Weakly bound nuclei:

A unified description of intrinsic and relative degrees of freedom

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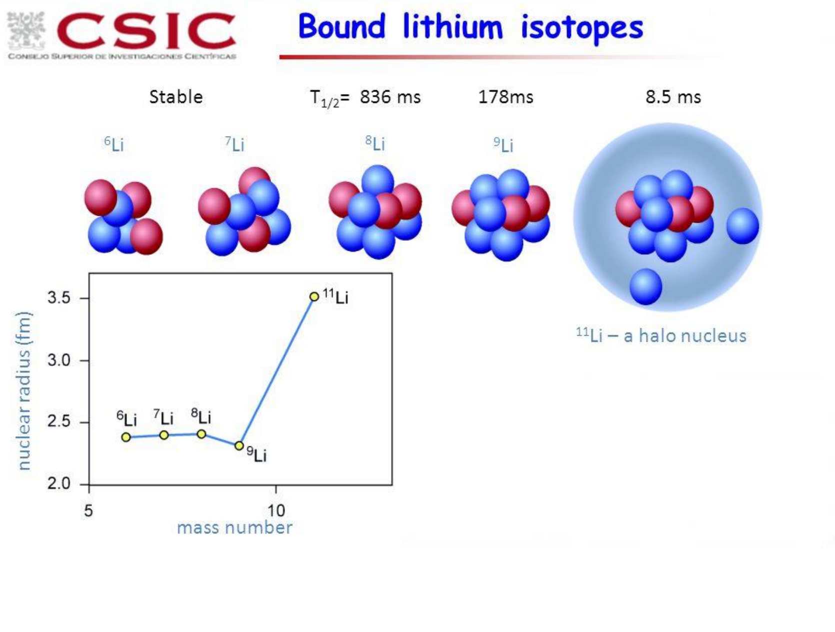

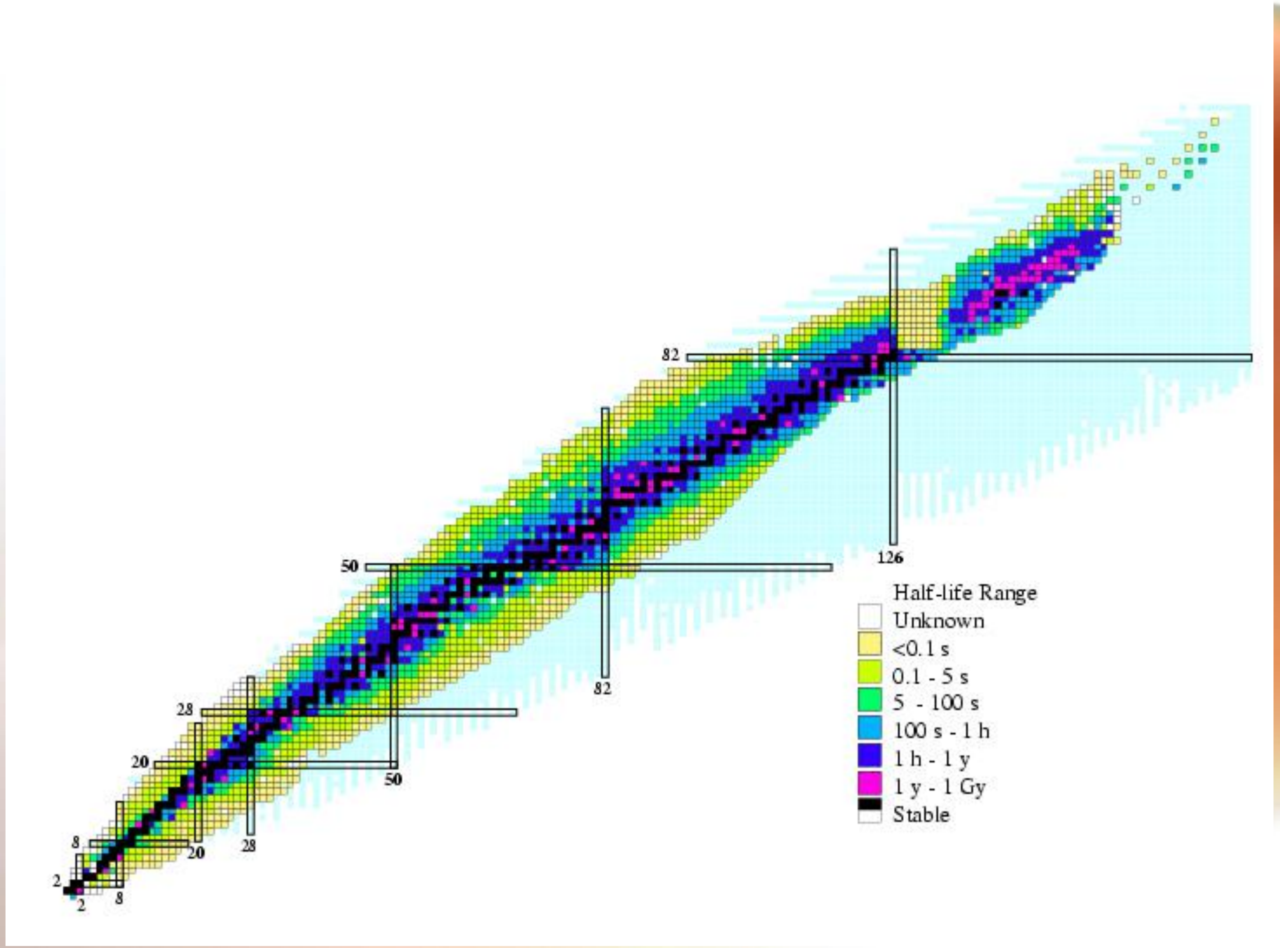

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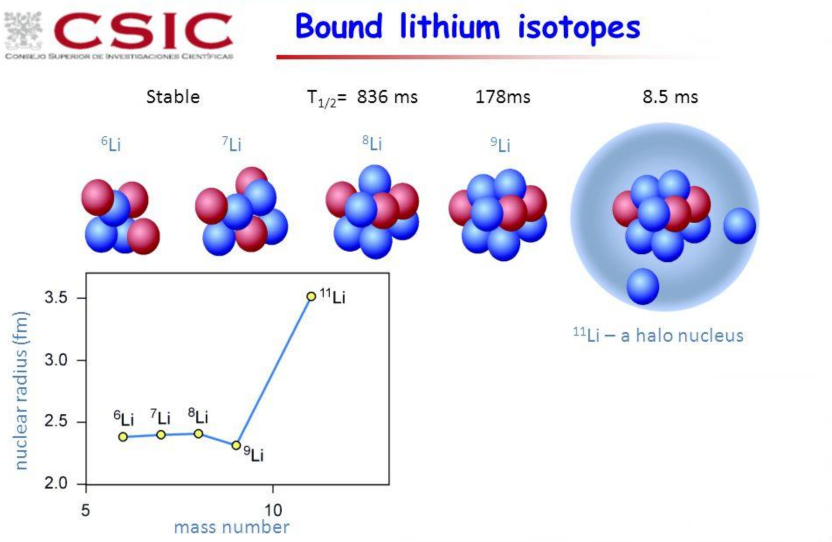
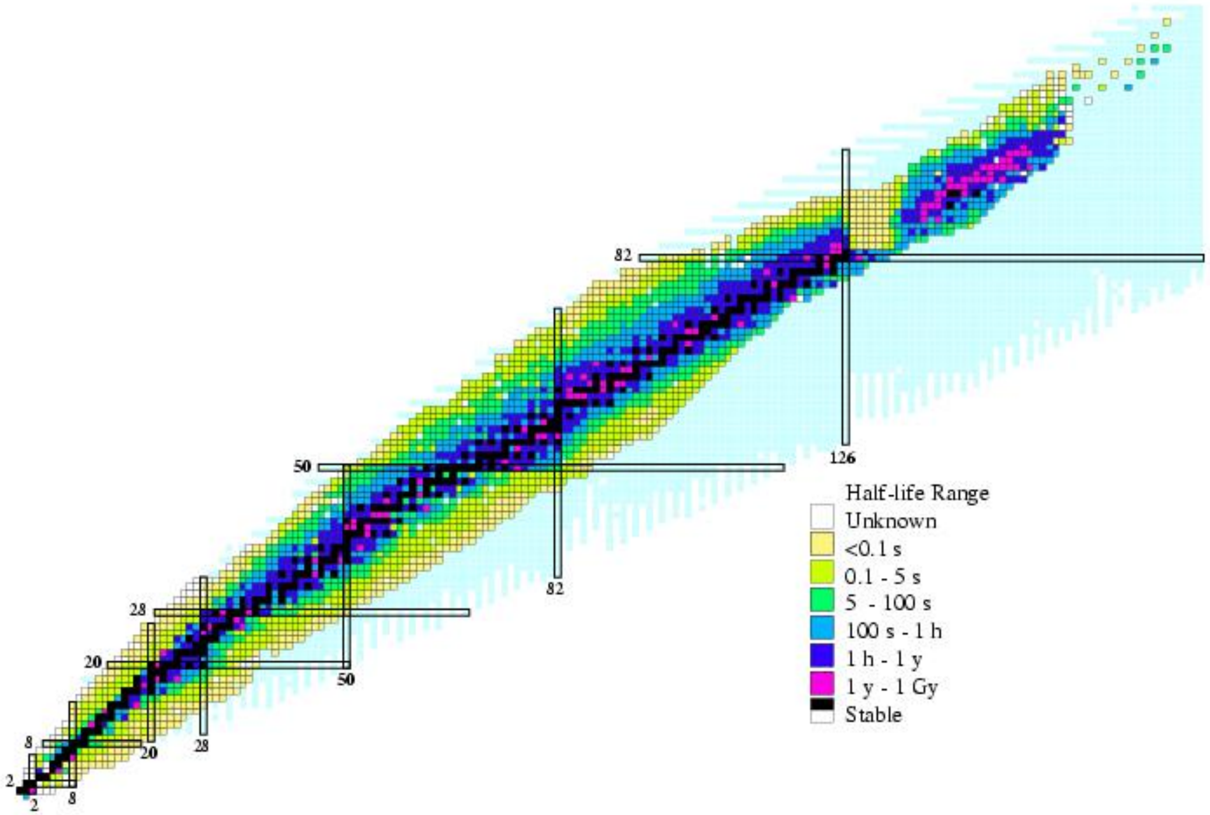


$$\left[T_x + T_y + V_{12}(r_{12}) + V_{13}(r_{13}) + V_{23}(r_{23})\right]\Psi = E\Psi$$



- ✓ The core is assumed to be an inert particle.
- ✓ What to do when experimental information is not available.

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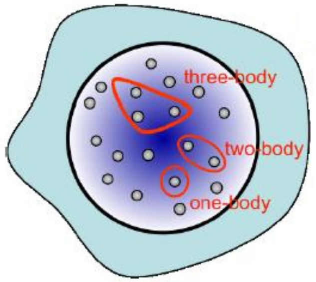

$$H = \sum_i T(i) + \sum_{i<j} V^{(2)}(i,j) + \sum_{i<j<k} V^{(3)}(i,j,k)$$

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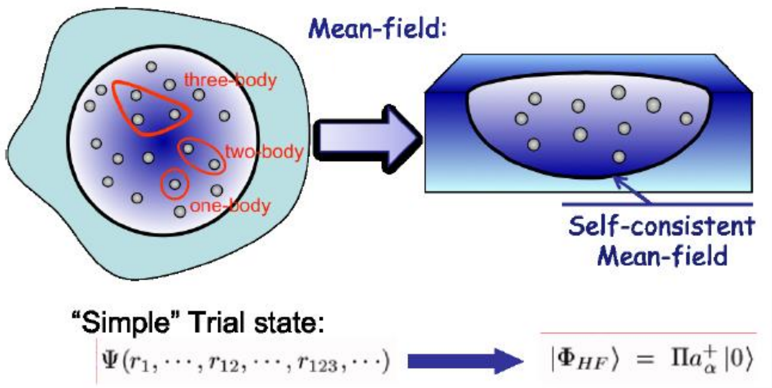

$$H = \sum_i T(i) + \sum_{i<j} V^{(2)}(i,j) + \sum_{i<j<k} V^{(3)}(i,j,k)$$

- ✓ The particles do not interact with each other, but through an average mean-field.
- ✓ The complex N-body wave function is replaced by a Slater determinant.

