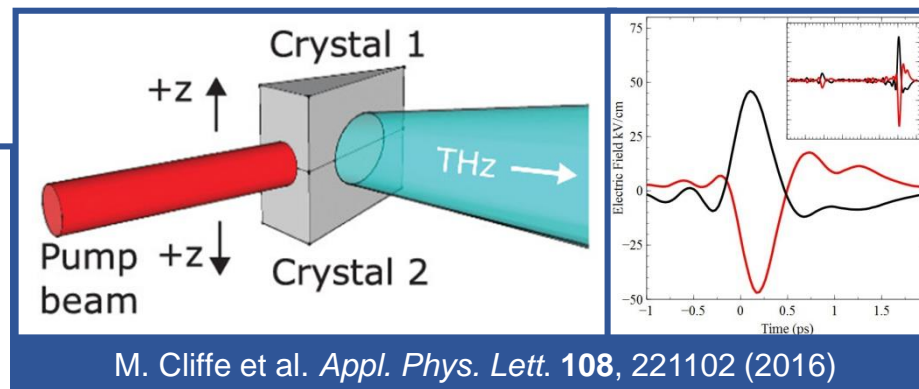
Our tilted-pulse-front setup

Our approach: Polarisation of the terahertz beam

To excite the accelerating mode of the DLW
 → THz beam with a TEM_{01} mode is required

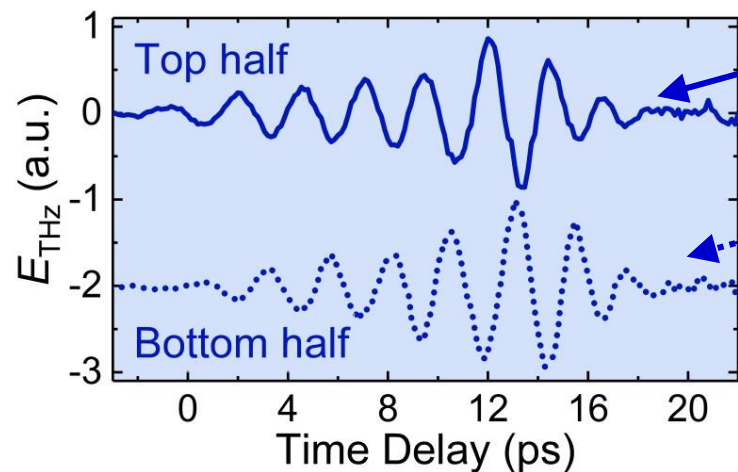


Previously demonstrated using the interferometric recombination of two polarity-inverted, linearly polarised THz beams

