

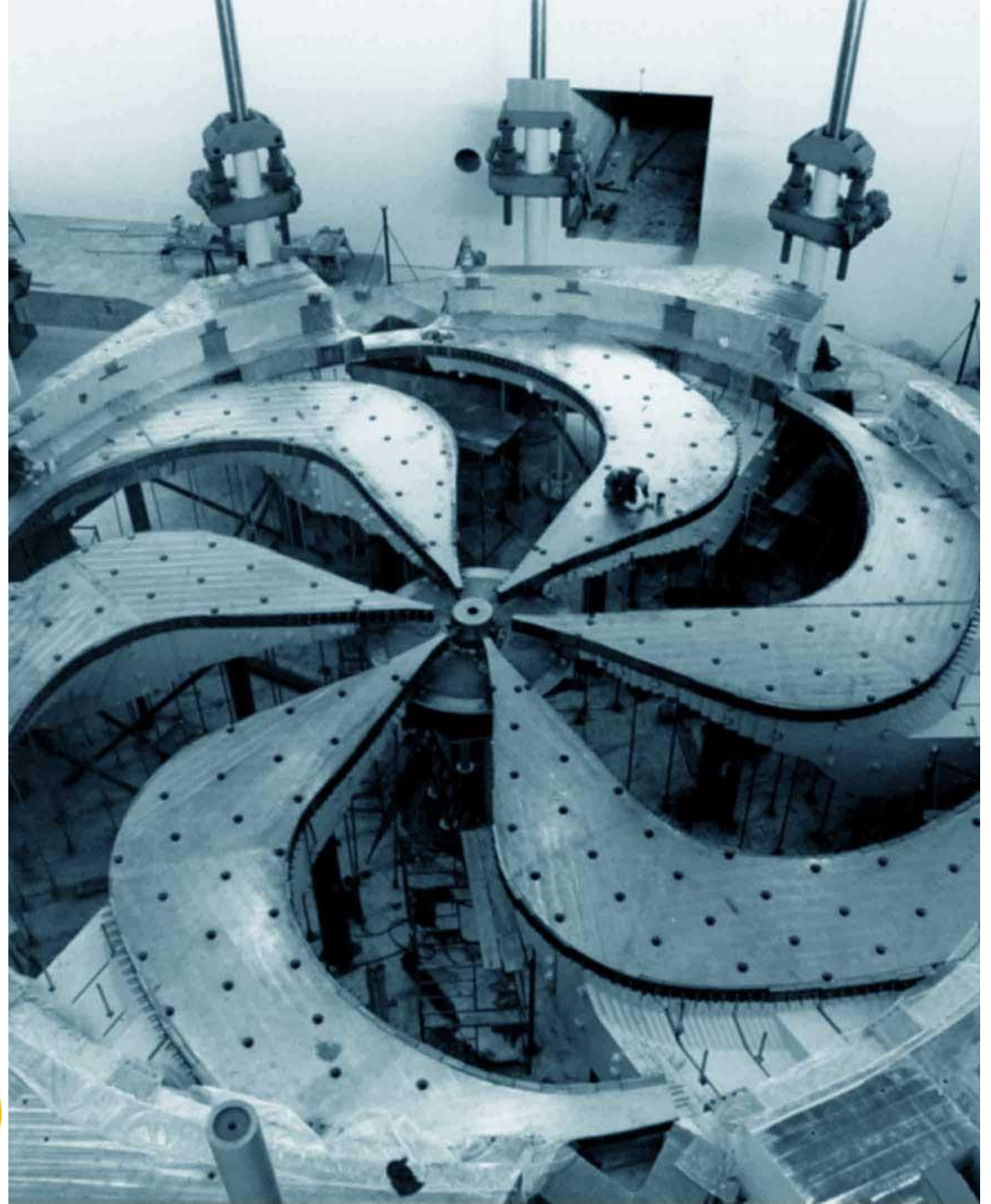


Ab Initio Correlation between Double Gamow-Teller Transitions and Neutrinoless Double Beta Decay

