

# Electroweak processes in few-nucleon systems

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# Outline

- 1 Introduction
- 2 Local  $\chi$ EFT interaction
- 3 EM charge & current
- 4 Weak current
- 5 Quantum computation & Nuclear Physics
- 6 Conclusions & perspectives

## Collaborators

- R. Schiavilla *Jefferson Lab. & ODU, Norfolk (VA, USA)*
- M. Piarulli & S. Pastore *GWU, St.Louis (MS, USA)*
- L. Girlanda *University of Salento & INFN-Lecce, Lecce (Italy)*
- A. Kievsky & L.E. Marcucci - *INFN-Pisa & Pisa University, Pisa (Italy)*
- A. Lovato, B. Wiringa, & S. Pieper, *ANL, Argonne (IL, USA)*

## Chiral symmetry

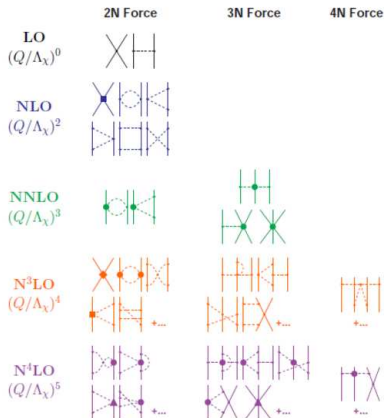
- $\mathcal{L}_{QCD}$  (almost) invariant under  $SU(2)$  “isospin” and “axial” rotations since  $m_u, m_d$  “small”
- Invariance also for *locals* rotations: Introducing external vector and axial currents
- Non-linear realization of the chiral symmetry for nucleons and nuclei [Weinberg, 1968, 1990], [CCWZ, 1969], [Gasser & Leutwyler, 1984], ...
- $\rightarrow \chi\text{EFT} = \text{Lagrangian written in terms of nucleonic and pionic d.o.f.} - \text{automatic inclusion of external vector and axial currents}$
- Infinite number of unknown coupling constants (low energy constants – LECs)
- All quantities can be written as powers of  $Q/\Lambda_\chi$ ,  $Q \sim$  small momenta or the pion mass,  $\Lambda_\chi \sim 1 \text{ GeV}$

## Potentials and currents from $\chi\text{EFT}$

- To obtain effective operators acting only on nucleons d.o.f.:
  - S-matrix: for a given process  $NN \rightarrow NN$  define  $V$  so that (on-shell)  
 $\langle NN | T_{\text{EFT}} | NN \rangle \equiv \langle NN | T_V | NN \rangle$
  - Unitary transformation: find  $U$  in order to decouple  $|NN\rangle$  Hilbert space from  $|NN\pi\rangle$ , etc.
- Alternative: Lattice  $\chi\text{EFT}$  [Lee *et al.*, 2010]

# NN & 3N interaction

For more information see for example  
 [Epelbaum, 2010], [Machleidt & Entem, 2011]



## NN interaction

- Jülich N4LO [Epelbaum, Krebs, & Meissner, 2014], [Reinert, Krebs, & Epelbaum, 2017]
- Idaho N4LO [Entem, Machleidt, & Nossyk, 2017]

LEC's fitted to the NN database or  $\pi N$  database

## 3N interaction

- N2LO [Epelbaum *et al.*, 2002]
- 3N force at N3LO & N4LO [Krebs *et al.*, 2012-2013]

- At N2LO there are two LECs  $c_D$  and  $c_E$ : fitted to some 3N expt. data (see later)
- At N3LO no new parameters
- At N4LO 10 new LECs [Girlanda *et al.*, 2011, 2019]

# Ingredients to study EW transitions

- Initial/final wave functions + transition operators (currents & charges)
- For transition to continuum: scattering wave functions
- Inclusive & semi-inclusive reactions: integral techniques (Lorentz integral transform – Laplace integral transform)

## Example

Transition  $|\alpha\rangle + \gamma \rightarrow |\beta\rangle$

$$\langle\beta|H_{e.m.}|\alpha; \mathbf{q}\lambda\rangle = \langle\Psi_\beta|\mathcal{K}_1|\Psi_\alpha\rangle \quad \mathcal{K}_1 = \frac{-e}{\sqrt{2\omega\Omega}} \int d\mathbf{x} e^{i\mathbf{q}\cdot\mathbf{x}} \hat{\epsilon}_{\mathbf{q}\lambda} \cdot \hat{\mathbf{J}}(\mathbf{x})$$

- $\mathcal{K}_1$  acts only on the nucleons' d.o.f.
- $|\alpha\rangle, |\beta\rangle$  initial & final nuclear states,  $\Psi_\alpha, \Psi_\beta$  corresponding w.f.
- $\mathbf{q}, \omega, \hat{\epsilon}_{\mathbf{q}\lambda}$  = momentum, energy, polarization of the absorbed photon
- for virtual photons, one needs also the m.e. of  $\hat{\mathbf{q}} \cdot \hat{\mathbf{J}}$  and  $\rho$

$$\mathbf{J}^\mu(\mathbf{q}) = \int d\mathbf{x} e^{i\mathbf{q}\cdot\mathbf{x}} \hat{\mathbf{J}}^\mu(\mathbf{x}) \quad \mu = 0, 1, 2, 3$$

## Current conservation (CC) $\nabla \cdot \widehat{\mathbf{J}}(\mathbf{x}) = -i[H, \rho(\mathbf{x})]$

- Strict interplay between  $H$ ,  $\widehat{\mathbf{J}}$  and  $\widehat{\rho}$  [Buchmann *et al.*, 1985], [Riska, 1989], [Schiavilla *et al.*, 1990]

$$\widehat{\rho}(\mathbf{x}) = \sum_{i=1}^A \frac{1 + \tau_z(i)}{2} \delta(\mathbf{r}_i - \mathbf{x}) \quad H = \sum_i T_i + \sum_{ij} V_{ij} + \dots \rightarrow \mathbf{J}(\mathbf{x}) = \sum_i \mathbf{J}_i^{(1)} + \sum_{ij} \mathbf{J}_{ij}^{(2)}(\mathbf{x}) + \dots$$

