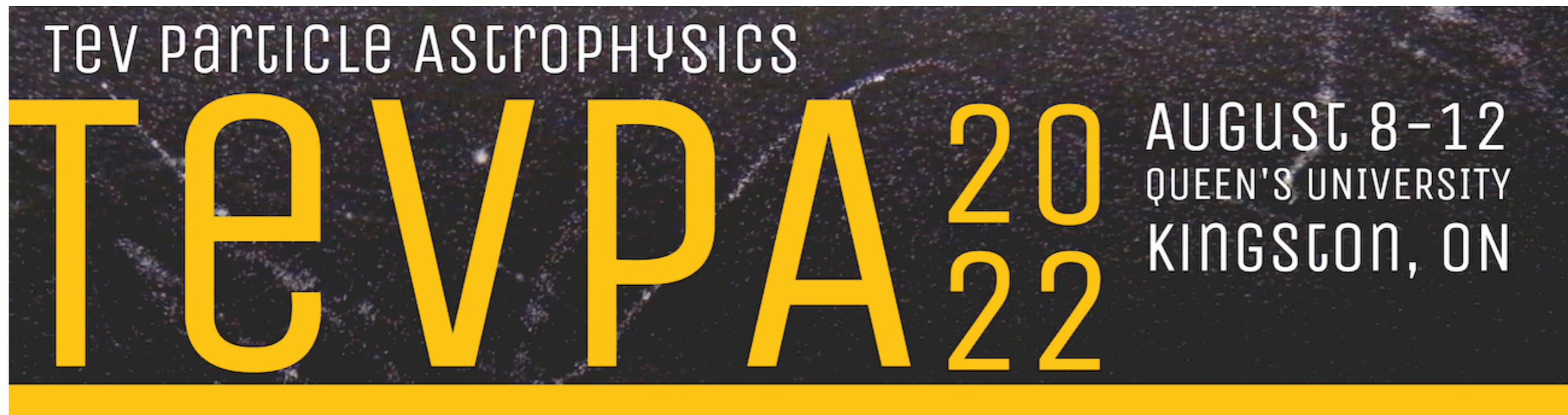


# Direct Detection of Low Mass Fast Moving Dark Matter

**Gopolang (Gopi) Mohlabeng**

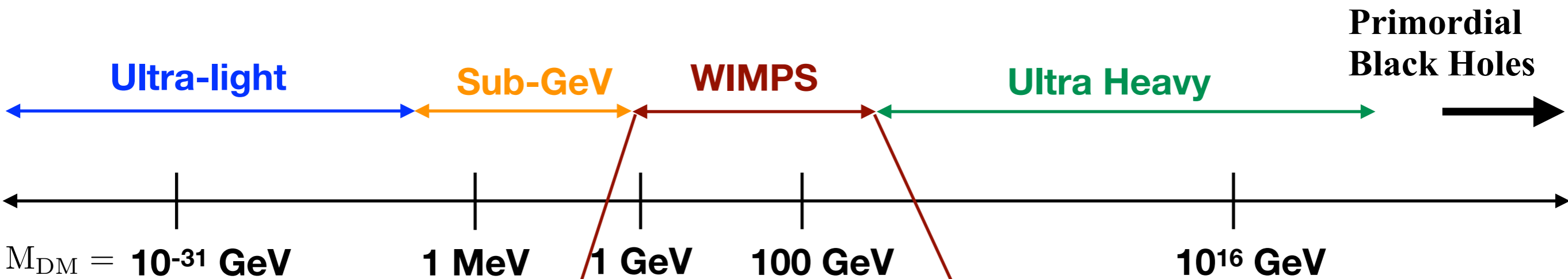
University of California, Irvine



10 August 2022



# Range of DM possibilities is VAST

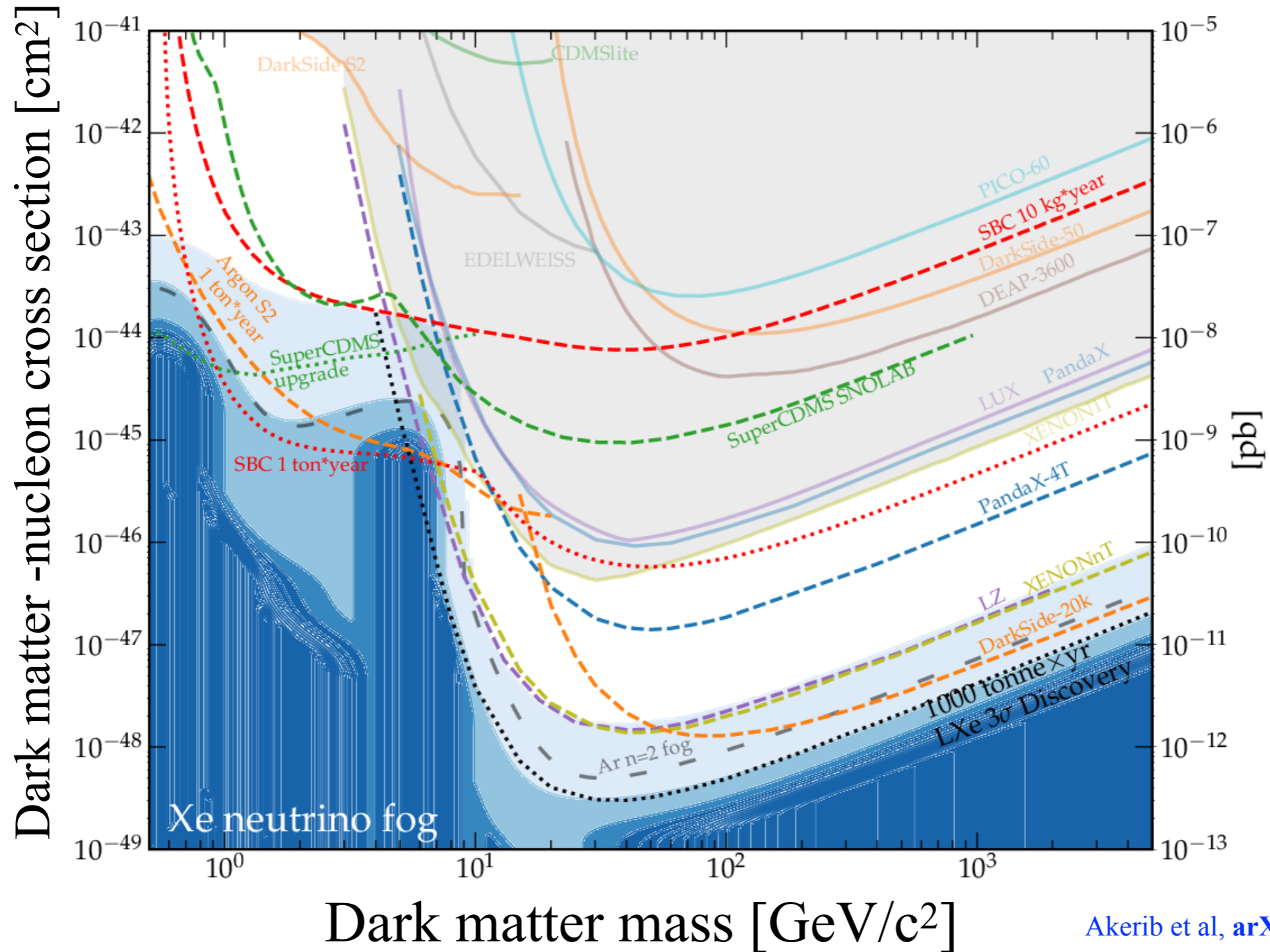


For the last decades,  
searches have been  
focussed on this  
region of space

WIMP Hypothesis is very motivated



# Current status of DM direct detection

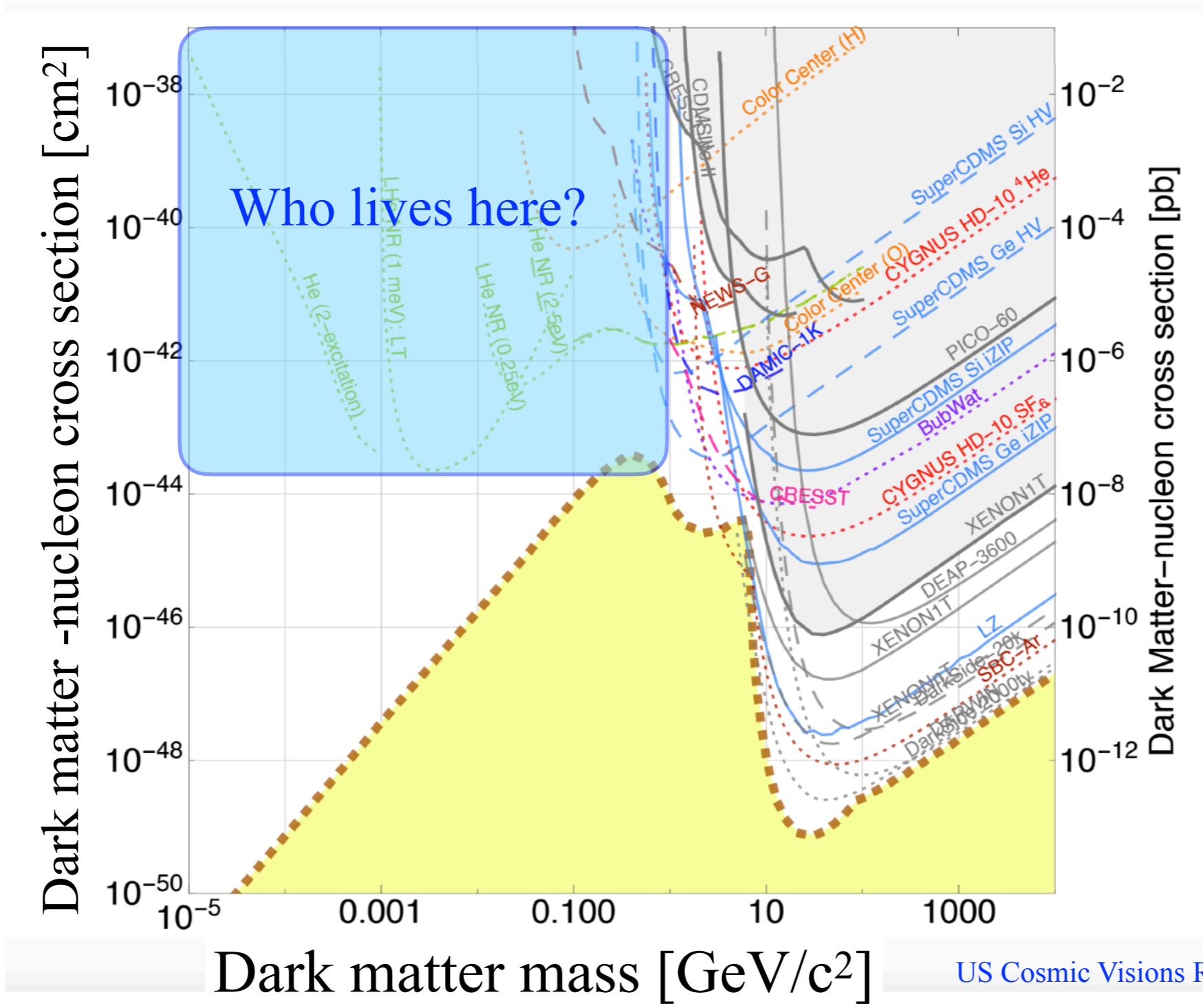


Very little parameter space left in *traditional* WIMP mass range (1 GeV - 100 TeV)

Larger detectors gives us more sensitivity, but “**Neutrino Floor/fog**” may be challenging

WIMPs remain highly motivated

We are compelled to move beyond the WIMP scale

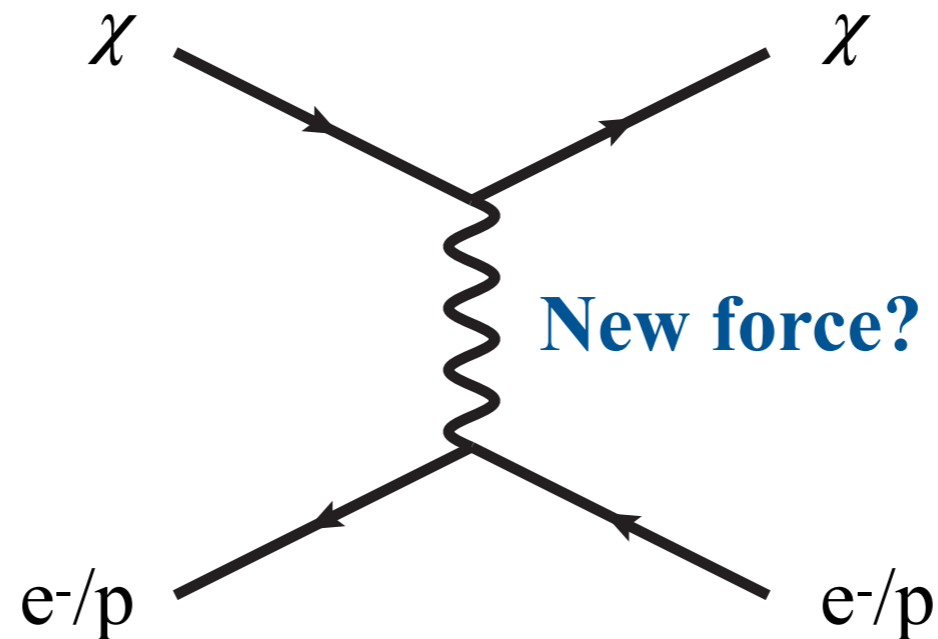


Many opportunities to explore here



Low mass (sub-GeV) DM searches are:

1. Complementary to WIMP searches
2. Very well motivated:



e.g. low mass DM coupled through new hidden sector mediator presents a good target to understand DM production in the early universe

- Thermal freeze-out
- Freeze-in
- Asymmetric



Big stumbling block is that light DM is **invisible** to current WIMP searches

Many nuclear recoil experiments have  $\sim$  keV recoil thresholds

Recoil energy caused by galactic DM with  $v \sim 10^{-3} c$

$$\begin{aligned} E_{nr} &\sim \frac{\mu_{\chi N}^2 v^2}{m_N} \\ &\sim 19 \text{ eV} \left( \frac{m_\chi}{500 \text{ MeV}} \right)^2 \left( \frac{100 \text{ GeV}}{m_N} \right) \\ &\ll \text{keV} \end{aligned}$$