Izzy Ginnett

Summer Research Student, Theory Department

23 August 2021

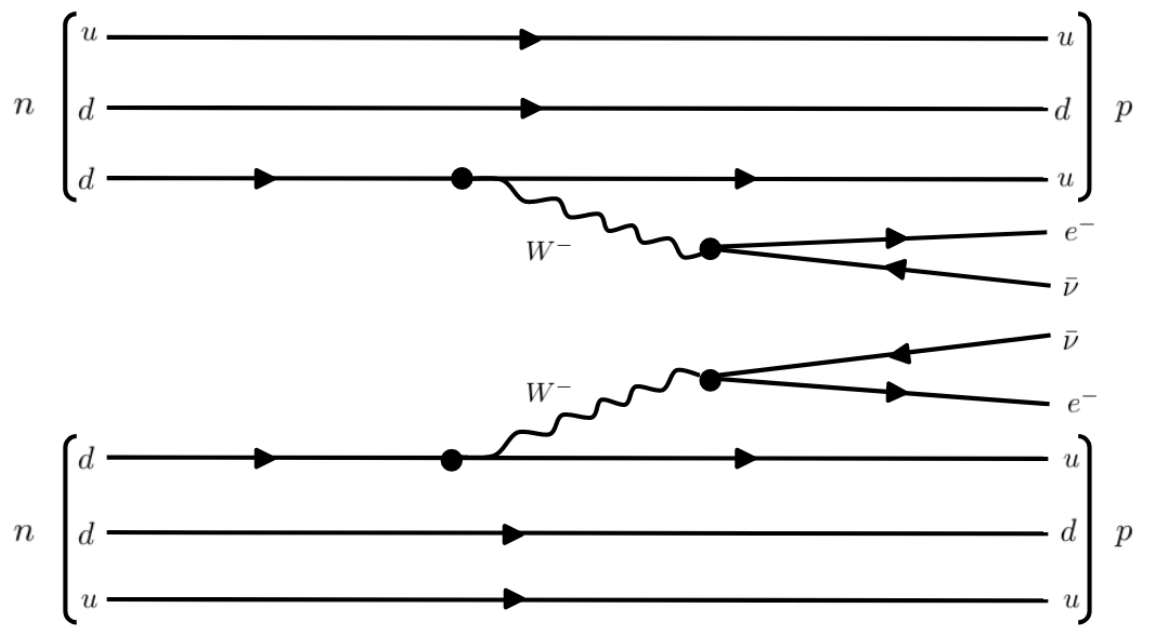


Arthur B. McDonald
Canadian Astroparticle Physics Research Institute

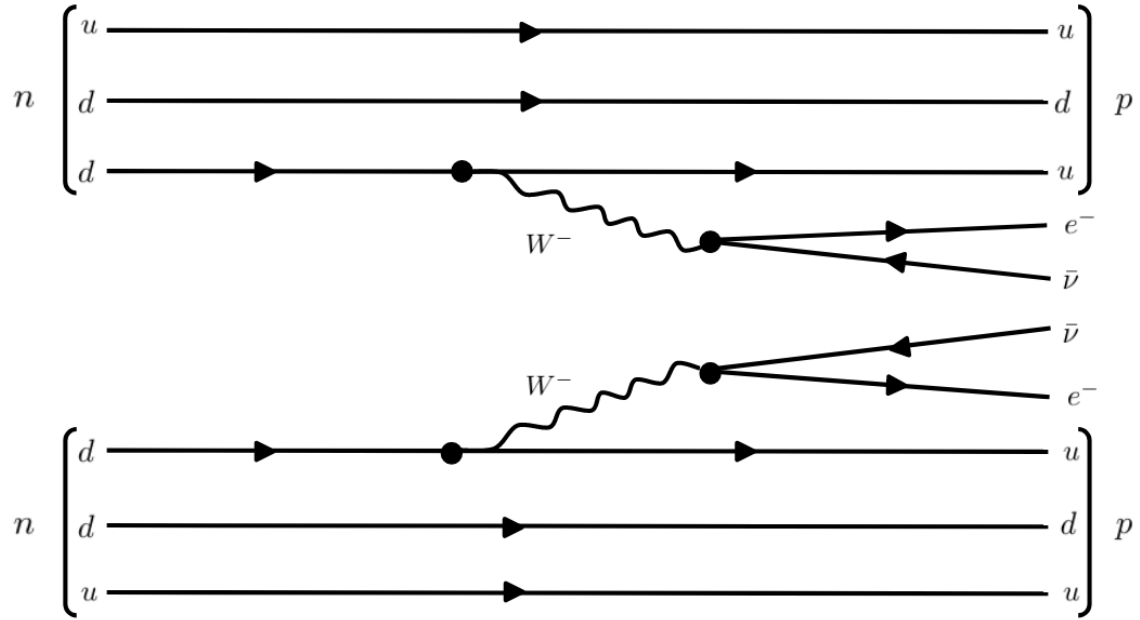


Discovery,
accelerated

$2\nu\beta\beta$

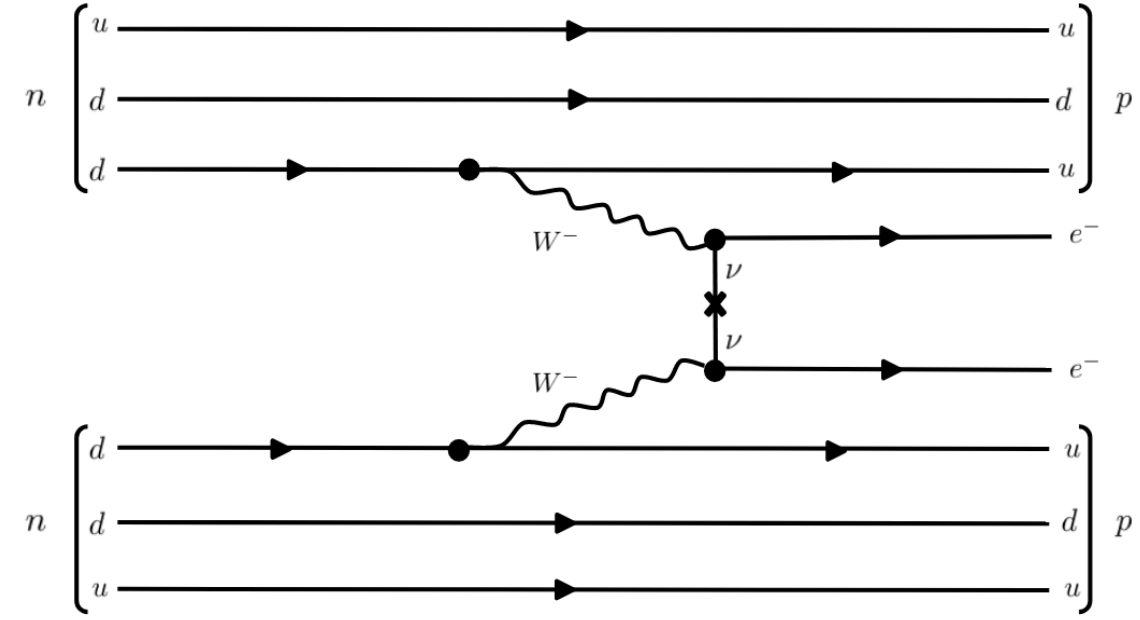


$2\nu\beta\beta$

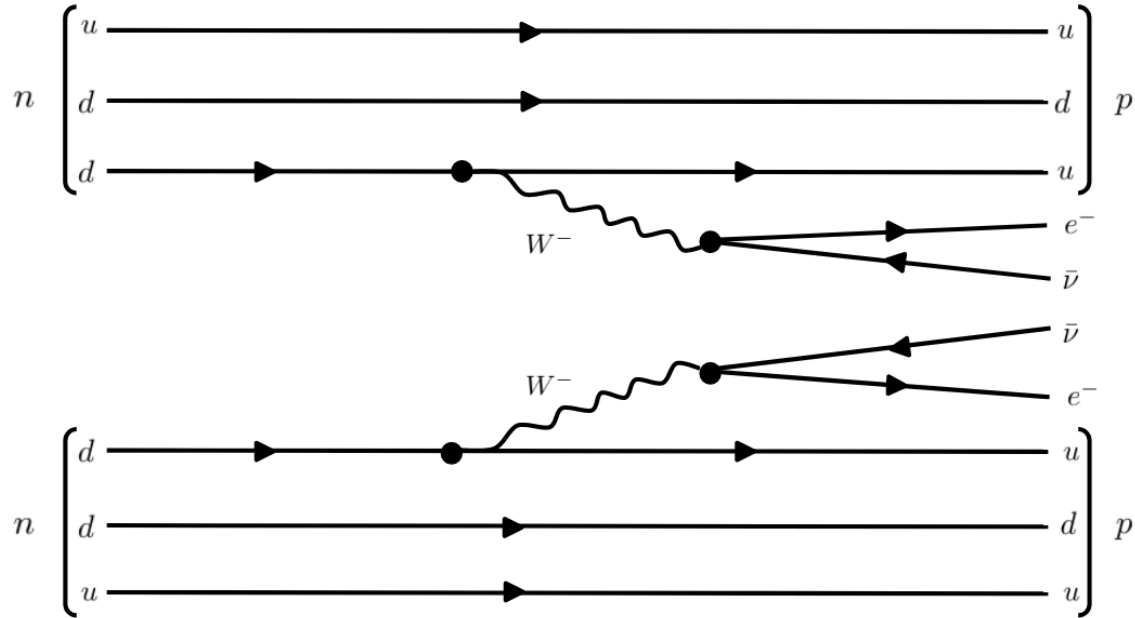


\Rightarrow

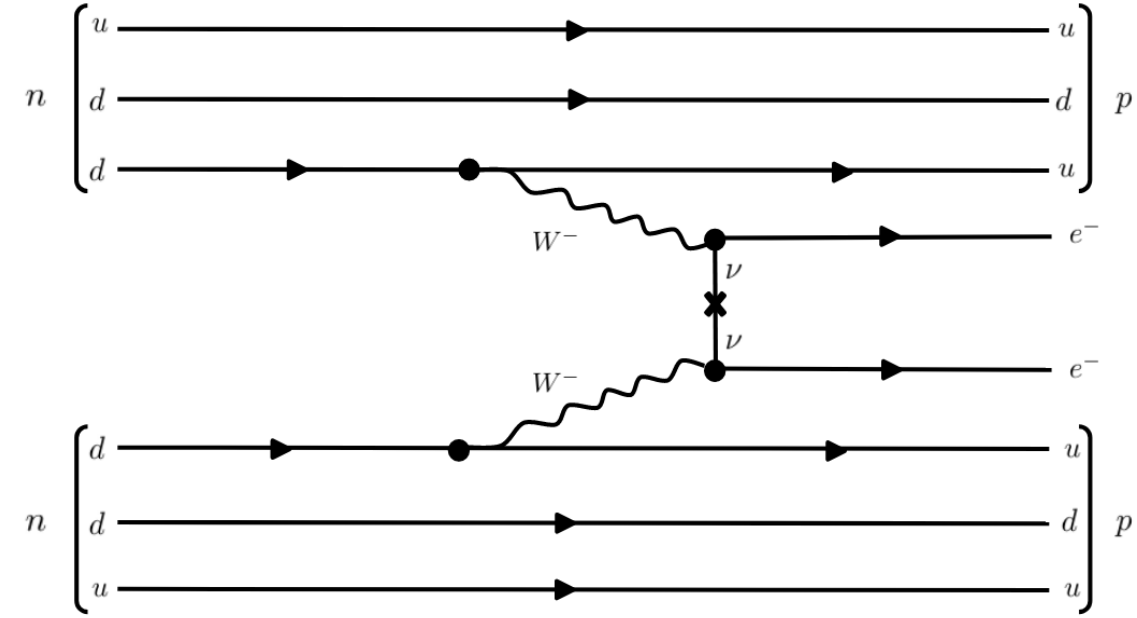
$0\nu\beta\beta$



$2\nu\beta\beta$



$0\nu\beta\beta$



Allowed by the current Standard Model

2 neutrinos in the final state

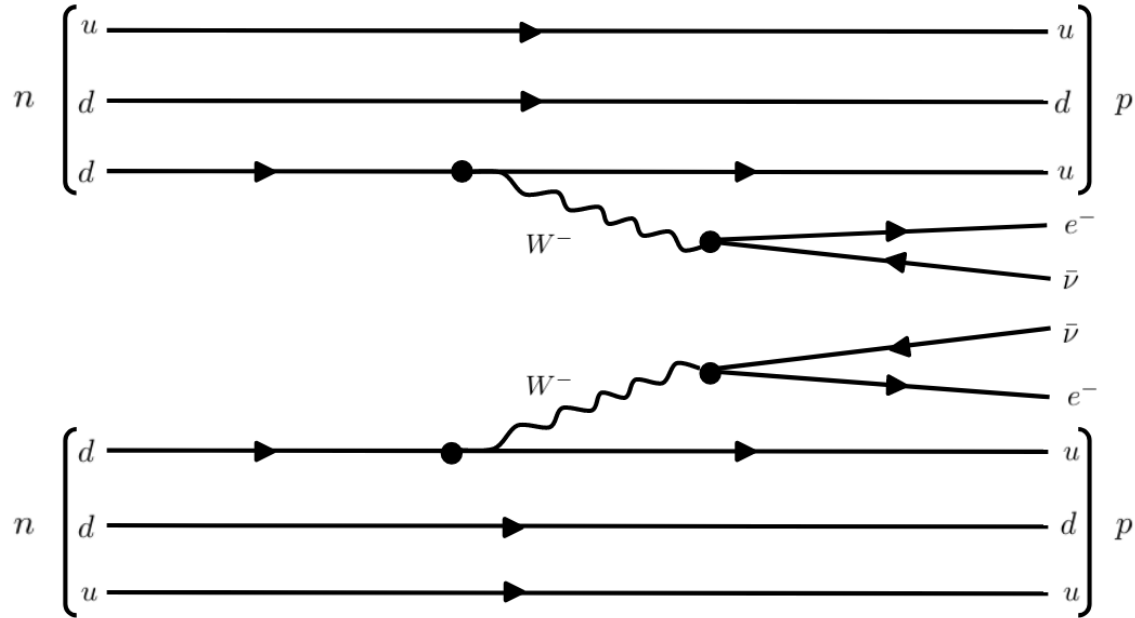
Neutrino is a Dirac fermion, i.e. particle and antiparticle are distinct