- $\mathbf{J}_{ij}^{(2)}(\mathbf{x}) =$  meson exchange currents
- Old approach:  $\mathbf{J}$  constructed not consistently with  $V$ ; CC verified “by hand” [Marcucci *et al.*, 2005]
- EFT approach:  $H$  and  $\mathbf{J}^\mu$  derived from the same Lagrangian [Park *et al.*, 1993], [Kolling *et al.*, 2009], [Pastore *et al.*, 2009], [Schiavilla *et al.*, 2018]
- However:
  - different cutoff used to regularize the short-range parts of  $V$  and  $\mathbf{J}$
  - different orders of chiral expansions of  $V$  and  $\mathbf{J}$
- In the present work: same regulators for  $V$  and  $\mathbf{J}$
- See also [Krebs *et al.*, 2019], for a more systematic approach

## Weak transitions

- Vector current: CVC hypothesis:  $V^\mu$  derived from  $J_{EM}^\mu$
- Axial current: PCAC (conservation in the limit  $m_\pi \rightarrow 0$ ) [Park *et al.*, 2003], [Baroni *et al.*, 2015–2016], [Krebs *et al.*, 2016]

# EFT potentials with $\Delta$ -particle d.o.f.

## Project

- NN/3N potentials and charge/currents derived from  $\chi$ EFT with contributions of the  $\Delta$ -particle d.o.f.
- $\rightarrow \chi$ EFT with nucleons, pions &  $\Delta$ 's: See, for example, [Bernard, Kaiser, & Meissner, 1995], [Hemmert, Holstein, & Kambor, 1998], [Krebs, Epelbaum, & Meissner, 2007]
- Further "condition": NN & 3N forces local in  $r$ -space [Piarulli *et al.*, 2014-2016]
- Another local  $r$ -space potential derived from  $\chi$ EFT [Gezerlis *et al.*, 2014], [Lynn *et al.*, 2017]

Case	NN	3N
LO $Q^0$		
NLO $Q^2$		
N2LO $Q^3$		
N3LO $Q^4$		

# Fit & Regularization

## Regularization in $r$ -space

- Diagrams with pion exchanges  $\sim \frac{1}{r^n}$   
[Epelbaum *et al.*, 2004], [Valderrama *et al.*, 2008]

- Regulated in  $r$  space with the function

$$C_{R_L}(r) = 1 - \frac{1}{(r/R_L)^6 e^{(r-R_L)/a_L} + 1}$$

- Contact terms
  - Using Fierz transformation to eliminate  $\nabla^2$  terms
  - Neglected some terms  $\sim \frac{d^2}{dr^2}$  giving small contributions
  - at the end:  $L^2$  and  $(L \cdot S)^2$  operators
  - Short-range part regularized with

$$\delta(\mathbf{r}) \longrightarrow C_{R_S}(r) = \frac{1}{\pi^{3/2} R_S^3} e^{-(r/R_S)^2}$$

## Fits & $\chi^2$

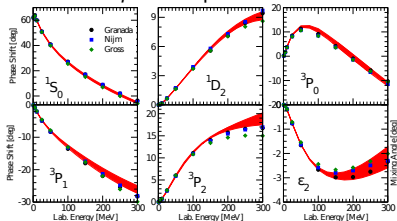
- $\sim 25$  LECs fitted to the NN database
- Interactions I**: Fit of data up to 125 MeV
- Interactions II**: Fit of data up to 200 MeV
- Choice  $(R_L, R_S) = (0.8, 1.2)$  fm:  
**models a**
- Choice  $(R_L, R_S) = (0.7, 1.0)$  fm:  
**models b**

model	$E_{\text{Lab}}$ (MeV)	$N_{pp+np}$	$\chi^2$
<i>la</i>	0–125	2668	1.05
<i>lb</i>	0–125	2665	1.07
<i>IIa</i>	0–200	3698	1.37
<i>IIb</i>	0–200	3695	1.37

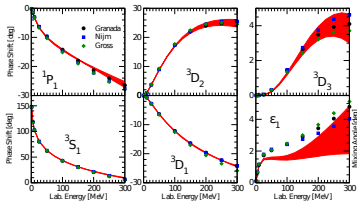


# $NN$ phase-shifts and deuteron radial components

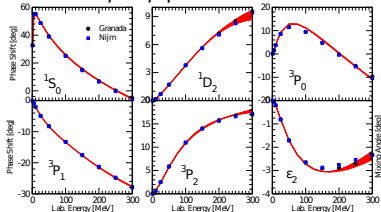
$n - p T = 1$  phase-shifts



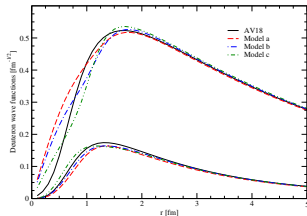
$n - p T = 0$  phase-shifts



$p - p$  phase-shifts



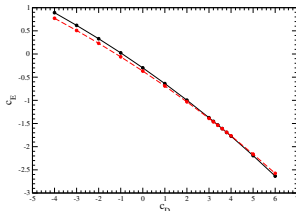
Deuteron radial components



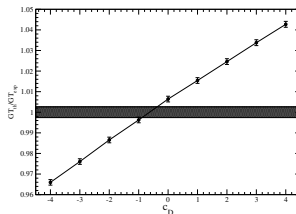
# 3N force – choice of $c_D$ & $c_E$

- For  $A > 2$  inclusion of the 3NF at N2LO [Epelbaum *et al.*, 2002] + Fujita-Miyazawa
- Two LECs to be fixed:  $c_D$  &  $c_E$
- Method 1:  $B(^3\text{H})$  and  $a_{n-d}^{(2)}$  Ia, Ib, IIa, IIb
- Method 2:  $B(^3\text{H})$  and Gamow-Teller (GT) matrix element Ia\*, Ib\*, IIa\*, IIb\*
- Bound- and continuum..states  $A = 3$  calculations performed using the HH technique [Kievsky *et al.*, 2008]

Fit of  $B(^3\text{H})$



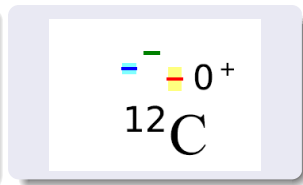
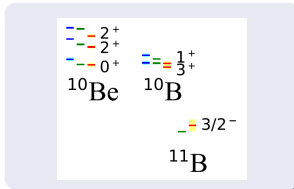
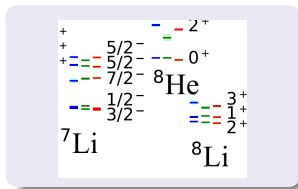
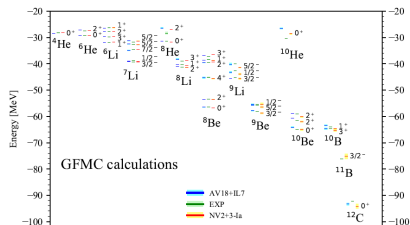
Fit of the GT matrix element of tritium beta decay