To probe lower masses need very low threshold nuclear recoil detectors

or

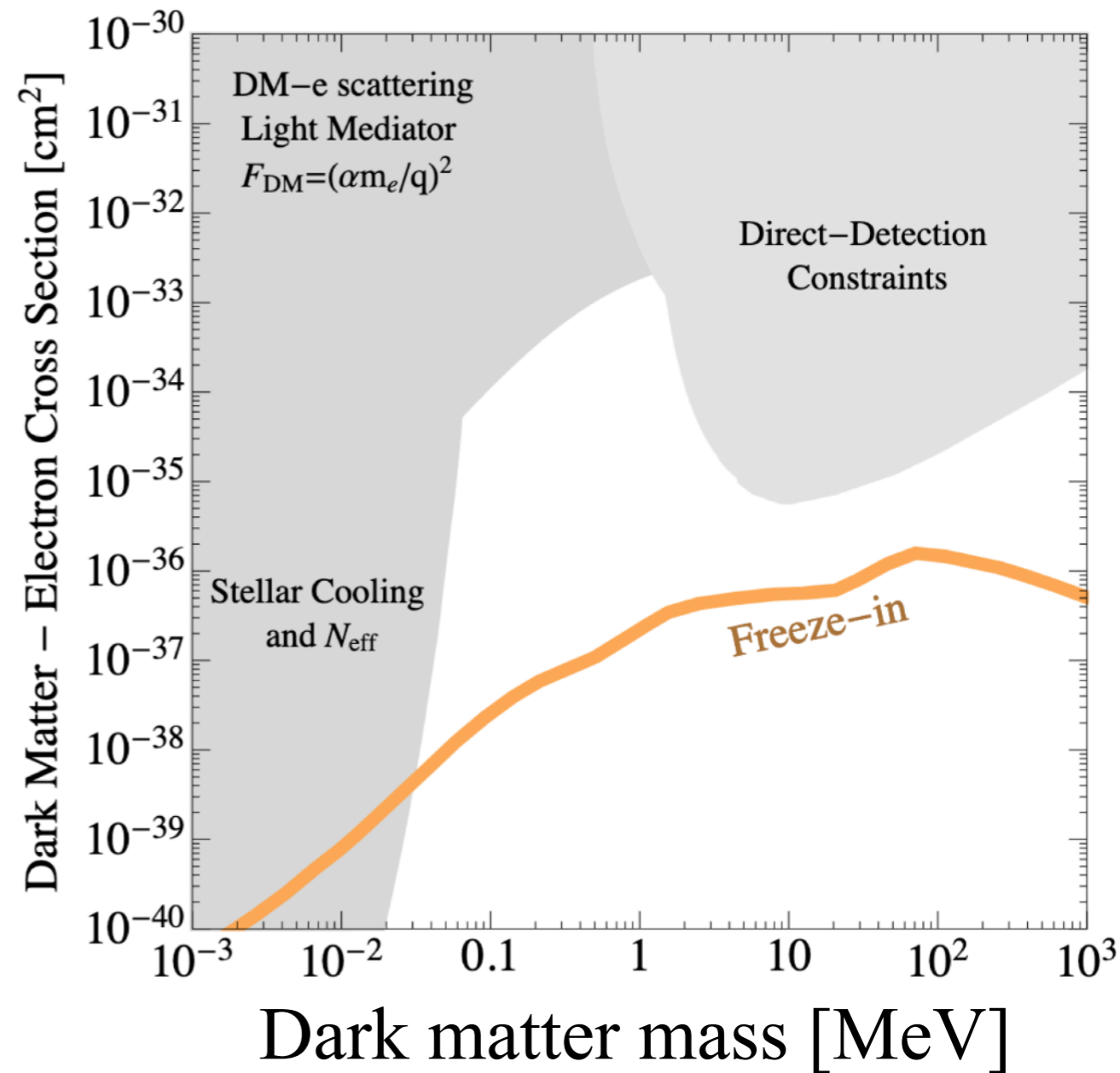
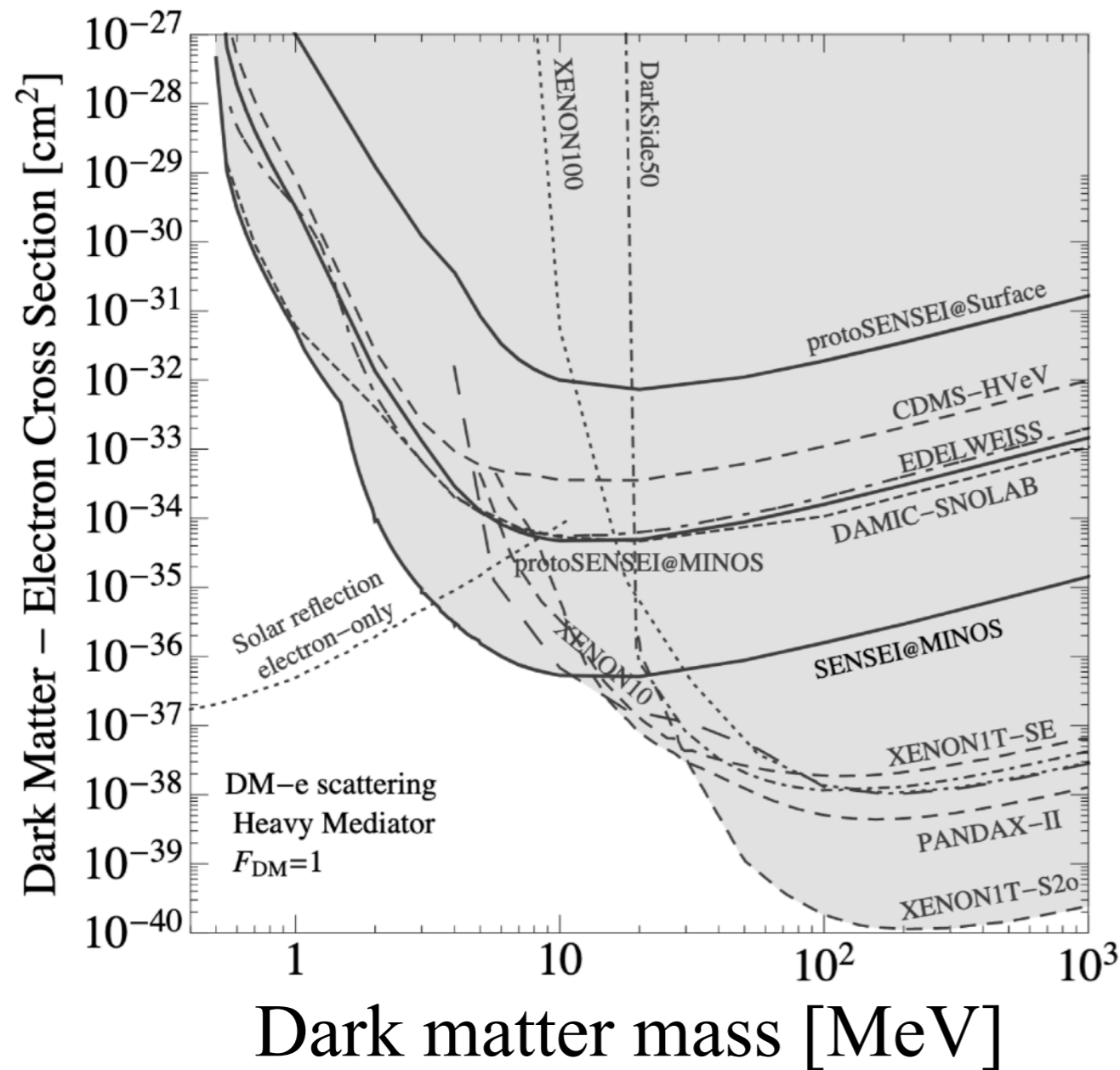
Other ways to maximize energy transfer to the target



Look for DM hitting lighter target to maximize energy transfer

## Dark matter - electron scattering

- Current searches include:
- Large volume noble liquid detectors
  - Small scale semi-conductor detectors



Complementary way to maximize energy transfer to target is

## Fast moving/accelerated dark matter

- Energy transferred to nucleus

$$E_{nr} \sim 50 \text{ keV} \left( \frac{m_\chi}{500 \text{ MeV}} \right)^2 \left( \frac{100 \text{ GeV}}{m_N} \right) \left( \frac{v}{0.1} \right)^2$$

- Max energy transferred to electron

$$E_{er} \lesssim 2 \text{ MeV} \left( \frac{m_\chi}{500 \text{ MeV}} \right) \left( \frac{v}{0.1} \right)^2$$

These are above  $\sim$  keV threshold energies in **current dark matter detectors** & actually **some large neutrino detectors**

**But**, galactic DM has to be **non-relativistic**

Accelerated DM must be small subcomponent total halo DM

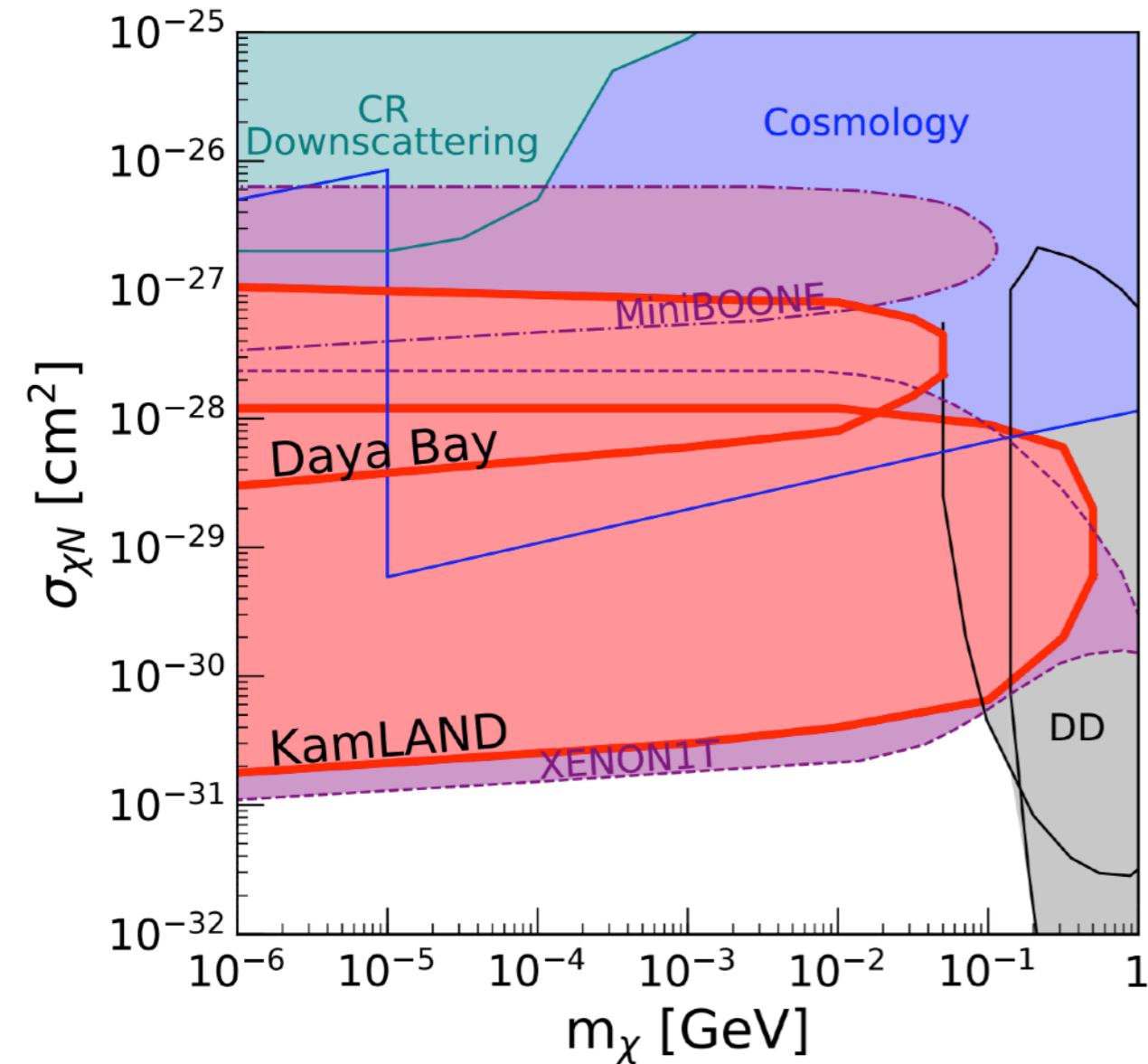
**Many ways to get accelerated dark matter**