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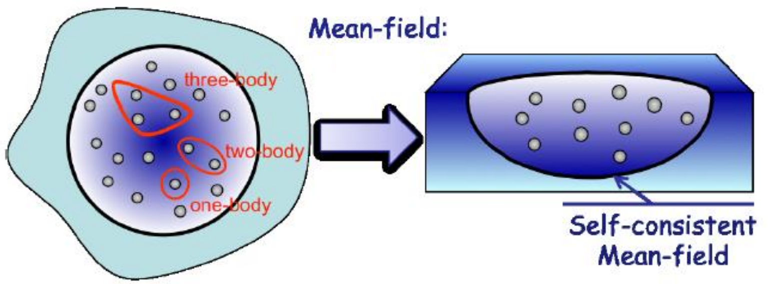

"Simple" Trial state:

$$\Psi(r_1, \cdots, r_{12}, \cdots, r_{123}, \cdots) \longrightarrow |\Phi_{HF}\rangle = \Pi a_\alpha^+ |0\rangle$$

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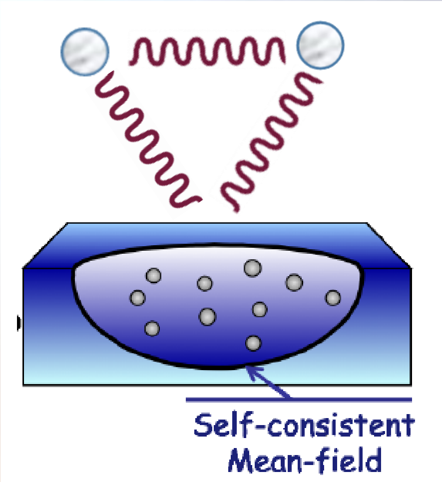

- ✓ The mean-field determines the interaction with the loosely bound nucleons.

Three-body problem

Weakly bound nuclei:

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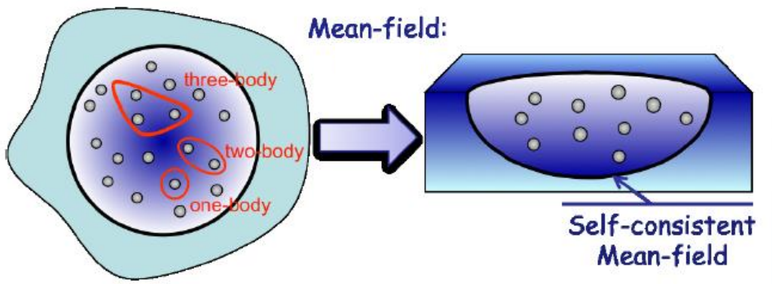

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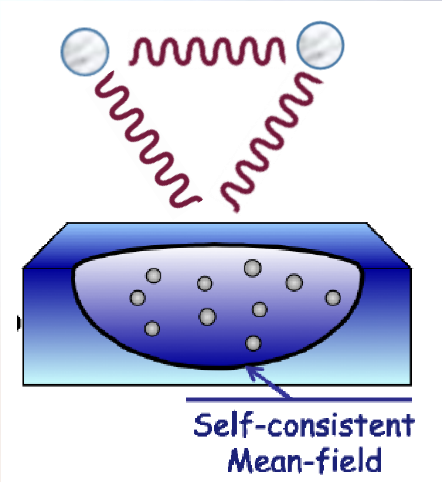

- ✓ The mean-field determines the interaction with the loosely bound nucleons.

Three-body problem

- ✓ The presence of the loosely bound nucleons distorts the mean-field.
- ✓ The distorted mean-field modifies the core-nucleon interaction and the wave function of the valence nucleons.

Self-consistent calculation

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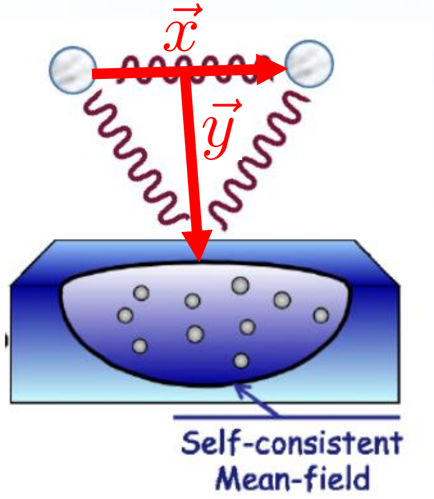
- ✓ Some formal hints about the formalism
- ✓ The case of ^{26}O
- ✓ Proton dripline: ^{70}Kr
- ✓ Approaching the dripline: Ca isotopes
- ✓ Summary and possible extensions

Formalism:



$$H = \underbrace{\frac{1}{4mA}\sum_{i=1}^{A}\sum_{j=1}^{A}(\vec{p}_i - \vec{p}_j)^2 + \sum_{i<j}^{A} V_{ij}}_{H_c} + \underbrace{\frac{p_x^2}{2\mu_x} + \frac{p_y^2}{2\mu_y} + V_{v_1 v_2}}_{H_3} + \underbrace{\sum_{i=1}^{A}(V_{iv_1} + V_{iv_2})}_{H_{coup}}$$

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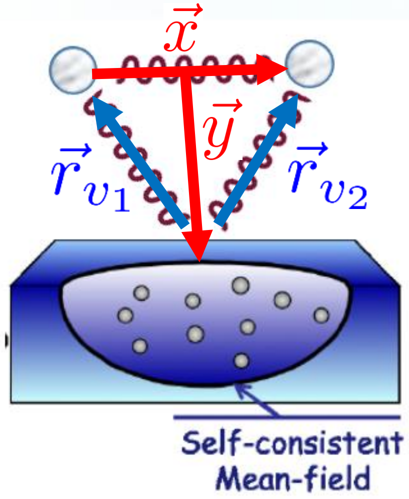

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$$\begin{aligned}
\Psi &= \mathcal{A}\{\Phi_c(\{\vec{r}_A\})\psi_3(\vec{r}_{v_1}, \vec{r}_{v_2})\} \\
&= \Phi_c(\{\vec{r}_A\})\psi_3(\vec{r}_{v_1}, \vec{r}_{v_2}) - \sum_{i=1}^{A}\Phi_c(\vec{r}_{v_1}, \{\vec{r}_{A-1}\})\psi_3(\vec{r}_i, \vec{r}_{v_2}) \\
&- \sum_{i=1}^{A}\Phi_c(\vec{r}_{v_2}, \{\vec{r}_{A-1}\})\psi_3(\vec{r}_{v_1}, \vec{r}_i) + \sum_{i<j}^{A}\Phi_c(\vec{r}_{v_1}, \vec{r}_{v_2}, \{\vec{r}_{A-2}\})\psi_3(\vec{r}_i, \vec{r}_j)
\end{aligned}$$

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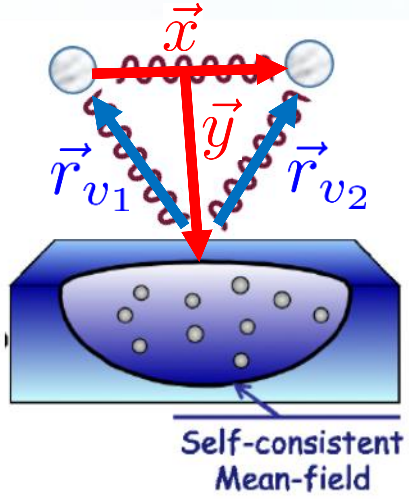

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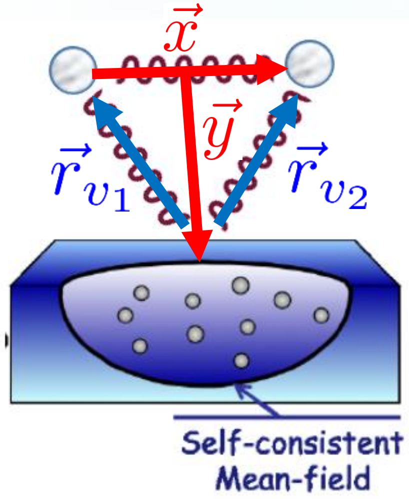
$$\Psi = \mathcal{A}\{\Phi_c(\{\vec{r}_A\})\psi_3(\vec{r}_{v_1}, \vec{r}_{v_2})\}$$

$$E = \langle\Psi|H|\Psi\rangle = \langle\Phi_c|H_c|\Phi_c\rangle + \langle\psi_3|H_3|\psi_3\rangle + \langle\Psi|H_{coup}|\Psi\rangle$$

Formalism:



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$$\Psi = \mathcal{A}\{\Phi_c(\{\vec{r}_A\})\psi_3(\vec{r}_{v_1}, \vec{r}_{v_2})\}$$

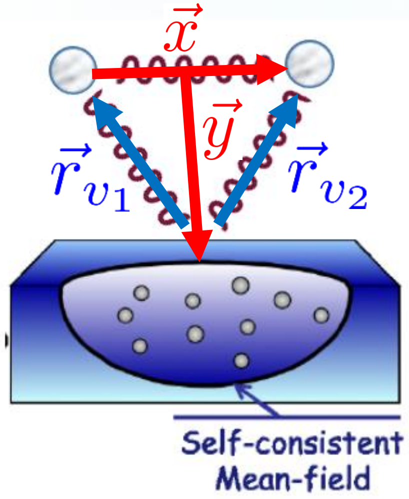
$$\langle\Psi|H'|\Psi\rangle = \langle\Psi|H|\Psi\rangle - E_c \int |\Phi_c(\{\vec{r}_A\})|^2 d\vec{r}_1 \cdots d\vec{r}_A - E_3 \int |\psi_3(\vec{r}_{v_1}, \vec{r}_{v_2})|^2 d\vec{r}_{v_1} d\vec{r}_{v_2}$$

$$0 = \frac{\delta}{\delta\Phi_c^*}\langle\Psi|H'|\Psi\rangle$$

$$0 = \frac{\delta}{\delta\psi_3^*}\langle\Psi|H'|\Psi\rangle$$

$$\Phi_c(\{\vec{r}_A\}) = \det(\{\phi_i^{q_i}(\vec{r}_i)\})$$

Formalism:



$$\begin{aligned}
\epsilon_i \phi_i^{q_i}(\vec{r}) &= -\frac{\hbar^2}{2mA}\sum_{k=1}^{A}\int \phi_k^{q_k *}(\vec{r'})\left(\vec{\nabla}_r - \vec{\nabla}_{r'}\right)^2\left(\phi_k^{q_k}(\vec{r'})\phi_i^{q_i}(\vec{r}) - \phi_k^{q_k}(\vec{r})\phi_i^{q_i}(\vec{r'})\delta_{q_i q_k}\right) d\vec{r'} \\
&+ \sum_{k=1}^{A}\int \phi_k^{q_k *}(\vec{r'})V_{ik}(\vec{r},\vec{r'})\left(\phi_k^{q_k}(\vec{r'})\phi_i^{q_i}(\vec{r}) - \phi_k^{q_k}(\vec{r})\phi_i^{q_i}(\vec{r'})\delta_{q_i q_k}\right) d\vec{r'} \\
&+ \int \psi_3^*(\vec{r}_{v_1},\vec{r}_{v_2})V_{v_1 i}(\vec{r}_{v_1},\vec{r})\left(\psi_3(\vec{r}_{v_1},\vec{r}_{v_2})\phi_i^{q_i}(\vec{r}) - \psi_3(\vec{r},\vec{r}_{v_2})\phi_i^{q_i}(\vec{r}_{v_1})\delta_{q_i q_{v_1}}\right) d\vec{r}_{v_1} d\vec{r}_{v_2} \\
&+ \int \psi_3^*(\vec{r}_{v_1},\vec{r}_{v_2})V_{v_2 i}(\vec{r}_{v_2},\vec{r})\left(\psi_3(\vec{r}_{v_1},\vec{r}_{v_2})\phi_i^{q_i}(\vec{r}) - \psi_3(\vec{r}_{v_1},\vec{r})\phi_i^{q_i}(\vec{r}_{v_2})\delta_{q_i q_{v_2}}\right) d\vec{r}_{v_1} d\vec{r}_{v_2}
\end{aligned}$$

$$\begin{aligned}
E_3\psi_3(\vec{r}_{v_1},\vec{r_{v_2}}) &= \left(\frac{p_x^2}{2\mu_x} + \frac{p_y^2}{2\mu_y} + V_{v_1 v_2}\right)\psi_3(\vec{r}_{v_1},\vec{r_{v_2}}) \\
&+ \sum_{i=1}^{A}\int \phi_i^{q_i *}(\vec{r})V_{v_1 i}(\vec{r}_{v_1},\vec{r})\left(\psi_3(\vec{r}_{v_1},\vec{r}_{v_2})\phi_i^{q_i}(\vec{r}) - \psi_3(\vec{r},\vec{r}_{v_2})\phi_i^{q_i}(\vec{r}_{v_1})\delta_{q_{v_1} q_i}\right) d\vec{r} \\
&+ \sum_{i=1}^{A}\int \phi_i^{q_i *}(\vec{r})V_{v_2 i}(\vec{r}_{v_2},\vec{r})\left(\psi_3(\vec{r}_{v_1},\vec{r}_{v_2})\phi_i^{q_i}(\vec{r}) - \psi_3(\vec{r}_{v_1},\vec{r})\phi_i^{q_i}(\vec{r}_{v_2})\delta_{q_{v_2} q_i}\right) d\vec{r}
\end{aligned}$$