	Ia*	Ib*	IIa*	IIb*
$c_D$	-0.635	-4.71	-0.61	-5.25
$c_E$	-0.09	0.55	-0.35	0.05

# GFMC spectrum of light nuclei

NN+3N interaction: *la* [Piarulli *et al.*, 2018]



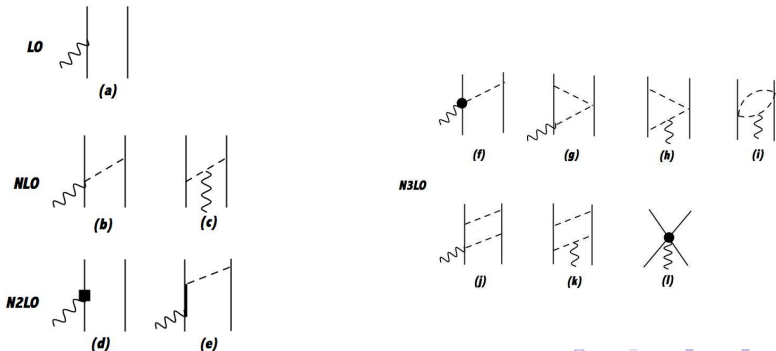
Nice reproduction of the energy levels – **3N force fitted using only  $A = 3$  data!!!**  
 but too soft neutron matter EOS – studies with  $la^* - llb^*$  in progress

# EM transitions

## EM charge & current from $\chi$ EFT

[Park *et al*, 1993], [Kolling *et al*, 2009], [Pastore *et al*, 2009]  
Including the  $\Delta$  d.o.f. [Schiavilla *et al.*, 2018]

Diagrams with the inclusion of the  $\Delta$  d.o.f. up to N2LO



# Determination of the LECs

## 3 new LECs $d_1^S$ , $d_2^S$ , $d_1^V$

- $d_1^S$ ,  $d_2^S$  multiply isoscalar operators
- $d_1^V$  multiplies an isovector operator
- Fitted to  $\mu_d$ ,  $\mu_{3H}$ , and  $\mu_{3He}$

	la*	lb*	lla*	llb*
$d_1^S$	-0.00999	-0.02511	-0.01170	-0.04955
$d_2^S$	-0.06571	-0.02384	-0.04714	-0.07947
$d_1^V$	-0.05120	-0.03509	-0.05128	-0.03880

Deuteron magnetic moment expt. 0.8574 n.m.

	la*	lb*	lla*	llb*
LO	0.8498	0.8485	0.8501	0.8501
N2LO(RC)	-0.0062	-0.0061	-0.0065	-0.0072
N3LO	0.0137	0.0151	0.0138	0.0145
TOT	0.8573	0.8575	0.8574	0.8574

Isoscalar magnetic moment expt. 0.4257 n.m.

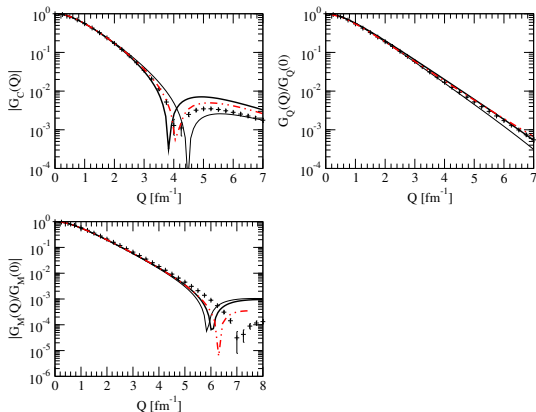
	la*	lb*	lla*	llb*
LO	0.4089	0.4075	0.4091	0.4089
NLO	0.0015	0.0020	0.0012	0.0018
N2LO	-0.0062	-0.0043	-0.0052	-0.0071
N3LO	0.0229	0.0215	0.0218	0.0231
TOT	0.4271	0.4267	0.4269	0.4267

Isvector magnetic moment expt -2.553 n.m.

	la*	lb*	lla*	llb*
LO	-2.1823	-2.1755	-2.1815	-2.1787
NLO	-0.1967	-0.2257	-0.1967	-0.2255
N2LO	-0.0388	-0.0657	-0.0395	-0.0617
N3LO	-0.1355	-0.0864	-0.1354	-0.0872
TOT	-2.5533	-2.5533	-2.5531	-2.5531

# Deuteron form factors

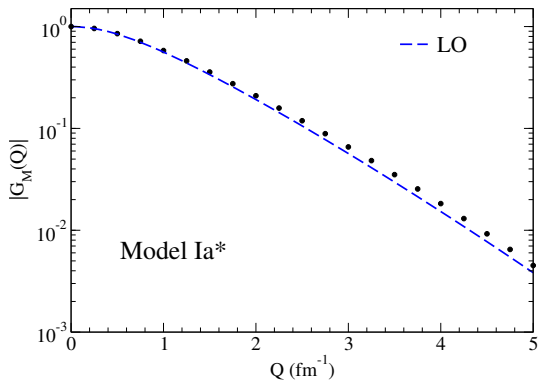
Conventional approach [Marcucci *et al.*, 2015]  
AV18/UIX interaction + meson exchange current [Marcucci *et al.*, 2005]



- Thin black line: IA results
- Thick black line: IA+MEC (FULL) results
- Dash-double-dotted red line: approximate inclusion of relativistic effects [Friar, 1975]
- Experimental data: see the review paper [Marcucci *et al.*, 2015]

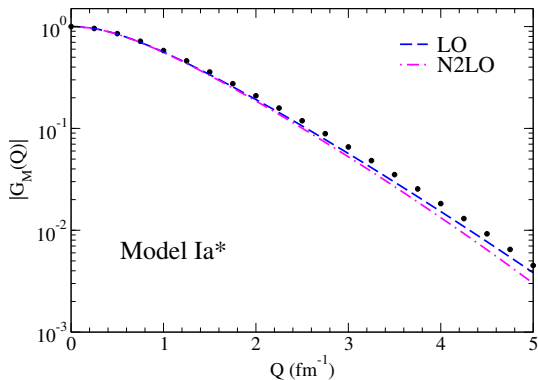
# Deuteron magnetic form factor (1)