Beyond-Standard-Model (BSM) process

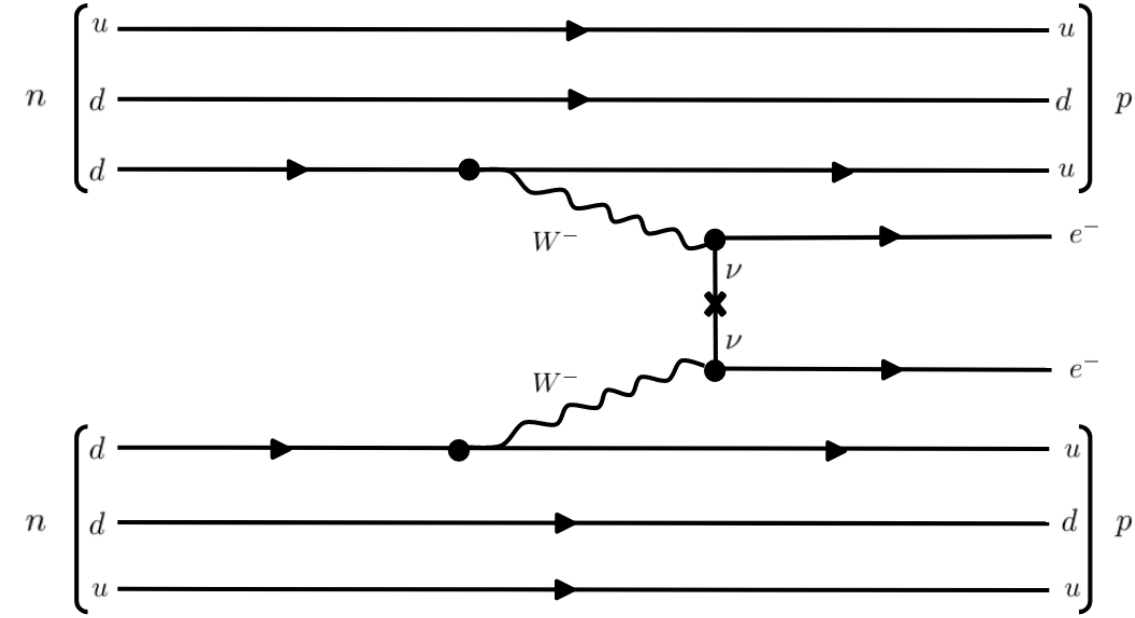
0 neutrinos in the final state

Neutrino is a Majorana fermion, i.e. its own antiparticle

$2\nu\beta\beta$



$0\nu\beta\beta$



Allowed by the current Standard Model

2 neutrinos in the final state

Neutrino is a Dirac fermion, i.e. particle and antiparticle are distinct

Beyond-Standard-Model (BSM) process

0 neutrinos in the final state

Neutrino is a Majorana fermion, i.e. its own antiparticle

Big impact on our understanding of BSM physics!

$$\left[T_{1/2}^{0\nu}\right]^{-1} = G^{0\nu} |M^{0\nu}|^2 \left(\frac{\langle m_{\beta\beta} \rangle}{m_e}\right)^2$$

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$0\nu\beta\beta$ half-life $0\nu\beta\beta$ phase factor Effective neutrino mass

$$\langle m_{\beta\beta} \rangle = \left| \sum_{k=1}^3 m_k U_{ek}^2 \right|$$

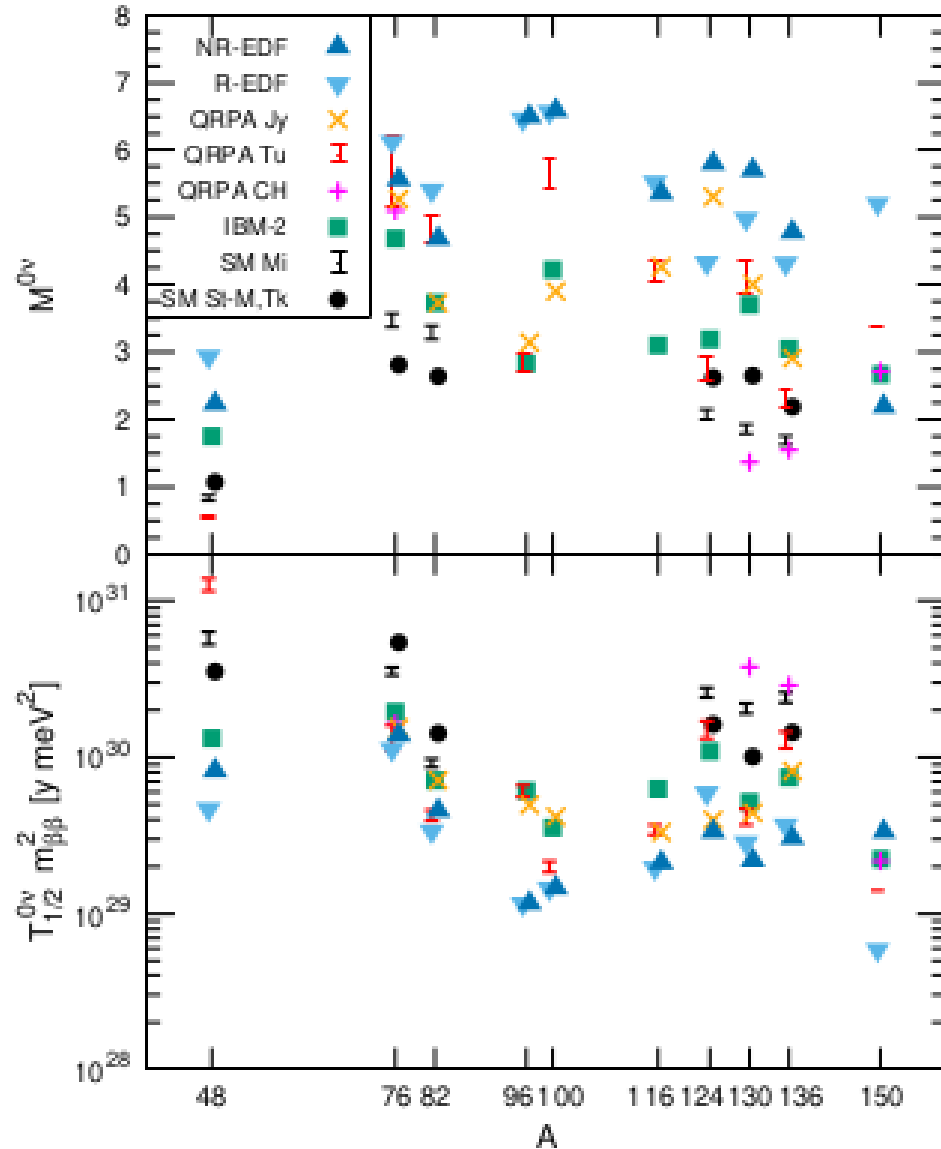
$$\boxed{\left[T_{1/2}^{0\nu} \right]^{-1}} = \boxed{G^{0\nu}} \boxed{\left| M^{0\nu} \right|^2} \left(\frac{\boxed{\langle m_{\beta\beta} \rangle}}{m_e} \right)^2$$

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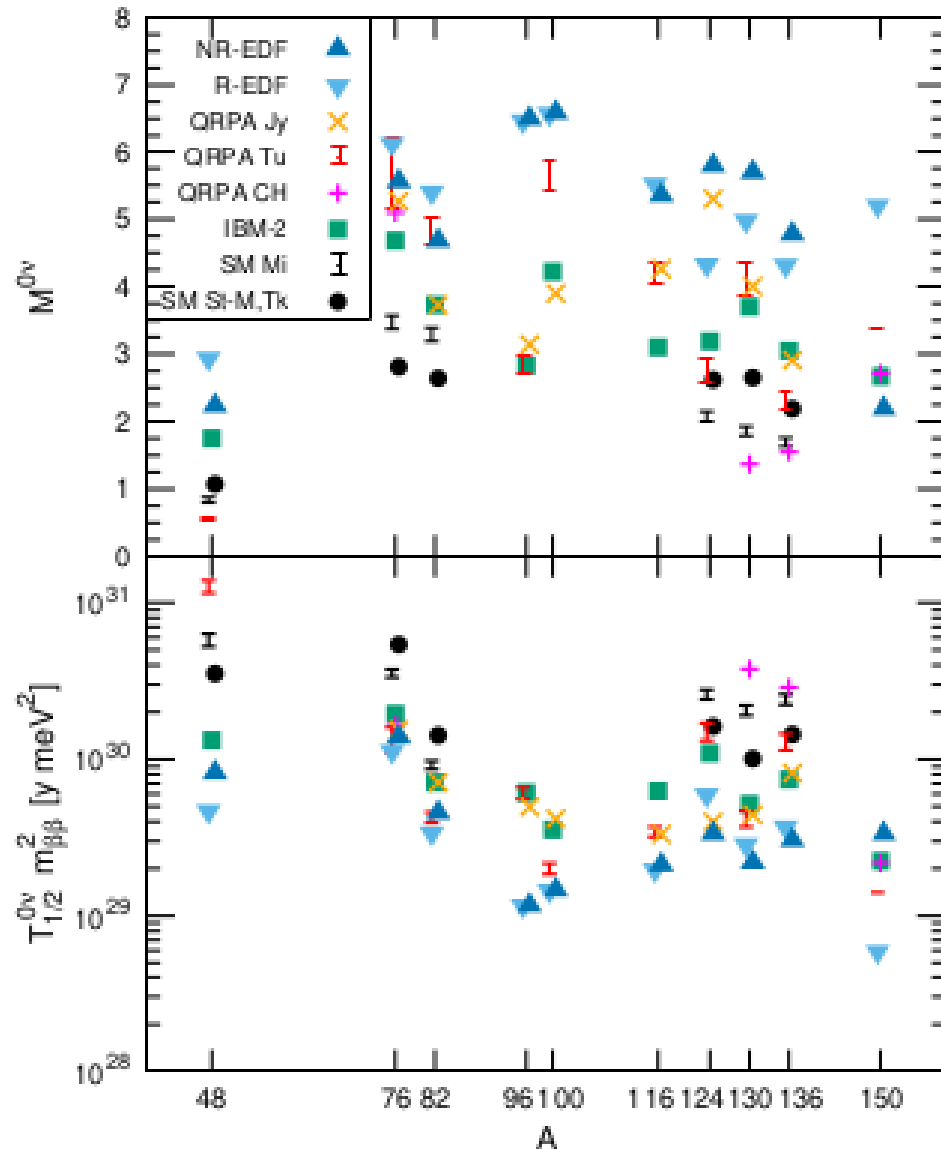
$$\langle m_{\beta\beta} \rangle = \left| \sum_{k=1}^3 m_k U_{ek}^2 \right|$$