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concludes that the single-particle wave functions ϕ_i have to satisfy the following set of equations (see Appendix C):

$$\left[-\vec{\nabla}\cdot\frac{\hbar^2}{2m_q^*(\vec{r})}\vec{\nabla}+U_q(\vec{r})+\vec{W}_q(\vec{r})\cdot(-i)(\vec{\nabla}\times\vec{\sigma})\right]\phi_i=e_i\phi_i\,, \tag{20}$$

where q stands for the charge of the single-particle state i. Equation (20) has the form of a local Schrödinger equation with an effective mass $m^*(\vec{r})$ which depends on the density only,

$$\frac{\hbar^2}{2m_q^*(\vec{r})}=\frac{\hbar^2}{2m}+\tfrac{1}{4}(t_1+t_2)\rho+\tfrac{1}{8}(t_2-t_1)\rho_q\,; \tag{21}$$

whereas, the potential $U(\vec{r})$ also depends on the kinetic energy density,

$$\begin{aligned}U_q(\vec{r})={}&t_0[(1+\tfrac{1}{2}x_0)\rho-(x_0+\tfrac{1}{2})\rho_q]+\tfrac{1}{4}t_3(\rho^2-\rho_q{}^2)\\&-\tfrac{1}{8}(3t_1-t_2)\nabla^2\rho+\tfrac{1}{16}(3t_1+t_2)\nabla^2\rho_q+\tfrac{1}{4}(t_1+t_2)\tau\\&+\tfrac{1}{8}(t_2-t_1)\tau_q-\tfrac{1}{2}W_0(\vec{\nabla}\cdot\vec{J}+\vec{\nabla}\cdot\vec{J}_q)+\delta_{q,+\frac{1}{2}}V_C(\vec{r})\,.\end{aligned} \tag{22a}$$

The form factor $\vec{W}$ of the one-body spin-orbit potential is

$$\vec{W}_q(\vec{r})=\tfrac{1}{2}W_0(\vec{\nabla}\rho+\vec{\nabla}\rho_q)+\tfrac{1}{8}(t_1-t_2)\vec{J}_q(\vec{r})\,. \tag{22b}$$

$$\begin{aligned}\epsilon_i\phi_i^{q_i}(\vec{r})={}&\Big[-\vec{\nabla}\cdot\frac{\hbar^2}{2m^*_{q_i}(\vec{r})}\vec{\nabla}+U_{q_i}(\vec{r})-i\vec{W}_{q_i}(\vec{r})\cdot(\vec{\nabla}\times\vec{\sigma})\\&-\vec{\nabla}\cdot\frac{1}{m'^*_{q_i}(\vec{r})}\vec{\nabla}+U'_{q_i}(\vec{r})-i\vec{W'}_{q_i}(\vec{r})\cdot(\vec{\nabla}\times\vec{\sigma})\Big]\phi_i^{q_i}(\vec{r})\end{aligned}$$

$$E_3\psi_3(\vec{r}_1,\vec{r}_2)=\left[\frac{p_x^2}{2\mu_x}+\frac{p_y^2}{2\mu_y}+V_{v_1v_2}+V_{cv_1}(\vec{r}_{cv_1})+V_{cv_2}(\vec{r}_{cv_2})\right]\psi_3(\vec{r}_1,\vec{r}_2)$$

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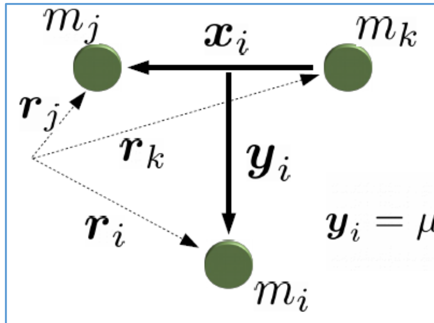
Adiabatic Expansion Method

$$\psi_3(\vec{r}_{v_1},\vec{r}_{v_2})=\frac{1}{\rho^{5/2}}\sum_n f_n(\rho)\Phi_n(\rho,\Omega)$$

$$\rho^2=x^2+y^2 \qquad \tan\alpha=\frac{x}{y}$$

$$\Omega\equiv\{\alpha,\Omega_x,\Omega_y\}$$



$$\boldsymbol{x}_i=\mu_{jk}(\boldsymbol{r}_j-\boldsymbol{r}_k)$$

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$$E_3\psi_3(\vec{r}_1,\vec{r}_2)=\left[\frac{p_x^2}{2\mu_x}+\frac{p_y^2}{2\mu_y}+V_{v_1v_2}+V_{cv_1}(\vec{r}_{cv_1})+V_{cv_2}(\vec{r}_{cv_2})\right]\psi_3(\vec{r}_1,\vec{r}_2)$$

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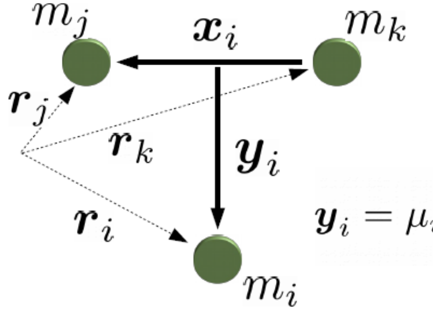

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$$\boldsymbol{y}_i=\mu_{i,jk}\left(\boldsymbol{r}_i-\frac{m_j\boldsymbol{r}_j+m_k\boldsymbol{r}_k}{m_j+m_k}\right)$$

$$\left[-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial\rho^2}+V_{eff}(\rho)-E_3\right]f_n-\frac{\hbar^2}{2m}\sum_m\left(2P_{nm}(\rho)\frac{\partial}{\partial\rho}+Q_{nm}(\rho)\right)f_m=0$$

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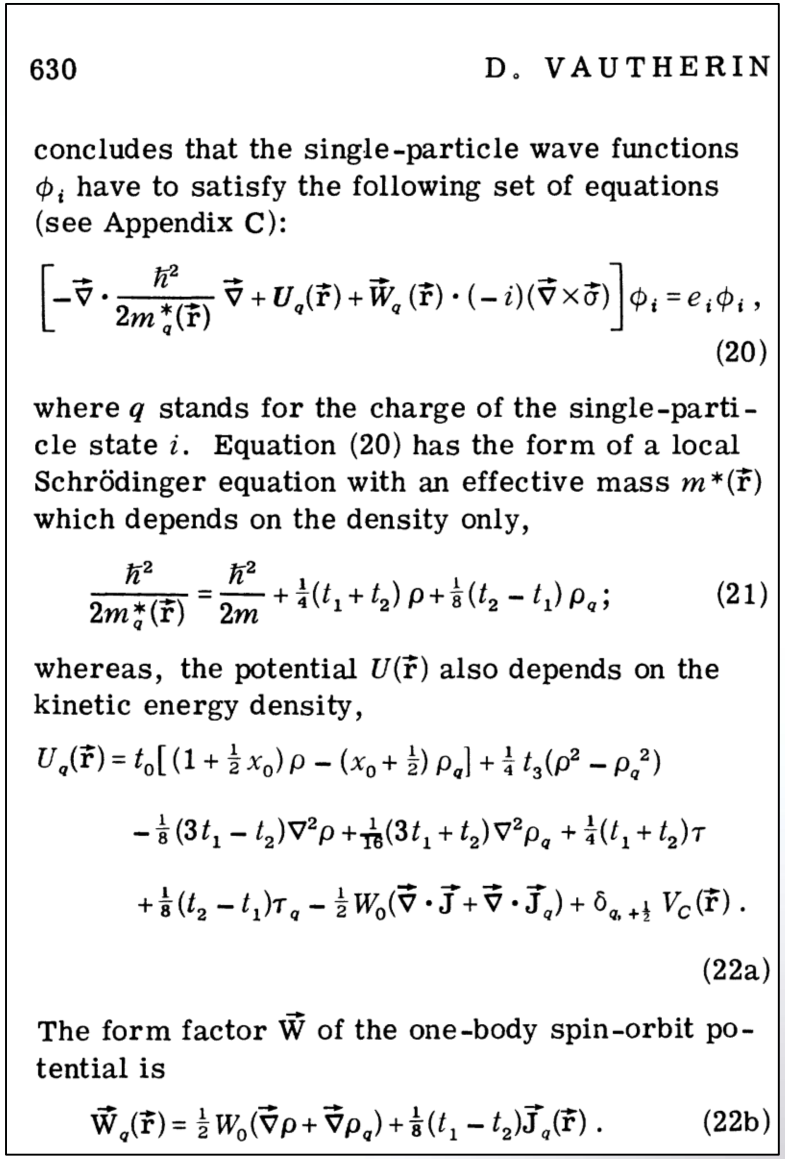

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$$\left[-\vec{\nabla}\cdot\frac{\hbar^2}{2m_q^*(\vec{r})}\vec{\nabla}+U_q(\vec{r})+\vec{W}_q(\vec{r})\cdot(-i)(\vec{\nabla}\times\vec{\sigma})\right]\phi_i=e_i\phi_i\,, \quad (20)$$

where q stands for the charge of the single-particle state i. Equation (20) has the form of a local Schrödinger equation with an effective mass $m^*(\vec{r})$ which depends on the density only,

$$\frac{\hbar^2}{2m_q^*(\vec{r})}=\frac{\hbar^2}{2m}+\tfrac{1}{4}(t_1+t_2)\rho+\tfrac{1}{8}(t_2-t_1)\rho_q\,; \quad (21)$$

whereas, the potential $U(\vec{r})$ also depends on the kinetic energy density,

$$U_q(\vec{r})=t_0[(1+\tfrac{1}{2}x_0)\rho-(x_0+\tfrac{1}{2})\rho_q]+\tfrac{1}{4}t_3(\rho^2-\rho_q^2)$$
$$-\tfrac{1}{8}(3t_1-t_2)\nabla^2\rho+\tfrac{1}{16}(3t_1+t_2)\nabla^2\rho_q+\tfrac{1}{4}(t_1+t_2)\tau$$
$$+\tfrac{1}{8}(t_2-t_1)\tau_q-\tfrac{1}{2}W_0(\vec{\nabla}\cdot\vec{J}+\vec{\nabla}\cdot\vec{J}_q)+\delta_{q,+\frac{1}{2}}V_C(\vec{r})\,. \quad (22a)$$

The form factor $\vec{W}$ of the one-body spin-orbit potential is

$$\vec{W}_q(\vec{r})=\tfrac{1}{2}W_0(\vec{\nabla}\rho+\vec{\nabla}\rho_q)+\tfrac{1}{8}(t_1-t_2)\vec{J}_q(\vec{r})\,. \quad (22b)$$

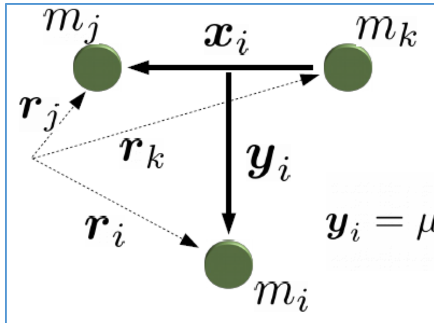
$$\epsilon_i\phi_i^{q_i}(\vec{r}) = \Bigg[-\vec{\nabla}\cdot\frac{\hbar^2}{2m^*_{q_i}(\vec{r})}\vec{\nabla}+U_{q_i}(\vec{r})-i\vec{W}_{q_i}(\vec{r})\cdot(\vec{\nabla}\times\vec{\sigma})$$
$$-\vec{\nabla}\cdot\frac{1}{m'^*_{q_i}(\vec{r})}\vec{\nabla}+U'_{q_i}(\vec{r})-i\vec{W}'_{q_i}(\vec{r})\cdot(\vec{\nabla}\times\vec{\sigma})\Bigg]\phi_i^{q_i}(\vec{r})$$

$$E_3\psi_3(\vec{r}_1,\vec{r}_2)=\left[\frac{p_x^2}{2\mu_x}+\frac{p_y^2}{2\mu_y}+V_{v_1v_2}+V_{cv_1}(\vec{r}_{cv_1})+V_{cv_2}(\vec{r}_{cv_2})\right]\psi_3(\vec{r}_1,\vec{r}_2)$$

Adiabatic Expansion Method

$$\psi_3(\vec{r}_{v_1},\vec{r}_{v_2})=\frac{1}{\rho^{5/2}}\sum_n f_n(\rho)\Phi_n(\rho,\Omega)$$

$$\rho^2=x^2+y^2 \qquad \tan\alpha=\frac{x}{y}$$
$$\Omega\equiv\{\alpha,\Omega_x,\Omega_y\}$$



$$\boldsymbol{x}_i=\mu_{jk}(\boldsymbol{r}_j-\boldsymbol{r}_k)$$
$$\boldsymbol{y}_i=\mu_{i,jk}\left(\boldsymbol{r}_i-\frac{m_j\boldsymbol{r}_j+m_k\boldsymbol{r}_k}{m_j+m_k}\right)$$

$$\left[-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial\rho^2}+V_{eff}(\rho)-E_3\right]f_n$$
$$-\frac{\hbar^2}{2m}\sum_m\left(2(P_{nm}(\rho)+P'_{nm}(\rho))\frac{\partial}{\partial\rho}+Q_{nm}(\rho)+Q'_{nm}(\rho)\right)f_m=0$$

D. Hove et al., JPG 45, 073001 (2018)

The case of ^{26}O:

PRL 109, 022501 (2012) PHYSICAL REVIEW LETTERS week ending 13 JULY 2012

$N = 16$ Spherical Shell Closure in ^{24}O

K. Tshoo,[1,*] Y. Satou,[1] H. Bhang,[1] S. Choi,[1] T. Nakamura,[2] Y. Kondo,[2] S. Deguchi,[2] Y. Kawada,[2] N. Kobayashi,[2] Y. Nakayama,[2] K. N. Tanaka,[2] N. Tanaka,[2] N. Aoi,[3] M. Ishihara,[3] T. Motobayashi,[3] H. Otsu,[3] H. Sakurai,[3] S. Takeuchi,[3] Y. Togano,[3] K. Yoneda,[3] Z. H. Li,[3] F. Delaunay,[4] J. Gibelin,[4] F. M. Marqués,[4] N. A. Orr,[4] T. Honda,[5] M. Matsushita,[5] T. Kobayashi,[6] Y. Miyashita,[7] T. Sumikama,[7] K. Yoshinaga,[7] S. Shimoura,[8] D. Sohler,[9] T. Zheng,[10] and Z. X. Cao[10]

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The unbound excited states of the neutron drip-line isotope ^{24}O have been investigated via the ^{24}O$(p, p')^{23}$O $+ n$ reaction in inverse kinematics at a beam energy of 62 MeV/nucleon. The decay energy spectrum of ^{24}O* was reconstructed from the momenta of ^{23}O and the neutron. The spin parity of the first excited state, observed at $E_x = 4.65 \pm 0.14$ MeV, was determined to be $J^\pi = 2^+$ from the angular distribution of the cross section. Higher-lying states were also observed. The quadrupole transition parameter β_2 of the 2_1^+ state was deduced, for the first time, to be 0.15 ± 0.04. The relatively high excitation energy and small β_2 value are indicative of the $N = 16$ shell closure in ^{24}O.