Dependence on the current orders (cumulative)  
Interaction Ia\*



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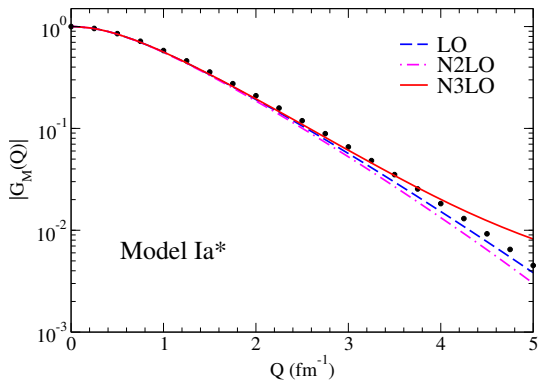
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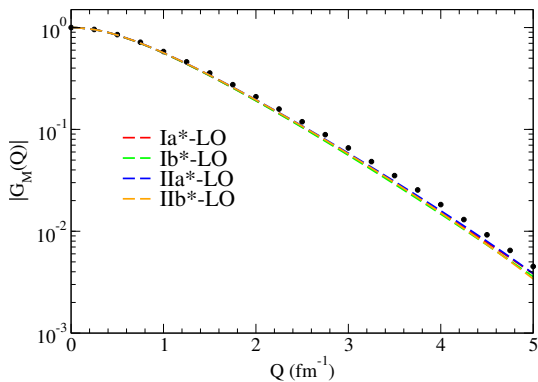
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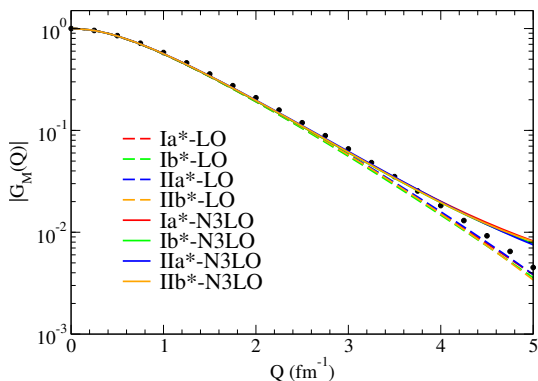
# Deuteron magnetic form factor (2)

Dependence on the NN interaction



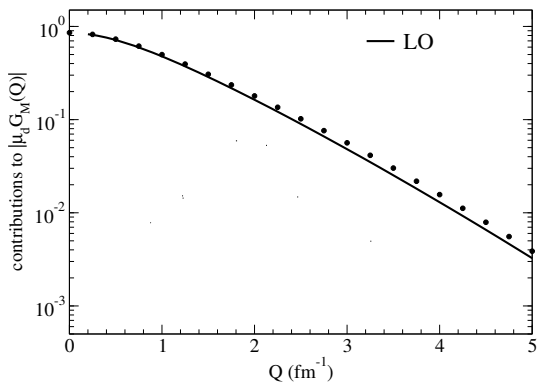
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# Deuteron magnetic form factor (3)

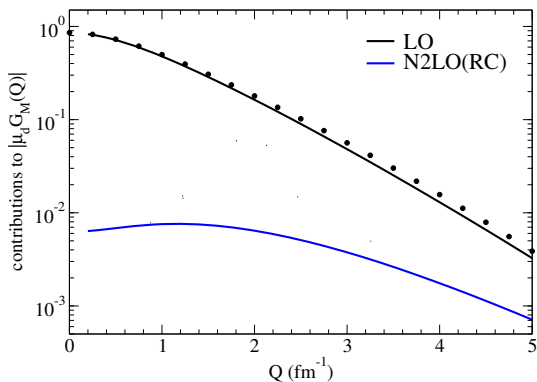
Contribution of the current orders  
Interaction Ia\*



- Problems with some N3LO terms
- "NM"= non minimal  $\sim d_i^S$

# Deuteron magnetic form factor (3)

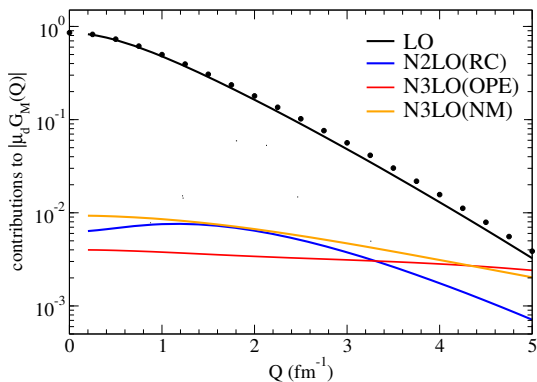
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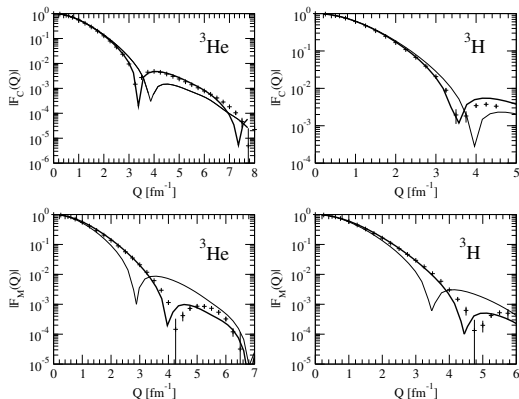
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# Trinucleon form factors

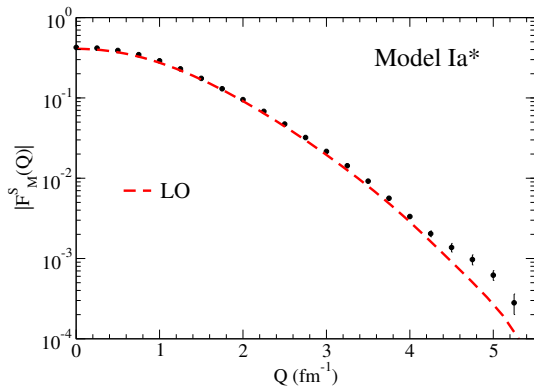
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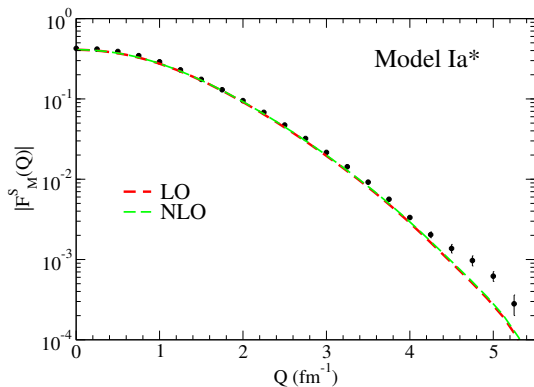
Calculation with the new interactions/currents  
Isoscalar magnetic form factor