$0\nu\beta\beta$ nuclear matrix element (NME)

This is where nuclear theory is important!



There is still a large spread in calculated $0\nu\beta\beta$ nuclear matrix elements.



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Question: what can be done to reduce NME uncertainties?

$$M^{DGT} = \langle 0_{gs,f}^+ | \sum_{j,k} [\sigma_j \tau_j^- \otimes \sigma_k \tau_k^-] | 0_{gs,i}^+ \rangle$$

- Gamow-Teller (GT) transition: neutron changes to proton and the spin from neutron to proton flips
- Double Gamow-Teller (DGT) transition: a strong nuclear process where two neutrons change to two protons with the same spin flipping behavior

$$M^{DGT} = \langle 0_{gs,f}^+ | \sum_{j,k} [\sigma_j \tau_j^- \otimes \sigma_k \tau_k^-]^0 | 0_{gs,i}^+ \rangle$$

$$M^{0\nu} = M_{GT}^{0\nu} - \left(\frac{g_V}{g_A} \right)^2 M_F^{0\nu} + M_T^{0\nu}$$

$$M_{GT}^{0\nu} = \langle 0_{gs,f}^+ | \sum_{j,k} [\sigma_j \tau_j^- \otimes \sigma_k \tau_k^-]^0 V_{GT}(r_{jk}) | 0_{gs,i}^+ \rangle$$

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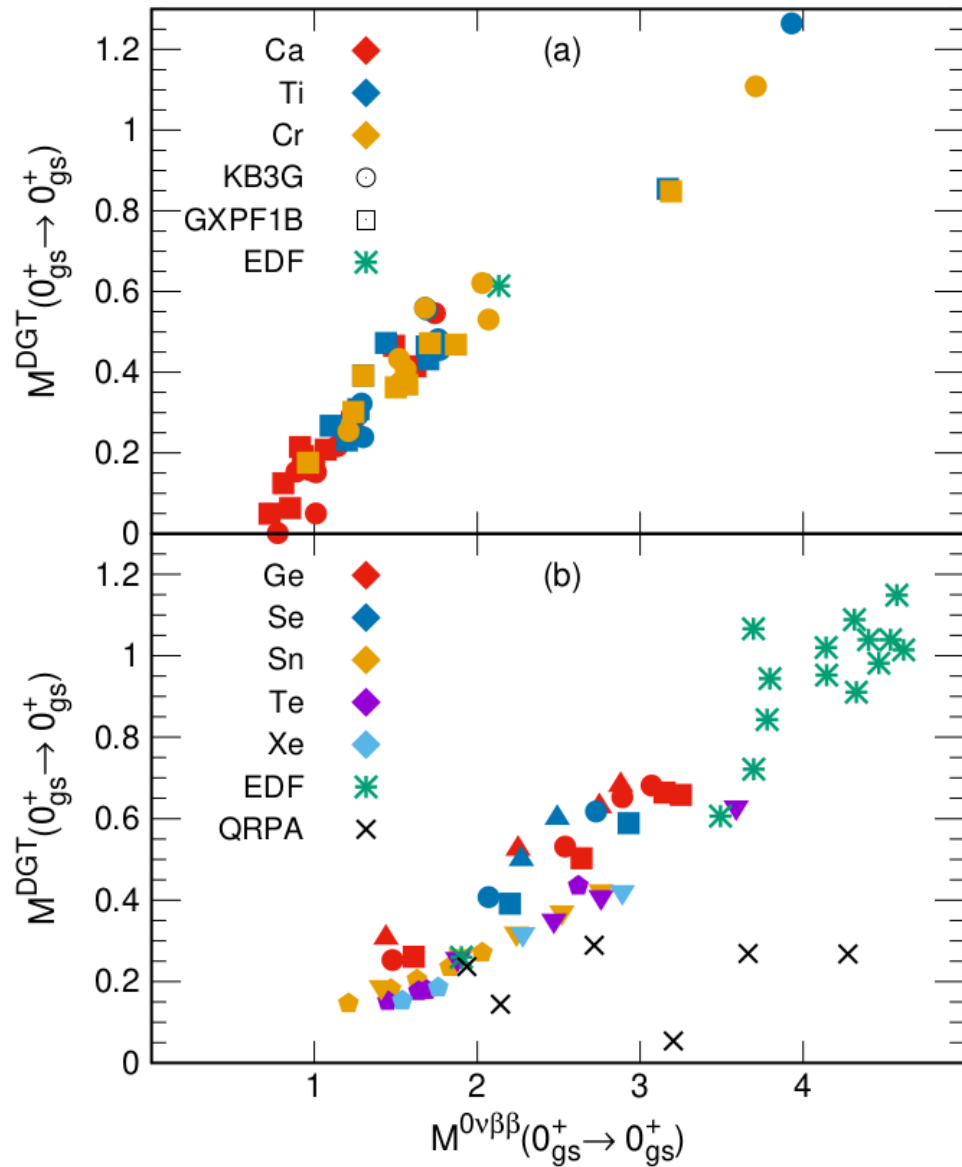
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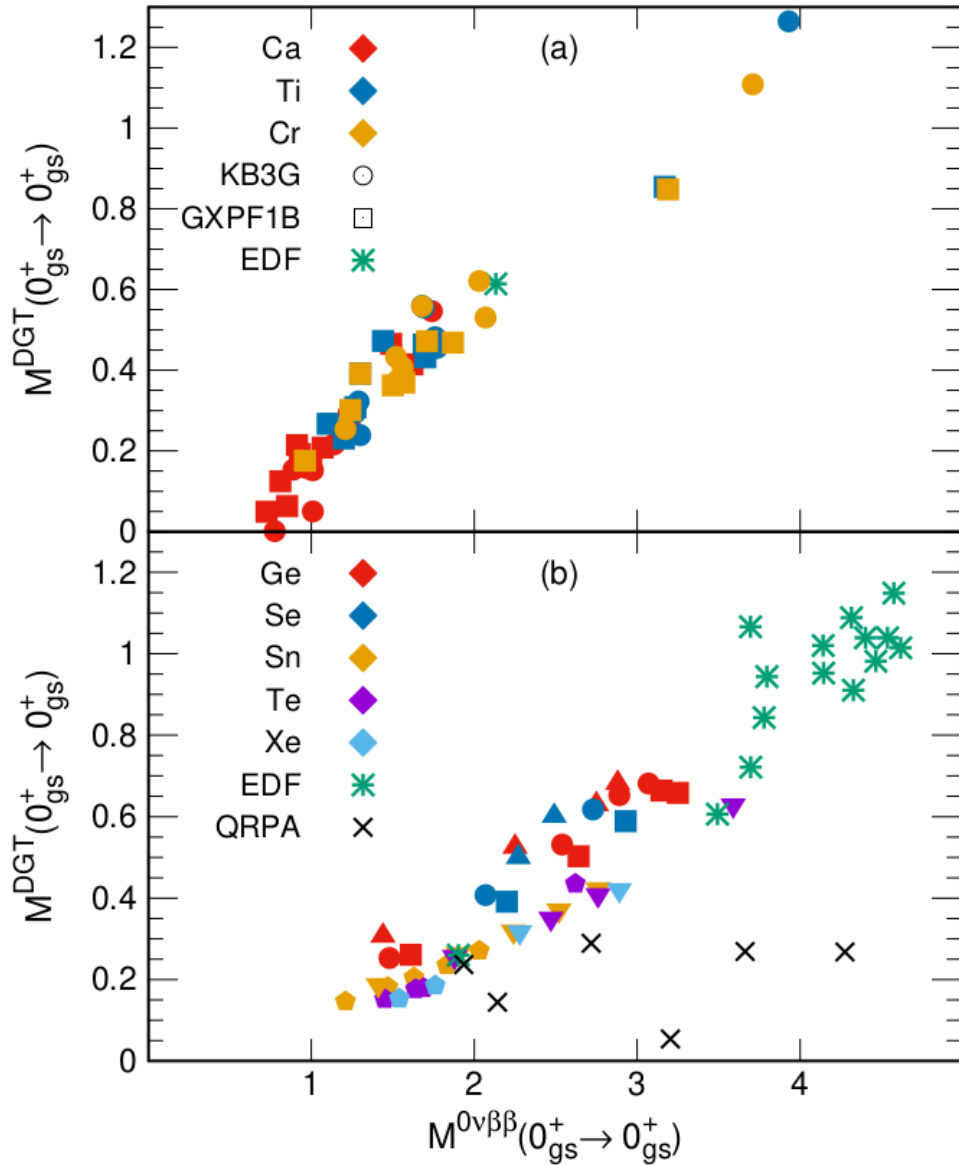
$$V_{GT}(r_{jk}) \approx \frac{1}{r_{jk}}$$

- Gamow-Teller (GT) transition: neutron changes to proton and the spin from neutron to proton flips
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Only difference is the neutrino potential!