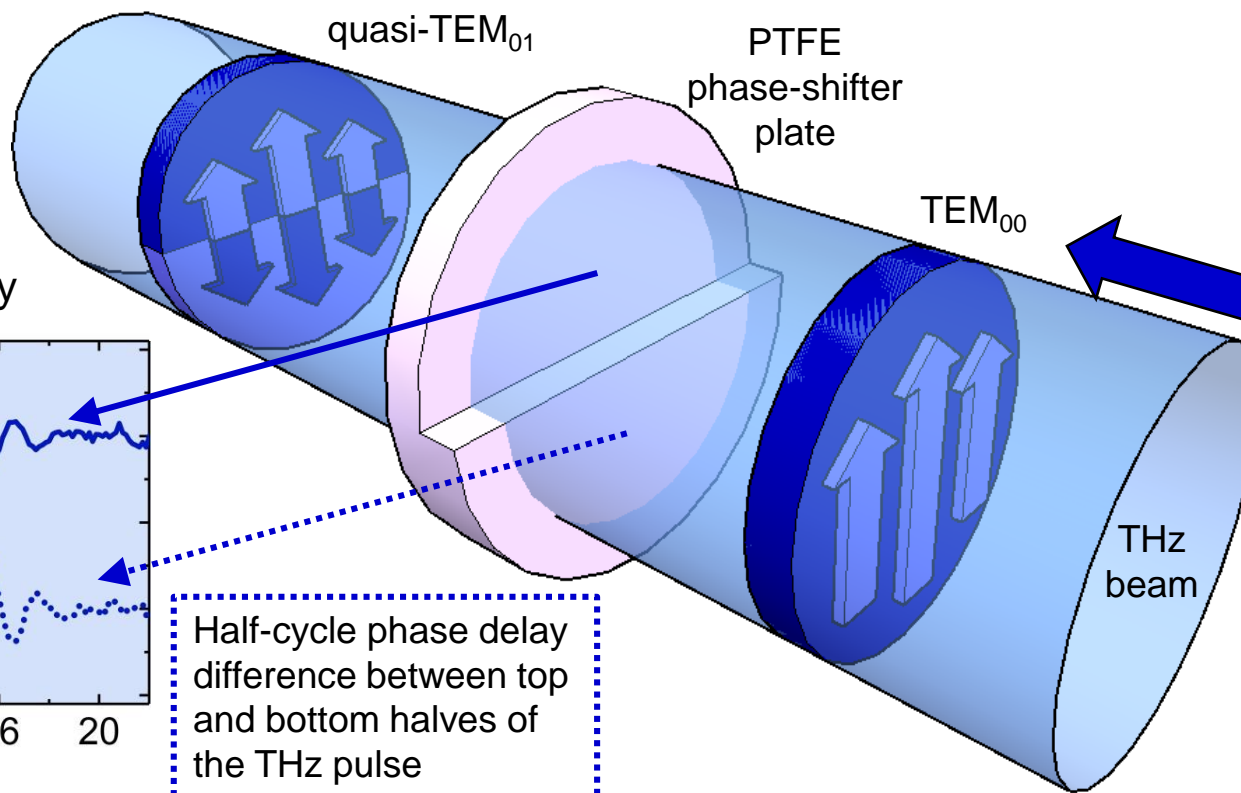


Phase-shifter plate

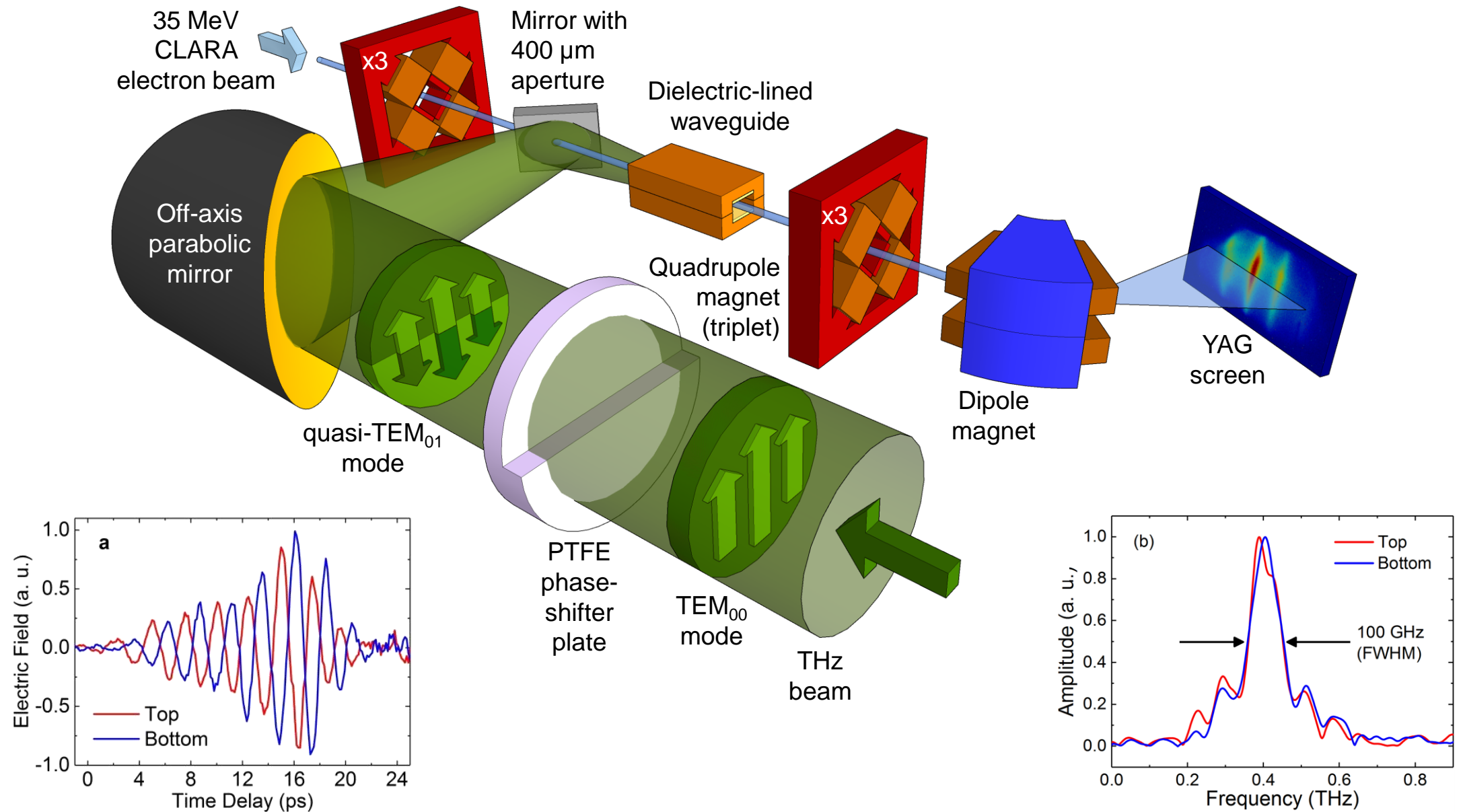
Achieve effective polarity inversion by using a phase delay