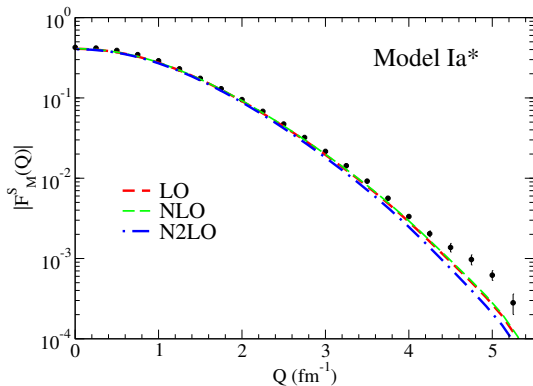
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Calculation with the new interactions/currents  
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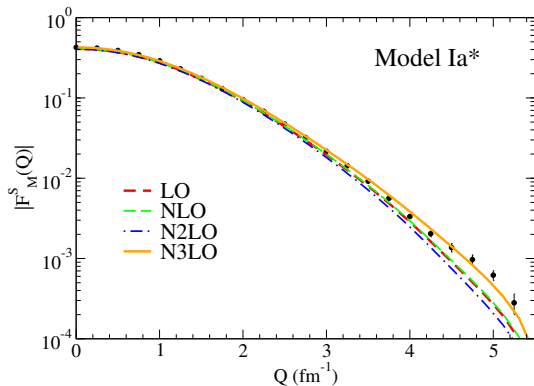
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Calculation with the new interactions/currents  
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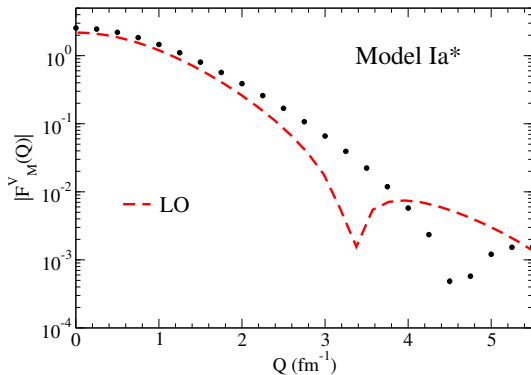
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Calculation with the new interactions/currents  
Isoscalar magnetic form factor



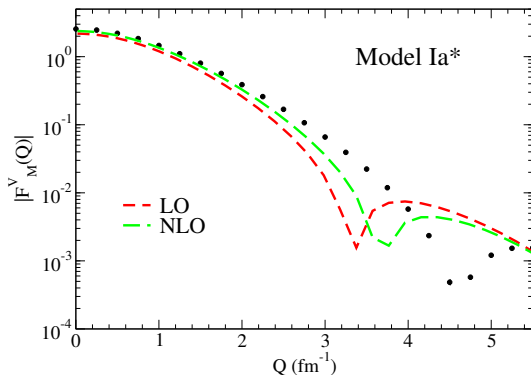
# Trinucleon magnetic form factors (2)

Calculation with the new interactions/currents  
Isovector magnetic form factor



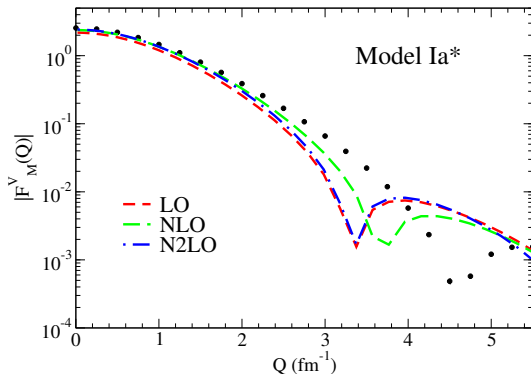
# Trinucleon magnetic form factors (2)

Calculation with the new interactions/currents  
Isovector magnetic form factor



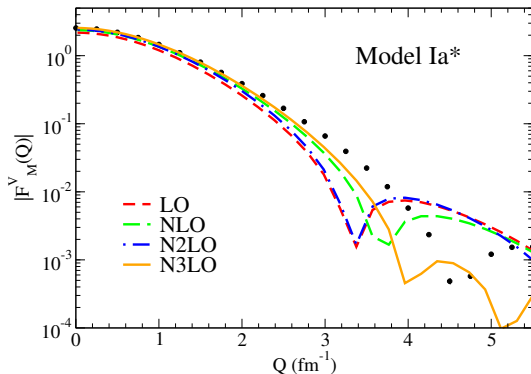
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Calculation with the new interactions/currents  
Isovector magnetic form factor



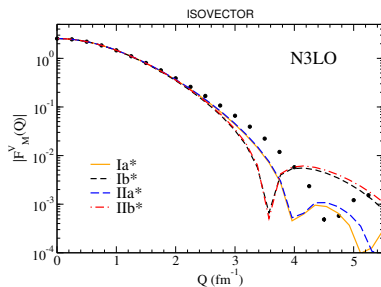
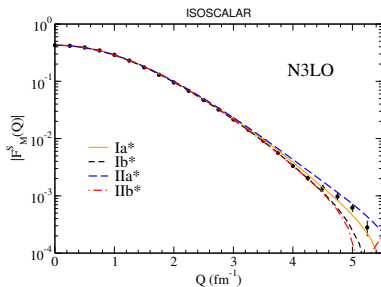
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Calculation with the new interactions/currents  
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# Trinucleon magnetic form factors (3)

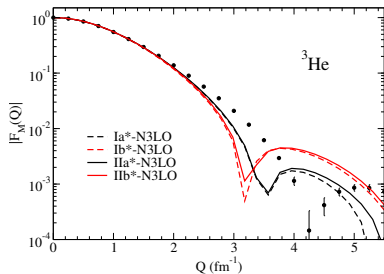
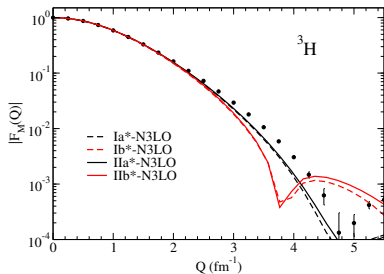
Calculation with the new interactions/currents  
Dependence on the interaction