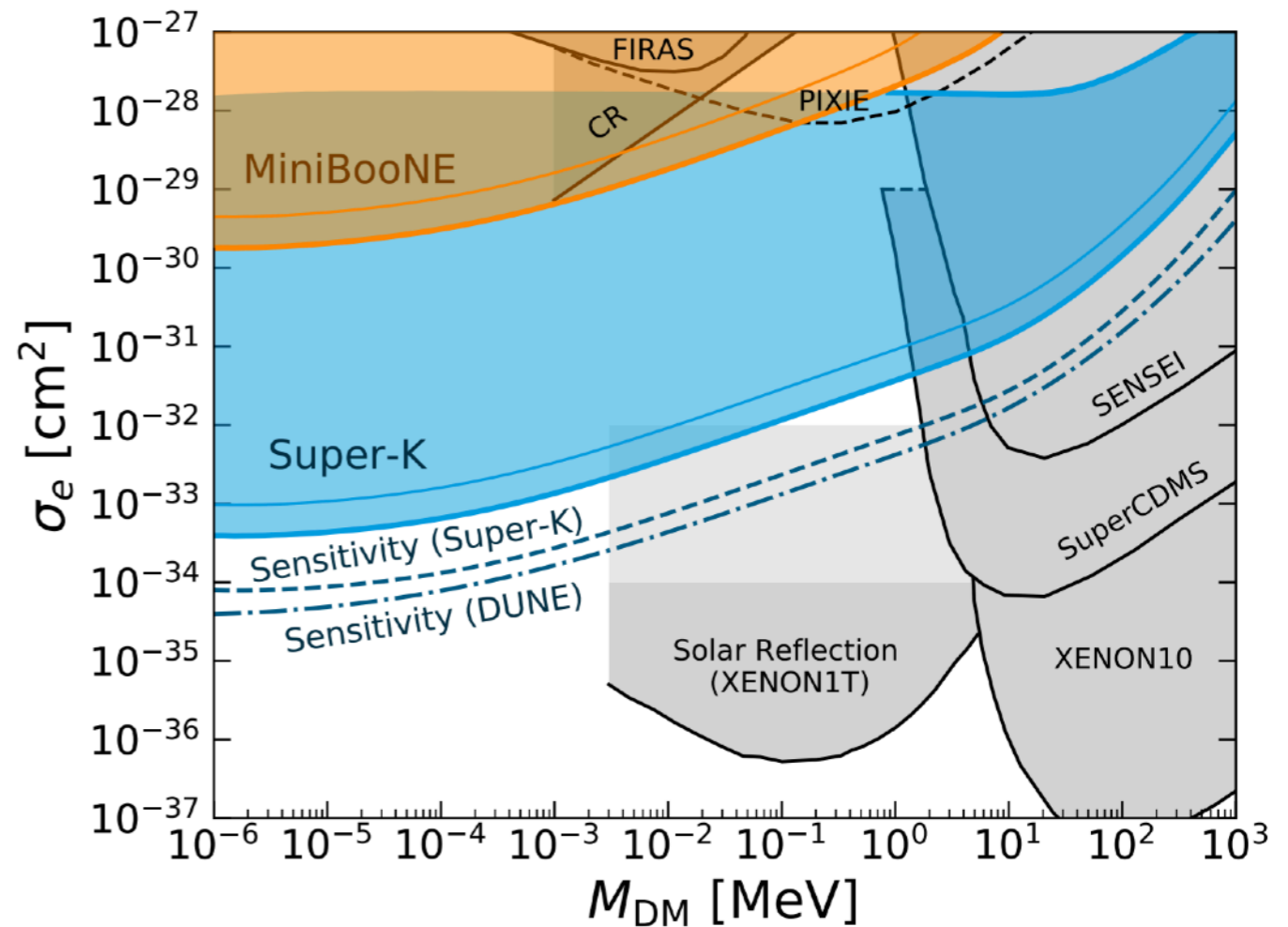


# Cosmic ray accelerated dark matter

- High energy CR can scatter off DM moving at  $v \sim 10^{-3} c$
- CR transfers energy to DM, accelerating it
- Accelerated DM reaches earth and scatters in detectors



Cappiello, Beacom : [Phys.Rev. D100 \(2019\) 10, 103011](#)



Ema, Sala, Sato : [Phys.Rev. Lett 122 \(2019\) 19, 181802](#)

Cappiello, Ng, Beacom : [Phys.Rev. D99 \(2019\) 6, 063004](#)

Dent et al : [Phys.Rev. D101 \(2020\) 11, 116007](#)

Krnjaic, McDermott : [Phys.Rev. D101 \(2020\) 12, 123022](#) + many others

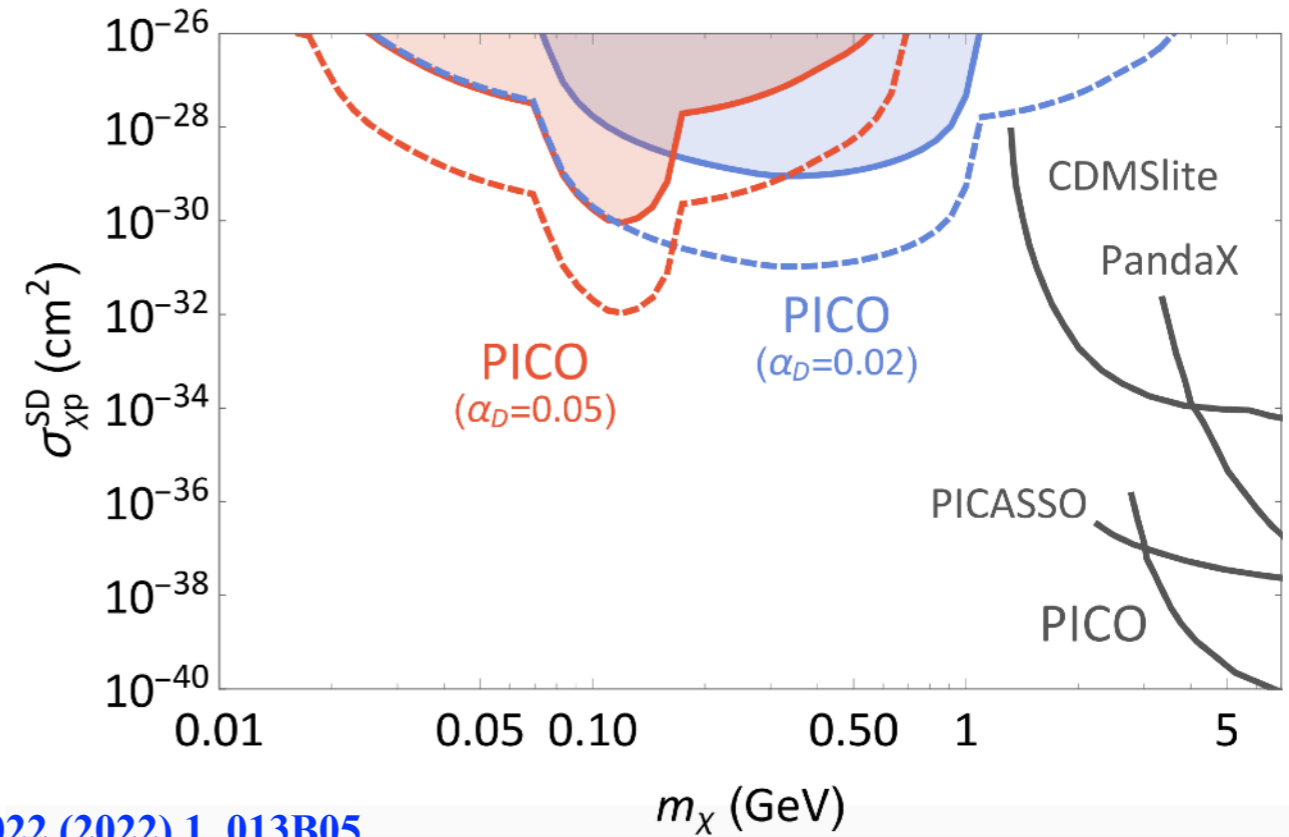
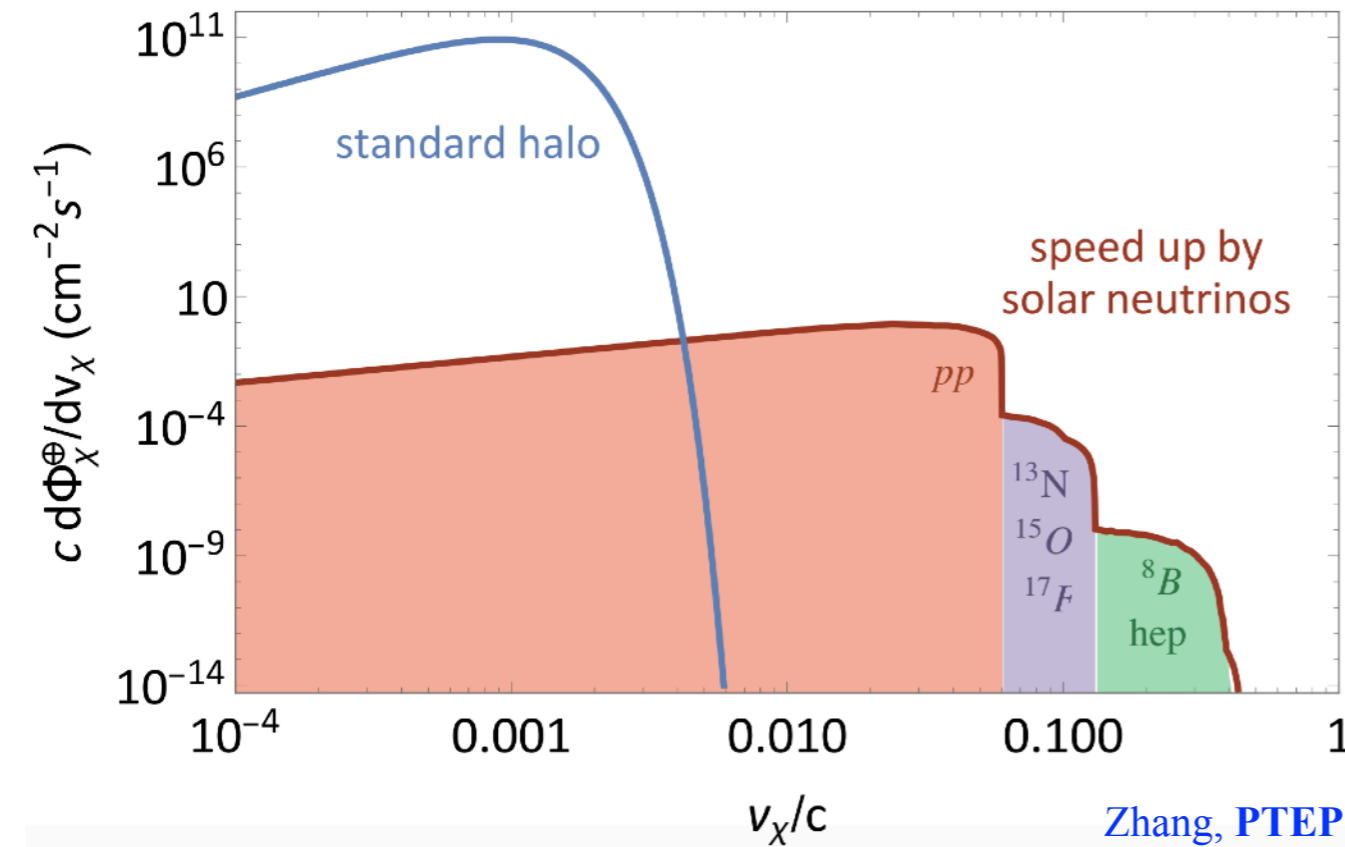


# Neutrino accelerated dark matter

- e.g. Solar neutrinos can scatter off DM moving at  $v \sim 10^{-3} c$

- After interaction, DM obtains velocity:  $v_\chi \sim \frac{2E_\nu}{m_\chi} \cos\theta$

- DM enters detector, transferring to target energy:  $E_{nr} \lesssim \frac{2E_\nu^2}{m_N}$



Zhang, PTEP 2022 (2022) 1, 013B05

Cosmic neutrino upscattering: [Jho et al, arXiv: 2101.11262](#)

See talk by Jong-Chul Park

DSNB neutrino upscattering: [Das, Sen, Phys.Rev.D 104 \(2021\) 7, 075029](#)



# Other interesting mechanisms

## Accelerated DM from evaporating primordial black holes:

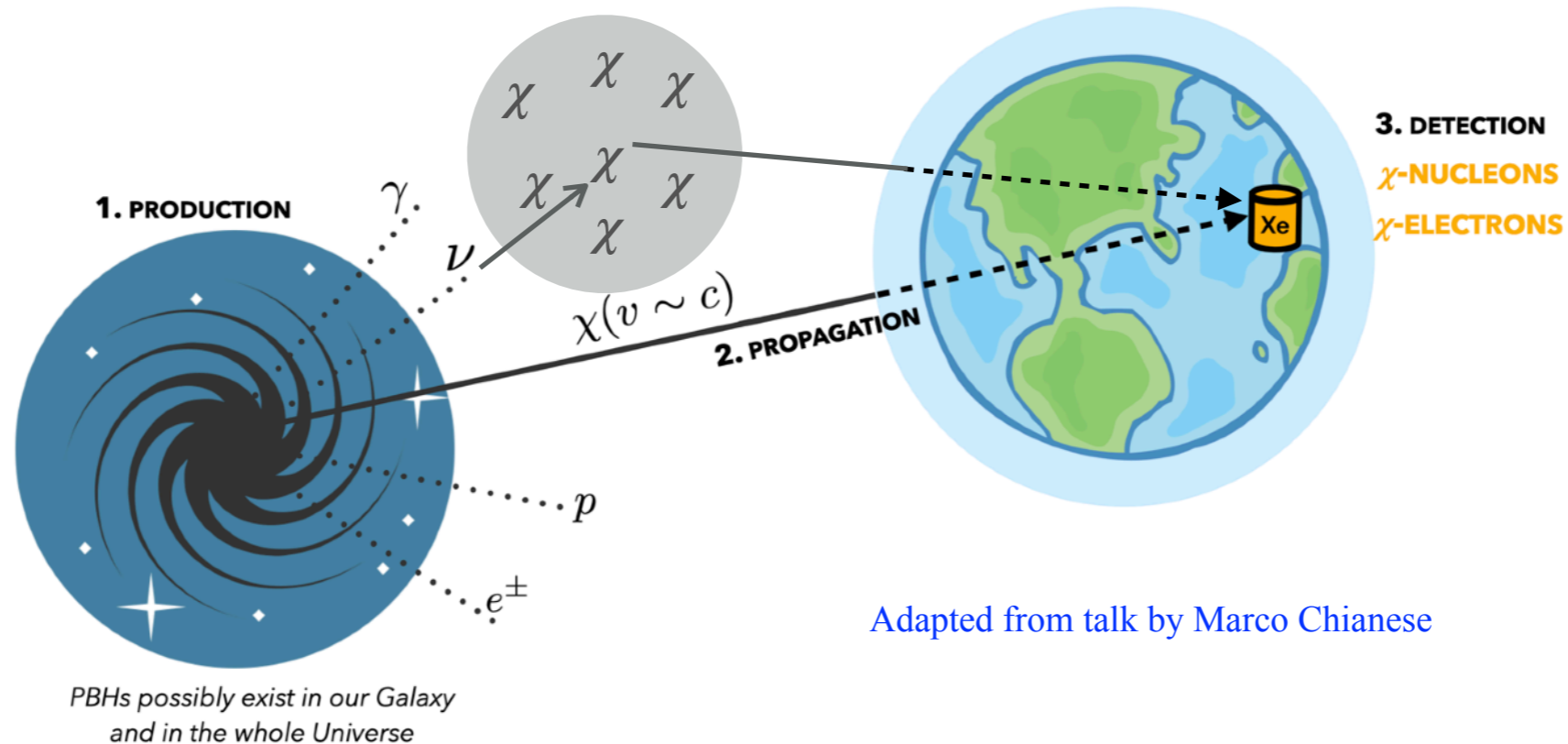
Calabrese et al, *Phys. Rev. D* (105) 2022 2