N. Shimizu et al. demonstrated that there is a linear correlation between double Gamow-Teller (DGT) transition NMEs and $0\nu\beta\beta$ decay NMEs using phenomenological techniques.



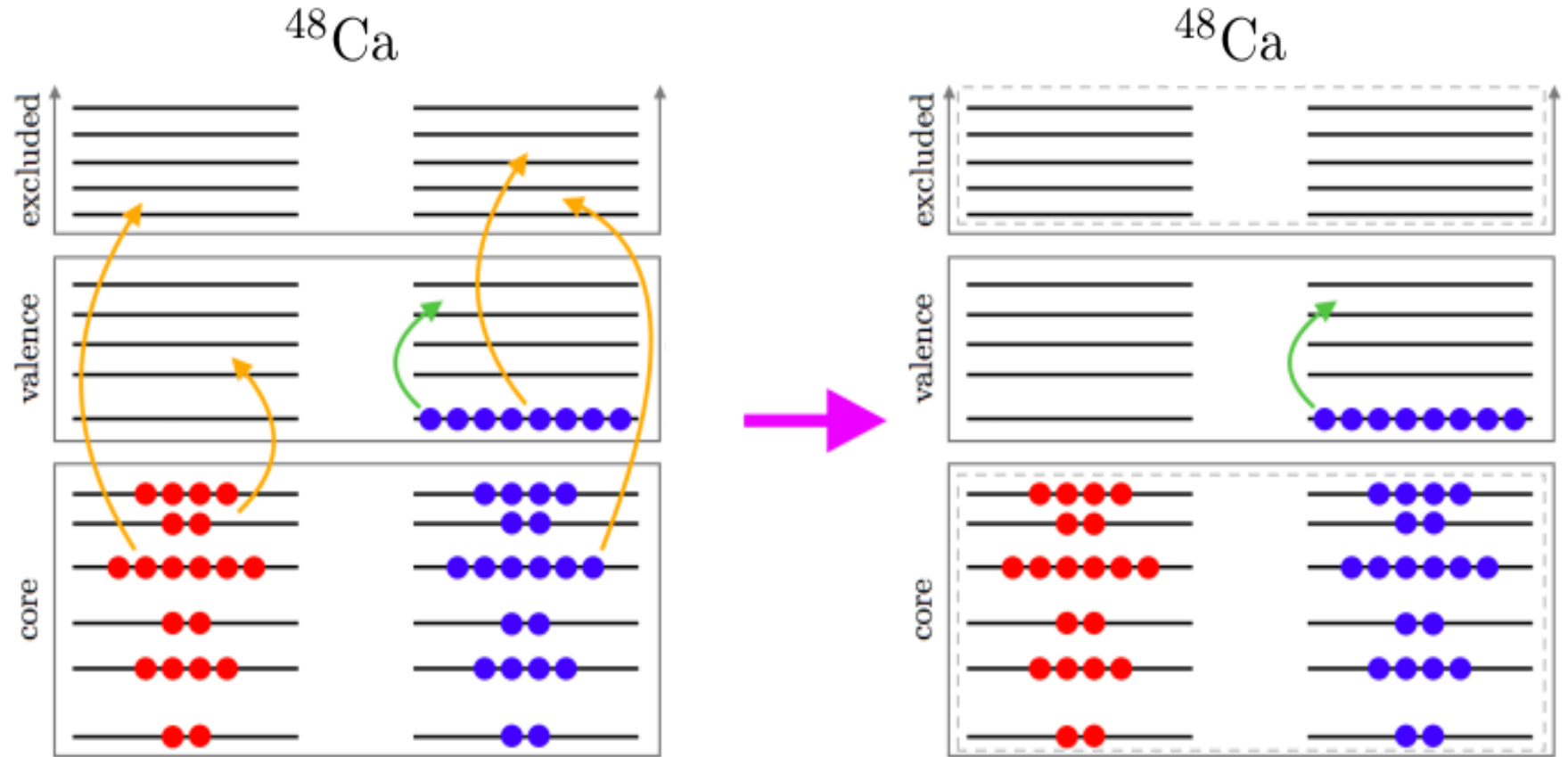
N. Shimizu et al. demonstrated that there is a linear correlation between double Gamow-Teller (DGT) transition NMEs and $0\nu\beta\beta$ decay NMEs using phenomenological techniques.

Question: does this correlation hold when using first principles computational methods?

Ab initio techniques solve the Schrödinger equation

$$\hat{H}\Psi = E\Psi$$

but the full-space Hamiltonian is difficult to solve...



Courtesy of C. G. Payne's thesis

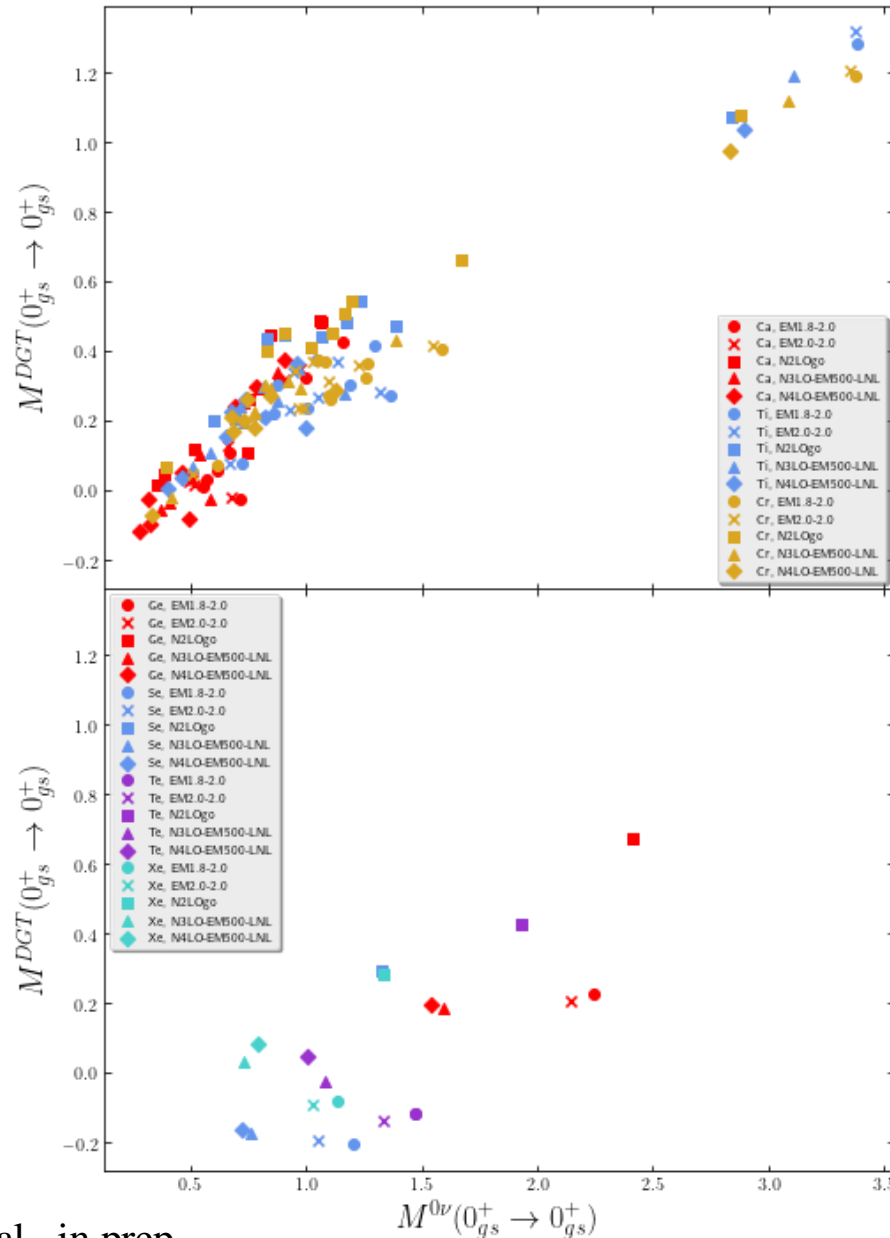
Steps of VS-IMSRG:

1. Decouple core energy and valence-space Hamiltonian from full-space Hamiltonian
2. Solve the simpler valence space Hamiltonian exactly

Isotopes considered in the plot (only even mass numbers):

- Ca: $44 < A < 58$,
- Ti: $44 < A < 60$,
- Cr: $46 < A < 62$,
- ^{76}Ge , ^{82}Se , ^{130}Te , ^{136}Xe

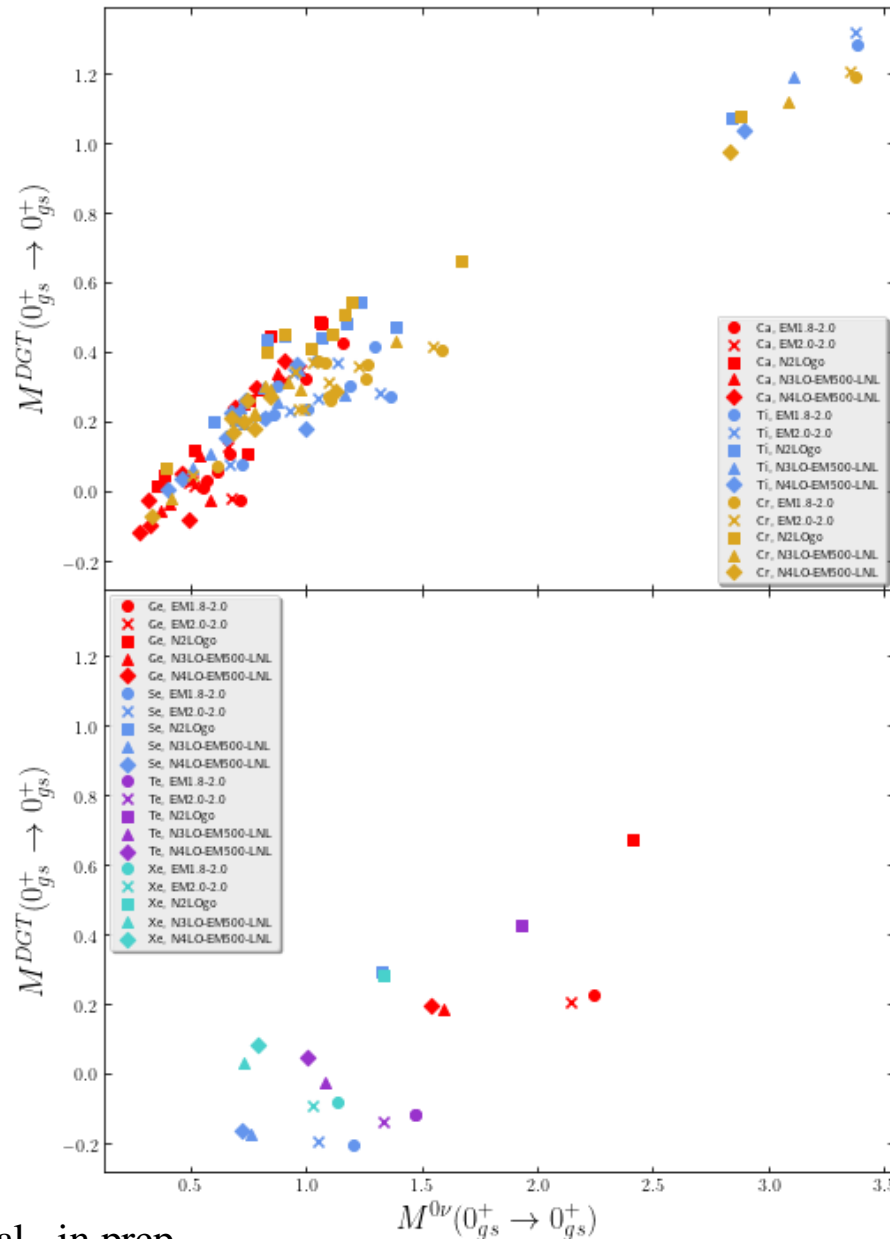
Used 5 different NN+3N chiral Hamiltonians



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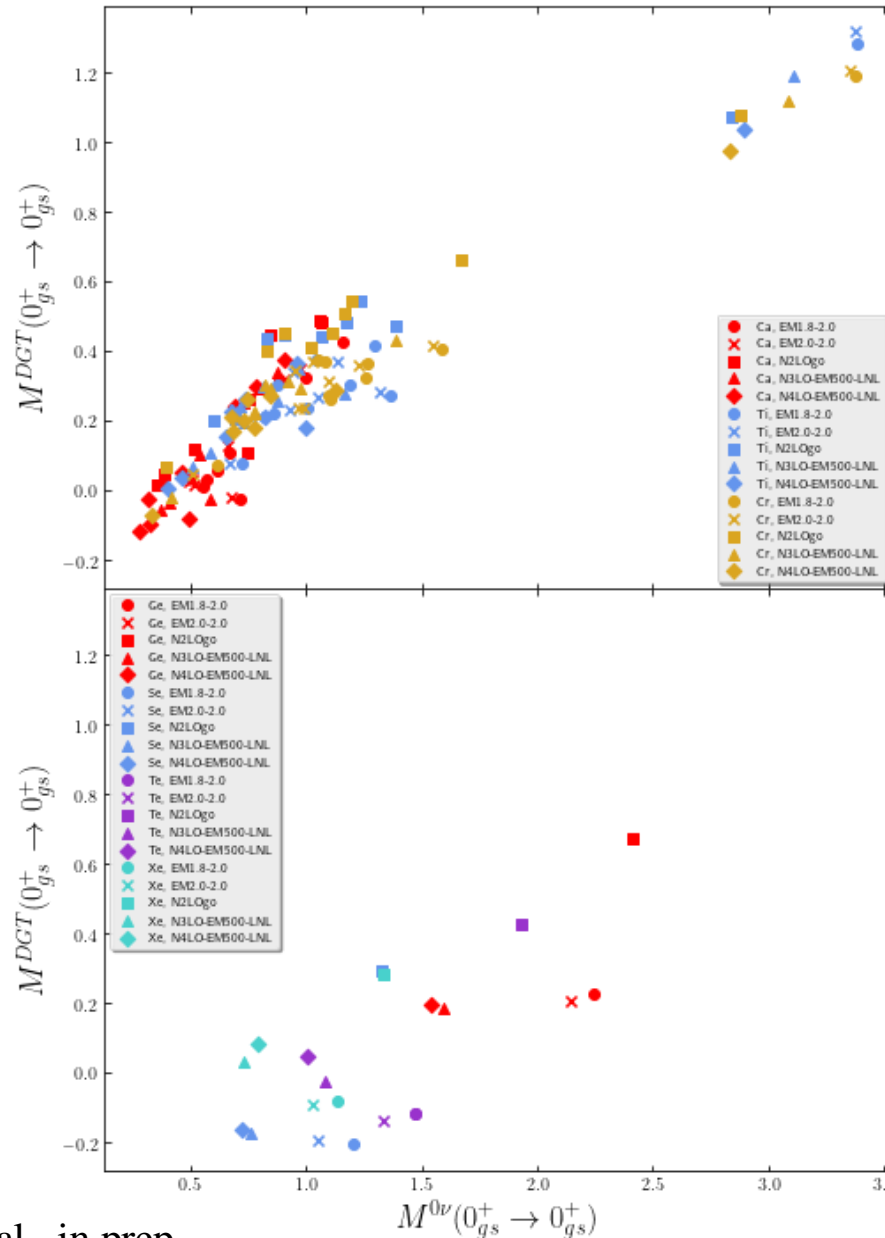


We still see a strong linear correlation between DGT NMEs and $0\nu\beta\beta$ NMEs!

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We still see a strong linear correlation between DGT NMEs and $0\nu\beta\beta$ NMEs!

This means that a process allowed by the Standard Model can constrain a Beyond-Standard-Model process.

- Experimental observation of $0\nu\beta\beta$ can have a big impact on our understanding of physics.
- Precise NME calculations are critical to constraining experimental searches.
- Previous results showed that there was a linear correlation between DGT transition NMEs and $0\nu\beta\beta$ decay NMEs.
- We showed that this linear correlation still holds when using the first principles computational technique known as VS-IMSRG.