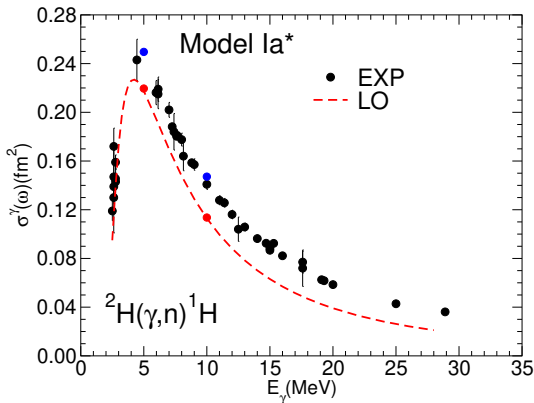
# Trinucleon magnetic form factors (4)

Calculation with the new interactions/currents  
 ${}^3\text{H}$  &  ${}^3\text{He}$  form factors



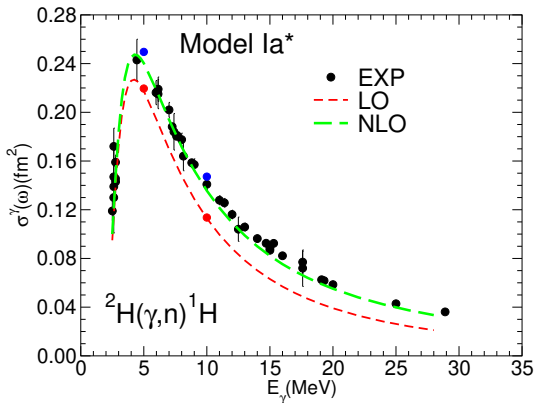
# Phodisintegration of the deuteron (1)

Calculation with the new interactions/currents  
Interaction Ia\*



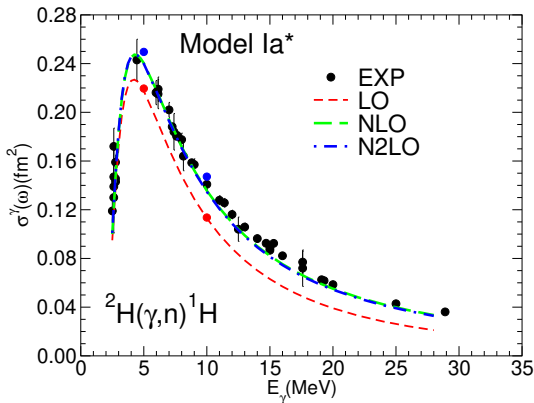
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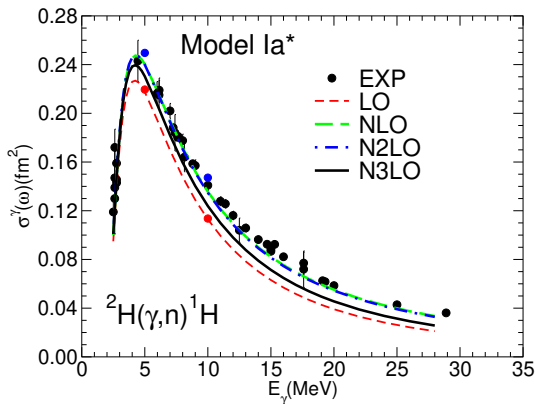
# Phodisintegration of the deuteron (1)

Calculation with the new interactions/currents  
Interaction Ia\*



# Phodisintegration of the deuteron (1)

Calculation with the new interactions/currents  
Interaction  $la^*$



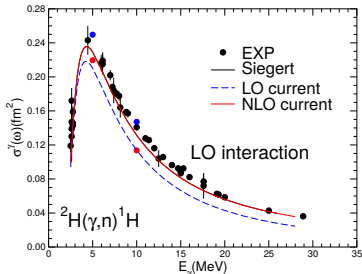
# Phodisintegration of the deuteron (2)

Multipole expansion of the EM current operator

$$T_{J,M}^E = -i \frac{q^{J-1}}{(2J+1)!!} \sqrt{\frac{J+1}{J}} \int d^3x \left[ \nabla \cdot \mathbf{J}_C(\mathbf{x}) + \frac{q^2}{J+1} \nabla \cdot (\mathbf{x} \times \boldsymbol{\mu}(\mathbf{x})) \right] x^J Y_{JM}(\hat{\mathbf{x}})$$

## Current conservation (CC)

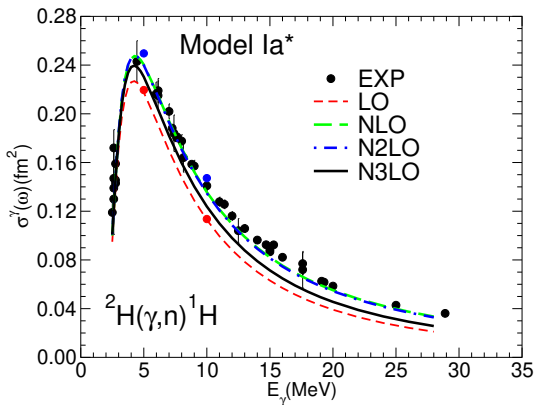
$$\nabla \cdot \mathbf{J}(\mathbf{x}) + i [H_0, \rho(\mathbf{x})] = 0 \quad \langle J_f, M_f | [H_0, \rho(\mathbf{x})] | J_i, M_i \rangle = (E_f - E_i) \langle J_f, M_f | \rho(\mathbf{x}) | J_i, M_i \rangle$$



- This observable is dominated by the  $E_1$  multipole
- $V$  at LO (one-pion exchange + LO contact terms)
- $J$  at LO+NLO (single nucleon current + one-pion exchange diagrams)
- In this case CC is exactly verified
- However, if  $V$  and  $J$  do not verify CC,  $E_\ell$  not well calculated
- [Siegert, 1937]: compute  $E_1$  forcing CC