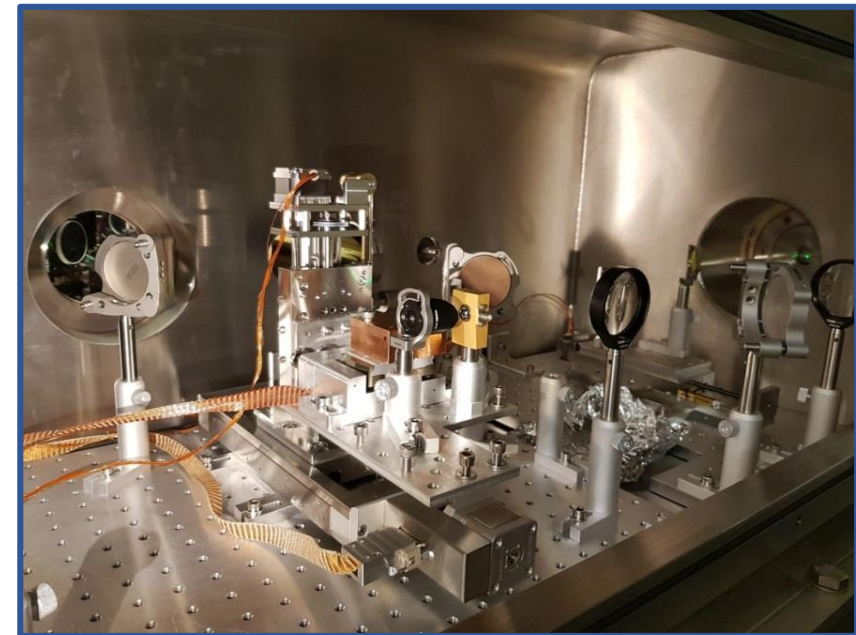
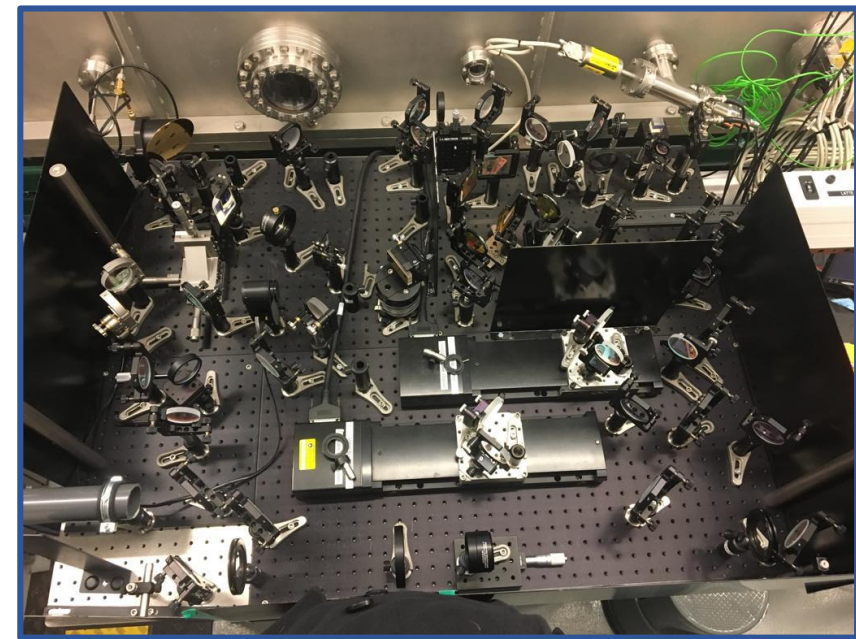
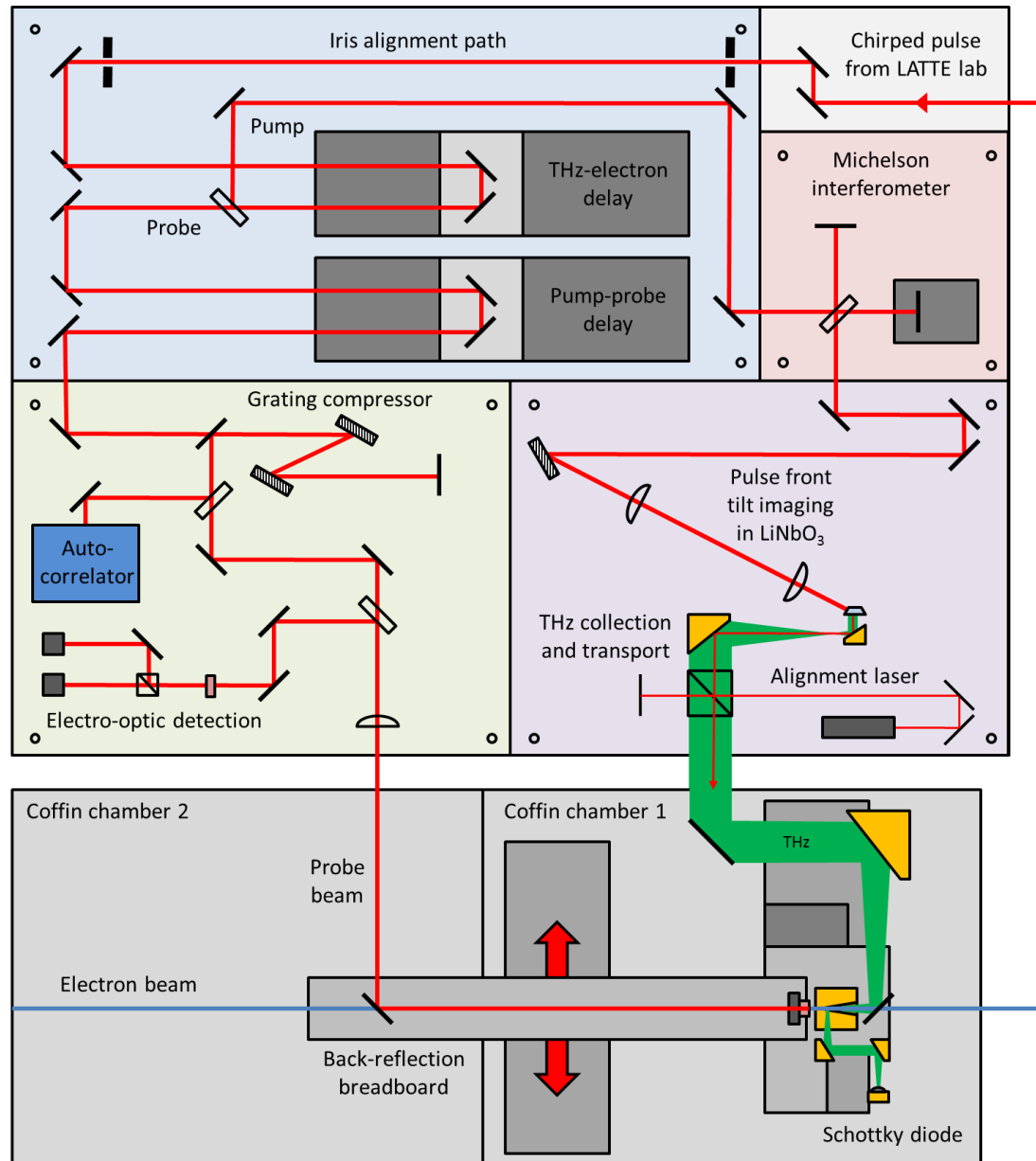
Half-cycle phase delay difference between top and bottom halves of the THz pulse