Calabrese et al, *Phys. Rev. D* (105) 2022 10

See talk by Marco Chianese

## Neutrinos from evaporating primordial black holes accelerate DM:

Chao et al, [arXiv: 2108.05608](https://arxiv.org/abs/2108.05608)



Adapted from talk by Marco Chianese

## Accelerated DM from supernova shockwaves:

See talk by Chris Cappiello

+ many others



# A different mechanism...

Consider non-minimal **Dark sector**

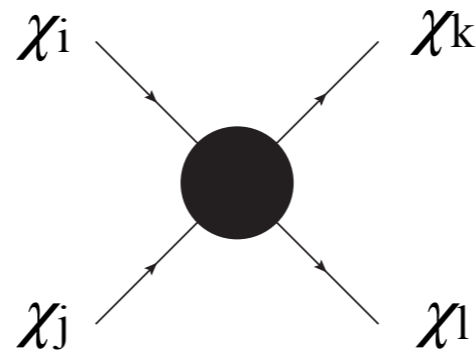
**Generic feature** of any non-minimal dark sector: small fraction of DM today may be **relativistic/semi-relativistic**

**Sources:**

- **Assisted Freeze-out:**

$$\chi_i \chi_j \rightarrow \chi_k \chi_l$$

$$\text{with } \chi_i \chi_j > \chi_k \chi_l$$



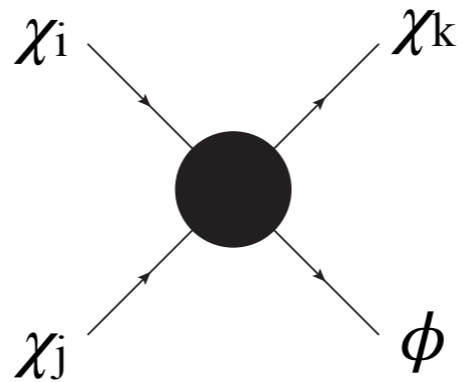
Belanger and Park : *JCAP* 03 (2012) 038

**See talk by Seodong Shin**

- **Semi-annihilation:**

$$\chi_i \chi_j \rightarrow \chi_k \phi$$

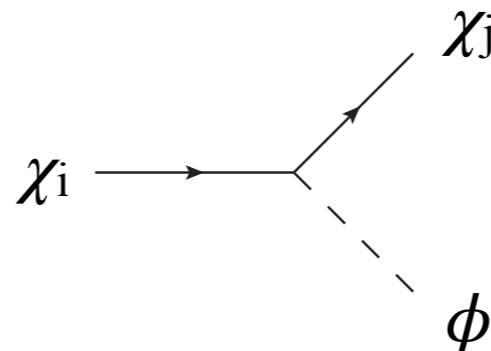
,  $Z_3$  DM symmetry



D'Eramo and Thaler : *JHEP* 06 (2010) 109

- **Decay:** (e.g inelastic DM)

$$\chi_i \rightarrow \chi_j \phi$$



Tucker-Smith & Weiner: *Phys.Rev.* D64 (2001) 043502

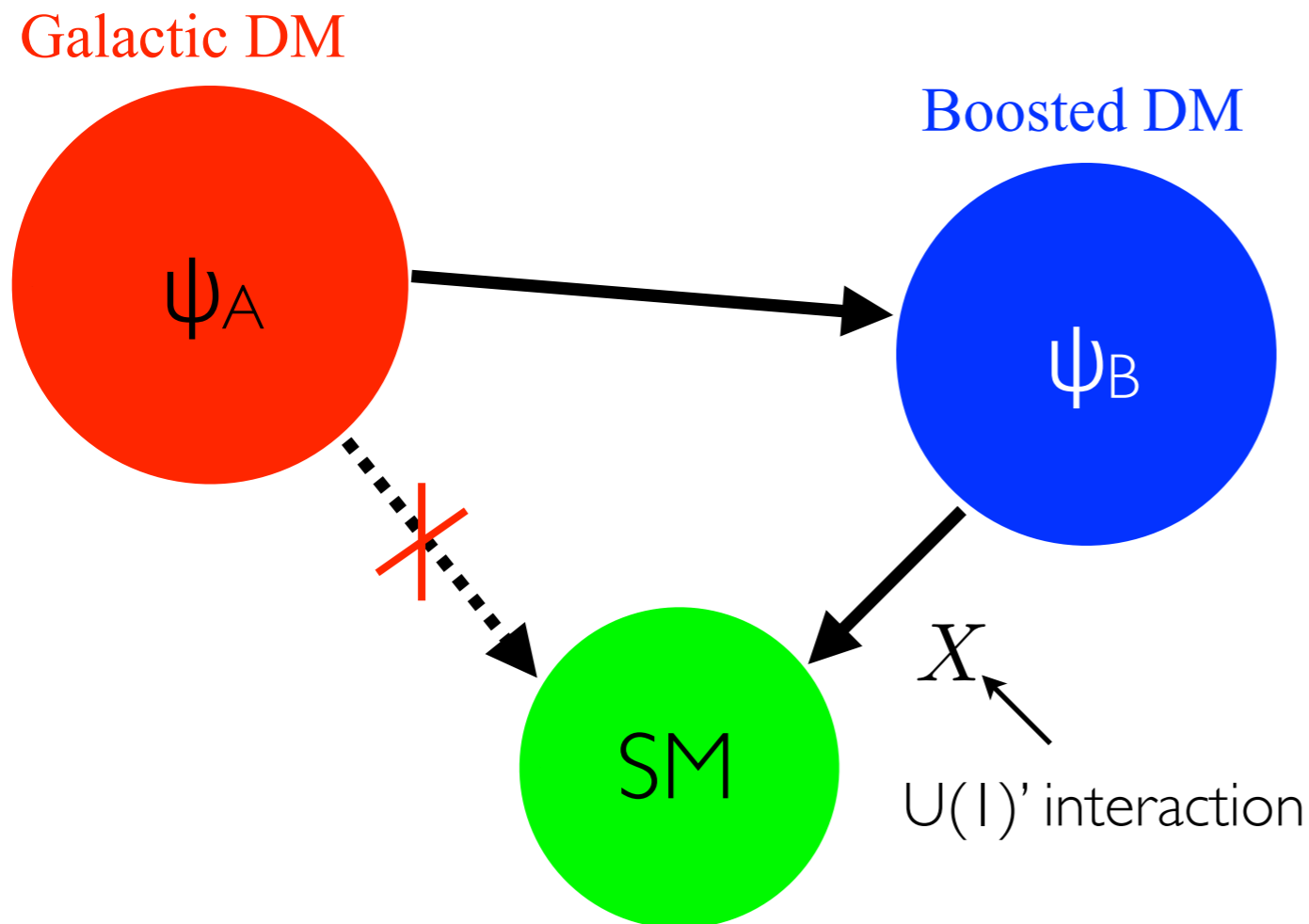
+ many others



# Boosted dark matter

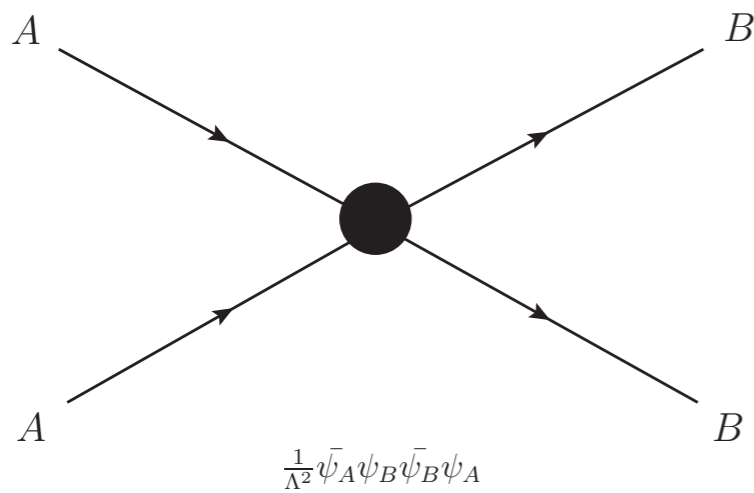
- Two stable DM particles,  $\psi_A$  &  $\psi_B$  with  $m_A > m_B$  (eg.  $U(1)' \otimes U(1)''$ )

For example:



- $\psi_A$  is the dominant DM component and has no direct coupling to SM

- $\psi_B$  is sub-dominant and couples to SM through new force

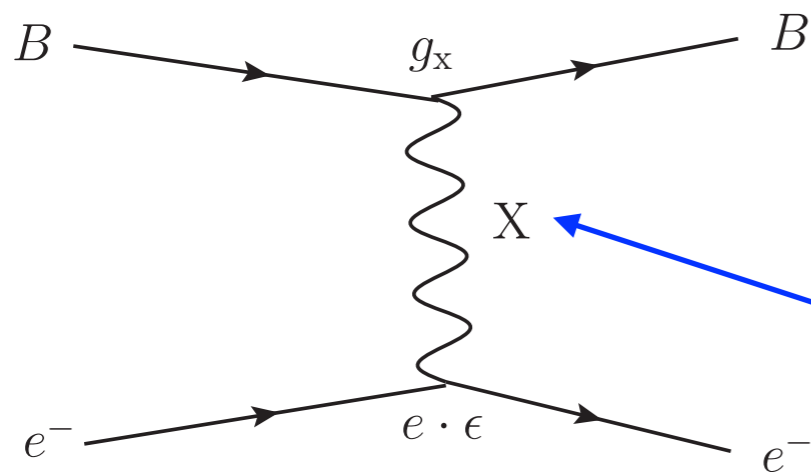


$$\frac{1}{\Lambda^2} \bar{\psi}_A \psi_B \bar{\psi}_B \psi_A$$

$$\frac{1}{\Lambda^2} \bar{\chi}_A \chi_B \bar{\chi}_B \chi_A$$

- 'A' particles self-annihilate producing accelerated 'B' particles with boost factor

$$\gamma = m_A/m_B$$



- boosted DM particles travel to Earth and scatter with SM in the detector

- Interacts through some light mediator particle  $X$

Agashe, et al : [JCAP 10 \(2014\) 062](#)

Alhazmi, Kim, Kong, [Mohlabeng](#), Park, Shin: [JHEP 05 \(2021\) 055](#)

Necib et al: [Phys.Rev. D95 \(2017\) 7, 075018](#)

Dutta et al: [JHEP 01 \(2022\) 144](#)

Kim et al: [JHEP 07 \(2020\) 057](#) + many others

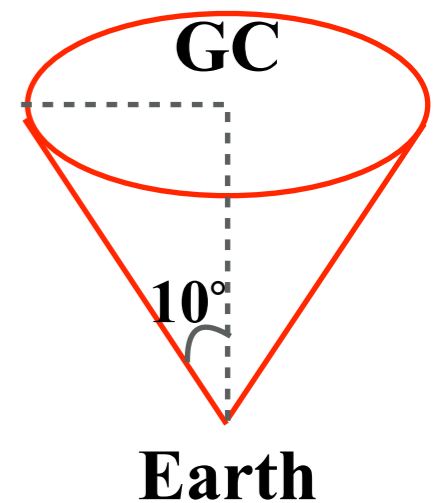


# Boosted dark matter from the Galactic Center

Annihilation of A to boosted B in the Galactic Center

**Flux:** NFW profile +  $10^\circ$  cone around GC

$$\Phi_{GC}^{10^\circ} = 9.9 \times 10^{-8} \text{ cm}^{-2} \text{ s}^{-1} \left( \frac{\langle \sigma_{A\bar{A} \rightarrow B\bar{B} \nu} \rangle}{5 \times 10^{-26} \text{ cm}^3/\text{s}} \right) \left( \frac{20 \text{ GeV}}{m_A} \right)^2$$



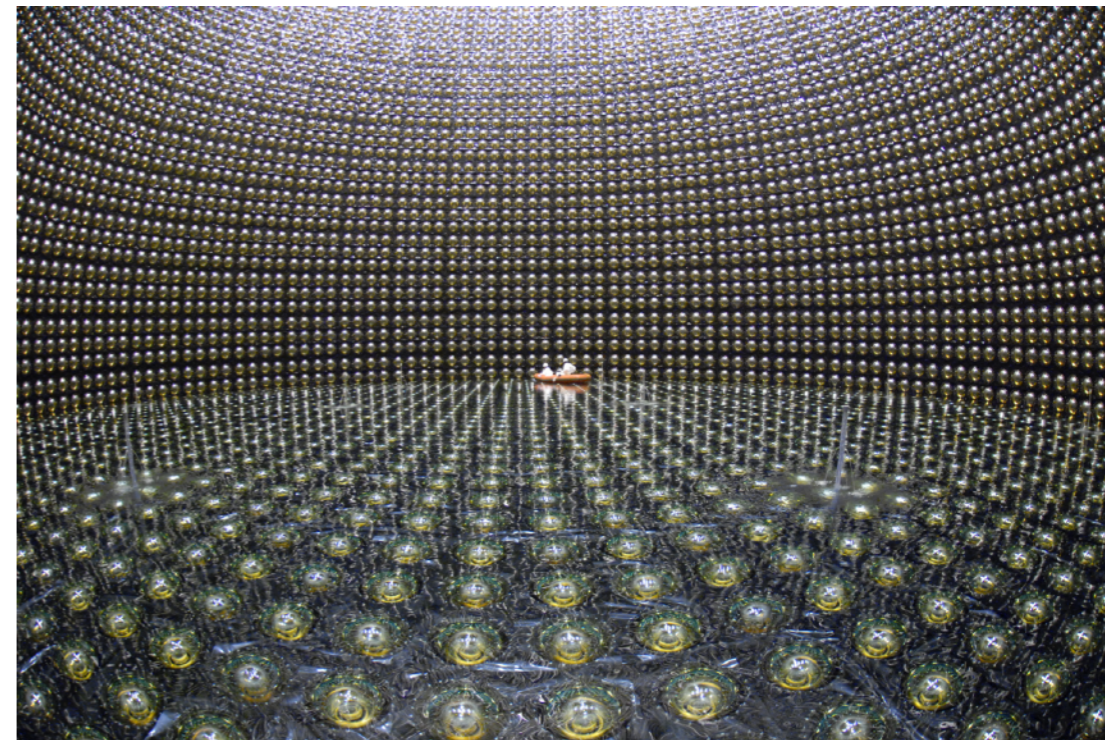
Lower flux means we need large volume detectors for sensitivity