I would like to thank the following people for their support and guidance:

- J. Holt
- A. Belley
- T. Miyagi
- S. R. Stroberg
- E. Love

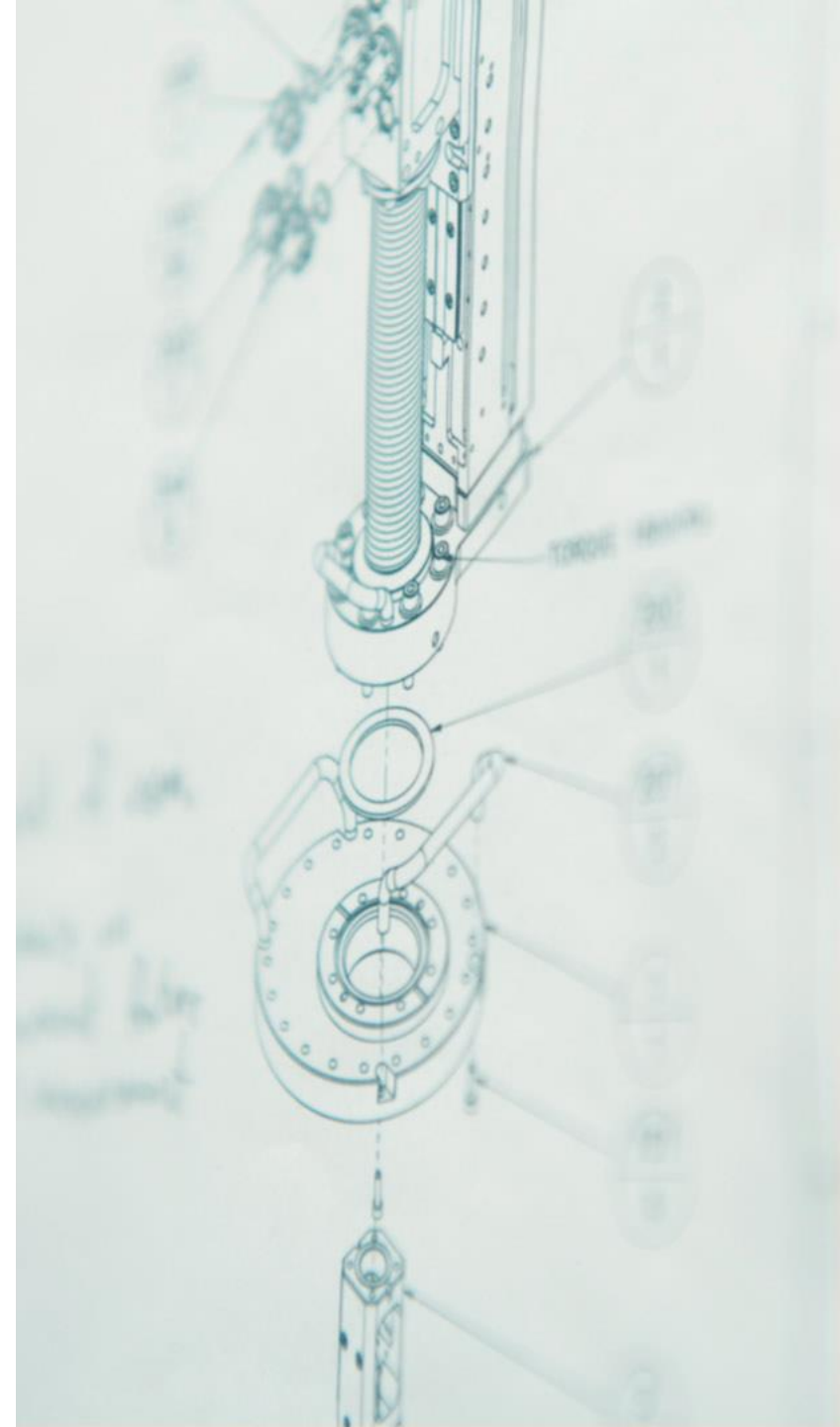
I would also like to thank Mitacs and Fulbright Canada for funding this experience and TRIUMF and the University of British Columbia for making this research experience possible.

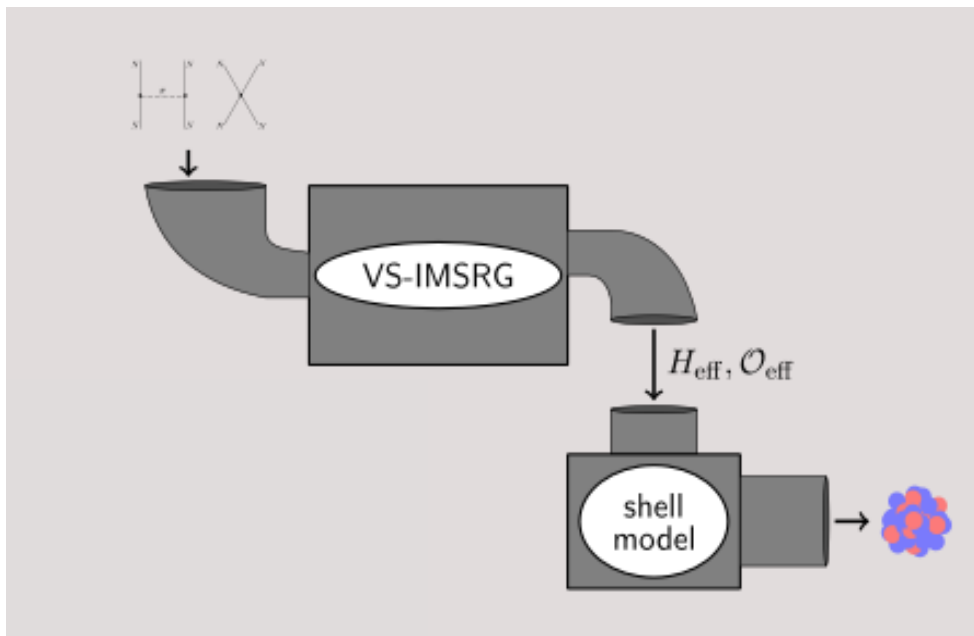


Thank you! Are there any questions?



Additional Slides





Courtesy of S. R. Stroberg

Important parameters:

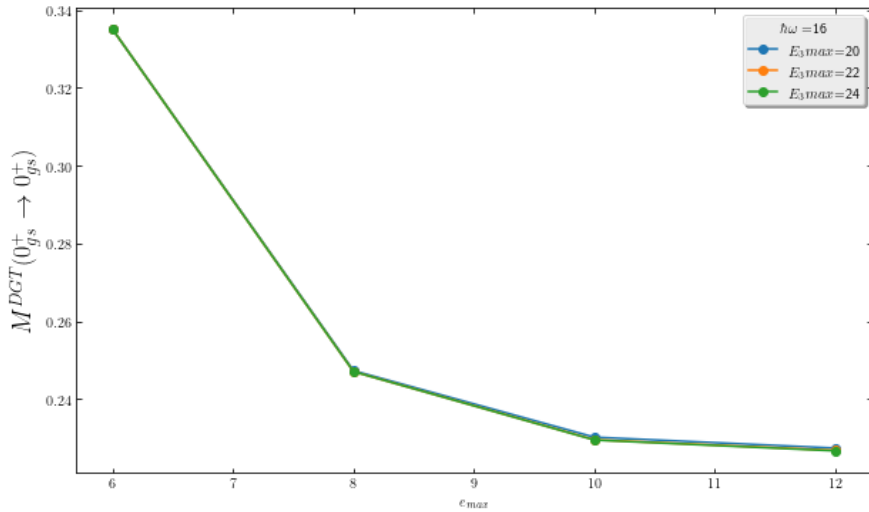
- e_{max} : harmonic oscillator basis size ($2n + l \leq e_{max}$)
- $E3_{max}$: 3-body force truncation
(ideally $E3_{max} = 3 \cdot e_{max}$)
- $\hbar\omega$: harmonic oscillator frequency