Our approach: Experimental approach



Our approach: Experimental setup



Results: Chirped electron bunch

CLARA operated in a long electron bunch configuration

- Bunch duration of 6 ps FWHM
 - Linear (approx.) chirp of 53 keV/ps
 - Bunch charge of 60 pC
- ← Estimated from beam dynamics simulations

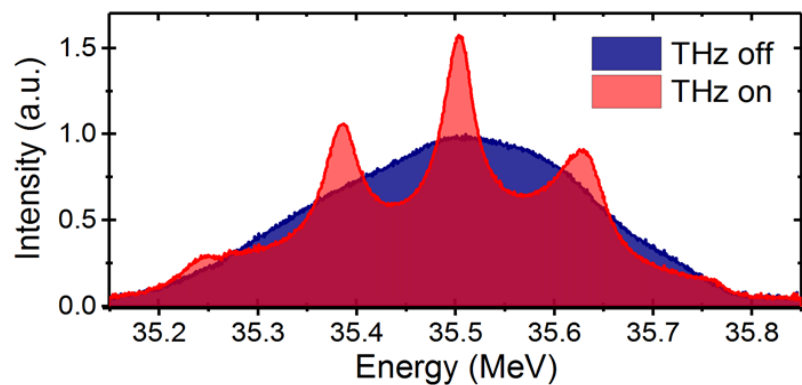
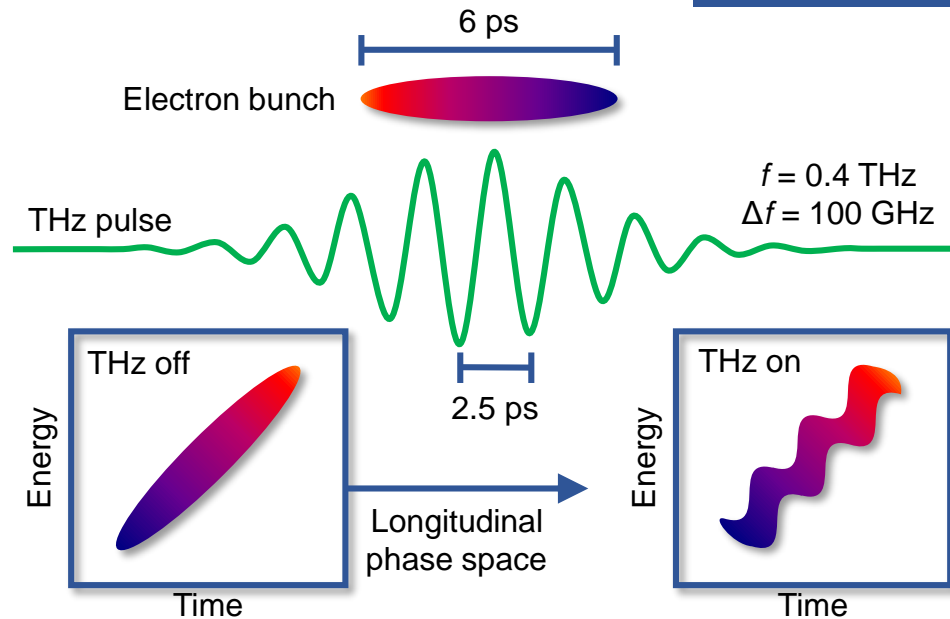
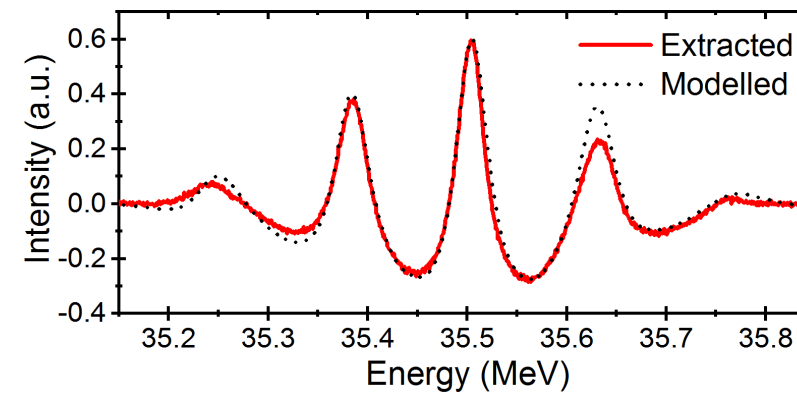
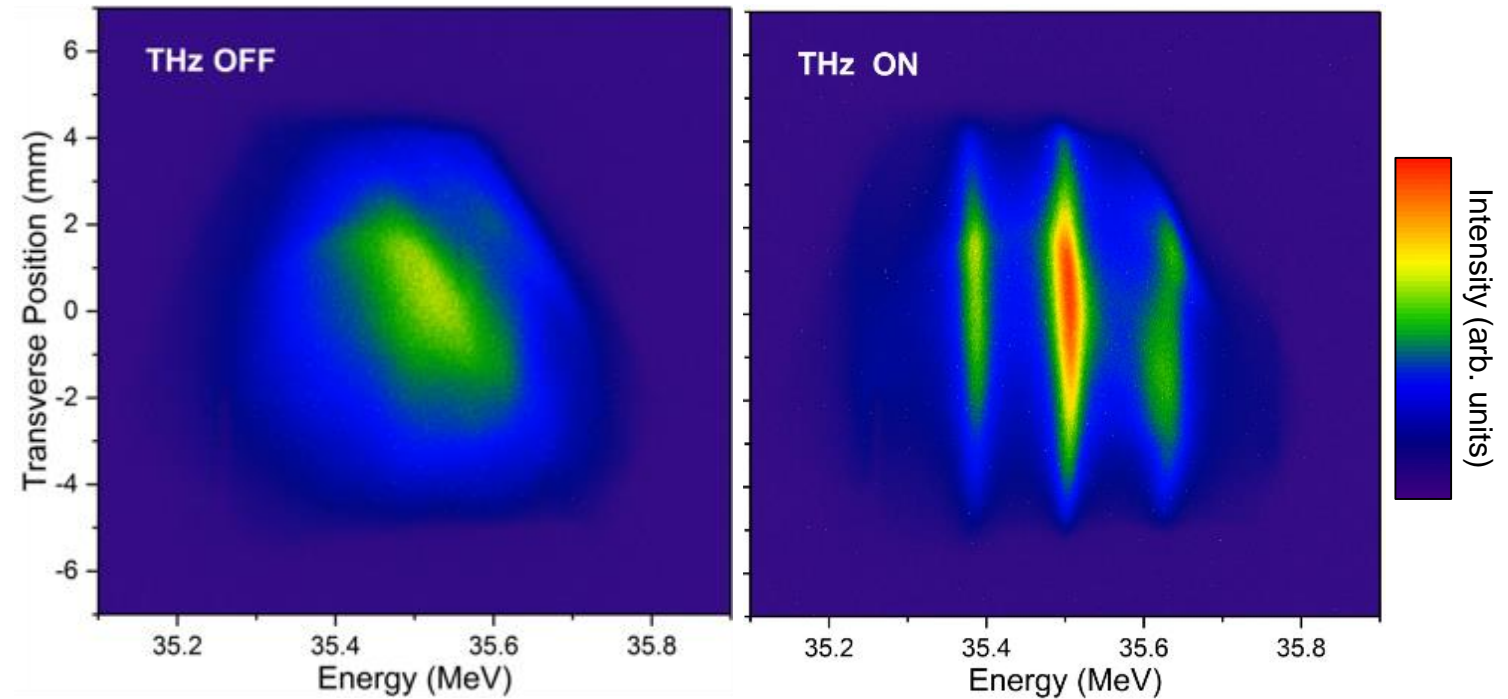


Image on electron energy spectrometer



Results: Chirped electron bunch

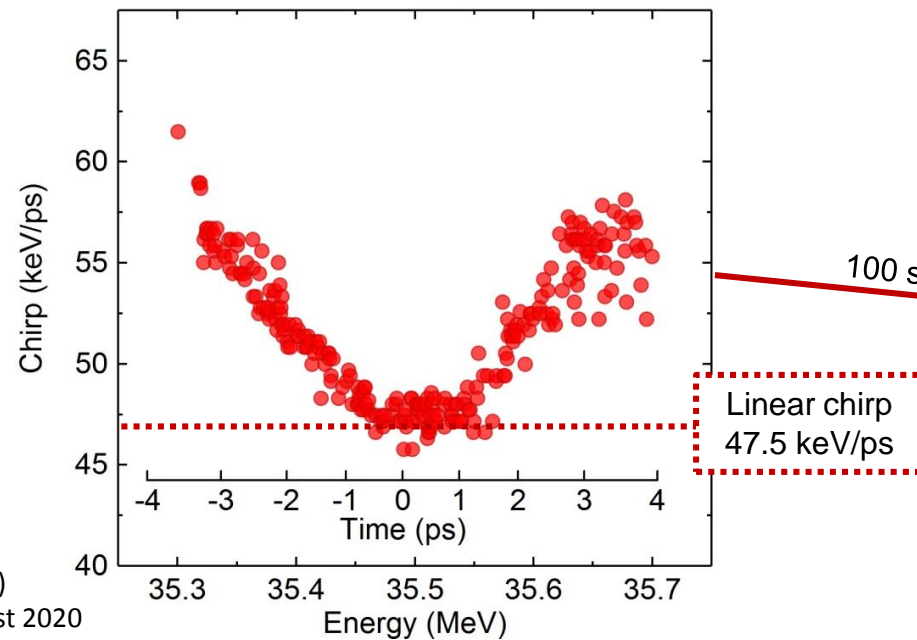
CLARA operated in a long electron bunch configuration

- Bunch duration of 6 ps FWHM
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 - Bunch charge of 60 pC
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Measured longitudinal phase-space distribution