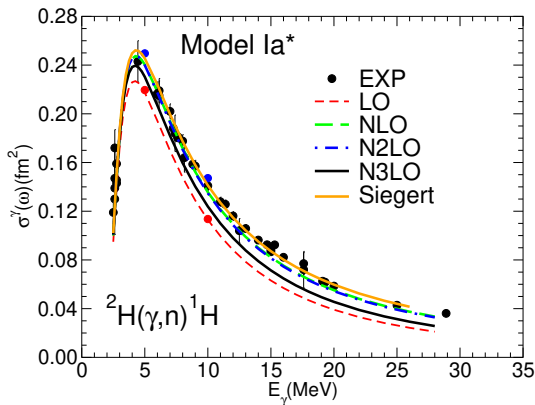
# Phodisintegration of the deuteron (3)

Calculation with the new interactions/currents  
Interaction  $1a^*$  Using Siegert



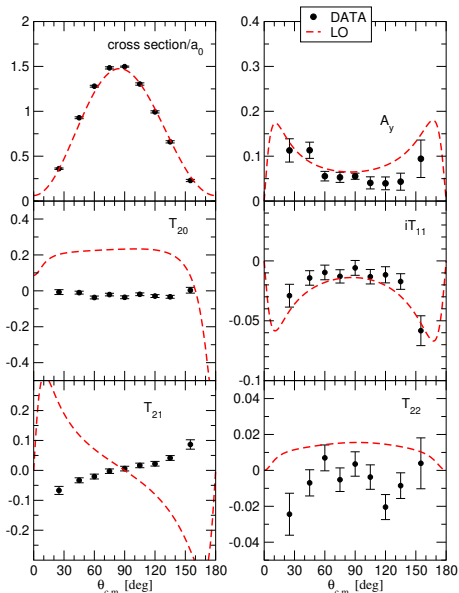
# Phodisintegration of the deuteron (3)

Calculation with the new interactions/currents  
Interaction  $la^*$  Using Siegert





# Proton-deuteron capture



PRELIMINARY

AV18 + new EM currents

$E_{cm} = 2 \text{ MeV}$

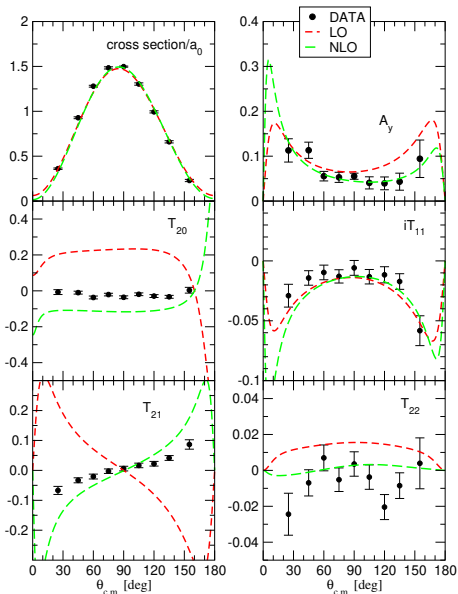
Data: [Smith & Knutson,  
1999]

Big effects from

$$\langle \Psi_{3\text{He}} | j_{EM} | \Psi_{pd}^{L=1, S=3/2, J} \rangle$$

Calculations with the new  
interactions in progress

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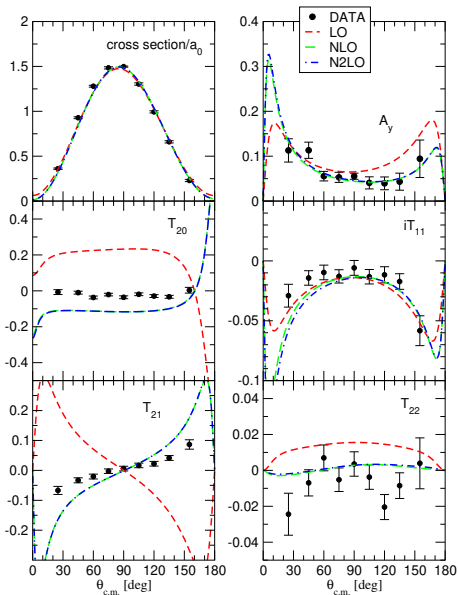
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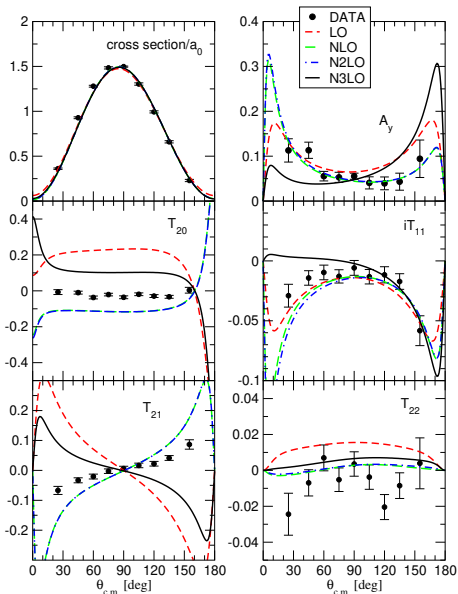
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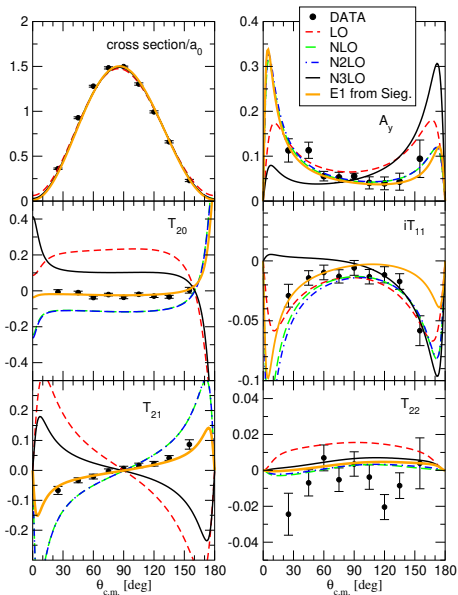
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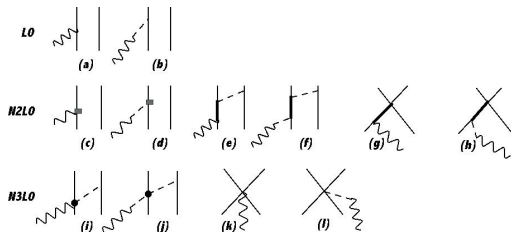
# Weak current