**Neutrino detectors:** Super-K, Hyper-K, Ice-Cube, DUNE

**Dark matter detectors:** XENON1T, DarkSide

Focus on electron scattering

For nucleon scattering:



Berger et al, *Phys.Rev.D* 103 (2021) 9, 095012

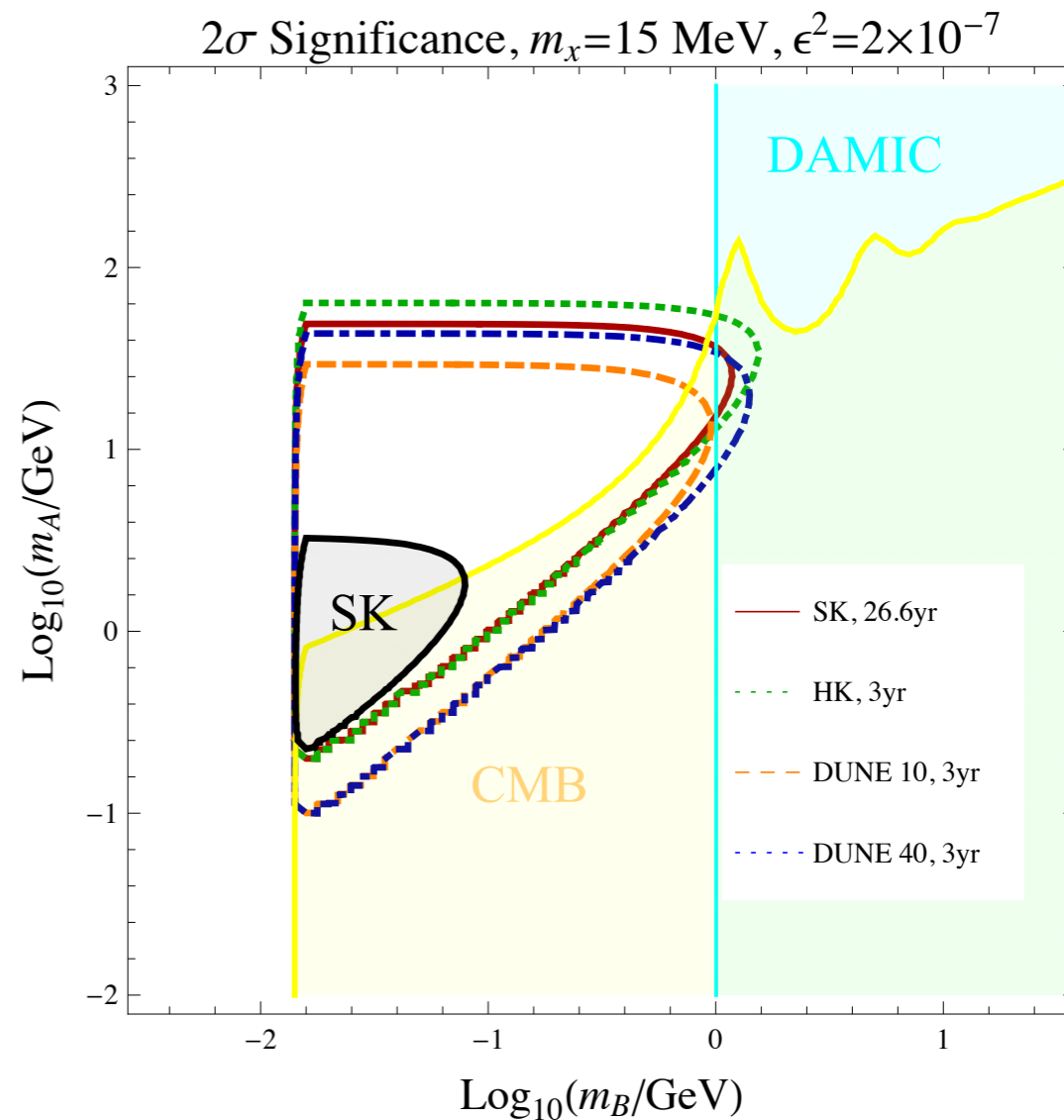
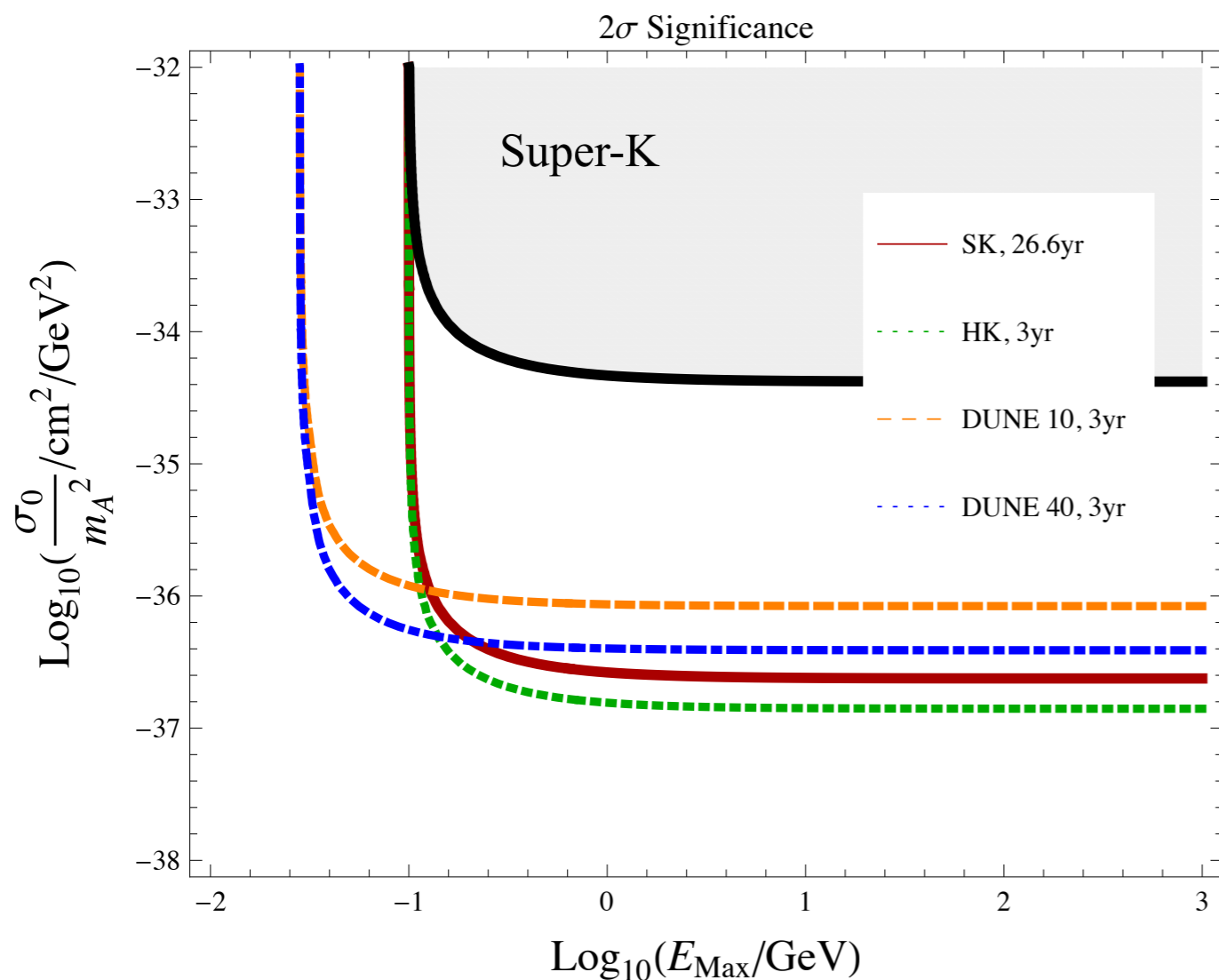


# Boosted DM from the Galactic Center

$$N_{\text{sig}} = \Delta T N_{\text{target}} \Phi_{\text{GC}}^{10^\circ} \int_{E_{\text{thres}}}^{E_{\text{max}}} dE_e \frac{d\sigma_{\text{Be}^- \rightarrow \text{Be}^-}}{dE_e}$$

2 $\sigma$  sensitivity:

$$S^{\theta_{\text{res}}} = \frac{N_{\text{sig}}}{\sqrt{N_{\text{BG}}^{\theta_{\text{res}}}}}$$