Note: all operators are currently truncated at the 2-body level

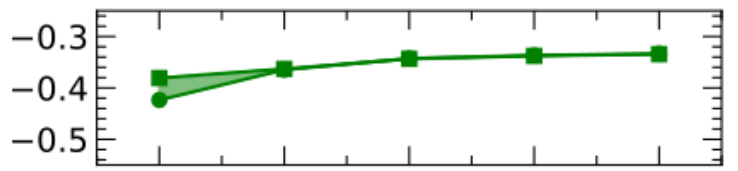
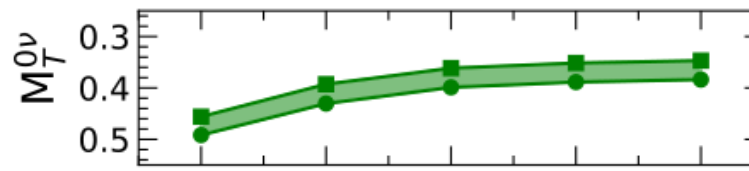
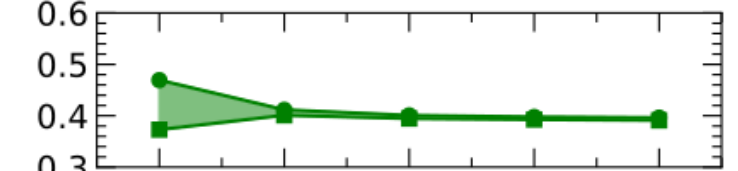
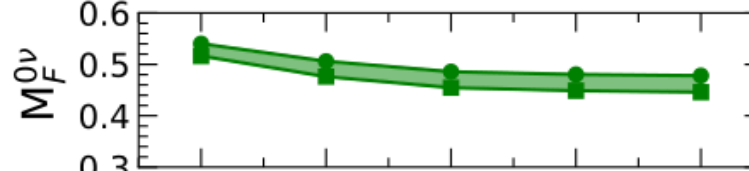
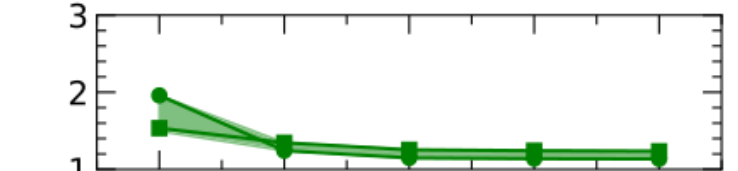
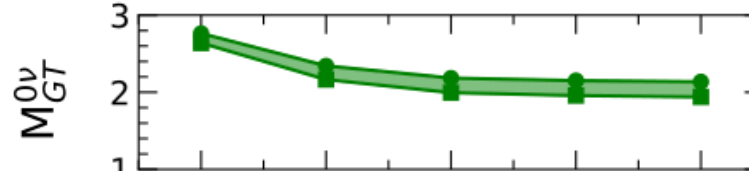
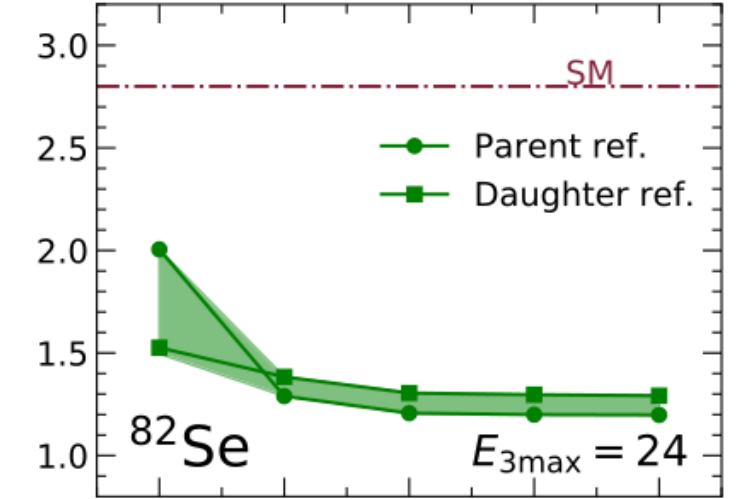
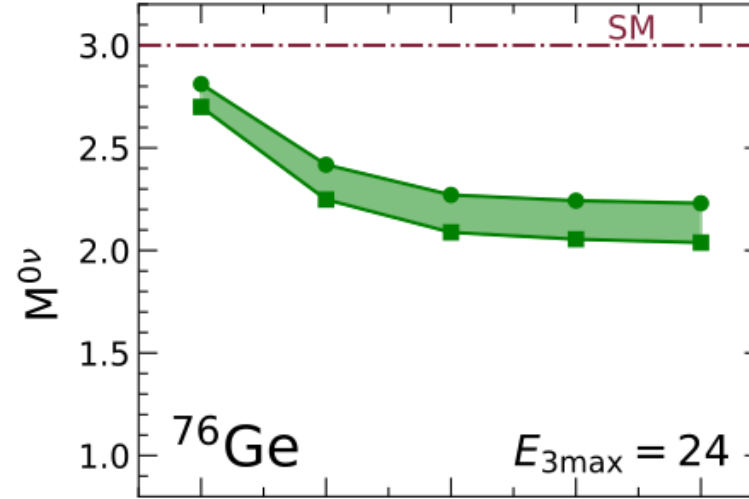
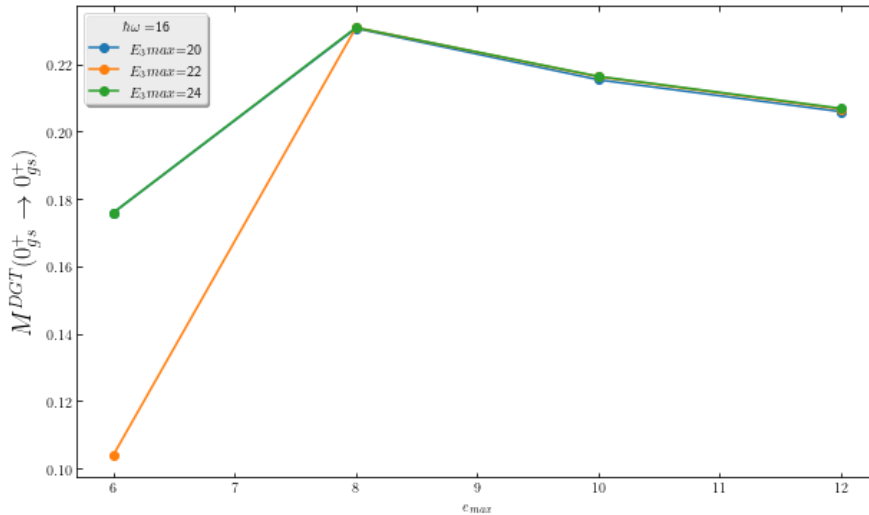
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Convergence of M^{DGT} for ^{76}Ge



Convergence of M^{DGT} for ^{82}Se



- Experimentally, DGT is yet to be observed.
- Lots of experimental interest!
- Current experimental searches use heavy ion double charge exchange interactions.
- M. Takaki et al. used $^{12}\text{C}(^{18}\text{O}, ^{18}\text{Ne})^{12}\text{Be}$ to probe DGT transitions [1].
- K. Takahisa et al. used and $^{13}\text{C}(^{11}\text{B}, ^{11}\text{Li})^{13}\text{O}$ and $^{56}\text{Fe}(^{11}\text{B}, ^{11}\text{Li})^{56}\text{Ni}$ [2].
- More experiments need to be done.

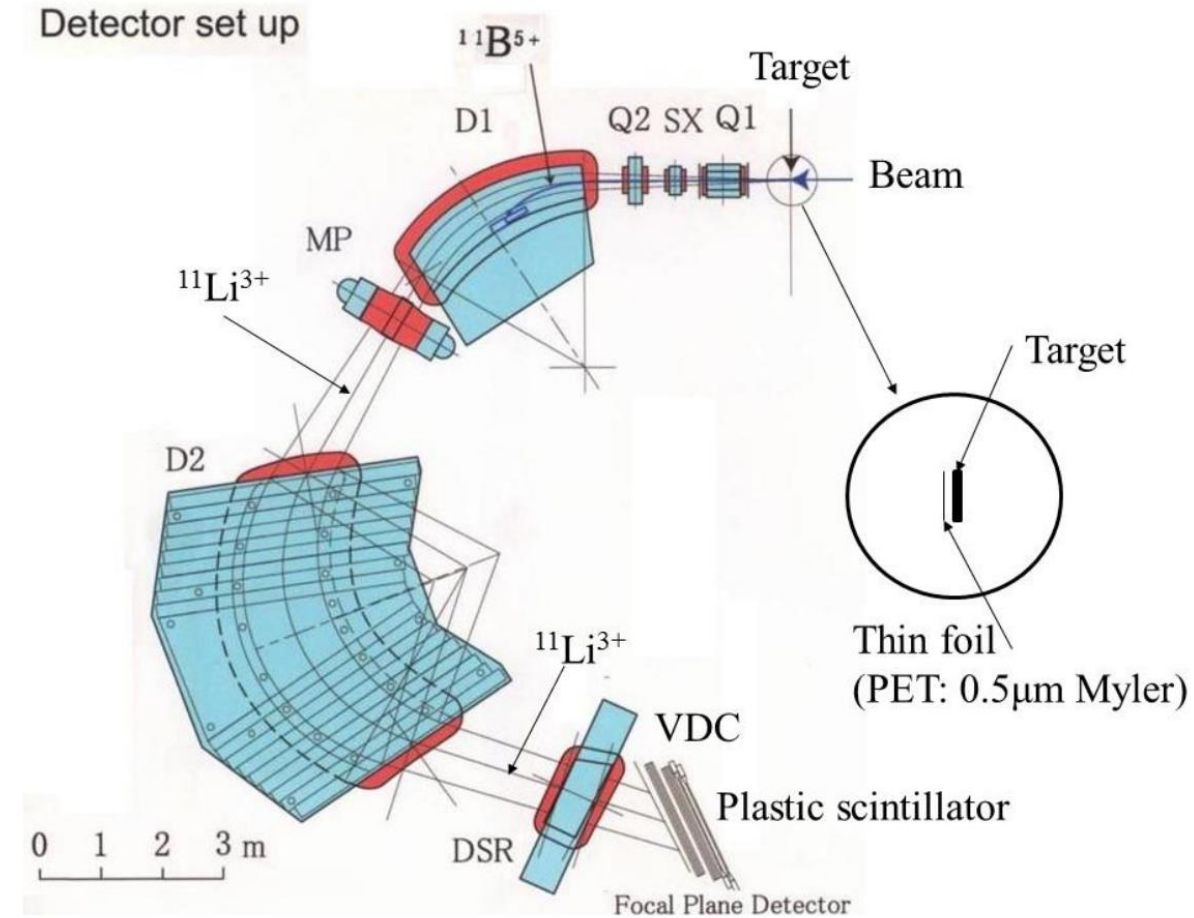


Diagram of the Grand Raiden spectrometer, Osaka University [2].

[1] M. Takaki et al., JPS Conf. Proc. **6**, 020038 (2015).

[2] K. Takahisa et al., arXiv:1703.08264, (2017).