## Axial current from $\chi$ EFT

[Park *et al.*, 2003], [Baroni *et al.*, 2015], [Krebs *et al.*, 2016]

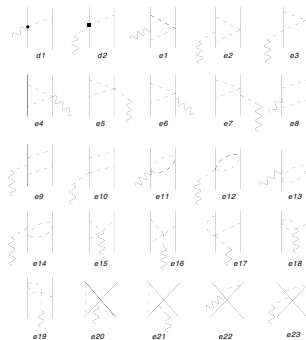
Including the  $\Delta$  d.o.f.: [Baroni *et al.*, 2018]

## N4LO diagrams



Single nucleon current diagram NLO  $\sim$  LEC

$d_R$



# Tritium beta decay

$$d_R = -\frac{m_\pi}{4 g_A \Lambda_\chi} c_D + \frac{m_\pi}{3} \left( c_3 + c_3^\Delta + 2 c_4 + 2 c_4^\Delta \right) + \frac{m_\pi}{6 m}$$

[Gardestig and Phillips, 2006], [Gazit *et al.*, 2009], [Schiavilla, 2018]

Fit of  $B(^3\text{H})$  and the GT matrix element in tritium  $\beta$  decay

la/b\* & lla/lb\* chiral Hamiltonians [Baroni *et al.*, 2018]

	la*	lb*	lla*	llb*
$c_D$	-0.635	-4.71	-0.61	-5.25
$c_E$	-0.09	0.55	-0.35	0.05
LO	0.9272	0.9247	0.9261	0.9263
N2LO	0.0345	0.0517	0.0345	0.0515
N3LO	-0.0108	-0.0261	-0.0102	-0.0272
TOT	0.9509	0.9503	0.9504	0.9506

Experimental value  $0.9511 \pm 0.0013$  (see [Baroni *et al.*, 2016])

Range of  $c_D$  and  $c_E$  values allowed by the experimental error on  $\text{GT}_{\text{exp}}$

	la*	lb*	lla*	llb*
$c_D$	(-0.89, -0.38)	(-4.99, -4.42)	(-0.89, -0.33)	(-5.56, -4.94)
$c_E$	(-0.01, -0.17)	(+0.70, +0.40)	(-0.25, -0.45)	(+0.23, -0.13)

# N4LO contribution

Contributions obtained with the [Baroni *et al.*, 2016] and [Krebs *et al.*, 2017] formulations of the N4LO axial current

	la*	lb*	lla*	llb*
N4LO(B)	-0.0672	-0.0732	-0.0671	-0.0716
N4LO(K)	-0.0364	-0.0540	-0.0365	-0.0543
B-K(OPE)	0.0141	0.0196	0.0142	0.0201
B-K(TPE)	0.0018	0.0024	0.0018	0.0025
B-K(CT)	-0.0467	-0.0412	-0.0466	-0.0399

- OPE = N4LO loop corrections to the OPE axial current
- TPE = N4LO genuine new TPE contributions
- CT = N4LO contact contributions induced by the regularization scheme in configuration space we have adopted
- Contributions at N4LO found to be relatively large and of opposite sign than those at LO
- Difference between the two formalisms at present not clarified



# Quantum computation & Nuclear Physics

A. Roggero [INT, Seattle, WA (USA)] & A. Baroni arXiv:1905:08383  
see also [Roggero & Carlson, 2018], [Dumitrescu *et al.*, 2018]

## Quantum computers

- Example: 1 qubit = spin of an electron
  - spin up =  $|1\rangle$ , spin down =  $|0\rangle$
  - but are also possible all the superposition states  $\alpha|1\rangle + \beta|0\rangle\dots$
- Gates: perform simple operation on the qubit
  - X,Y,Z gates = Pauli matrix operations
  - Hadamard gate

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \sim X + Z$$

- “Ancillary” qubit = qubit in a definite known state

arXiv:1905:08383

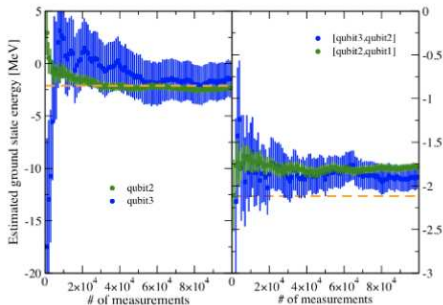
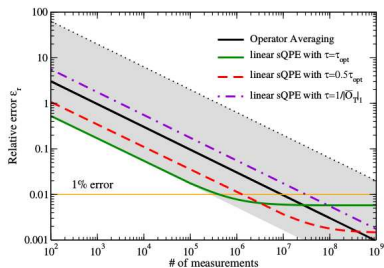
- Deuteron  $\Psi = \alpha|S\rangle + \beta|D\rangle$
- The deuteron can be described by a single qubit!
- Using a “normal” computer, compute the matrix elements

$$\begin{bmatrix} \langle S|H|S\rangle & \langle S|H|D\rangle \\ \langle D|H|S\rangle & \langle D|H|D\rangle \end{bmatrix} = aI + bX + cZ$$

- Potential used: Argonne V6
- Quantum gates  $\rightarrow H\Psi$
- Run on an IBM simulator
- Problems: number of iterations, noise, ...

# The deuteron on a quantum computer

A. Roggero [INT, Seattle, WA (USA)] & A. Baroni arXiv:1905:08383



## Conclusions

- Ongoing effort to construct interactions & EW currents from  $\chi$ EFT but local in configuration space
- Structure and EW processes in light nuclei using the HH & GFMC techniques
- Presented results for
  - Magnetic moments and FFs up to  $A = 3$
  - Photo-disintegration of the deuteron
  - GT matrix element of tritium  $\beta$  decay
- Still something to be clarified ...
  - Better quantification of “theoretical error”
  - Benchmark calculations

## Perspectives

- Work in progress
  - FF of  ${}^4\text{He}$
  - $p - d$  &  $d - d$  radiative capture at BBN energies (LUNA experiment)
  - $p - p$  &  $p - {}^3\text{He}$  astrophysical factors
  - ${}^6\text{He}$  beta decay
  - ...
- Further advances:
  - Test of the interactions in  $p - {}^3\text{He}$  elastic scattering
  - Inclusion of the N4LO contact interaction in the 3N force
- Quantum computing of scattering processes  $n - p$ ,  $n - d$ , etc.

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Thank you for your attention!