Physics Letters B 672 (2009) 17–21

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Evidence for a doubly magic ^{24}O

C.R. Hoffman[a,*], T. Baumann[b], D. Bazin[b], J. Brown[c], G. Christian[b,d], D.H. Denby[e], P.A. DeYoung[e], J.E. Finck[f], N. Frank[b,d,1], J. Hinnefeld[g], S. Mosby[h], W.A. Peters[b,d,2], W.F. Rogers[h], A. Schiller[b,3], A. Spyrou[b], M.J. Scott[f], S.L. Tabor[a], M. Thoennessen[b,d], P. Voss[f]

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A R T I C L E I N F O

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The decay energy spectrum for neutron unbound states in ^{24}O ($Z = 8$, $N = 16$) has been observed for the first time. The resonance energy of the lowest lying state, interpreted as the 2^+ level, has been observed at a decay energy above 600 keV. The resulting excitation energy of the 2^+ level above 4.7 MeV, supplies strong evidence that ^{24}O is a doubly magic nucleus. The data is also consistent with the presence of a second excited state around 5.33 MeV which can be interpreted as the 1^+ level.

^{24}O is, to a large extent, a spherical nucleus

The case of ^{26}O:

PRL 116, 102503 (2016) PHYSICAL REVIEW LETTERS week ending 11 MARCH 2016

Nucleus ^{26}O: A Barely Unbound System beyond the Drip Line

Y. Kondo,[1] T. Nakamura,[1] R. Tanaka,[1] R. Minakata,[1] S. Ogoshi,[1] N. A. Orr,[2] N. L. Achouri,[2] T. Aumann,[3,4] H. Baba,[5] F. Delaunay,[2] P. Doornenbal,[5] N. Fukuda,[5] J. Gibelin,[2] J. W. Hwang,[6] N. Inabe,[5] T. Isobe,[5] D. Kameda,[5] D. Kanno,[1] S. Kim,[6] N. Kobayashi,[1] T. Kobayashi,[7] T. Kubo,[5] S. Leblond,[2] J. Lee,[5] F. M. Marqués,[2] T. Motobayashi,[5] D. Murai,[8] T. Murakami,[9] K. Muto,[7] T. Nakashima,[1] N. Nakatsuka,[9] A. Navin,[10] S. Nishi,[1] H. Otsu,[5] H. Sato,[5] Y. Satou,[6] Y. Shimizu,[5] H. Suzuki,[5] K. Takahashi,[7] H. Takeda,[5] S. Takeuchi,[5] Y. Togano,[4,1] A. G. Tuff,[11] M. Vandebrouck,[12] and K. Yoneda[5]

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[2]*LPC Caen, ENSICAEN, Université de Caen, CNRS/IN2P3, F-14050 Caen, France*
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The unbound nucleus ^{26}O has been investigated using invariant-mass spectroscopy following one-proton removal reaction from a ^{27}F beam at 201 MeV/nucleon. The decay products, ^{24}O and two neutrons, were detected in coincidence using the newly commissioned SAMURAI spectrometer at the RIKEN Radioactive Isotope Beam Factory. The ^{26}O ground-state resonance was found to lie only $18 \pm 3(\text{stat}) \pm 4(\text{syst})$ keV above threshold. In addition, a higher lying level, which is most likely the first 2^+ state, was observed for the first time at $1.28^{+0.11}_{-0.08}$ MeV above threshold. Comparison with theoretical predictions suggests that three-nucleon forces, pf-shell intruder configurations, and the continuum are key elements to understanding the structure of the most neutron-rich oxygen isotopes beyond the drip line.

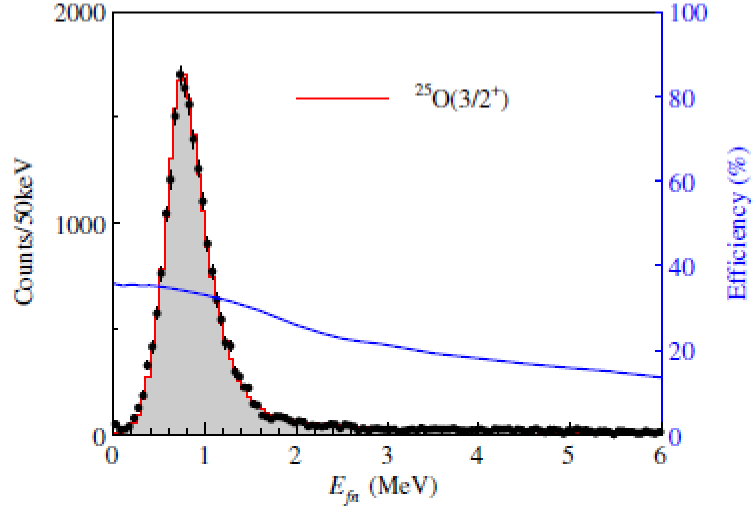
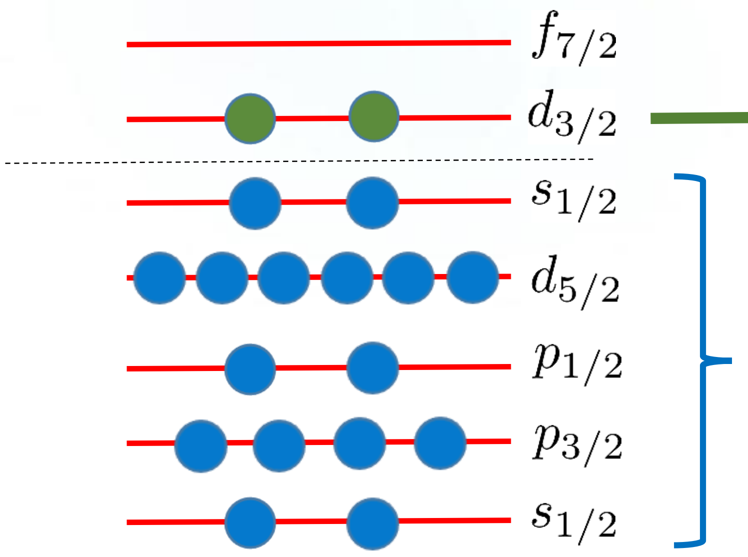


FIG. 1. Decay-energy spectrum of ^{24}O + n observed in one-proton removal from ^{26}F. The red-shaded histogram shows the fit, after accounting for the experimental response of the setup, assuming population of the ground state of ^{25}O. The blue curve represents the overall detection efficiency.

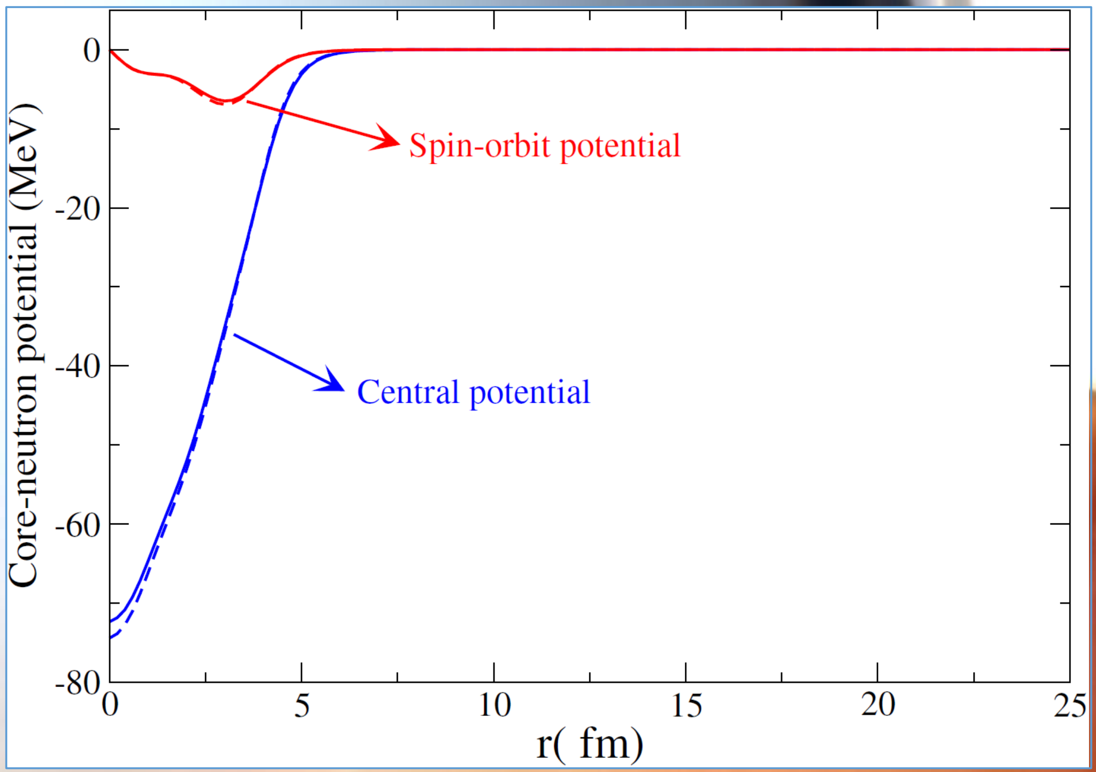
Experimental information about ^{26}O is available

The case of ^{26}O:



Valence neutrons

16 neutron states occupied by the neutrons in the ^{24}O core.



Adiabatic Expansion Method

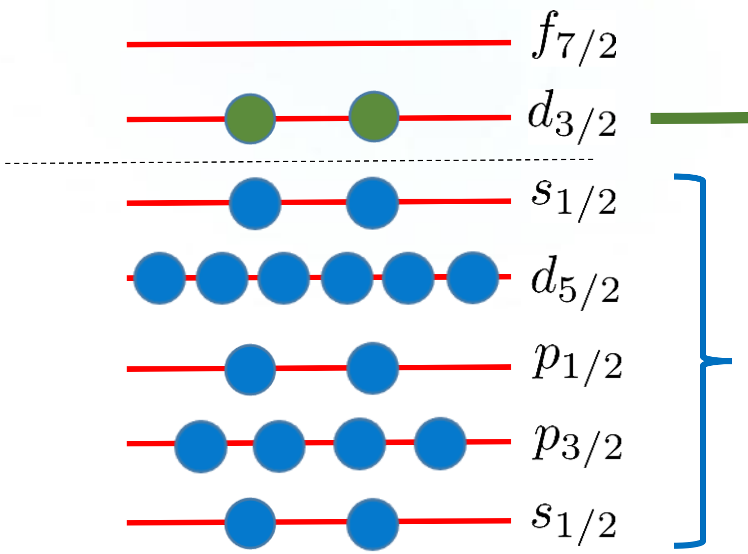
$$\psi_3(\vec{r}_{v_1}, \vec{r}_{v_2}) = \frac{1}{\rho^{5/2}} \sum_n f_n(\rho) \Phi_n(\rho, \Omega)$$

$$\left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \rho^2} + V_{eff}(\rho) - E_3 \right] f_n - \frac{\hbar^2}{2m} \sum_m \left(2P_{nm}(\rho) \frac{\partial}{\partial \rho} + Q_{nm}(\rho) \right) f_m = 0$$