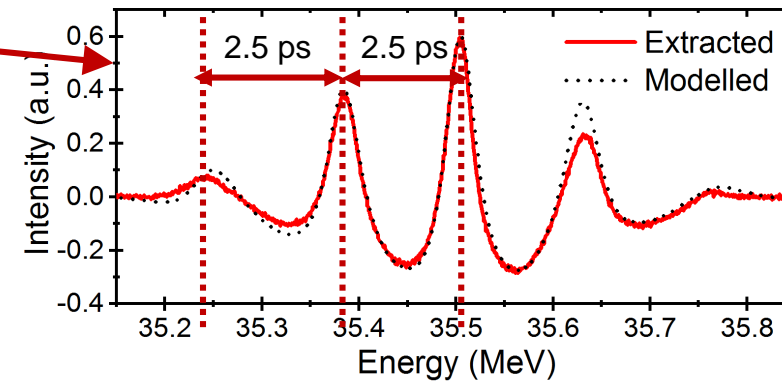
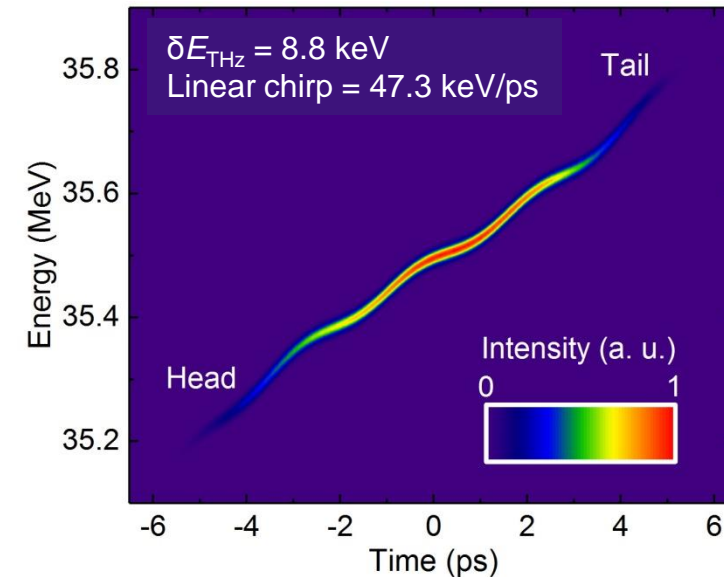
Obtained from 100 single-shot spectra with varying THz-electron bunch timing, using:

- The measured energy modulation period $\delta E(E)$
- The known THz period $\tau = 2.5$ ps (0.4 THz)



Determining the peak terahertz-driven acceleration

Calculation imposing a sinusoidal THz-driven modulation on to a model electron bunch, including:
→ Bunch emittance, energy spread and chirp

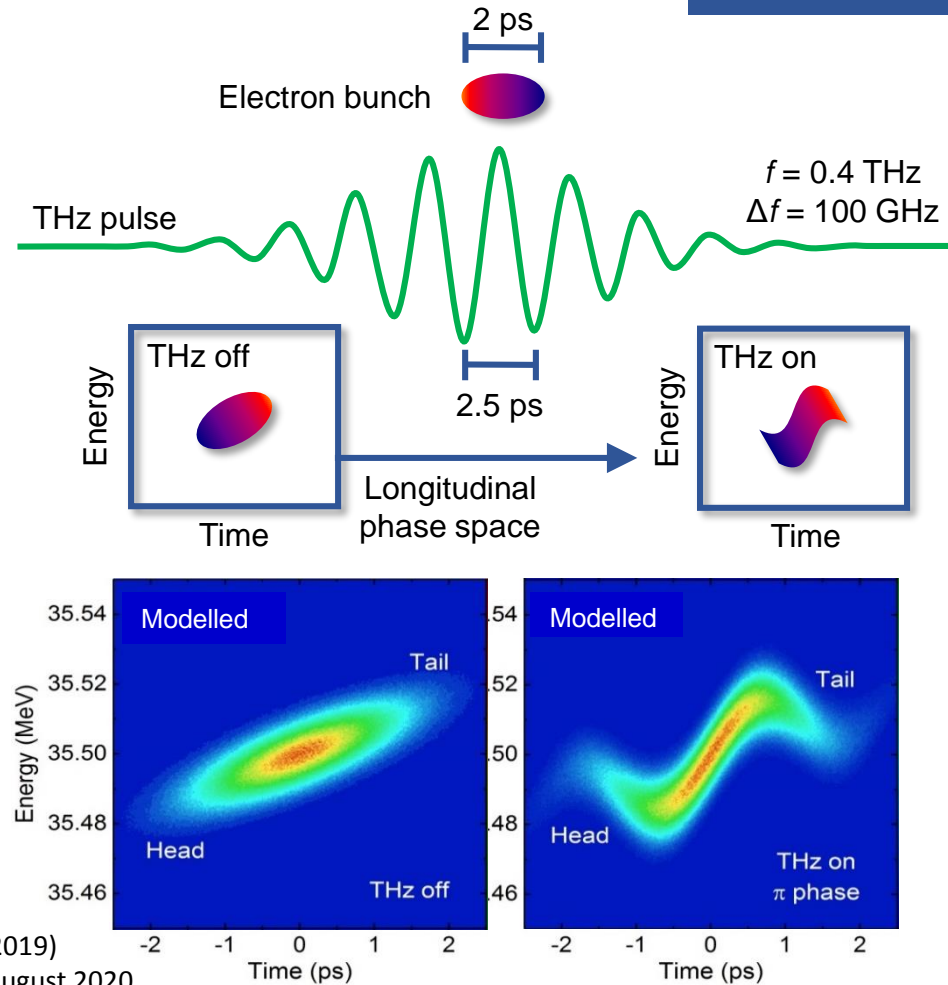


Results: Short electron bunch

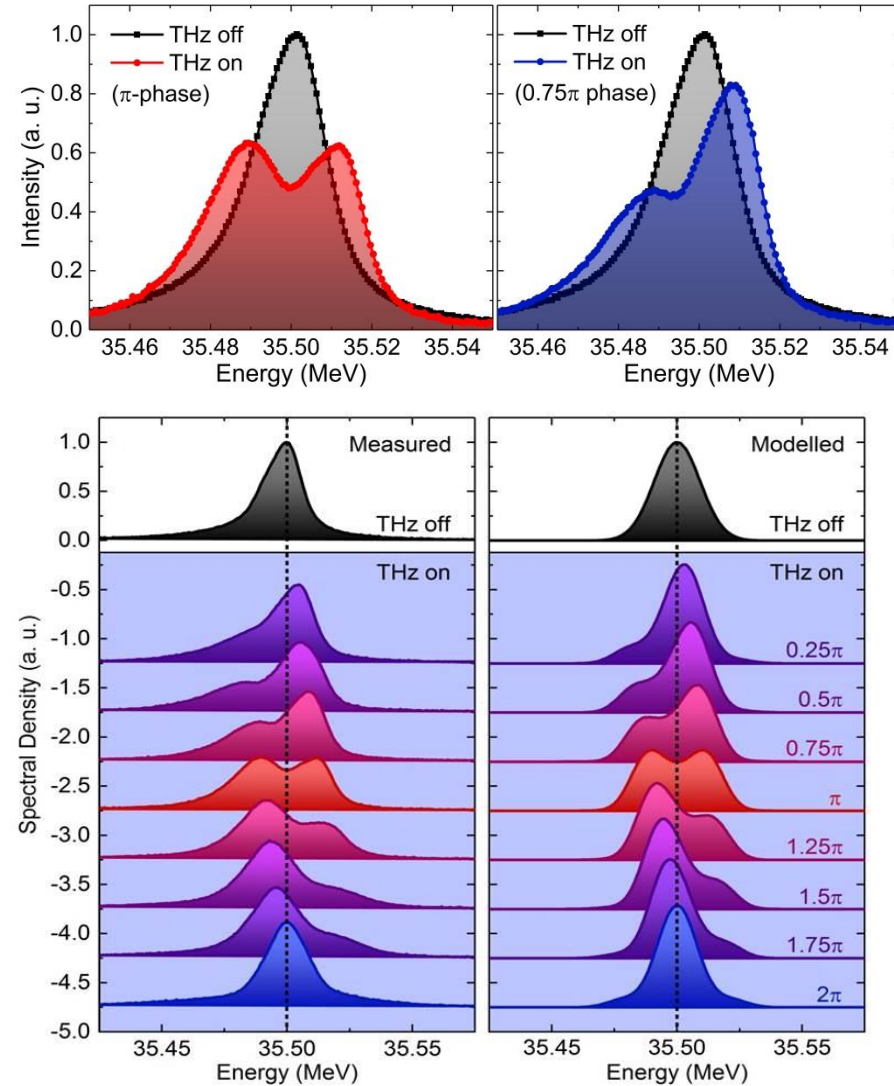
CLARA operated in a short electron bunch configuration