Alhazmi, Kong, Mohlabeng, Park: JHEP 04 (2017) 158





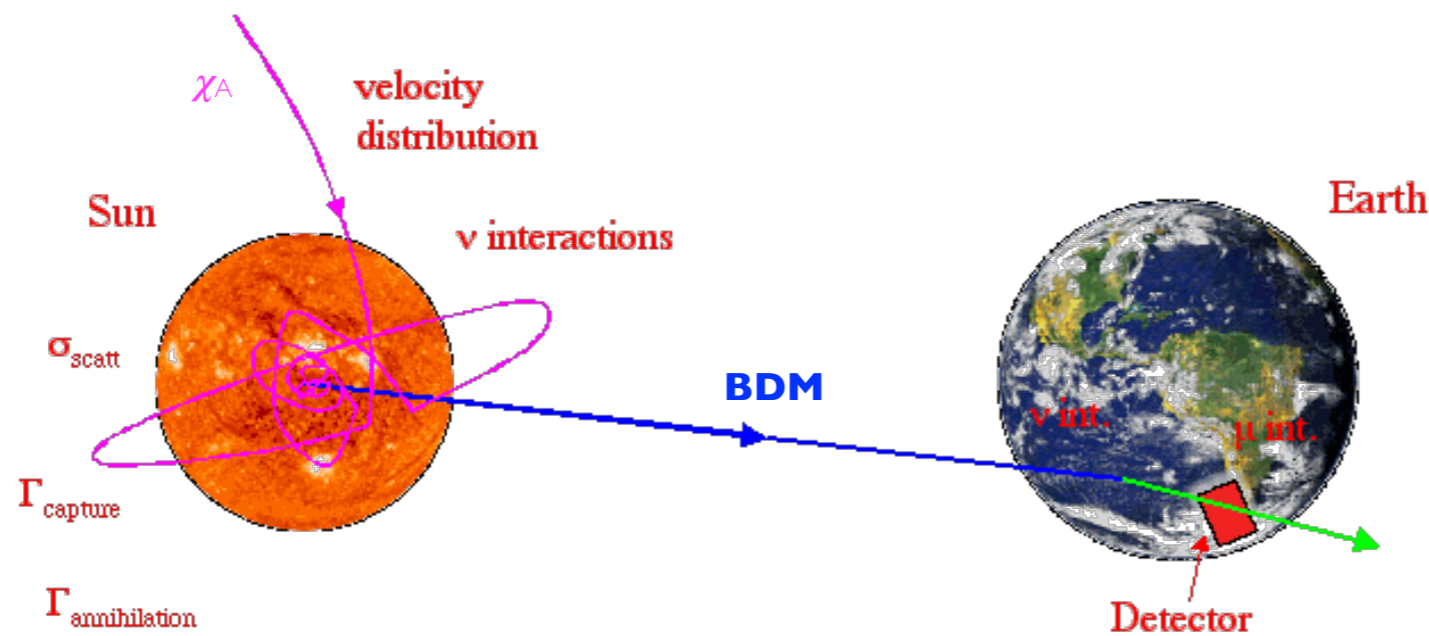
# Boosted DM from the Sun

$\chi_A$  can get captured in the Sun and annihilate to  $\chi_B$

$\chi_B$  travel to earth and scatter in the detector

The Sun is a point-like source so we don't consider an observation angle

Including self-interactions enhances the capture rate in the Sun



Kong, Mohlabeng, Park: *Phys. Lett. B* 743 (2015) 256-266

Berger et al: *JCAP* 02 (2015) 005

Alhazmi, Kong, Mohlabeng, Park: *JHEP* 04 (2017) 158

Bhattacharya et al: *JCAP* 05 (2017) 002

## Boosted DM from the Sun

Time evolution of  $\chi_A$  number density in the Sun is

$$\frac{dN_\chi}{dt} = C_c + (C_s - C_e)N_\chi - (C_a + C_{se})N_\chi^2$$

**$C_c$** : capture rate by nuclei inside Sun

**$C_s$** : capture rate by DM already captured in Sun

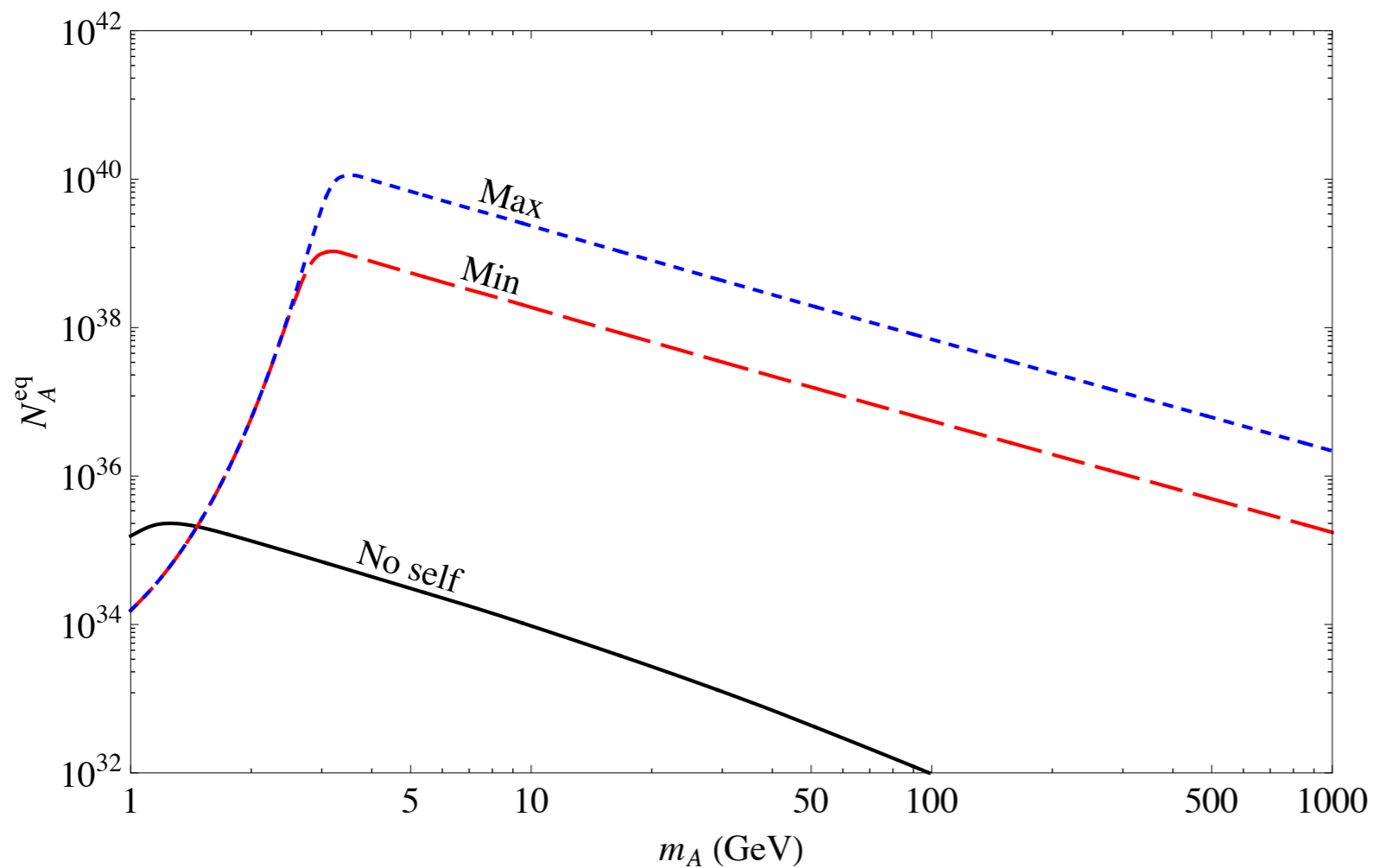
**$C_e$** : Evaporation rate due to DM-nuclei scattering

**$C_{se}$** : evaporation rate due to DM-self interaction

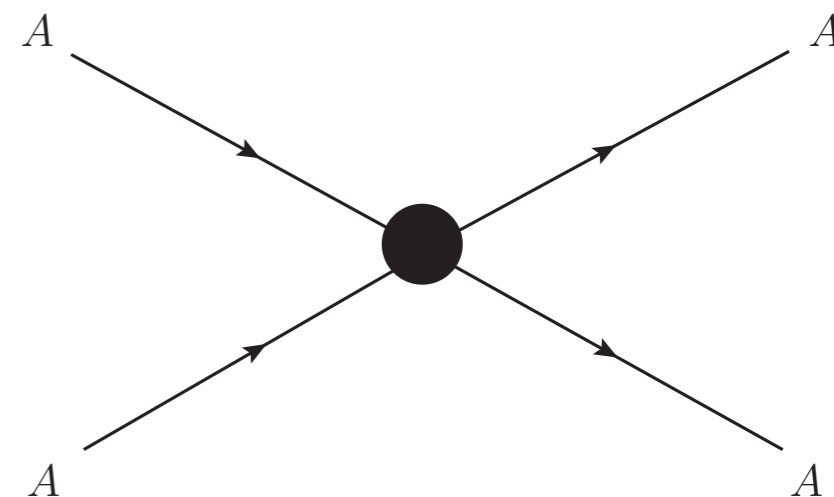
**$C_a$** : annihilation rate



# Importance of DM Self-interactions



Kong, Mohlabeng, Park: *Phys. Lett. B* 743 (2015) 256-266



Kinematics of  
dwarf spheroidal  
galaxies

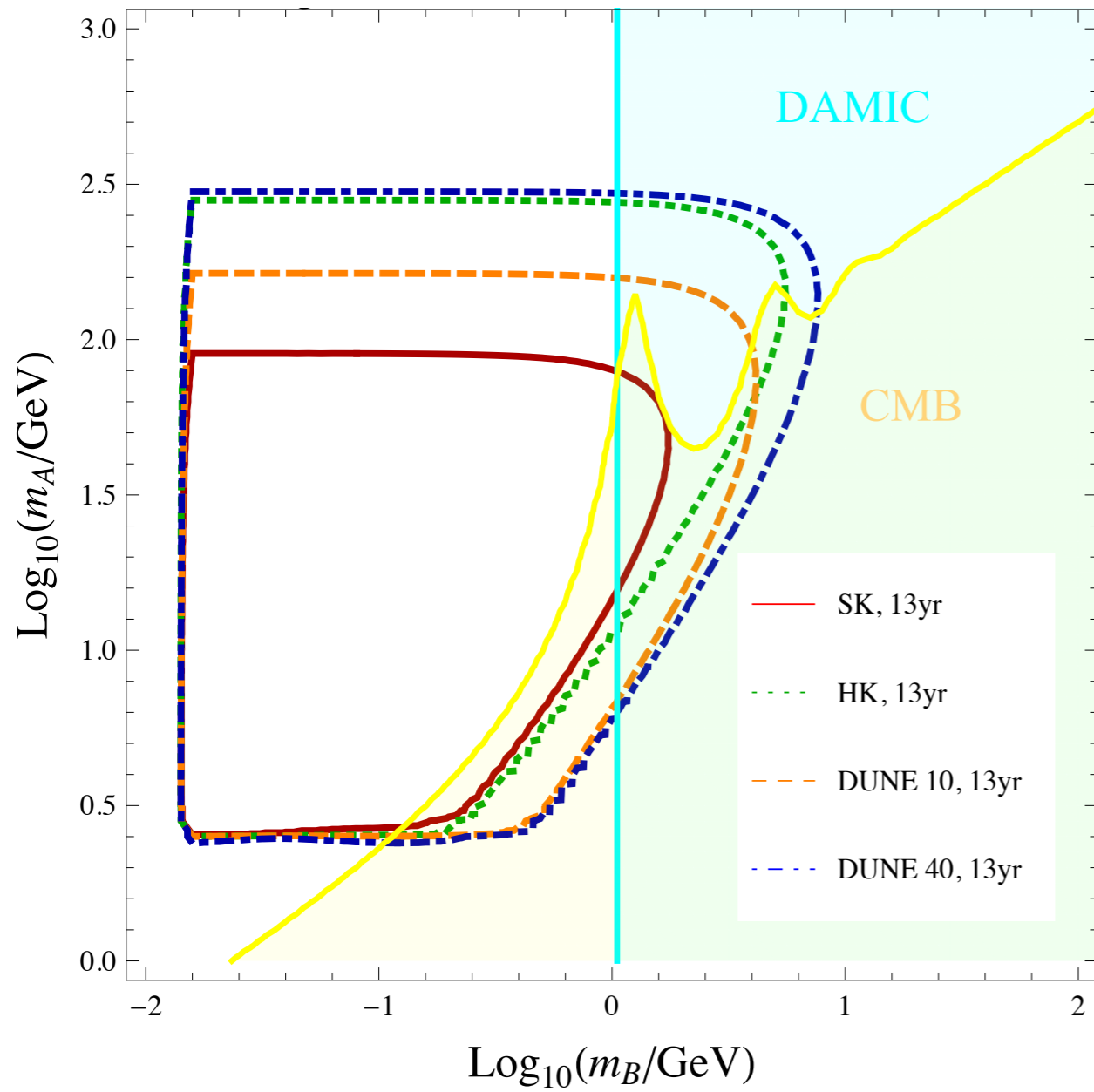
$$0.1 \text{ cm}^2/g < \frac{\sigma_{AA}}{m_A} < 1.25 \text{ cm}^2/g$$

↑  
min
↑  
max

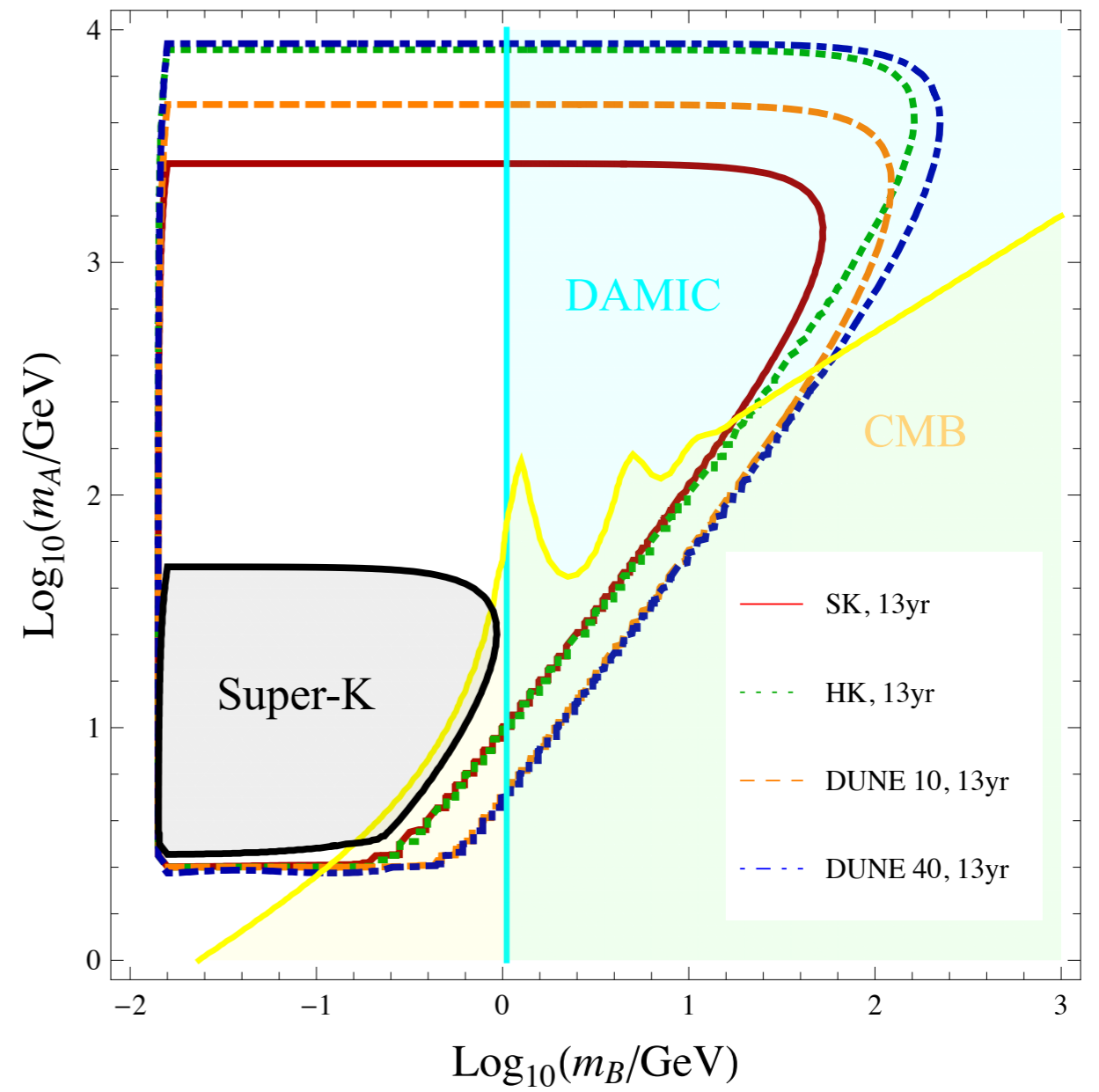
Matter  
distribution of  
bullet cluster



90% CL, Min-self



90% CL, Max-self



Alhazmi, Kong, Mohlabeng, Park: *JHEP* 04 (2017) 158



# Boosted DM at direct detection experiments

- Large volume DD experiments can look for lower  $A$  masses

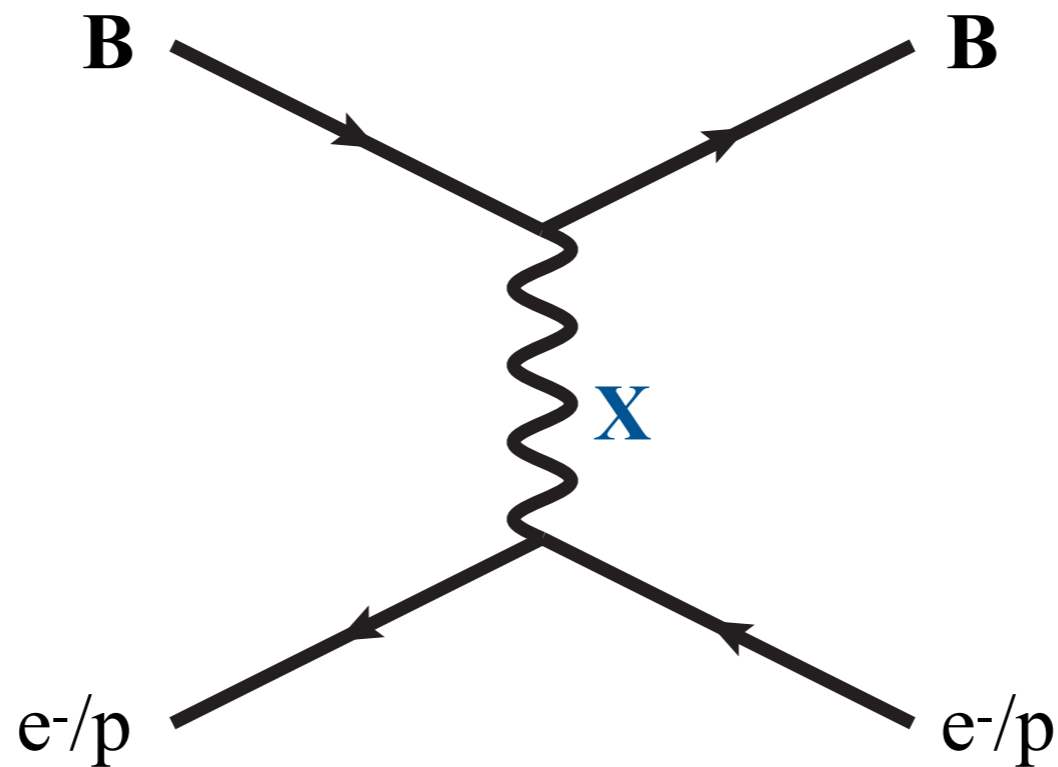
XENON1T

DarkSide

...