	SLy4	SkM*	Sk3
$E_{d_{3/2}}$	0.85	0.83	1.23
$E_{d_{3/2}}$(HF)	-0.96	-1.15	-0.53

s, p, d, and f waves are included in the calculation

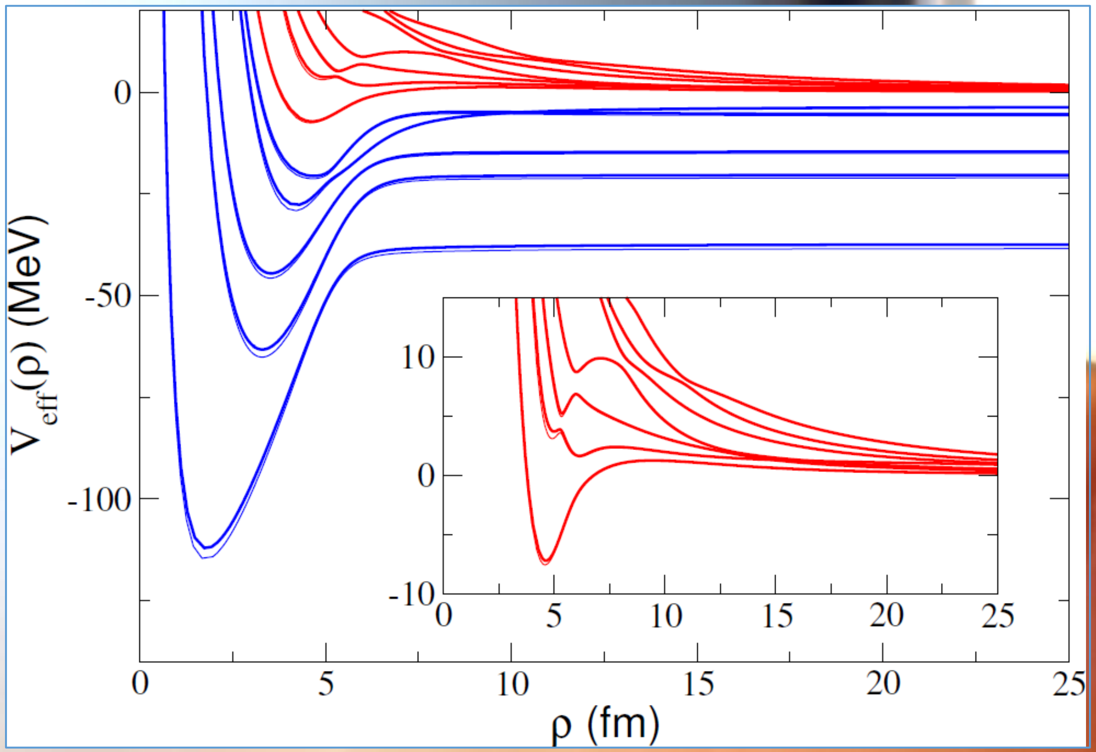
The case of ^{26}O:



Valence neutrons

16 neutron states occupied by the neutrons in the ^{24}O core.

Pauli forbidden to the valence neutrons !!!



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The adiabatic channels associated to Pauli forbidden states are removed from the calculation

s, p, d, and f waves are included in the calculation

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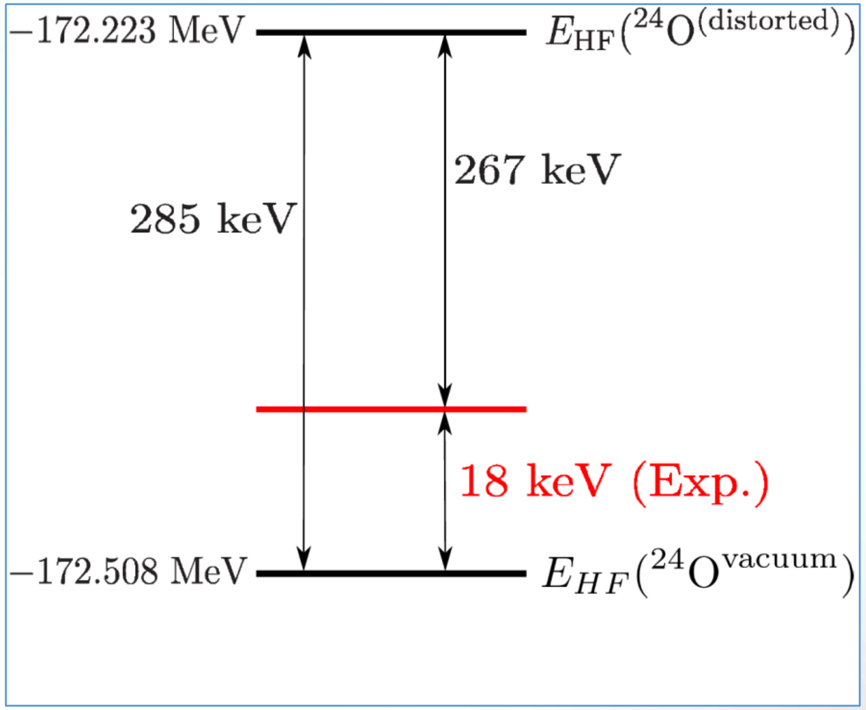

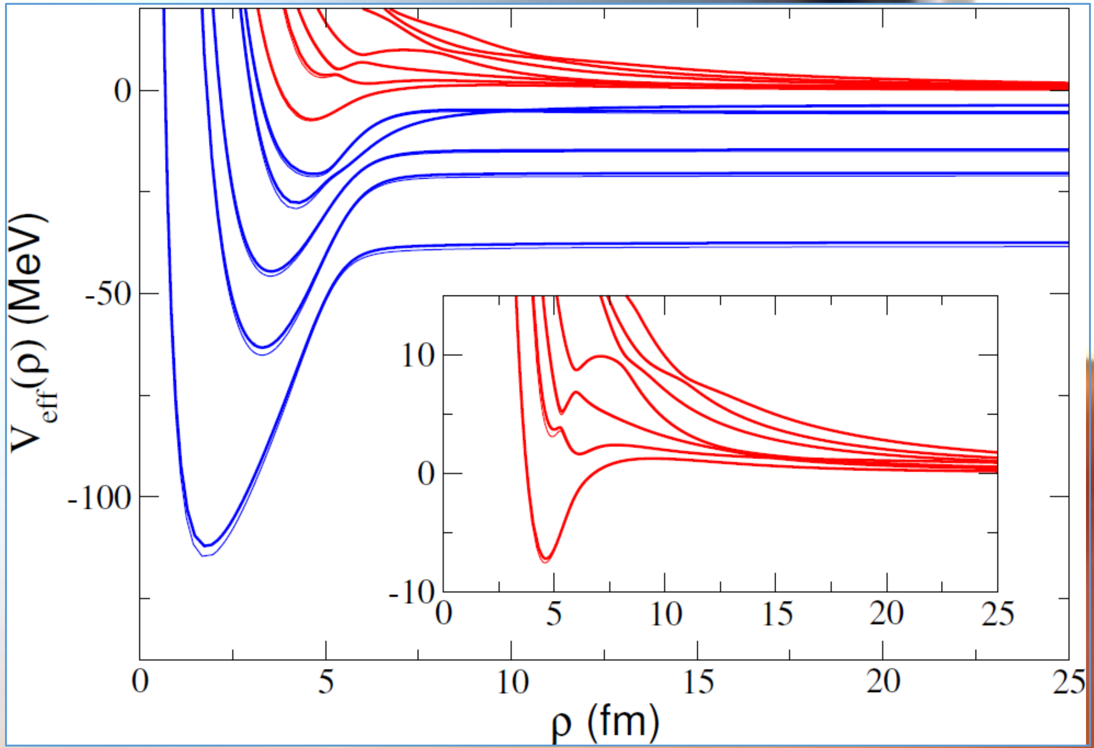

$$\left[-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial\rho^2} + V_{eff}(\rho) + V_{3b}(\rho) - E_3\right] f_n$$
$$-\frac{\hbar^2}{2m}\sum_m \left(2P_{nm}(\rho)\frac{\partial}{\partial\rho} + Q_{nm}(\rho)+\right) f_m = 0$$

The adiabatic channels associated to Pauli forbidden states are removed from the calculation

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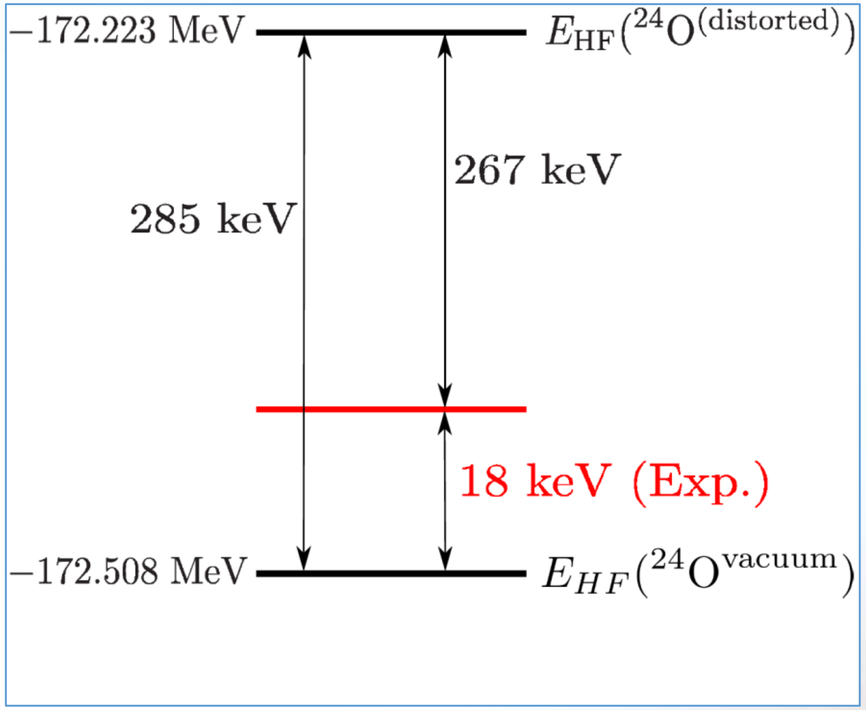

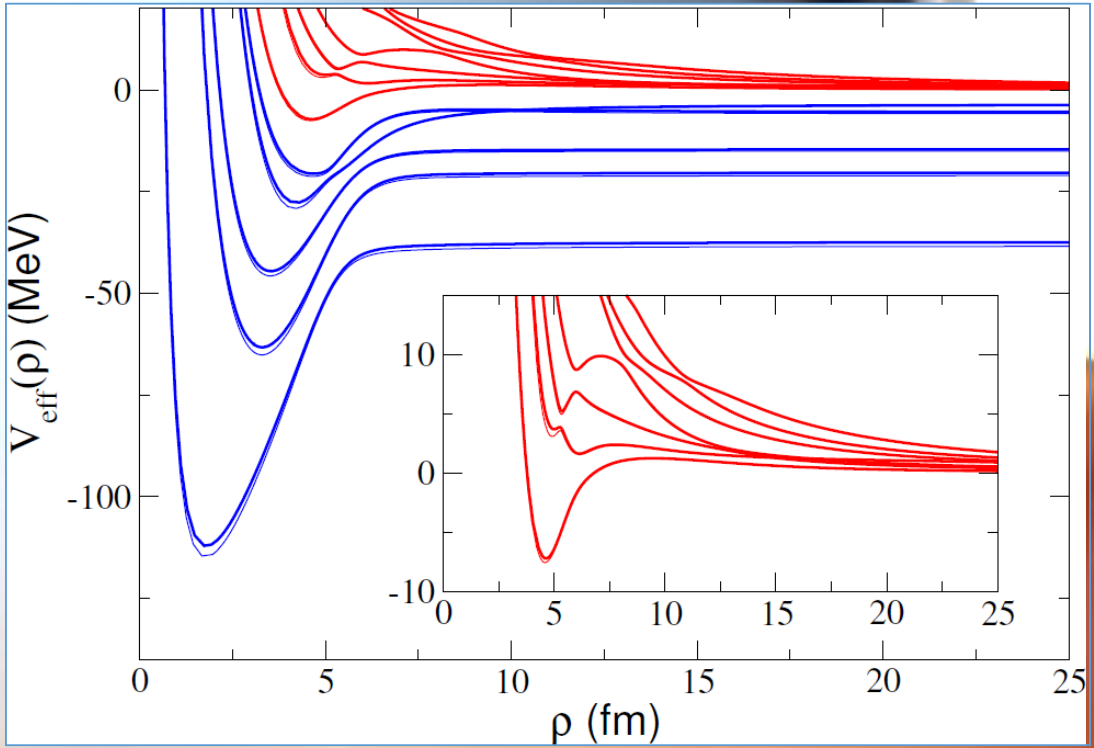

	$(d_{3/2},d_{3/2})$	$(f_{7/2},f_{7/2})$	$(p_{3/2},p_{3/2})$
% of the norm	90.1	3.7	2.1

$$1 = \sum_{\ell_x} \int_0^{\infty} \left(|f_{\ell_x=0}(\rho)|^2 + |f_{\ell_x=1}(\rho)|^2 + |f_{\ell_x=2}(\rho)|^2 + \cdots \right)$$