- Bunch duration of 2 ps FWHM
- Residual linear chirp of <math><10\text{ keV/ps}</math>

Estimated from
beam dynamics
simulations



“Single-bucket” acceleration

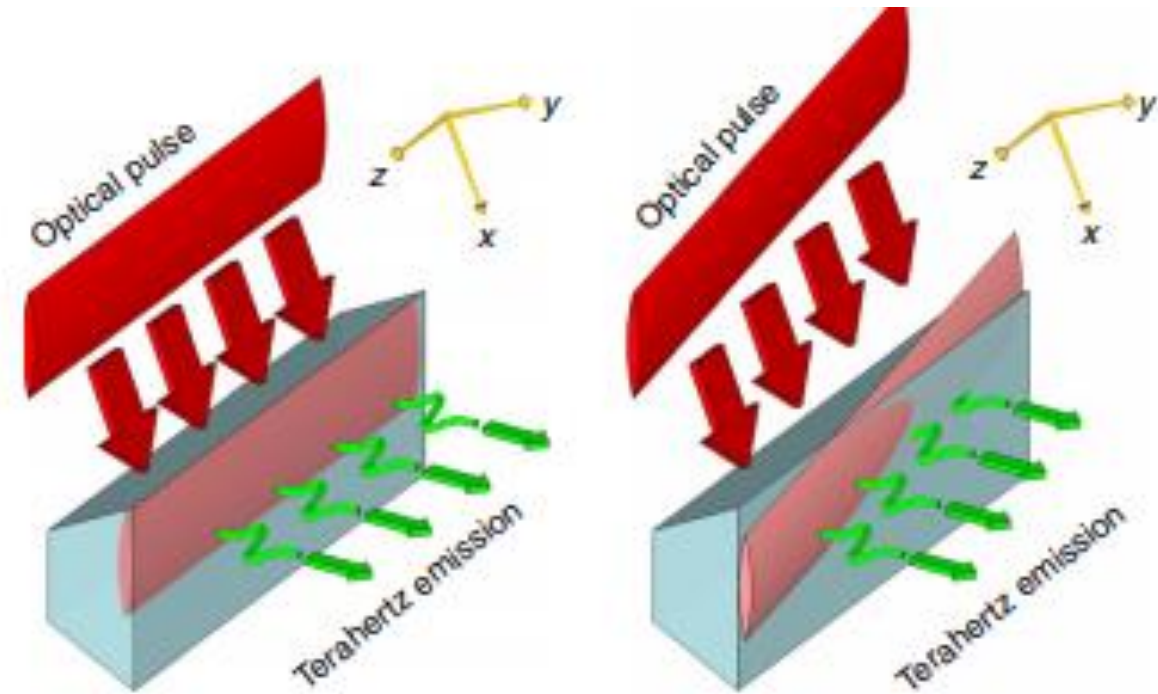
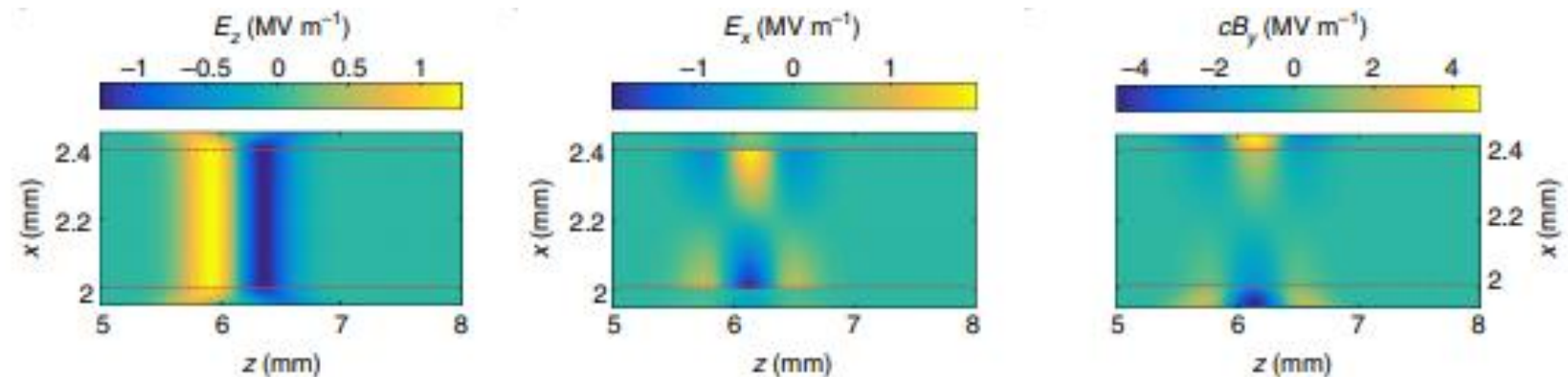
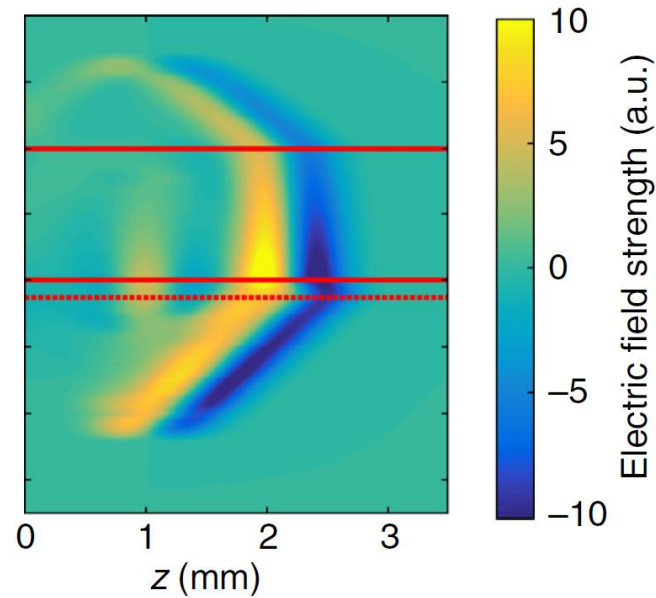