} Lower threshold than neutrino detectors

Scattering with either nucleon or electrons



Focus on electron scattering



To obtain recoil rates:

$$N_{sig} = \Phi_B \sigma_{Be} N_e^{eff} t_{exp}$$

BDM flux → Scattering cross-section → Effective # of electrons → exposure time

However, it is important to include atomic effects related to DM - e scattering

$$\frac{dN_{sig}}{dE_{eR}} = \Phi_B t_{exp} N_{Xe} \frac{d\sigma_{Be} v_{rel}}{dE_{eR}}$$

# of Xenon atoms in volume

Ionization form factor

Differential scattering cross-section

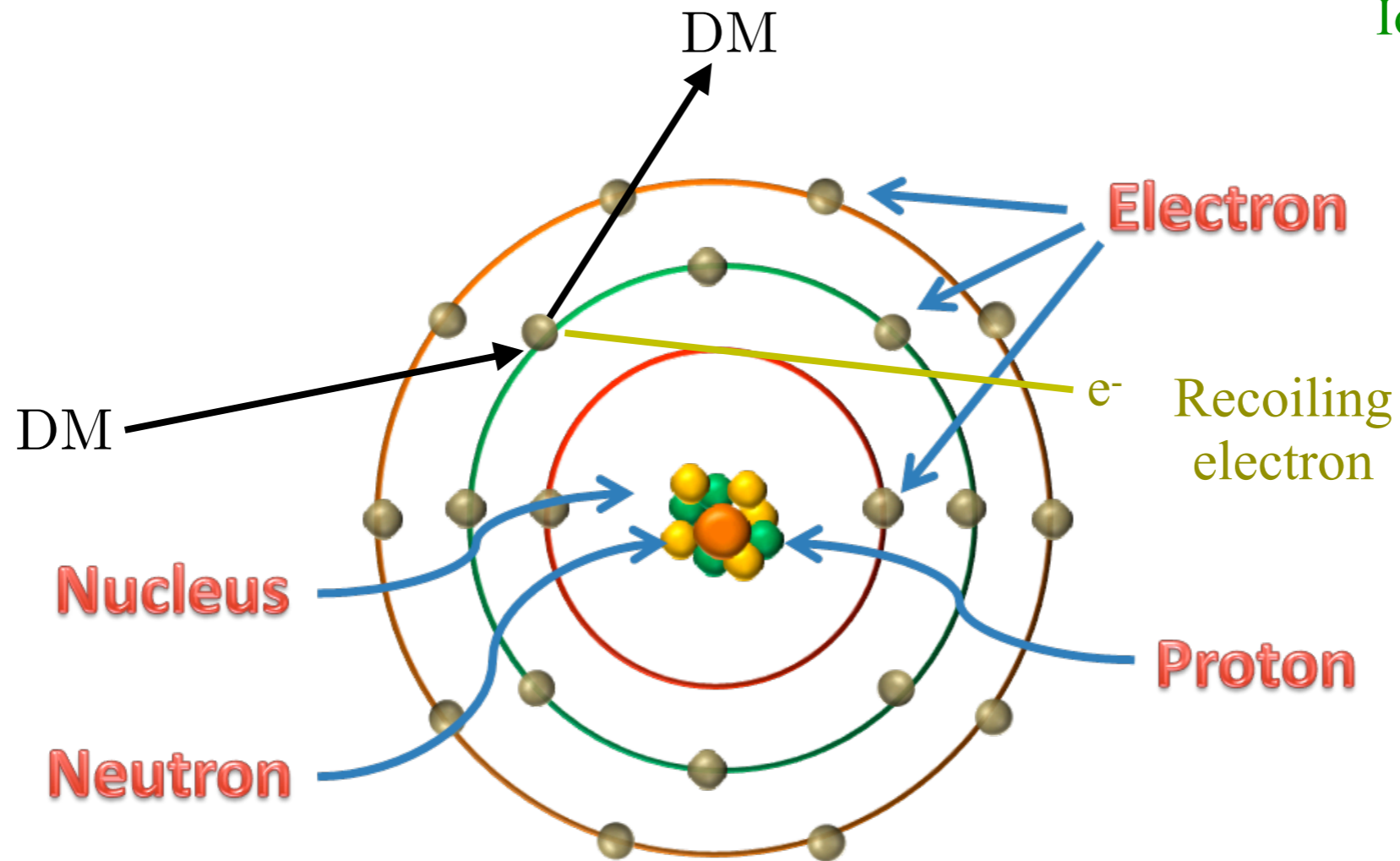
$$\frac{d\sigma_{Be} v_{rel}}{dE_{eR}} = \frac{1}{64\pi} \frac{1 - v_{rel}^2}{v_{rel}} \frac{1}{m_B^2 E_{eR} (2m_e + E_{eR}) (m_e - |E_{nl}^B|)} \int_{q_{min}}^{q_{max}} dq q |\mathcal{M}|^2 \underbrace{|f_{ion}(E_{eR}, q)|^2}_{\text{Ionization form factor}}$$



# Ionization effects:

$$\frac{d\sigma_{Be}v_{rel}}{dE_{eR}} = \frac{1}{64\pi} \frac{1 - v_{rel}^2}{v_{rel}} \frac{1}{m_B^2 E_{eR} (2m_e + E_{eR}) (m_e - |E_{nl}^B|)} \int_{q_{min}}^{q_{max}} dq q |\mathcal{M}|^2 \underbrace{|f_{ion}(E_{eR}, q)|^2}_{\text{Ionization form factor}}$$

Ionization form factor



- Bound electrons have non-negligible momentum dependence
- Ionization function takes into account electron momentum dependence
- Electrons are bound in different orbitals with binding energies, ionization function accounts mom transfer required to ionize electron from orbitals

$$|f_{ion}(E_{eR}, q)|^2 = \frac{2k'^3}{(2\pi)^3} \int dr^3 \psi_{ef}^*(\mathbf{r}) e^{i\mathbf{q}\cdot\mathbf{r}} \psi_{ei}(\mathbf{r})$$

↑
↑  
free electron
bound electron  
wave-function
wave-function

- Different functions considered

**Plane - Wave:** bound electron wave function is described by **Roothaan-Hartree-Fock** wavefunctions

Bunge et al: *Atom. Data Nucl. Data Tabl.* 53 (1993) 113-162

Outgoing electron wave function is described by **plane wave**

Essig et al: *Phys. Rev. D* 85 (2012) 076007

Kopp et al: *Phys. Rev. D* 80 (2009) 083502

Cao et al: *Chin. Phys. C* 45 (2021) 4, 045002 + many others



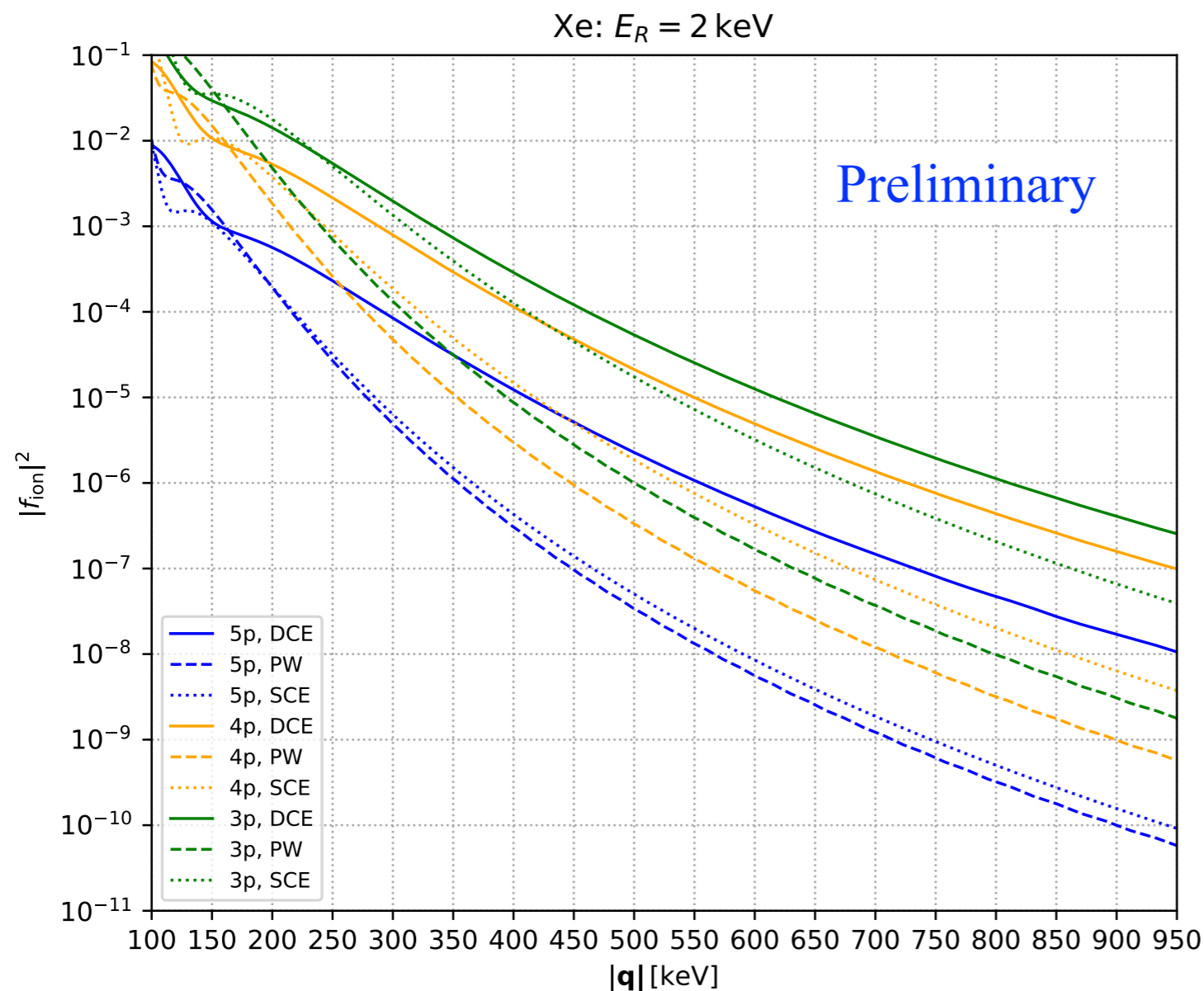


**Relativistic ionization function:** - bound and ionized electron wave functions are obtained by solving relativistic Dirac equation

i.e. solve  $\hat{h}\psi_{nk} = E_{nk}\psi_{nk}$

using Dirac hamiltonian  $\hat{h} = \alpha \cdot \mathbf{p} + m_e(\beta - 1) + V_{eff}(r)$