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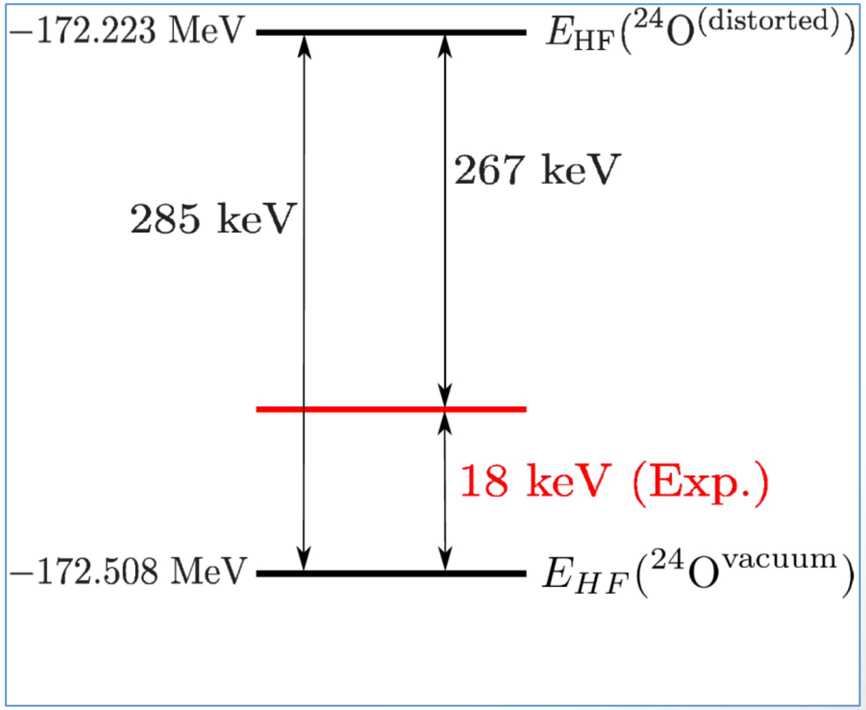

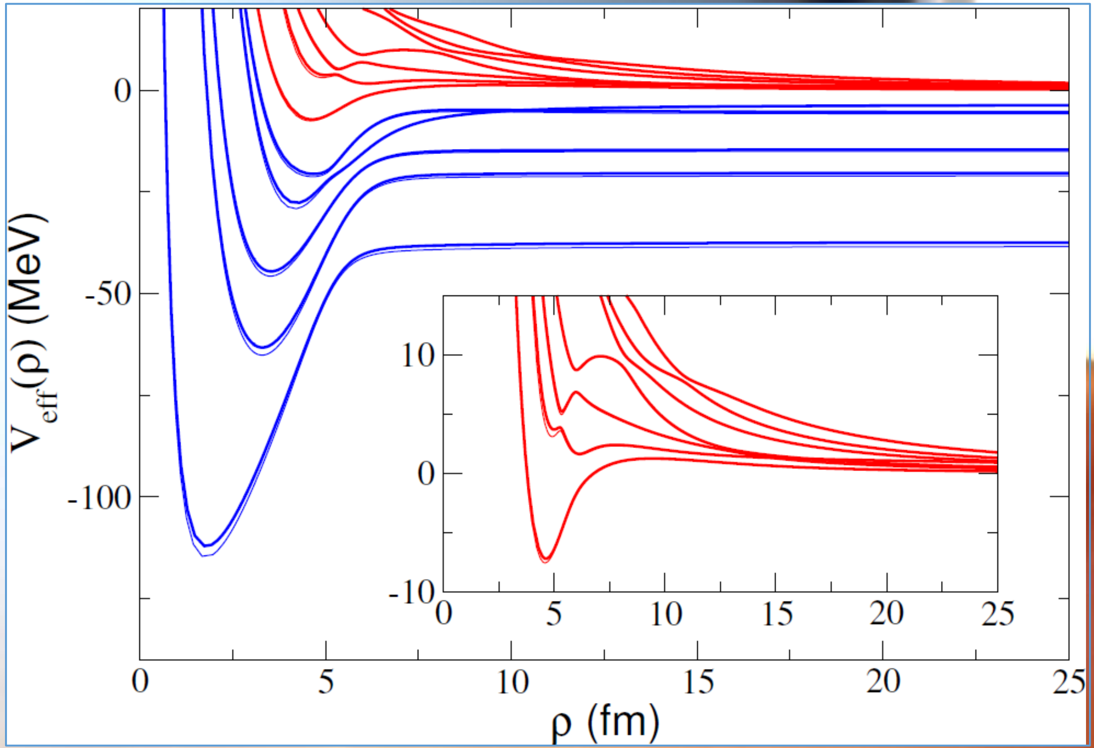

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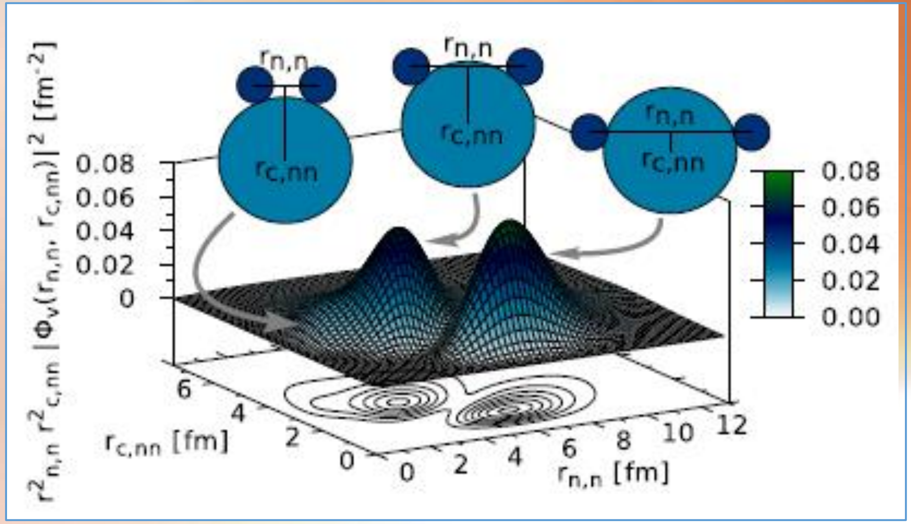

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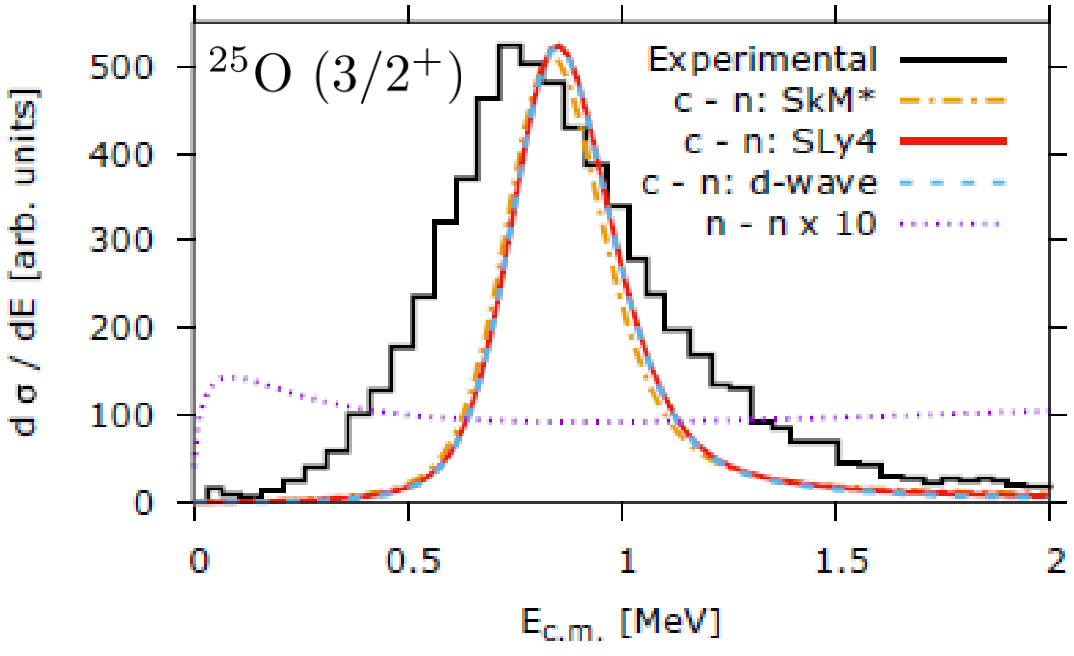
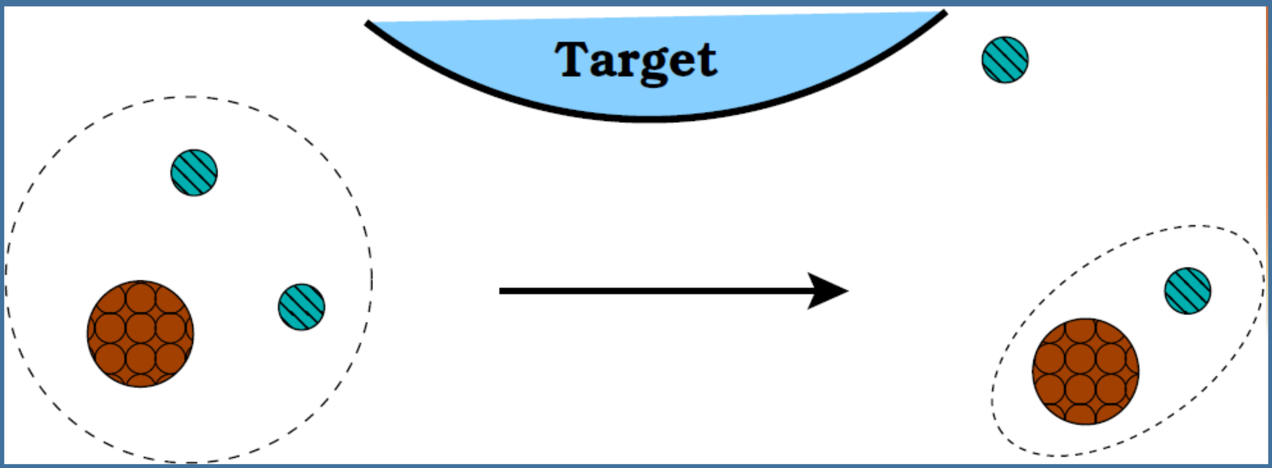


FIG. 4. The invariant mass spectra of core neutron for the SkM* (dash-dotted, orange) and SLy4 (solid, red) Skyrme parameters. The SLy4 core-neutron d-wave contribution (dashed, blue) and neutron-neutron (dotted, purple) invariant mass spectrum is also included. The black step curve is the measurements from Ref. [26].

Sudden approximation

$$\frac{d^6\sigma}{d\boldsymbol{k}_x\boldsymbol{k}_y} \propto \sum_M \sum_{s_x\sigma_x\sigma_y} \left| \langle e^{i\boldsymbol{k}_x\cdot\boldsymbol{k}_y} \chi_{s_y}^{\sigma_y} w_{s_x}^{\sigma_x}(\boldsymbol{k}_x, x) | \Psi^{JM}(\boldsymbol{x}, \boldsymbol{y})\rangle \right|^2$$

$$\frac{d\sigma}{dE_{nc}} = \frac{E_c E_n}{E_c + E_n} \frac{m(M_c + M_n)}{M_c M_n} \frac{1}{k_x} \frac{d\sigma}{dk_x}$$



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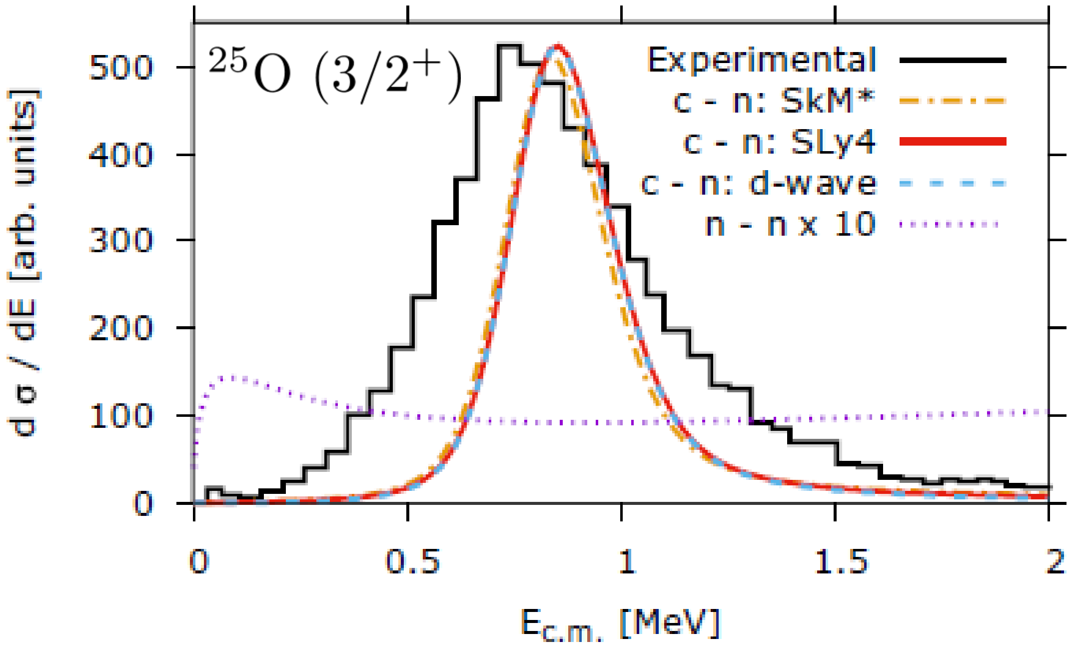

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$$\frac{d\sigma}{dE_{nc}} = \frac{E_c E_n}{E_c + E_n}\frac{m(M_c + M_n)}{M_c M_n}\frac{1}{k_x}\frac{d\sigma}{dk_x}$$

Once the NN interaction has been chosen, the invariant mass spectrum is fully determined.