M.T. Hibberd et al,
arXiv 1908.04055 (2019)
Nature Photonics, August 2020

Another Approach: Travelling Wave THz Source

Use evanescent wave formed at crystal-vacuum boundary

- Additional pulse-front-tilt gives time delay across THz pulse
- Effective phase velocity matched to particles by tuning angle of extra tilt
- Can use single cycle pulse – No need for narrowband
- No need for a ‘structure’



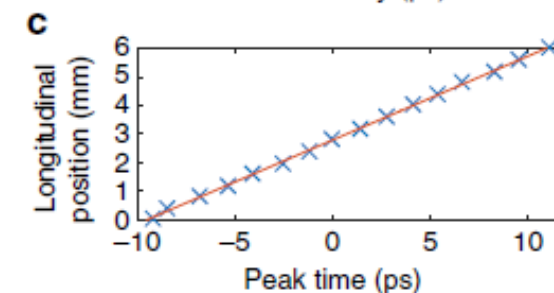
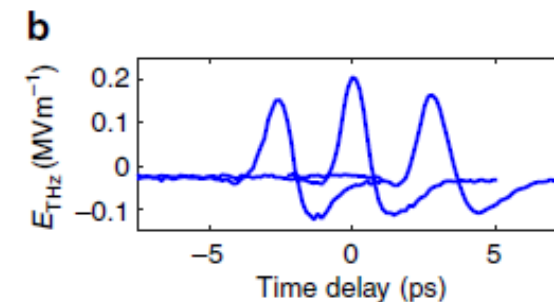
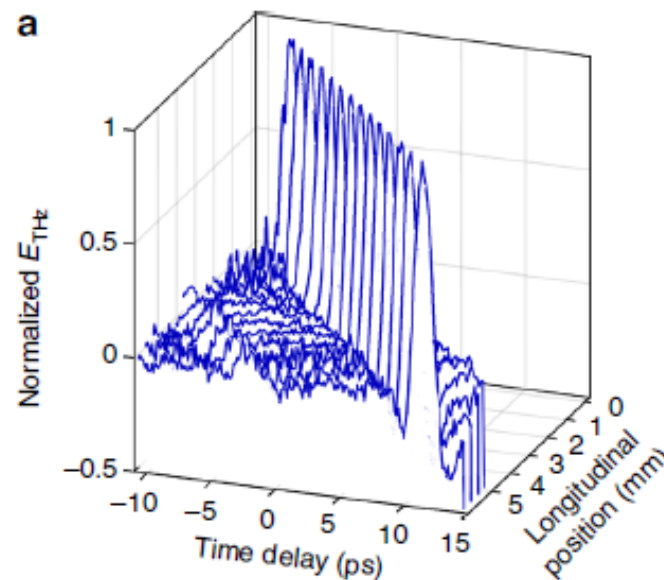
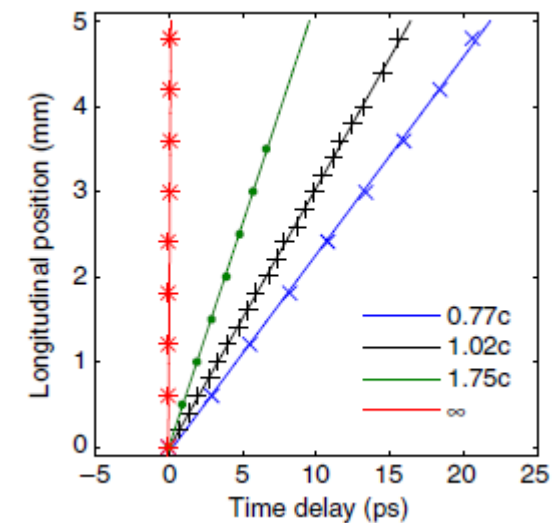
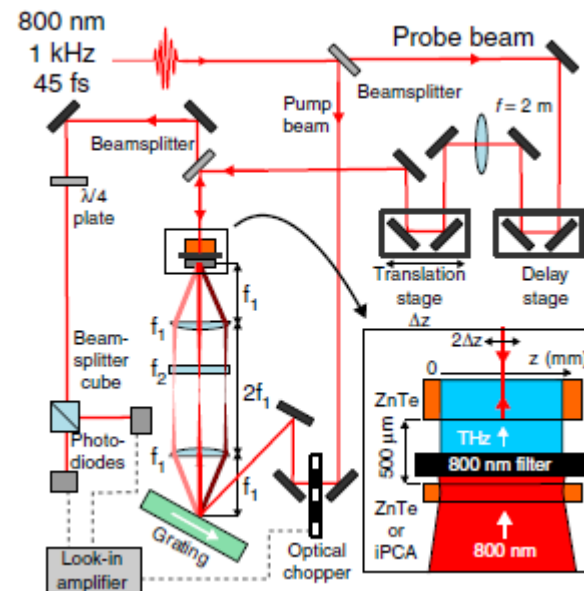
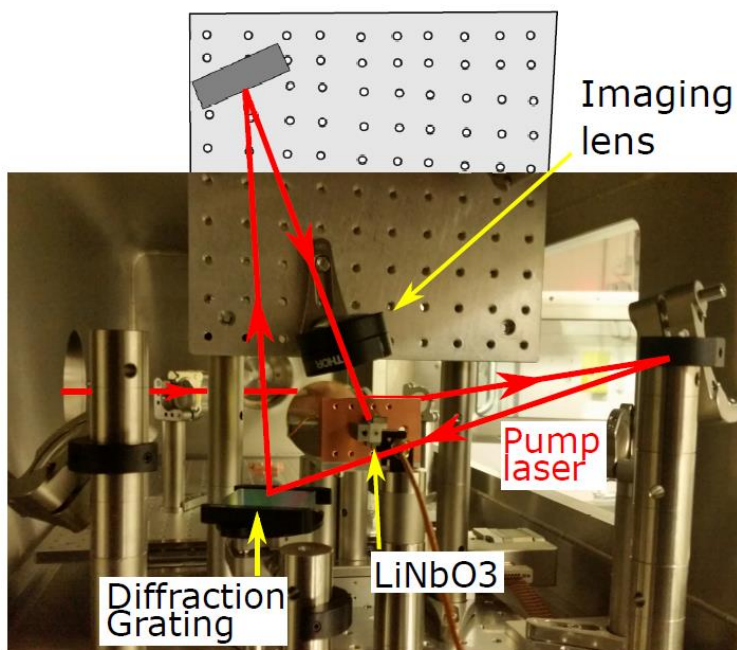
Our Approach: Travelling Wave THz Source

Proof of principle experiment

Showed Velocity matching possible over wide range

Accelerating fields behave as dispersion free

Requires complex 3d geometries to generate the THz when compared to waveguide schemes.