- accounts for Lorentz structure of DM - e interactions



Roberts et al: *Phys. Rev. D* 93 (2016) 115037

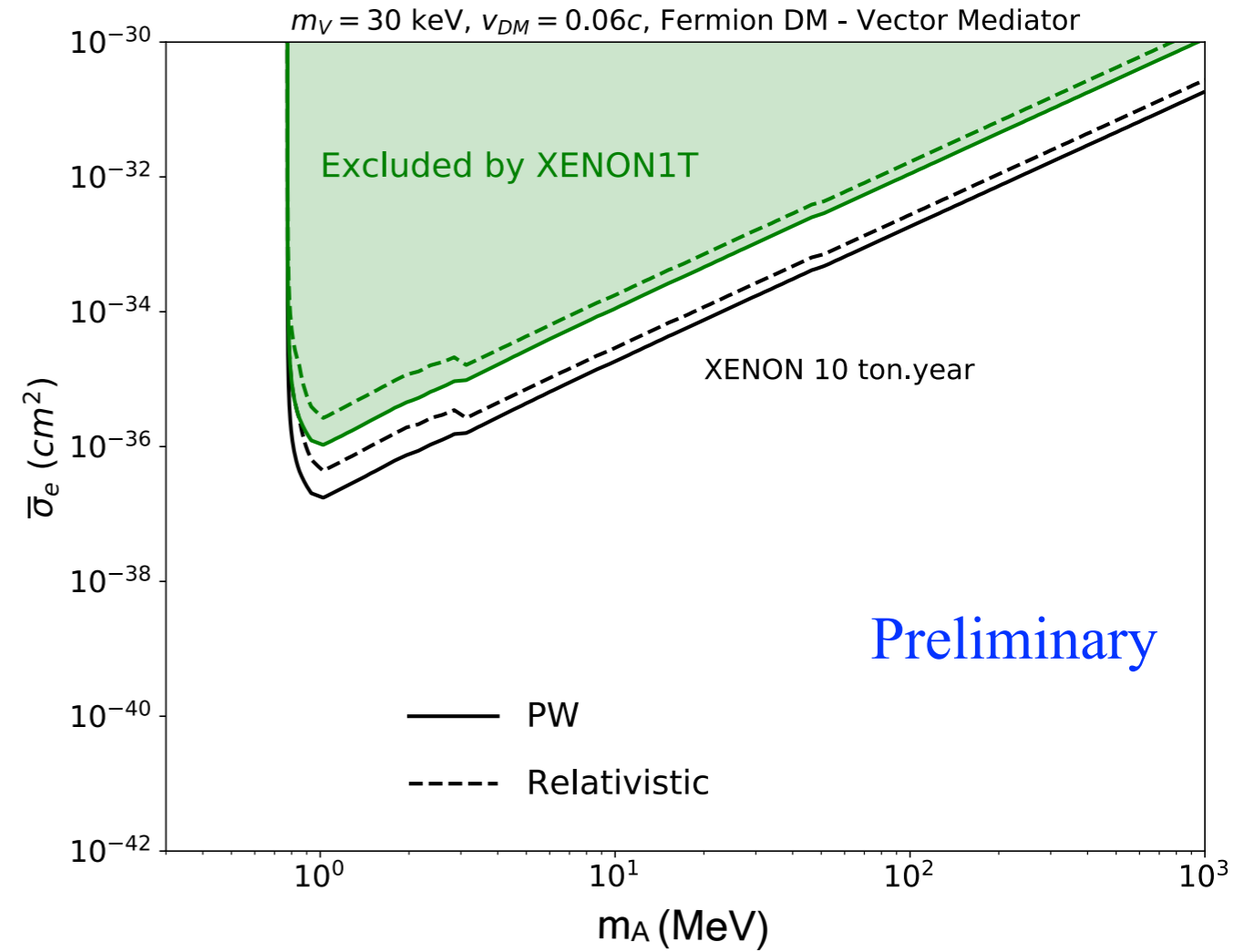
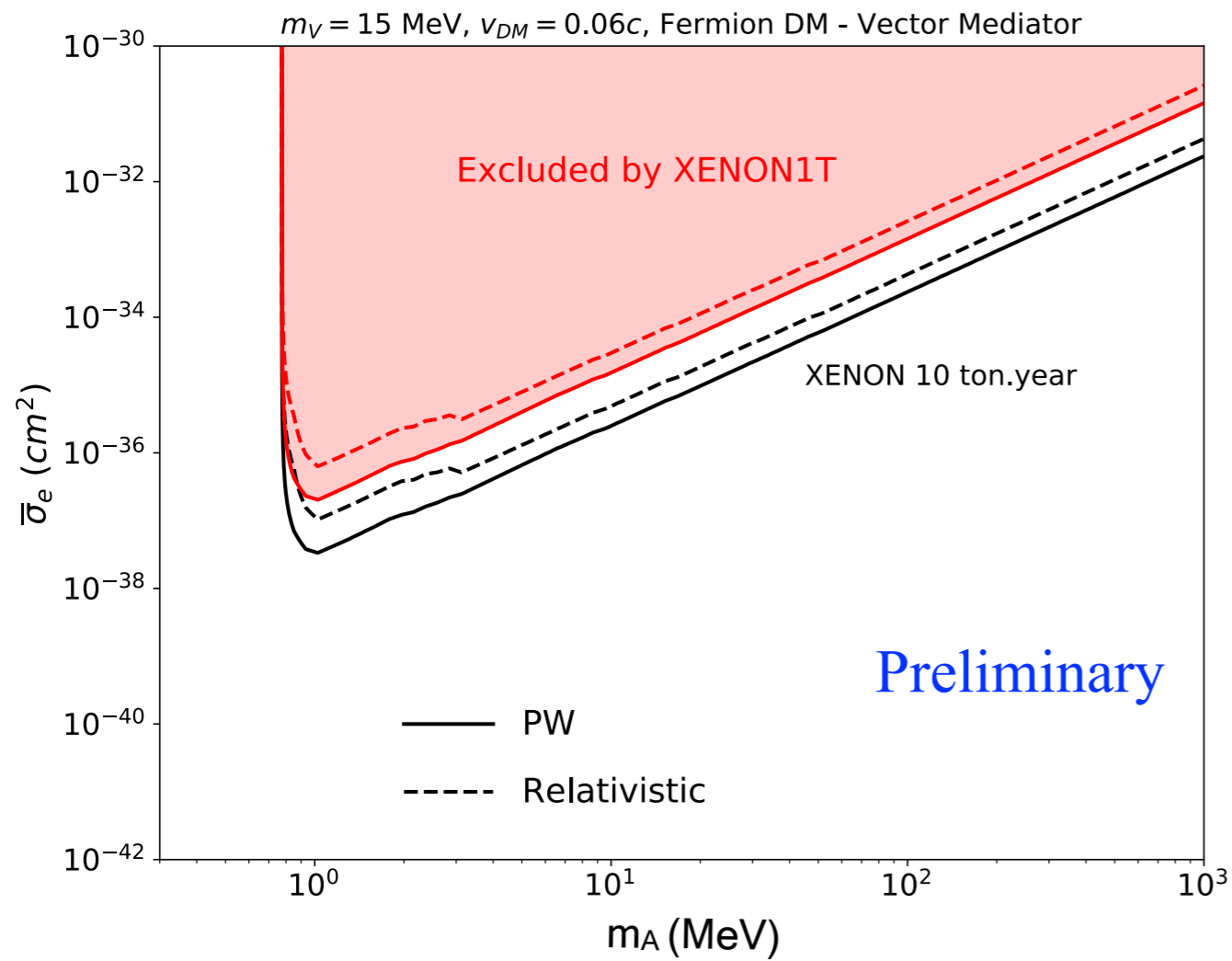
Roberts, Flambaum: *Phys. Rev. D* 100 (2019) 063017

At high momentum transfer,  
relativistic & PW are different

Alhazmi, Kim, Kong, Mohlabeng, Park, Shin:  
In Progress

# Some Limit plots

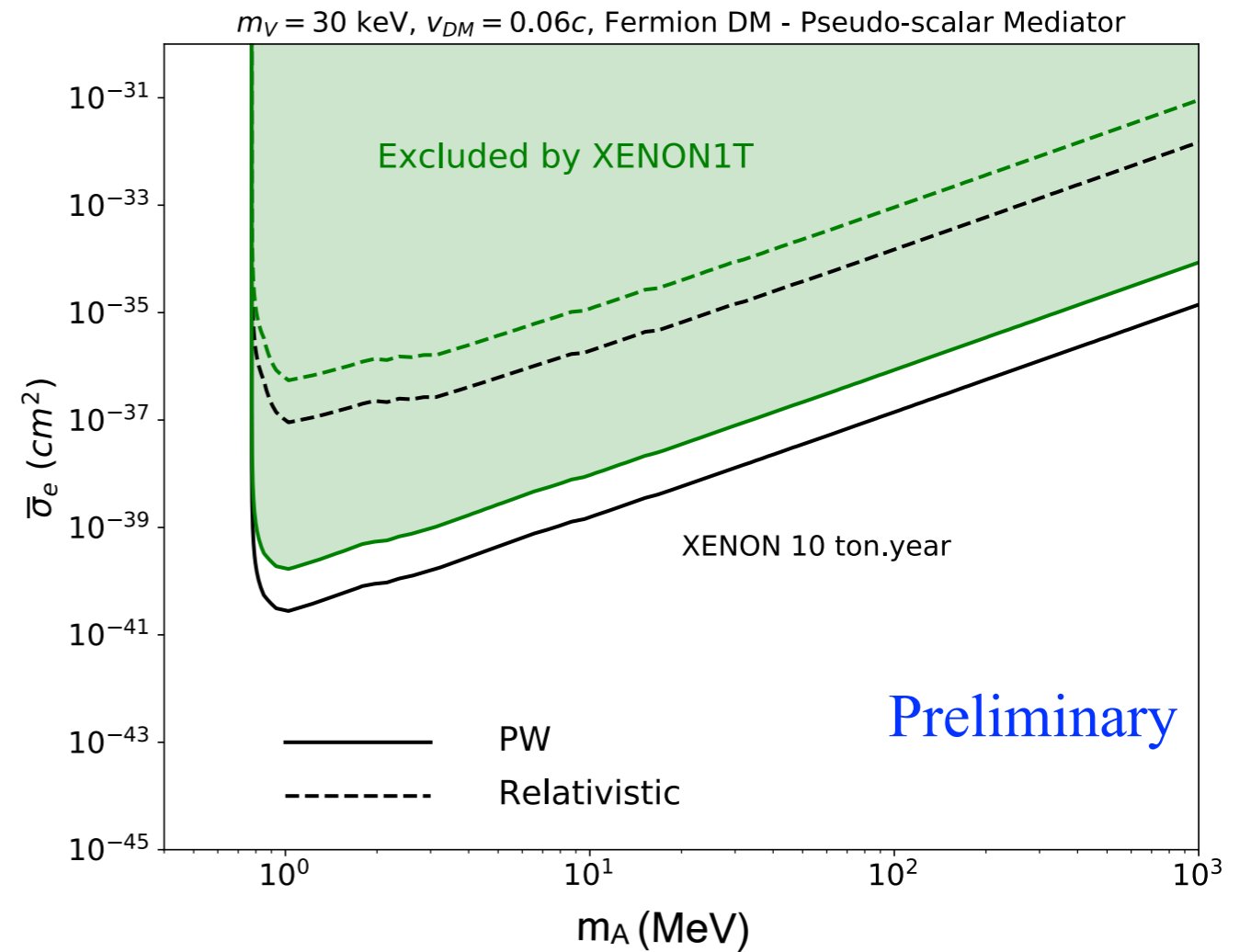
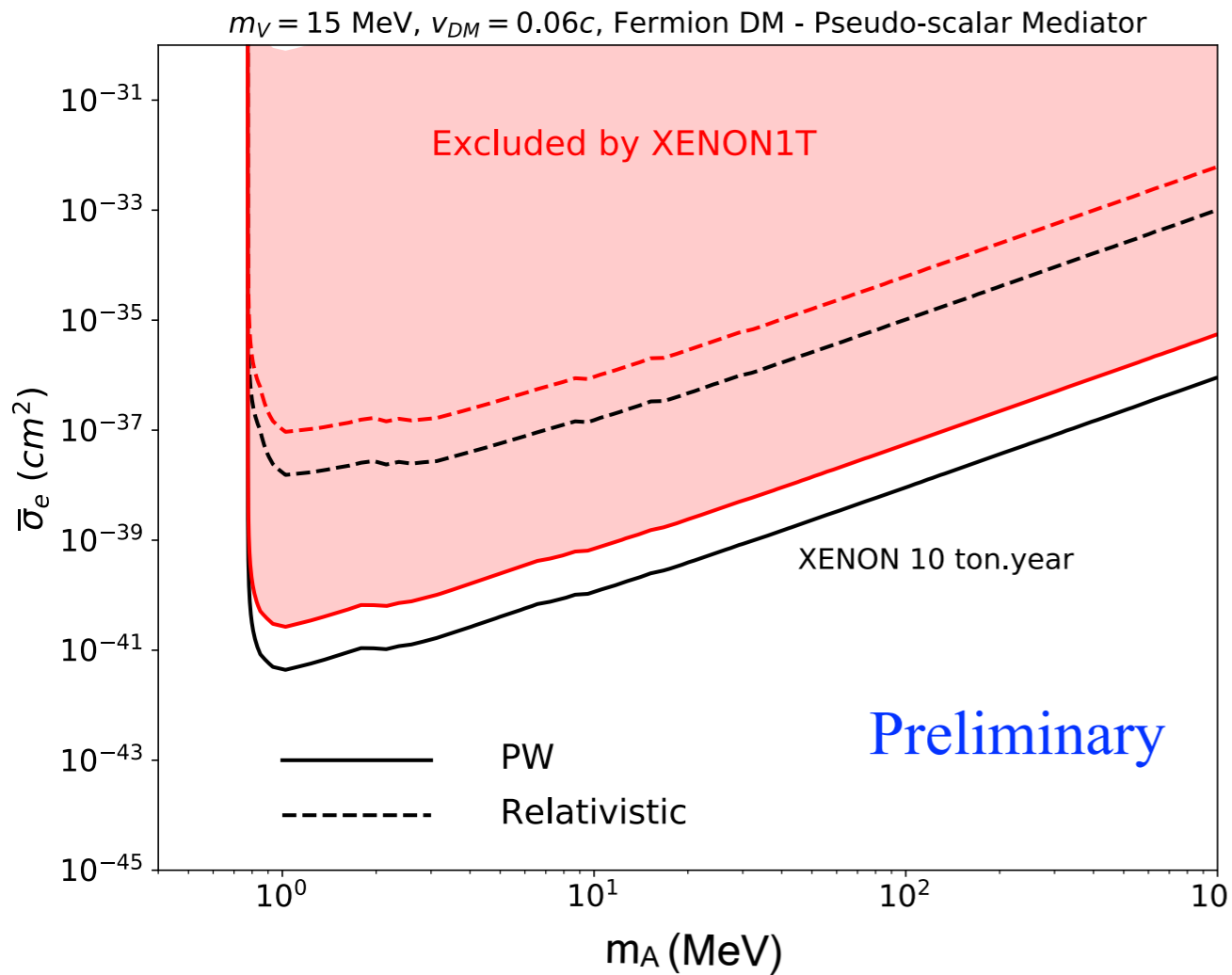
## Fermion Dark Matter with a vector mediator



Alhazmi, Kim, Kong, Mohlabeng, Park, Shin: In Progress



# Fermion Dark Matter with a pseudo-scalar mediator



Alhazmi, Kim, Kong, Mohlabeng, Park, Shin: In Progress



# Outlook

Accelerated DM is interesting phenomenological prospect

Can give striking signals at large volume neutrino detectors

Can produce interesting signals in direct detection experiments

Atomic effects can be important for accelerated DM-e scattering

Need to be careful when considering atomic effects

---

Ongoing work, need to include other limits

**Xenon100, Xenon10, Darkside50, Super-K, BBN?**

Working on examining ionization effects from general accelerated DM scenarios

**Cosmic-ray up-scattered DM**

**Neutrino up-scattered DM + others**



Thank you