Two-proton capture: ^{70}Kr

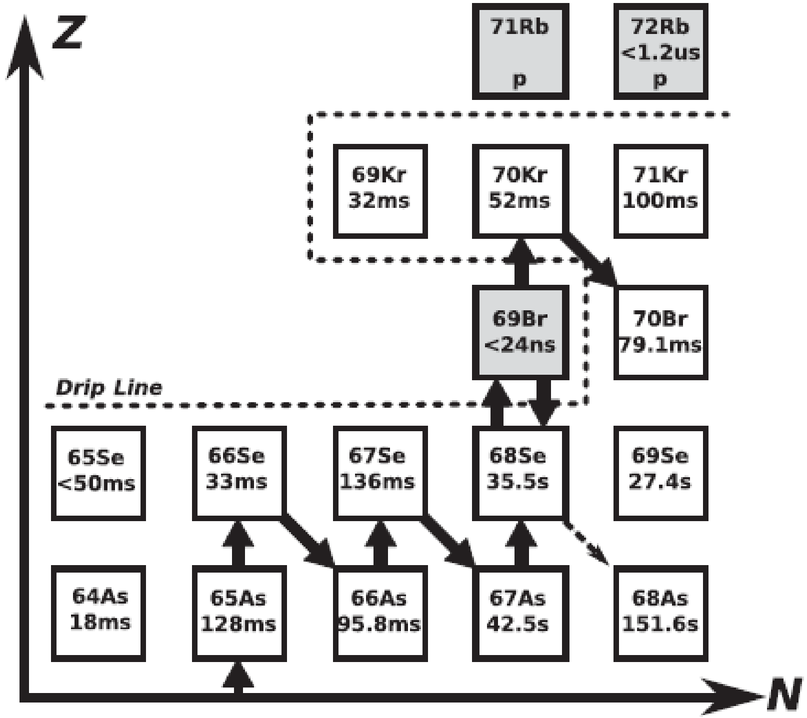

FIG. 1. Illustration of $2p$-capture reactions through ^{69}Br bypassing the ^{68}Se waiting point. The slow β decay of ^{68}Se restricts the rp-process reaction flow in type I x-ray bursts.

A.M. Rogers et al., PRL 106, 252503 (2011)

$$^{68}\mathrm{Se} + p + p \rightarrow {}^{70}\mathrm{Kr} + \gamma$$

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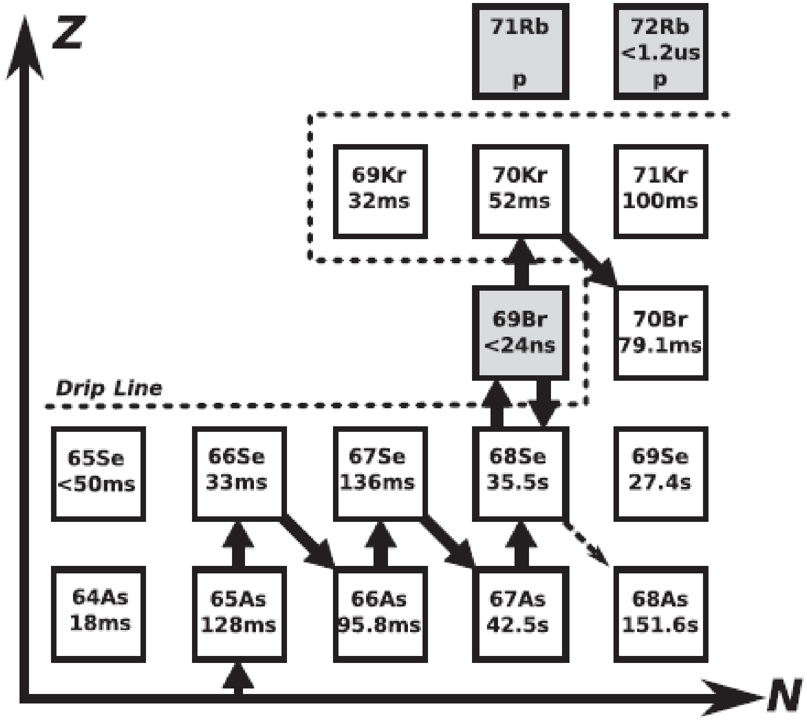

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$$R_{ppc}(E) = \frac{8\pi}{(\mu_{cp}\mu_{cp,p})^{3/2}} \frac{\hbar^3}{c^2} \left(\frac{E_\gamma}{E}\right)^2 \sigma_\gamma^\lambda(E_\gamma)$$

$$\sigma_\gamma^\lambda(E_\gamma) = \frac{(2\pi)^3(\lambda+1)}{\lambda((2\lambda+1)!!)^2} \left(\frac{E_\gamma}{\hbar c}\right)^{2\lambda-1} \frac{d}{dE}\mathcal{B}(E\lambda, 0^+ \rightarrow \lambda^\pi)$$

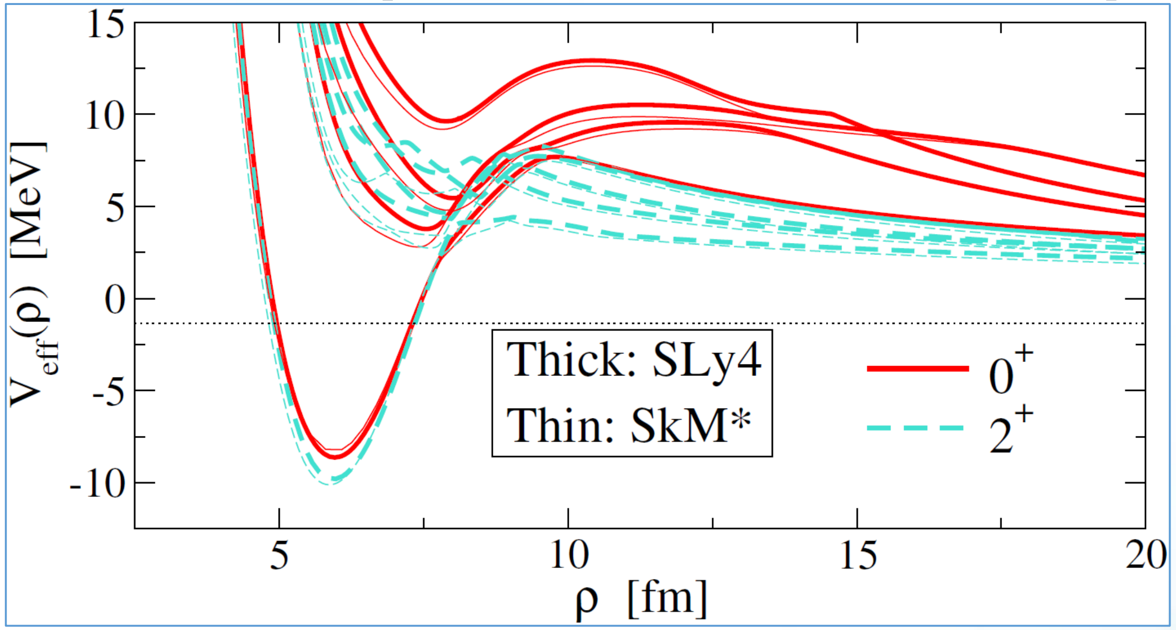
$$E_\gamma = E + |E_{g.s}|$$

$$\frac{d}{dE}\mathcal{B}(E\lambda, 0^+ \rightarrow \lambda^\pi) = \sum_i \left|\langle \psi_{\lambda^\pi}^{(i)} \parallel \hat{\mathcal{O}}_\lambda \parallel \Psi_{0^+}\rangle\right|^2 \delta(E - E_i)$$

$$\langle R_{ppc}(E)\rangle = \frac{1}{2(KT)^3} \int E^2 R_{ppc}(E) e^{-E/KT} dE$$

✓ Dominated by the E2 transition $2^+ \rightarrow 0^+$

Two-proton capture: ^{70}Kr



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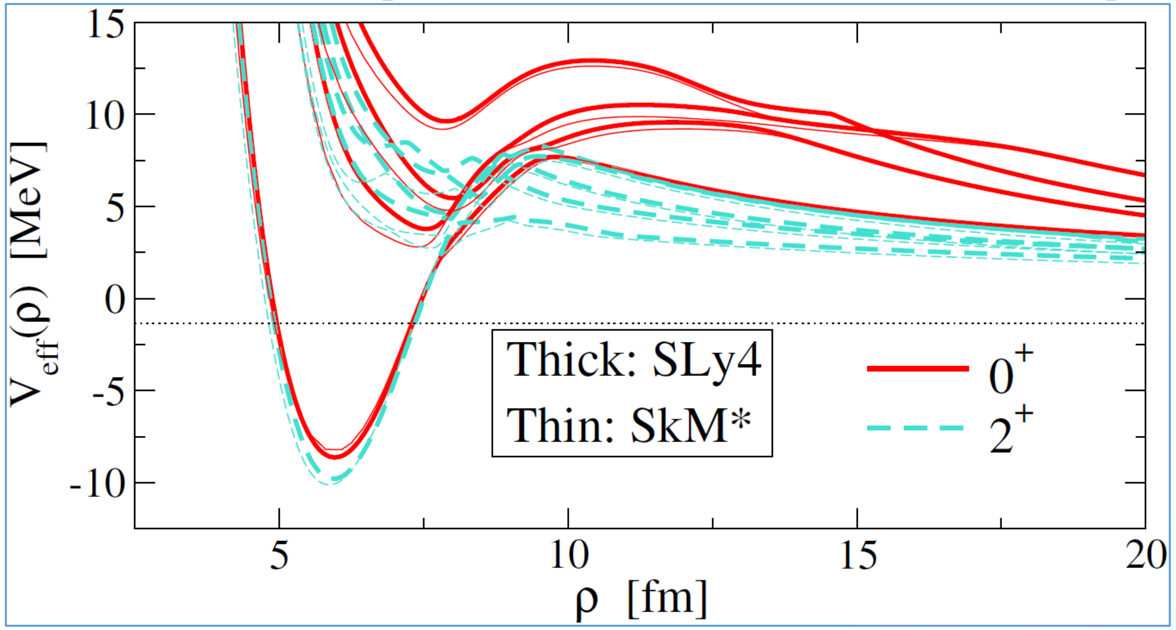
Two-proton capture: ^{70}Kr

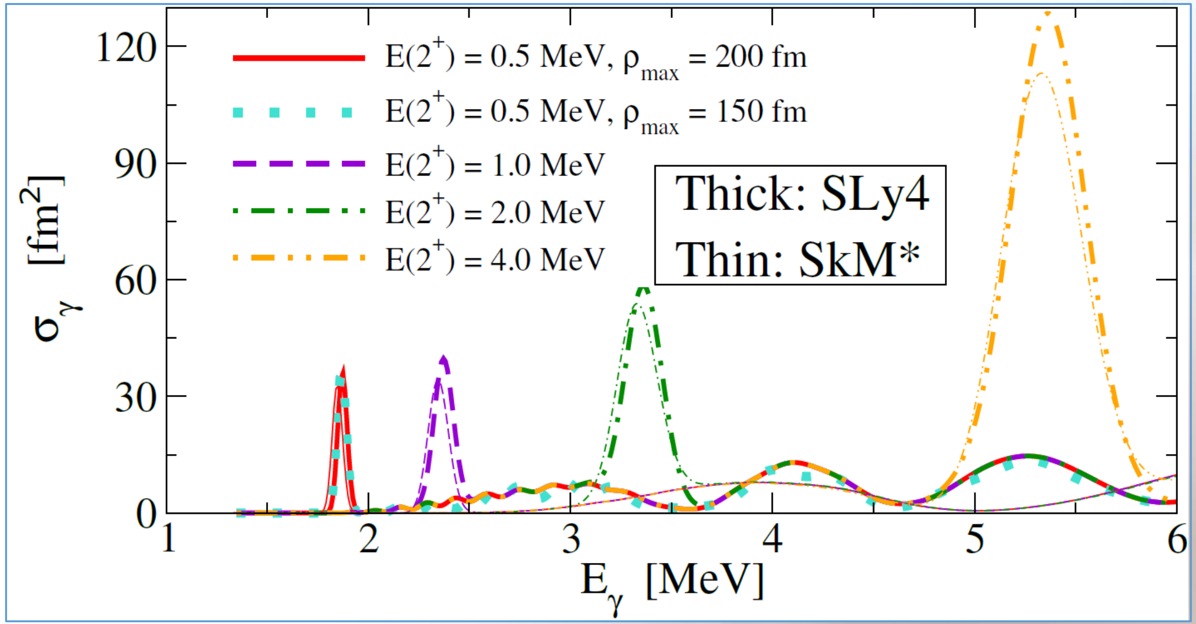


$$^{68}\mathrm{Se} + p + p \to {}^{70}\mathrm{Kr} + \gamma$$

$$R_{ppc}(E) = \frac{8\pi}{(\mu_{cp}\mu_{cp,p})^{3/2}} \frac{\hbar^3}{c^2} \left(\frac{E_\gamma}{E}\right)^2 \sigma_\gamma^\lambda(E_\gamma)$$