D.A. Walsh et al,
Nature Communications, September 2017

D.S. Lake, IOP Particle & Beams Group Annual Meeting 2020, September 2020

Our Approach: 100 keV Gun Experiment

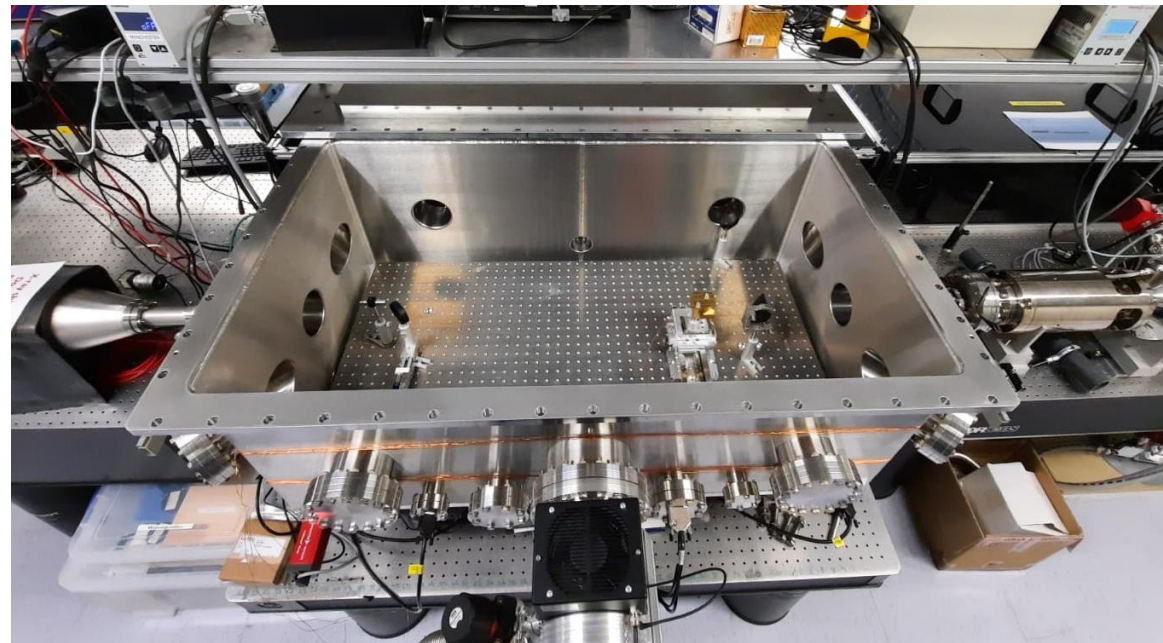
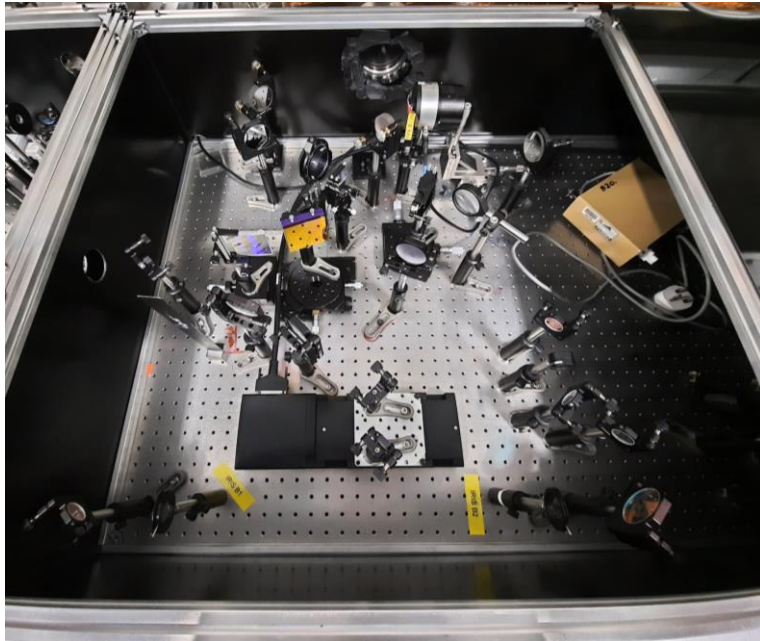
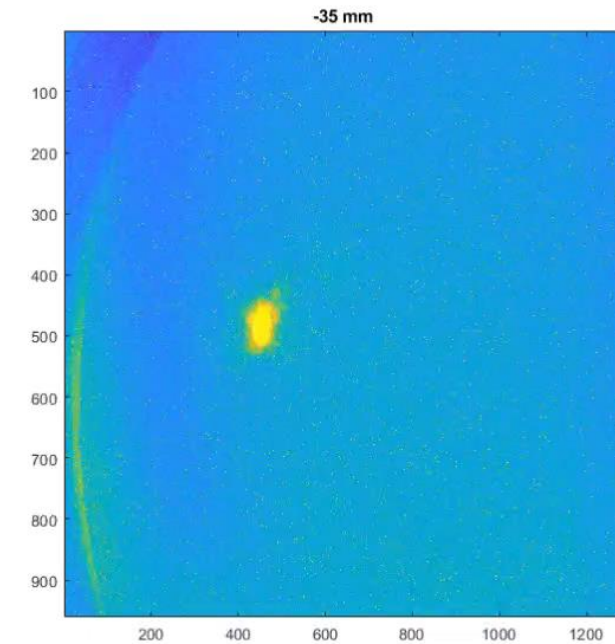
Similar to the experiment with the DLW and relativistic electrons

- 100 keV photo-electron gun
- Vacuum chamber designed for two interaction points
- MCP detector

→ Different waveguide matched to 100 keV electrons

Short term goals

- Terahertz driven electron deflection
- Terahertz driven compression
- Terahertz driven acceleration of compressed bunch



Summary

Demonstrated THz-driven linear acceleration of a 35 MeV relativistic electron beam

- Measurements performed at the CLARA test facility, Daresbury Laboratory
- Generation of narrowband THz pulses with required polarity-inversion
- Energy spectra modulation of chirped electron beams,
- Near-single bucket acceleration/deceleration of short bunches

Progress towards THz-driven compression and acceleration of non-relativistic bunches

- Early stages of experiments
- THz deflection observed
- Eventual goal of acceleration to from 100 keV to MeV and above