$$\sigma_\gamma^\lambda(E_\gamma) = \frac{(2\pi)^3(\lambda+1)}{\lambda((2\lambda+1)!!)^2} \left(\frac{E_\gamma}{\hbar c}\right)^{2\lambda-1} \frac{d}{dE}\mathcal{B}(E\lambda, 0^+ \to \lambda^\pi)$$

$$E_\gamma = E + |E_{g.s}|$$

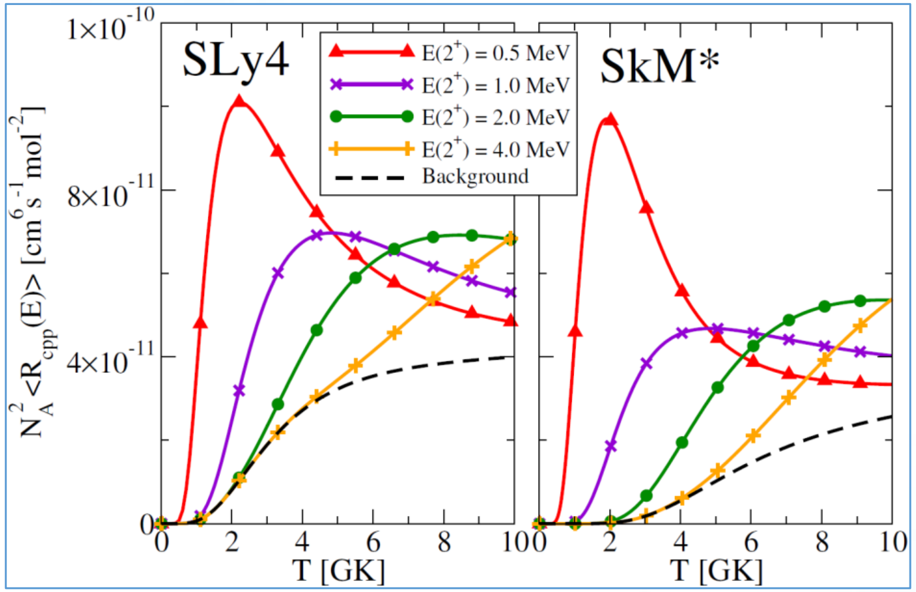


$$\frac{d}{dE}\mathcal{B}(E\lambda, 0^+ \to \lambda^\pi) = \sum_i \left| \langle \psi_{\lambda^\pi}^{(i)} \parallel \hat{\mathcal{O}}_\lambda \parallel \Psi_{0^+} \rangle \right|^2 \delta(E - E_i)$$

$$\langle R_{ppc}(E) \rangle = \frac{1}{2(KT)^3} \int E^2 R_{ppc}(E) e^{-E/KT} dE$$

- ✓ Dominated by the E2 transition $2^+ \to 0^+$

Two-proton capture: ^{70}Kr



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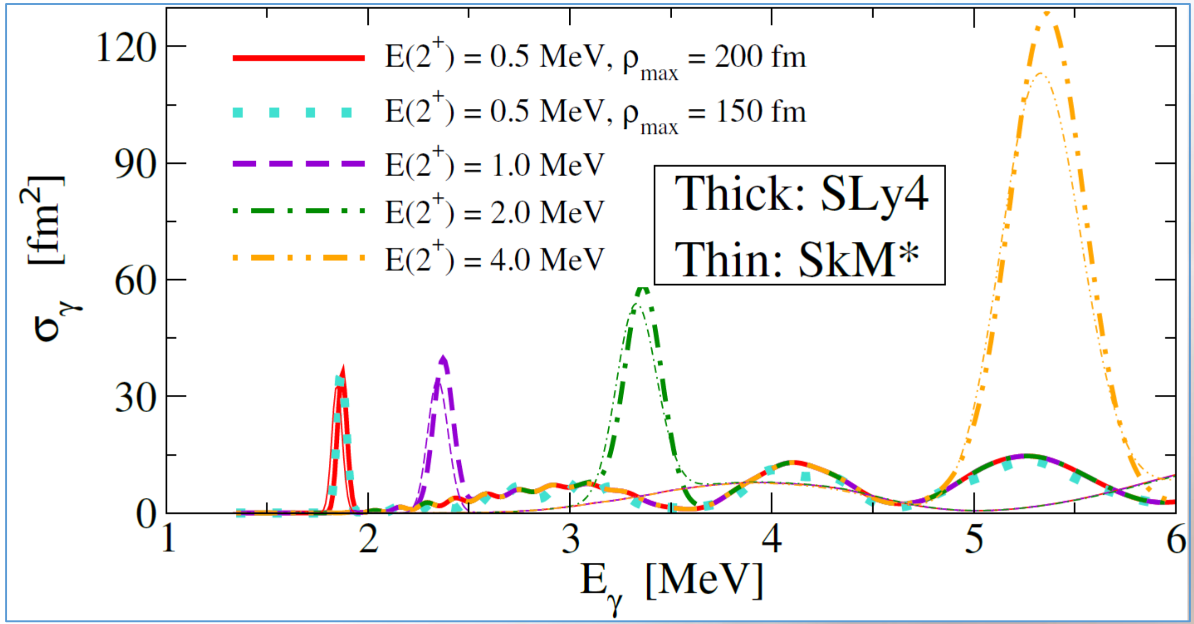

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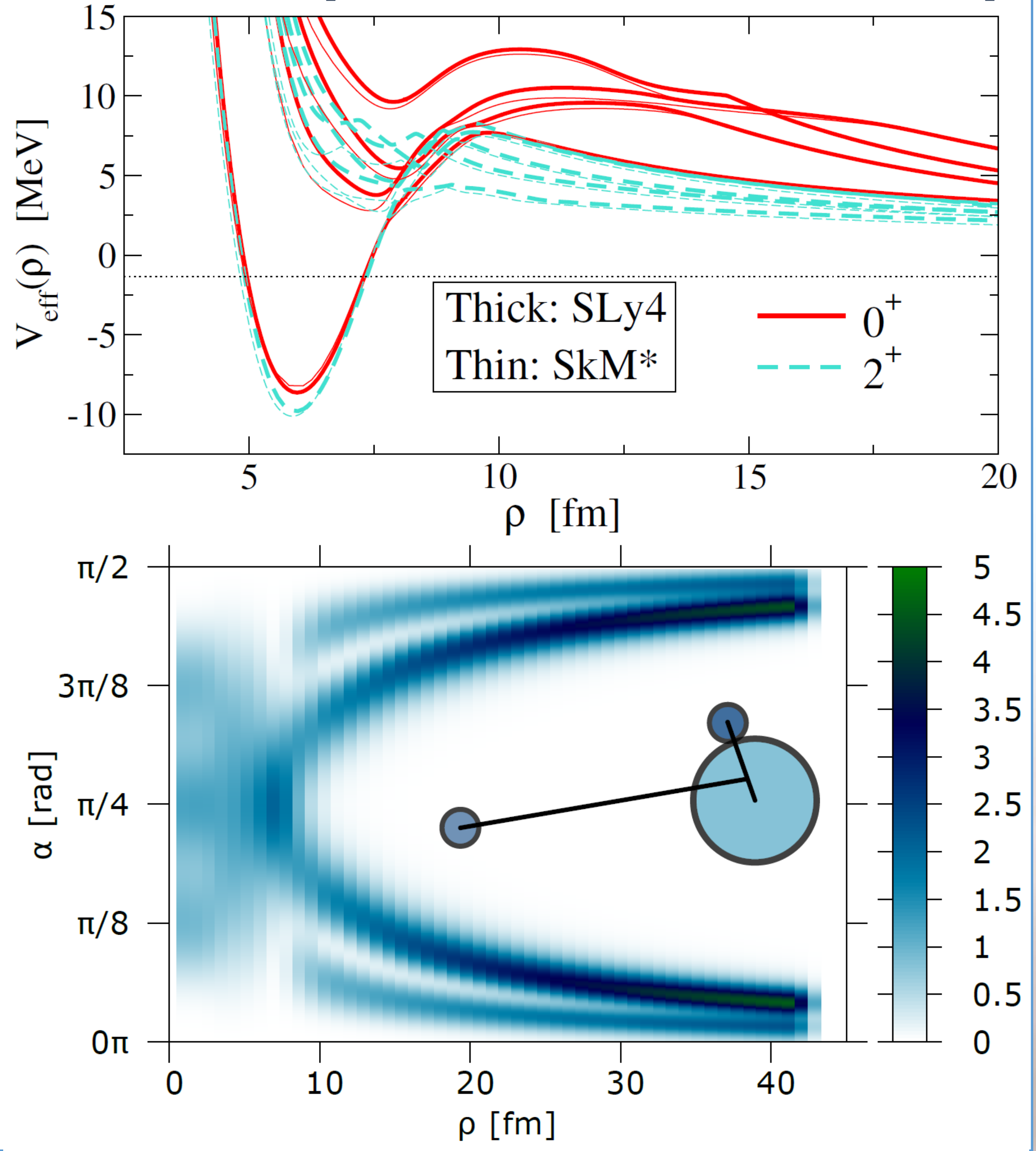
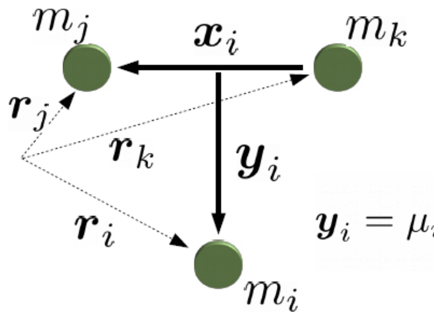
Two-proton capture: ^{70}Kr



Capture mechanism

$$P(\alpha, \rho) = \sin^2\alpha \cos^2\alpha \int |\Phi(\rho, \alpha, \Omega_x, \Omega_y)|^2 d\Omega_x d\Omega_y$$



$$x_i = \mu_{jk}(r_j - r_k)$$

$$y_i = \mu_{i,jk}\left(r_i - \frac{m_j r_j + m_k r_k}{m_j + m_k}\right)$$

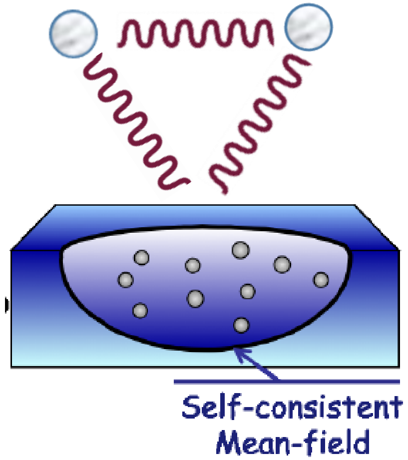
$$x_i = \rho \sin\alpha_i$$
$$y_i = \rho \cos\alpha_i$$

✓ Sequential capture through the $f_{5/2}$ resonance in ^{69}Br at 0.6 MeV.

D. Hove et al., PLB 782, 42 (2018)

Summary:

The model presented here treats the many-body core and the two valence particles self-consistently:



- ✓ The mean-field determines the interaction with the loosely bound nucleons.

Three-body problem

- ✓ The presence of the loosely bound nucleons distorts the mean-field.
- ✓ The distorted mean-field modifies the core-nucleon interaction and the wave function of the valence nucleons.

Self-consistent calculation

Core: Spherical Skyrme Hartree-Fock

- ✓ Core deformation?
- ✓ Finite range NN interactions?
- ✓ odd-odd nuclei?

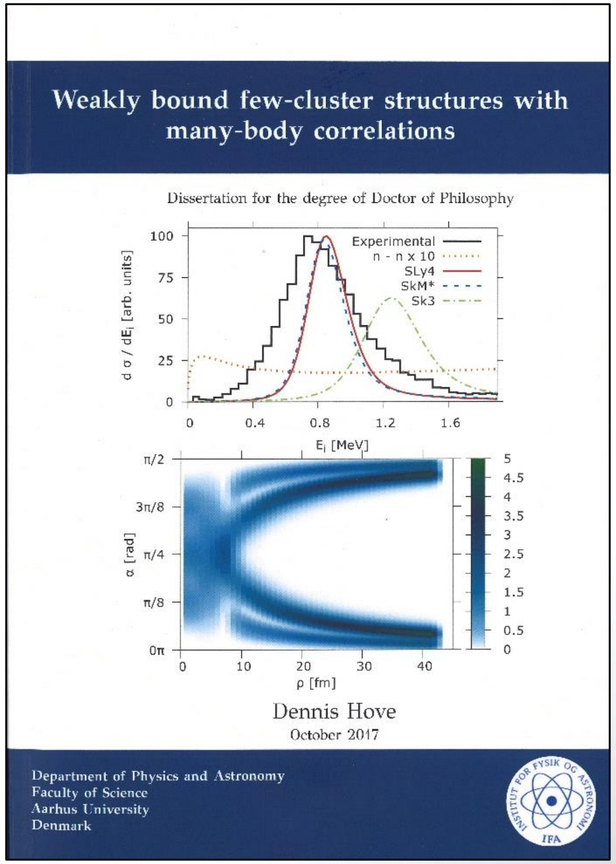
Three-body calculation: Adiabatic expansion

- ✓ Three-body force from the mean field?

- ✓ Generalization to more than one cluster and more than two valence nucleons

Weakly bound nuclei:

A unified description of intrinsic and relative degrees of freedom



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